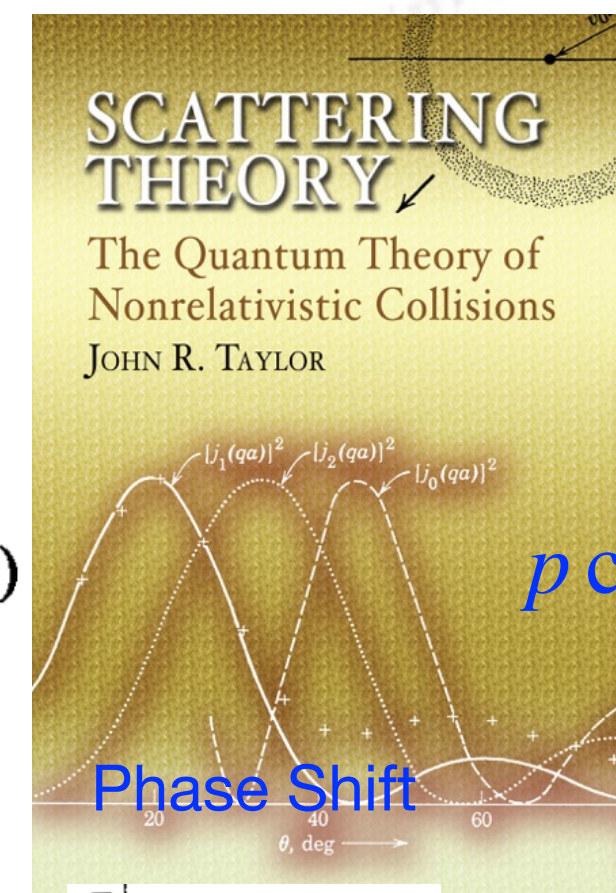
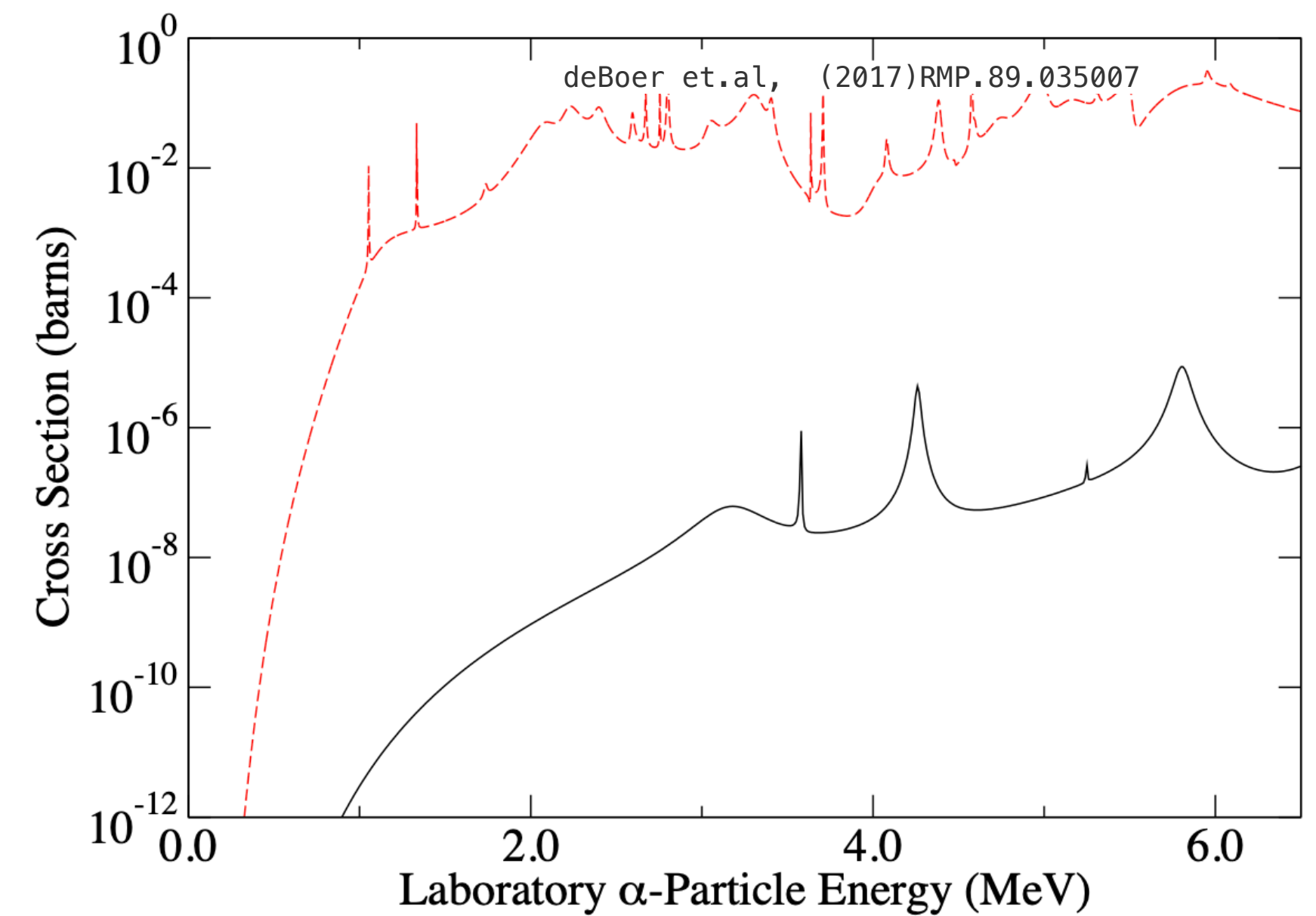
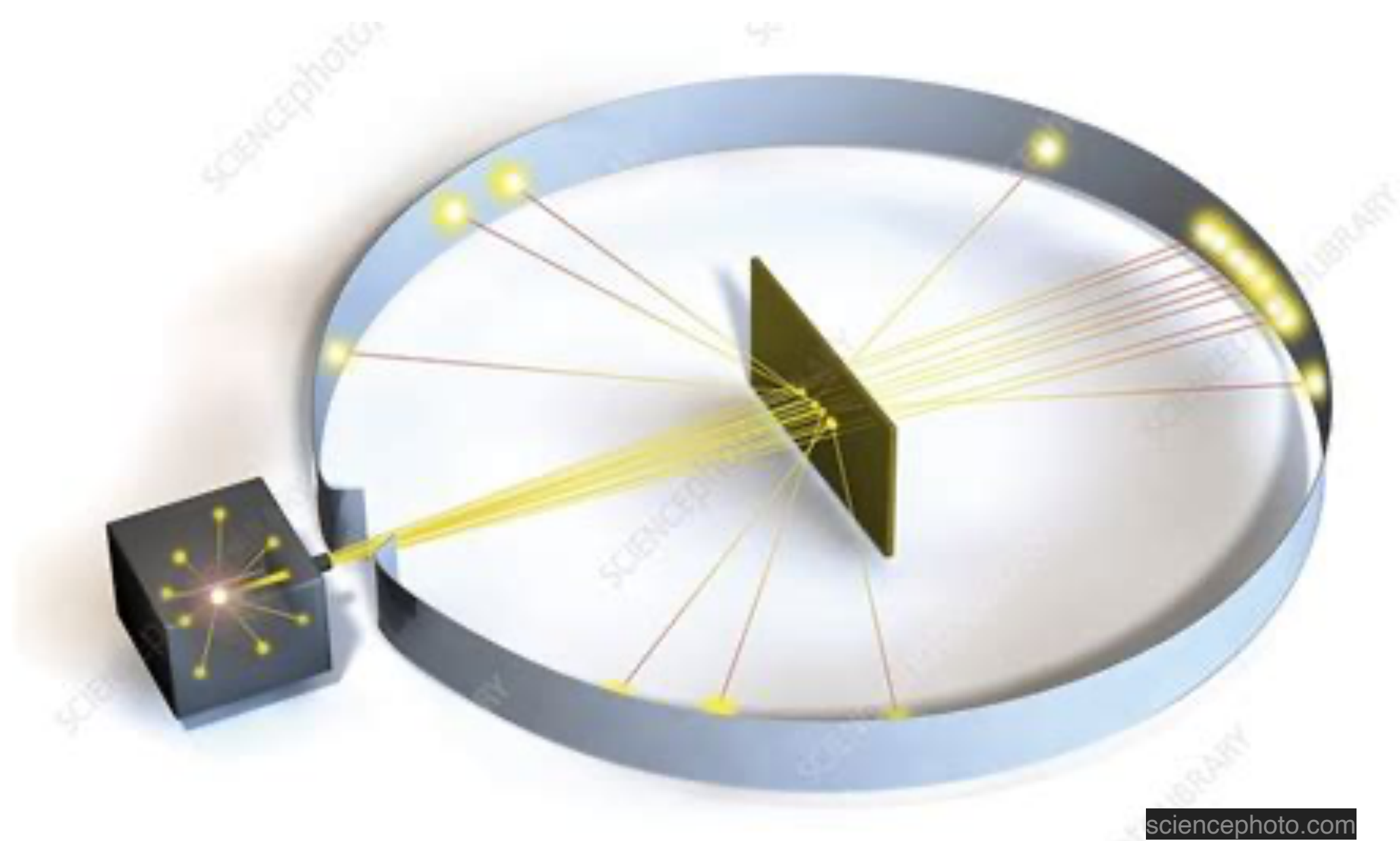


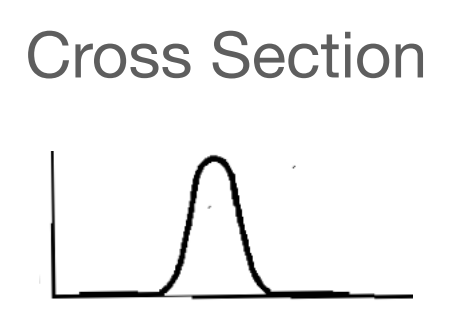
Nucleons in a finite volume

from ground states to the continuum

From Scattering Experiments to EFT



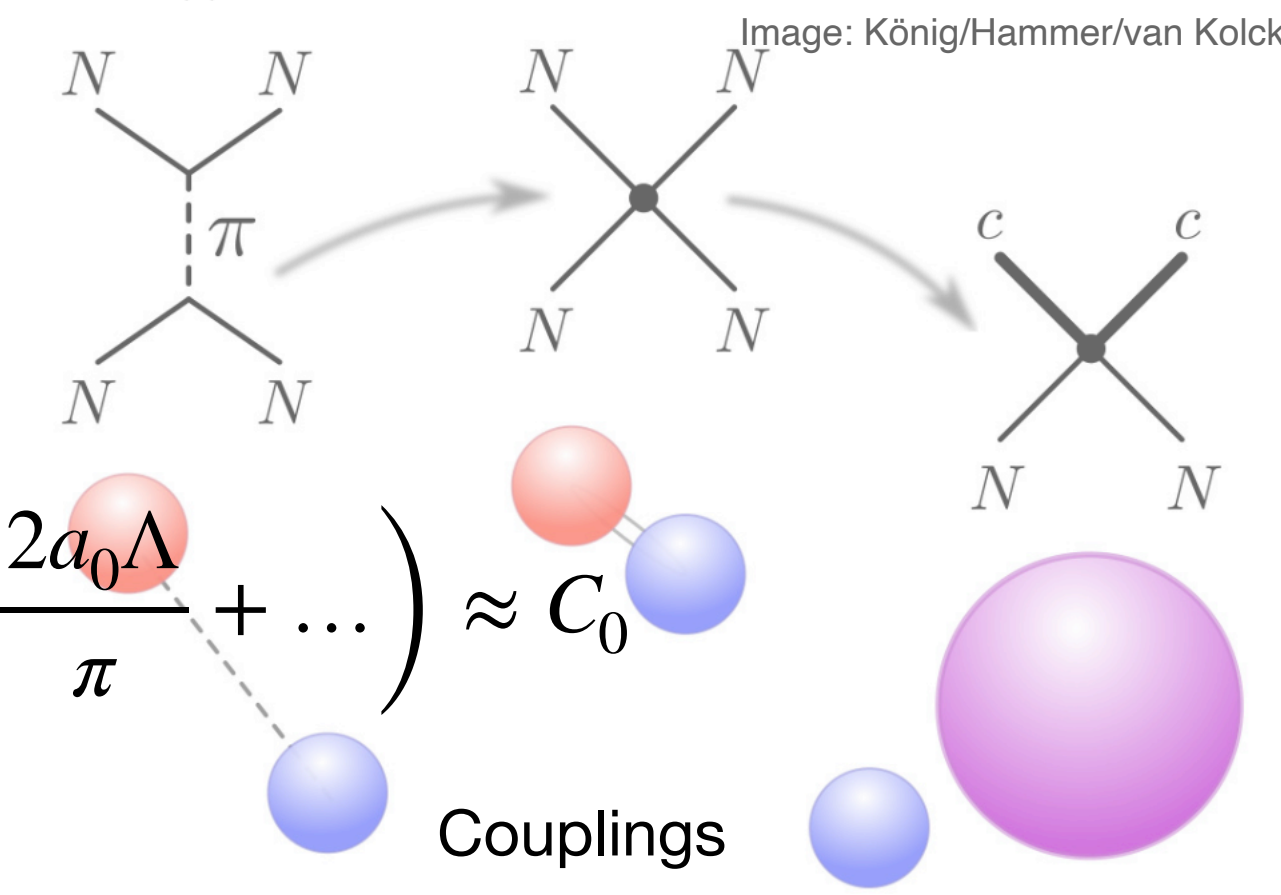
$$\sigma_l(p) = \frac{4\pi(2l + 1)}{p^2} \sin^2 \delta_l(p)$$



$$p \cot(\delta_0(p)) = -a_0^{-1} + \frac{1}{2}r_0p^2 + \dots$$

ERE

$$\frac{8\pi}{m} \left(1 + \frac{2a_0\Lambda}{\pi} + \dots \right) \approx C_0$$



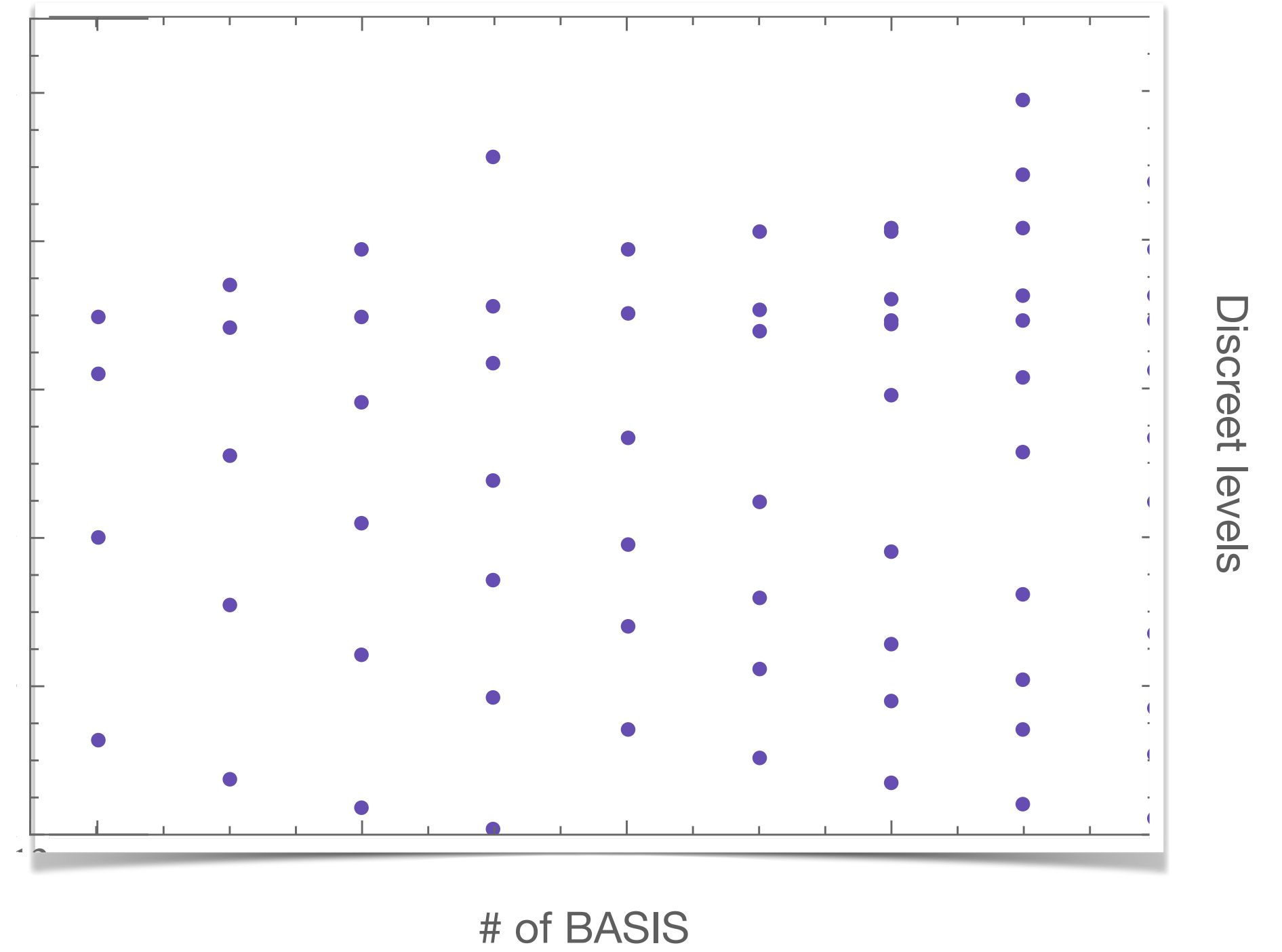
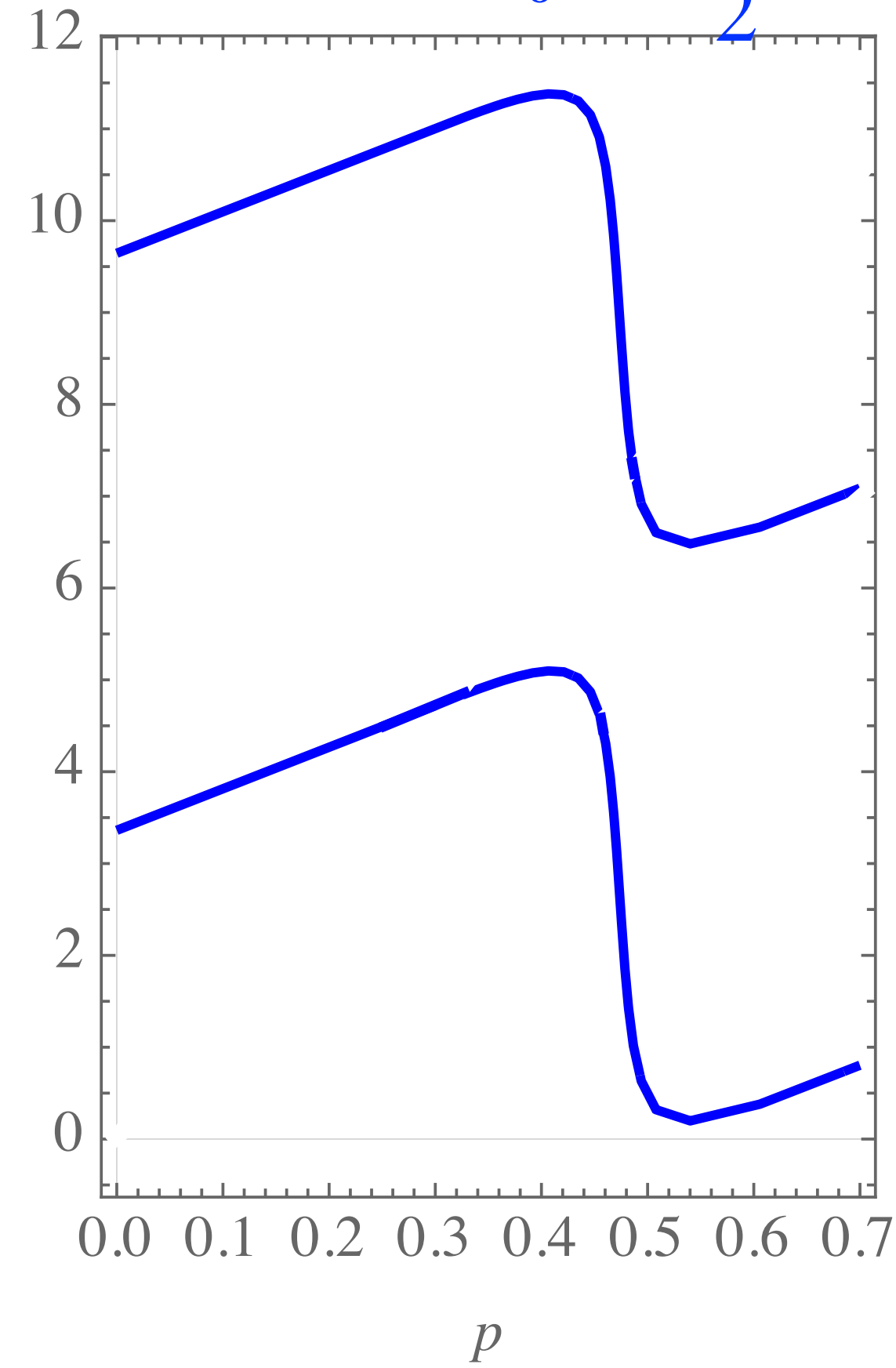
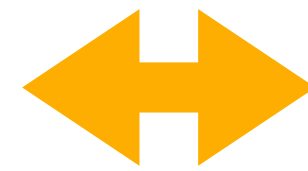
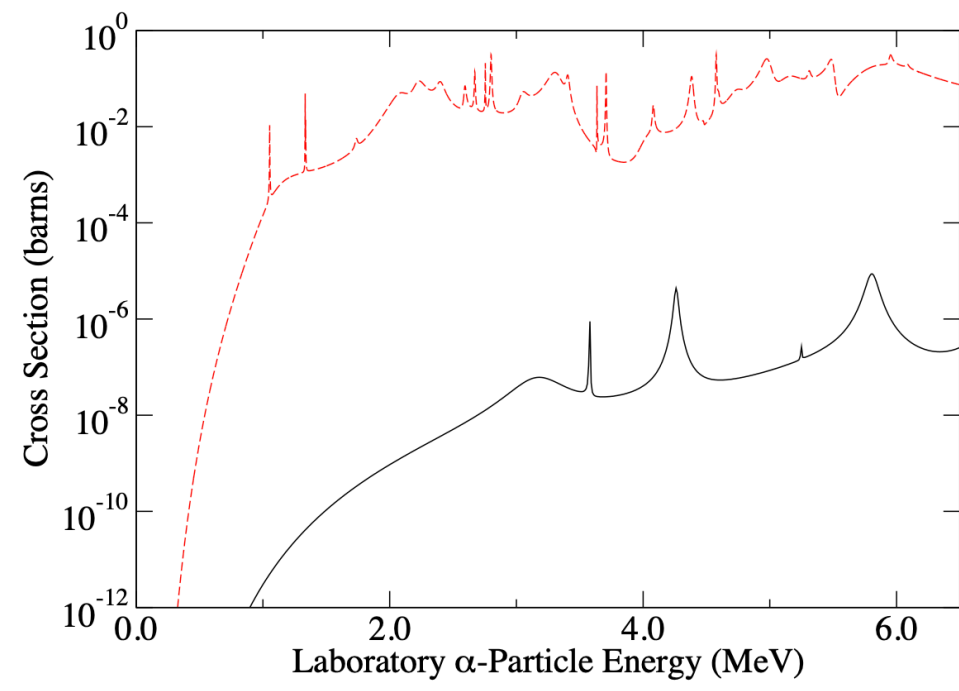
Couplings

From Theory to Experiments?

$$p \cot(\delta_0(p)) = -a_0^{-1} + \frac{1}{2}r_0 p^2 + \dots$$

$H =$

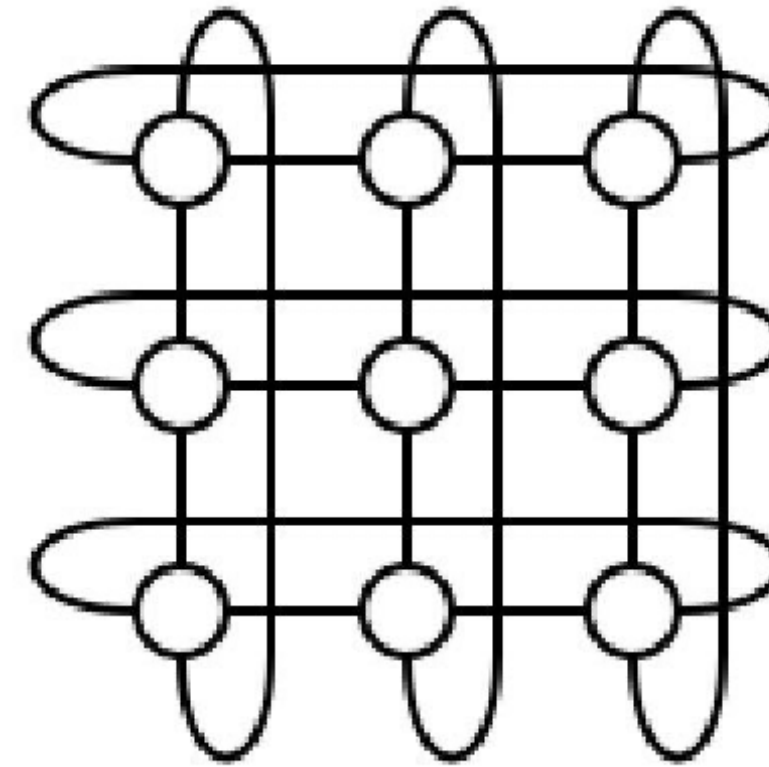
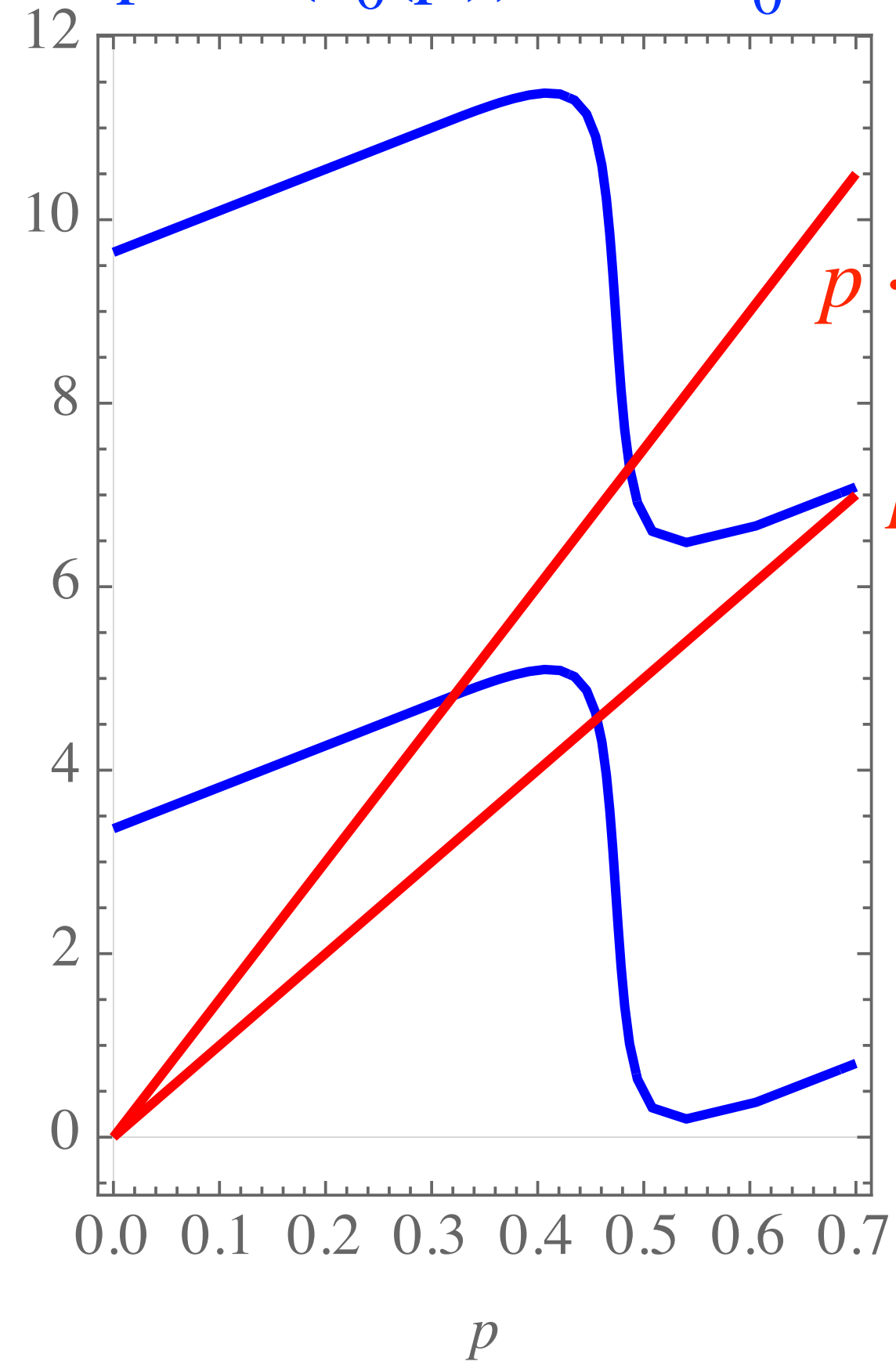
$$\sum_{pq} K_{pq} a_p^\dagger a_q + \frac{1}{4} \sum_{pqrs} V_{pqrs} a_p^\dagger a_q^\dagger a_s a_r$$



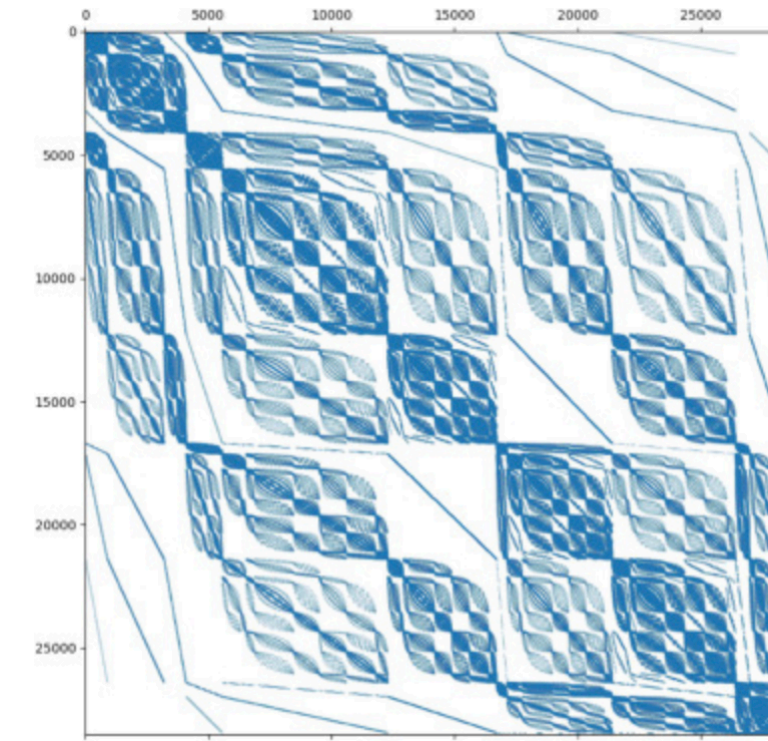
Analytical tools: Finite Volume effects



$$p \cot(\delta_0(p)) = -a_0^{-1} + \frac{1}{2}r_0p^2 + \dots$$



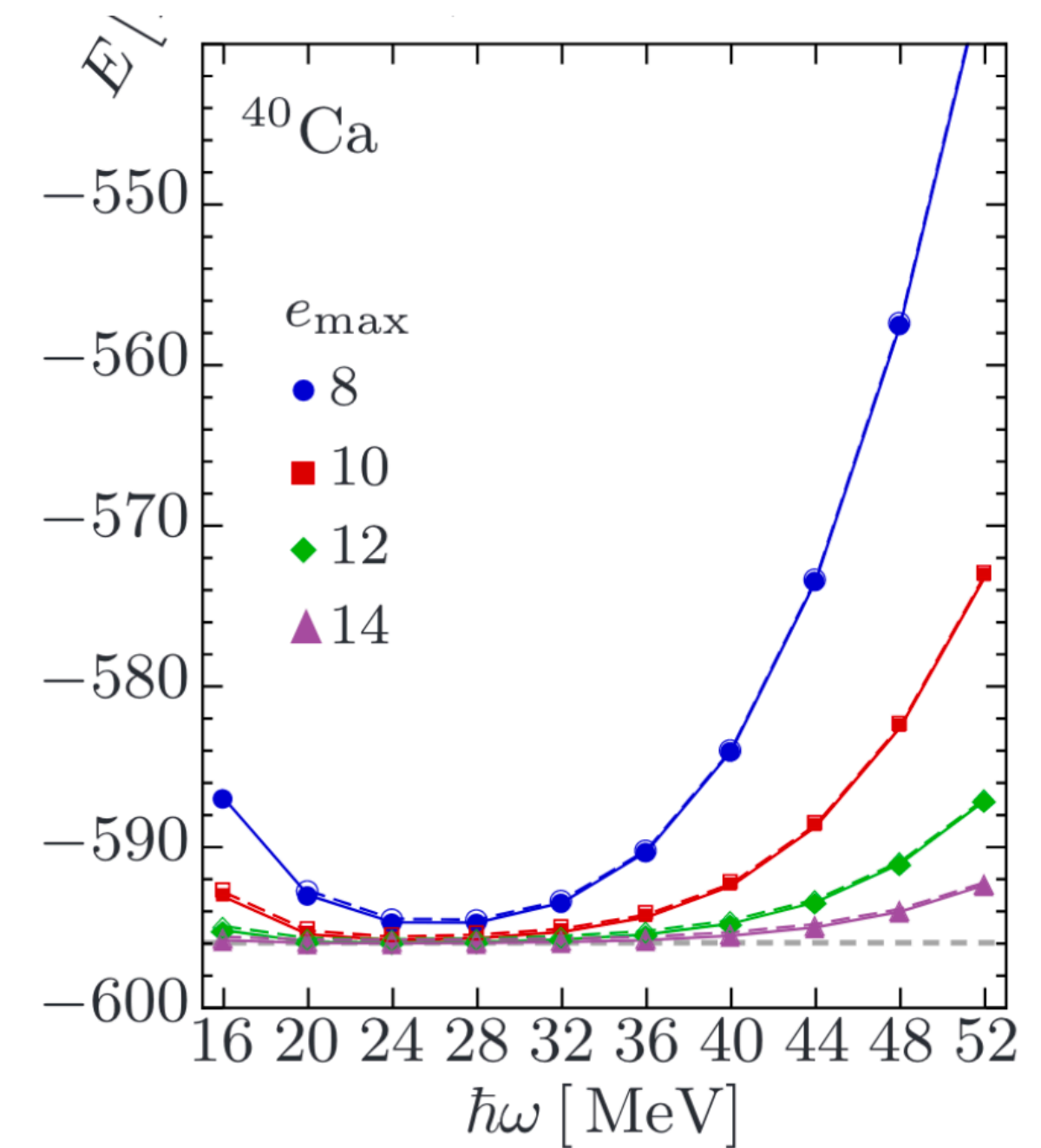
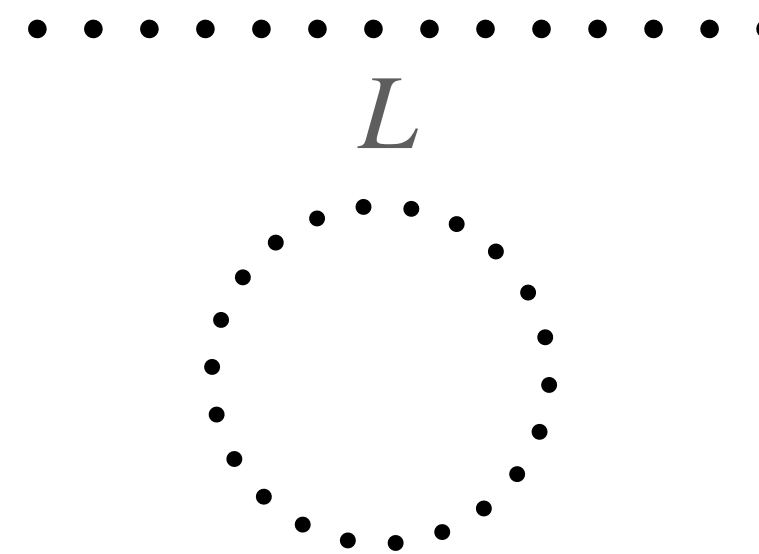
$$H =$$



N. Shimizu Comp. Phys. Comm. 244, 372 (2019)

Boundary Condition!

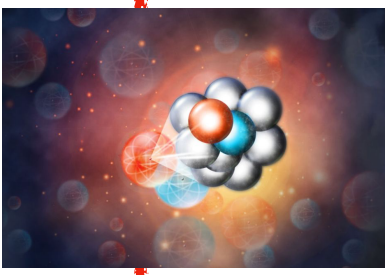
$$\alpha \cos(p \cdot L/2) + \beta \sin(p \cdot L/2) = 0$$



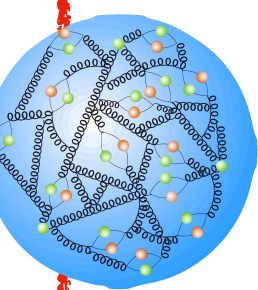
H. Hergert j.physrep.2015.12.007



Λ_{LO}



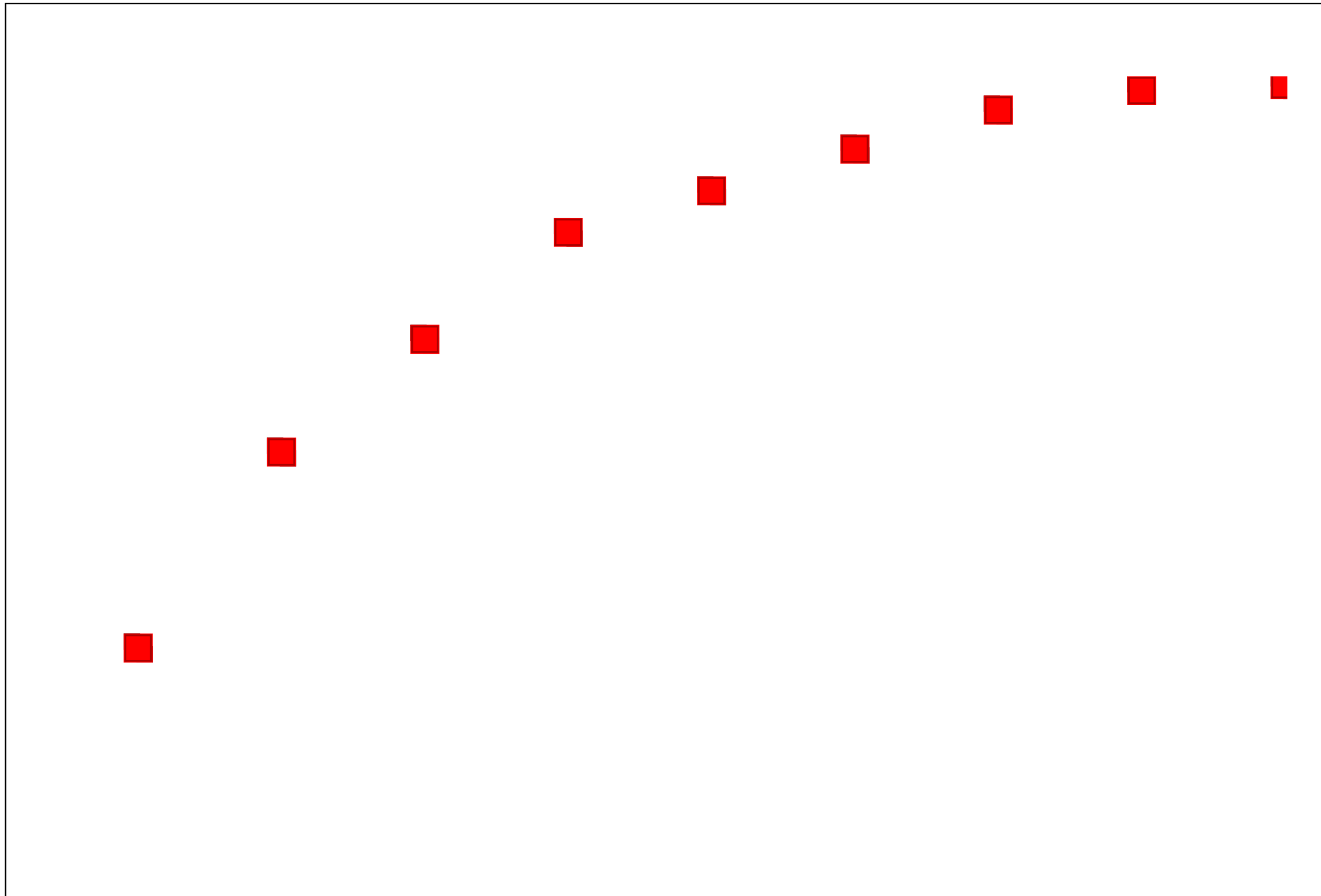
Λ_{Hi}



$$e^{-i\delta_0(p)} = S_0(p) \propto \frac{C_0^2}{ip - \kappa} - \frac{1}{a_0^*} + \frac{1}{2}r_0^*(ip - \kappa)^2 + \dots$$

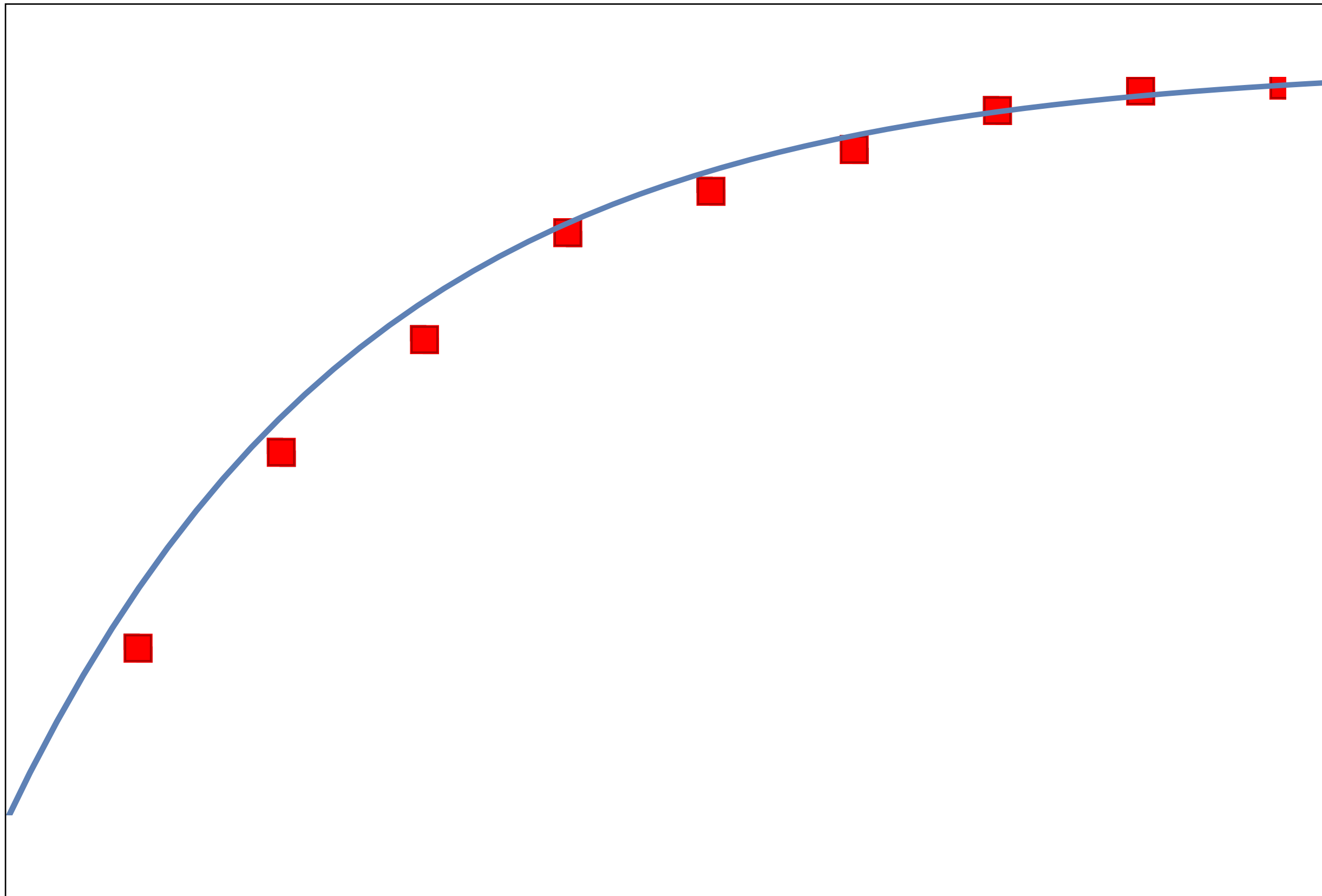
$$e^{-i\delta_0(p)} = S_0(p) \propto \frac{C_0^2}{ip - \kappa} - \frac{1}{a_0^*} + \frac{1}{2}r_0^*(ip - \kappa)^2 + \dots$$

L

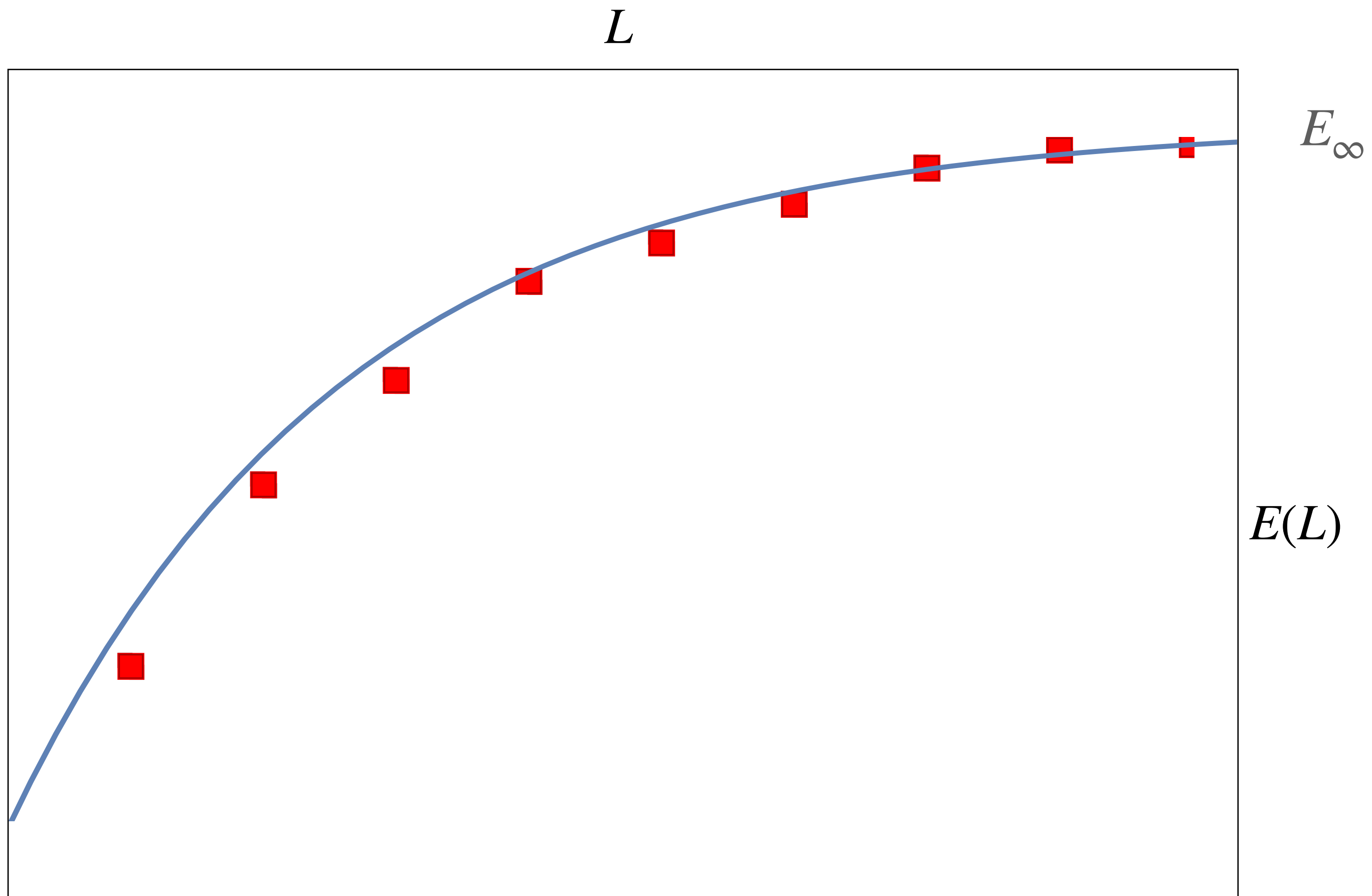


$$e^{-i\delta_0(p)} = S_0(p) \propto \frac{C_0^2}{ip - \kappa} - \frac{1}{a_0^*} + \frac{1}{2}r_0^*(ip - \kappa)^2 + \dots$$

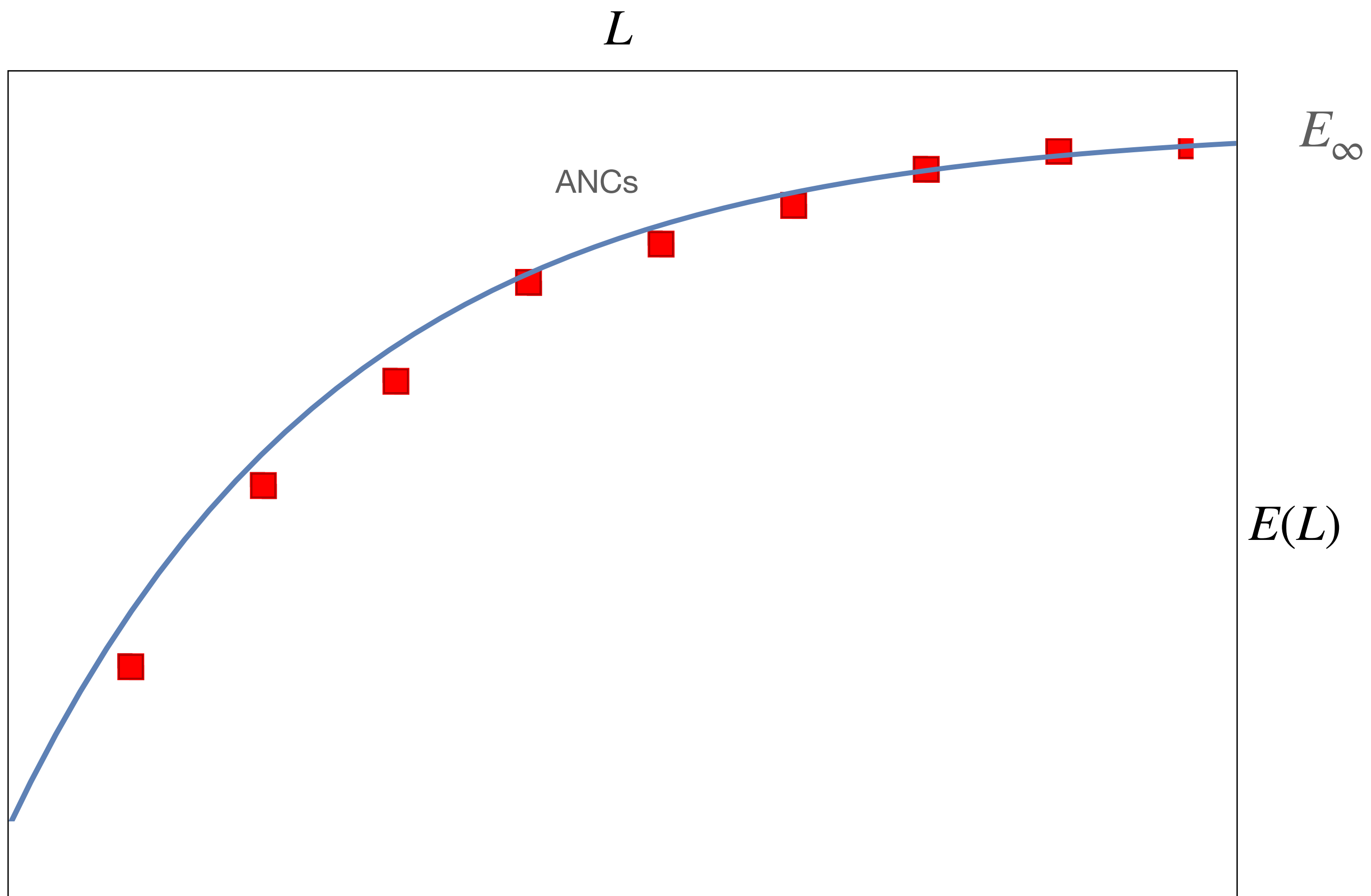
L



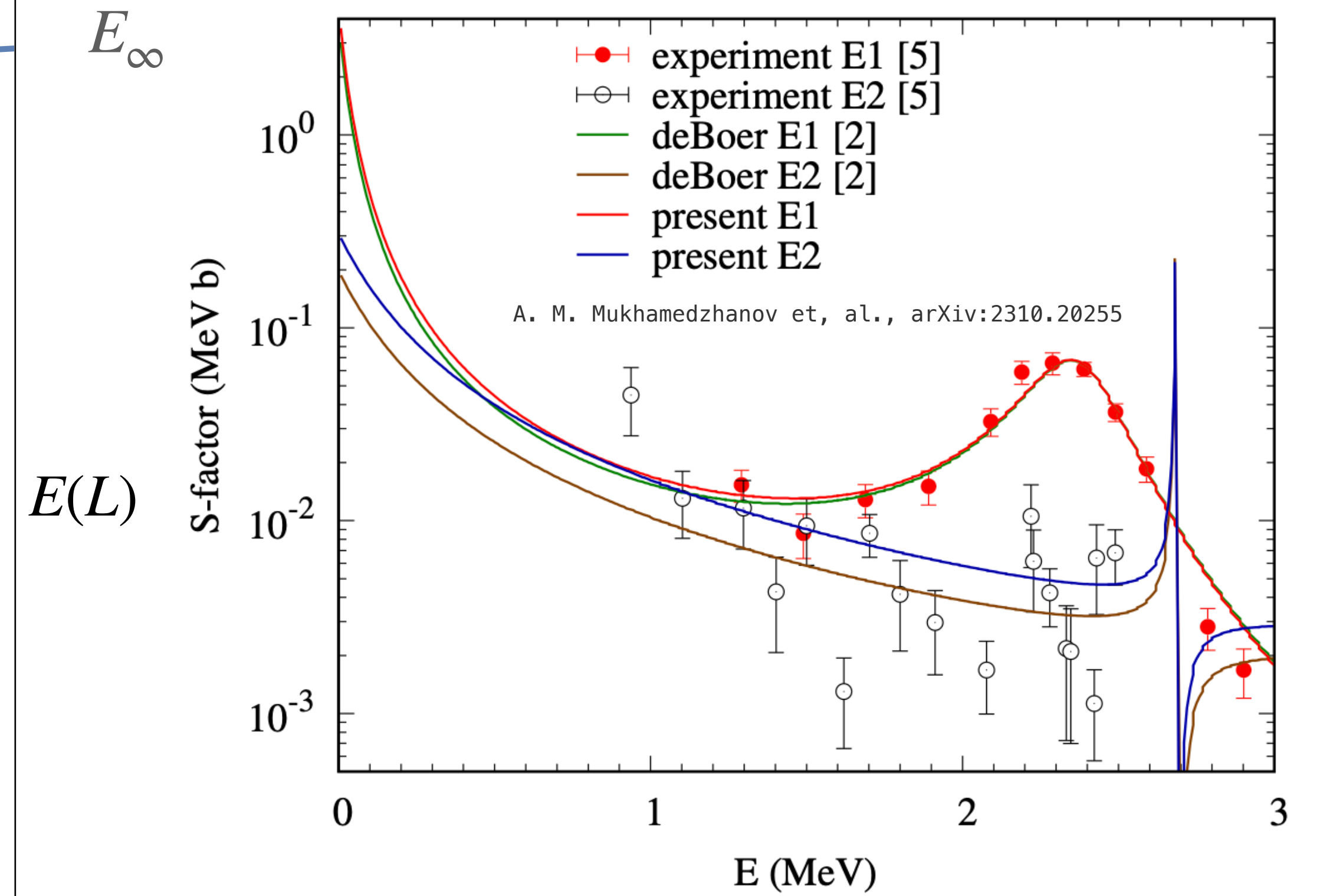
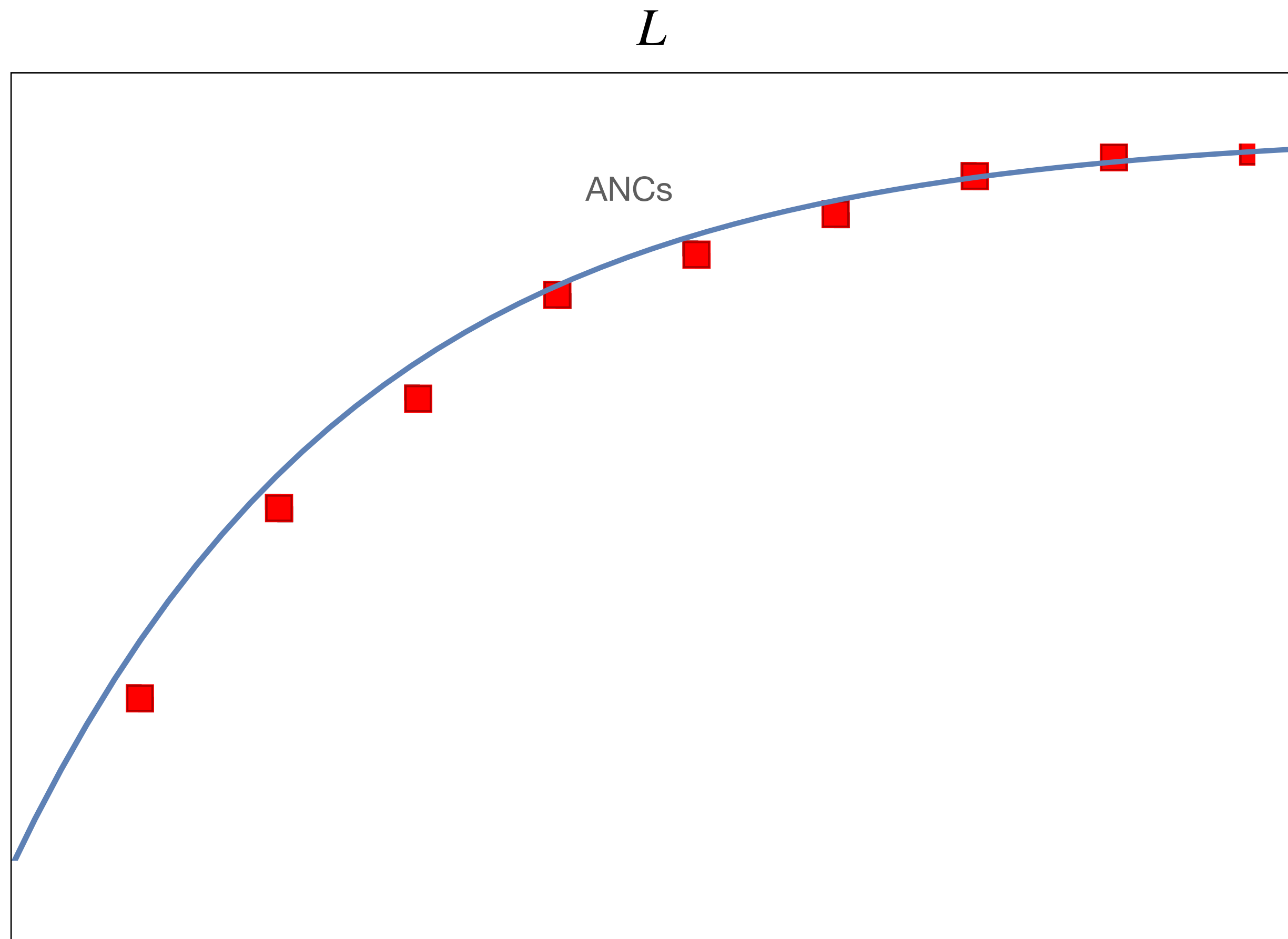
$$e^{-i\delta_0(p)} = S_0(p) \propto \frac{C_0^2}{ip - \kappa} - \frac{1}{a_0^*} + \frac{1}{2}r_0^*(ip - \kappa)^2 + \dots$$



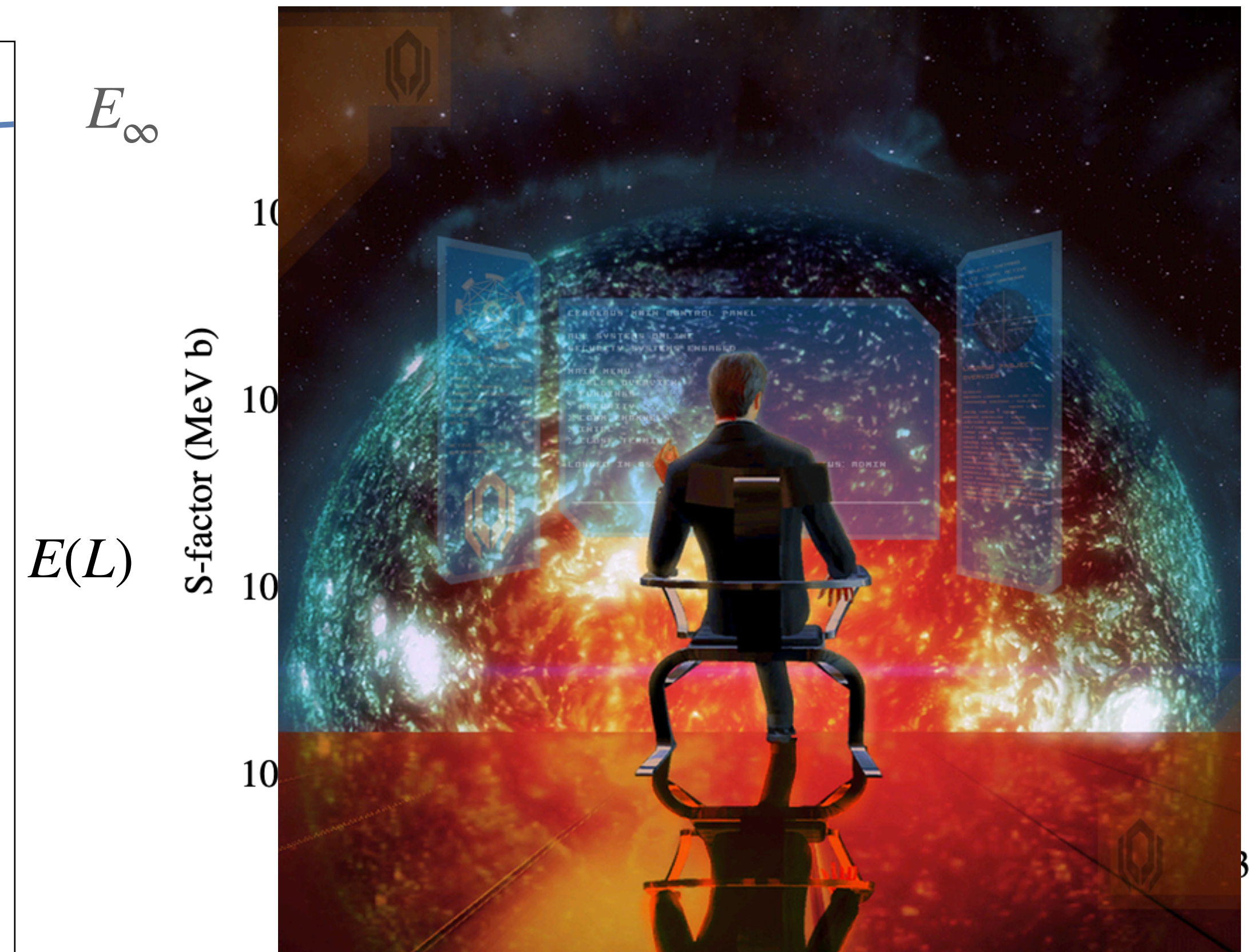
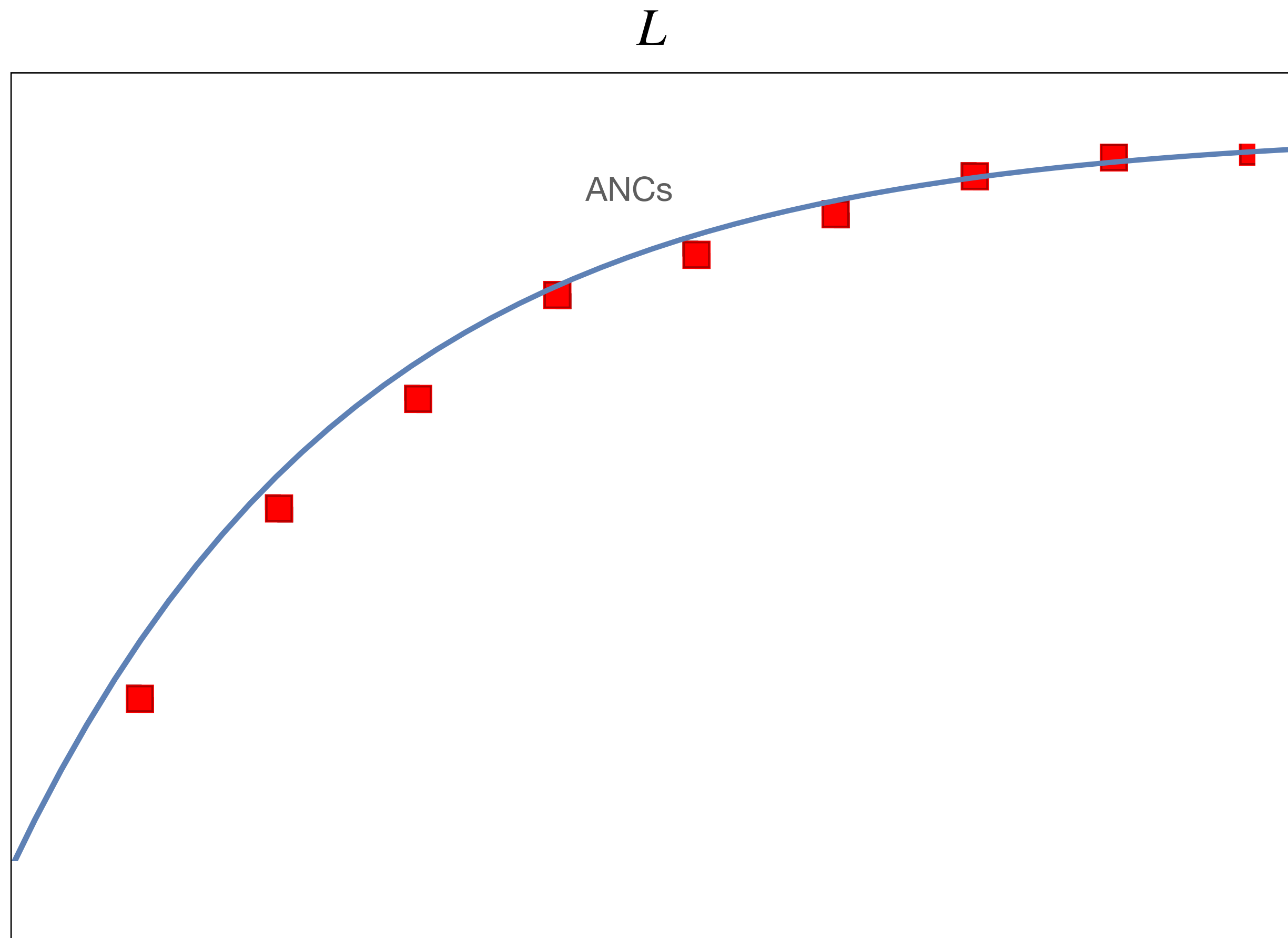
$$e^{-i\delta_0(p)} = S_0(p) \propto \frac{C_0^2}{ip - \kappa} - \frac{1}{a_0^*} + \frac{1}{2}r_0^*(ip - \kappa)^2 + \dots$$



$$e^{-i\delta_0(p)} = S_0(p) \propto \frac{C_0^2}{ip - \kappa} - \frac{1}{a_0^*} + \frac{1}{2}r_0^*(ip - \kappa)^2 + \dots$$



$$e^{-i\delta_0(p)} = S_0(p) \propto \frac{C_0^2}{ip - \kappa} - \frac{1}{a_0^*} + \frac{1}{2}r_0^*(ip - \kappa)^2 + \dots$$



Summary: Why Finite Volume?

- Consequence of lattice regularization

Discreet levels of spectra do not reproduce scattering info

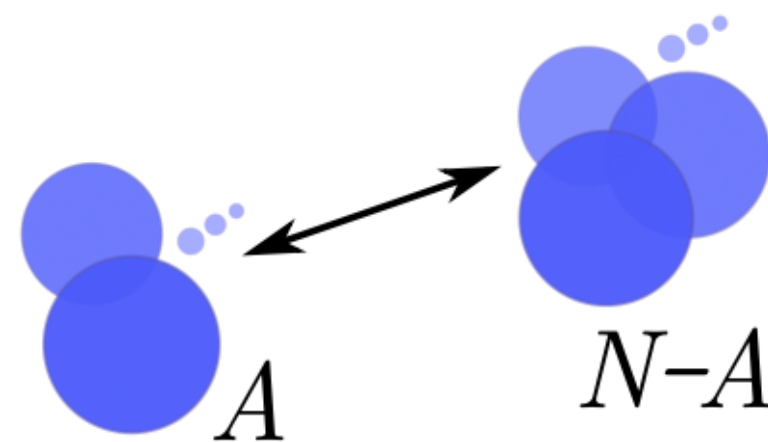
- Scattering matrices

Lüscher, Commun.Math.Phys.**104** 153(1986)

Lüscher, Commun.Math.Phys.**104** 177(1986)

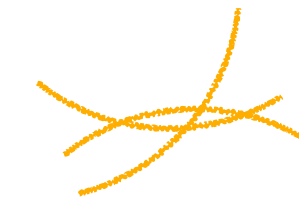
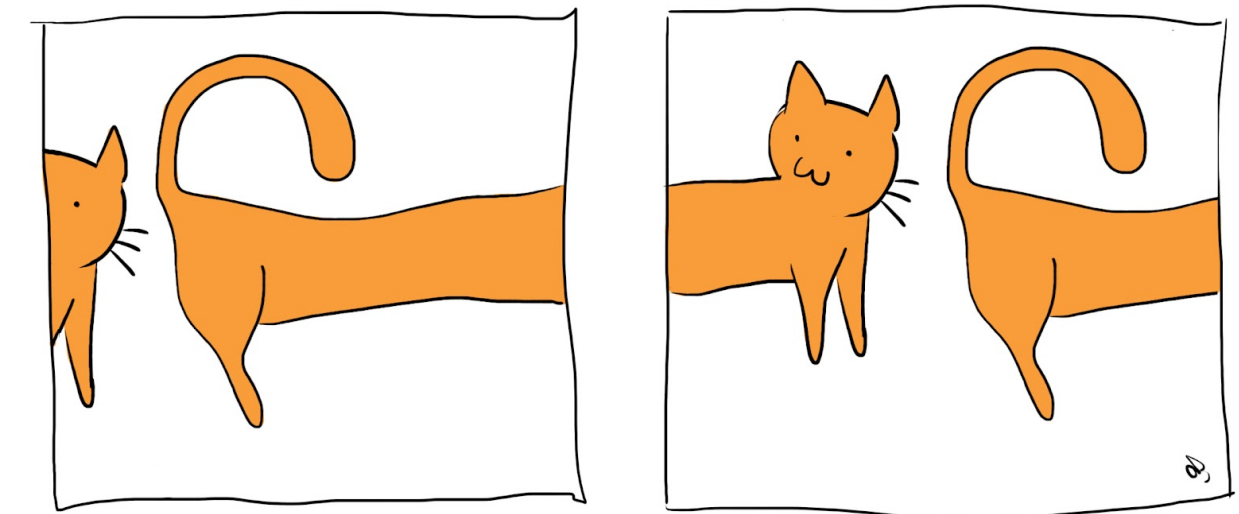
- Nuclear EFT and reaction theory

S.R. Beane et al., PLB 585 106 (2004)
 S. Koenig et al.,PRL, 107 112001(2011),
 Meißner et.al.,PRL, 114 091602(2015).
 S. Koenig , D. Lee, PLB 779 9 (2018),



dingercatadventures.blogspot.com

PERIODIC BOUNDARY CONDITIONS

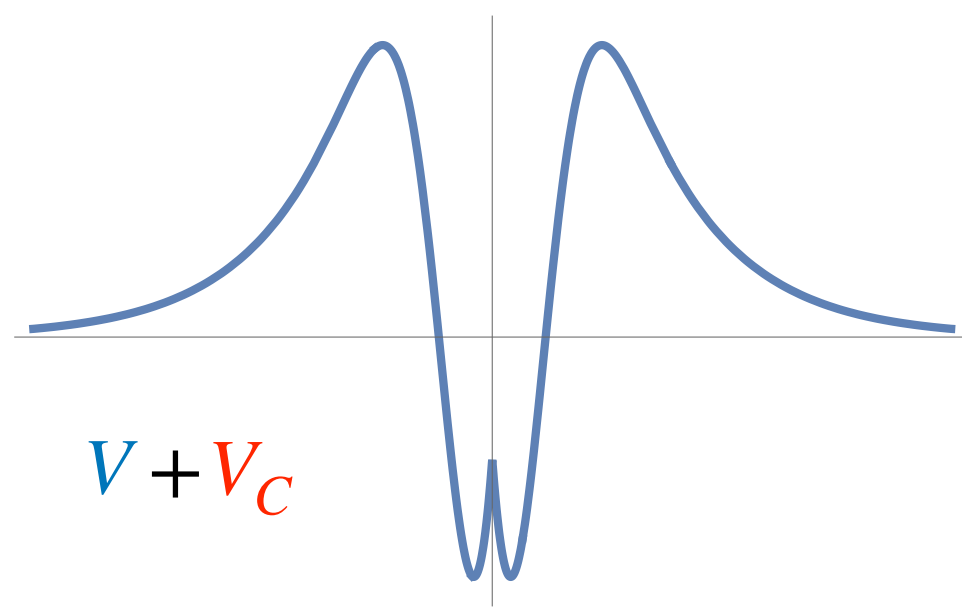


Assumption: Interactions are Short-Ranged

Our setup

HY, S. König, D. Lee, PRL, **131**, 212 502 (2023)

We are considering a non-relativistic system **short-range interactions** + **repulsive Coulomb interactions**



With COM frame coordinates:

$$\mathbf{r} = \mathbf{x}_1 - \mathbf{x}_2$$

$$V_C(r) = \frac{\gamma}{2\mu r}, \quad \gamma = 2\mu\alpha Z_1 Z_2 > 0$$

$$H = -\frac{1}{2\mu} \nabla^2 + V + V_C$$

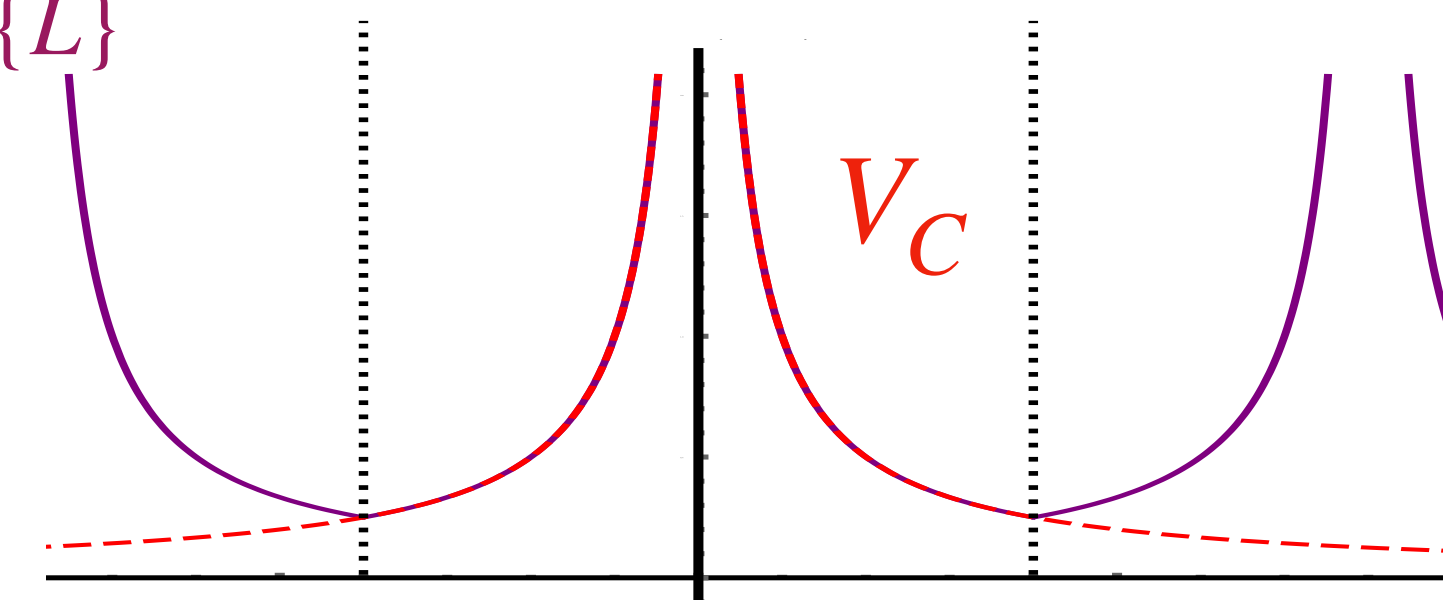
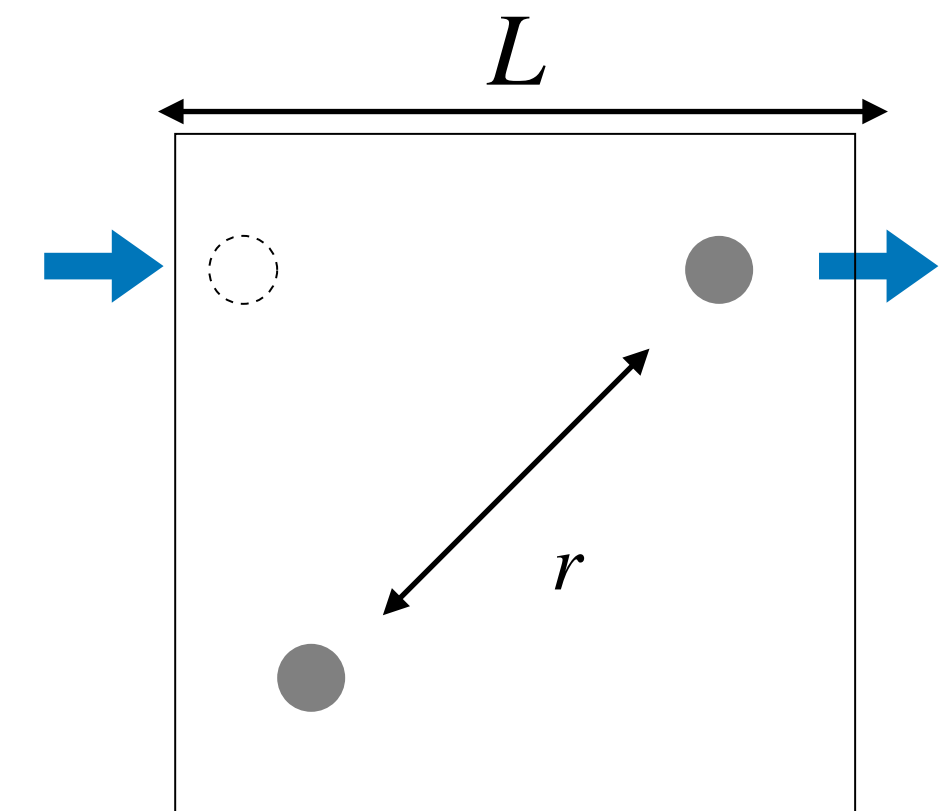
$$V_L(x) = \sum_{\mathbf{n}} V(x + \mathbf{n}L)$$

$$H_L = -\frac{1}{2\mu} \nabla^2 + V_L + V_{C,\{L\}}$$

$$\psi_\infty(r) \rightarrow \frac{C_0}{\sqrt{4\pi r}} W_{-\bar{\eta}, 1/2}(2\kappa r)$$

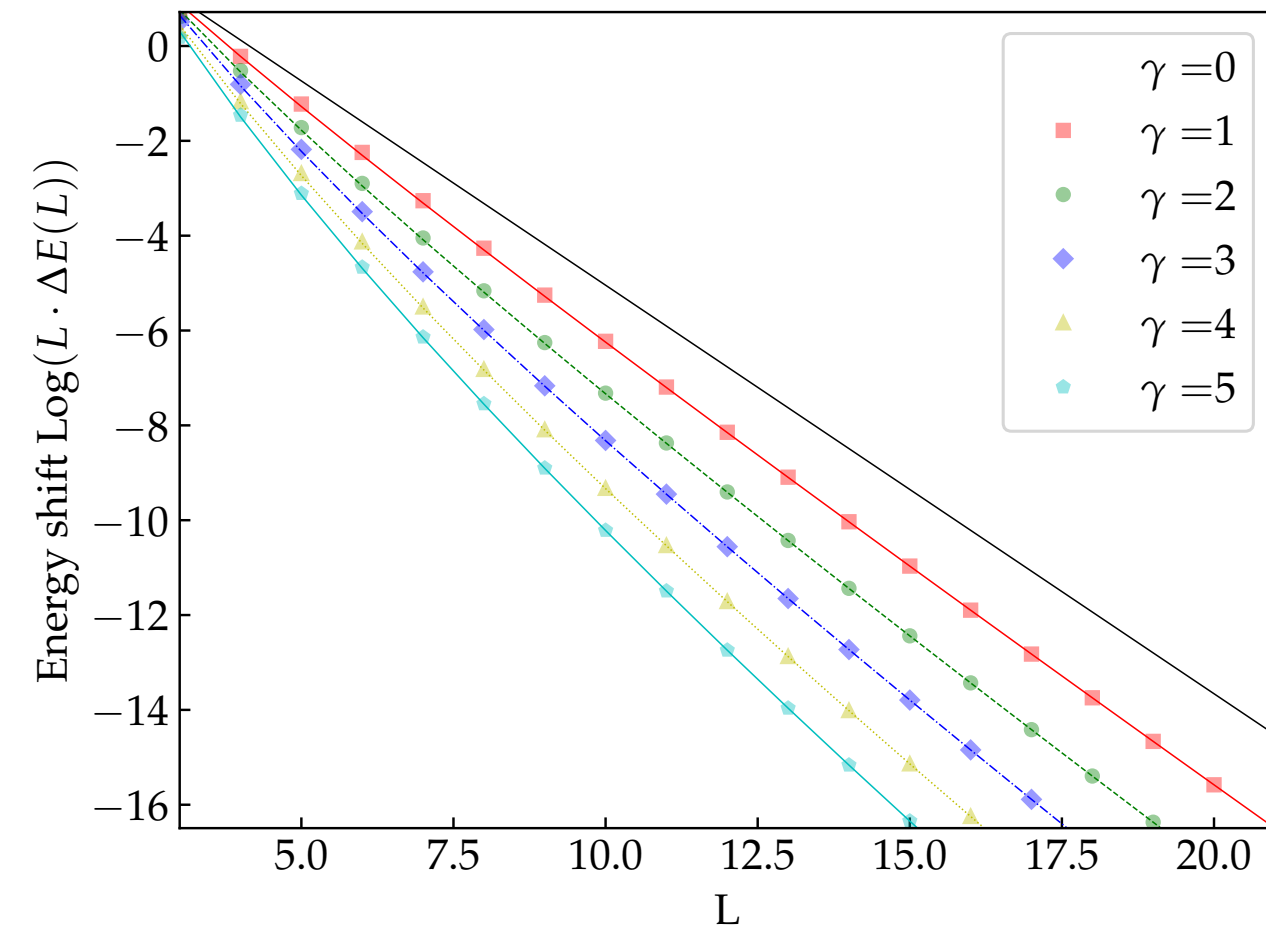
$$\bar{\eta} = \frac{\gamma}{2\kappa}$$

In finite volume

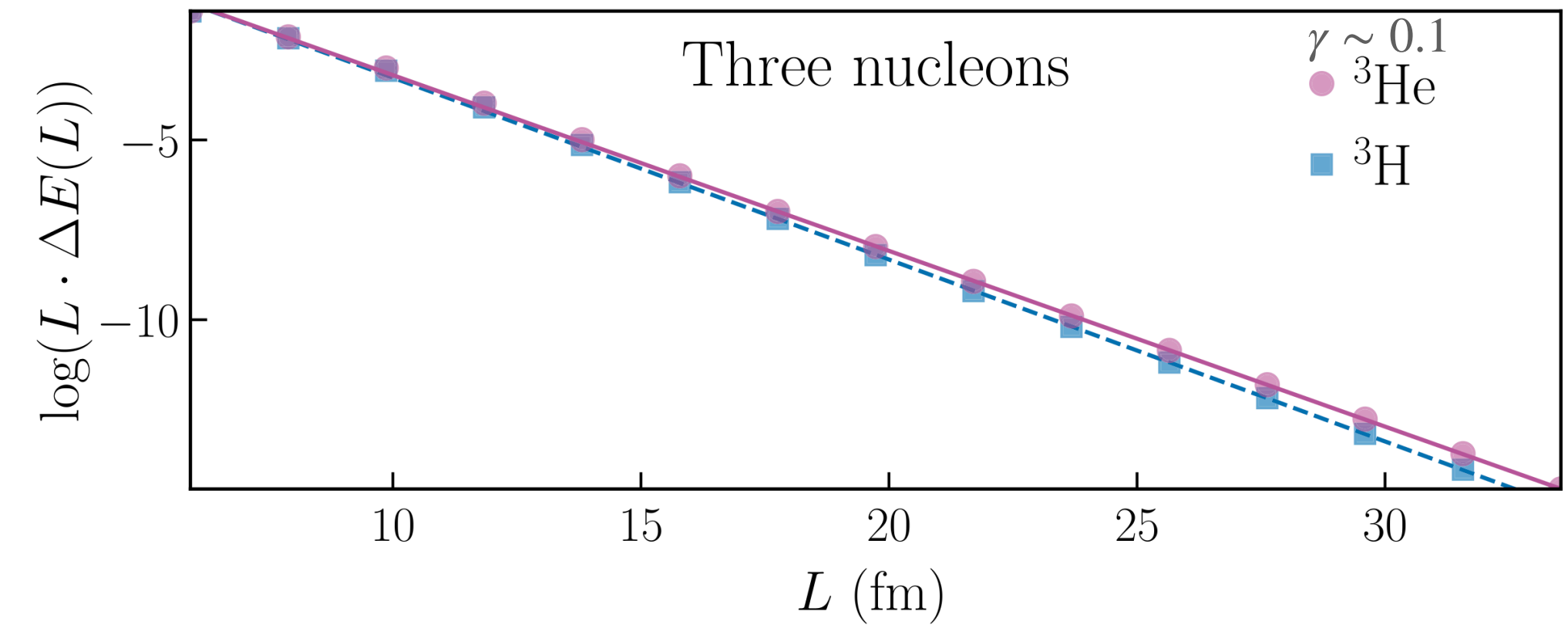


*P. Guo, PRC 103 064611 (2021),
P. Guo et. al., PRD 103 094520 (2021).*

Main result



γ	Finite-volume fit			Continuum result	
	κ_∞	C_0	L range	κ_∞	C_0
1.0	0.8610(3)	5.039(2)	17–28	0.861	5.049
2.0	0.8607(3)	11.71(4)	15–26	0.860	11.79
3.0	0.8605(7)	29.95(20)	14–24	0.859	30.31
4.0	0.8604(1)	83.14(10)	14–22	0.858	84.76
5.0	0.8604(2)	247.9(5)	14–18	0.857	255.4

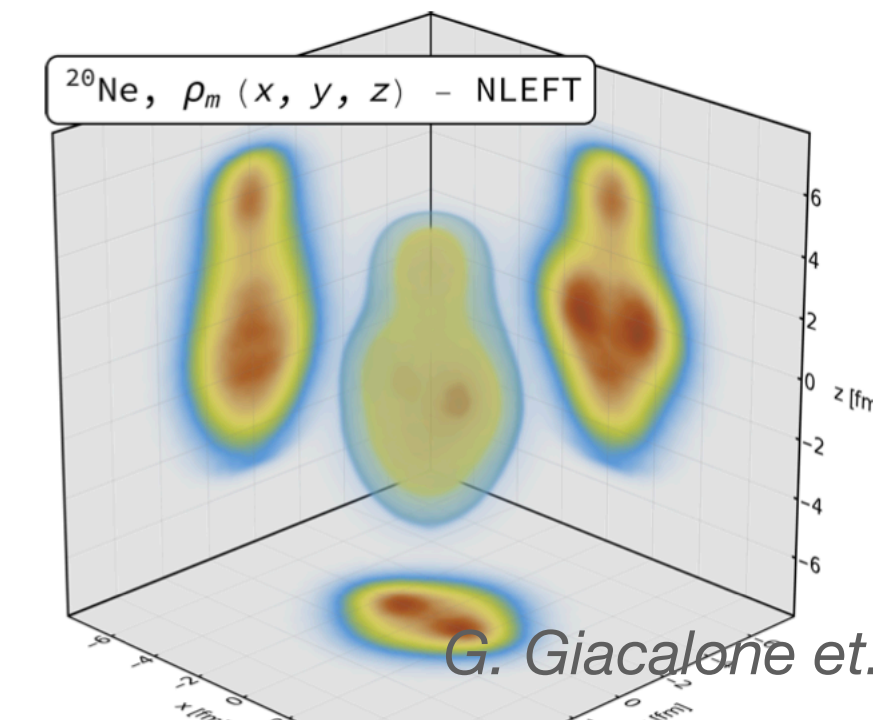
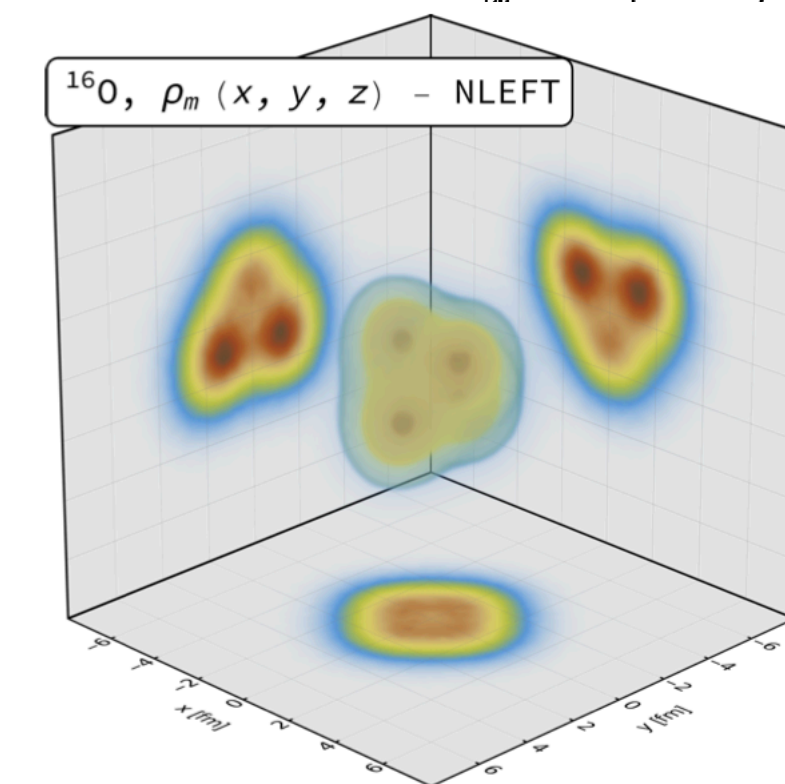
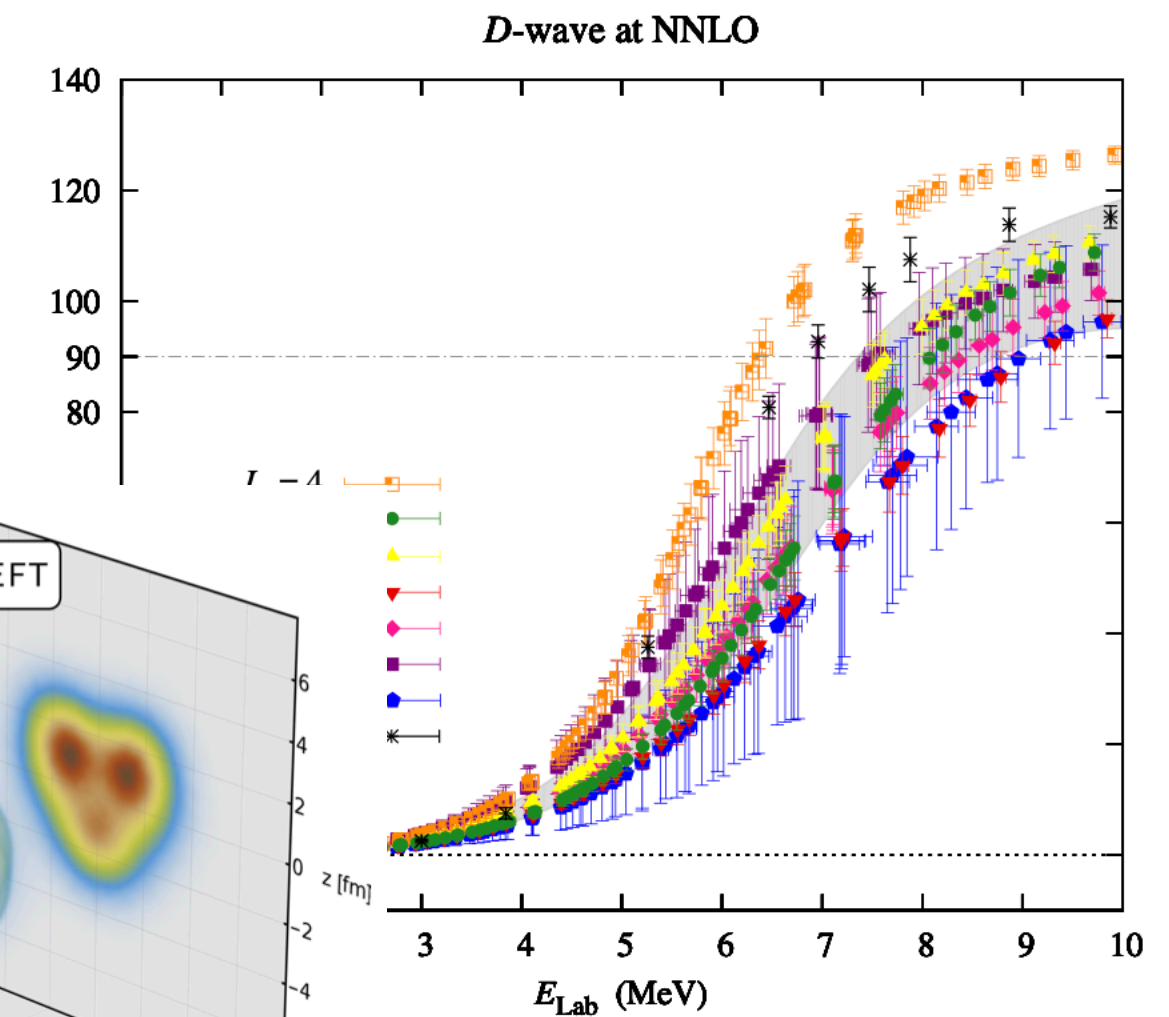
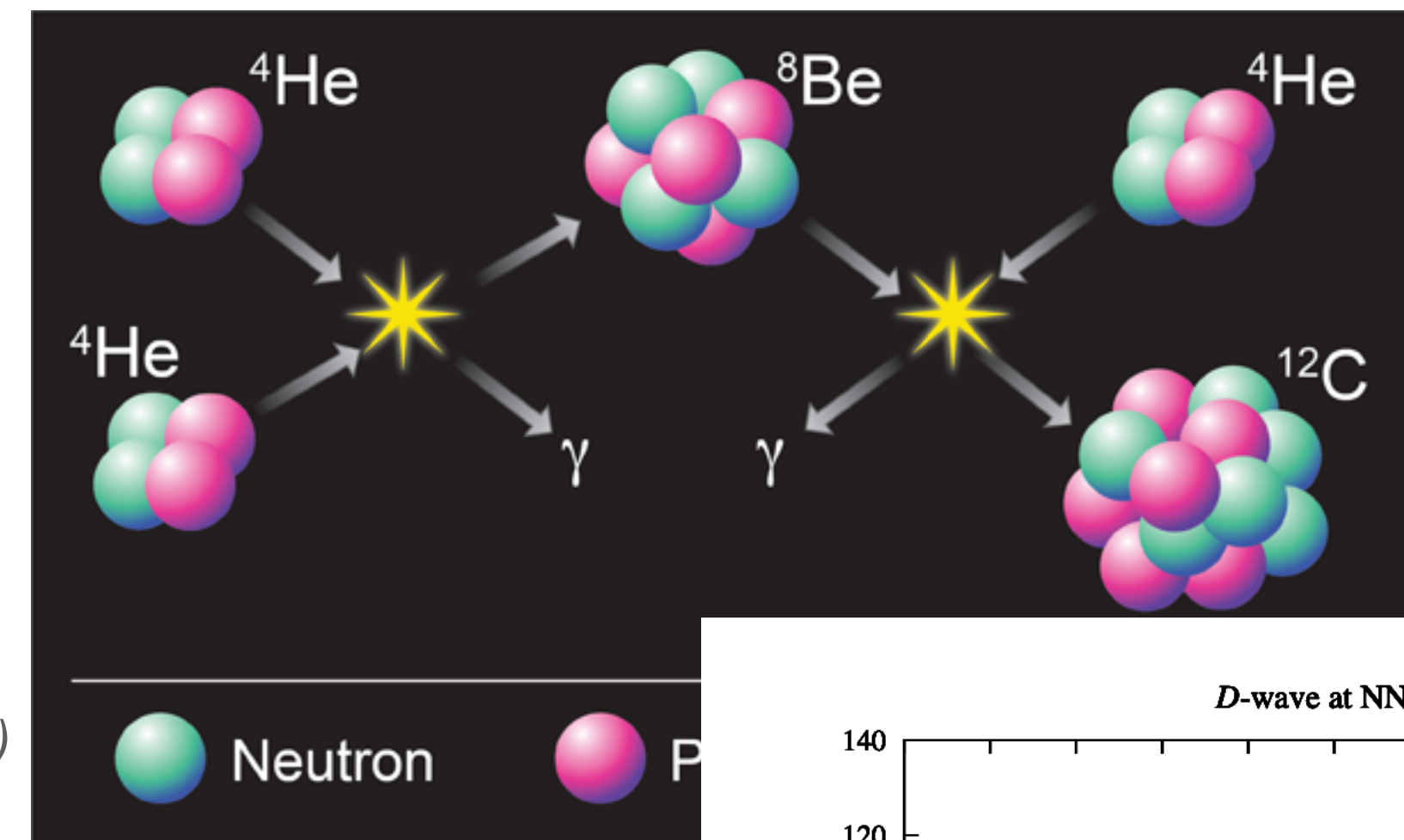


$$\Delta E(L) = -\frac{3C_0^2}{\mu L} \left[W_{-\bar{\eta}, \frac{1}{2}}(\kappa L) \right]^2 + \Delta \tilde{E}(L) + \Delta \tilde{E}'(L) + \mathcal{O}(e^{-\sqrt{2}\kappa L})$$

HY, S. König, D. Lee, PRL, **131**, 212 502 (2023)

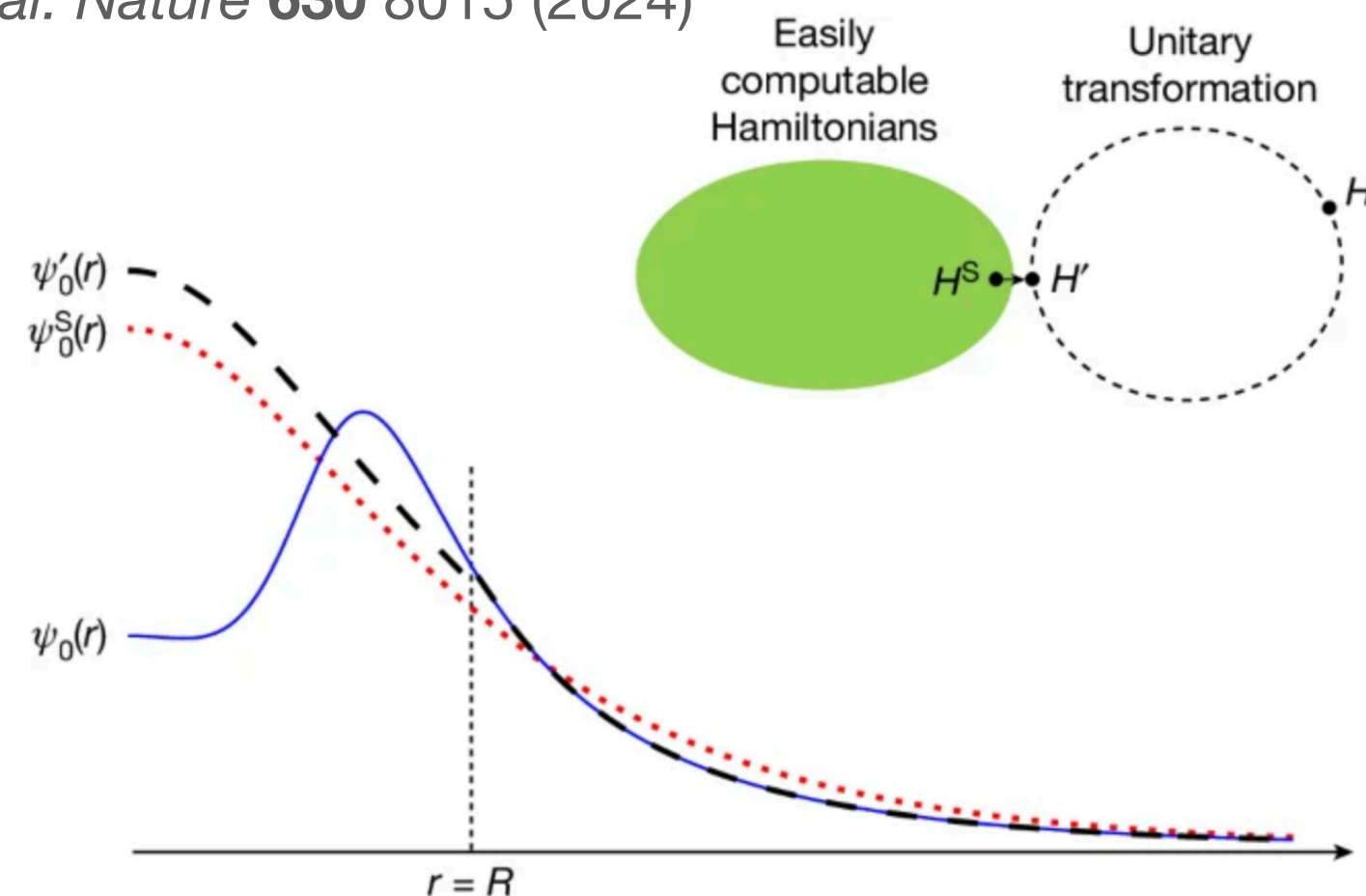
Milestones of NLEFT

- First ab initio Calculation of ^{12}C Hoyle state
E. Epelbaum, et. al. PRL 106 192501 (2011)
- First ab initio alpha - alpha scattering
S. Elhatisari, et. al. Nature 528 111 (2015)
- Accessing correlated densities
S. Elhatisari, et. al. PRL119, 222505 (2017)
- Now: Free from sign problem
 (Wave function matching + $\text{SU}(4)$ invariant interactions)
S. Elhatisari, et. al. Nature 630 8015 (2024)



B. Lu, et. al. PLB, 797 134863(2019)

S. Elhatisari, et. al. Nature 630 8015 (2024)

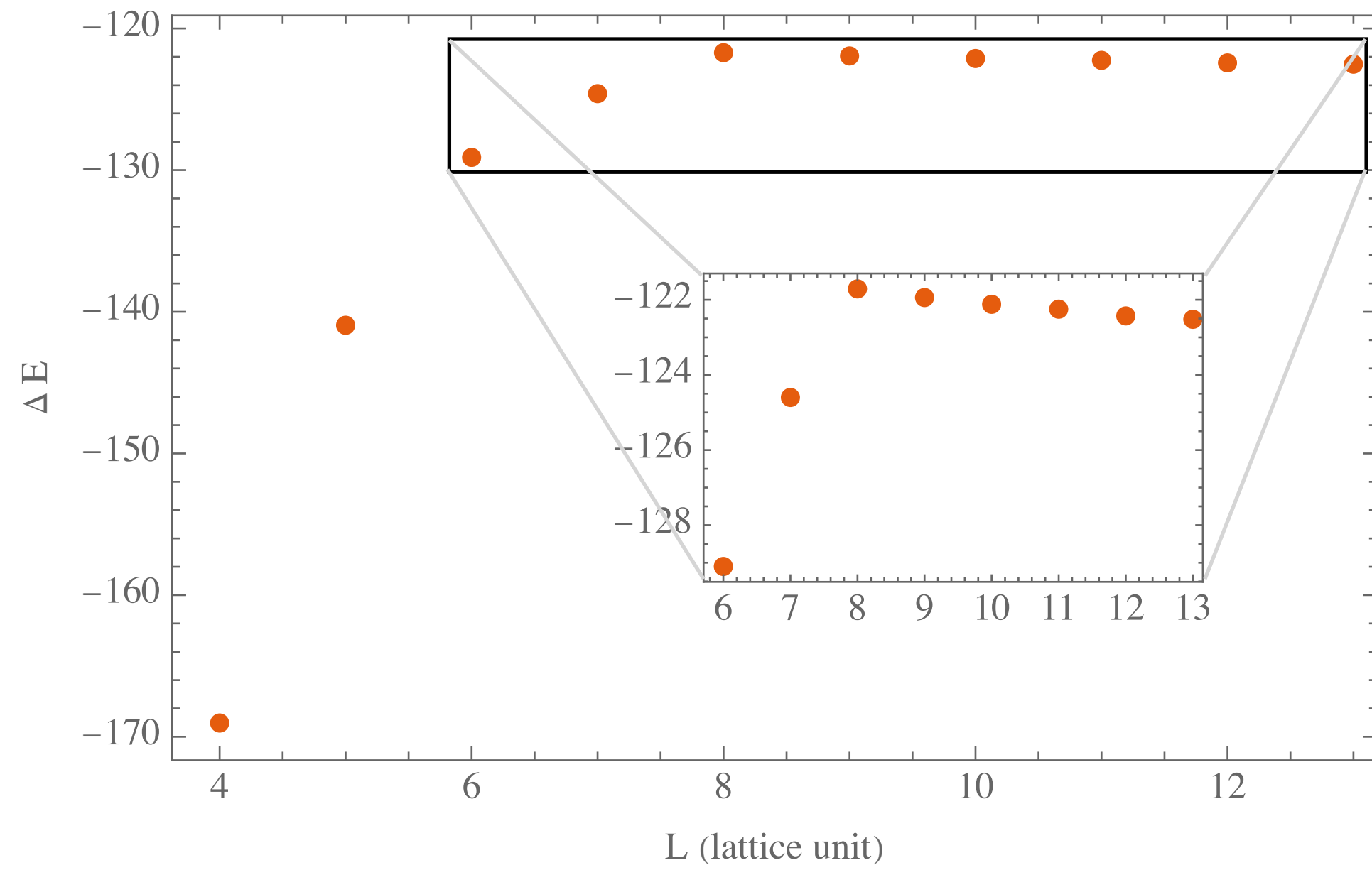


160: a Multi-Channel Case

Prelim results & PhD Thesis

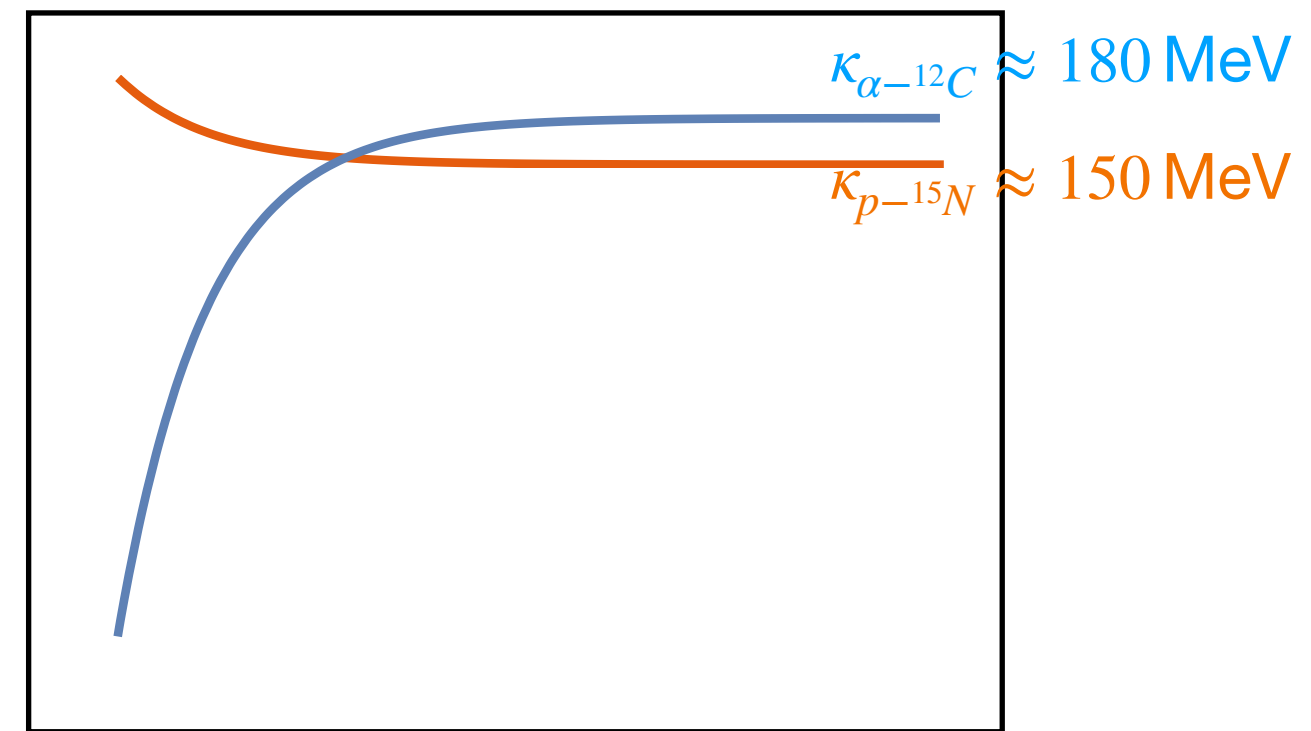
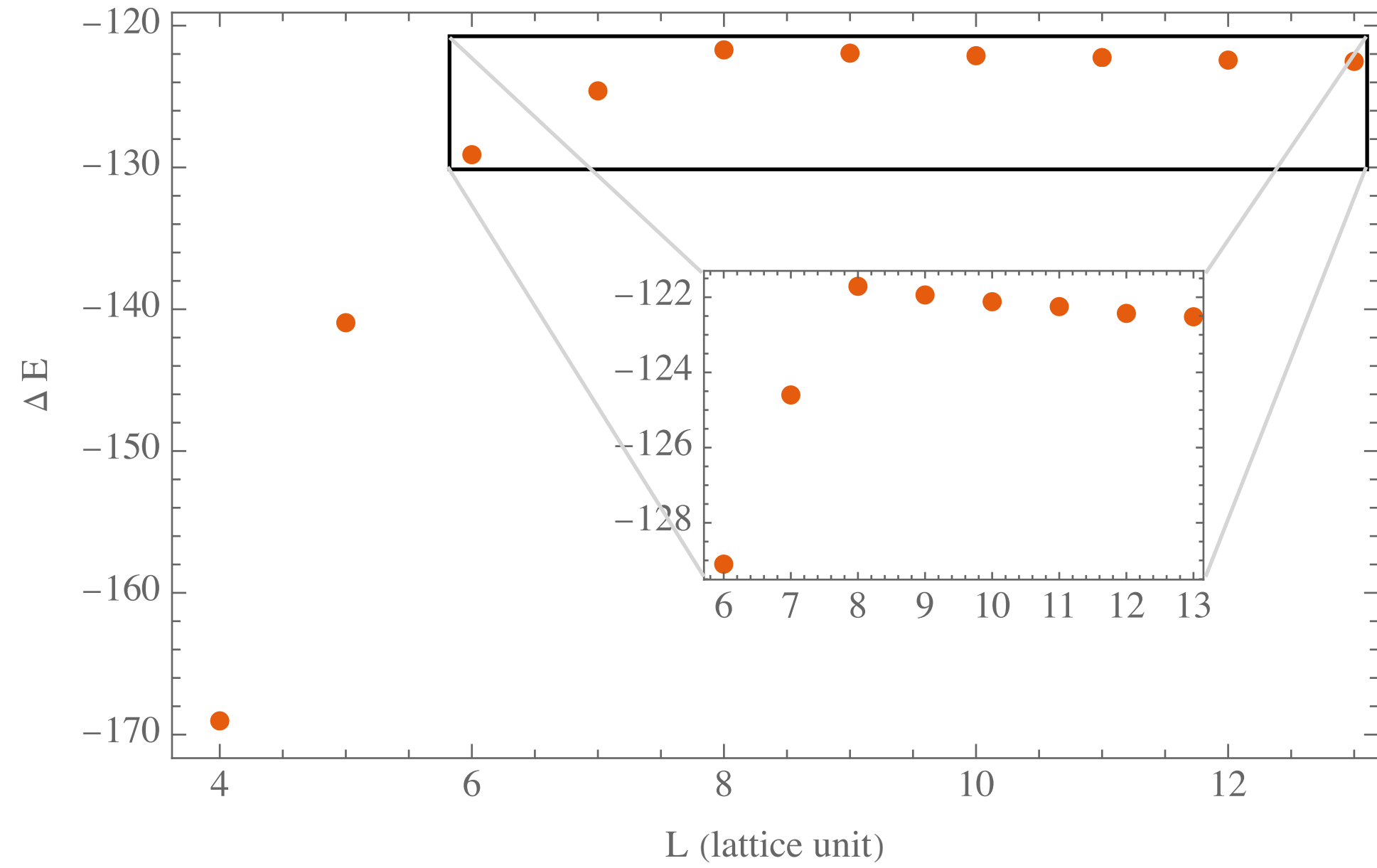
16O: a Multi-Channel Case

Prelim results & PhD Thesis



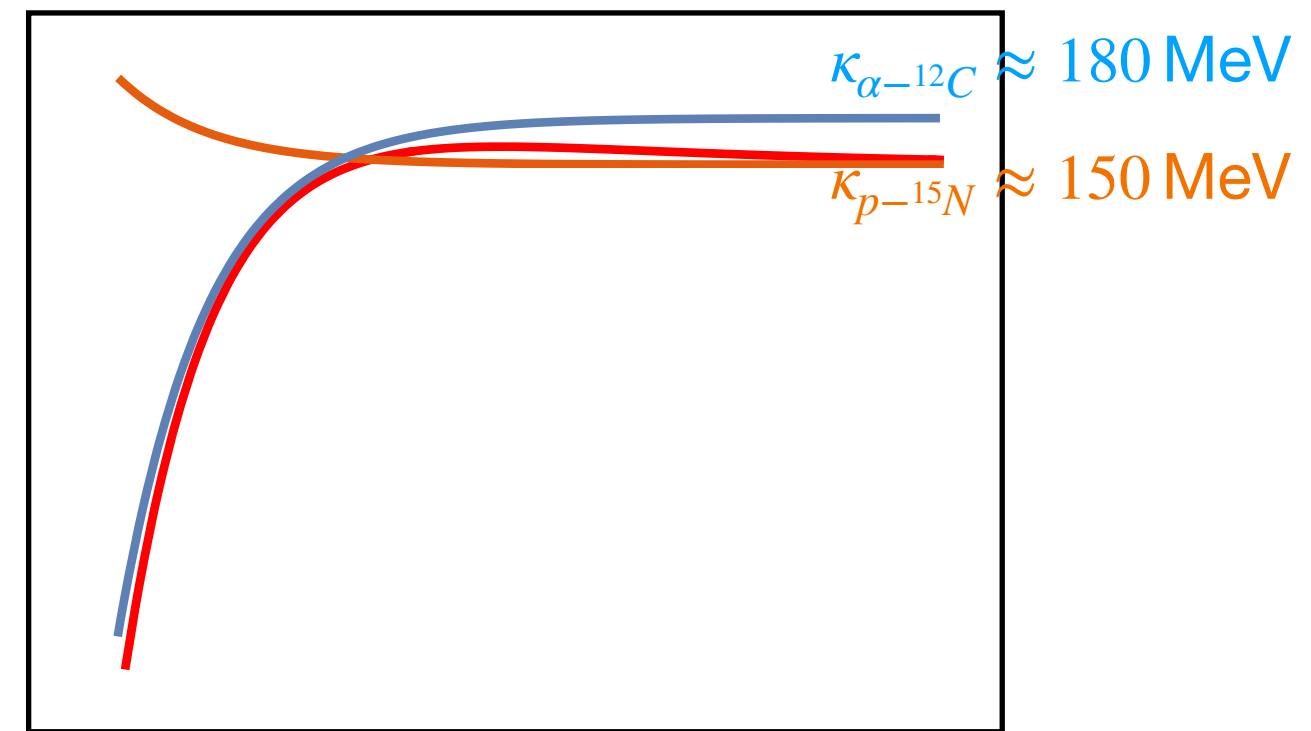
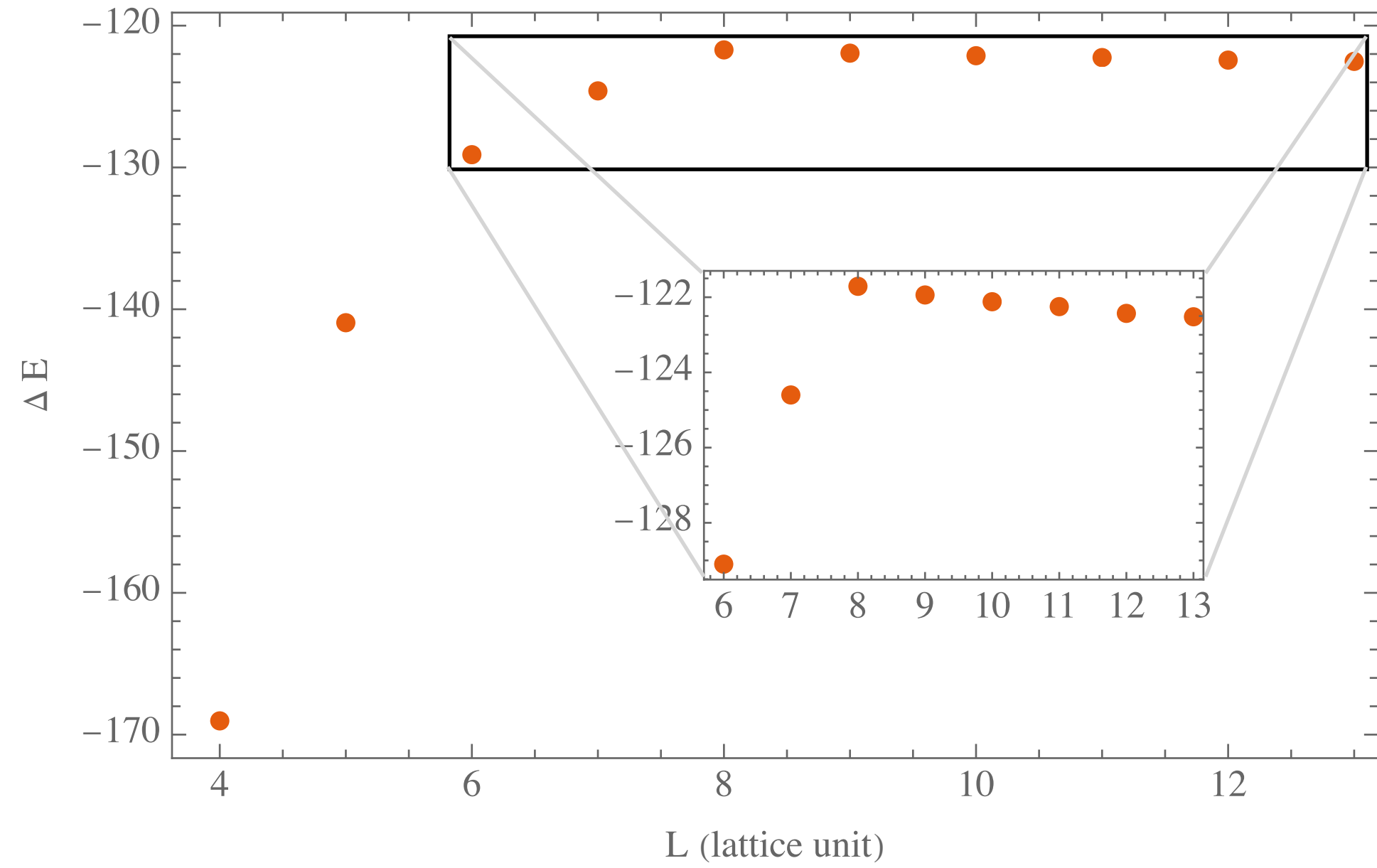
16O: a Multi-Channel Case

Prelim results & PhD Thesis



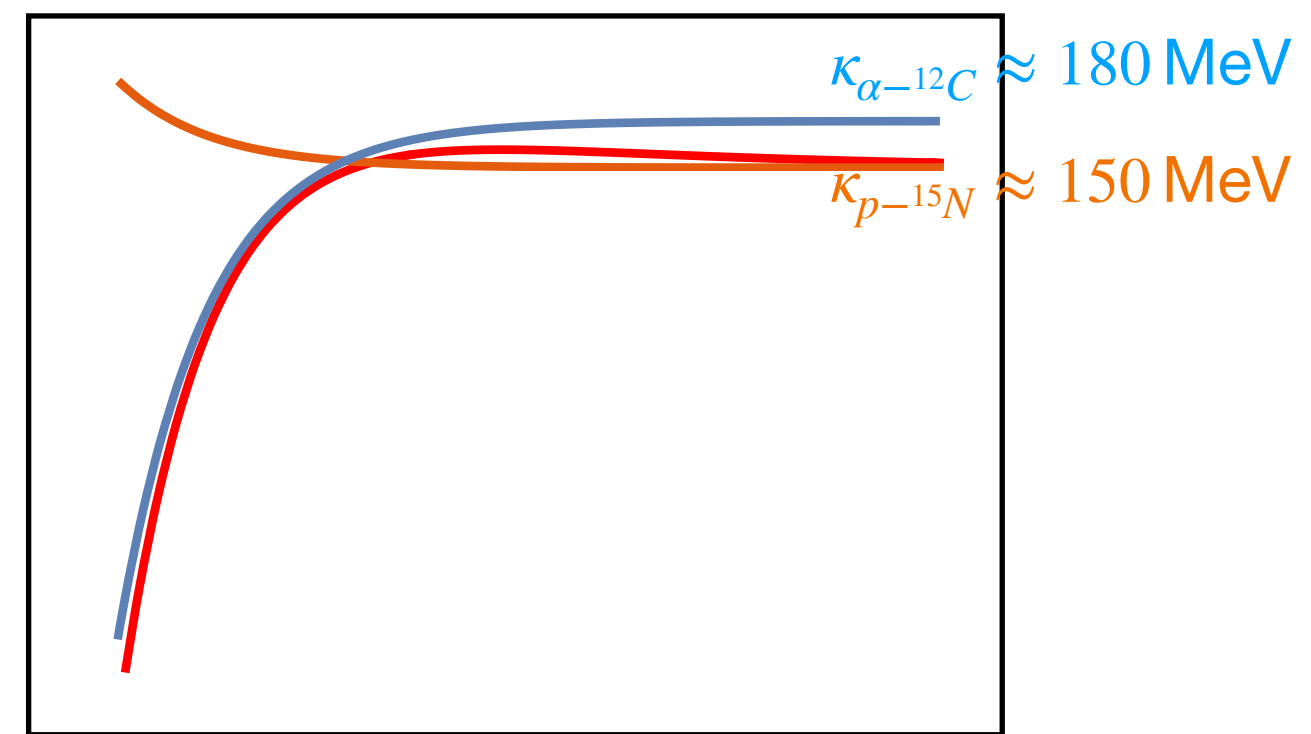
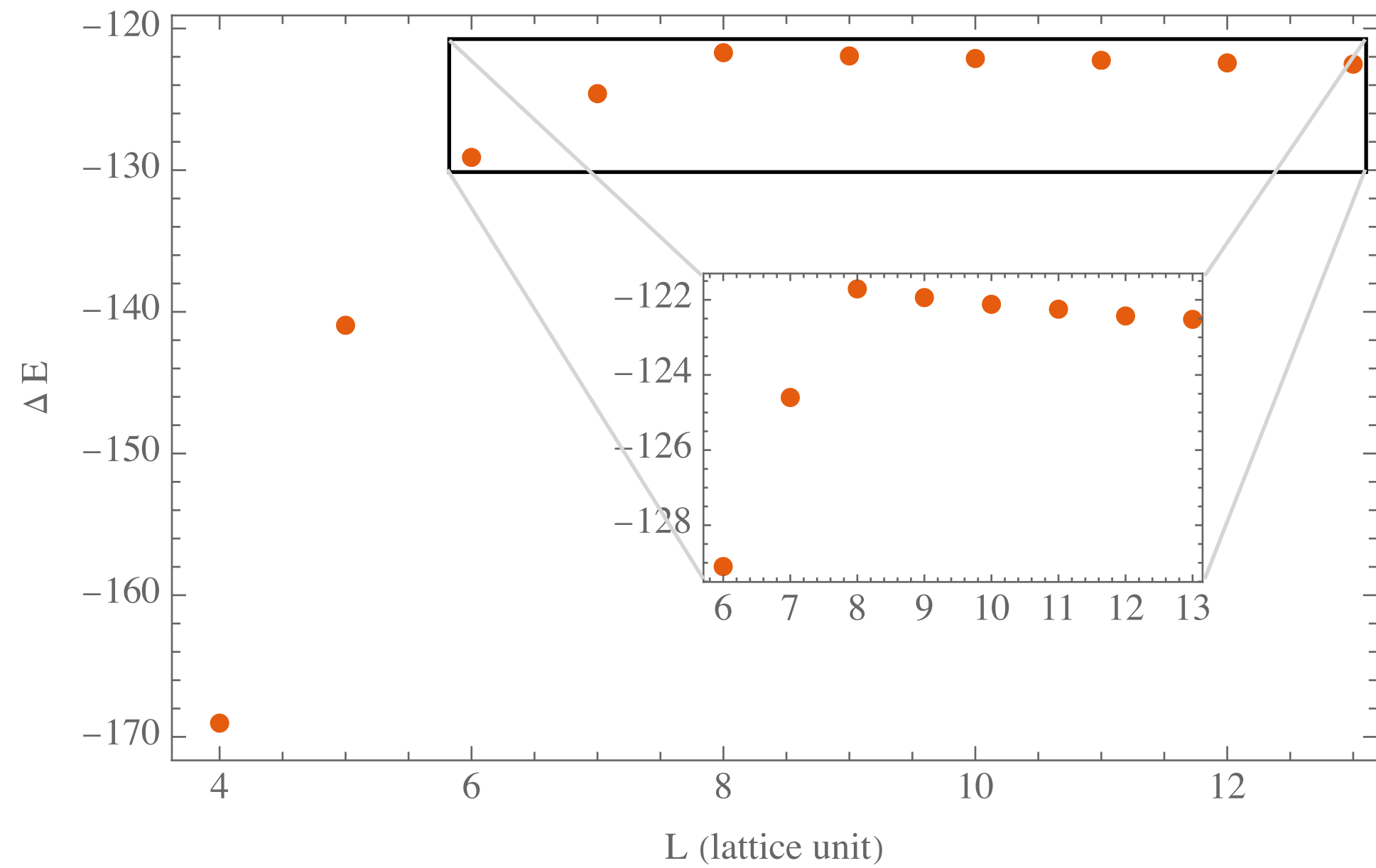
16O: a Multi-Channel Case

Prelim results & PhD Thesis



16O: a Multi-Channel Case

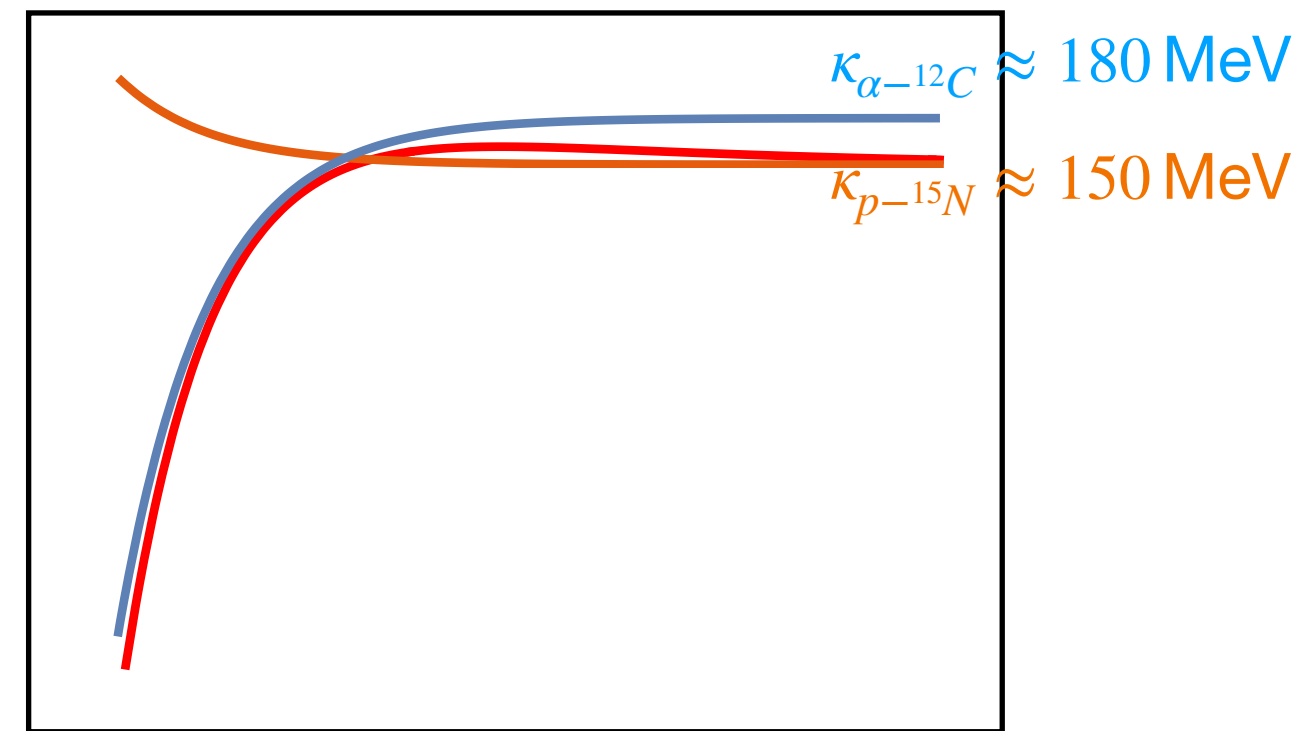
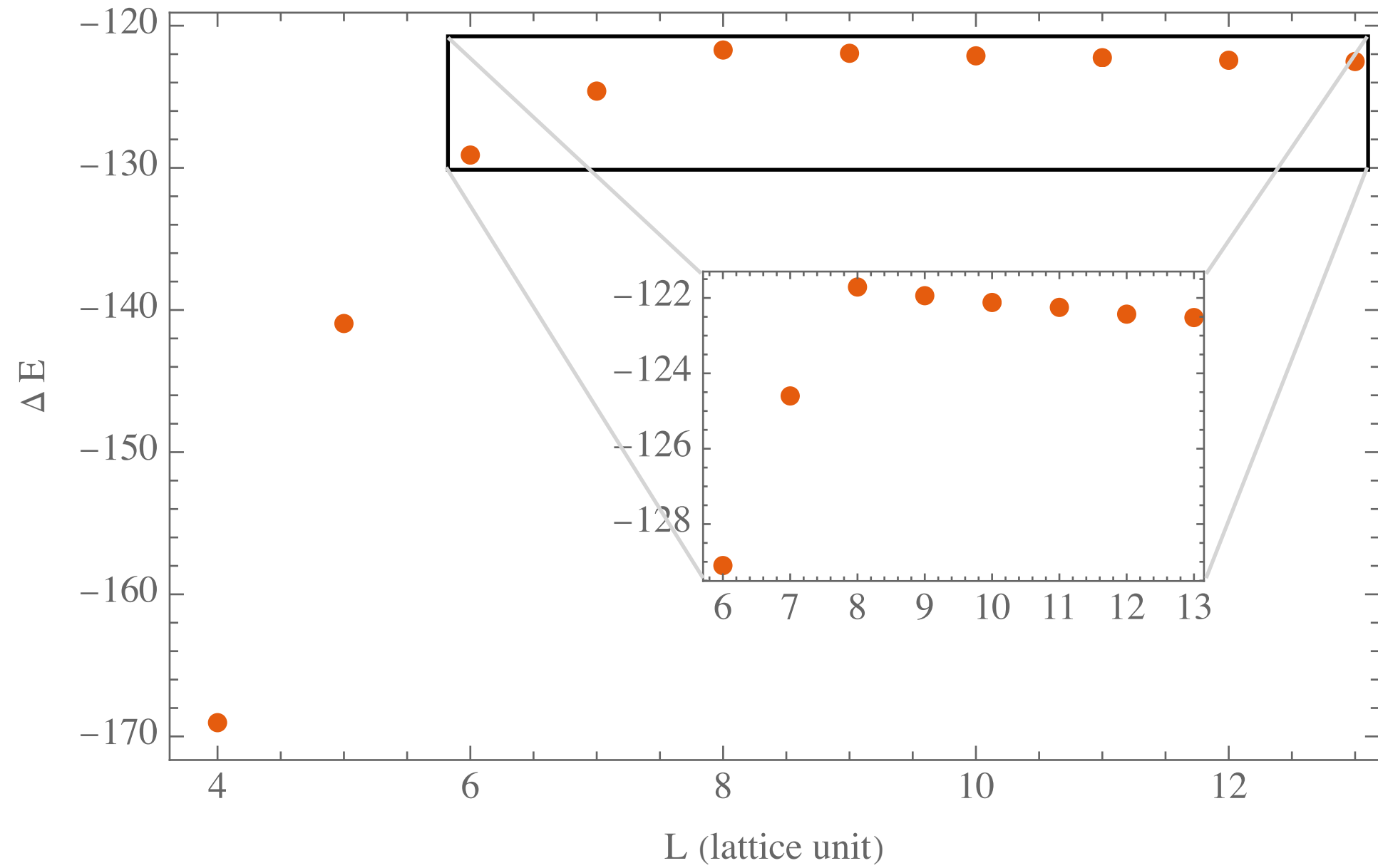
Prelim results & PhD Thesis



10% difference in Binding Energy!

16O: a Multi-Channel Case

Prelim results & PhD Thesis

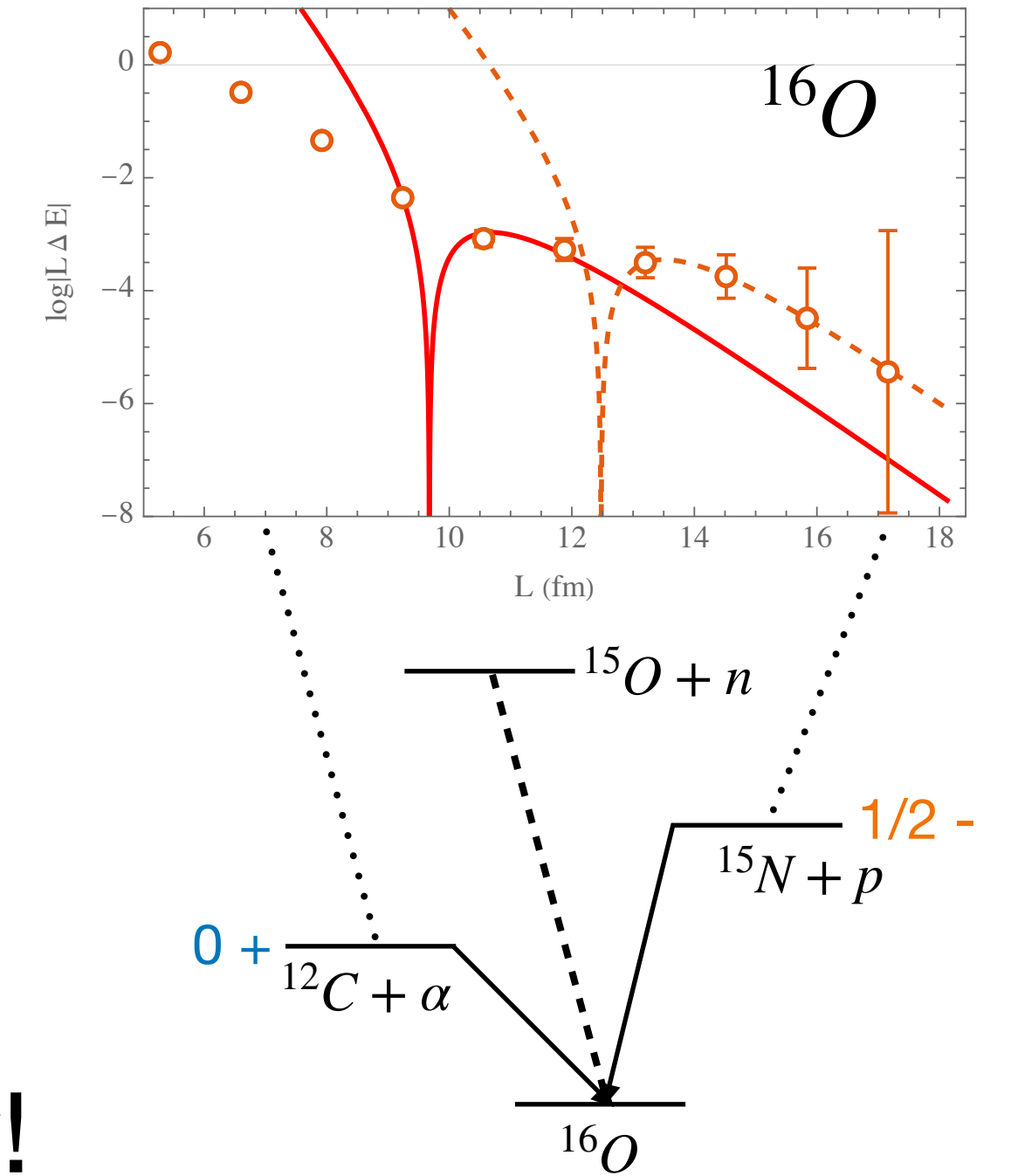
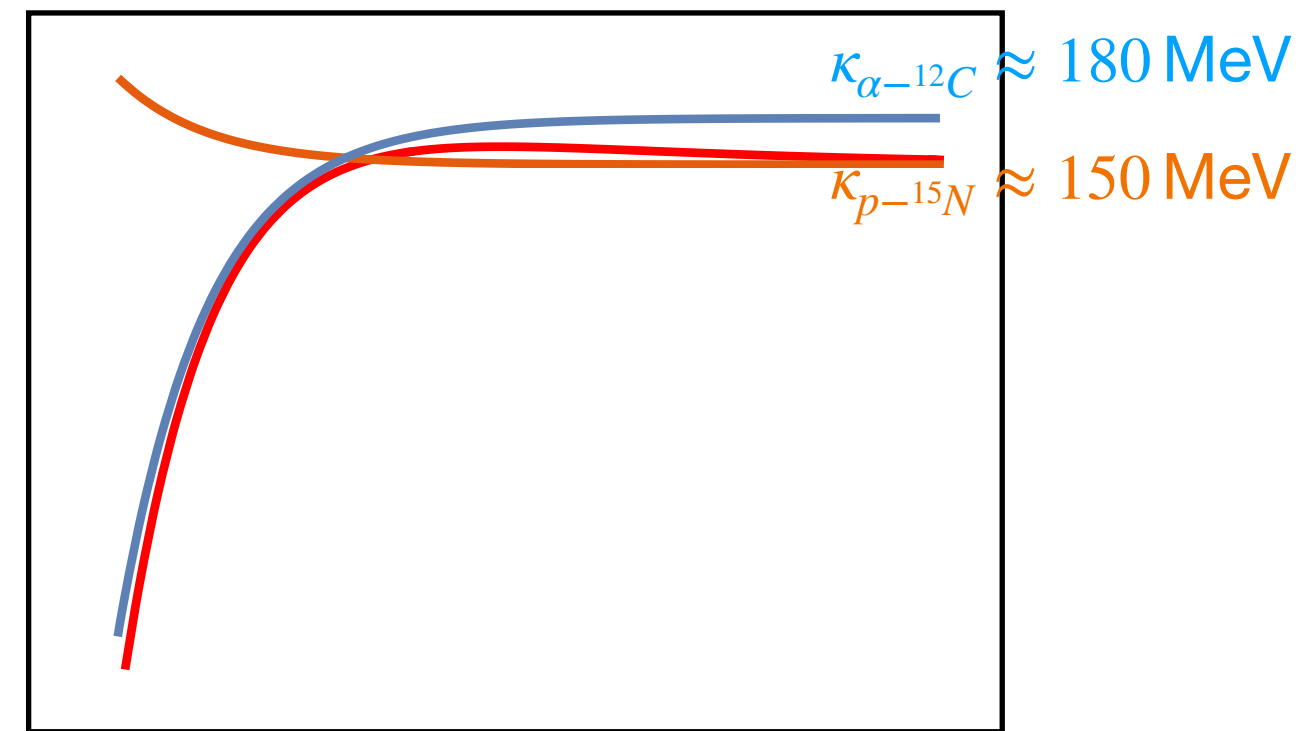
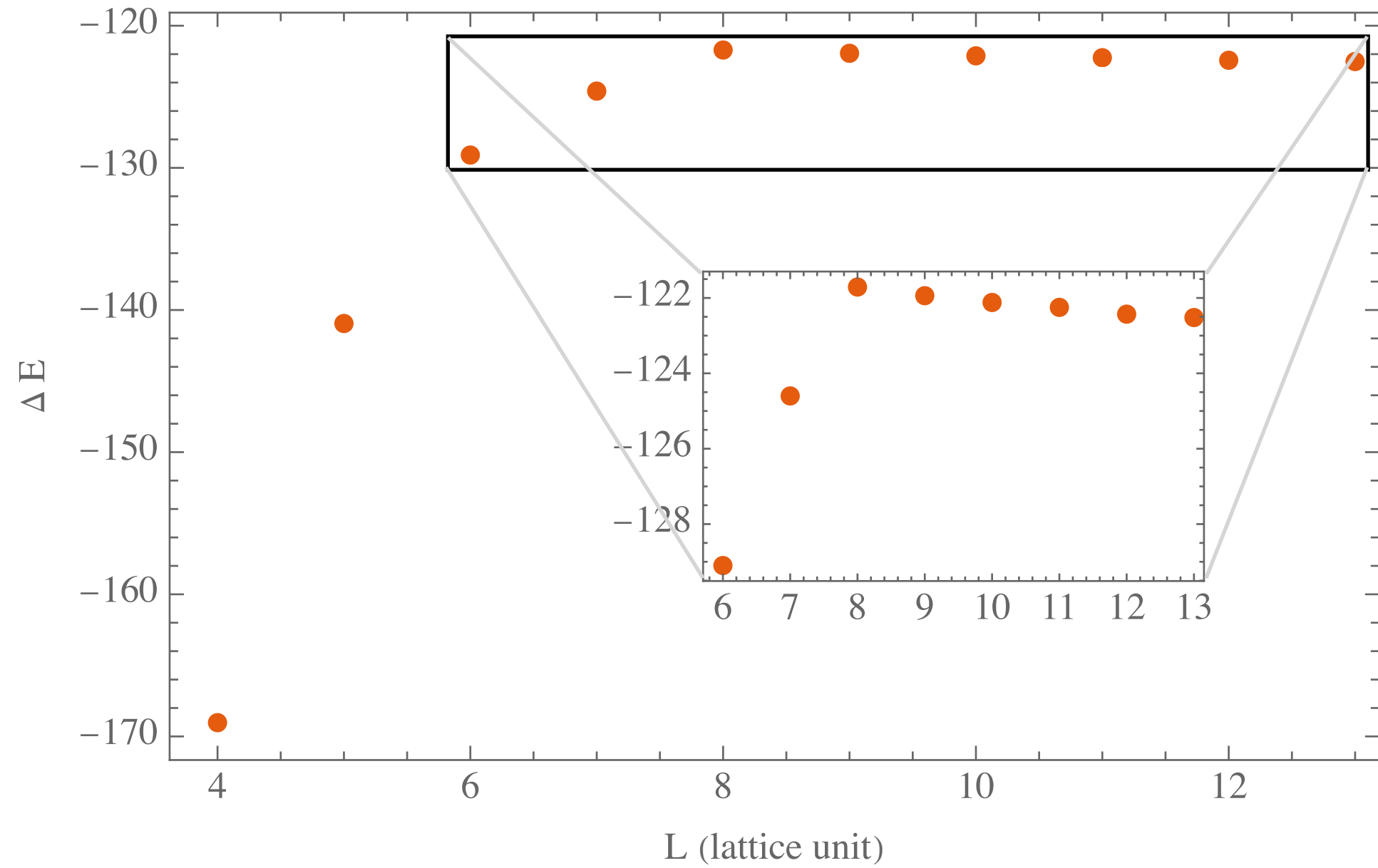


10% difference in Binding Energy!

Problem: overfitting

16O: a Multi-Channel Case

Prelim results & PhD Thesis

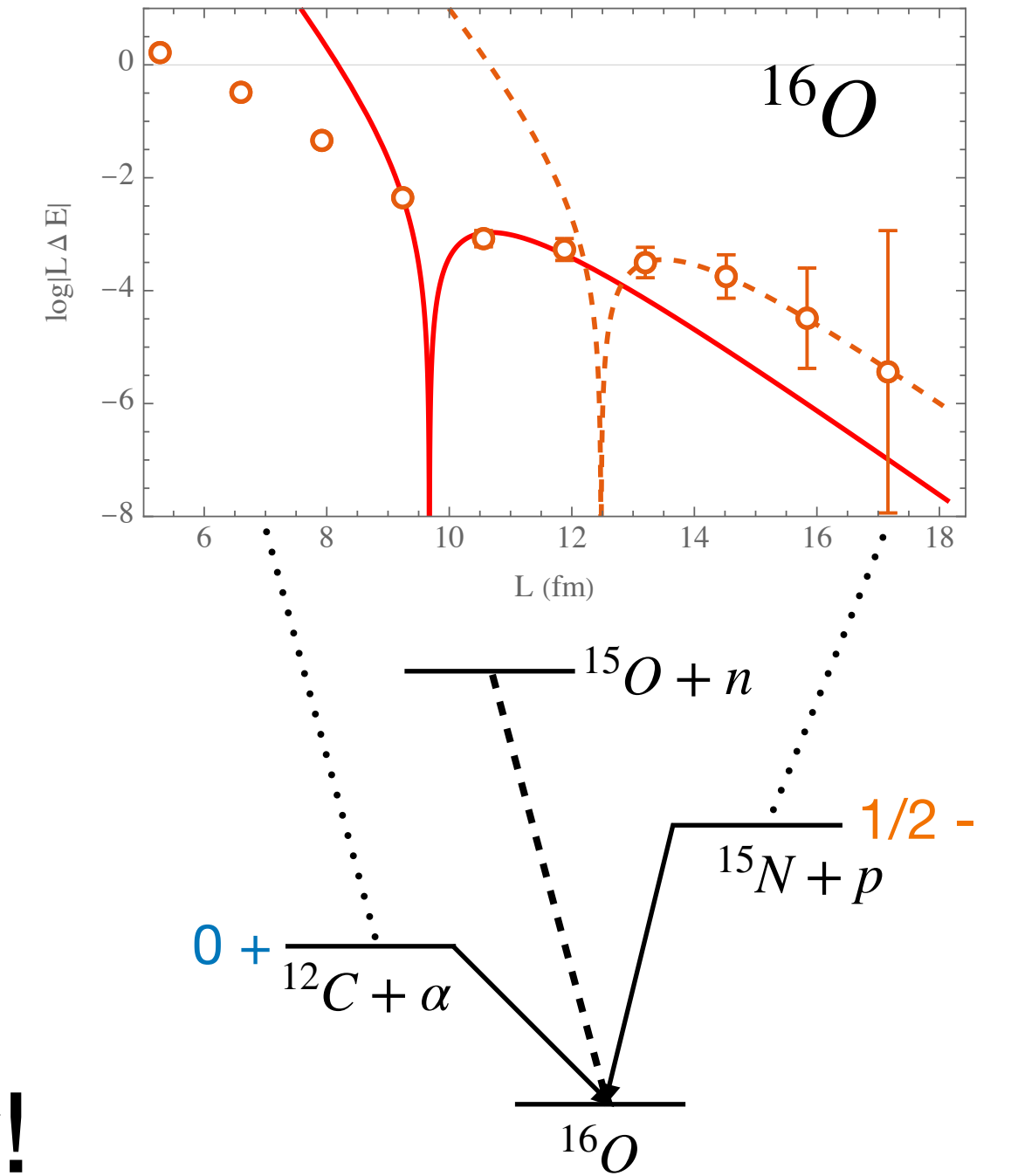
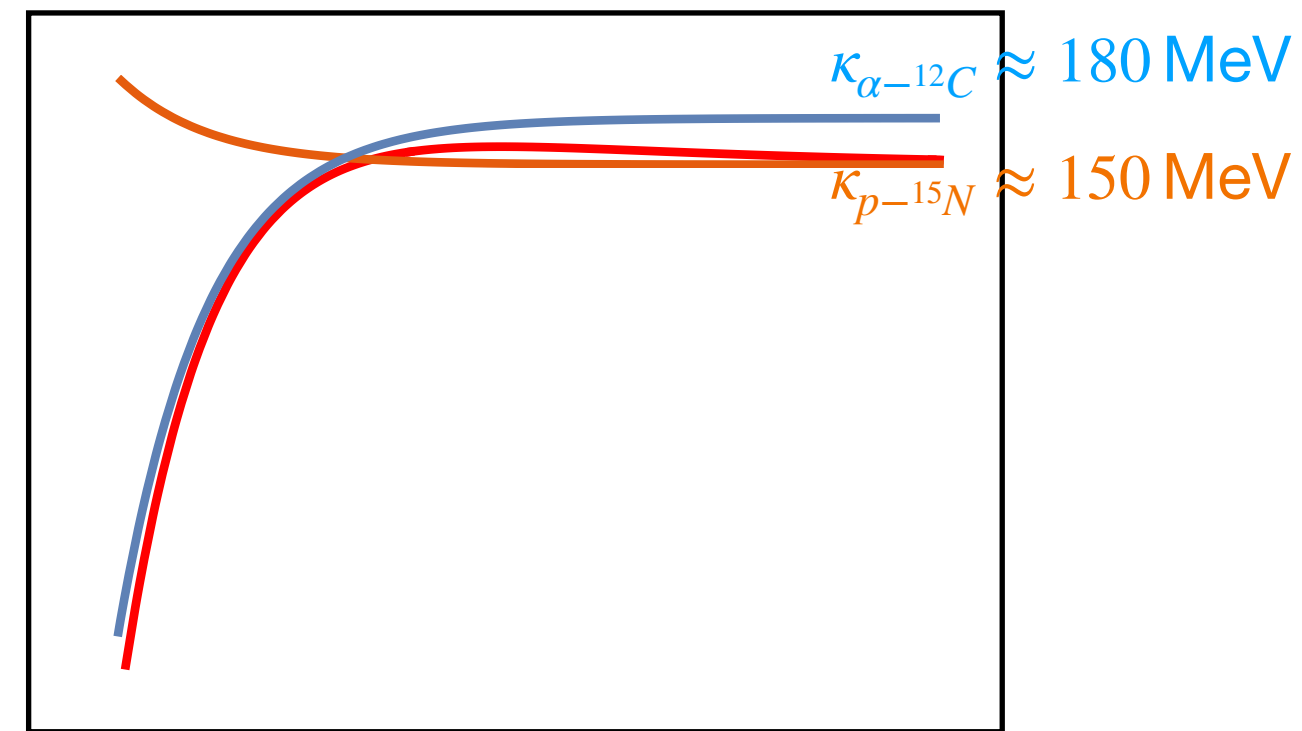
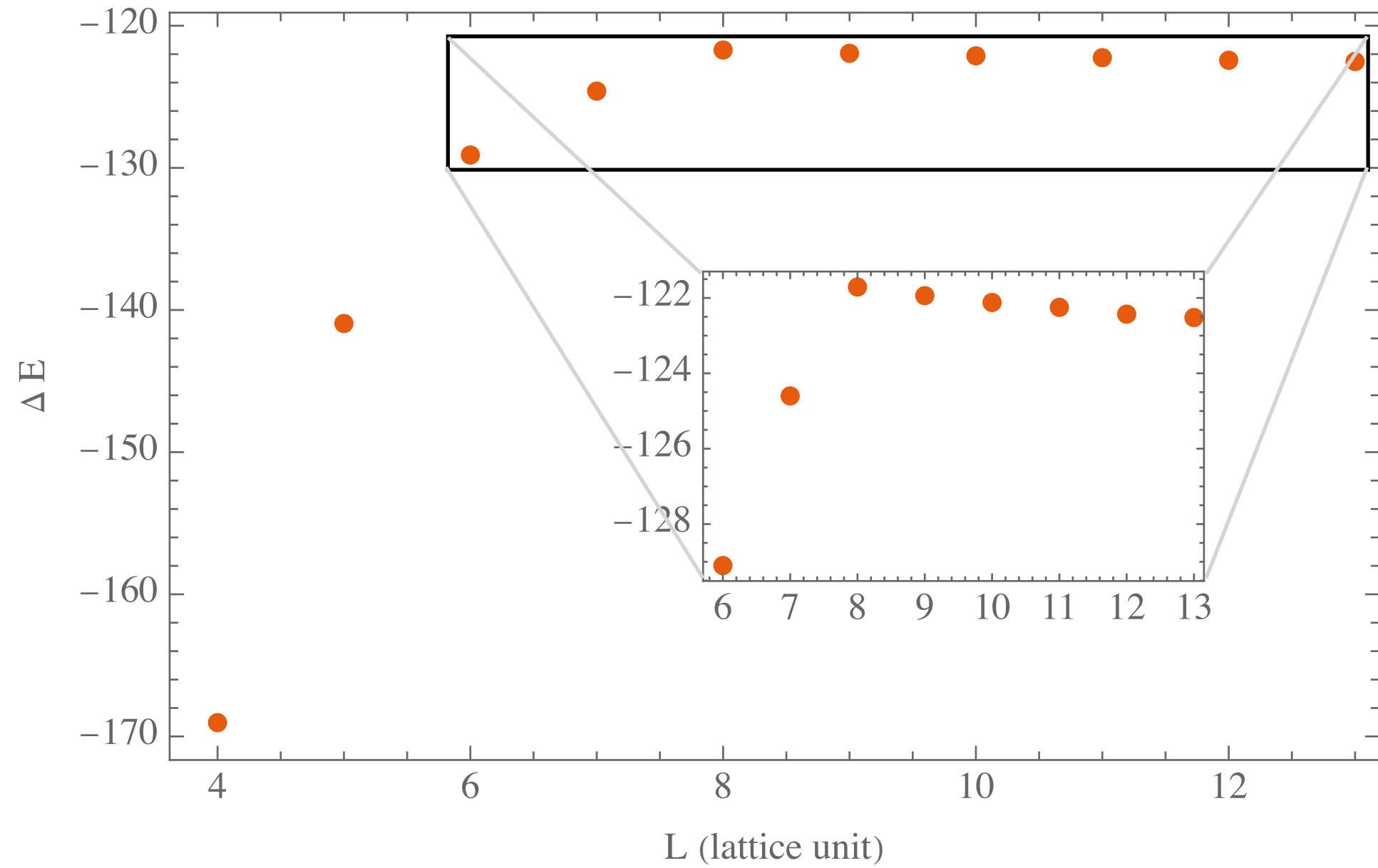


10% difference in Binding Energy!

Problem: overfitting

16O: a Multi-Channel Case

Prelim results & PhD Thesis



10% difference in Binding Energy!

Problem: overfitting

Solution: Applying for grants

ANCs of ^{16}O and ^{20}Ne

Isotope States	ANCs ($\text{fm}^{-1/2}$)	Reference
$^{16}\text{O}(0^+; \text{G.S.})$	11.01	Orlov [2021]
	58	deBoer et al. [2017]
	709	Sayre et al. [2012]
	13.9(24)	Adhikari and Basu [2009]
	750-4000	Morais and Lichtenthäler [2011]
	380 ± 80	Pinhole algorithm
	210 ± 20	This work (w/o fitting separation energy)
$^{20}\text{Ne}(0^+; \text{G.S.})$	980 ± 100	This work (w/ fitting separation energy)
	$3.3(1) \times 10^3$	Orlov [2021]
	2500 (1150)	Costantini et al. [2010]
	3400(700)	Costantini et al. [2010]; Mao et al. [1996]
	$3.80(95) \times 10^3$	Pinhole algorithm
	$3.31(2) \times 10^3$	This work

ANCs of ^{16}O and ^{20}Ne

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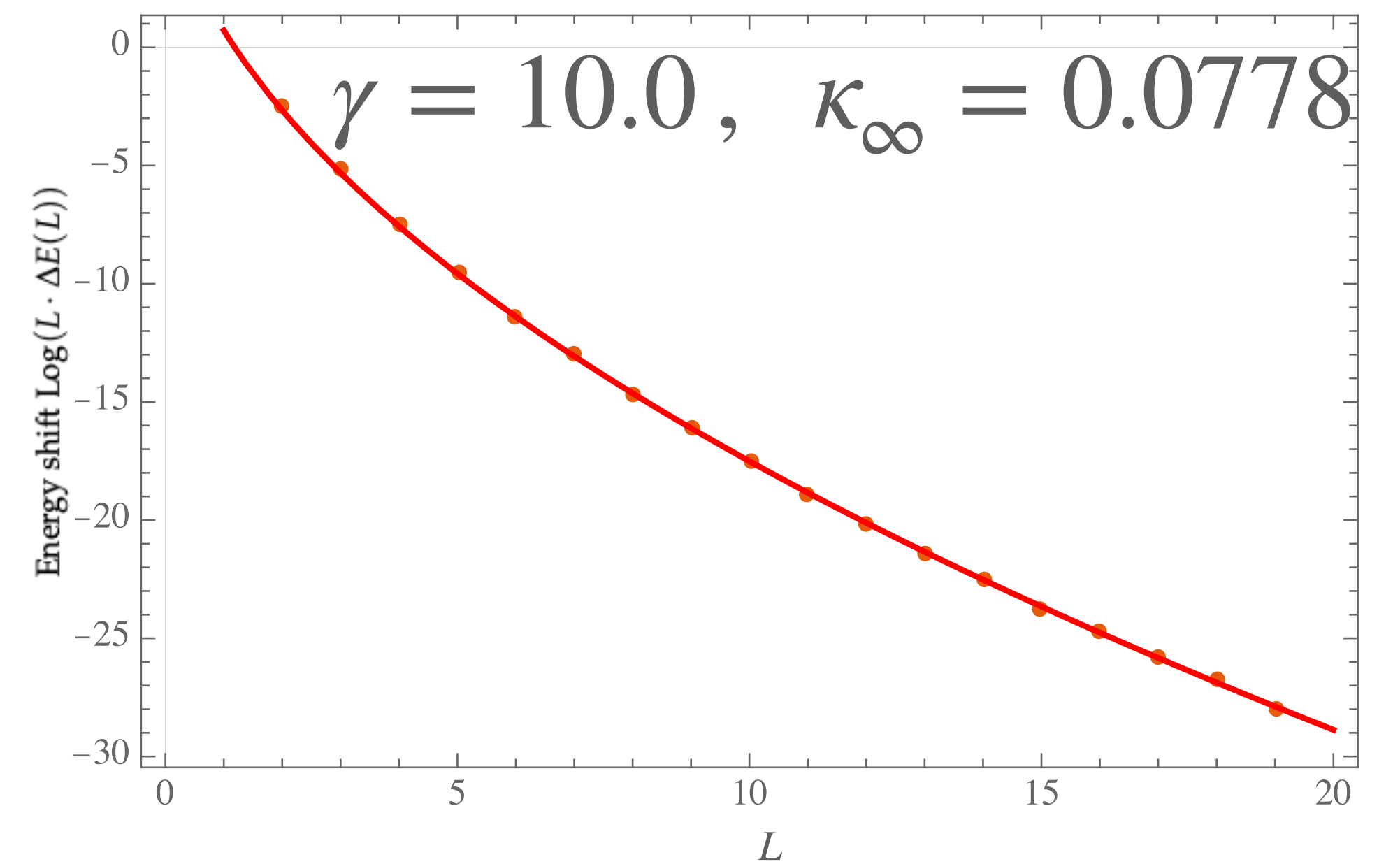
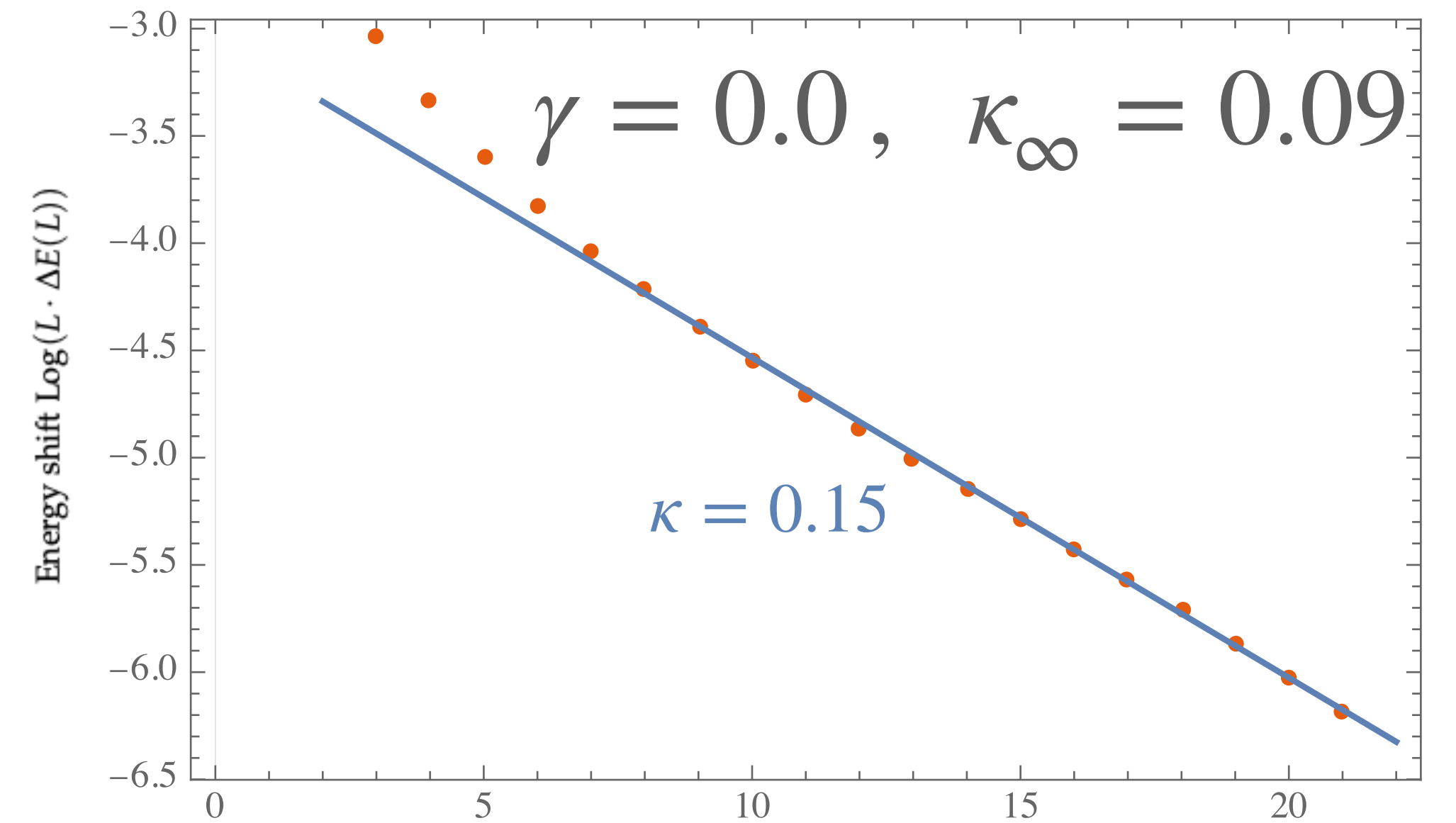
+ Geometric Factors

HY, et. al. 2024 (WIP)

SOON™

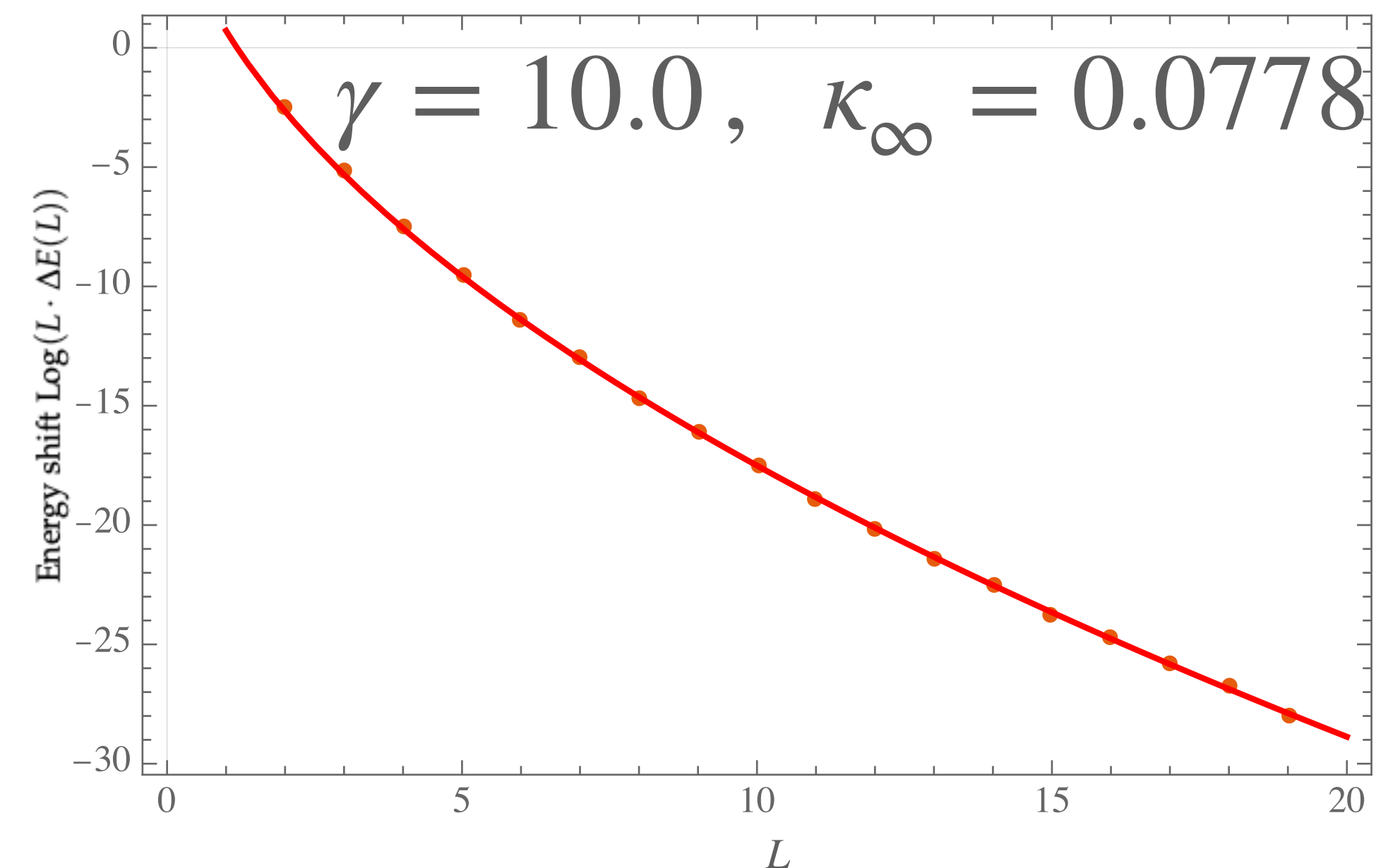
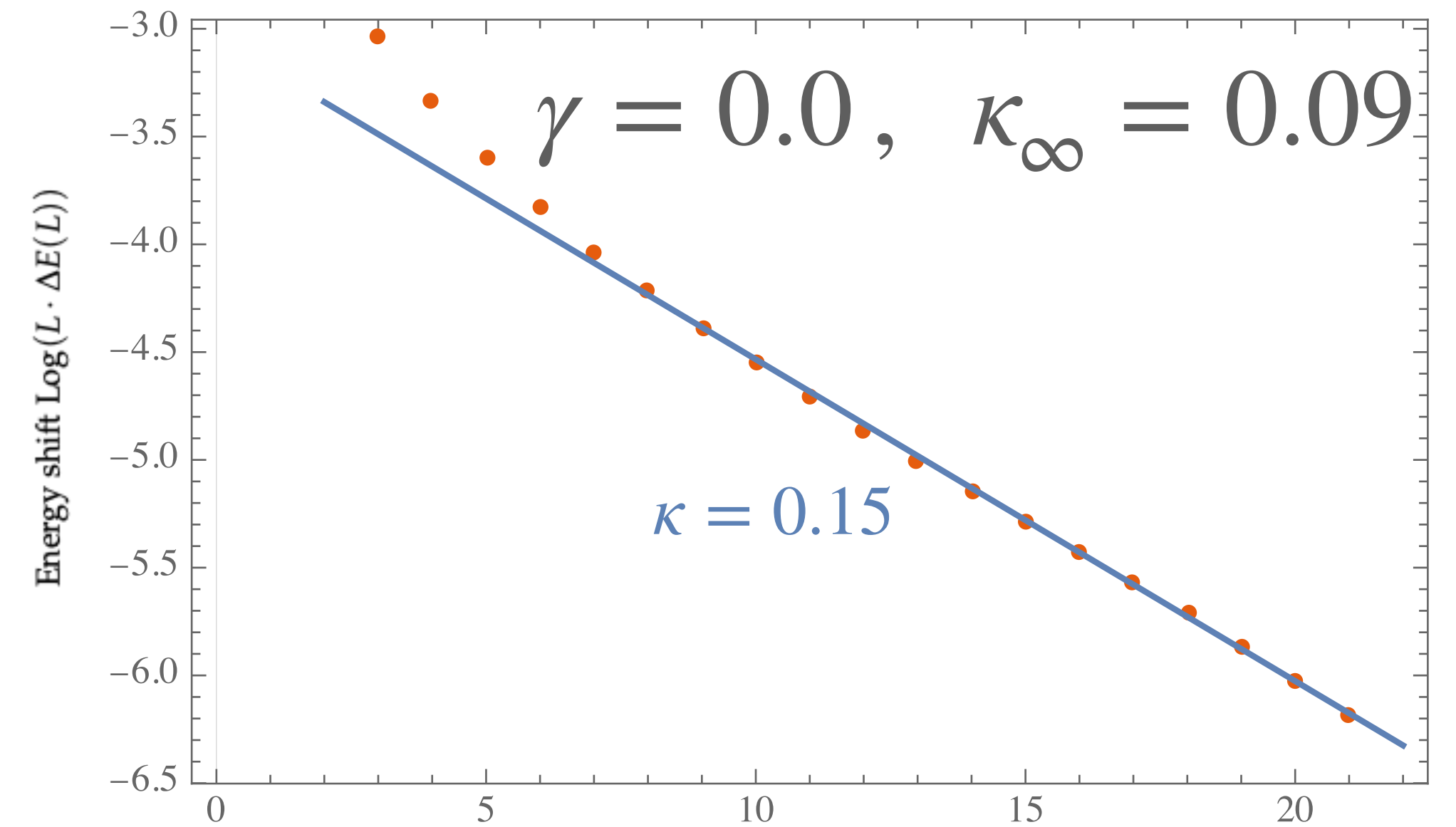
So, it is working. And...

Exotic bound-states with coulomb:



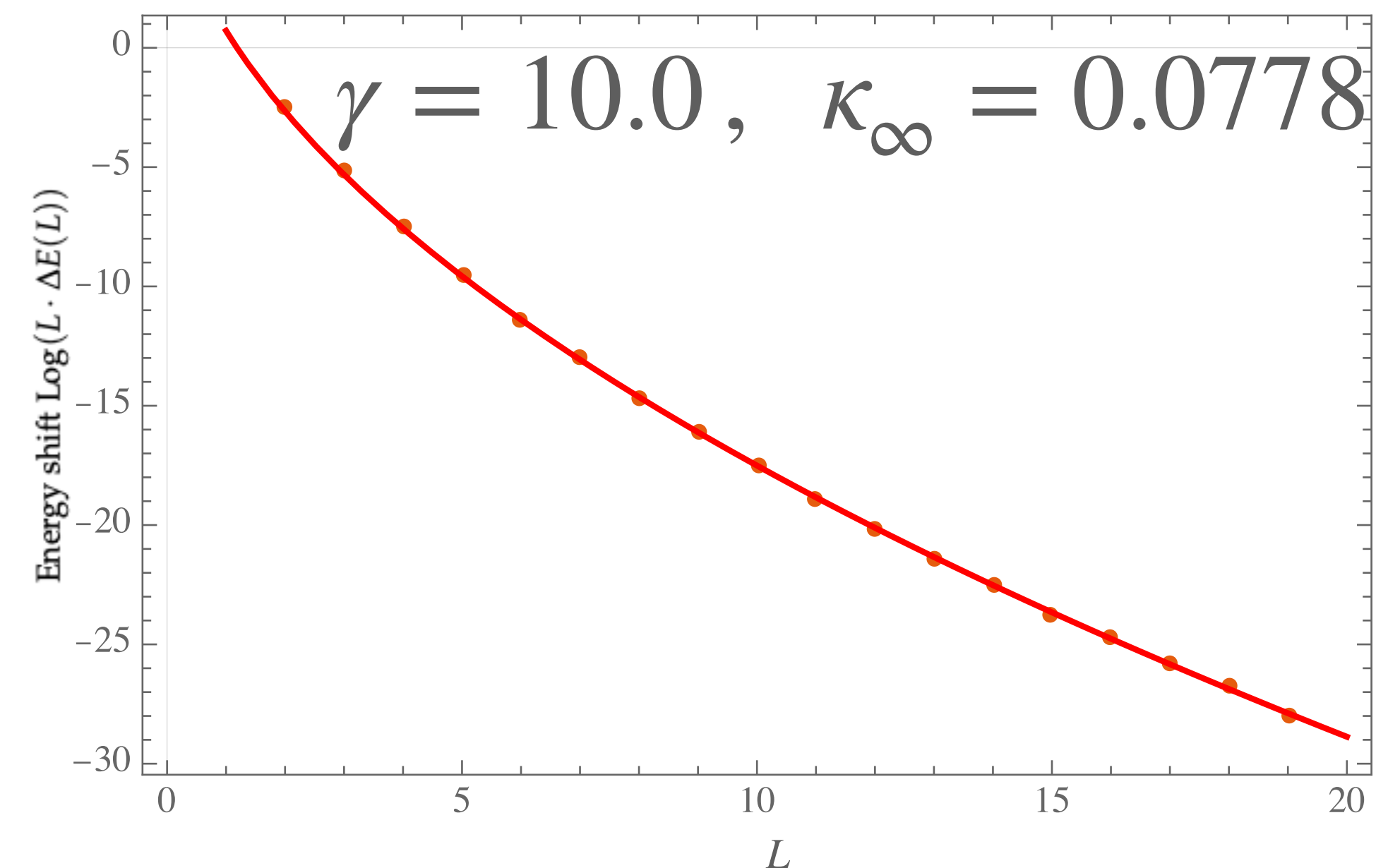
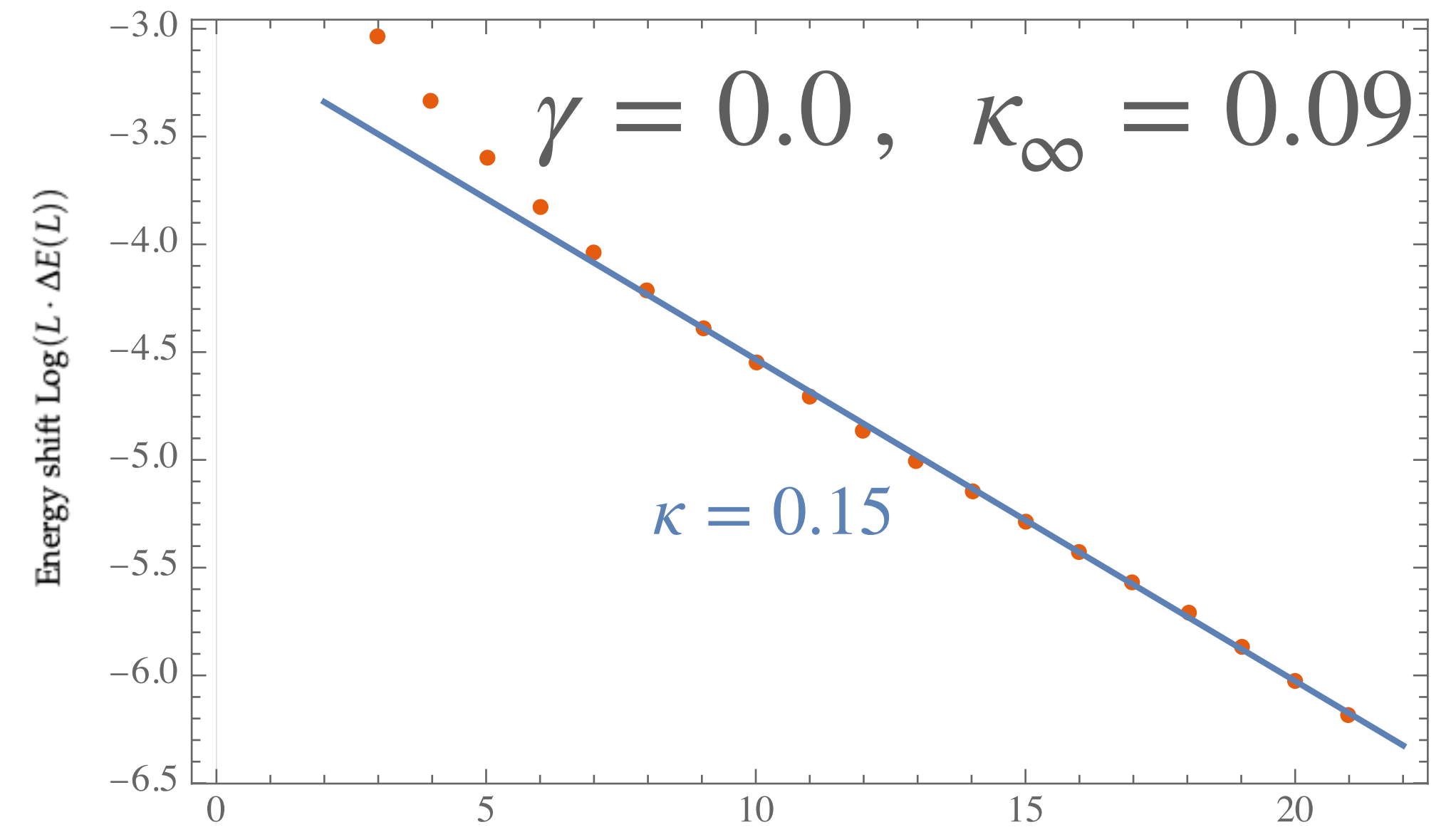
Exotic bound-states with coulomb:

- Exponential dependences (no-coulomb) do not apply to very small κ



Exotic bound-states with coulomb:

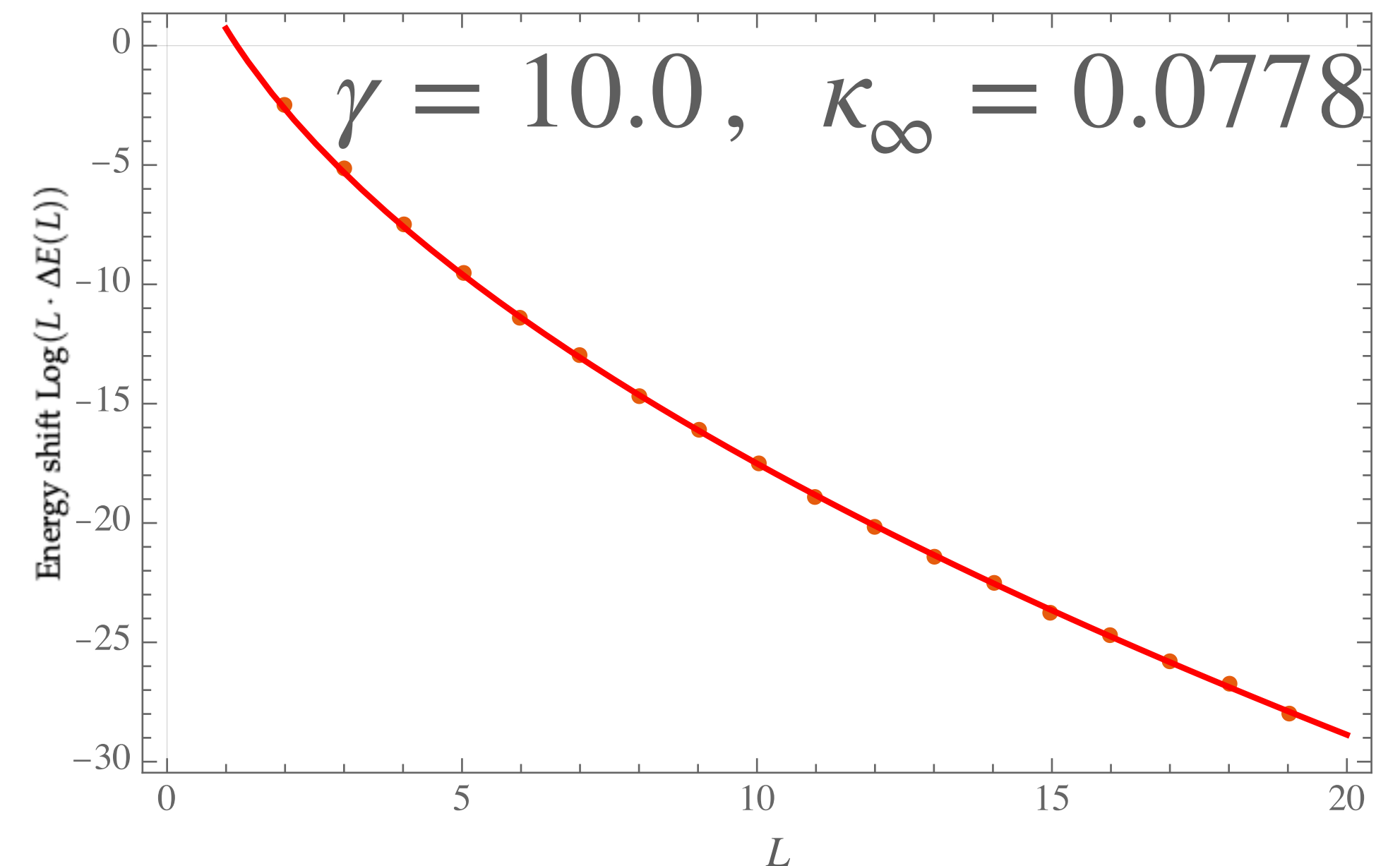
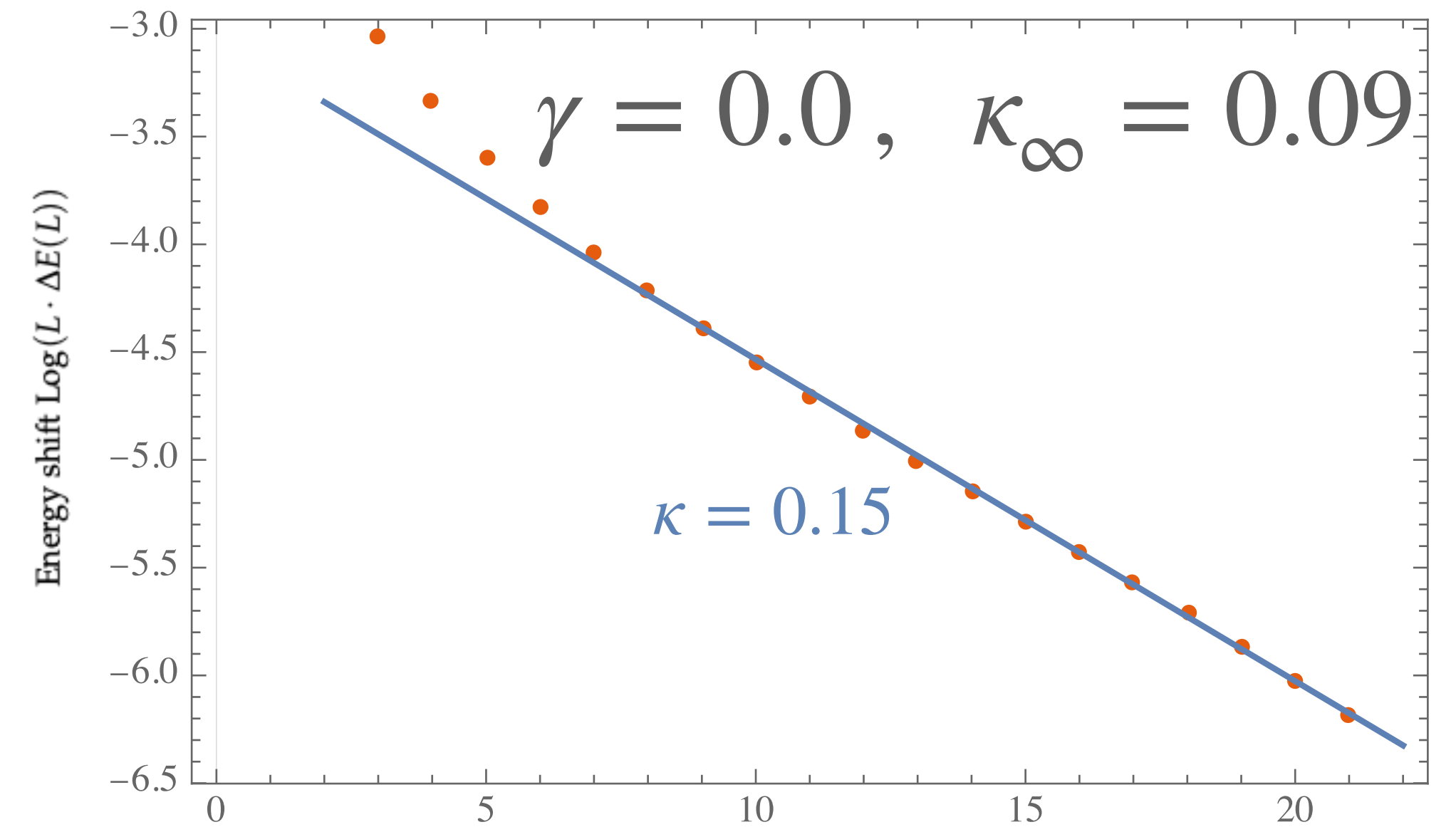
- Exponential dependences (no-coulomb) do not apply to very small κ
- But turning on coulomb force changes everything



Exotic bound-states with coulomb:

- Exponential dependences (no-coulomb) do not apply to very small κ
- But turning on coulomb force changes everything

	$\log C_0$	κ
FV	209.008	0.0772
WF	209.075	0.0770

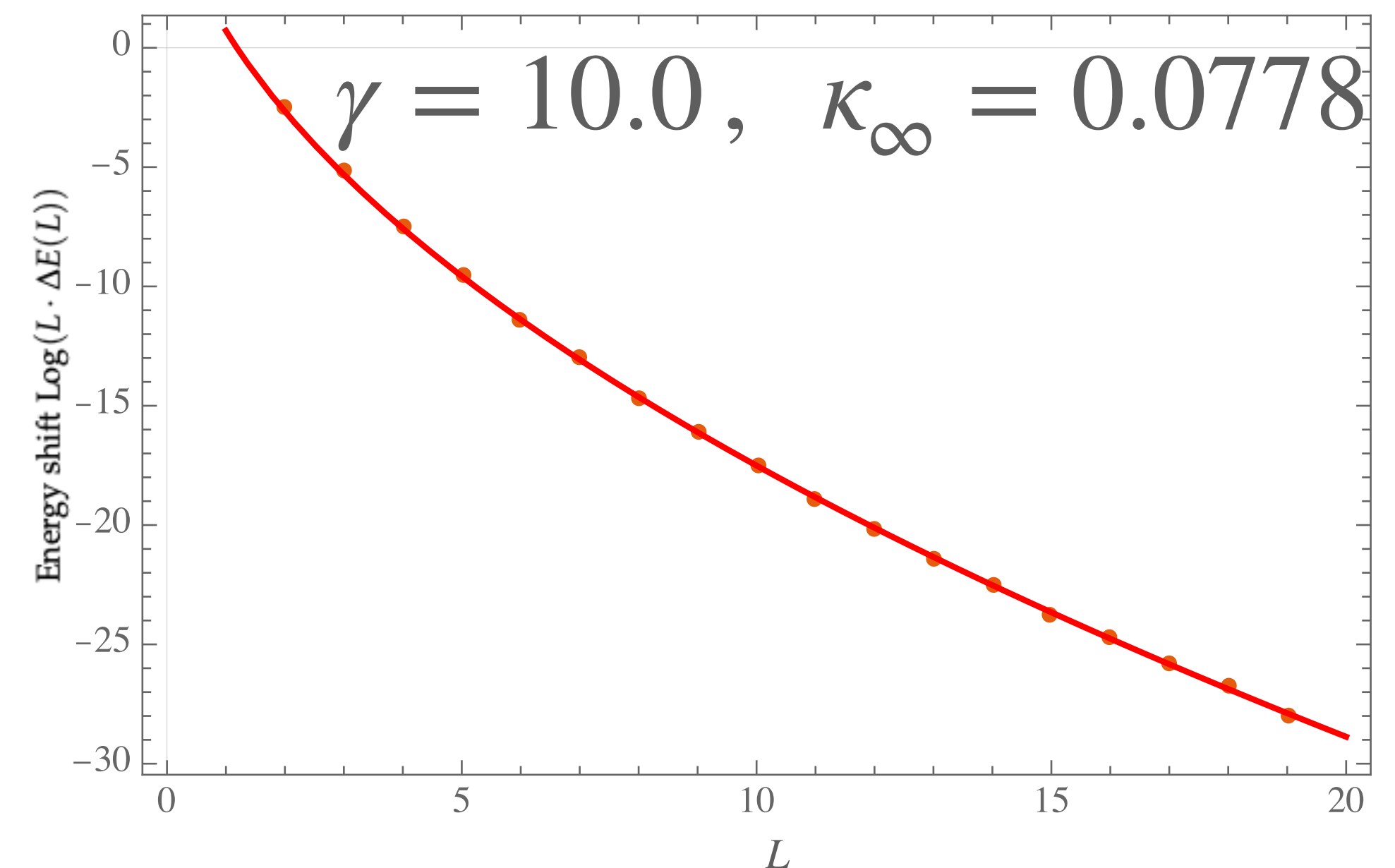
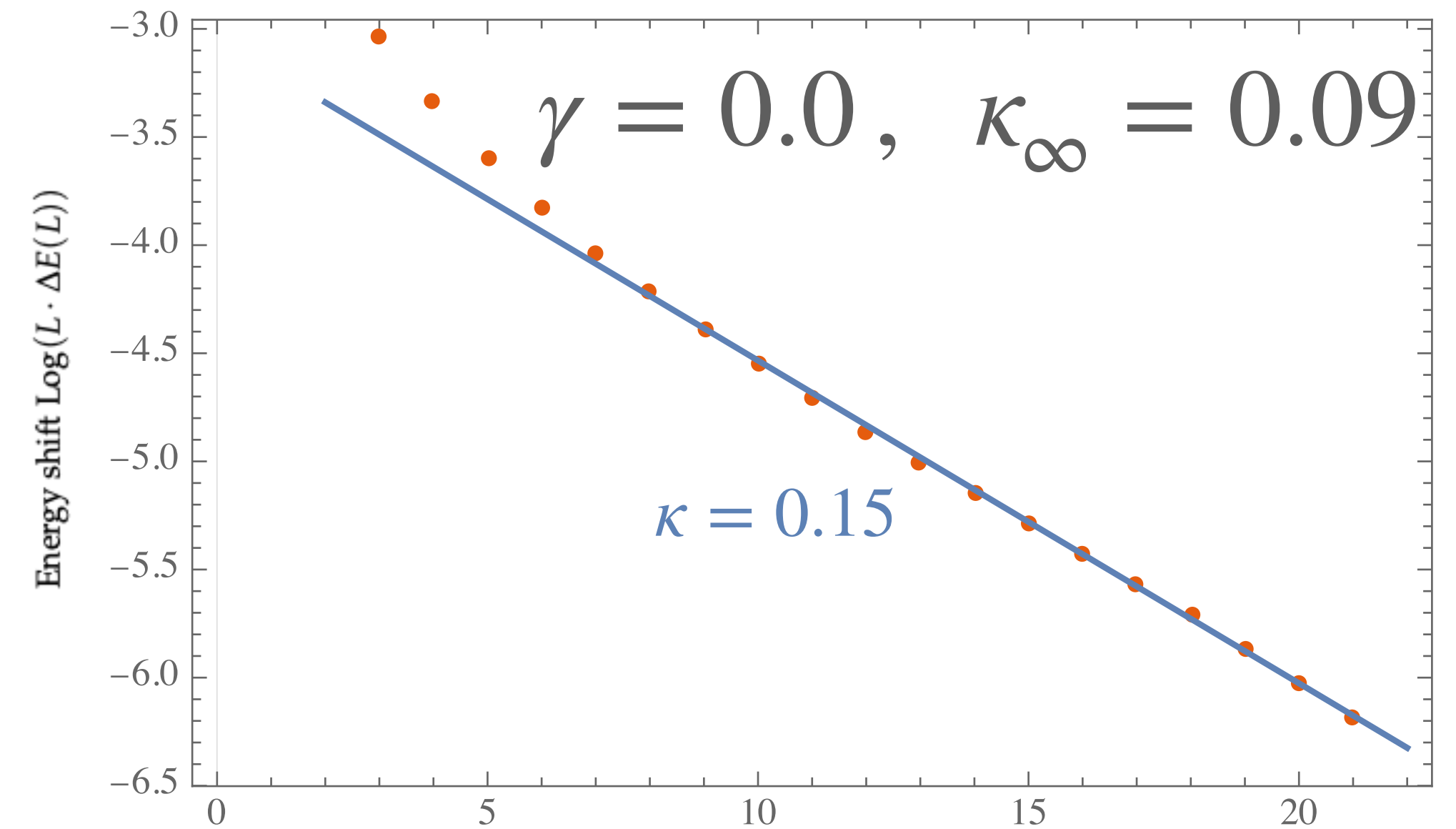


Exotic bound-states with coulomb:

- Exponential dependences (no-coulomb) do not apply to very small κ
- But turning on coulomb force changes everything

	$\log C_0$	κ
FV	209.008	0.0772
WF	209.075	0.0770

- Why?



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An important distinction

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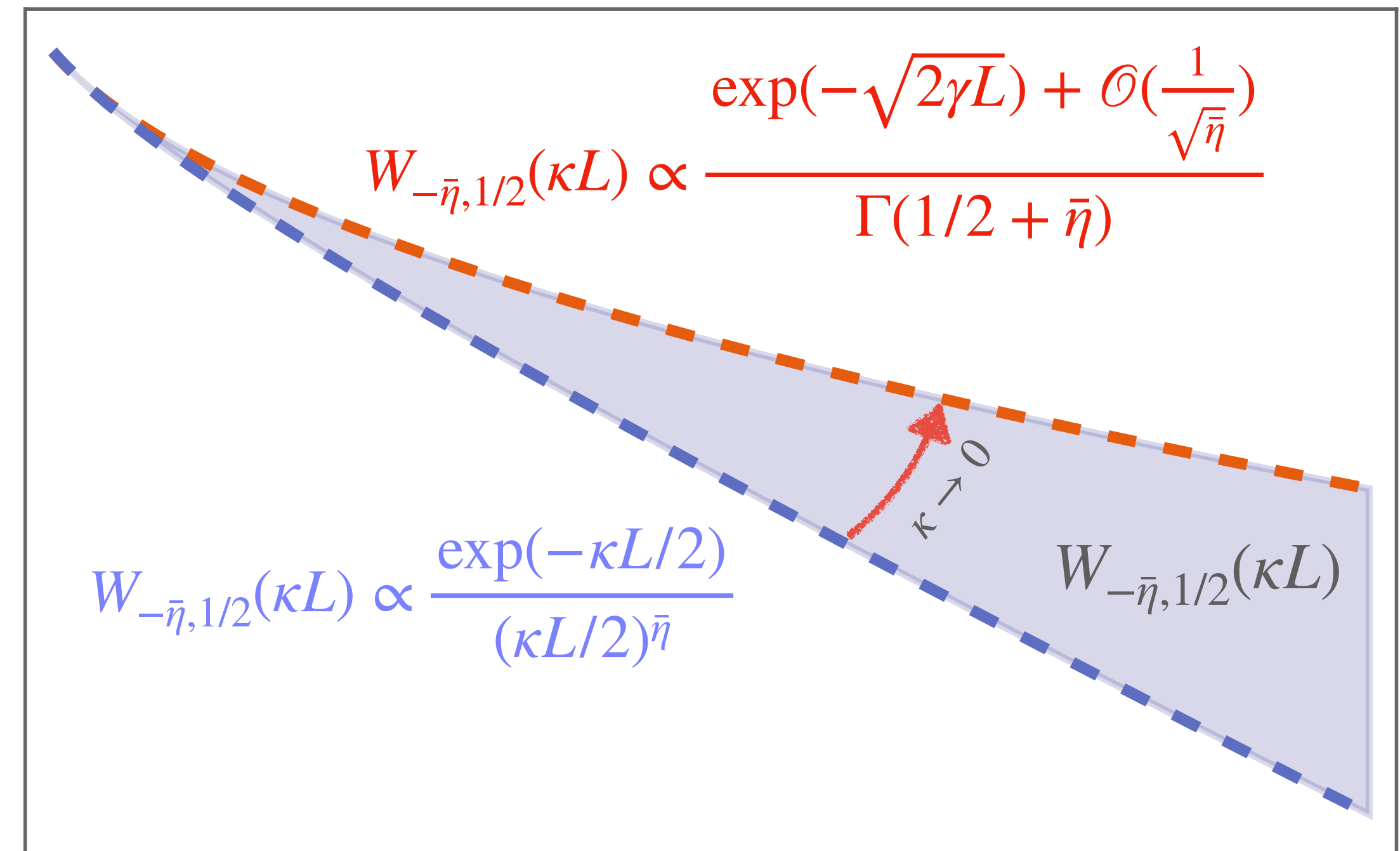
$$W_{-\bar{\eta}, 1/2}(\kappa L) \propto \frac{\exp(-\sqrt{2\gamma L}) + \mathcal{O}\left(\frac{1}{\sqrt{\bar{\eta}}}\right)}{\Gamma(1/2 + \bar{\eta})}$$

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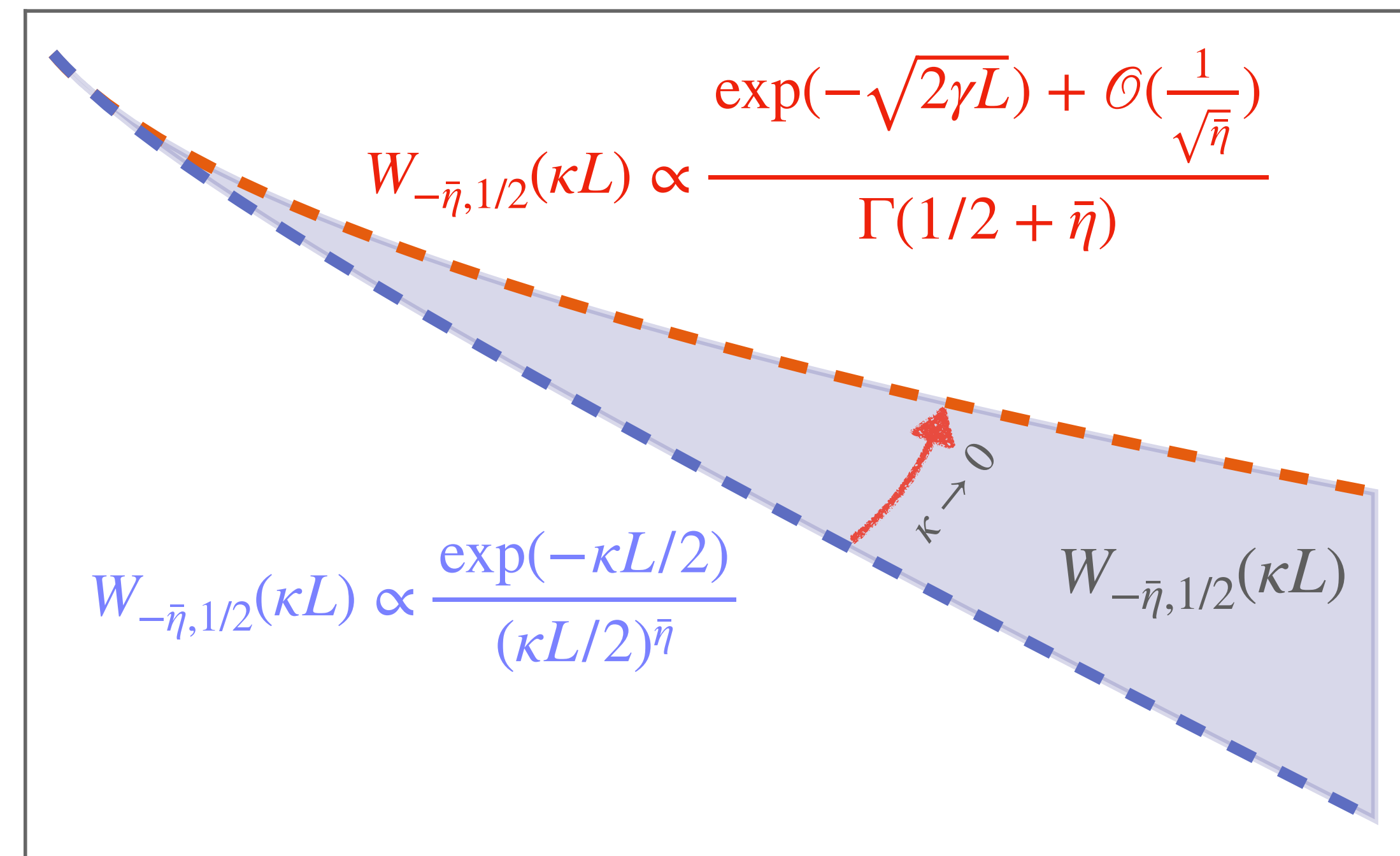
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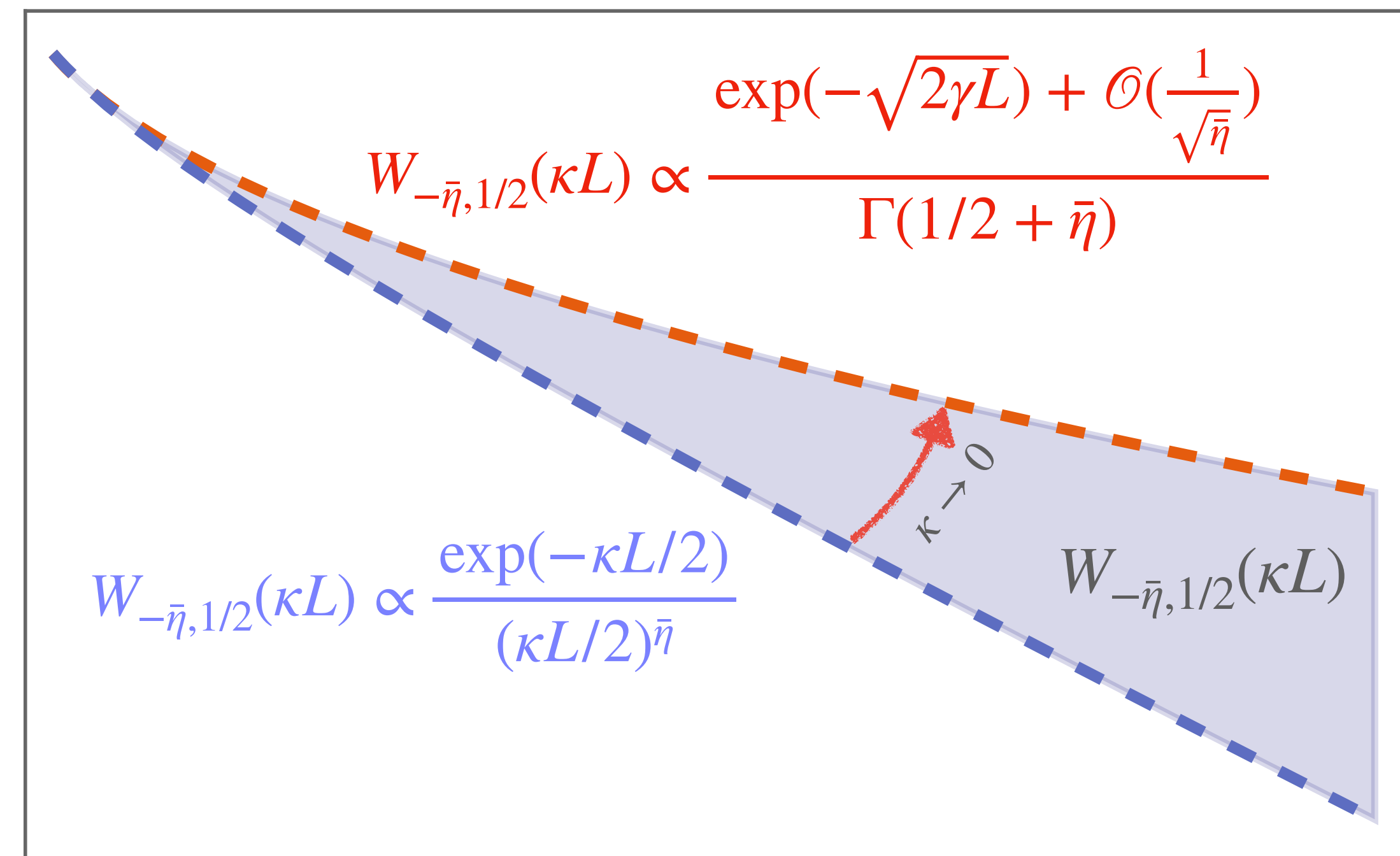
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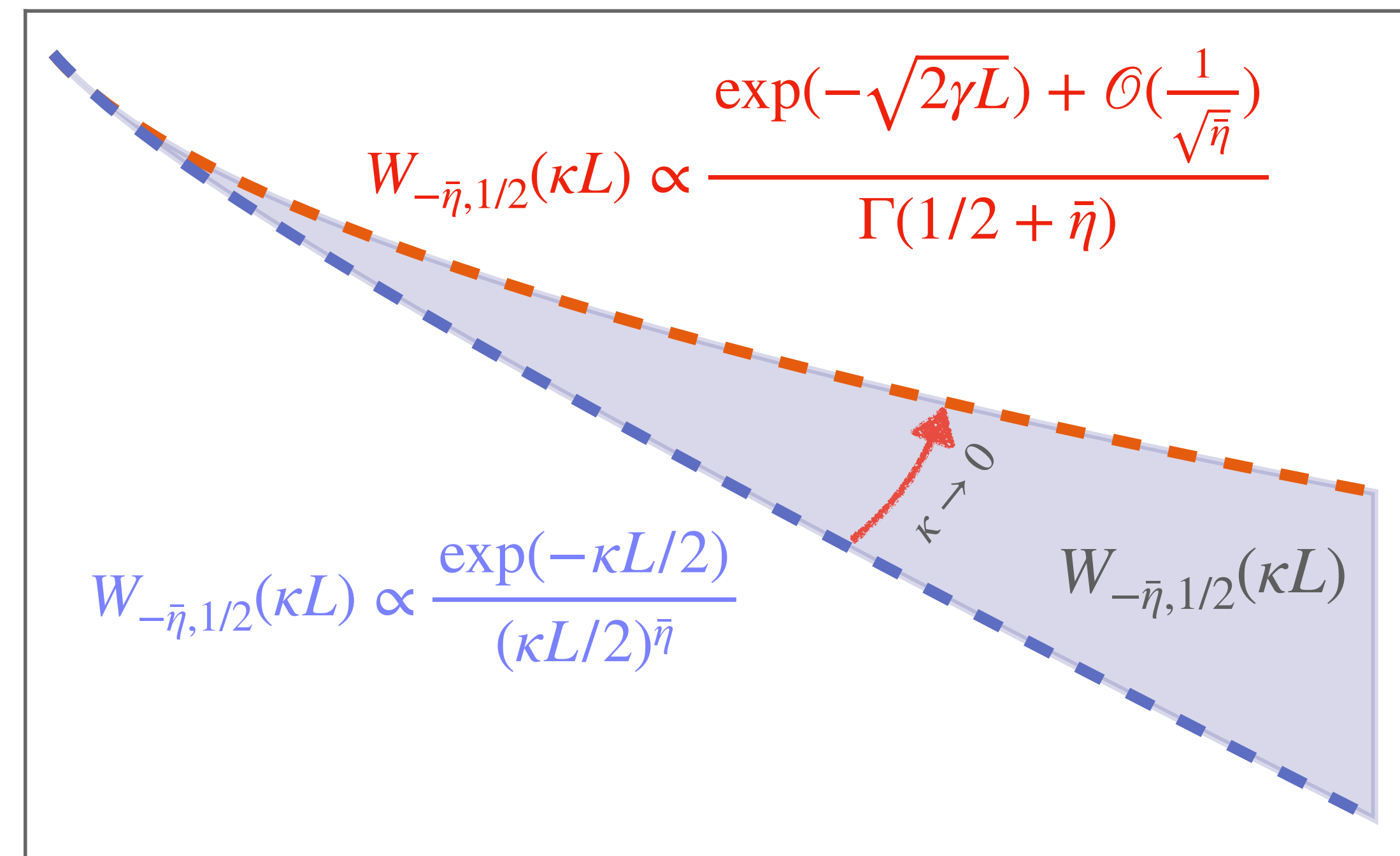
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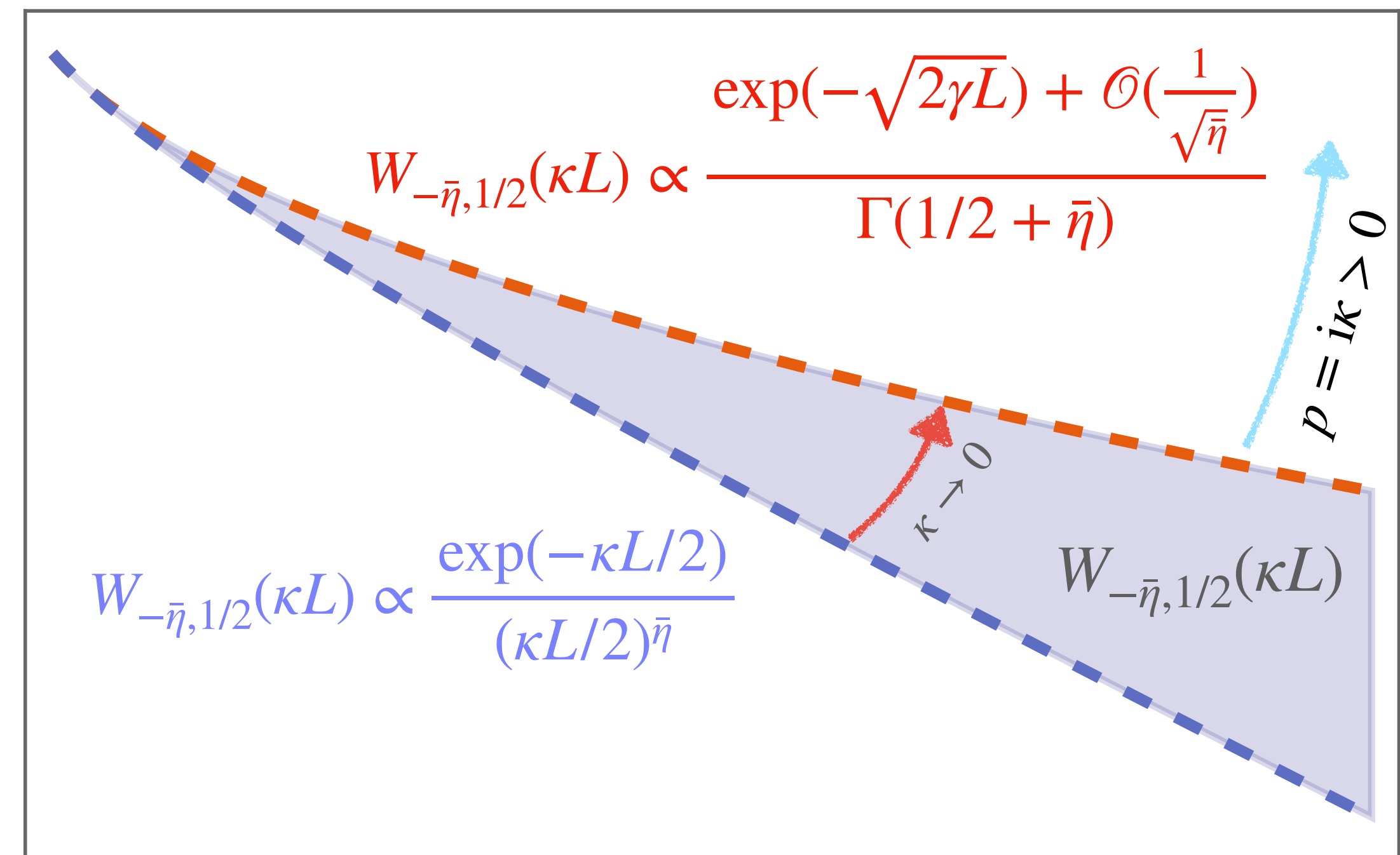
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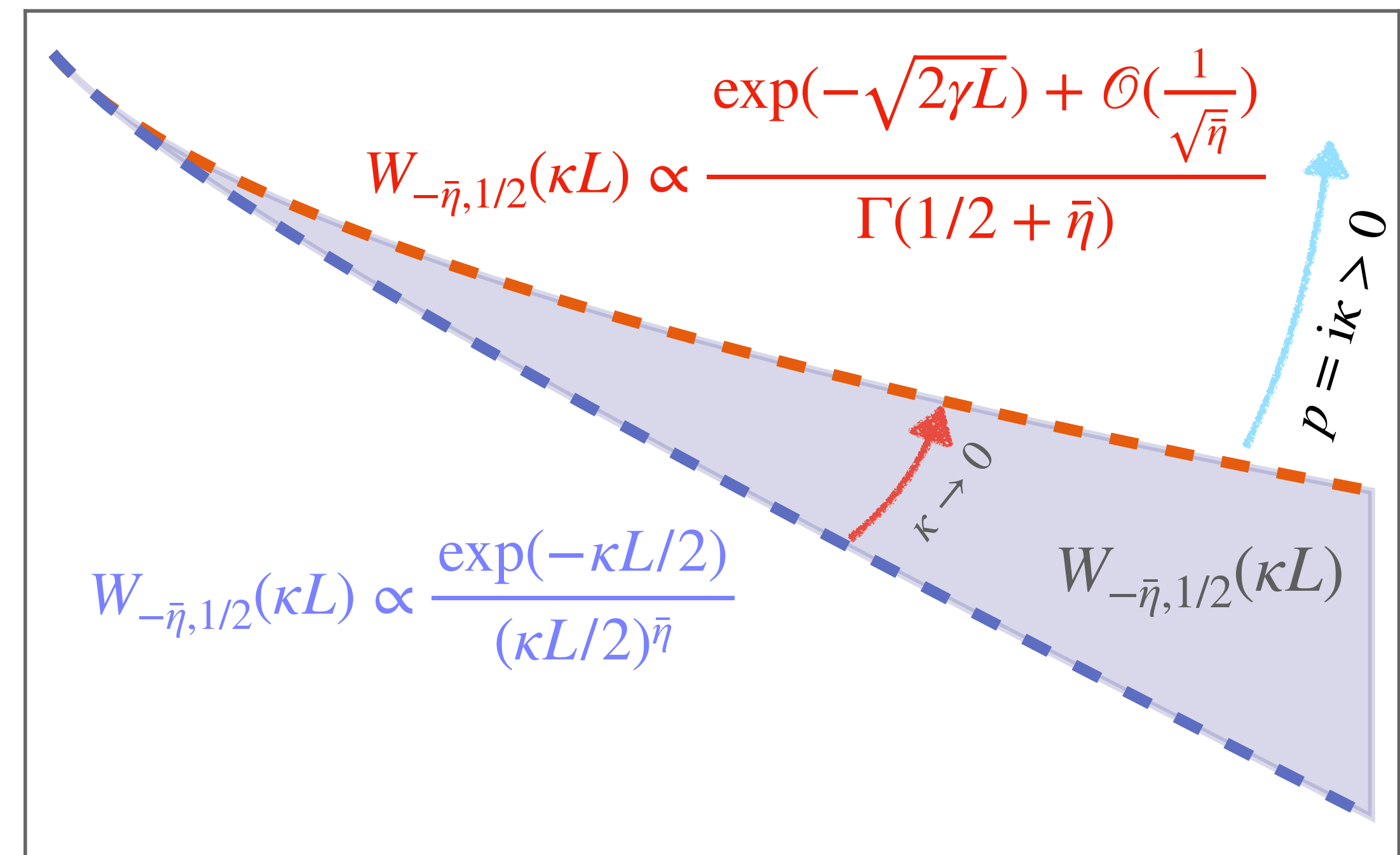
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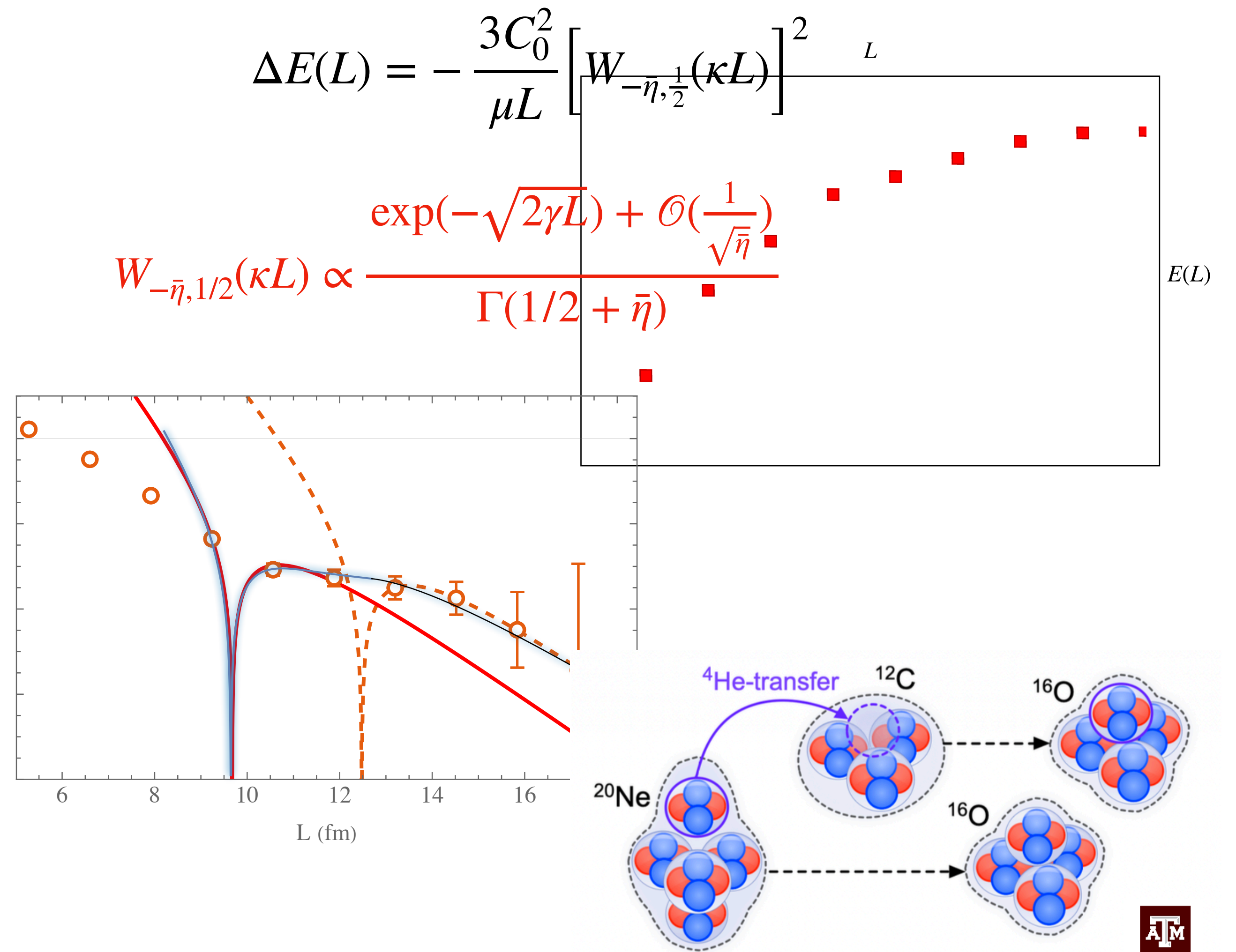
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To the Continuum!

Wrapping up...

Volume dependences:
from ground states to the continuum

- We include the difficult **coulomb interactions** (no-PT) for **bound states** in the volume dependencies
- This difficult interaction opens doors to the **continuum**
- Application: *ab initio* calculations (work in progress)
- ...



Thank you!

