

Low energy neutrino scattering in semiconductor detectors

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Yu-Feng Li and Shuo-yu Xia, JHEP 10 (2023) 021, Arxiv: 2211.11582

Yu-Feng Li and Shuo-yu Xia, Nucl.Phys.B 1006 (2024) 116632, Arxiv: 2310.05704

23rd Sep, 2024



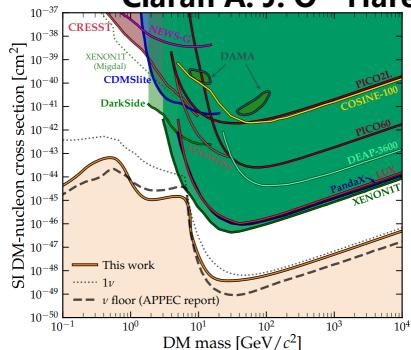
Outline

- Background
- Neutrino scattering in non-relativistic effective field theory (NR EFT)
- Low energy neutrino excitation in semiconductor
 - In-medium screening effect for neutrino-electron excitation
 - Migdal effect induced by neutrino-nucleus scattering
- Summary

Background

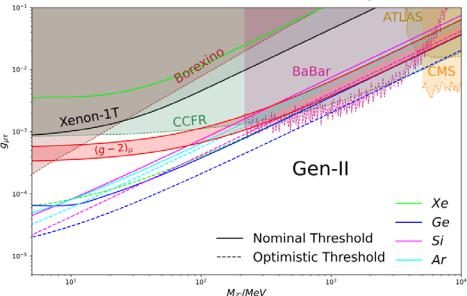
Rich physics

Ciaran A. J. O' Hare, 2022

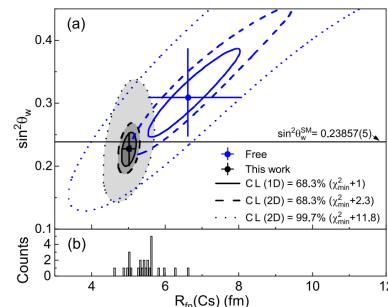
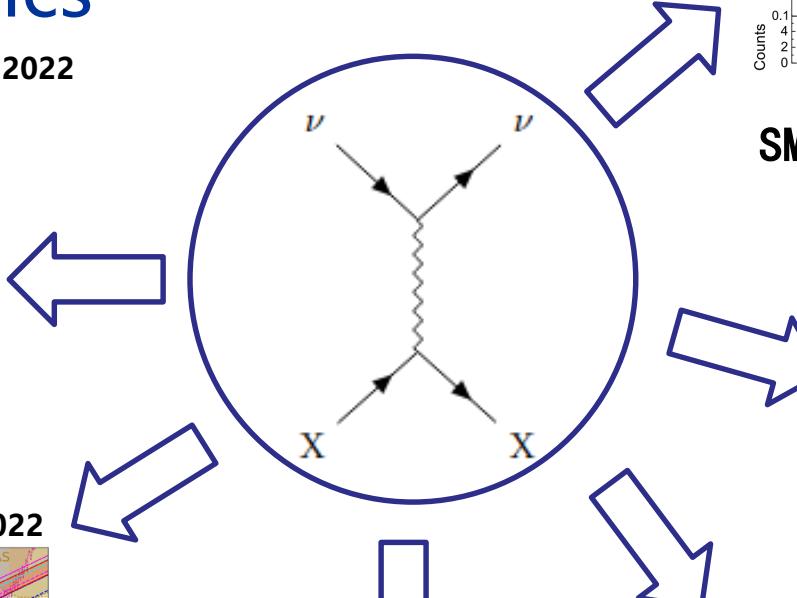


Neutrino Fog

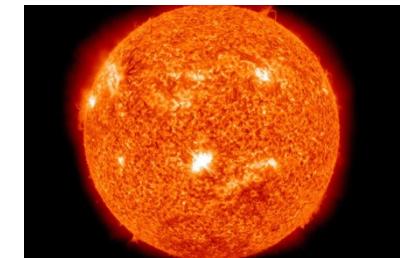
Y. F. Li and S. Y. Xia, 2022



Neutrino Non-Standard Interaction



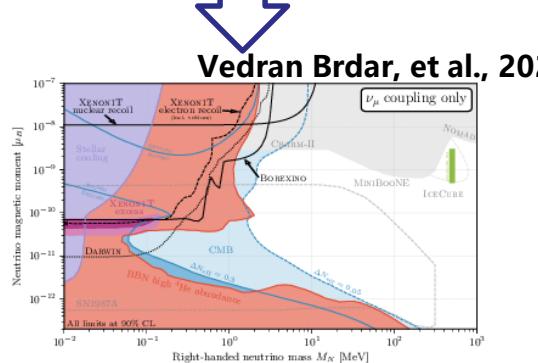
SM Precision Test



Solar Physics



Neutrino Magnetic Moments



Sterile Neutrino



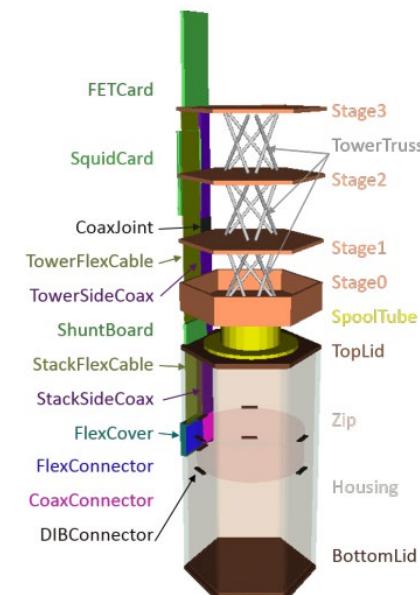
Background

■ Neutrino detection in semiconductor detectors

- Very low threshold
- Detection via phonons and electrons
- Several electron excitations



Barak, Liron and others,
Phys.Rev.Lett. 125 (2020) 17, 171802



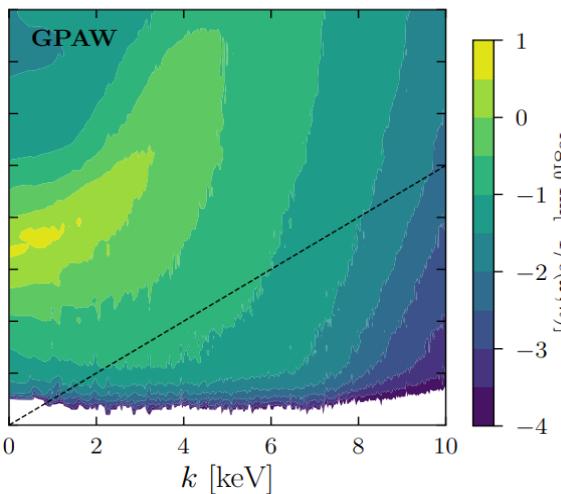
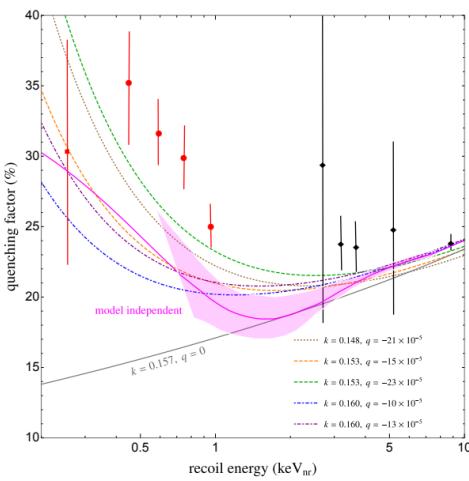
Agnese, R. and others,
Phys.Rev.D 95 (2017) 8, 082002



Background

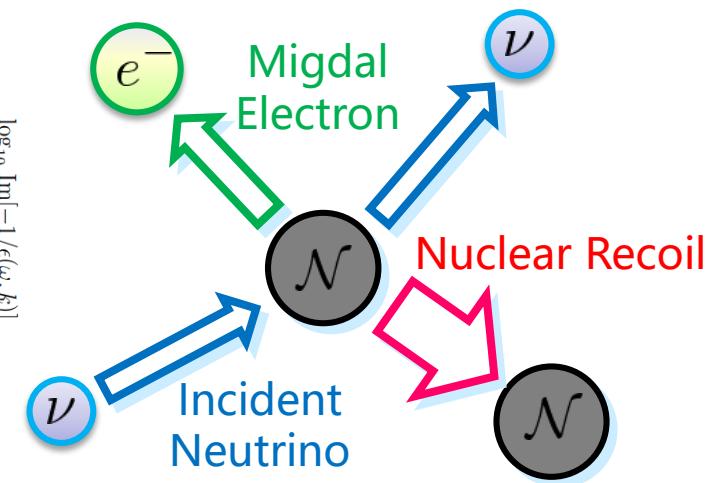
■ low energy excitations in semiconductor

- Quenching models at low energy
- Migdal effect in nuclear recoil
- Screening effect and other collective behavior
-



Liao, Jiajun and others ,
Phys.Rev.D 104 (2021) 1, 015005

Knapen, Simon and others ,
Phys.Rev.D 105 (2022) 1, 015014



Migdal Effect



Neutrino scattering in NR EFT

Appendix of Arxiv: 2211.11582

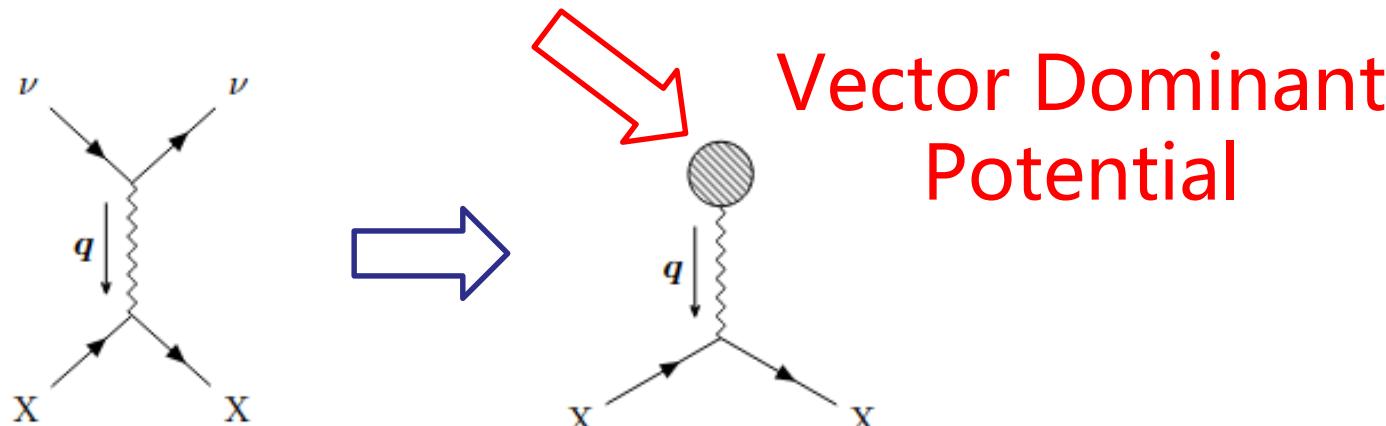
■ Neutrino electron scattering Lagrangian in SM

$$\mathcal{L}_{E\nu\text{ES}} = -i \frac{G_F}{\sqrt{2}} [\bar{\nu}_\alpha \gamma^\rho (1 - \gamma^5) \nu_\alpha] [\bar{\phi}_e \gamma_\rho (g_{\alpha;V} - g_{\alpha;A} \gamma^5) \phi_e]$$

■ Neutrino electron excitation in NR EFT

Andrea Mitridate, et al., JHEP 09 (2021) 123, Arxiv: 2106.12586

$$\mathcal{L}_{E\nu\text{ES}} = -i \sqrt{2} G_F g_{\alpha;V} \bar{\nu}_{\alpha,L} \gamma^0 \nu_{\alpha,L} \phi_+^\dagger \phi_+$$





Neutrino scattering in NR EFT

■ Contribution of neutrino magnetic moment

$$\mathcal{L}_{\text{mag}} = -i\mu_\nu \frac{m_\nu}{m_e} \frac{4\pi\alpha}{q^2} \left[\bar{\nu}_\alpha (1 - \gamma^5) i \frac{\sigma^{\mu\nu} q_\nu}{2m_\nu} \nu_\alpha \right] [\bar{\phi}_e \gamma_\mu \phi_e]$$

■ Neutrino magnetic moment (ER) in NR EFT

- Isotropic structure of the crystal
- Leading order interaction

$$\mathcal{L}_{\text{mag}} = i\mu_\nu \frac{E_\nu}{m_e} \frac{8\pi\alpha}{q^2} \bar{\nu}_{\alpha,\text{L}} \nu_{\alpha,\text{L}} \phi_+^\dagger \phi_+$$

Leading Order Potential



Neutrino scattering in NR EFT

■ Neutrino nucleus scattering in EFT

- Enhancement from coherent effect

↗ Coherent Effect

$$\mathcal{L}_{\text{eff}} = -i\sqrt{2}G_F [g_n N F_N^2(\vec{q}) + g_p Z F_Z^2(\vec{q})] \bar{\nu}_{e,L} \gamma^0 \nu_{e,L} \phi_+^\dagger \phi_+$$

■ Neutrino magnetic moment (NR) in NR EFT

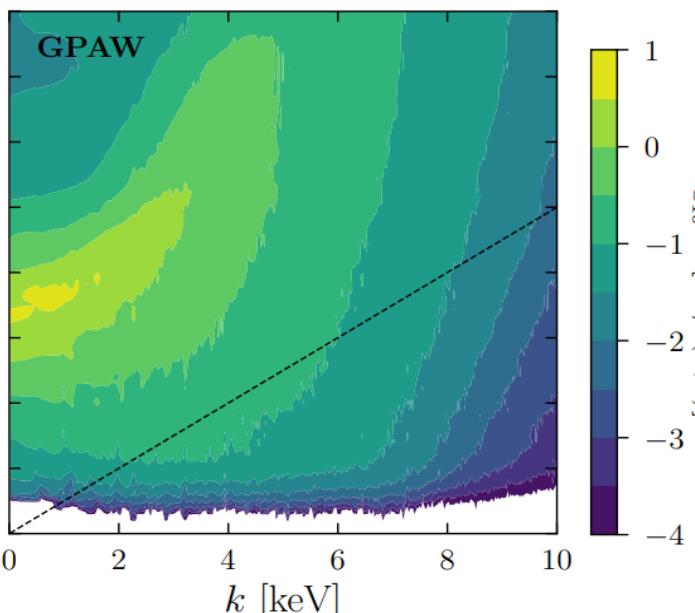
- Coherently coupled to Z

$$\mathcal{L}_{\text{mag}} = i Z F_Z^2(\vec{q}) \mu_\nu \frac{E_\nu}{m_e} \frac{8\pi\alpha}{q^2} \bar{\nu}_{\alpha,L} \nu_{\alpha,L} \phi_+^\dagger \phi_+$$

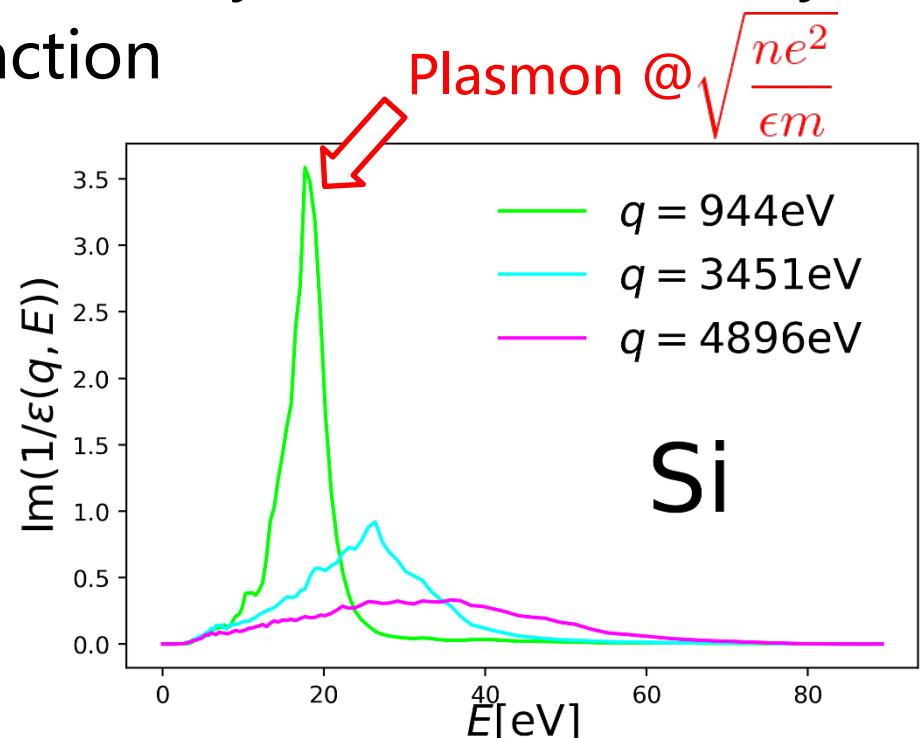


Screened neutrino excitation

- Describe excitation with energy loss function
 - Calculation based on Density Functional Theory
 - Inverse dielectric function



Knapen, Simon and others ,
Phys.Rev.D 105 (2022) 1, 015014



Data for silicon



Screened neutrino excitation

Yu-Feng Li and Shuo-yu Xia, JHEP 10 (2023) 021, Arxiv: 2211.11582

■ Correlation function in EELS

$$S^{\hat{H}_I^\dagger \hat{H}_I}(\omega) = 2V |V_{\text{cou}}(\mathbf{Q})| \text{Im} \left[\frac{-1}{\epsilon(\mathbf{Q}, \omega)} \right]$$

Coulomb Potential **Inverse Dielectric Function**

■ Dielectric response of neutrino

$$S^{\hat{H}_I^\dagger \hat{H}_I}(\omega) = 2V \frac{|V_{E\nu ES}|^2}{|V_{\text{cou}}(\mathbf{Q})|} \text{Im} \left[\frac{-1}{\epsilon(\mathbf{Q}, \omega)} \right]$$

Effective Potential of Neutrino



Screened neutrino excitation

■ Event rate with in-medium screening effect

$$R_{\text{SM}} = \frac{N_{\text{cell}} \Omega G_F^2 g_{\alpha;V}^2}{\pi \alpha} \int \Phi(E_\nu) dE_\nu$$

Not Sensitive to Plasmon

$$\times \sum_G \int_{1\text{BZ}} \frac{d^3q d\omega}{(2\pi)^3} |\mathbf{G} + \mathbf{q}| \text{Im} \left[\frac{1}{\epsilon_{G,G}(q, \omega)} \right] \times [\dots]$$

$$R_{\text{mag}} = 32 N_{\text{cell}} \Omega \mu_\nu^2 \frac{\pi \alpha}{m_e^2} \int \Phi(E_\nu) E_\nu^2 dE_\nu$$

Sensitive to Plasmon

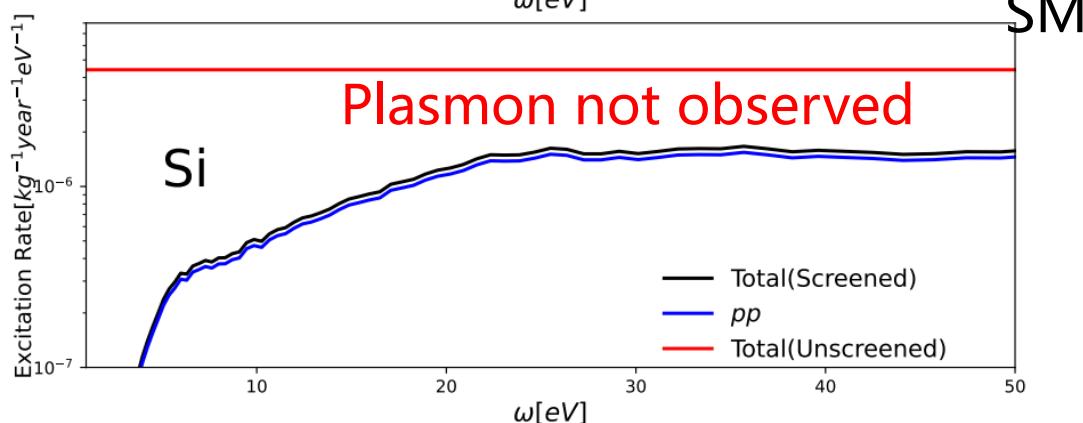
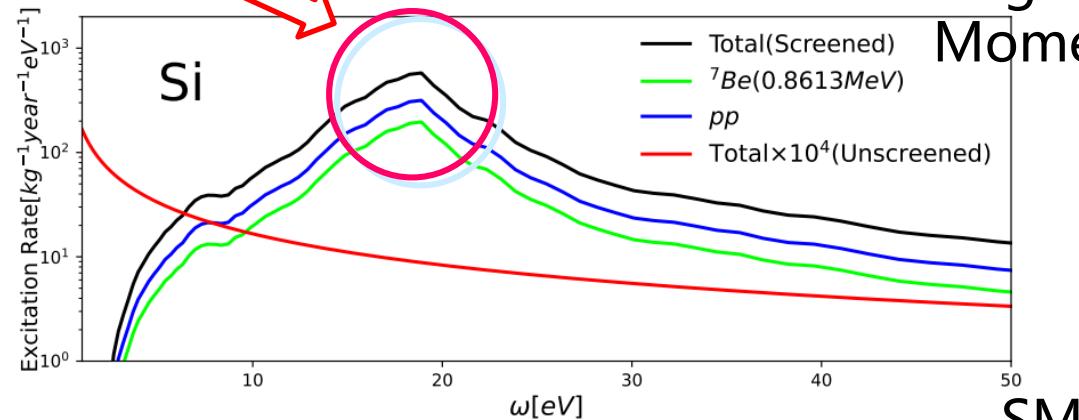
$$\times \sum_G \int_{1\text{BZ}} \frac{d^3q d\omega}{(2\pi)^3} \frac{1}{|\mathbf{G} + \mathbf{q}|^3} \text{Im} \left[\frac{1}{\epsilon_{G,G}(q, \omega)} \right] \times [\dots]$$



Screened neutrino excitation

■ Event rate with in-medium screening effect Significant

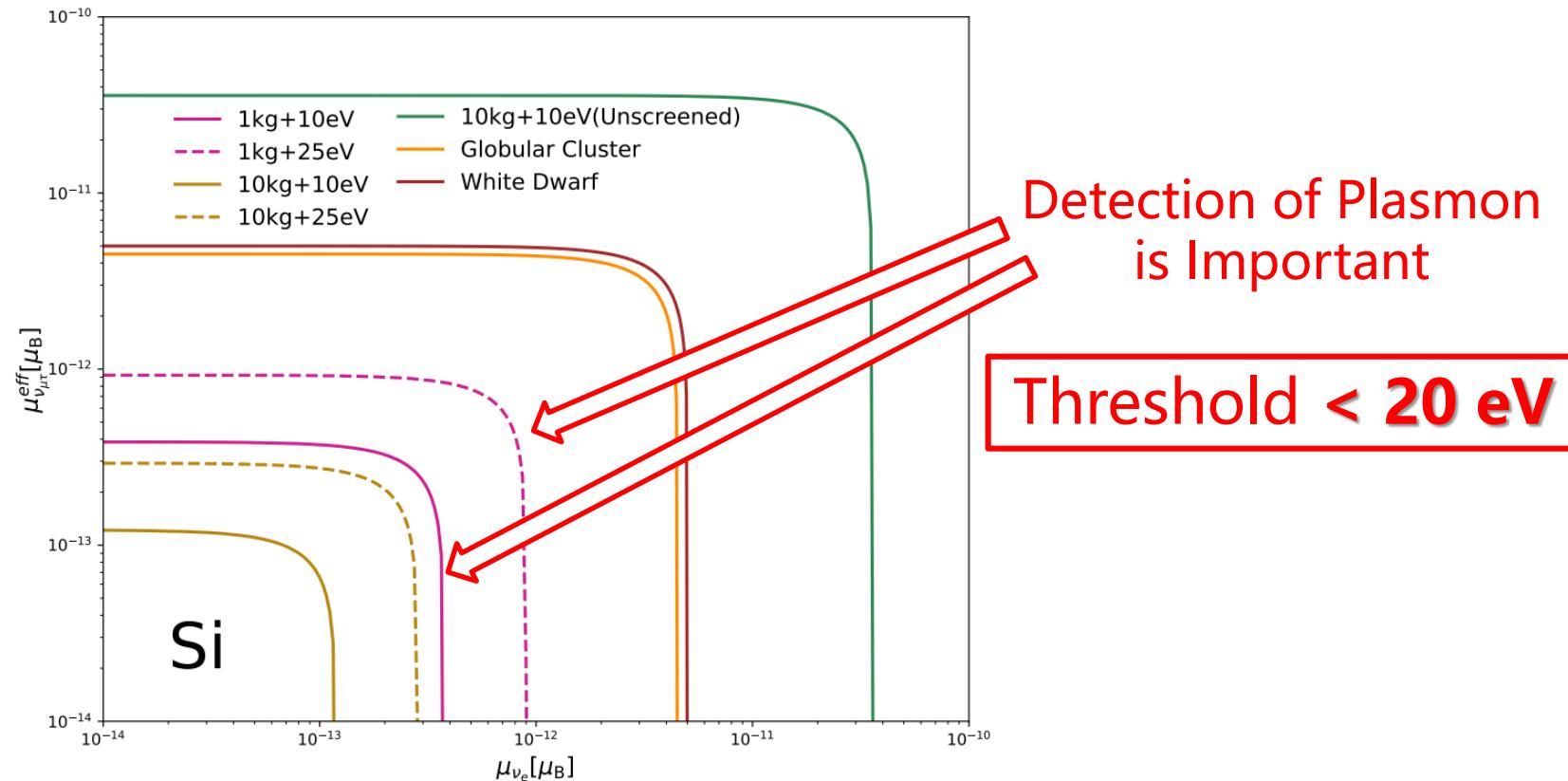
Plasmon Enhancement → Magnetic Moment





Screened neutrino excitation

- Constraints on neutrino magnetic moment
 - Better than astrophysical constraints at near future





Neutrino induced Migdal effect

■ The physical potential in Migdal effect

- Additional electron emission for detection
- Benefit from the coherent effect in CEvNS
- Possible explanation for quenching model

■ Experiments /Proposals

- Neutron nucleus scattering

Peter Cox, et al., Arxiv: 2208.12222

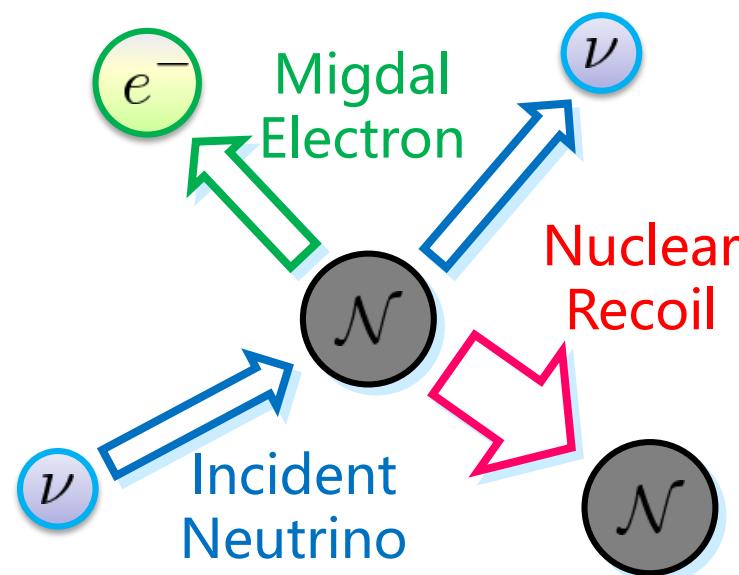
Nicole F. Bell, et al., Arxiv: 2112.08514

H.M. Araújo, et al., Arxiv: 2207.08284

- Dark matter direct detection

EDELWEISS collaboration, Arxiv: 2203.03993

XENON Collaboration, Arxiv: 1907.12771





Neutrino induced Migdal effect

Yu-Feng Li and Shuo-yu Xia, Nucl.Phys.B 1006 (2024) 116632, Arxiv: 2310.05704

■ Migdal effect as a multi-phonon process

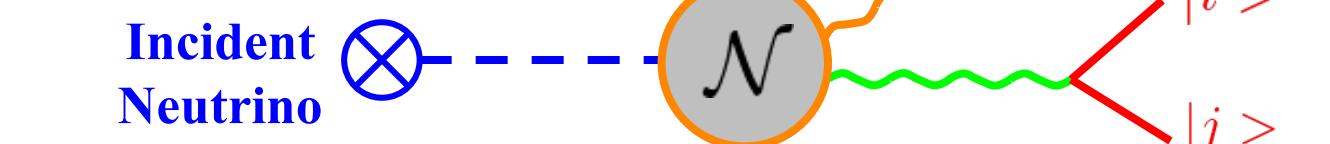
$$i\mathcal{M} = -iV_{CE\nu NS}(\mathbf{q}) Ne^{-W(\mathbf{q})} \prod_{s=1}^n \left(\frac{-i\mathbf{q} \cdot \epsilon_{\mathbf{k}_s, \alpha_s}}{\sqrt{2Nm_N\omega_{\mathbf{k}_s, \alpha_s}}} \right)$$

$$\times \sum_{\mathbf{G}'} \sum_{\mathbf{k}} \left\{ \left(\frac{\mathbf{q} \cdot (\mathbf{k} + \mathbf{G}')}{m_N} \right) \left[\frac{1}{(\varepsilon_i - \varepsilon_j)^2} \right] \right.$$

$$\times \left[\frac{Z_{ion} 4\pi \alpha_e}{V |\mathbf{k} + \mathbf{G}'|^2} \right] \langle i | e^{i(\mathbf{k} + \mathbf{G}') \cdot \hat{x}} | j \rangle \left. \right\}$$

$$\times \sum_{\mathbf{G}} \delta_{\sum_s \mathbf{k}_s + \mathbf{G} + \mathbf{k}, \mathbf{G}}$$

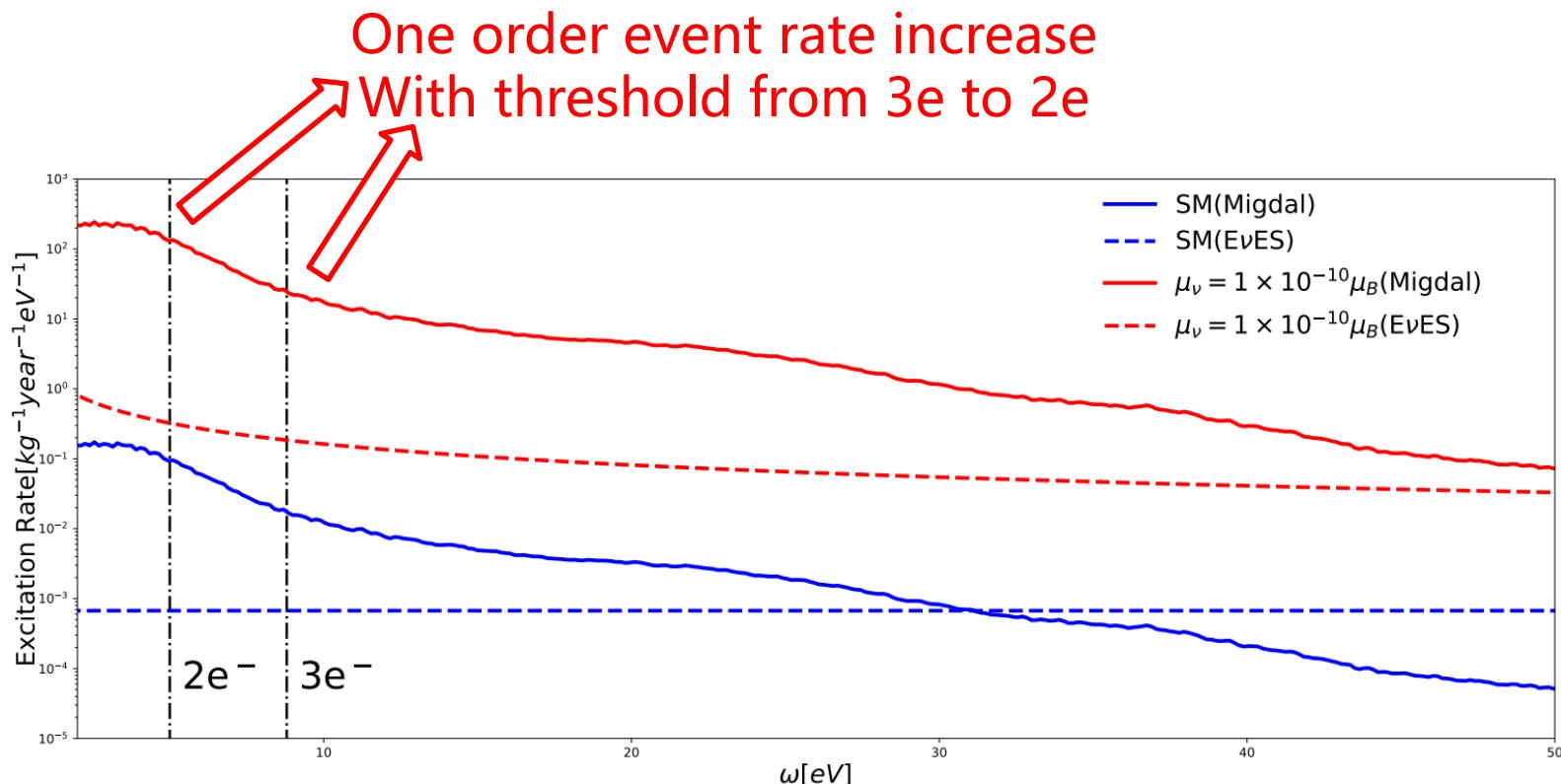
Incident
Neutrino





Neutrino induced Migdal effect

■ Event rate with the Migdal effect

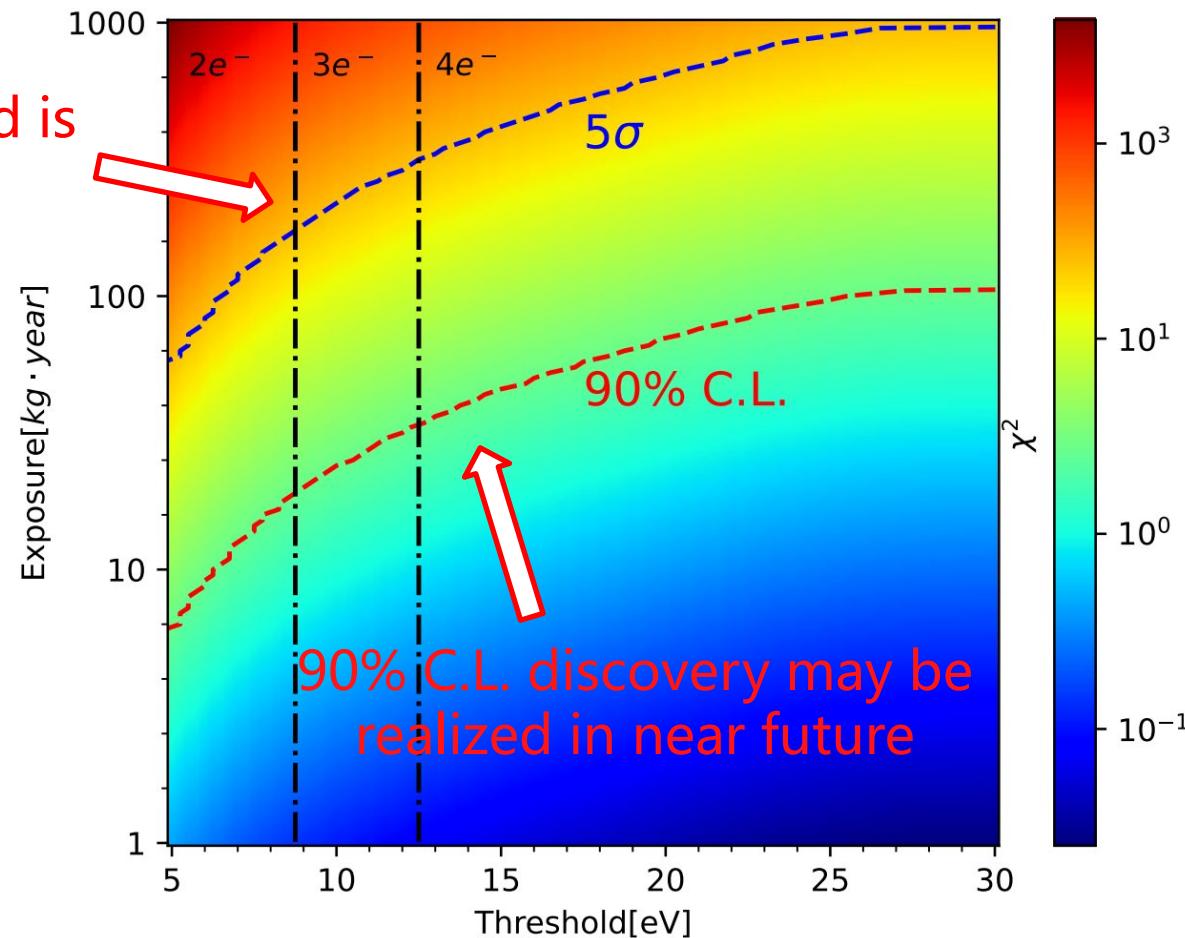




Neutrino induced Migdal effect

■ Sensitivity of the Migdal effect

Low threshold is important



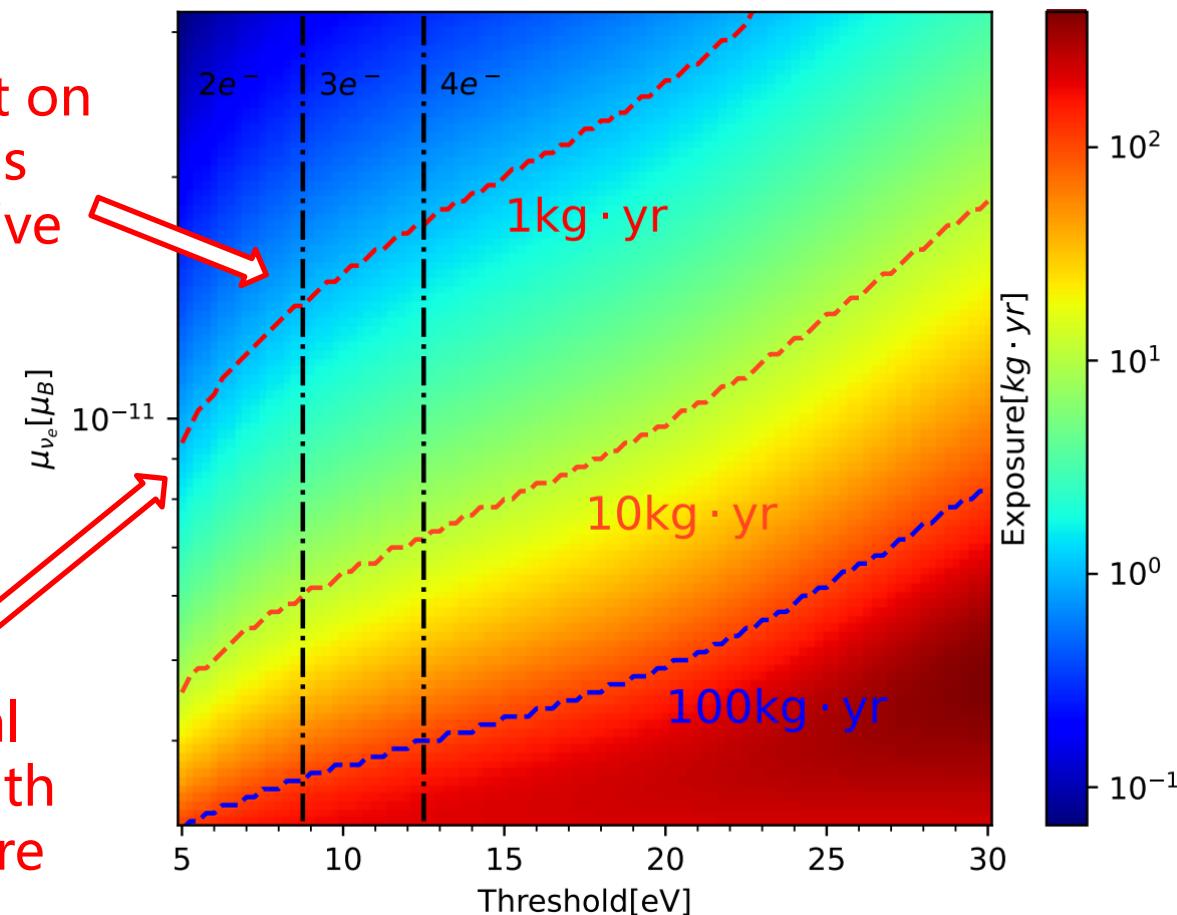


Neutrino induced Migdal effect

■ Constraints on neutrino magnetic moment

Improvement on threshold is more effective

Similar to astrophysical constraints with small exposure





Summary

- Plasmon can enhance the observation of neutrino magnetic moment
- Migdal effect may improve the detection of low energy nuclear recoil
- Low threshold is important to detect low energy neutrino excitation
- Challenges
 - Background, resolution, experimental scale,
 - Dielectric response, excitation models,

Thank You For Your Attention