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Tianjin University

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Exploration for Hyperon Halo with Density Functional Theory

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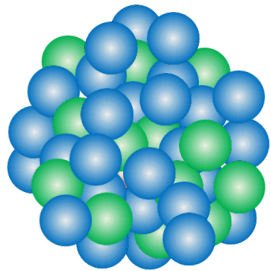


Outline

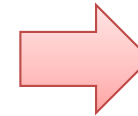
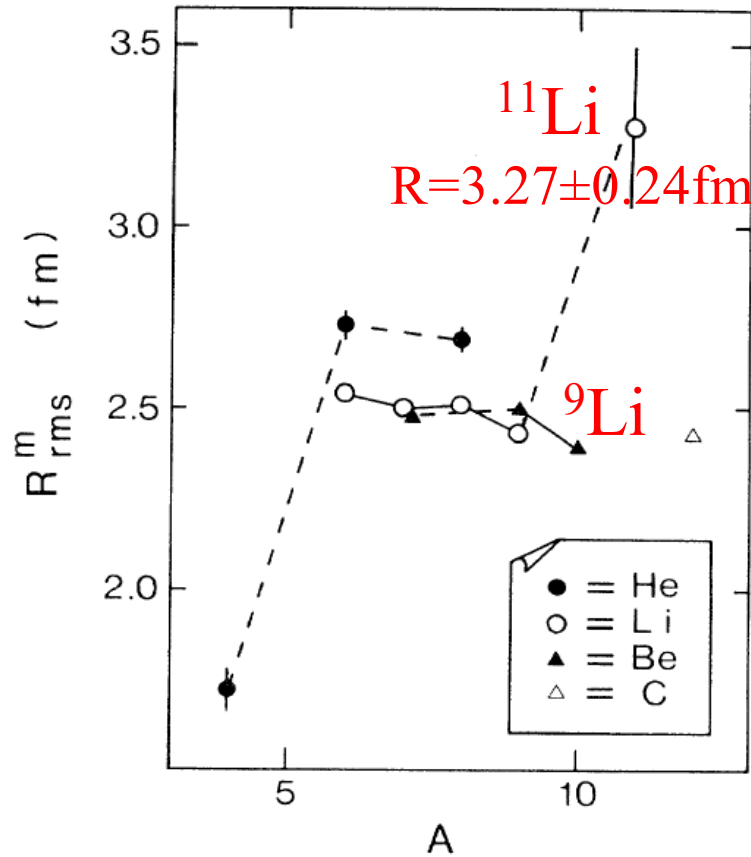
- Introduction
- Density Functional theory (DFT) for hypernucleus
- Hyperon halo orbits in C hypernuclei
- Hyperon halo orbits in Zr hypernuclei
- Possible hyperon halo in O hypernuclei
- Summary

Introduction

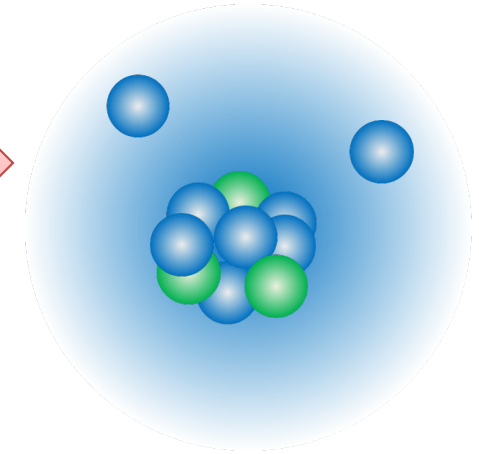
Stable nucleus



Neutron-rich nuclei

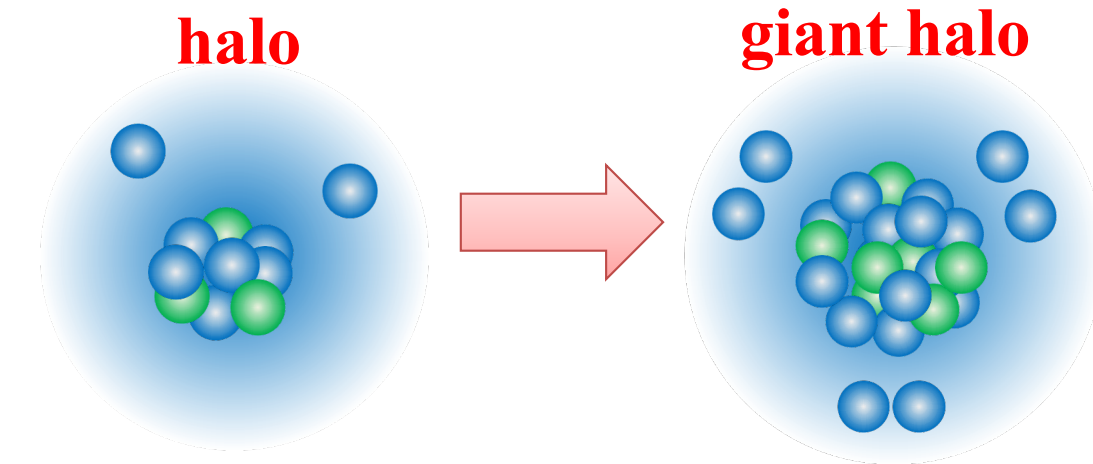


HALO

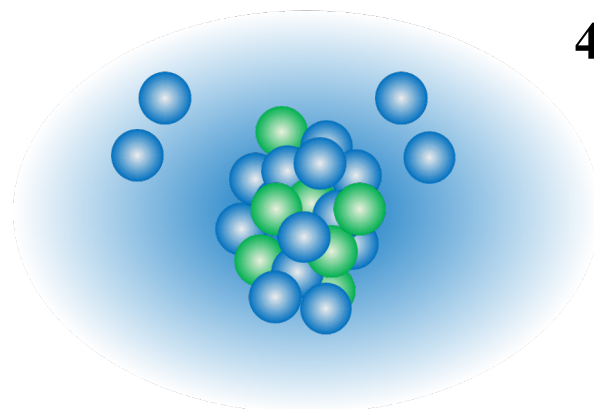


Tanihata, et al., PRL55, 2676 (1985)

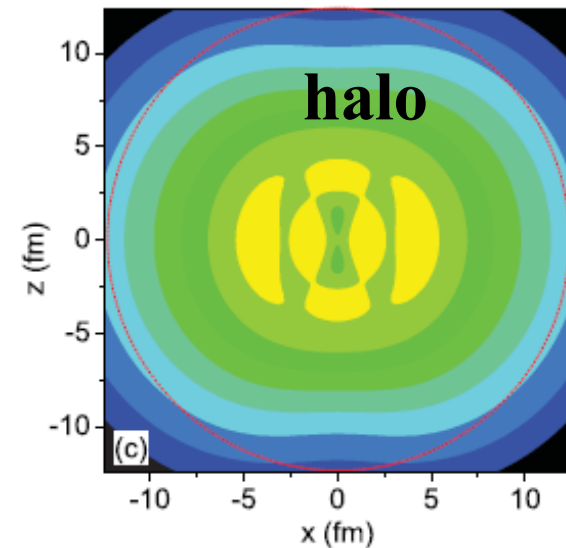
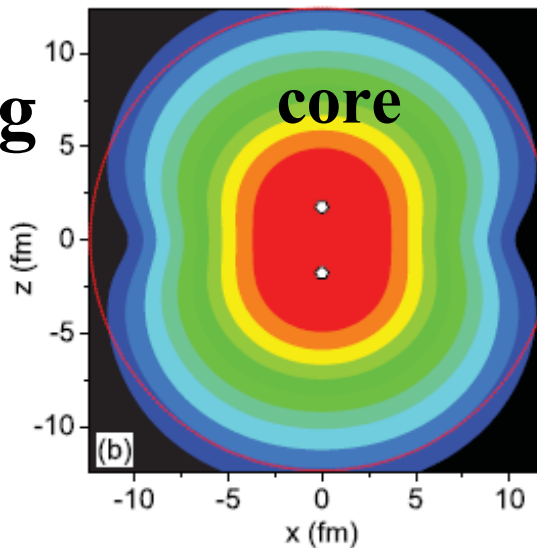
Introduction



J. Meng, and P. Ring. PRL. 80:460 (1998).
J. Meng, H. Toki, J. Y. Zeng, S.Q. Zhang,
S.G. Zhou. PRC, 65:041302(R) (2002).
J. Terasaki, S. Q. Zhang, S. G. Zhou, J. Meng,
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M. Grasso, S. Yoshida, N. Sandulescu, N.
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J. Meng, and S. G. Zhou, J. Phys. 42,
093101(2015)

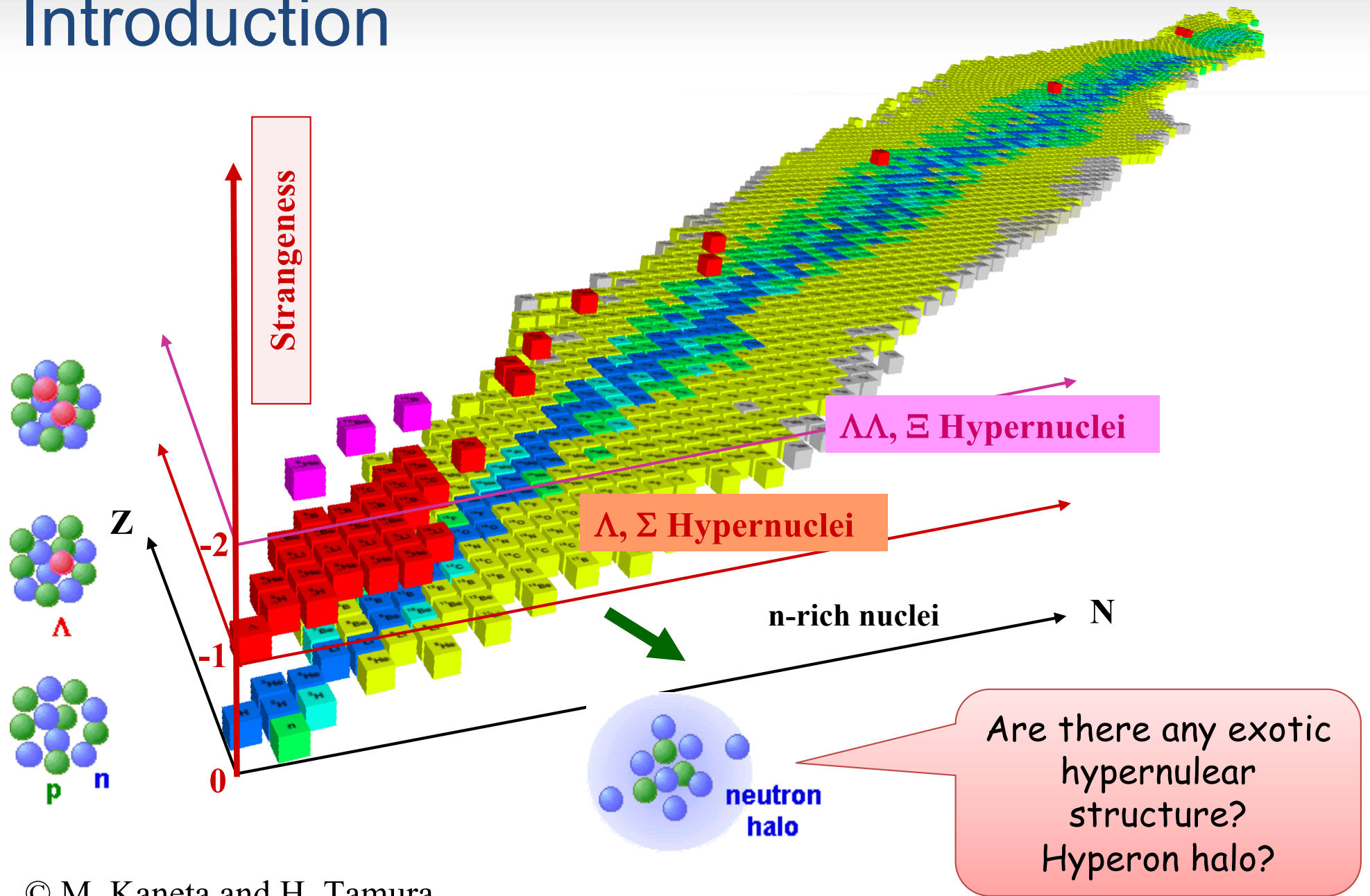


^{44}Mg



S.G. Zhou, S. G. Zhou, J. Meng, P. Ring, E. G. Zhao, PRC 82, 011301(R) (2010)
Y. Chen. L.L. Lu, H.Z. Liang, J. Meng, PRC.85.067301(2012)

Introduction



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<http://lambda.phys.tohoku.ac.jp/strangeness/index-j.html>

Introduction

➤ Hyperon effect on the neutron drip line

- Few body model

E. Hiyama, M. Kamimura, T. Motoba, T. Yamada, Y. Yamamoto, Phys. Rev. C 53, 2075 (1996).

- Shell model

R. Wirth, R. Roth, Phys. Lett. B 779 (2018) 336–341

- Skyrme Hartree-Fock model

X. R. Zhou, A. Polls, H. J. Schulze, and I. Vidana, Phys. Rev. C 78, 054306 (2008).

- Relativistic Hartree-Bogoliubov

D. Vretenar, W. Pöschl, G. A. Lalazissis, and P. Ring, Phys. Rev. C 57, R1060(R) (1998).

H. F. Lv, J. Meng, S. Q. Zhang, and S. G. Zhou, Eur. Phys. J. A 17, 19 (2003).

➤ Hyperon drip line (multistrangeness)

- Skyrme Hartree-Fock model

J. Margueron, E. Khan, and F. Gulminelli, Phys. Rev. C 96, 054317 (2017).

- Hartree-Fock-Bogoliubov

H. Güven, K. Bozkurt, E. Khan, and J. Margueron, Phys. Rev. C 98, 014318 (2018).

- Relativistic Hartree-Bogoliubov

H. F. Lv, J. Meng, Chinese Phys. Lett. 19, No. 12, 1775 (2002)

H. F. Lv, Chinese Phys. Lett. 25, 3613 (2008).

Y. T. Rong, P. Zhao, and S. G. Zhou, Phys. Lett. B 807, 135533 (2020).

DFT for hypernuclei

★ Skyrme-like interaction

✓ NN interaction [D. Vautherin and D.M. Brink, PRC 5, 626 \(1972\).](#)

$$v_{NN}(\mathbf{r}_1 - \mathbf{r}_2) = t_0 (1 + x_0 P_\sigma) \delta(\mathbf{r}_1 - \mathbf{r}_2) + \frac{1}{2} t_1 (1 + x_1 P_\sigma) [\mathbf{k}'^2 \delta(\mathbf{r}_1 - \mathbf{r}_2) + \delta(\mathbf{r}_1 - \mathbf{r}_2) \mathbf{k}^2] \\ + t_2 (1 + x_2 P_\sigma) \mathbf{k}' \cdot \delta(\mathbf{r}_1 - \mathbf{r}_2) \mathbf{k} + iW_0 (\boldsymbol{\sigma}_1 + \boldsymbol{\sigma}_2) \cdot \mathbf{k}' \delta(\mathbf{r}_1 - \mathbf{r}_2) \times \mathbf{k}, \quad (1)$$

$$v_{den-NN}(\mathbf{r}_1, \mathbf{r}_2, \mathbf{r}_3) = \frac{1}{6} t_3 (1 + x_3 P_\sigma) \delta(\mathbf{r}_1 - \mathbf{r}_2) \rho^\alpha \left(\frac{\mathbf{r}_1 + \mathbf{r}_2}{2} \right), \quad (2)$$

✓ $N\Lambda$ interaction [D. E. Lanskoy and Y. Yamamoto, PRC 55, 2330 \(1997\).](#)

$$v_{\Lambda N}(\mathbf{r}_\Lambda - \mathbf{r}_N) = t_0^\Lambda (1 + x_0^\Lambda P_\sigma) \delta(\mathbf{r}_\Lambda - \mathbf{r}_N) + \frac{1}{2} t_1^\Lambda [\mathbf{k}'^2 \delta(\mathbf{r}_\Lambda - \mathbf{r}_N) + \delta(\mathbf{r}_\Lambda - \mathbf{r}_N) \mathbf{k}^2] \\ + t_2^\Lambda \mathbf{k}' \delta(\mathbf{r}_\Lambda - \mathbf{r}_N) \cdot \mathbf{k} + iW_0^\Lambda \mathbf{k}' \delta(\mathbf{r}_\Lambda - \mathbf{r}_N) \cdot (\boldsymbol{\sigma}_N + \boldsymbol{\sigma}_\Lambda) \times \mathbf{k} \quad (3)$$

$$v_{den-\Lambda N}(\mathbf{r}_\Lambda, \mathbf{r}_N, \rho) = \frac{3}{8} t_3^\Lambda (1 + x_3^\Lambda P_\sigma) \delta(\mathbf{r}_\Lambda - \mathbf{r}_N) \rho^\gamma \left(\frac{\mathbf{r}_\Lambda + \mathbf{r}_N}{2} \right), \quad (4)$$

✓ $\Lambda\Lambda$ interaction [D. E. Lanskoy, PRC 58, 3351 \(1998\).](#)

$$V_{\Lambda\Lambda} = \lambda_0 \delta(\mathbf{r}_1 - \mathbf{r}_2) + \frac{1}{2} \lambda_1 [\mathbf{k}'^2 \delta(\mathbf{r}_1 - \mathbf{r}_2) + \delta(\mathbf{r}_1 - \mathbf{r}_2) \mathbf{k}^2] + \lambda_2 \mathbf{k}' \delta(\mathbf{r}_1 - \mathbf{r}_2) \mathbf{k} + \lambda_3 \delta(\mathbf{r}_1 - \mathbf{r}_2) \rho_N^\alpha \left(\frac{\mathbf{r}_1 + \mathbf{r}_2}{2} \right)$$

Hyperon halo orbits in C hypernuclei

★Energy density functional

✓ NN interaction SIII, SLy4, SkM*

✓ NΛ interaction LY5 $W_0^\Lambda = 62 \text{ MeV fm}^5$ D. E. Lansky and Y. Yamamoto, PRC 55, 2330 (1997)

to reproduce the spin-orbit splitting of 1p states (0.152 MeV) in ${}^{13}_\Lambda\text{C}$

H. Kohri et al., Phys. Rev.C 65, 034607 (2002)

LY5r $W_0^\Lambda = 4.7 \text{ MeV fm}^5$

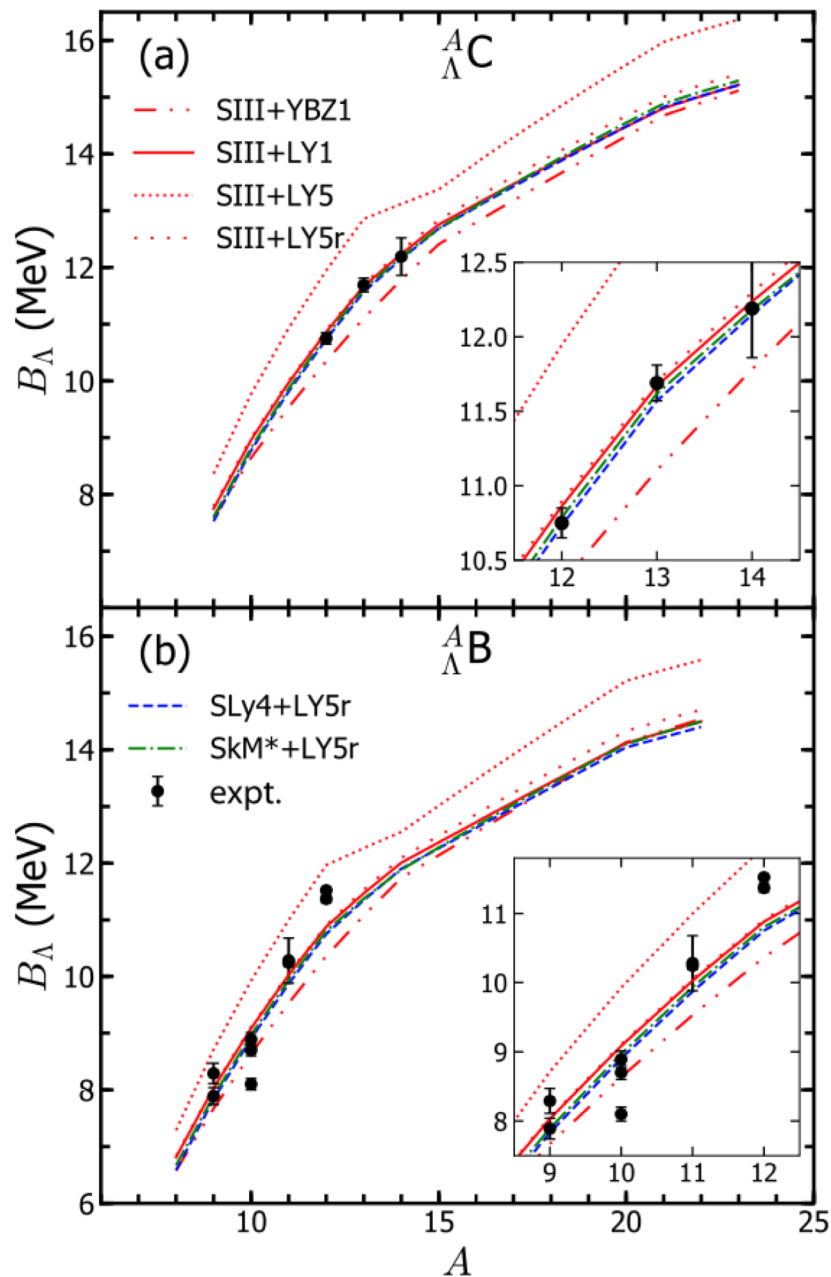
Table 3.6: The calculated binding energies of ${}^{13}_\Lambda\text{C}$ obtained by SkM* and Set V in Ref. [15] with $W_0^\Lambda = 4.7 \text{ MeV fm}^5$. The experiment data are taken from Ref. [24] [H. Kohri, et al., Phys. Rev. C 65, 9 (2002)]. All the values are in MeV.

Λ orbit	e_Λ	B_Λ	B_Λ^{exp}	$B_\Lambda(s_{1/2}) - B_\Lambda$	$E_{\text{ex}}^{\text{exp}}$
$s_{1/2}$	-13.156	11.618	11.69 ± 0.12	--	--
$p_{1/2}$	-1.782	0.273	} 0.155	11.344	$10.982 \pm 0.031(\text{stat}) \pm 0.056(\text{syst})$
$p_{3/2}$	-1.936	0.428		11.190	$10.830 \pm 0.031(\text{stat}) \pm 0.056(\text{syst})$

Cal. **Exp.**

$$B_\Lambda = B({}^{A+1}_\Lambda Z) - B({}^A Z)$$

Hyperon halo orbits in C hypernuclei

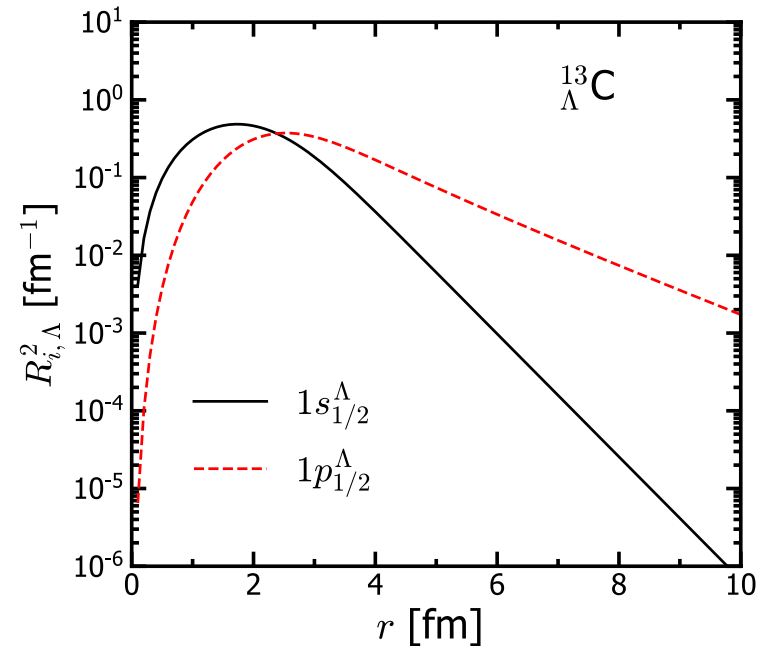
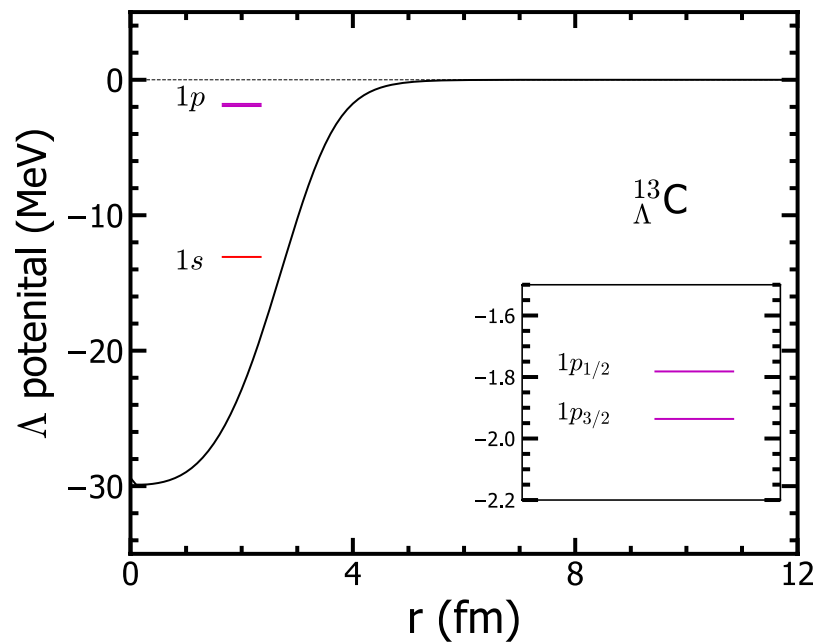


Good agreement with
the present data!

Y. Z., H. Sagawa, and E. Hiyama,
Phys. Rev. C **103**, 034321 (2021).

Hyperon halo orbits in C hypernuclei

Y. Z., H. Sagawa, and E. Hiyama, Phys. Rev. C **103**, 034321 (2021)



Λ orbit	s.p. energy	rms radius
$1s_{1/2}$	-13.156	2.144
$1p_{1/2}$	-1.782	3.410

60% increase
Halo orbits

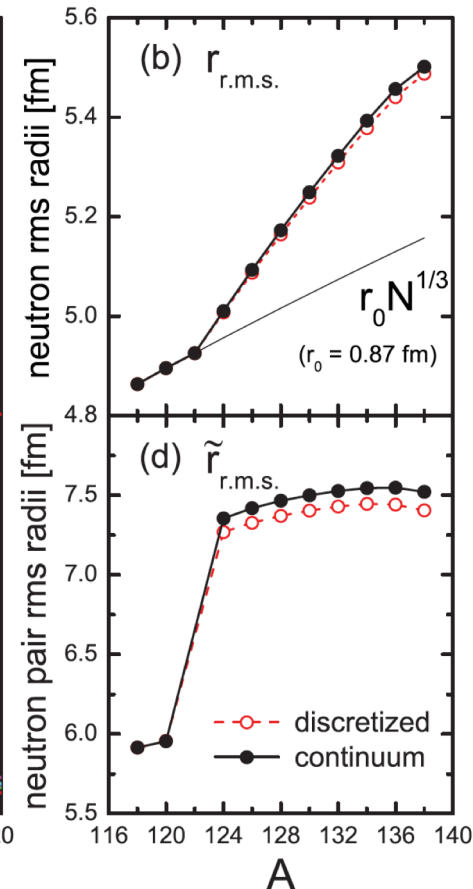
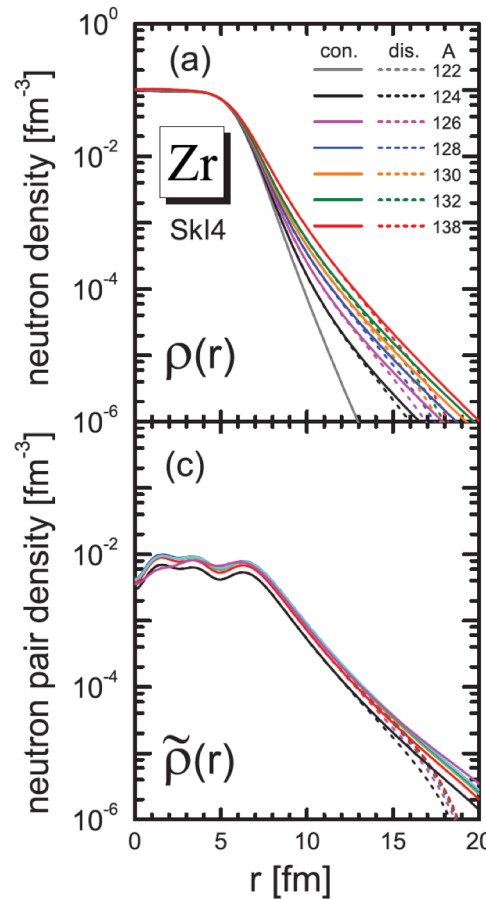
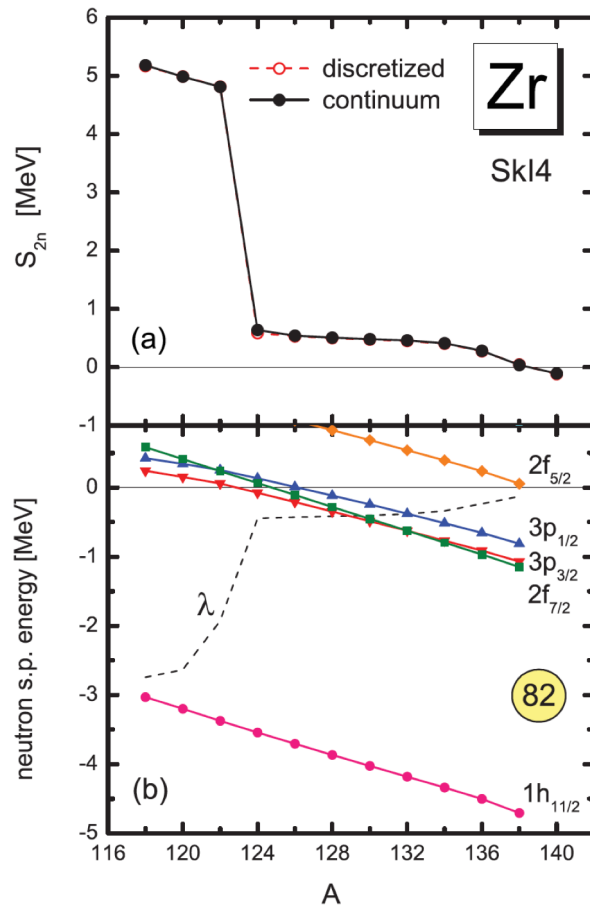
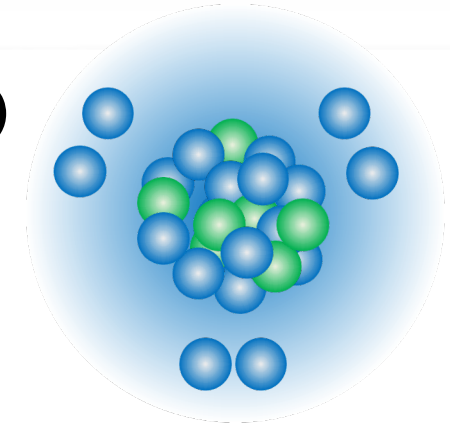
Hyperon halo orbits in Zr hypernuclei

□ Zr isotopes

Giant halo Zr ($A > 124$)

J. Meng and P. Ring, PRL 80, 460 (1998).

Y. Z., M. Matsuo, and J. Meng, PRC 86, 54318 (2012).



How about 1 or 2 Λ particles are added?

Hyperon halo orbits in Zr hypernuclei

■ Hartree-Fock-Bogoliubov (HFB) model

✓ NN interaction: SkI4 P.-G. Reinhard and H. Flocard, NPA 584, 467 (1995).

✓ NN pairing: DDDI

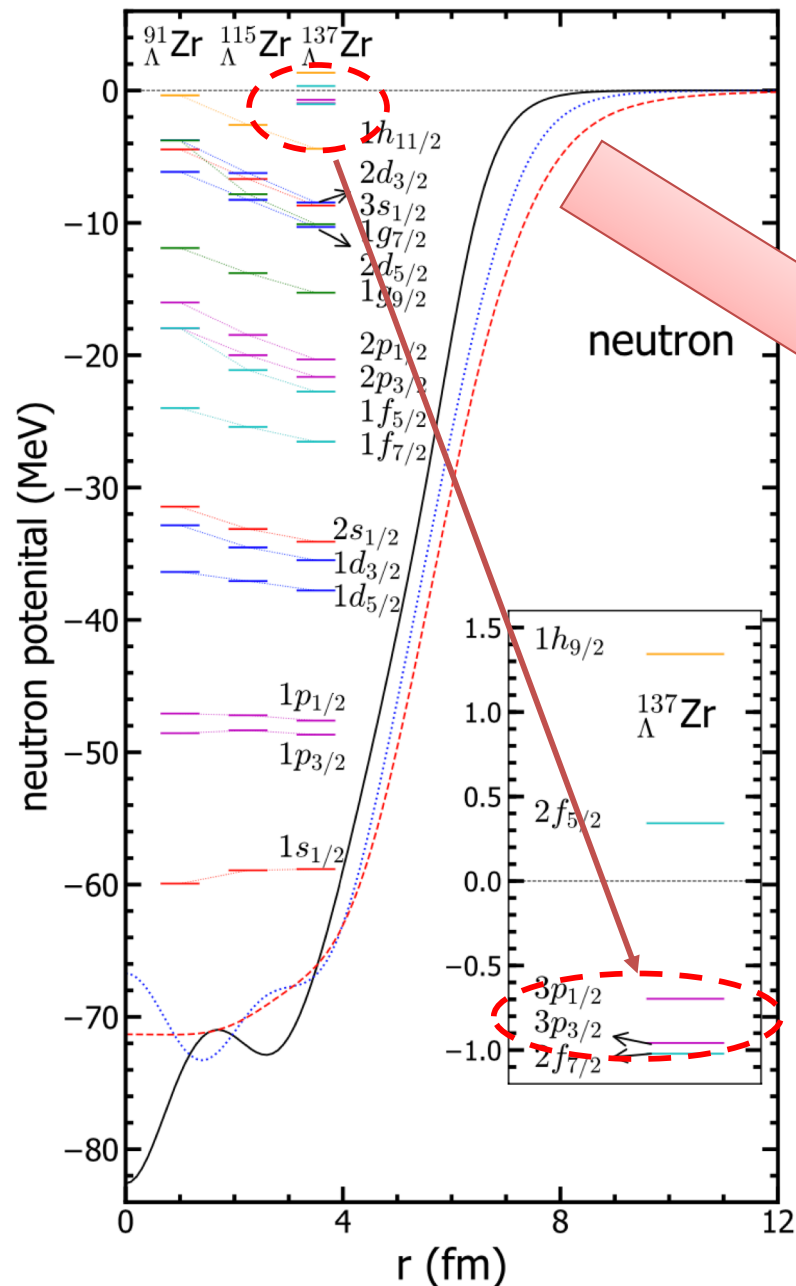
$$\Delta(\mathbf{r}) = \frac{1}{2} V_0 \left[1 - \eta \left(\frac{\rho_q(\mathbf{r})}{\rho_0} \right)^\alpha \right] \tilde{\rho}(\mathbf{r}), \quad q = n \text{ or } p.$$

$$V_0 = -300 \text{ MeV fm}^3, \quad \eta = 0.5, \quad \alpha = 0.5, \quad \rho_0 = 0.16 \text{ fm}^{-3}, \quad E_{\text{cut}} = 70 \text{ MeV}$$

M. Grasso et al. PRC 74(2006)64317.

✓ N Λ interaction: LY5r Y. Z., H. Sagawa, and E. Hiyama, Phys. Rev. C 103, 034321 (2021)

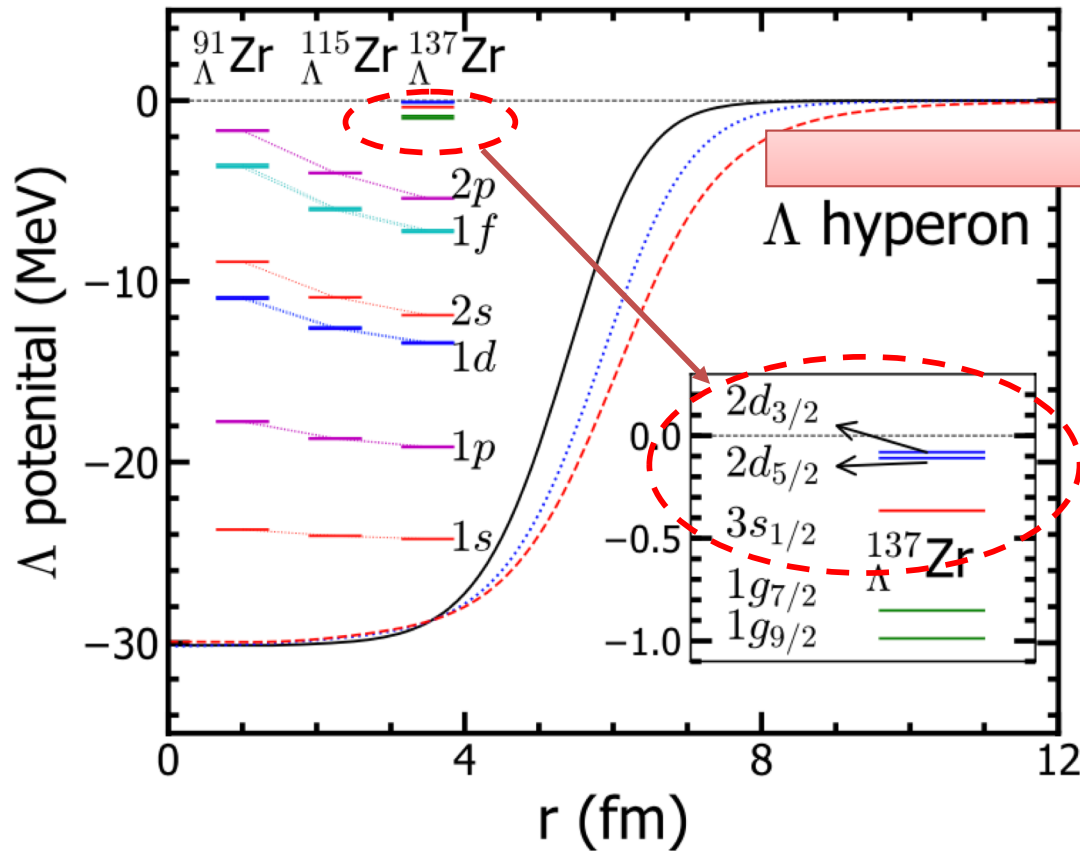
Hyperon halo orbits in Zr hypernuclei



In ${}^{136}\text{Zr}$, the diffusive neutron potential can hold the weakly bound p orbits, which can hold up to 6 neutrons \rightarrow giant neutron halo

Y. Z., H. Sagawa, and E. Hiyama, Prog. Theor. Exp. Phys. **2022** 023D01 (2022)

Hyperon halo orbits in Zr hypernuclei



In ^{136}Zr , the diffusive neutron density leads to a diffusive hyperon potential also.

	state	s.p.e.	v^2	r_{rms}
n	$3p_{1/2}$	-0.726	0.872	8.7611
	$3p_{3/2}$	-0.986	0.926	8.3981
	$2f_{5/2}$	0.292	0.155	7.2496
	$2f_{7/2}$	-1.047	0.886	6.9517
	$1h_{9/2}$	1.271	0.050	6.1079
	$1h_{11/2}$	-4.432	0.992	5.7730
total				5.4624
Λ	$3s_{1/2}$	-0.370	0.500	9.9822
	$2d_{3/2}$	-0.090	0.0	8.5794
	$2d_{5/2}$	-0.118	0.0	8.4148
	$1g_{7/2}$	-0.870	0.0	5.8904
	$1g_{9/2}$	-1.006	0.0	5.8721

5 almost degenerate weakly bound low-l orbits can hold up to 12 Λ hyperons

Possible hyperon halo in O hypernuclei

■ Hartree-Fock-Bogoliubov (HFB) model

✓ NN interaction: SIII

✓ NN pairing: DDDI $\Delta(\mathbf{r}) = \frac{1}{2} V_0 \left[1 - \eta \left(\frac{\rho_q(\mathbf{r})}{\rho_0} \right)^\alpha \right] \tilde{\rho}(\mathbf{r}), \quad q = n \text{ or } p.$

$V_0 = -458.4 \text{ MeV fm}^3, \eta = 0.83, \alpha = 0.51, \rho_0 = 0.08 \text{ fm}^{-3}, E_{\text{cut}} = 60 \text{ MeV}$ neutron drip line: ^{24}O

✓ $\text{N}\Lambda$ interaction: LY5r $W_0^\Lambda = 4.7 \text{ MeV fm}^5$ YZ et al., PRC 103, 034321 (2021)

✓ $\Lambda\Lambda$ pairing: DDDI volume-type

$V_0 = -139 \text{ MeV fm}^3, \eta = 0, \alpha = 0.51, E_{\text{cut}} = 60 \text{ MeV}$ H. Güven, et al., PRC 98, 014318 (2018)

✓ $\Lambda\Lambda$ interaction:

$$V_{\Lambda\Lambda} = \lambda_0 \delta(\mathbf{r}_1 - \mathbf{r}_2) + \frac{1}{2} \lambda_1 \left[\mathbf{k}'^2 \delta(\mathbf{r}_1 - \mathbf{r}_2) + \delta(\mathbf{r}_1 - \mathbf{r}_2) \mathbf{k}^2 \right] + \lambda_2 \mathbf{k}' \delta(\mathbf{r}_1 - \mathbf{r}_2) \mathbf{k} + \lambda_3 \delta(\mathbf{r}_1 - \mathbf{r}_2) \rho_N^\alpha \left(\frac{\mathbf{r}_1 + \mathbf{r}_2}{2} \right)$$

$$\lambda_0 = -76.2 \text{ MeV fm}^3$$

$$B_{\Lambda\Lambda}^{\text{exp}} = 19.07 \pm 0.08 \pm 0.07 \text{ MeV of } ^{11}_{\Lambda\Lambda}\text{Be}$$

$$\lambda_1 = 57.5 \text{ MeV fm}^5$$

Ekawa et al., PTEP. 2019, 021D02 (2019)

$$\lambda_2 = ?$$

Try: $\lambda_2 = 0$ or -100 MeV fm^5

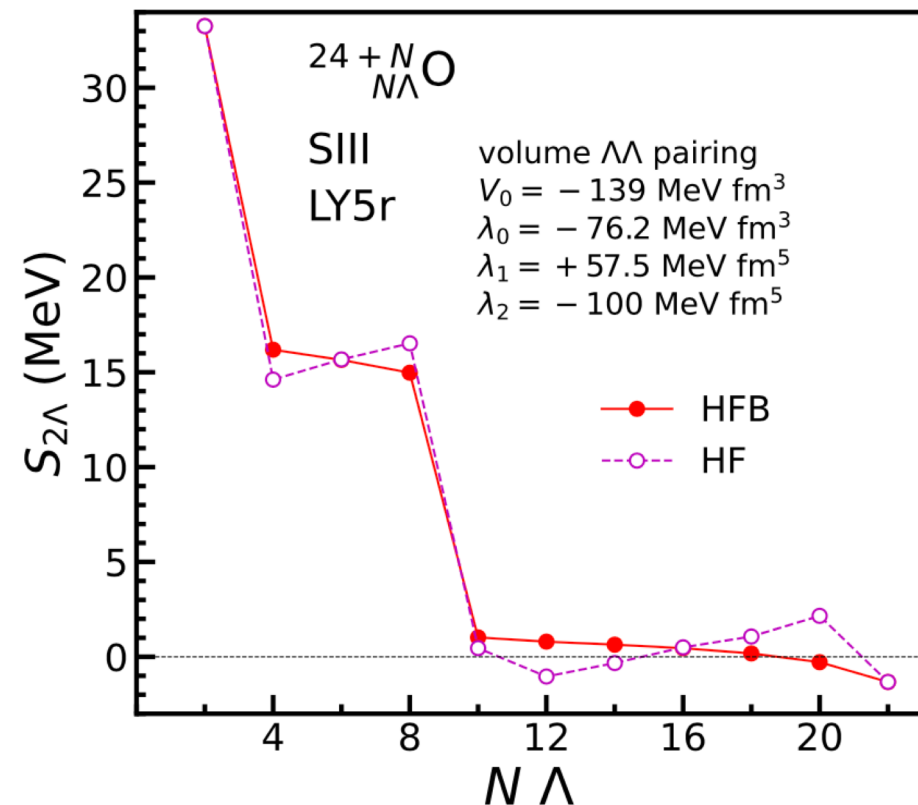
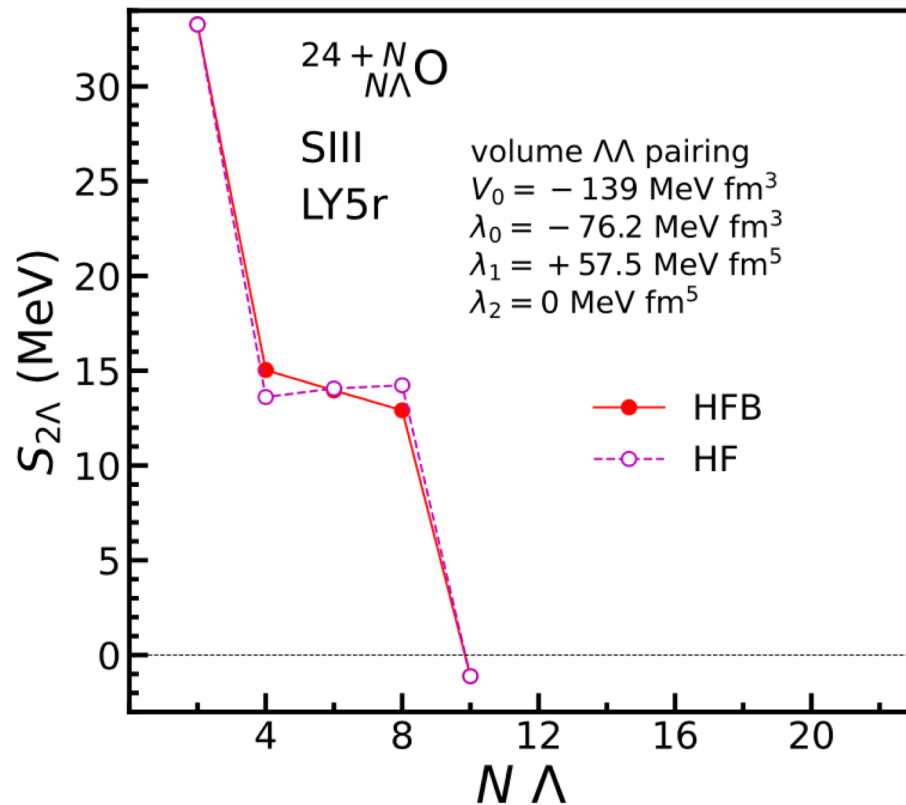
Possible hyperon halo in O hypernuclei

✓ $\Lambda\Lambda$ interaction

$$V_{\Lambda\Lambda} = \lambda_0 \delta(\mathbf{r}_1 - \mathbf{r}_2) + \frac{1}{2} \lambda_1 [\mathbf{k}'^2 \delta(\mathbf{r}_1 - \mathbf{r}_2) + \delta(\mathbf{r}_1 - \mathbf{r}_2) \mathbf{k}^2] + \lambda_2 \mathbf{k}' \delta(\mathbf{r}_1 - \mathbf{r}_2) \mathbf{k} + \lambda_3 \delta(\mathbf{r}_1 - \mathbf{r}_2) \rho_N^\alpha \left(\frac{\mathbf{r}_1 + \mathbf{r}_2}{2} \right)$$

$\lambda_2 = 0 \text{ MeV fm}^5$

$\lambda_2 = -100 \text{ MeV fm}^5$



! With more attractive p-wave interaction, the hyperon drip line could be extended.

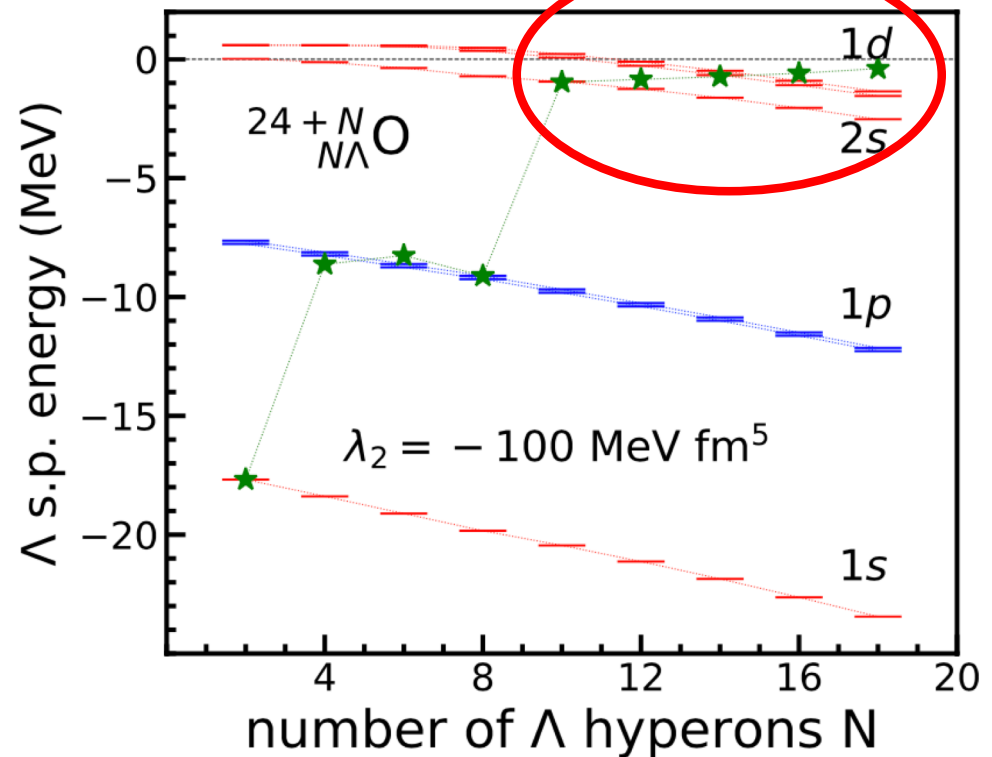
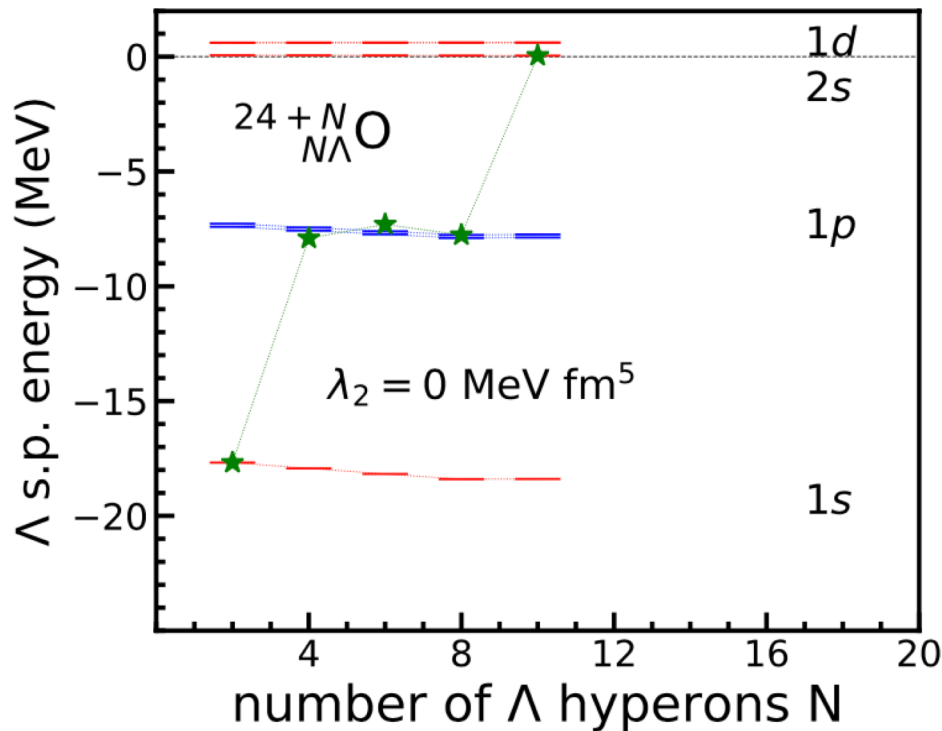
Possible hyperon halo in O hypernuclei

✓ $\Lambda\Lambda$ interaction

$$V_{\Lambda\Lambda} = \lambda_0\delta(\mathbf{r}_1-\mathbf{r}_2) + \frac{1}{2}\lambda_1 [\mathbf{k}'^2\delta(\mathbf{r}_1-\mathbf{r}_2) + \delta(\mathbf{r}_1-\mathbf{r}_2)\mathbf{k}^2] + \lambda_2\mathbf{k}'\delta(\mathbf{r}_1-\mathbf{r}_2)\mathbf{k} + \lambda_3\delta(\mathbf{r}_1-\mathbf{r}_2)\rho_N^\alpha\left(\frac{\mathbf{r}_1+\mathbf{r}_2}{2}\right)$$

$$\lambda_2 = 0 \text{ MeV fm}^5$$

$$\lambda_2 = -100 \text{ MeV fm}^5$$



! 2s and 1d orbits are pulled down from the continuum to the weakly bound states as hyperon number increases.

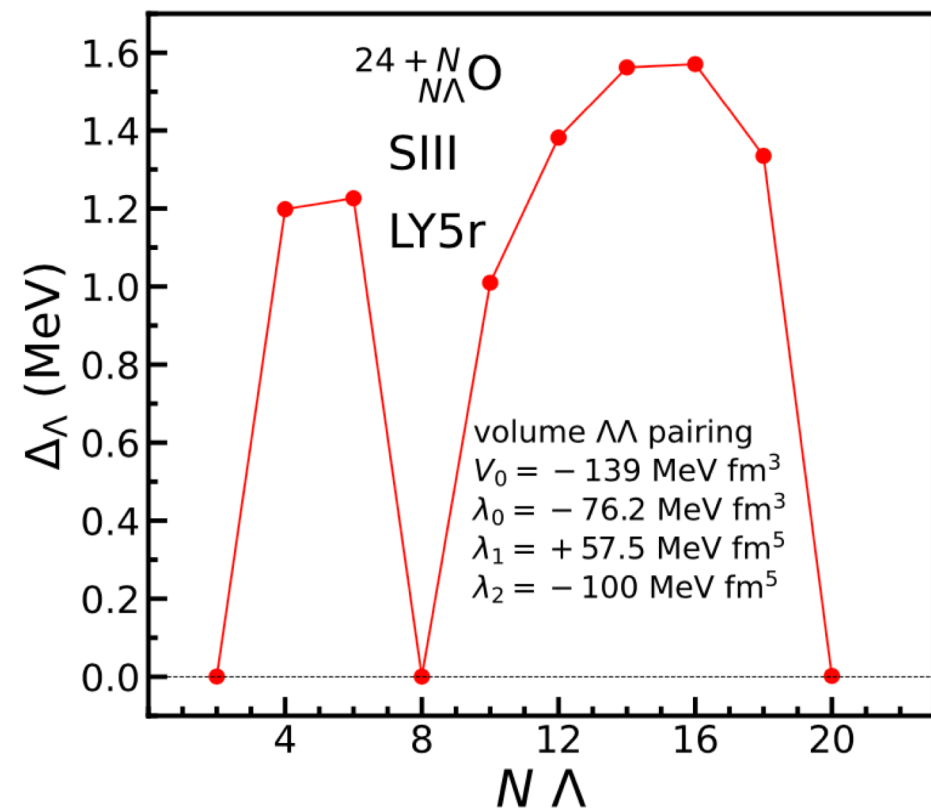
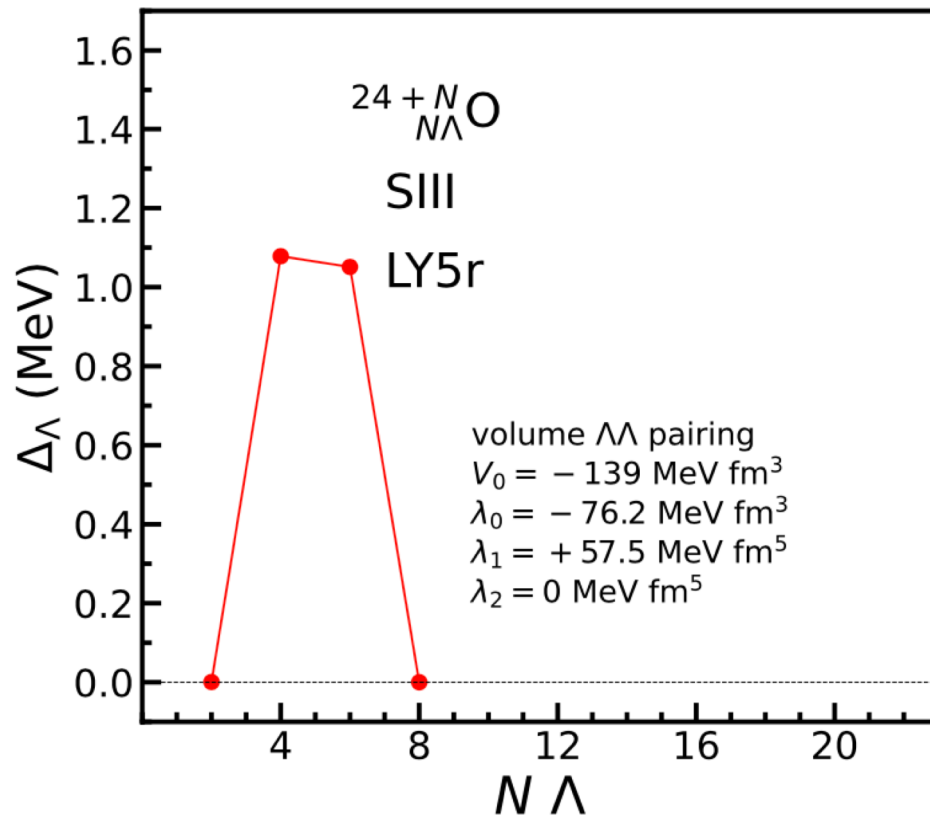
Possible hyperon halo in O hypernuclei

✓ $\Lambda\Lambda$ interaction

$$V_{\Lambda\Lambda} = \lambda_0 \delta(\mathbf{r}_1 - \mathbf{r}_2) + \frac{1}{2} \lambda_1 [\mathbf{k}'^2 \delta(\mathbf{r}_1 - \mathbf{r}_2) + \delta(\mathbf{r}_1 - \mathbf{r}_2) \mathbf{k}^2] + \lambda_2 \mathbf{k}' \delta(\mathbf{r}_1 - \mathbf{r}_2) \mathbf{k} + \lambda_3 \delta(\mathbf{r}_1 - \mathbf{r}_2) \rho_N^\alpha \left(\frac{\mathbf{r}_1 + \mathbf{r}_2}{2} \right)$$

$$\lambda_2 = 0 \text{ MeV fm}^5$$

$$\lambda_2 = -100 \text{ MeV fm}^5$$

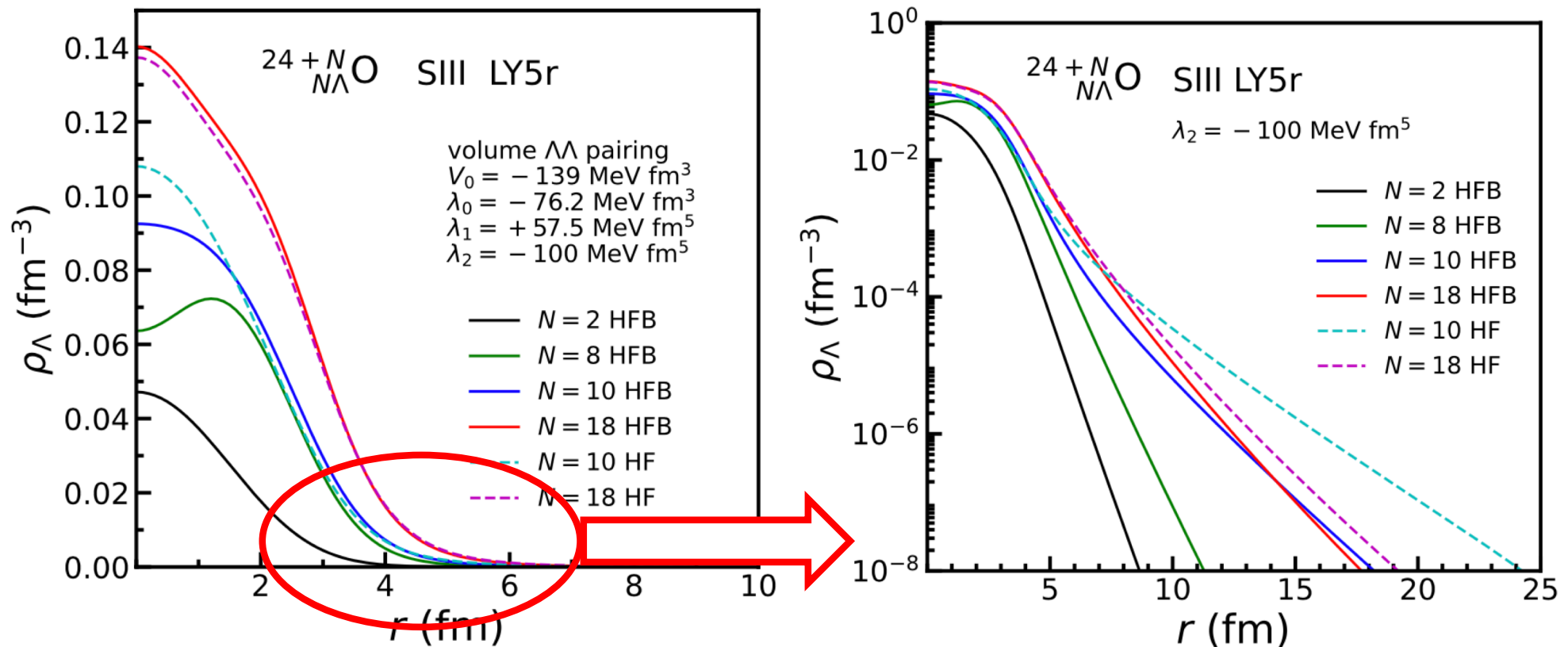


! Obvious hyperon shell closure at $N=2, 8, 20$

Possible hyperon halo in O hypernuclei

$$\lambda_2 = -100 \text{ MeV fm}^5$$

➤ Hyperon density distribution

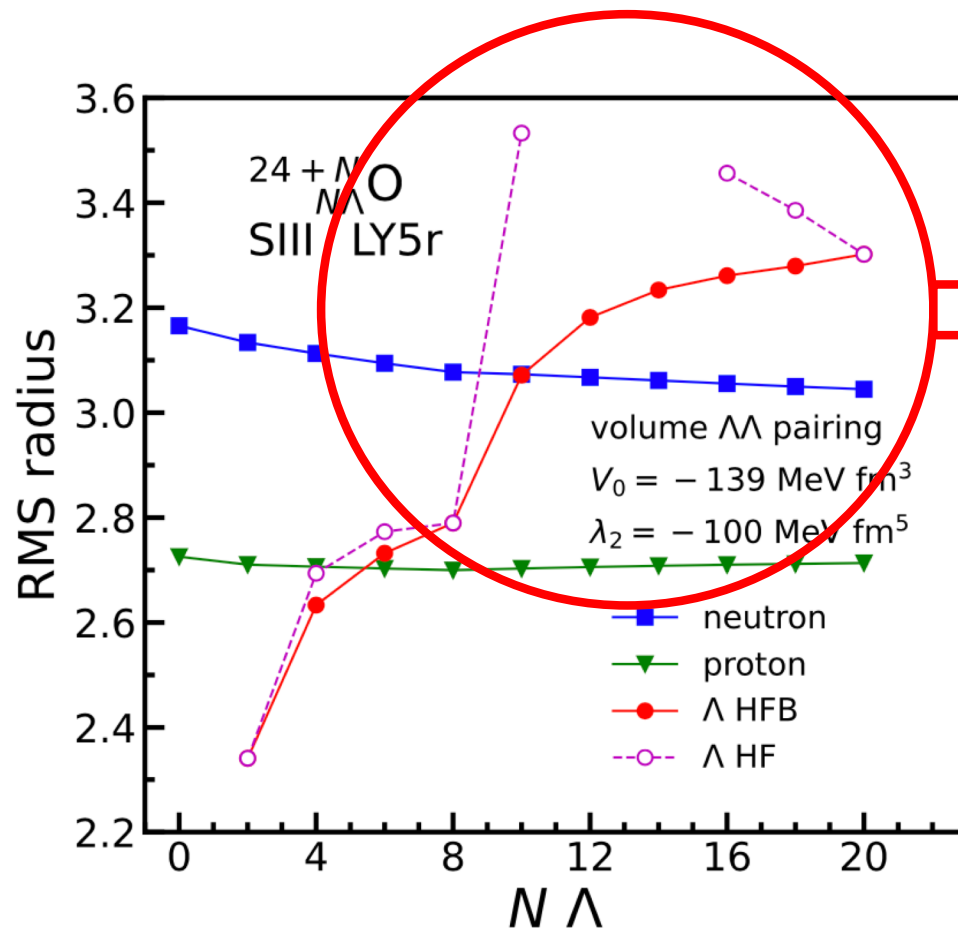


- ! Obvious increase of the hyperon density from $N=8$ to $N=10$
- ! "pairing anti-halo" effect

Possible hyperon halo in O hypernuclei

$$\lambda_2 = -100 \text{ MeV fm}^5$$

- RMS radius of neutron, proton and hyperon



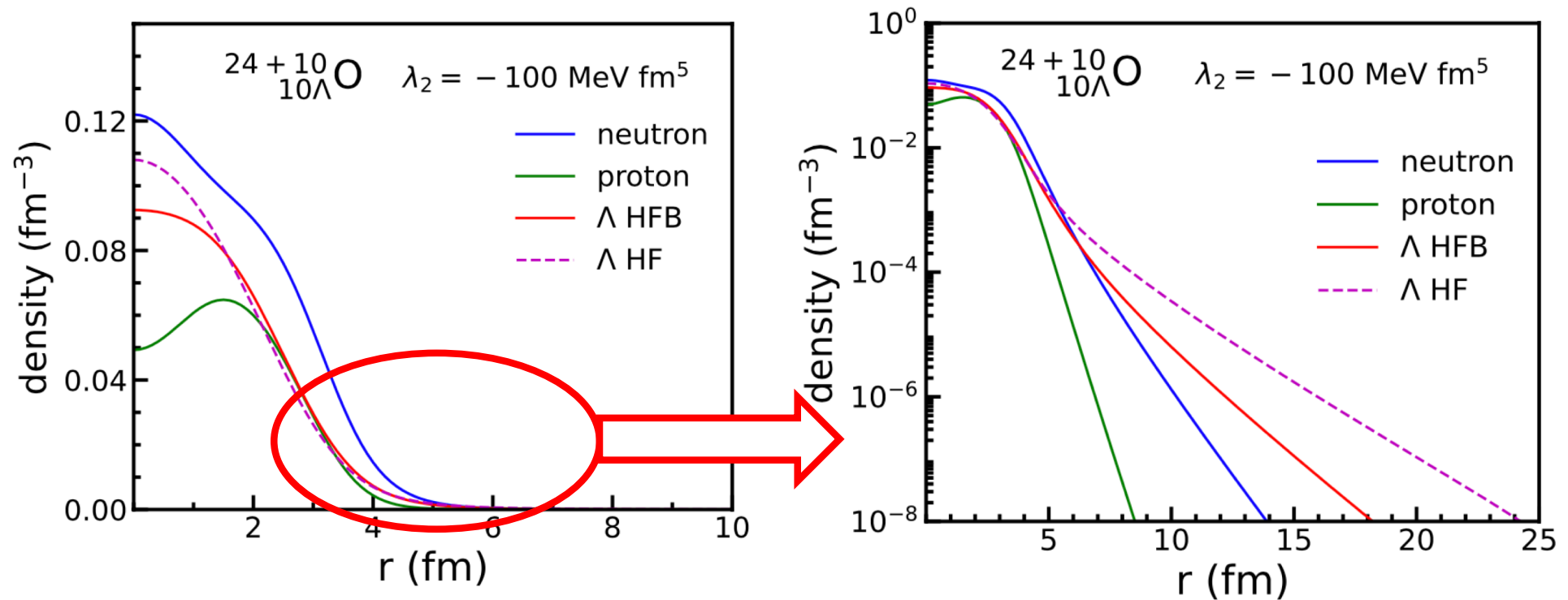
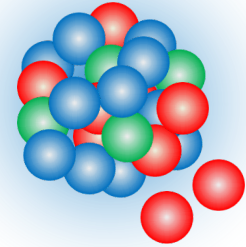
! quick increase of hyperon RMS radius, even larger than neutron after $N=10$

! "pairing anti-halo" effect

Possible hyperon halo in $^{24}_{10}\Lambda\text{O}$ hypernuclei

$$\lambda_2 = -100 \text{ MeV fm}^5$$

➤ Density of neutron, proton and hyperon in $^{24}_{10}\Lambda\text{O}$

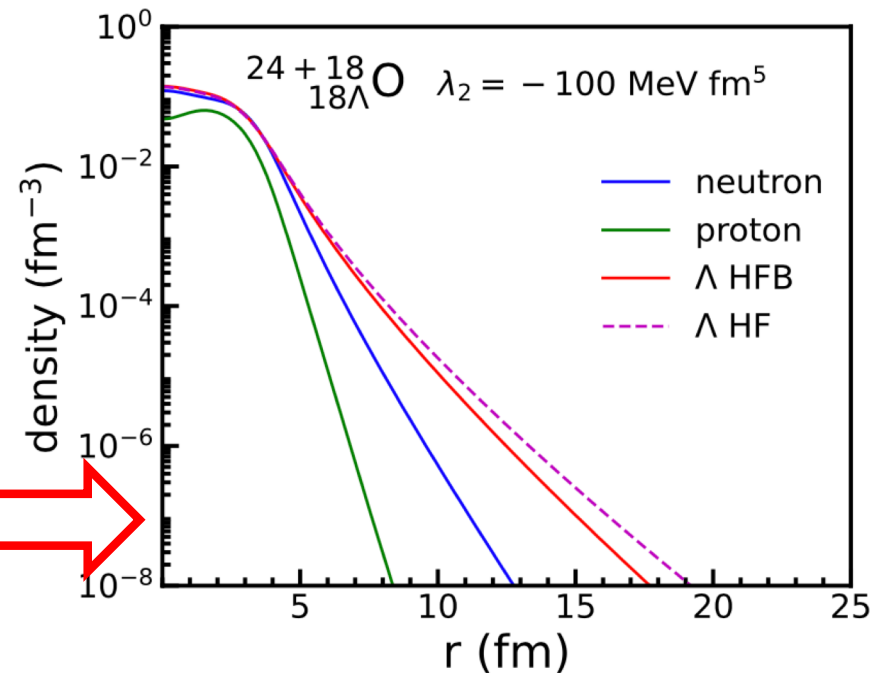
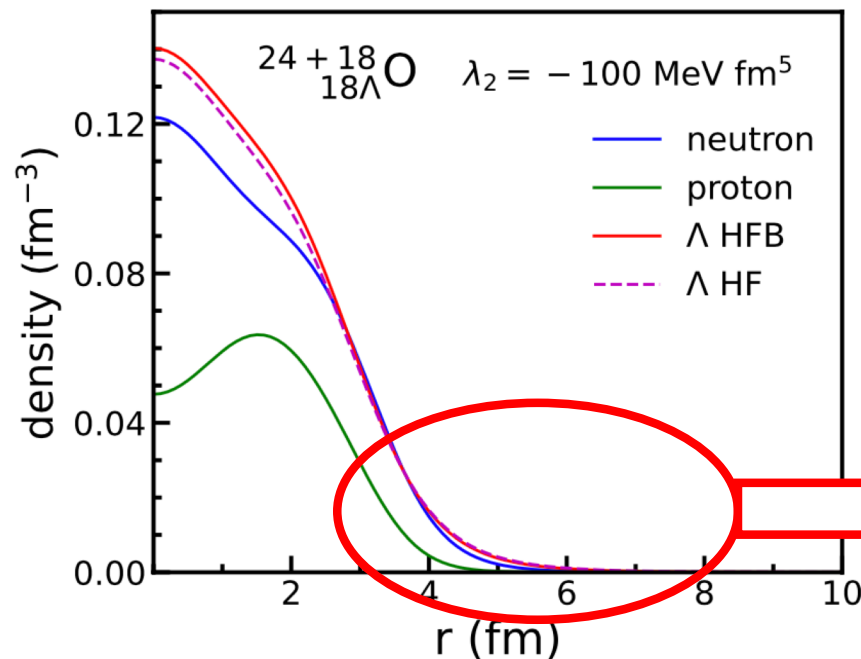
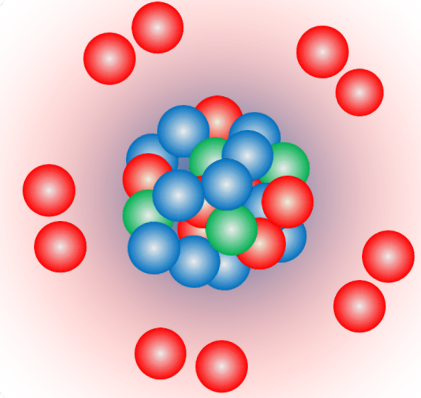


! Large tail of hyperon density in HF cal. \rightarrow large hyperon RMS radius
! "pairing anti-halo" effect

Possible hyperon halo in O hypernuclei

$$\lambda_2 = -100 \text{ MeV fm}^5$$

➤ Density of neutron, proton and hyperon in $^{24+18}_{18\Lambda}\text{O}$



! The weakly bound levels can hold up to $N=10$ hyperons, which may form the hyperon halo.

Summary

- Skyrme energy density functional theory for hypernucleus including ΛN interaction (with spin-orbit interaction) and $\Lambda\Lambda$ interaction was developed. The spin-orbit splitting and the $\Lambda\Lambda$ binding energy data was reproduced.
- We studied the possible halo orbits of Λ hyperon in C, Zr and O isotopes.
 - ✓ C, halo orbits: 1p, 2s (neutron-rich)
 - ✓ Zr, halo orbits: 3s, 2d, 1g, almost degenerate. They can hold even more hyperons than neutrons in the giant neutron halo.
 - ✓ $^{24}\text{O} + \Lambda$ hyperon: $\Lambda\Lambda$ interaction and pairing are important to determine the Λ drip line and the Λ halo.

Thank you!

Appendix

Introduction

Nuclear Landscape

stable nuclei

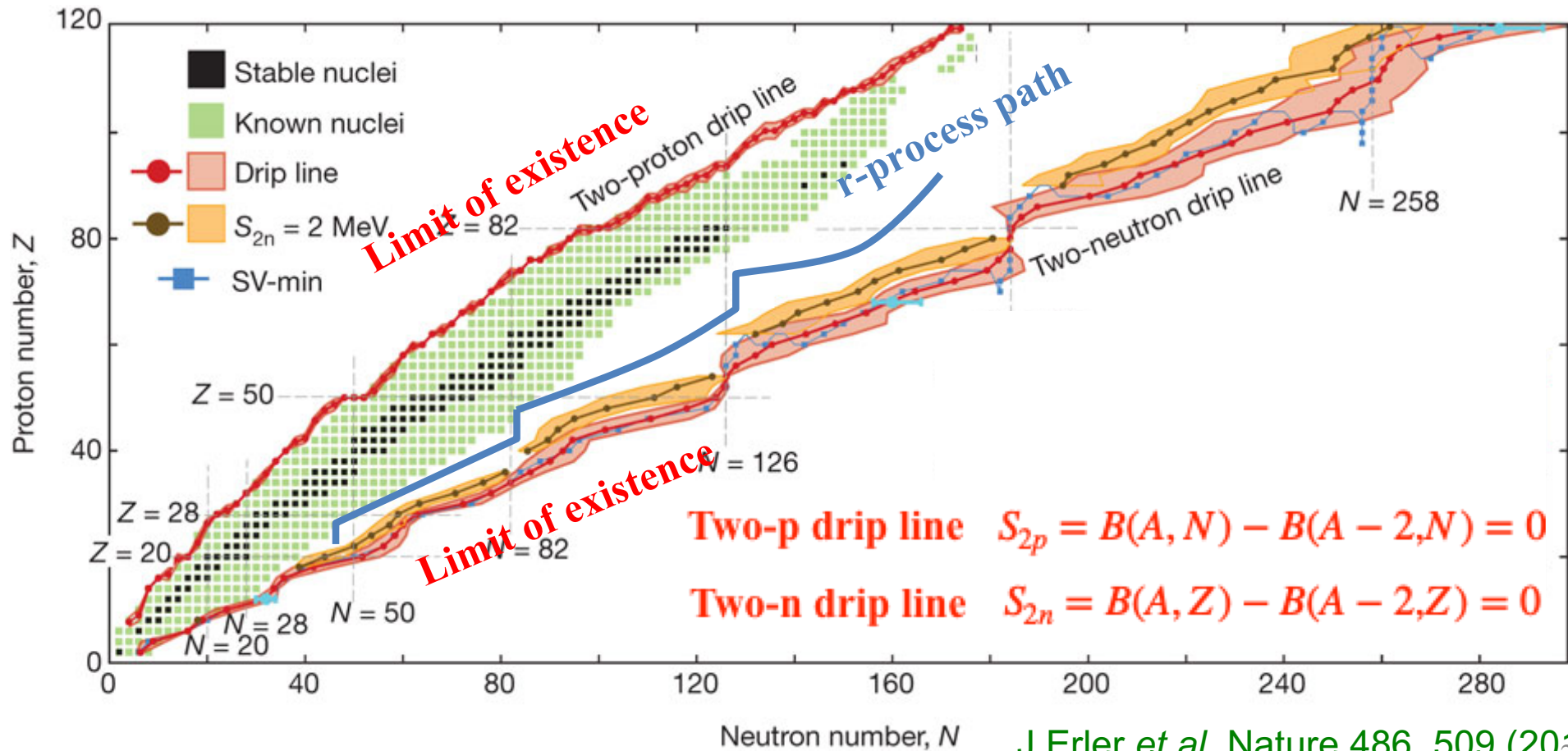
~300 nuclei

unstable nuclei observed so far

~2700 nuclei

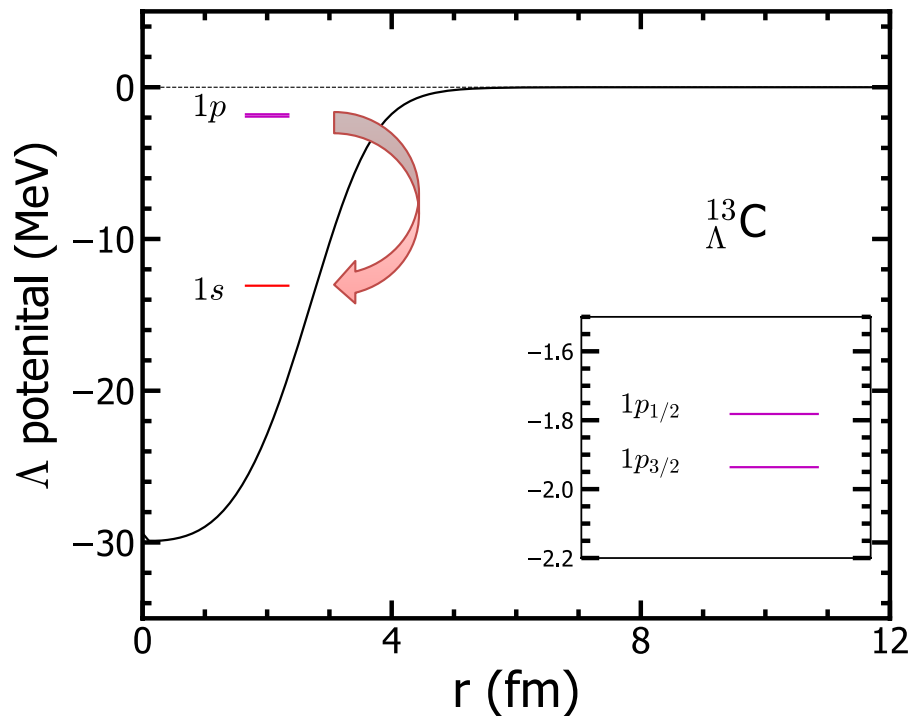
drip-lines (limit of existence) (theoretical predictions)

~8000 nuclei



Hyperon halo orbits in C hypernuclei

★ E1 transition from $1p$ to $1s$ state



$$B(EI; J_i \rightarrow J_f) = \frac{1}{2J_i + 1} |\langle f || \hat{Q}_I || i \rangle|^2$$

$$\langle f || \hat{Q}_I || i \rangle = \langle f || e r^I Y_I(\theta, \varphi) || i \rangle$$

$$B(E1; J_i \rightarrow J_f) = \frac{3e_{\Lambda}^2}{4\pi} \langle f | r | i \rangle^2 (2j_f + 1) \begin{pmatrix} j_f & 1 & j_i \\ -\frac{1}{2} & 0 & \frac{1}{2} \end{pmatrix}^2$$

$$\langle f | r | i \rangle = \int_0^{\infty} R_{f,\Lambda}(r) r R_{i,\Lambda}(r) dr.$$

$$e_{\Lambda}^{(E1)} = -Z M_{\Lambda} e / (A M_N + M_{\Lambda})$$

$$^{13}_{\Lambda}\text{C} \quad \text{Cal. } \Lambda (1p_{1/2} \rightarrow 1s_{1/2}): \quad B(E1) = 0.103 \text{ e}^2 \text{ fm}^2$$

$$^{11}\text{Be} \quad \text{Exp. } n \left(\frac{1}{2}^+ \rightarrow \frac{1}{2}^- \right): \quad B(E1) = 0.099 \pm 0.010 \text{ e}^2 \text{ fm}^2$$

T. Nakamura et al., Phys. Lett. B 394, 11 (1997).

Hyperon halo orbits in Zr hypernuclei

■ Hartree-Fock-Bogoliubov (HFB) model

NN interaction: SkI4 P.-G. Reinhard and H. Flocard, NPA 584, 467 (1995).

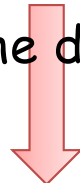
NN pairing:
$$\Delta(\mathbf{r}) = \frac{1}{2} V_0 \left[1 - \eta \left(\frac{\rho_q(\mathbf{r})}{\rho_0} \right)^\alpha \right] \tilde{\rho}(\mathbf{r}), \quad q = n \text{ or } p.$$

$V_0 = -300 \text{ MeV fm}^3$, $\eta = 0.5$, $\alpha = 0.5$, $\rho_0 = 0.16 \text{ fm}^{-3}$, $E_{\text{cut}} = 70 \text{ MeV}$ M. Grasso et al. PRC 74(2006)64317.

N Λ interaction: LY5r Y. Z., H. Sagawa, and E. Hiyama, Phys. Rev. C 103, 034321 (2021)

$\Lambda\Lambda$ interaction: S $\Lambda\Lambda$ 1 $\lambda_0 = -312.6 \text{ MeV fm}^3$ D. E. Lanskoy, PRC 58, 3351 (1998).

to reproduce the data of NAGARA event: ${}_{\Lambda\Lambda}^6\text{He}$



K. Nakazawa and H. Takahashi, PTPS 185, 335 (2010).

S $\Lambda\Lambda$ 1r $\lambda_0 = -50.0 \text{ MeV fm}^3$ $\lambda_1 = 57.5 \text{ MeV fm}^5$ $\lambda_2 = 0$

Table 1. The calculated double- Λ binding energy $B_{\Lambda\Lambda}^{\text{cal.}}$ and the gain energy $\Delta B_{\Lambda\Lambda}^{\text{cal.}}$ for several light double- Λ hypernuclei together with the experimental data [35,37]. All the energies are in MeV.

Nuclide	$B_{\Lambda\Lambda}^{\text{cal.}}$	$B_{\Lambda\Lambda}^{\text{expt.}}$	$\Delta B_{\Lambda\Lambda}^{\text{cal.}}$	$\Delta B_{\Lambda\Lambda}^{\text{expt.}}$
${}_{\Lambda\Lambda}^6\text{He}$	7.564	6.91 ± 0.16 (NAGARA)	0.536	0.67 ± 0.17
${}_{\Lambda\Lambda}^{11}\text{Be}$	18.594	20.86 ± 3.06 (MIKAGE)	0.478	2.64 ± 3.09
		22.12 ± 2.67 (MIKAGE)		3.90 ± 2.71
		20.83 ± 1.27 (HIDA)		2.61 ± 1.34
		19.07 ± 0.11 (MINO)		1.87 ± 0.37
${}_{\Lambda\Lambda}^{13}\text{B}$	22.221	23.3 ± 0.7 (KEK-E176)	0.435	0.6 ± 0.8

$$\Delta B_{\Lambda\Lambda} = B_{\Lambda\Lambda} - 2B_{\Lambda}$$

$$B_{\Lambda} = B_{\Lambda}^{(A+1)Z} - B^{(A)Z}$$

$$B_{\Lambda\Lambda} = B_{\Lambda\Lambda}^{(A+2)Z} - B^{(A)Z}$$

Hyperon halo orbits in Zr hypernuclei

Table 2 Bulk properties of Zr isotopes with one or two Λ hyperons: rms radius $r_{\text{rms}}^n, r_{\text{rms}}^p, r_{\text{rms}}^\Lambda$ (fm), the average pairing gap Δ_n, Δ_p (MeV), single- or double- Λ binding energies $B_{\Lambda(\Lambda)}$ (MeV), and double- Λ gain energy $\Delta B_{\Lambda\Lambda}$ (MeV). The Λ hyperon(s) is placed in the ground state $1s_{1/2}$ orbit.

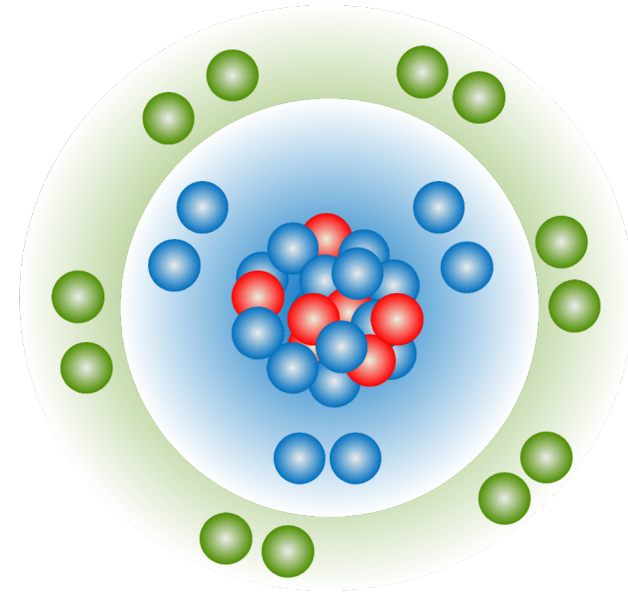
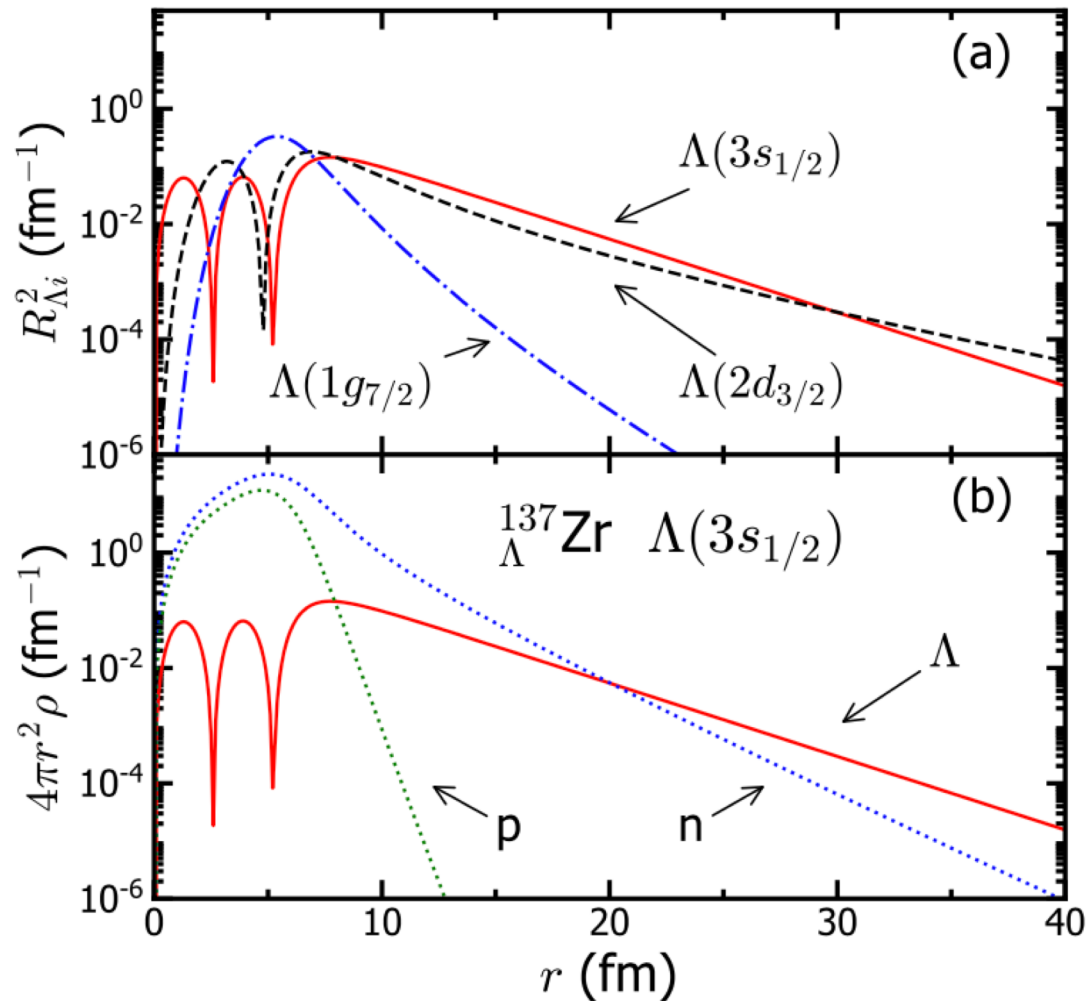
	r_{rms}^n	r_{rms}^p	r_{rms}^Λ	Δ_n	Δ_p	$B_{\Lambda(\Lambda)}$	$\Delta B_{\Lambda\Lambda}$
^{90}Zr	4.2674	4.1588		0.000	0.915		
$^{91}_{\Lambda}\text{Zr}$	4.2608	4.1515	3.1346	0.000	0.901	23.468	
$^{92}_{\Lambda\Lambda}\text{Zr}$	4.2543	4.1445	3.1197	0.000	0.889	47.123	0.188
^{114}Zr	4.8000	4.4112		0.953	0.796		
$^{115}_{\Lambda}\text{Zr}$	4.7934	4.4036	3.3059	0.953	0.785	23.868	
$^{116}_{\Lambda\Lambda}\text{Zr}$	4.7870	4.3962	3.2897	0.952	0.773	47.905	0.169
^{136}Zr	5.4649	4.5415		0.620	0.602		
$^{137}_{\Lambda}\text{Zr}$	5.4585	4.5328	3.4003	0.593	0.579	24.074	
$^{138}_{\Lambda\Lambda}\text{Zr}$	5.4520	4.5244	3.3823	0.568	0.556	48.315	0.168

much
smaller
than
 ${}_{\Lambda\Lambda}^6\text{He}$

DDDI nn (pp) pairing are included by the HFB calculation

$$\Delta(\mathbf{r}) = \frac{1}{2} V_0 \left[1 - \eta \left(\frac{\rho_q(\mathbf{r})}{\rho_0} \right)^\alpha \right] \tilde{\rho}(\mathbf{r}), \quad q = n \text{ or } p. \quad V_0 = -300 \text{ MeV fm}^3, \eta = 0.5, \alpha = 0.5, \\ \rho_0 = 0.16 \text{ fm}^{-3}, E_{\text{cut}} = 70 \text{ MeV}$$

Hyperon halo orbits in Zr hypernuclei



Giant Hyperon halo outside the giant neutron halo

Y. Z., H. Sagawa, and E. Hiyama, Prog. Theor. Exp. Phys. **2022** 023D01 (2022)

Possible hyperon halo in O hypernuclei

$$\lambda_2 = -100 \text{ MeV fm}^5$$

