

The 23rd International Conference on  
Few-Body Problems in Physics (FB23)

# Microwave Shielded Polar Molecules

Tao Shi

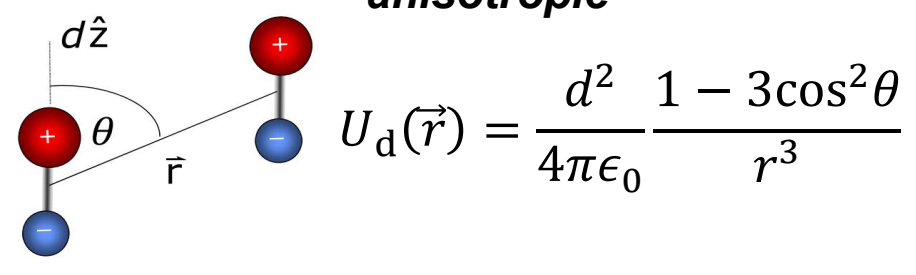
*23<sup>rd</sup> Sep. 2024, Beijing, China*



PRL **130**, 183001 (2023); Nature **626**, 283 (2024);  
arXiv:2405.13645; arXiv:2406.06412

# Dipolar interactions

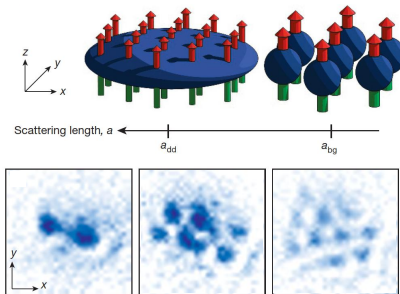
**Dipolar interaction: long-range and anisotropic**



$$U_d(\vec{r}) = \frac{d^2}{4\pi\epsilon_0} \frac{1 - 3\cos^2\theta}{r^3}$$

## Magnetic Atoms

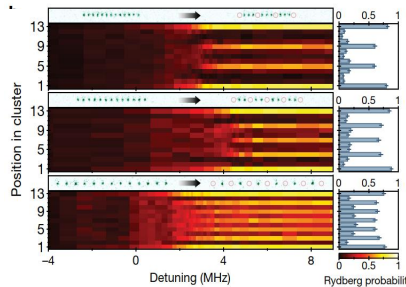
Weak dipoles  $\sim 10 \mu_B$ , stable (10 s)



H. Kadau et al.,  
Nature 530, 194 (2016)

## Rydberg Atoms

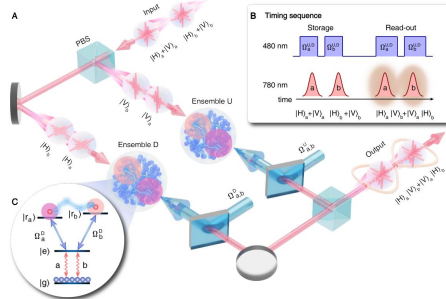
Strong dipoles  $\sim 104$  Debye, lifetime  $\sim 100 \mu s$



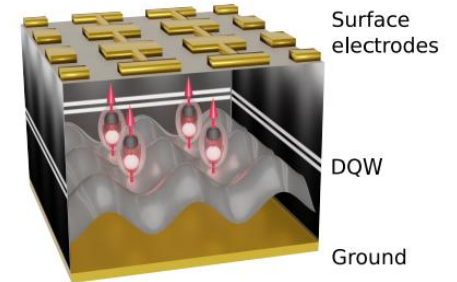
H. Bernien et al.,  
Nature 551, 579 (2017)

## Dipolar Excitons

Strong dipoles  $\sim 103$  Debye, lifetime  $\sim 1 \mu s$



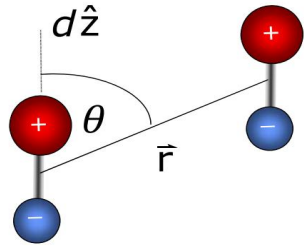
G. Ye, et al.,  
Nature Photonics (2023)



C. Lagoin et al.,  
Nature 609, 485 (2022)

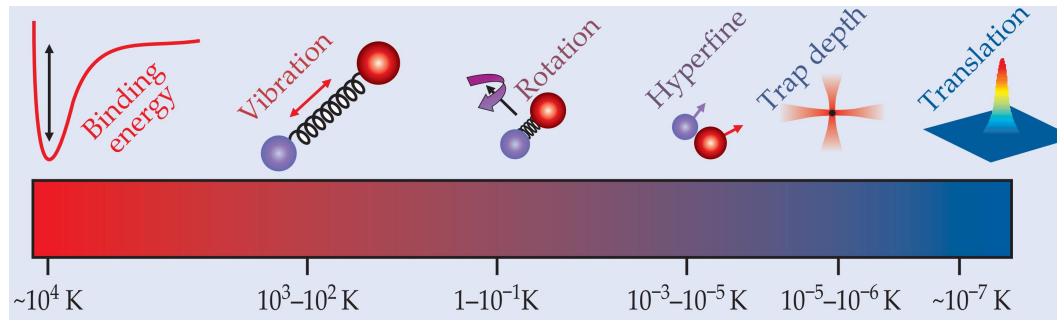
# Dipolar interactions

**Dipolar interaction: long-range and anisotropic**


$$U_d(\vec{r}) = \frac{d^2}{4\pi\epsilon_0} \frac{1 - 3\cos^2\theta}{r^3}$$

## Dipolar Molecules

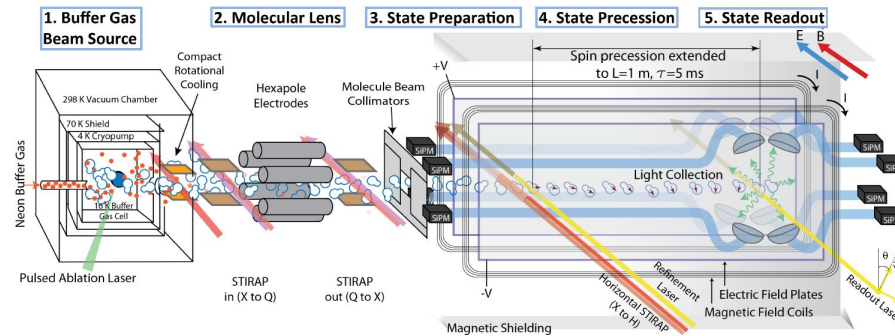
Medium dipoles  $\sim 3$  Debye, alone stable (10 s)



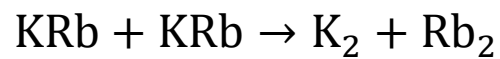
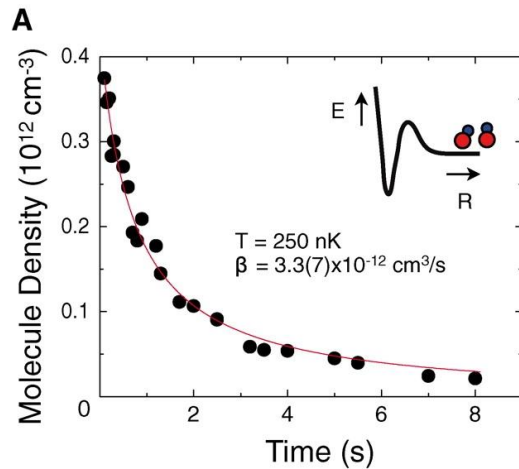
# Polar Molecules

## Precision measurement

ACME III (Harvard), JILA, ICL, Caltech, Amsterdam...

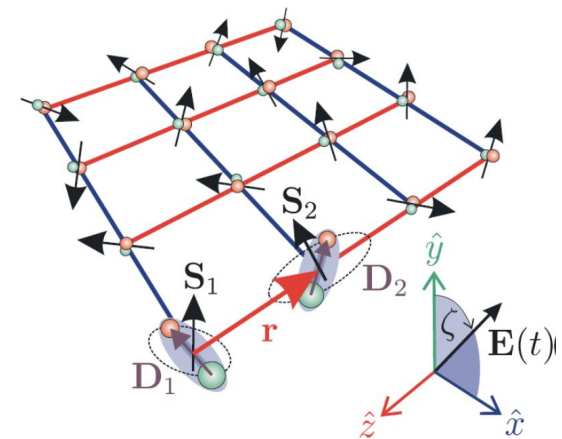


## Ultracold chemistry



JILA, 2010

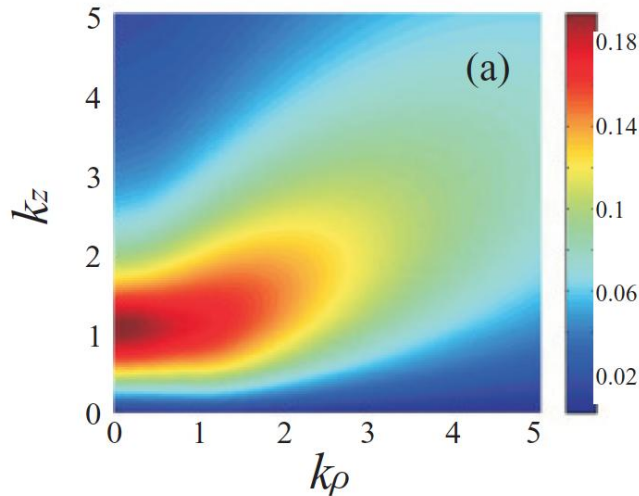
## Spin systems



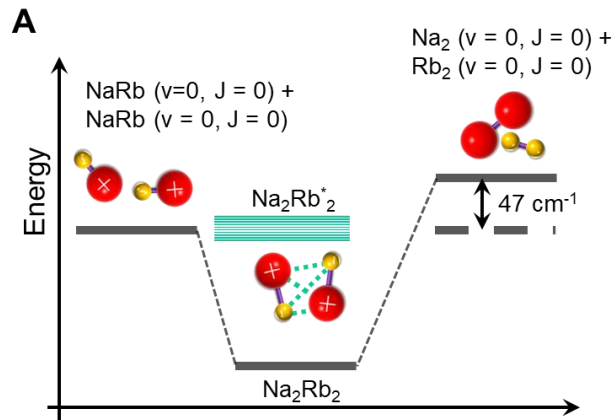
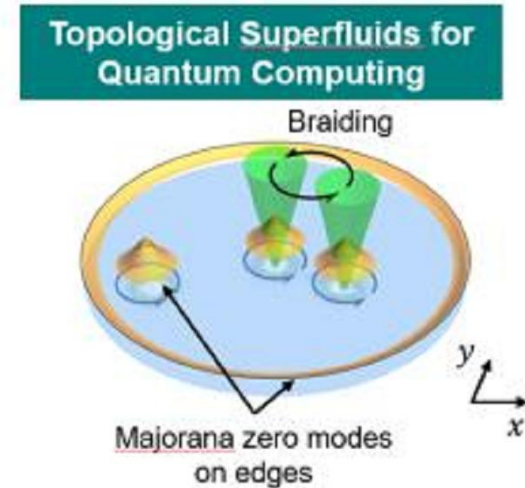
A. Micheli, Nat. Phys. 2, 341 (2006).

# BCS-BEC in Polar Molecules

## Dipolar superfluids

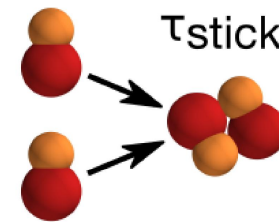


- You and Marinescu, PRA 60, 2324 (1999).
- Baranov et al., PRA 66, 013606 (2002).
- Cooper and Shlyapnikov, PRL 103, 155302 (2009).
- Shi, Zhang, Sun, and Yi, PRA **82**, 033623 (2010).
- Wu and Hirsch, PRB **81**, 020508 (2010).
- Shi, Zou, Hu, Sun, and Yi, PRL 110, 045301 (2013).
- Qi, Shi, and Zhai, PRL 110, 045302 (2013).



X. Ye, et al., Sci. Adv. 4, eaaq0083 (2018).

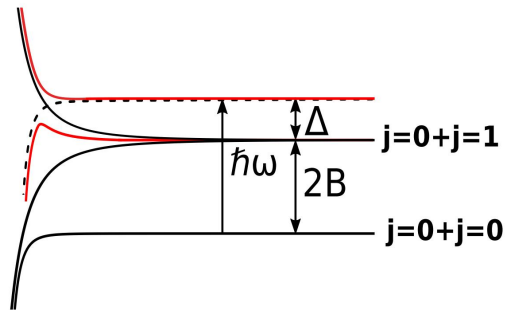
**unstable at the short range**



Bause et al., PRR 3, 033013 (2021).

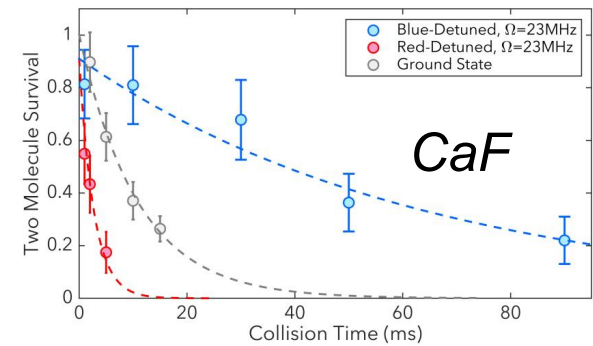
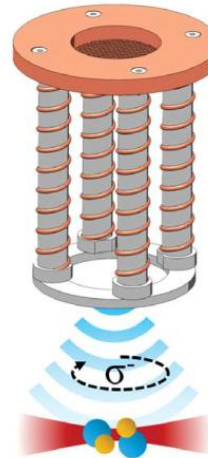
# Microwave Shielding

## Theoretical proposal:



Karman and Hutson, PRL 121, 163401 (2018)  
Lassablière and Quéméner, PRL 121, 163402 (2018)

## Suppression of inelastic scatterings:



Anderegg et al., Science 373, 779 (2021).

# Realization of fermionic MSPMs

nature

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## Evaporation of microwave-shielded polar molecules to quantum degeneracy

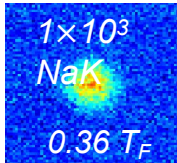
[Andreas Schindewolf](#), [Roman Bause](#), [Xing-Yan Chen](#), [Marcel Duda](#), [Tijs Karman](#), [Immanuel Bloch](#) & [Xin-Yu Luo](#) 

[Nature](#) **607**, 677–681 (2022) | [Cite this article](#)



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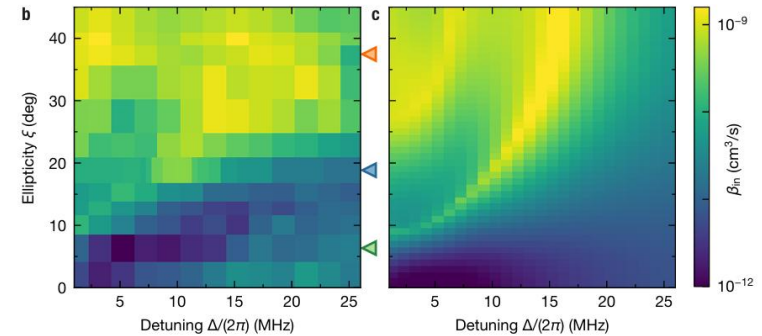
nature > articles > article

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## Field-linked resonances of polar molecules

[Xing-Yan Chen](#), [Andreas Schindewolf](#), [Sebastian Eppelt](#), [Roman Bause](#), [Marcel Duda](#), [Shrestha Biswas](#), [Tijs Karman](#), [Timon Hilker](#), [Immanuel Bloch](#) & [Xin-Yu Luo](#) 

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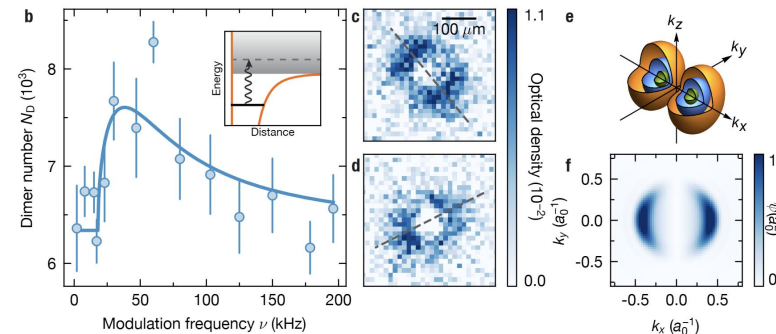
nature > articles > article

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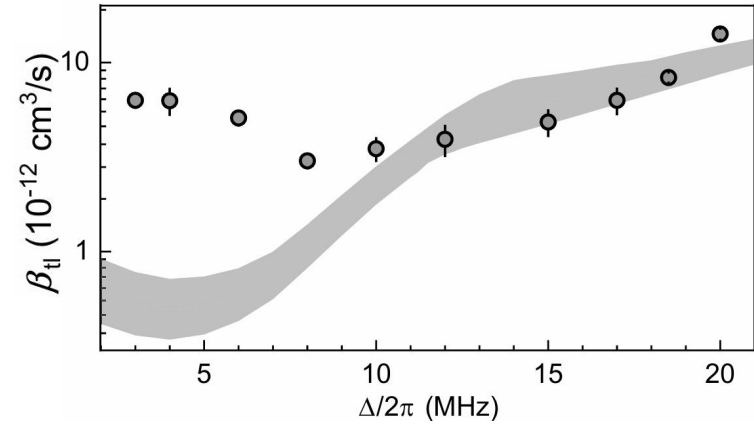
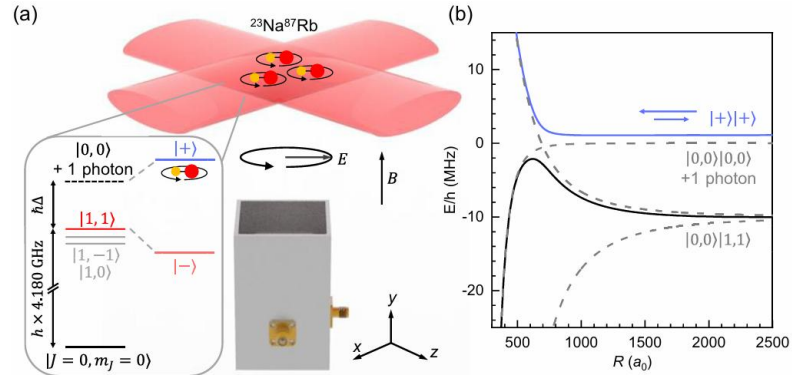
## Ultracold field-linked tetratomic molecules

[Xing-Yan Chen](#), [Shrestha Biswas](#), [Sebastian Eppelt](#), [Andreas Schindewolf](#), [Fulin Deng](#), [Tao Shi](#)  [Su Yi](#), [Timon A. Hilker](#), [Immanuel Bloch](#) & [Xin-Yu Luo](#) 

[Nature](#) **626**, 283–287 (2024) | [Cite this article](#)



# Realization of Bosonic MSPMs



J. Lin et al., Phys. Rev. X **13**, 031032 (2023)

nature

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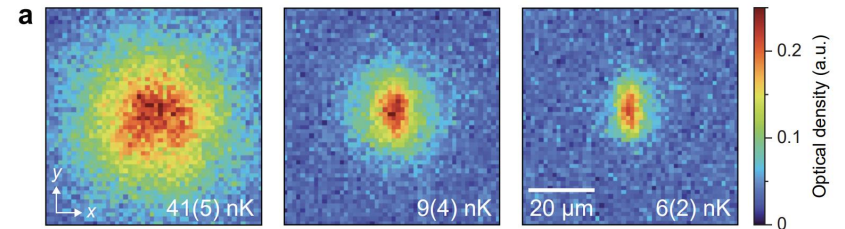
[nature](#) > [articles](#) > [article](#)

Article | Published: 03 June 2024

## Observation of Bose–Einstein condensation of dipolar molecules

[Niccolò Bigagli](#), [Weijun Yuan](#), [Siwei Zhang](#), [Boris Bulatovic](#), [Tijs Karman](#), [Ian Stevenson](#) & [Sebastian Will](#)

[Nature](#) **631**, 289–293 (2024) | [Cite this article](#)





# Outline

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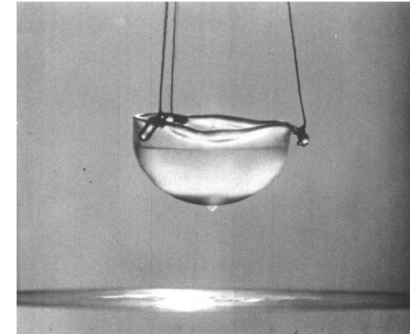
Introduction

Effective potentials for MSPMs

Fermionic MSPMs

Bosonic MSPMs

Conclusion and Outlook



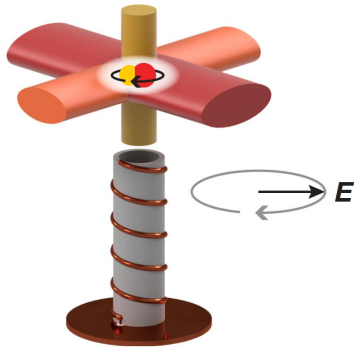
Liquid Helium

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Effective potentials for MSPMs

# Microwave shielding

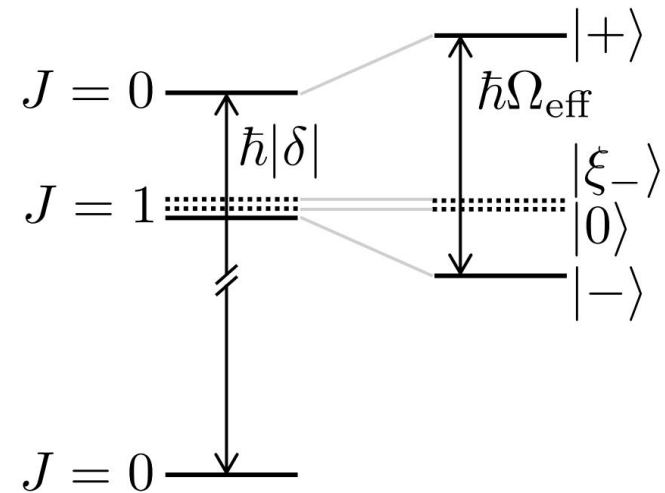
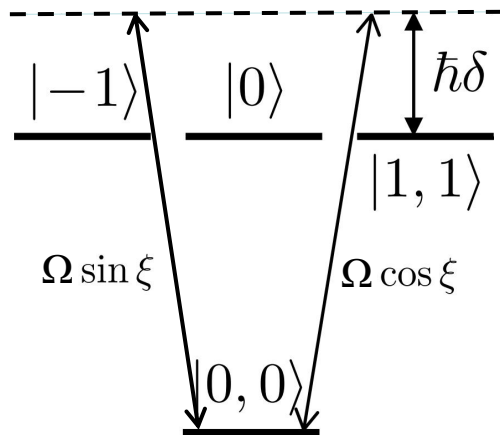
## Single molecule:



$$\frac{\hbar\Omega}{2} e^{-i\omega_0 t} |\xi_+\rangle \langle 0,0| + \text{h.c.}$$

$$|\xi_+\rangle \equiv \cos \xi |1,1\rangle + \sin \xi |1,-1\rangle$$

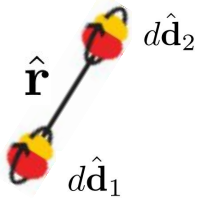
$$|\xi_-\rangle \equiv \cos \xi |1,-1\rangle - \sin \xi |1,1\rangle$$



# Microwave shielding

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**Two molecules:**



Two-body Hamiltonian:

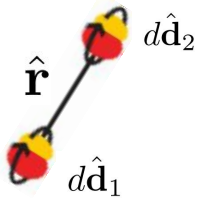
$$\hat{H}_2 = \sum_{j=1,2} \hat{h}_j + V(\mathbf{r}_1 - \mathbf{r}_2)$$

Kinetic energy:

$$\hat{h}_j = -\hbar^2 \nabla_j^2 / (2M) + \hat{h}_{\text{in}}(j)$$

# Microwave shielding

## Two molecules:



Two-body Hamiltonian:

$$\hat{H}_2 = \sum_{j=1,2} \hat{h}_j + V(\mathbf{r}_1 - \mathbf{r}_2)$$

Kinetic energy:

$$\hat{h}_j = -\hbar^2 \nabla_j^2 / (2M) + \hat{h}_{\text{in}}(j)$$

Dipolar interaction:

$$\begin{aligned} V(\mathbf{r}) &= \frac{d^2}{4\pi\epsilon_0 r^3} \left[ \hat{\mathbf{d}}_1 \cdot \hat{\mathbf{d}}_2 - 3(\hat{\mathbf{d}}_1 \cdot \hat{\mathbf{r}})(\hat{\mathbf{d}}_2 \cdot \hat{\mathbf{r}}) \right] \\ &= -8\sqrt{\frac{2}{15}} \pi^{3/2} \frac{d^2}{4\pi\epsilon_0 r^3} \sum_{m=-2}^2 Y_{2m}^*(\hat{\mathbf{r}}) \Sigma_{2,m}, \end{aligned}$$

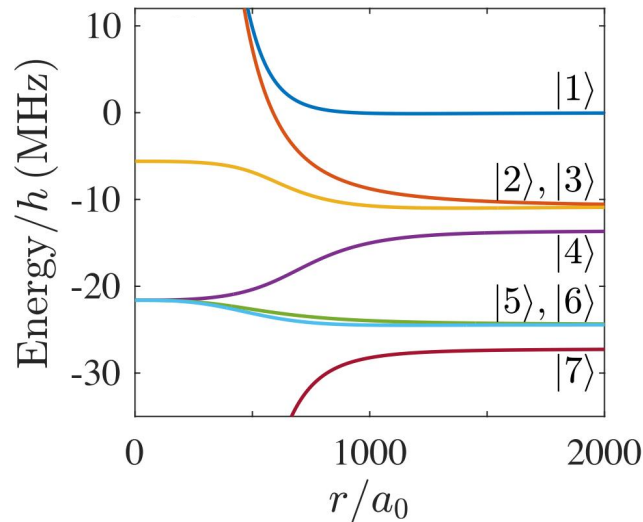
Spherical tensors:

$$\Sigma_{2,0} = \frac{1}{\sqrt{6}} (\hat{d}_1^+ \hat{d}_2^- + \hat{d}_1^- \hat{d}_2^+ + 2\hat{d}_1^0 \hat{d}_2^0), \quad \Sigma_{2,\pm 1} = \frac{1}{\sqrt{2}} (\hat{d}_1^\pm \hat{d}_2^0 + \hat{d}_1^0 \hat{d}_2^\pm), \quad \text{and } \Sigma_{2,\pm 2} = \hat{d}_1^\pm \hat{d}_2^\pm$$

The parity symmetry reduces the Hamiltonian to a 7D matrix in the symmetric space.

# Microwave shielding

**BO approximation:**



$$|1\rangle = |+, +\rangle,$$

$$|2\rangle = |+, 0\rangle_s, |3\rangle = |+, \xi_-\rangle_s,$$

$$|4\rangle = |+, -\rangle_s,$$

$$|5\rangle = |-, 0\rangle_s, |6\rangle = |-, \xi_-\rangle_s,$$

$$|7\rangle = |-, -\rangle$$

$$|1\rangle \text{ mostly couples to } |c\rangle = \frac{1}{\sqrt{\cos^2 \theta + 1}} (\sqrt{2} \cos \theta |2\rangle_s + \sin \theta e^{i\varphi} |3\rangle_s).$$

The 7D Hamiltonian is further reduced to a 2D matrix,  
which can be solved analytically!

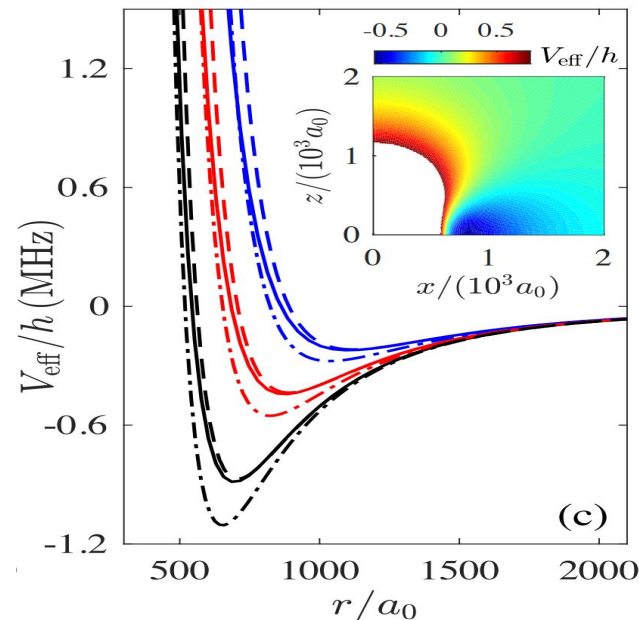
# Microwave shielding

## Effective potential:

$$V_{\text{eff}}(\mathbf{r}) = \frac{C_3}{r^3} [3 \cos^2 \theta - 1 + 3\mathcal{F}_\xi(\varphi) \sin^2 \theta] \\ + \frac{C_6}{r^6} \sin^2 \theta \{1 - \mathcal{F}_\xi^2(\varphi) + [1 - \mathcal{F}_\xi(\varphi)]^2 \cos^2 \theta\}.$$

$$\mathcal{F}(\xi, \varphi) = \sin 2\xi \cos 2\varphi$$

$$C_3 = d^2 / [48\pi\epsilon_0(1 + \delta_r^2)] \quad C_6 = d^4 / [128\pi^2\epsilon_0^2\Omega(1 + \delta_r^2)^{3/2}]$$



# Algorithm

## 7-Channel Scatterings:

$$\sum_{\nu'=1}^7 \left( -\frac{\hbar^2 \nabla^2}{M} \delta_{\nu\nu'} + V_{\nu\nu'} \right) \psi_{\nu'}(\mathbf{r}) = \frac{\hbar^2 k_\nu^2}{M} \psi_\nu(\mathbf{r}),$$

Angular momentum basis:

$$\psi_\nu(\mathbf{r}) = \sum_{lm} Y_{lm}(\hat{\mathbf{r}}) \phi_{\nu lm}(r)/r$$

Log-derivative method:

B. Johnson, Journal of Computational Physics 13, 445 (1973).

$$\partial_r \mathcal{Y}(r) = -\mathcal{V}(r) - \mathcal{Y}^2(r) \quad \mathcal{Y}(r) = \partial_r \phi(r) \phi^{-1}(r)$$

Matching the asymptotic solution:

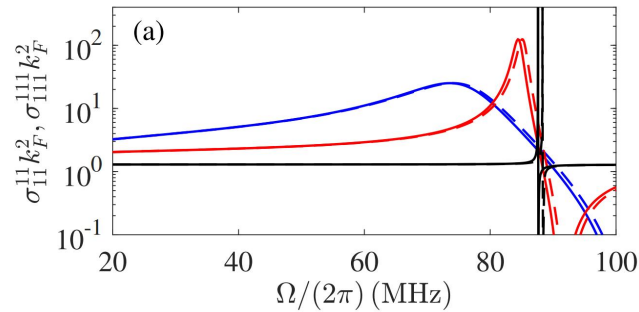
$$\phi_{\nu'l'm',\nu lm}^a(r) = \mathbf{J}_l(k_\nu r) \delta_{\nu\nu'} \delta_{ll'} \delta_{mm'} + \mathbf{N}_{l'}(k_{\nu'} r) \boxed{K_{\nu'l'm',\nu lm}},$$

$$f_{\nu lm}^{\nu'l'm'} = i \frac{1}{\sqrt{k_{\nu'}}} \left( \frac{1}{K + i} K \right)_{\nu'l'm',\nu lm} \frac{1}{\sqrt{k_\nu}} \quad \sigma_{\nu lm}^{\nu'l'm'} = 4\pi \left| f_{\nu lm}^{\nu'l'm'} \right|^2$$

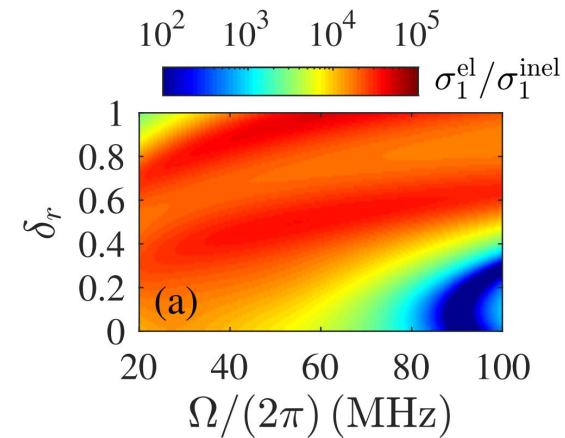
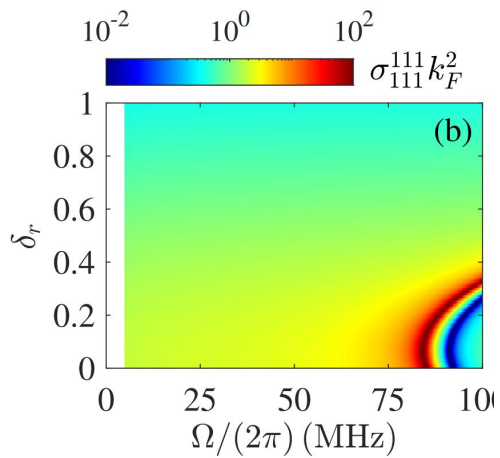
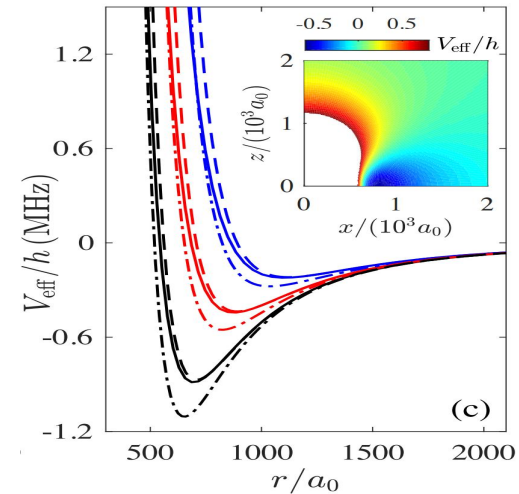


# Scattering cross sections

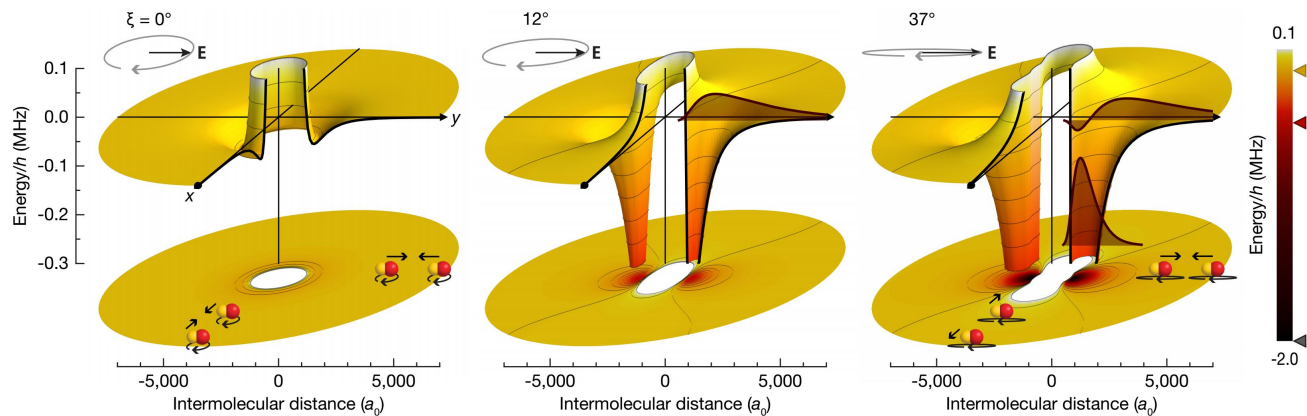
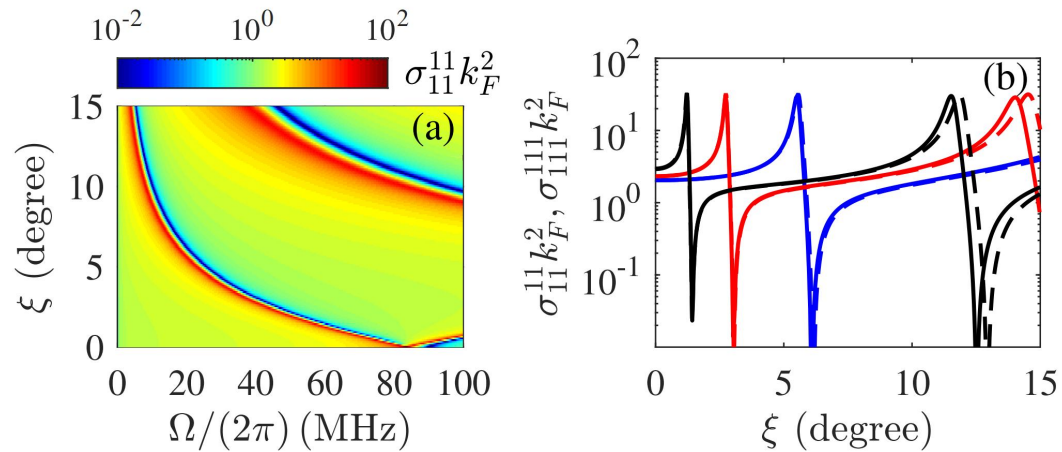
**p-wave:**



$k_1/k_F = 0.04$  (black lines),  
 $0.45$  (red lines), and  $1$  (blue lines).

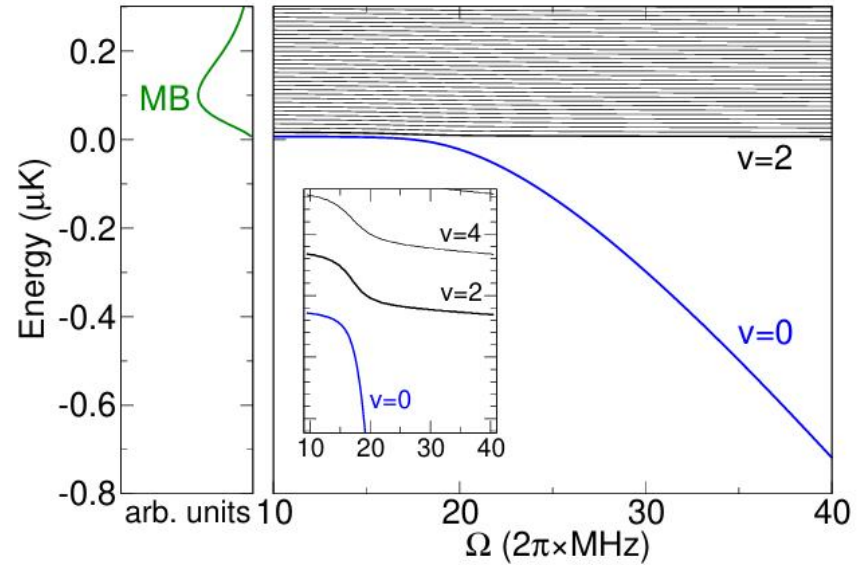
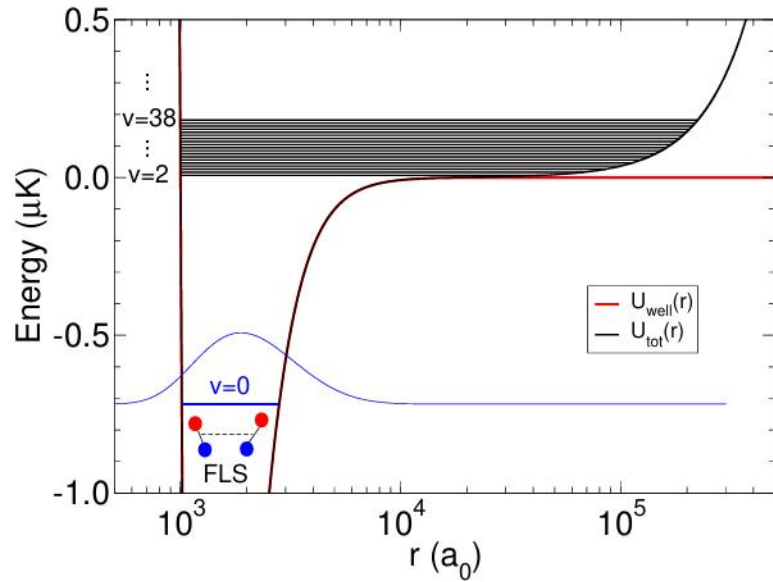


# Finite elliptic angles



$$V_{\text{eff}}(\mathbf{r}) \sim -C_3(1 + 3 \sin 2\xi)/r^3$$

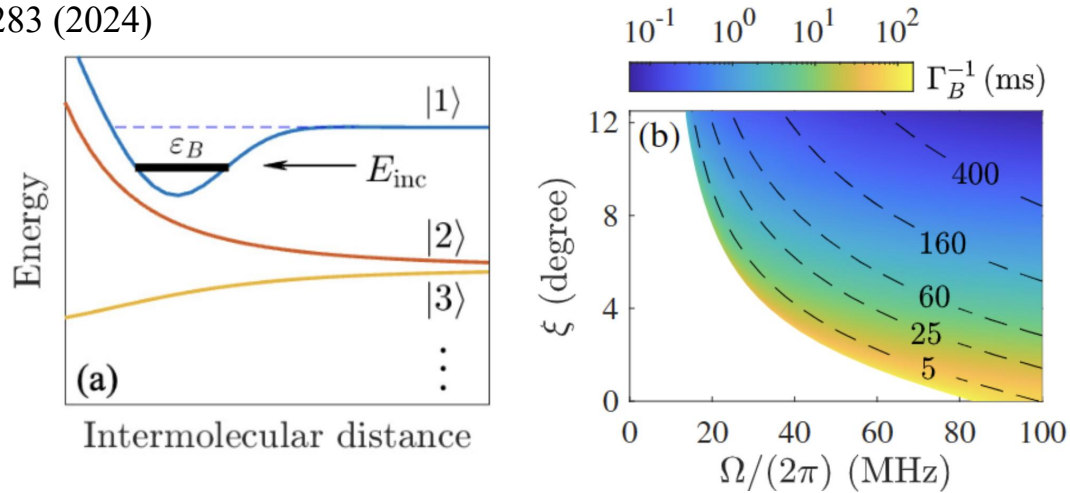
# Tetramer bound states



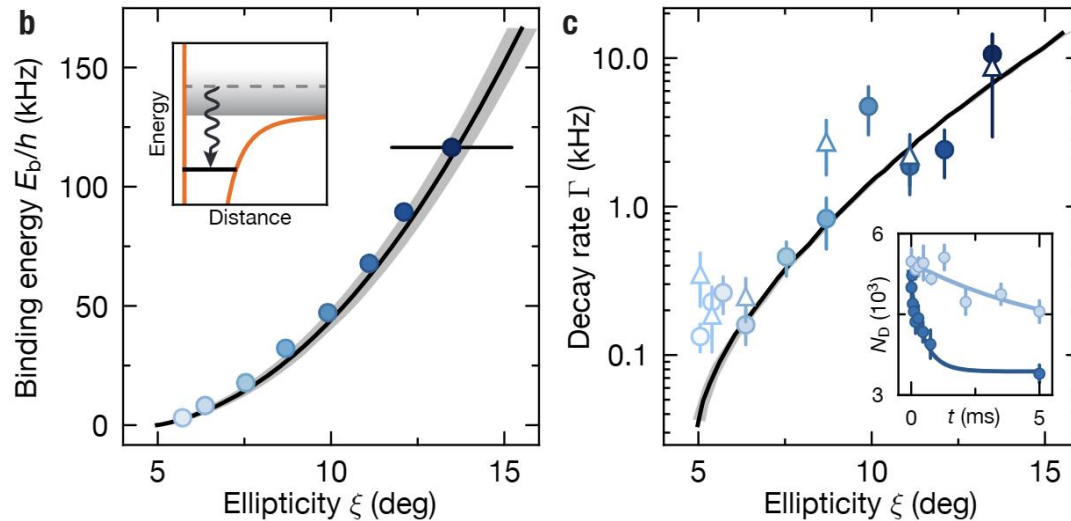
Goulven Quémener, John L. Bohn, and James F. E. Croft,  
Phys. Rev. Lett. **131**, 043402 (2023)

# Tetramer bound states

X. Y. Chen et al., Nature **626**, 283 (2024)



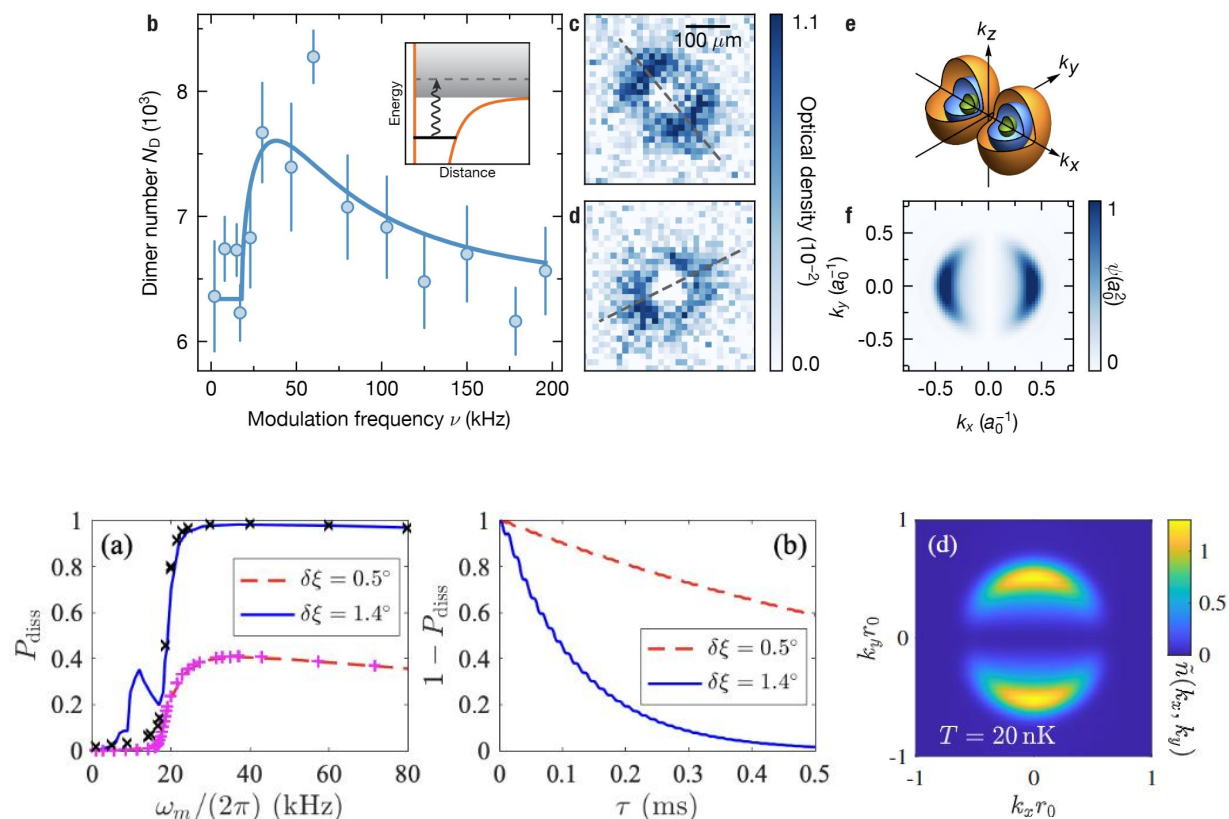
$$\sigma(E_{\text{inc}}) = \frac{2\pi}{k_2^2} |ig^2G(E_{\text{inc}}) + S_{\text{bg}} - 1|^2$$



# Modulational dissociation

$$V_{\text{eff}}(\mathbf{r}) = \frac{C_3}{r^3} [3 \cos^2 \theta - 1 + 3 \mathcal{F}_\xi(\varphi) \sin^2 \theta] + \frac{C_6}{r^6} \sin^2 \theta \{1 - \mathcal{F}_\xi^2(\varphi) + [1 - \mathcal{F}_\xi(\varphi)]^2 \cos^2 \theta\}.$$

$$H_{\text{MD}} \approx (\varepsilon_B + \omega_m) |\psi_B\rangle \langle \psi_B| + \int d\mathbf{k} \frac{k^2}{M} |\psi_{\mathbf{k}}\rangle \langle \psi_{\mathbf{k}}| + \int \frac{d\mathbf{k}}{(2\pi)^{3/2}} g_{\mathbf{k}} |\psi_{\mathbf{k}}\rangle \langle \psi_B| + \text{h.c.},$$

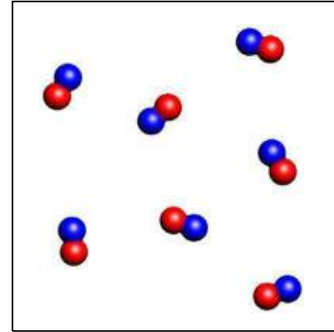
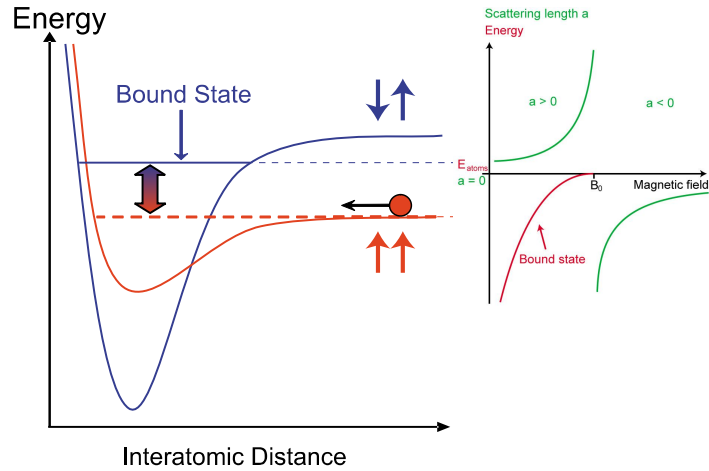


X. Y. Chen et al., Nature 626, 283 (2024)  
 F. Deng *et al.*, arXiv:2405.13645

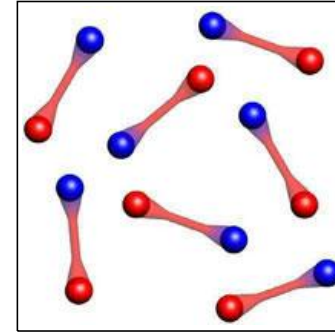
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## Fermionic MSPMs

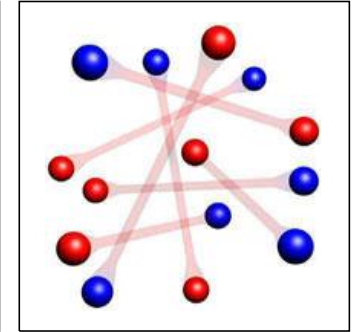
# BCS-BEC crossover in atomic gases



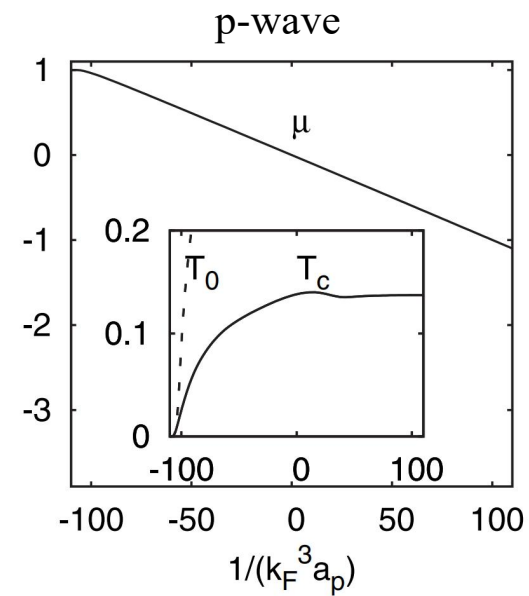
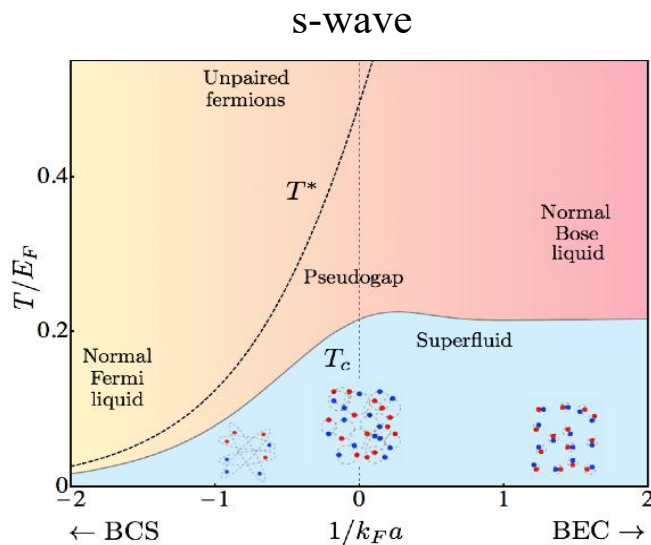
BEC of Molecules



Crossover Superfluid

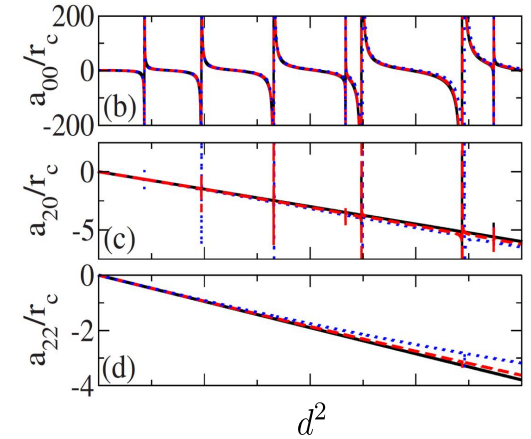


BCS state

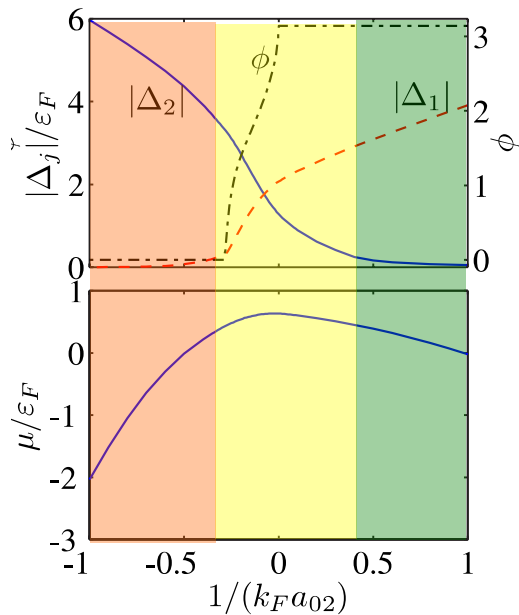


# BCS-BEC crossover in dipolar gases

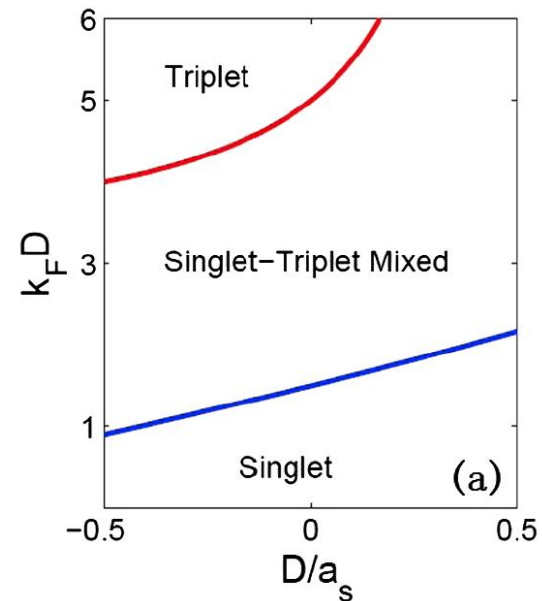
$$-\frac{M\Delta_j}{16\pi^2\lambda_j} = \sum_{j'} \Delta_{j'} \int \frac{d^3p}{(2\pi)^3} W_{j'}^*(\hat{\mathbf{p}}) \left( \frac{1}{2E_{\mathbf{p}}} - \frac{M}{p^2} \right) W_{j'}(\hat{\mathbf{p}})$$



Kanjilal and Blume, PRA 78, 040703(R) (2008)



Shi, Zou, Hu, Sun, & Yi, PRL 110, 045301 (2013).



Qi, Shi, and Zhai, PRL 110, 045302 (2013).



# BCS superfluid in MS molecules

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**Many-body Hamiltonian:**

$$\hat{H} = \int d^3\mathbf{r} \hat{\psi}^\dagger(\mathbf{r}) \left( -\frac{\hbar^2 \nabla^2}{2M} - \mu \right) \hat{\psi}(\mathbf{r}) + \frac{1}{2} \int d\mathbf{r} d\mathbf{r}' \hat{\psi}^\dagger(\mathbf{r}) \hat{\psi}^\dagger(\mathbf{r}') V_{\text{eff}}(\mathbf{r} - \mathbf{r}') \hat{\psi}(\mathbf{r}') \hat{\psi}(\mathbf{r}),$$

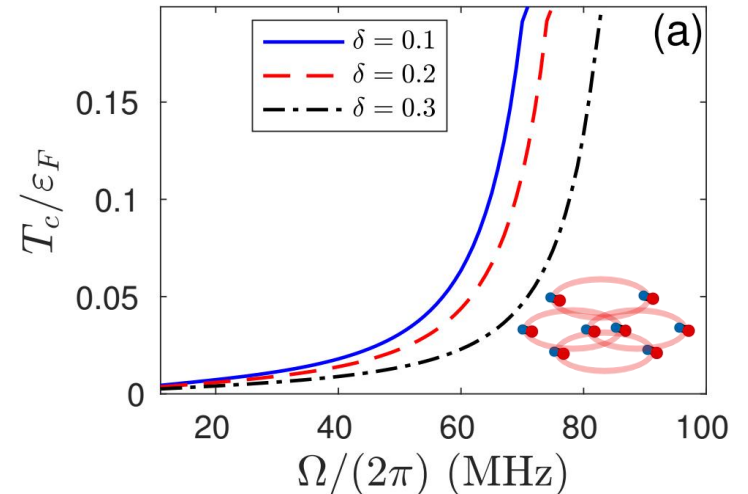
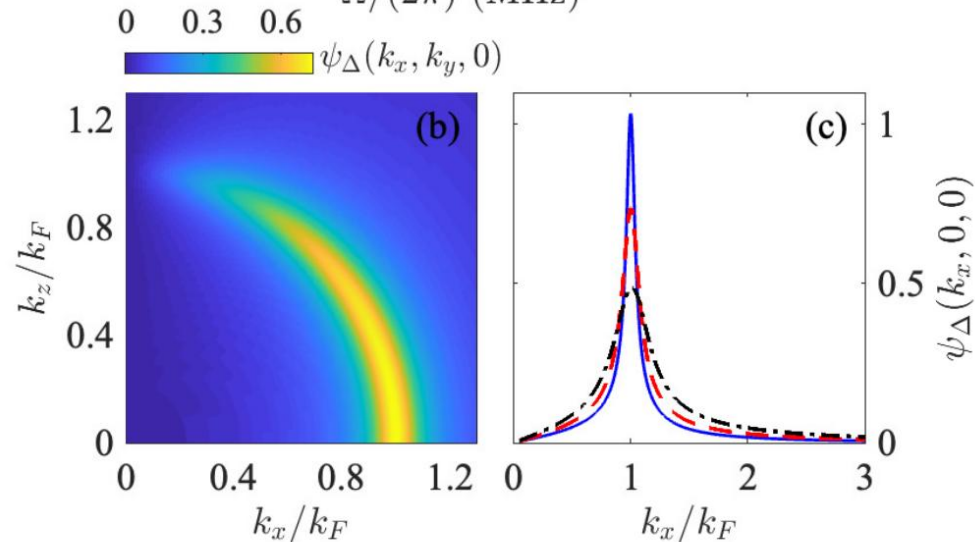
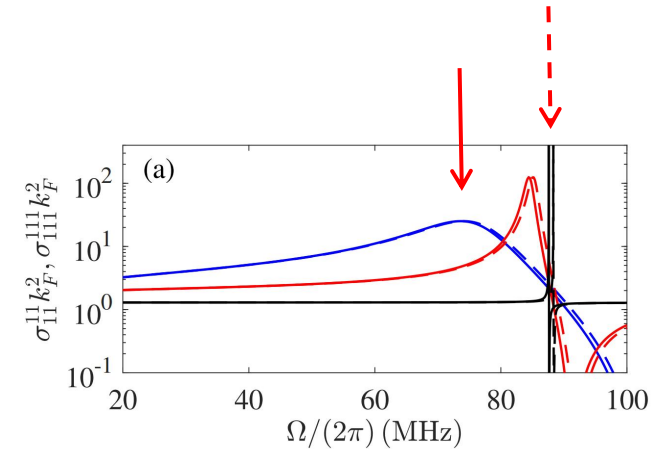
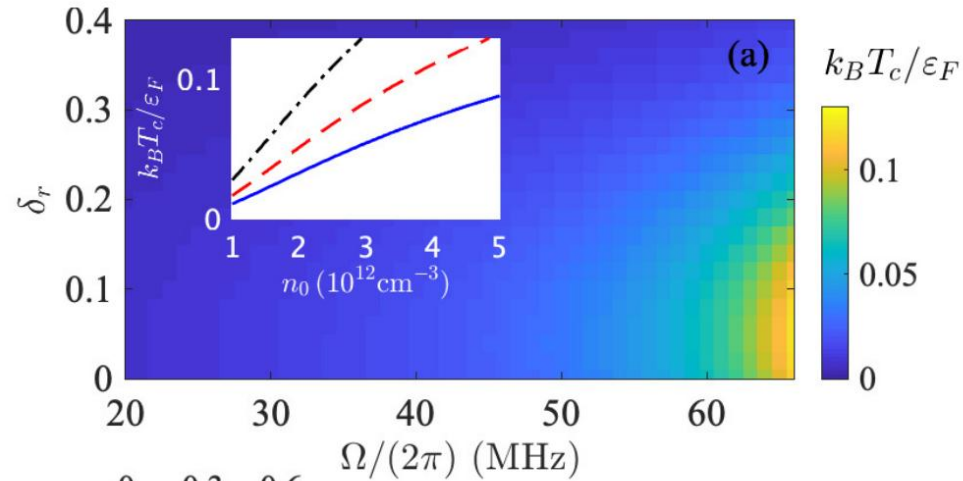
**Gap equation:**

$$\Delta_{lm}(k) = -\frac{2}{\pi} \sum_{l'} i^{l'-l} \int_0^\infty p^2 dp \tilde{V}_{ll',m}(k, p) \frac{\tanh(\beta_c \varepsilon_{\mathbf{p}}/2)}{2\varepsilon_{\mathbf{p}}} \Delta_{l'm}(p)$$

- (1) Renormalization free
- (2) Logarithmic discretization

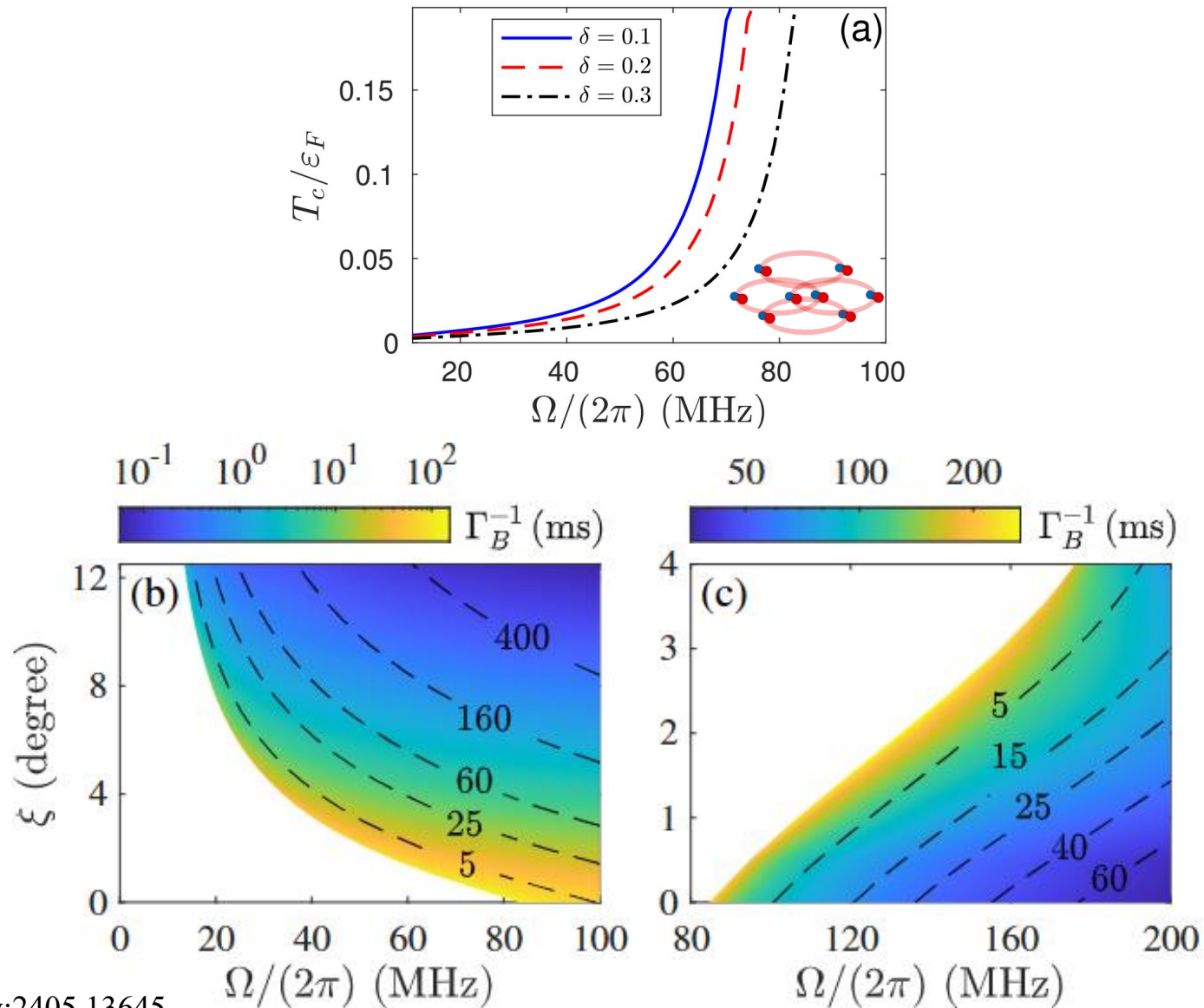
# BCS superfluid in MS molecules

$\Omega/(2\pi) = 28$  (solid line), 38 (dashed line), and 48 MHz (dash-dotted line)



$\Omega/(2\pi) = 48$  (solid line), 58 (dashed line), and 66 MHz (dash-dotted line)

# Extending lifetimes of tetramers

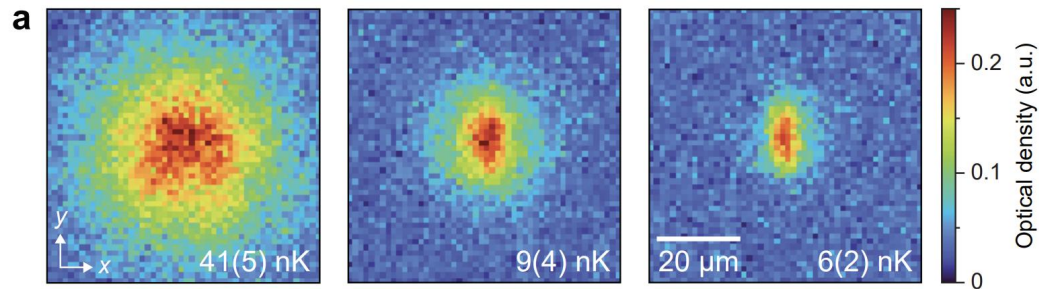


**LiRb: 10.4s**

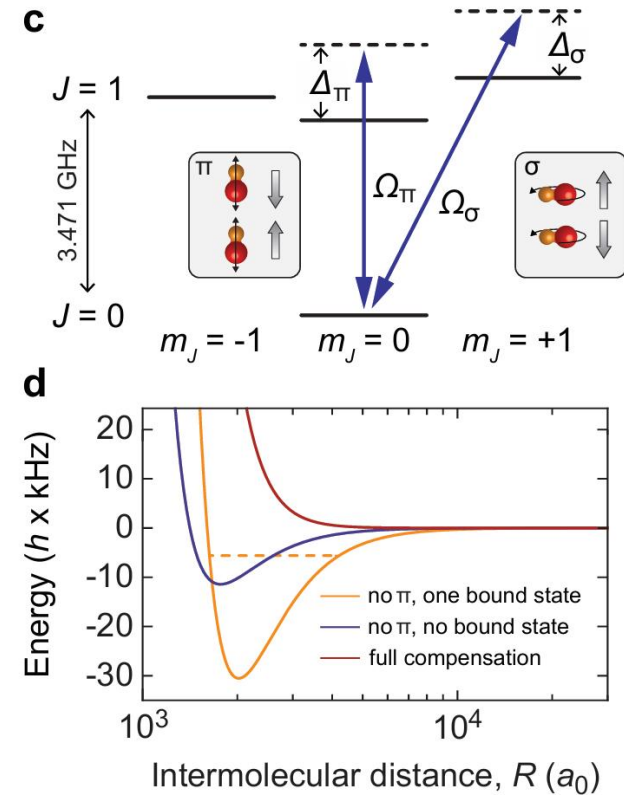
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## Bosonic MSPMs

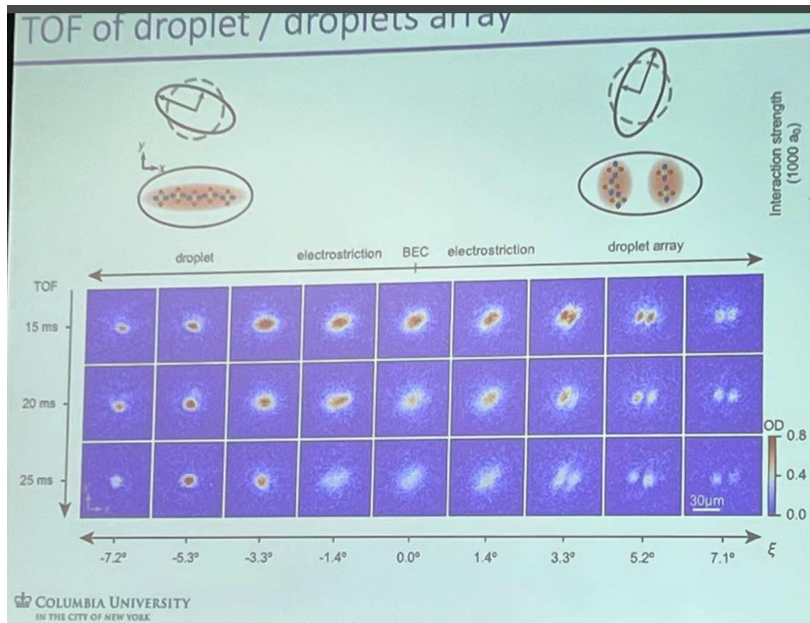
# Bose molecules



N. Bigagli *et al.*, Nature 631, 289 (2024).



# Bose molecules



## Conclusions

From molecular BEC, entering the dipolar phases of molecules

- Stronger dipole-dipole interaction
- Control over the interaction anisotropy

Open questions:

- Validity of first order quantum fluctuation term?
- The role of inner hard-wall interaction? 3D crystal?
- droplet at  $a_s < 0$ ? Is  $a_s$  still a good parameter?

$$g_{qf} = \frac{32}{3\sqrt{\pi}} g \sqrt{a_s^3} Q_5(\epsilon_{dd})$$

Gorshkov, P. *et al.* PRL 101, 073201 (2008)

More exciting quantum phases with dipolar molecules are within reach!

COLUMBIA UNIVERSITY  
IN THE CITY OF NEW YORK

**DAMOP S. Will's talk**

# Bose molecules

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## **Hamiltonian:**

W. Jin *et al.*, arXiv:2406.06412

$$H = H_0 + H_{\text{int}},$$

$$H_0 = \int d\mathbf{r} \left[ \frac{1}{2M} \nabla \hat{\psi}^\dagger(\mathbf{r}) \nabla \hat{\psi}(\mathbf{r}) + V(\mathbf{r}) \hat{\psi}^\dagger(\mathbf{r}) \hat{\psi}(\mathbf{r}) \right]$$

$$H_{\text{int}} = \frac{1}{2} \int d\mathbf{r} d\mathbf{r}' U(\mathbf{r} - \mathbf{r}') \hat{\psi}^\dagger(\mathbf{r}) \hat{\psi}^\dagger(\mathbf{r}') \hat{\psi}(\mathbf{r}') \hat{\psi}(\mathbf{r}),$$

$$U(\mathbf{r}) = \frac{C_3}{r^3} (3 \cos^2 \theta - 1) + \frac{C_6}{r^6} \sin^2 \theta (1 + \cos^2 \theta),$$

$$U_{\text{pp}}(\mathbf{r}) = \frac{4\pi \hbar^2 a_s}{M} \delta(\mathbf{r}) + \frac{C_3}{r^3} (3 \cos^2 \theta - 1).$$

# Bose molecules

## Hamiltonian:

W. Jin *et al.*, arXiv:2406.06412

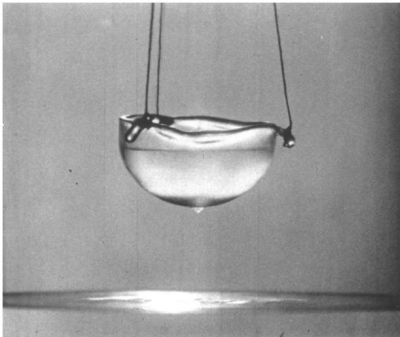
$$H = H_0 + H_{\text{int}},$$

$$H_0 = \int d\mathbf{r} \left[ \frac{1}{2M} \nabla \hat{\psi}^\dagger(\mathbf{r}) \nabla \hat{\psi}(\mathbf{r}) + V(\mathbf{r}) \hat{\psi}^\dagger(\mathbf{r}) \hat{\psi}(\mathbf{r}) \right]$$

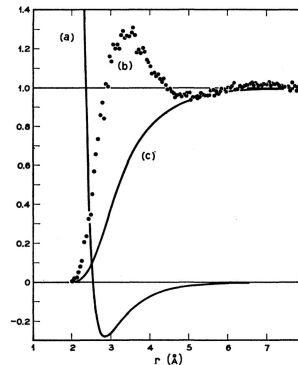
$$H_{\text{int}} = \frac{1}{2} \int d\mathbf{r} d\mathbf{r}' U(\mathbf{r} - \mathbf{r}') \hat{\psi}^\dagger(\mathbf{r}) \hat{\psi}^\dagger(\mathbf{r}') \hat{\psi}(\mathbf{r}') \hat{\psi}(\mathbf{r}),$$

$$U(\mathbf{r}) = \frac{C_3}{r^3} (3 \cos^2 \theta - 1) + \frac{C_6}{r^6} \sin^2 \theta (1 + \cos^2 \theta),$$

$$U_{\text{PP}}(\mathbf{r}) = \frac{4\pi\hbar^2 a_s}{M} \delta(\mathbf{r}) + \frac{C_3}{r^3} (3 \cos^2 \theta - 1). \quad \textbf{Inapplicable!!!}$$



Liquid Helium



W. L. McMillan,  
Phys. Rev. **138** A442 (1965)

$$g_2(\mathbf{r}, \mathbf{r}') = \frac{\langle \Psi | \hat{\psi}^\dagger(\mathbf{r}) \hat{\psi}^\dagger(\mathbf{r}') \hat{\psi}(\mathbf{r}') \hat{\psi}(\mathbf{r}) | \Psi \rangle}{n(\mathbf{r})n(\mathbf{r}')},$$



# Bose molecules

W. Jin *et al.*, arXiv:2406.06412

$$U(\mathbf{r}) = \frac{C_3}{r^3} (3 \cos^2 \theta - 1) + \frac{C_6}{r^6} \sin^2 \theta (1 + \cos^2 \theta),$$

## Variational ansatz:

$$|\Psi\rangle = e^{-\alpha^2/2} e^{\alpha \hat{b}^\dagger} |0\rangle$$

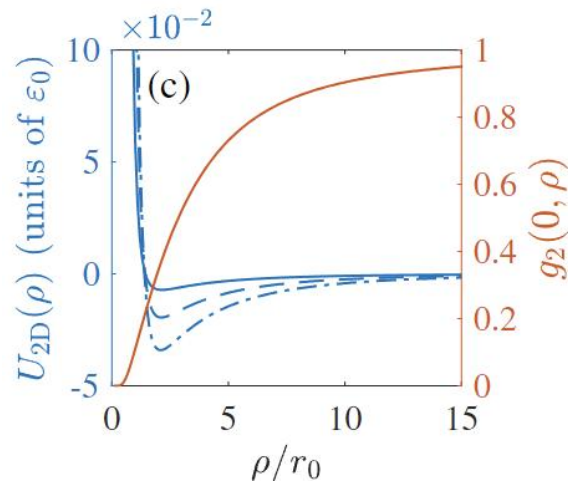
$$|\Psi\rangle = S |\alpha\rangle / \sqrt{\mathcal{N}},$$

$$\hat{b} = \int d\mathbf{r} \phi_0(\mathbf{r}) \hat{\psi}(\mathbf{r})$$

$$S = \exp\left[\frac{1}{2} \int d\mathbf{r} d\mathbf{r}' \chi(\mathbf{r}, \mathbf{r}') \hat{\psi}^\dagger(\mathbf{r}) \hat{\psi}^\dagger(\mathbf{r}') \hat{\psi}(\mathbf{r}') \hat{\psi}(\mathbf{r})\right]$$

TS, E. Demler, and J. I. Cirac, *Annals of Physics* **390**, 245 (2018).

TS, E. Demler, and J. I. Cirac, *PRL* **125**, 180602 (2020).



$$|\Psi_N\rangle = \frac{1}{\sqrt{\mathcal{N}}} \int D[\mathbf{r}] \prod_{i<j(=1)}^N J(\mathbf{r}_i, \mathbf{r}_j) \prod_{j=1}^N \phi_0(\mathbf{r}_j) \hat{\psi}^\dagger(\mathbf{r}_j) |0\rangle$$

R. Jastrow, *Phys. Rev.* **98**, 1479 (1955).

# Bose molecules

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## **Ground state energy:**

$$E = \frac{1}{2M} \int d\mathbf{r} \nabla \phi(\mathbf{r}) \nabla \phi(\mathbf{r}) - \int d\mathbf{r} d\mathbf{r}' \frac{1}{4M} \nabla \ln J(\mathbf{r}, \mathbf{r}') \nabla \ln n(\mathbf{r}) G_2(\mathbf{r}, \mathbf{r}') \\ - \int d\mathbf{r} d\mathbf{r}' \frac{1}{4M} \nabla^2 \ln J(\mathbf{r}, \mathbf{r}') G_2(\mathbf{r}, \mathbf{r}') + \int d\mathbf{x} V(\mathbf{r}) n(\mathbf{r}) + \frac{1}{2} \int d\mathbf{r} d\mathbf{r}' U(\mathbf{r} - \mathbf{r}') G_2(\mathbf{r}, \mathbf{r}'),$$

## **Cluster expansions:**

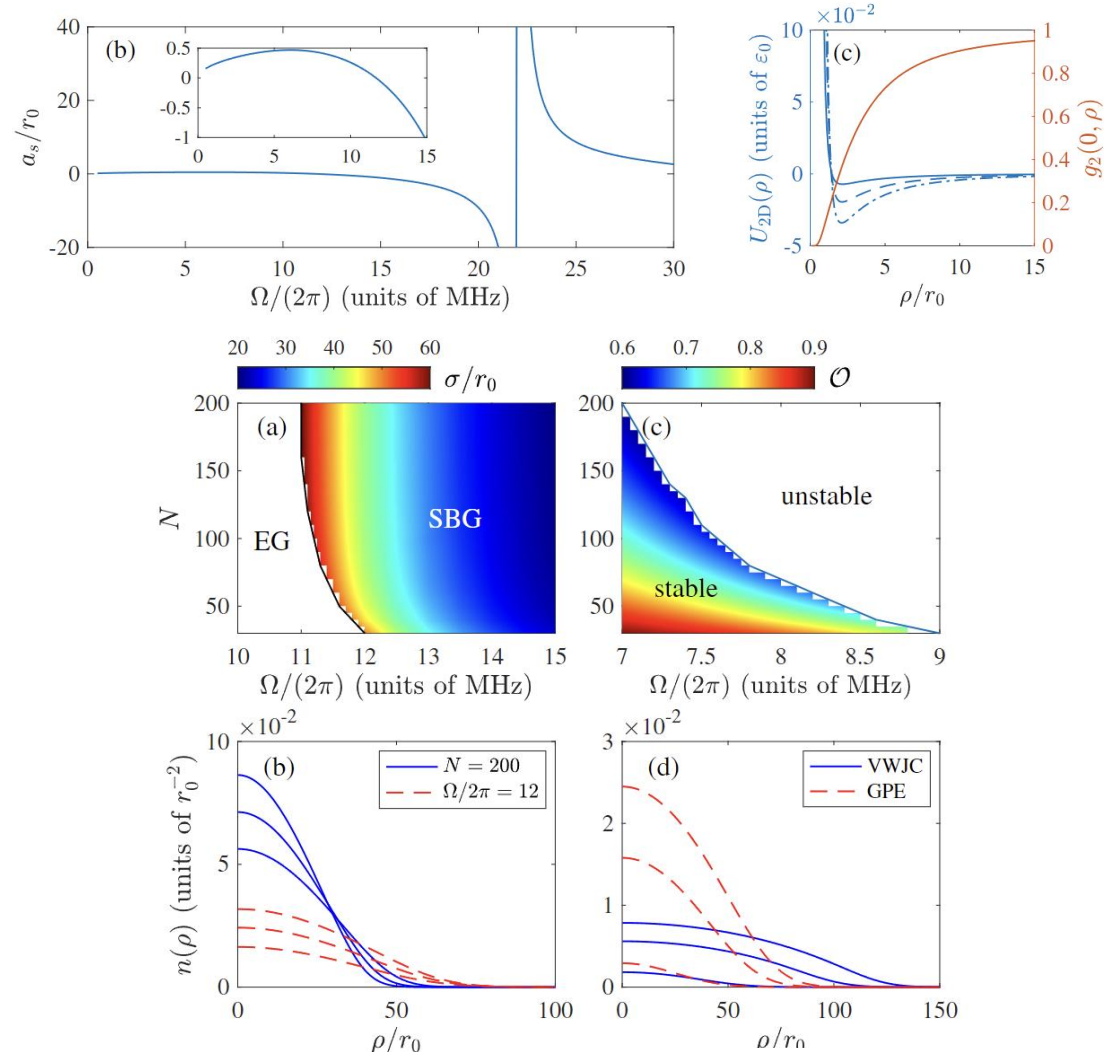
*J. B. Aviles, Annals of Physics 5, 251 (1958)*

$$G_2(\mathbf{r}, \mathbf{r}') = n(\mathbf{r}) n(\mathbf{r}') J^2(\mathbf{r}, \mathbf{r}') \left[ 1 + \int d\mathbf{r}_1 f(\mathbf{r}, \mathbf{r}_1) n(\mathbf{r}_1) F(\mathbf{r}_1, \mathbf{r}') \right].$$

$$f(\mathbf{r}_i, \mathbf{r}_j) = J^2(\mathbf{r}_i, \mathbf{r}_j) - 1$$

# Bose molecules

## Phase diagrams:

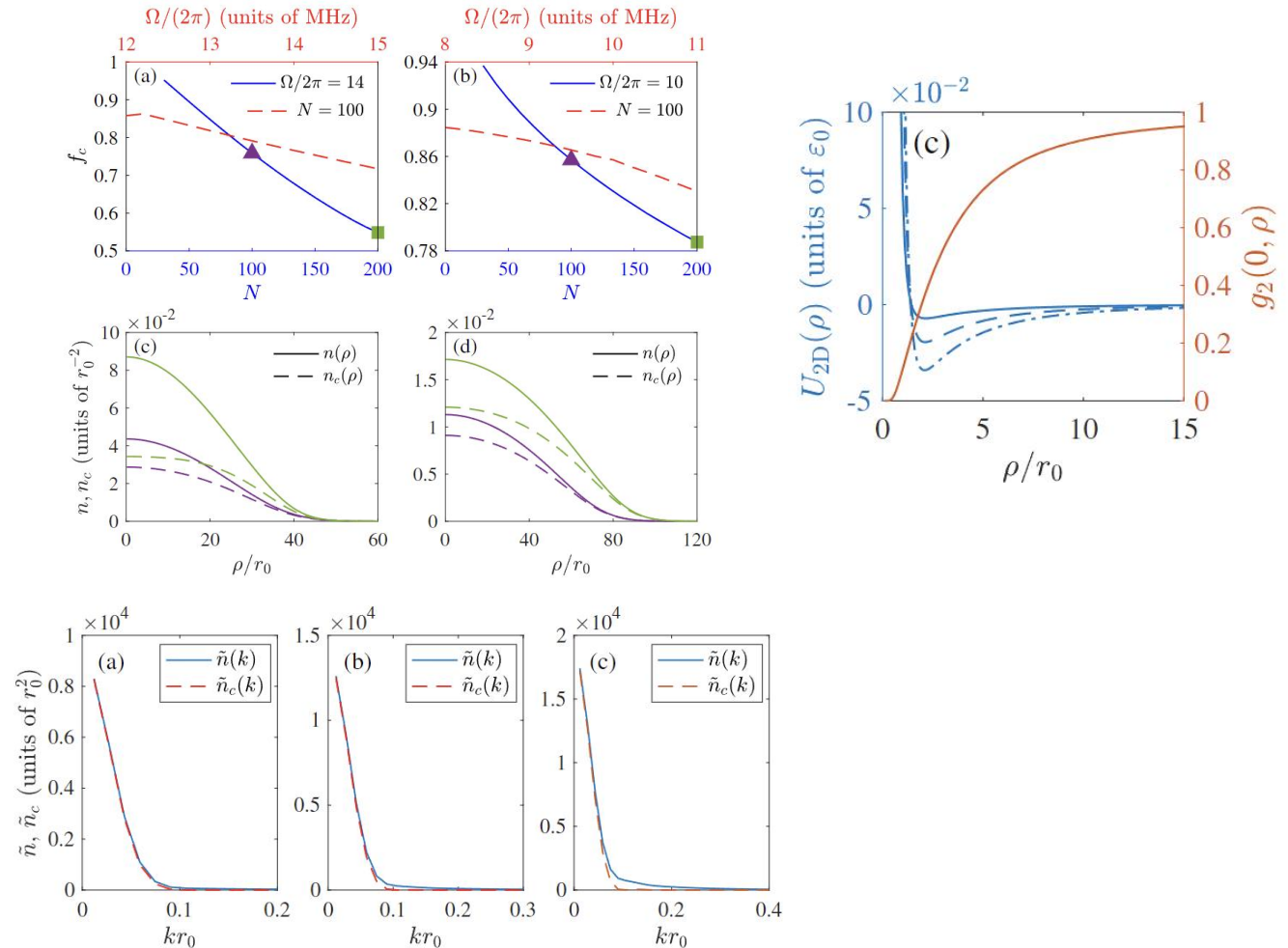


$$U(\mathbf{r}) = \frac{C_3}{r^3}(3 \cos^2 \theta - 1) + \frac{C_6}{r^6} \sin^2 \theta (1 + \cos^2 \theta), \quad U_{\text{pp}}(\mathbf{r}) = \frac{4\pi\hbar^2 a_s}{M} \delta(\mathbf{r}) + \frac{C_3}{r^3}(3 \cos^2 \theta - 1).$$

# Bose molecules

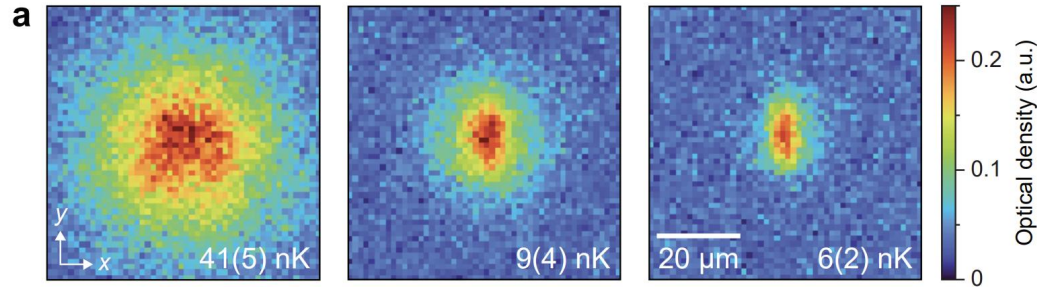
## Condensate fractions and momentum distributions

$$G_1(\rho, \rho') = \langle \Psi | \psi^\dagger(\rho') \psi(\rho) | \Psi \rangle = \sum_{\ell \geq 0} N_\ell \bar{\varphi}_\ell(\rho) \bar{\varphi}_\ell^*(\rho')$$

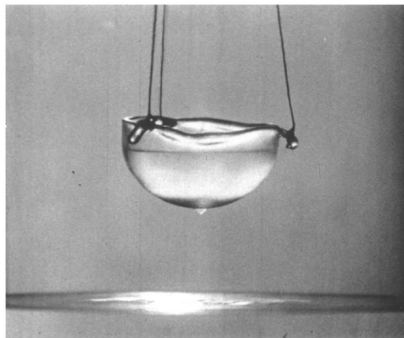
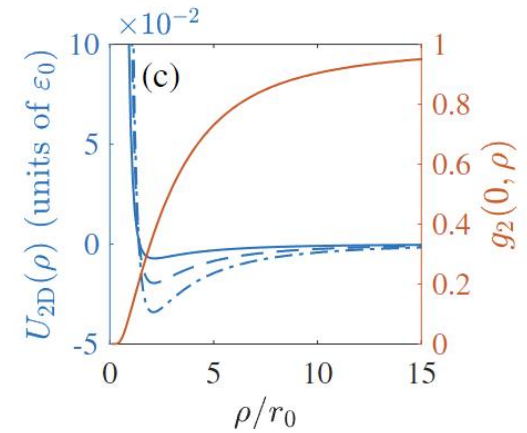


# Bose molecules

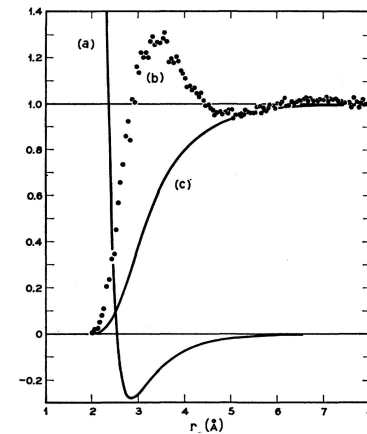
## Helium 4 v.s. MSPMs



N. Bigagli *et al.*, Nature 631, 289 (2024).



Liquid Helium



W. L. McMillan,  
Phys. Rev. **138** A442 (1965)

# Conclusion and Outlook

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## New paradigm in MS polar molecules beyond atoms and Helium

F. Deng, et al., PRL 130, 183001 (2023)

J. Lin et al., Phys. Rev. X **13**, 031032 (2023)

N. Bigagli *et al.*, Nature 631, 289 (2024).

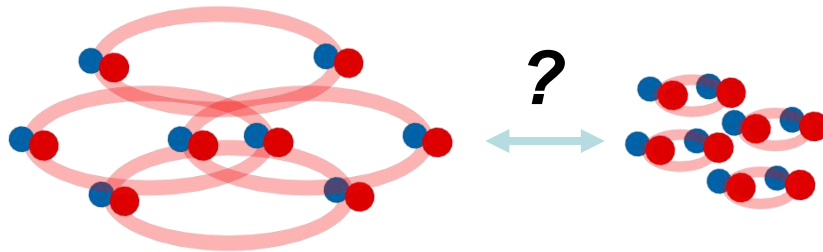
W. Jin *et al.*, arXiv:2406.06412

T. Langen *et al.*, arXiv:2407.09391

## Realization of tetramer BEC and BCS-BEC crossover

X. Y. Chen et al., Nature 626, 283 (2024)

F. Deng *et al.*, arXiv:2405.13645



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---

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Dr. Xinyu Luo (MPQ)

Dr. Xingyan Chen (MPQ)

NaK group in MPQ



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