



**NNU · 南京师范大学**  
NANJING NORMAL UNIVERSITY

正德厚生  
為學敏行

# 矢量底粲介子衰变的有效理论与 极化分析

朱瑞林

南京师范大学

Tao-Xiao-Zhu: 2209.15521; 2303.07220

Geng-Cao-Zhu: 2310.03425

2023年11月24-26日

重味物理前沿论坛@华中师大



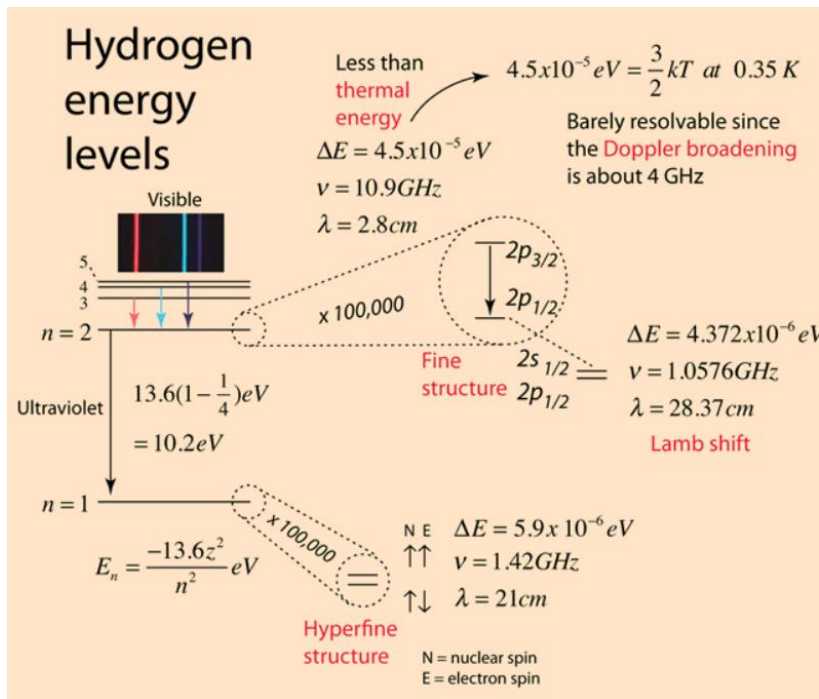
# Contents

---

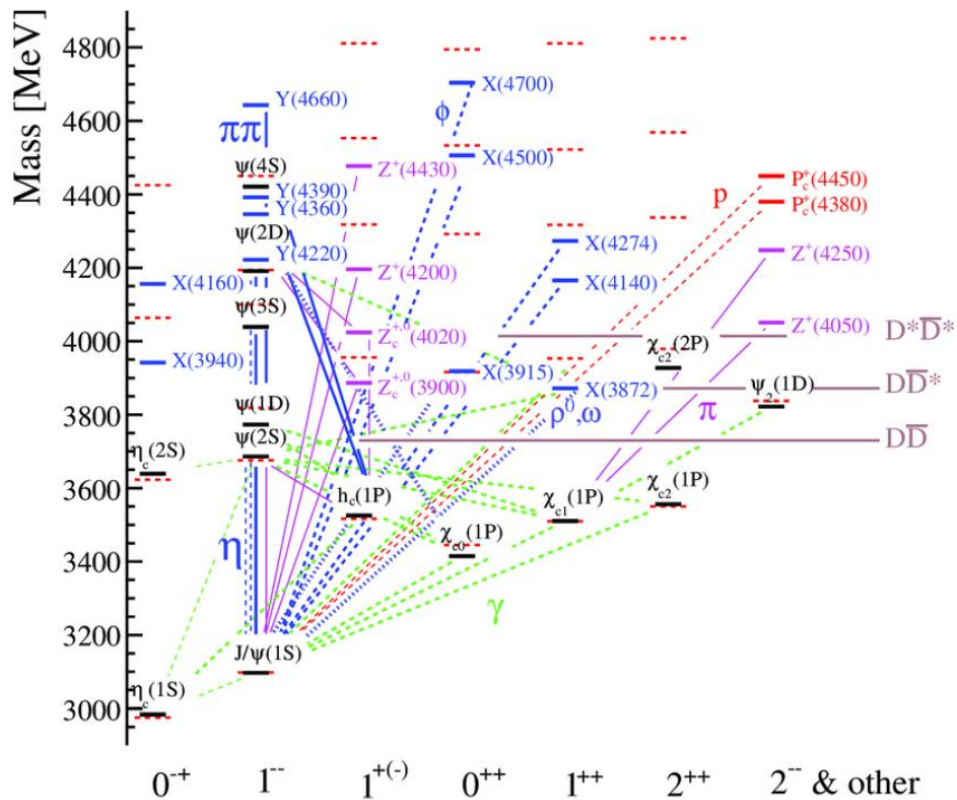
- **Research Background**
- **EFT (NRQCD/pNRQCD) Frameworks for  $B_c^*$  Decays**
- **Polarization Analysis in  $B_c^*$  Decays**
- **Summary and Outlook**

# 1、 Research Background

## Spectroscopy at angstrom scale and femto-scale



Hydrogen Spectrum  
QM  $\rightarrow$  QED

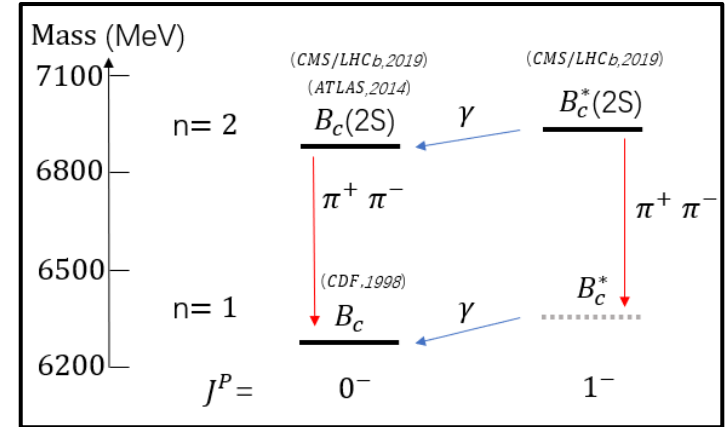


Charm region Spectrum  
QCD

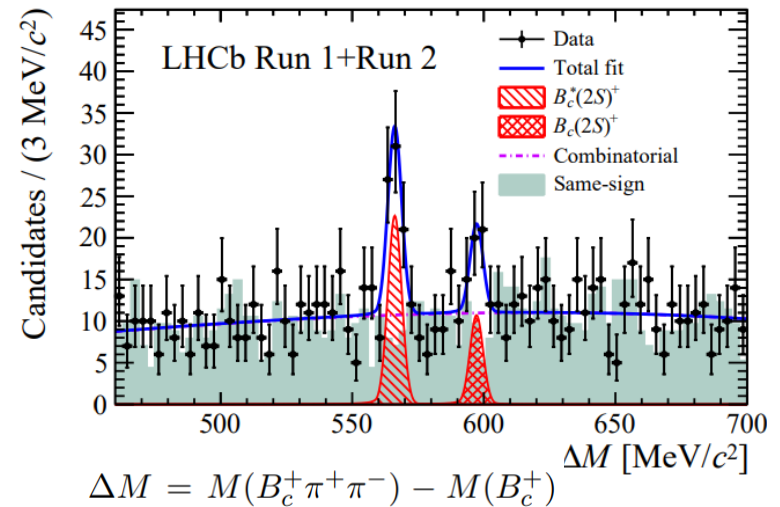
# Conventional Meson Spectroscopy

## Quark-antiquark S-wave meson puzzle *The Bc family is not complete*

$[u\bar{b}]$ $B^+, B^{*+}$	$[u\bar{c}]$ $\bar{D}^0, \bar{D}^{*0}$	$[u\bar{s}]$ $K^+, K^{*+}$	$[u\bar{d}]$ $\pi^+, \rho^+$	$[u\bar{u}]$ $\pi^0, \eta, \eta'$ $\rho^0, \omega, \phi$
$[d\bar{b}]$ $B^0, B^{*0}$	$[d\bar{c}]$ $D^-, D^{*-}$	$[d\bar{s}]$ $K^0, K^{*0}$	$[d\bar{d}]$ $\pi^0, \eta, \eta'$ $\rho^0, \omega, \phi$	
$[s\bar{b}]$ $B_s^0, B_s^{*0}$	$[s\bar{c}]$ $D_s^-, D_s^{*-}$	$[s\bar{s}]$ $\eta, \eta'$ $\omega, \phi$		
$[c\bar{b}]$ $B_c^+, B_c^{*+}$	$[c\bar{c}]$ $\eta_c, J/\psi$			
$[b\bar{b}]$ $\eta_b, \Upsilon$				

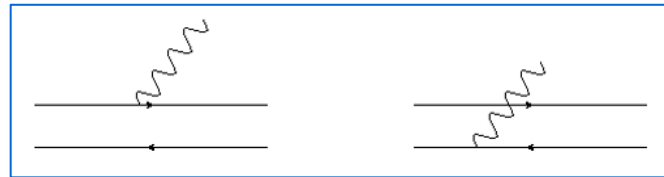
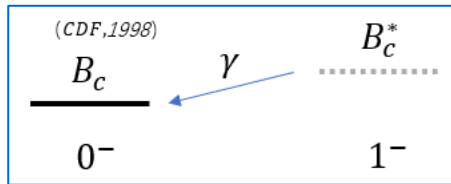


## Beauty-charm family



# Vector $B_c^*(1^-)$ Electromagnetic and Weak Decays

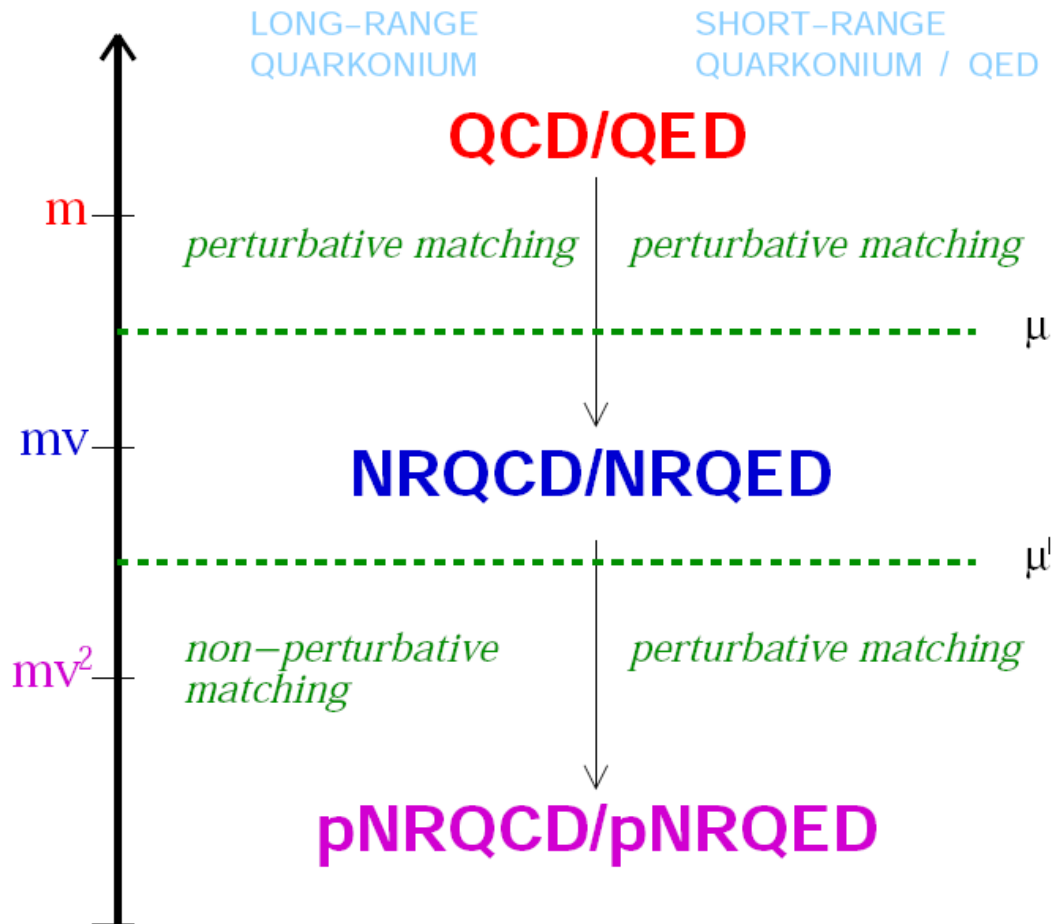
- Hyperfine splitting between  $B_c^*$  and  $B_c$ :  $\sim 60\text{MeV}$
- $B_c^*$  major (99.99%) electromagnetic decays to ground  $B_c$



- *However, 60MeV photon is hard to detect at LHC environment; current  $e^+ e^-$  colliders can not create two beauty and charm pairs*
- *Solid theoretical analysis and new observables are required*

## 2、 EFT Frameworks for $B_c^*$ Decays

### ➤ Nonrelativistic Effective Theory in QCD/QED



$$v^2 \approx 0.1 \text{ for the } \Upsilon$$

$$\alpha_s(mv) \sim v$$

Bodwin-Braaten-Lapage  
1995

Pineda-Soto-Brambilla-Vairo  
2000

# NRQCD/pNRQCD/ $\gamma$ pNRQCD in unequal mass case

## ➤ Heavy quark field in QCD

$$(i\gamma^\mu D_\mu - M)\Psi = 0, \quad (i\gamma^\mu D_\mu - M')\Psi' = 0$$

## ➤ Rewrite heavy quark field in QCD

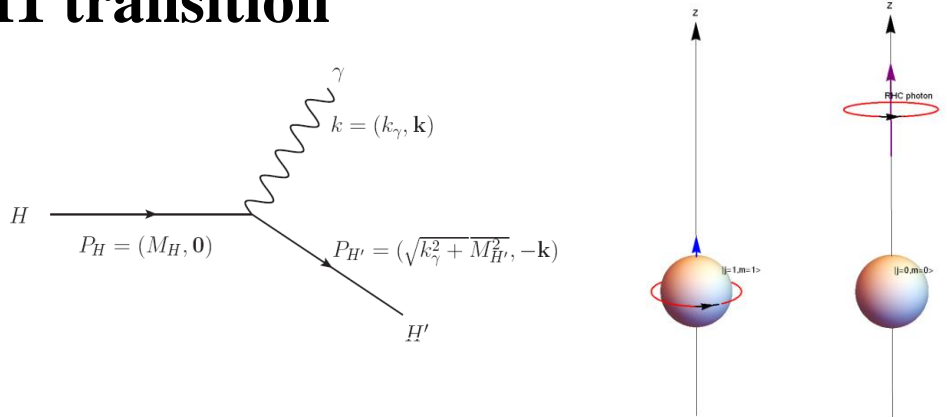
$$\Psi = e^{-iMt}\tilde{\Psi} = e^{-iMt} \begin{pmatrix} \psi \\ \chi \end{pmatrix}, \quad \Psi' = e^{iM't}\tilde{\Psi}' = e^{iM't} \begin{pmatrix} \psi' \\ \chi' \end{pmatrix},$$

## ➤ Obtain NRQCD Lagrangian

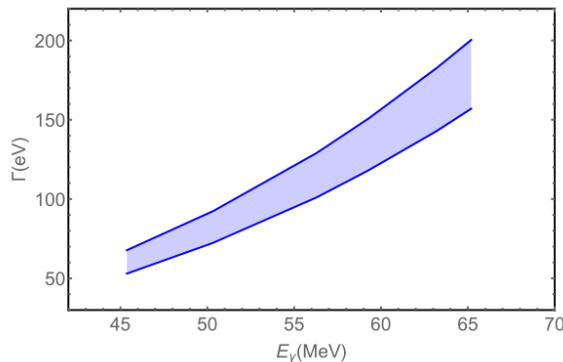
$$\begin{aligned} \mathcal{L}_{NRQCD} = & \psi^\dagger \left( iD_t - \frac{1}{2M}(i\mathbf{D})^2 \right) \psi + \frac{c_F}{2M} \psi^\dagger \boldsymbol{\sigma} \cdot g\mathbf{B} \psi \\ & + \frac{c_D}{8M^2} \psi^\dagger (\mathbf{D} \cdot g\mathbf{E} - g\mathbf{E} \cdot \mathbf{D}) \psi^\dagger + \frac{c_S}{8M^2} \psi^\dagger (i\boldsymbol{\sigma} \cdot \mathbf{D} \times g\mathbf{E} - i\boldsymbol{\sigma} \cdot g\mathbf{E} \times \mathbf{D}) \psi^\dagger \\ & + \frac{c_4}{8M^3} \psi'^\dagger (\mathbf{D}^2)^2 \psi' + \mathcal{O}(1/M^3) \\ & + [\psi \rightarrow i\sigma^2 \chi'^*, A_\mu \rightarrow -A_\mu^T, M \rightarrow M'] + \mathcal{L}_{light}. \end{aligned}$$

# Vector $B_c^*(1^-)$ Electromagnetic Decay

- $B_c^*(1S)$  major (99.99%) electromagnetic decays to  $B_c(1S)$ : M1 transition



- We generalize the pNRQCD to unequal mass case and obtain the effective Lagrangian



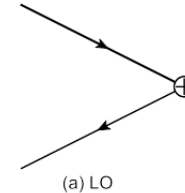
$$\begin{aligned}
 \mathcal{L}_{\gamma\text{pNRQCD}} = & \int d^3r \text{Tr} \left[ e \frac{e_Q - e'_Q}{2} V_A^{\text{em}} S^\dagger \mathbf{r} \cdot \mathbf{E}^{\text{em}} S \right. \\
 & + e \left( \frac{e_Q m'_Q - e'_Q m_Q}{4m_Q m'_Q} \right) \left[ V_S^{\frac{\sigma \cdot \mathbf{B}}{m}} \left\{ S^\dagger, \boldsymbol{\sigma} \cdot \mathbf{B}^{\text{em}} \right\} S \right. \\
 & + \frac{1}{8} V_S (r \cdot \nabla)^2 \frac{\sigma \cdot \mathbf{B}}{m} \left\{ S^\dagger, \mathbf{r}^i \mathbf{r}^j (\nabla^i \nabla^j \boldsymbol{\sigma} \cdot \mathbf{B}^{\text{em}}) \right\} S \\
 & \left. \left. + V_O^{\frac{\sigma \cdot \mathbf{B}}{m}} \left\{ O^\dagger, \boldsymbol{\sigma} \cdot \mathbf{B}^{\text{em}} \right\} O \right] \right. \\
 & + e \left( \frac{e_Q m_Q^2 - e'_Q m_Q^2}{32m_Q^2 m_Q'^2} \right) \left[ 4 \frac{V_S^{\frac{\sigma \cdot \mathbf{B}}{m^2}}}{r} \left\{ S^\dagger, \boldsymbol{\sigma} \cdot \mathbf{B}^{\text{em}} \right\} S \right. \\
 & + 4 \frac{V_S^{\frac{\boldsymbol{\sigma} \cdot (\mathbf{r} \times \mathbf{r} \times \mathbf{B})}}{m^2}}}{r} \left\{ S^\dagger, \boldsymbol{\sigma} \cdot [\hat{\mathbf{r}} \times (\hat{\mathbf{r}} \times \mathbf{B}^{\text{em}})] \right\} S \\
 & - V_S^{\frac{\boldsymbol{\sigma} \cdot \nabla \times \mathbf{E}}{m^2}} \left[ S^\dagger, \boldsymbol{\sigma} \cdot [-i \nabla \times, \mathbf{E}^{\text{em}}] \right] S \\
 & \left. - V_S^{\frac{\boldsymbol{\sigma} \cdot \nabla_r \times \mathbf{r} \cdot \nabla \mathbf{E}}{m^2}} \left[ S^\dagger, \boldsymbol{\sigma} \cdot [-i \nabla_r \times, \mathbf{r}^i (\nabla^i \mathbf{E}^{\text{em}})] \right] S \right] \\
 & + e \left( \frac{e_Q m_Q^3 - e'_Q m_Q^3}{8m_Q^3 m_Q'^3} \right) \left[ V_S^{\frac{\nabla_r^2 \sigma \cdot \mathbf{B}}{m^3}} \left\{ S^\dagger, \boldsymbol{\sigma} \cdot \mathbf{B}^{\text{em}} \right\} \nabla_r^2 S \right. \\
 & \left. \left. + V_S^{\frac{(\nabla_r \cdot \boldsymbol{\sigma})(\nabla_r \cdot \mathbf{B})}{m^3}} \left\{ S^\dagger, \boldsymbol{\sigma}^i \mathbf{B}^{\text{em} j} \right\} \nabla_r^i \nabla_r^j S \right] \right],
 \end{aligned}$$



# Bc\* decay constants in EFT

## ➤ Bc\* decay constants in QCD

$$\langle 0 | \bar{b} \gamma^\mu c | B_c^*(P, \varepsilon) \rangle = f_{B_c^*}^v m_{B_c^*} \varepsilon^\mu,$$



## ➤ Bc\* decay constants in NRQCD

$$f_{B_c^*}^v = \sqrt{\frac{2}{m_{B_c^*}}} \underbrace{C_v(m_b, m_c, \mu_f)}_{\text{matching coefficients}} \underbrace{\langle 0 | \chi_b^\dagger \sigma \cdot \varepsilon \psi_c | B_c^*(\mathbf{P}) \rangle}_{\text{NRQCD LDMEs}}(\mu_f) + O(v^2)$$

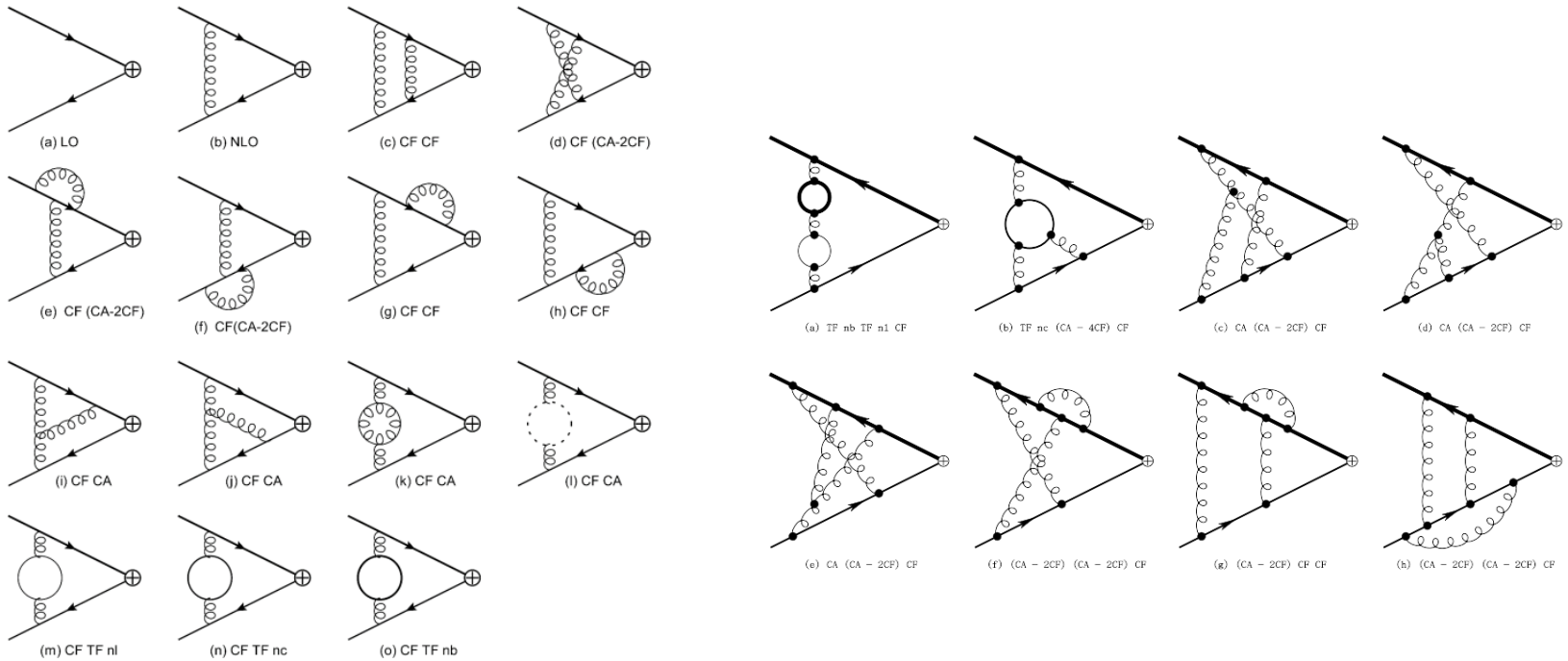
## ➤ Matching Formulae

Braaten-Fleming, PRD52,181(1995);  
Lee-Sang-Kim, JHEP01,113(2011)

$$Z_J Z_{2,b}^{\frac{1}{2}} Z_{2,c}^{\frac{1}{2}} \Gamma_J = C_J \tilde{Z}_J^{-1} \tilde{Z}_{2,b}^{\frac{1}{2}} \tilde{Z}_{2,c}^{\frac{1}{2}} \tilde{\Gamma}_J$$

$\tilde{Z}_J$ : NRQCD  $\overline{\text{MS}}$  current renormalization constants

# Typical diagrams up to three-loop



LO:1, NLO:1, NNLO:11, N<sup>3</sup>LO:268

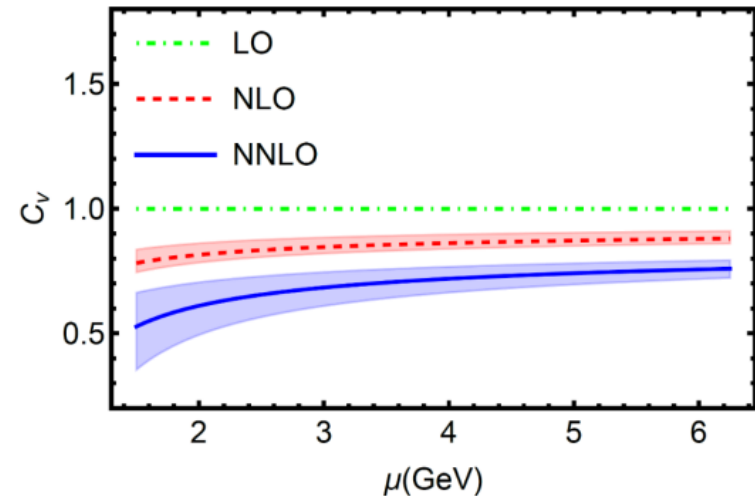
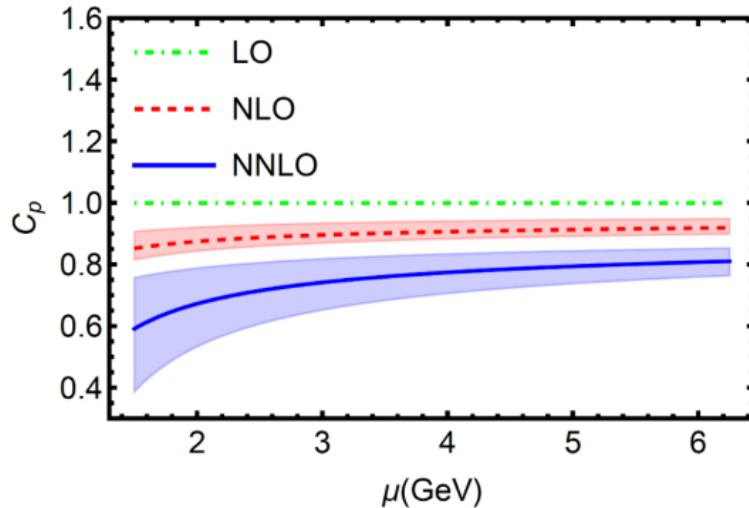
# Two-loop results

## ➤ Matching coefficients at two loop

$\mu_f \in [1.5, 1.2, 1] \text{ GeV}$ ,  $\mu \in [6.25, 4.75, 3] \text{ GeV}$ ,  $m_b \in [5.25, 4.75, 4.25] \text{ GeV}$ ,  $m_c \in [2, 1.5, 1] \text{ GeV}$

	LO	NLO	NNLO
$C_p$	1	$0.9117^{+0.0072+0.0061-0.0156}_{-0-0.0160-0.0064+0.0263}$	$0.7897^{+0.0206+0.0119+0.0149}_{-0.0310-0.0482-0.0133-0.0141}$
$C_v$	1	$0.8697^{+0.0107+0.0061-0.0156}_{-0-0.0236-0.0064+0.0263}$	$0.7363^{+0.0230+0.0106+0.0117}_{-0.0234-0.0526-0.0117-0.0121}$

$\mu$  dependence for matching coefficients



# Three loop results for vector and pseudoscalar currents

## ➤ Matching coefficients for vector current

$$C = 1 - 2.29 \left( \frac{\alpha_s^{(n_l)}}{\pi} \right) - 35.44 \left( \frac{\alpha_s^{(n_l)}}{\pi} \right)^2 - 1686.27 \left( \frac{\alpha_s^{(n_l)}}{\pi} \right)^3 + \mathcal{O}(\alpha_s^4),$$

for  $n_l = 3, n_c = 1, n_b = 0,$

Sang-Zhang-Zhou, arXiv:2210.02979

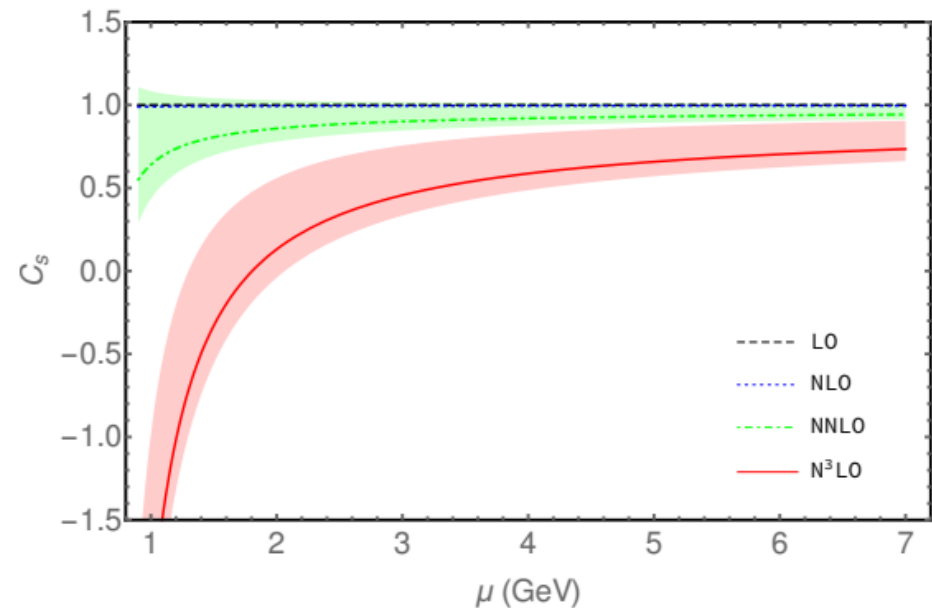
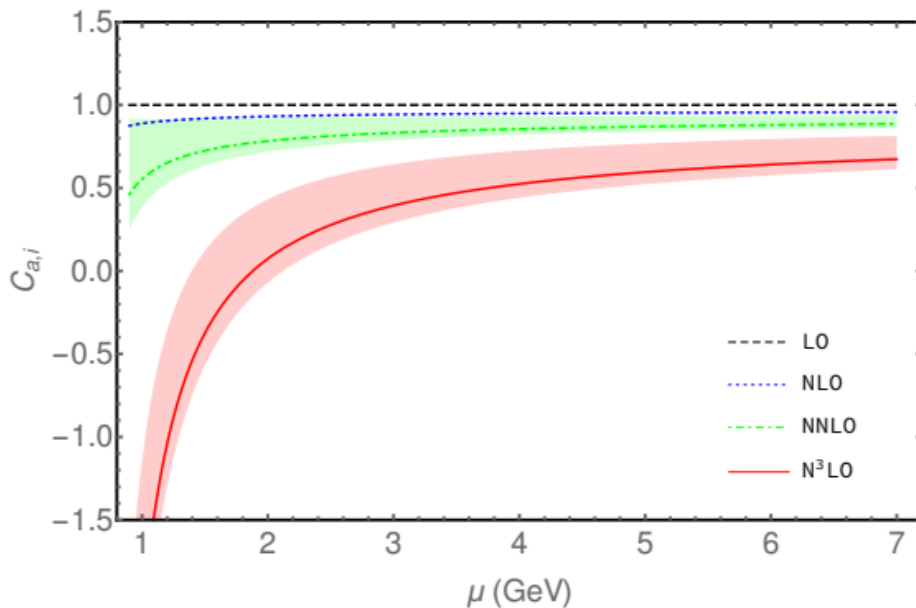
## ➤ Matching coefficients for pseudoscalar current

$$C(x_{\text{phys}}) = 1 - 1.62623 \left( \frac{\alpha_s^{(n_l)}(m_r)}{\pi} \right) - 6.51043 \left( \frac{\alpha_s^{(n_l)}(m_r)}{\pi} \right)^2 - 1520.59 \left( \frac{\alpha_s^{(n_l)}(m_r)}{\pi} \right)^3 + \mathcal{O}(\alpha_s^4)$$

Feng-Jia-Mo-Pan-Sang-Zhang, arXiv:2208.04302

# Results for axial-vector and scalar currents

## ➤ Matching coefficients for axial-vector and scalar up to three loop



Nonconvergence behaviors also in other two currents

Multi-loop integral calculation performed by AMFlow (Ma et al)

## Sub-leading Contribution

### ➤ Relativistic corrections

$$\begin{aligned} & \langle 0 | \bar{Q}_1 \gamma^5 Q_2 | Q_2 \bar{Q}_1 \rangle_{\text{QCD}} \\ &= \sqrt{2M_H} \left[ C_0^P \langle 0 | \chi_1^\dagger \psi_2 | Q_2 \bar{Q}_1(\mathbf{p}) \rangle_{\text{NRQCD}} + C_2^P \langle 0 | (\mathbf{D}\chi_1)^\dagger \cdot \mathbf{D}\psi_2 | Q_2 \bar{Q}_1(\mathbf{p}) \rangle_{\text{NRQCD}} + \dots \right] \end{aligned}$$

Employing EOM:  $\langle 0 | (\mathbf{D}\chi_1)^\dagger \cdot \mathbf{D}\psi_2 | Q_2 \bar{Q}_1(\mathbf{p}) \rangle = -2m_r E \langle 0 | \chi_1^\dagger \psi_2 | Q_2 \bar{Q}_1(\mathbf{p}) \rangle.$

$$f_{B_c^*} = 2 \sqrt{\frac{N_c}{m_{B_c^*}}} \left[ C_v + \frac{d_v E_{B_c^*}}{12} \left( \frac{8}{M} - \frac{3}{m_r} \right) \right] |\Psi_{B_c^*}(0)|,$$

$$f_{B_c} = 2 \sqrt{\frac{N_c}{m_{B_c}}} \left[ C_p - \frac{d_p E_{B_c}}{4m_r} \right] |\Psi_{B_c}(0)|,$$

# Wave function scale dependence

## ➤ Wave function at origin

For Power-law potential  $V(r) = Ar^a + C$ .

Exact solution  $|\psi_\mu^n(0)|^2 = f(n, a)(A\mu)^{3/(2+a)}$

Scale relation  $|\Psi_{B_c^*}(0)| = |\Psi_{J/\psi}(0)|^{1-y} |\Psi_\Upsilon(0)|^y,$

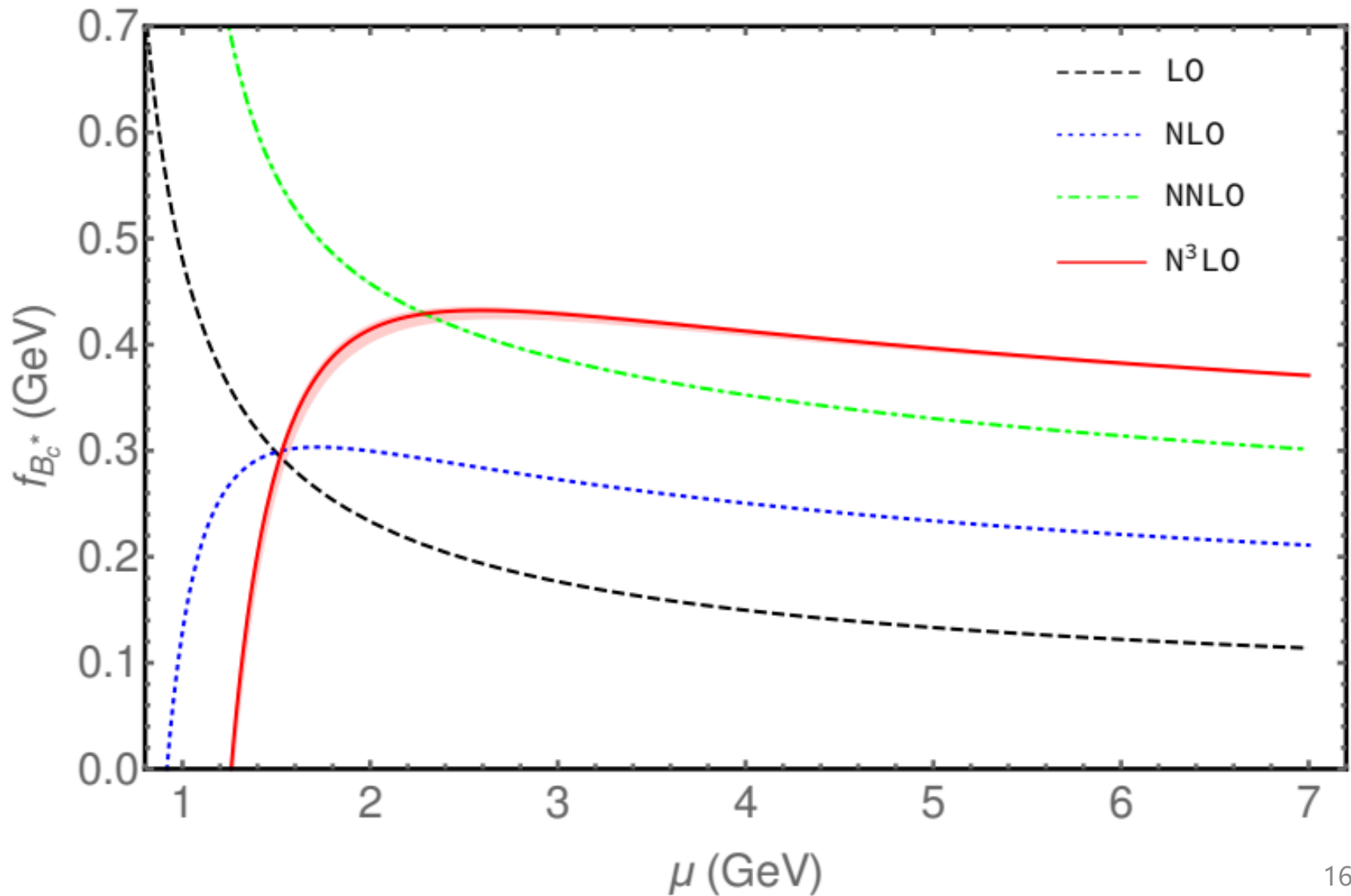
$$y = y_c = \ln((1 + m_c/m_b)/2) / \ln(m_c/m_b)$$

Collins-Imbo-King-Martell, PLB 393 (1997) 155–160

$$|\psi_1(0)|^2 = |\psi_1^{(0)}(0)|^2 \left( 1 + \sum_{k=1}^n f_k a_s^k \right). \quad \begin{aligned} |\psi_1^{(0)}(0)|^2 &= \frac{(m_b C_F \alpha_s)^3}{8\pi}, \\ E_1^{(0)} &= -\frac{1}{4} m_b (C_F \alpha_s)^2, \end{aligned}$$

Beneke et al., PRL. 112, 151801 (2014)

# Convergent vector $B_c^*$ decay constant





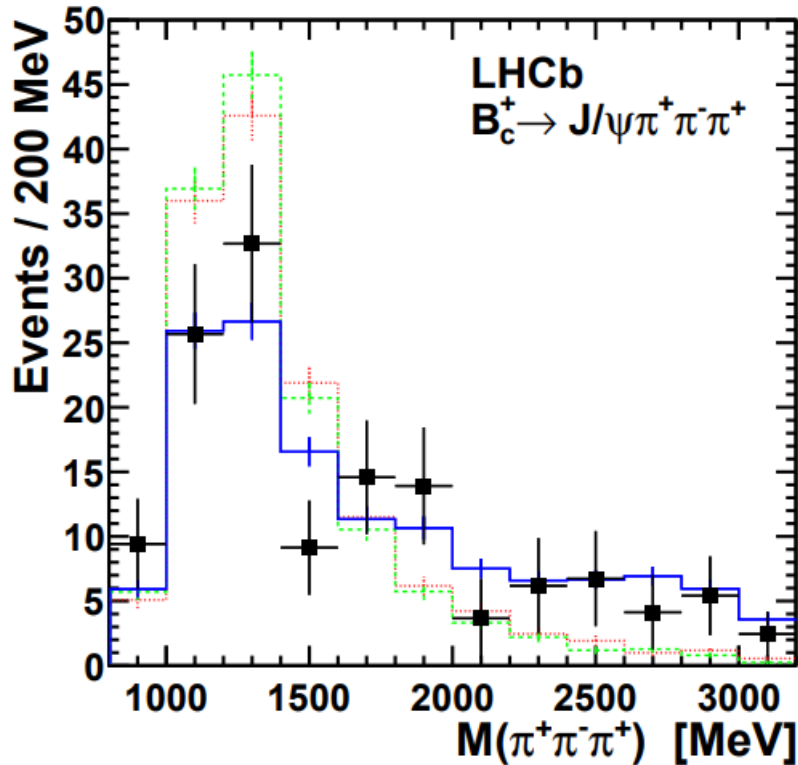
# Leptonic decay branching ratios

Branching ratios	N <sup>3</sup> LO
$\mathcal{B}(B_c^{*+} \rightarrow e^+ \nu_e)$	$(3.85_{-0.46+0.03+0.37}^{+0.29-0.07-1.35}) \times 10^{-6}$
$\mathcal{B}(B_c^{*+} \rightarrow \mu^+ \nu_\mu)$	$(3.85_{-0.46+0.03+0.37}^{+0.29-0.07-1.35}) \times 10^{-6}$
$\mathcal{B}(B_c^{*+} \rightarrow \tau^+ \nu_\tau)$	$(3.40_{-0.41+0.03+0.33}^{+0.25-0.06-1.19}) \times 10^{-6}$
$\mathcal{B}(B_c^+ \rightarrow e^+ \nu_e)$	$(1.91_{-0.23+0.12+0.22}^{+0.15-0.19-0.70}) \times 10^{-9}$
$\mathcal{B}(B_c^+ \rightarrow \mu^+ \nu_\mu)$	$(8.18_{-1.00+0.52+0.94}^{+0.63-0.83-2.99}) \times 10^{-5}$
$\mathcal{B}(B_c^+ \rightarrow \tau^+ \nu_\tau)$	$(1.96_{-0.24+0.12+0.23}^{+0.15-0.20-0.72}) \times 10^{-2}$

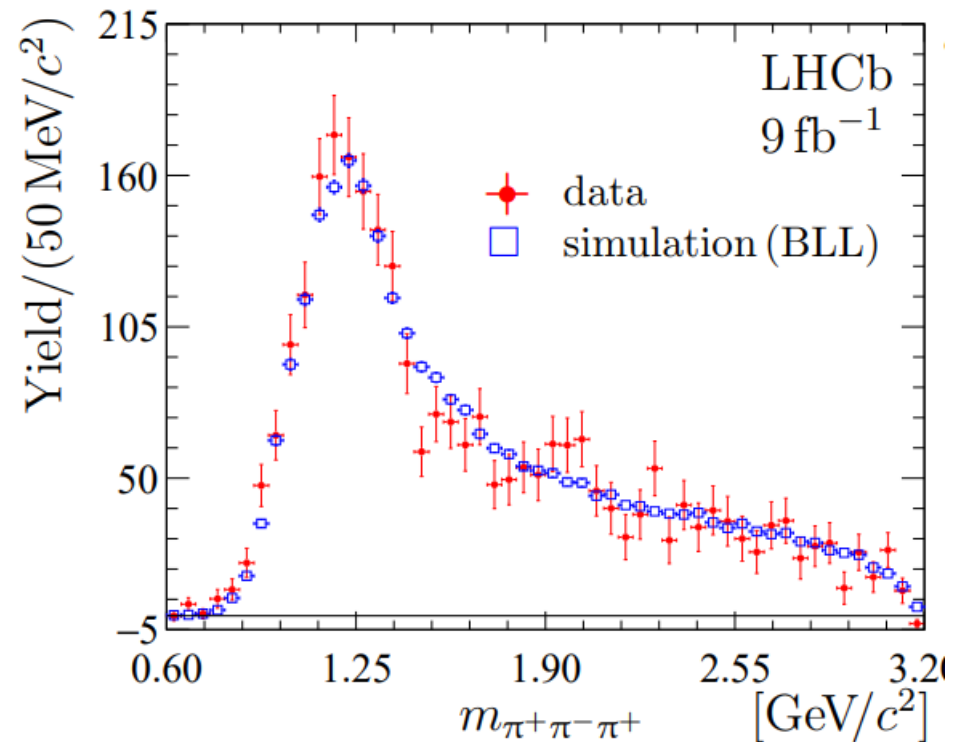
$$\Gamma(B_c^*(\lambda = \pm 1) \rightarrow \ell \nu_\ell) = \frac{|V_{cb}|^2}{12\pi} G_F^2 f_{B_c^*}^2 \left(1 - \frac{m_\ell^2}{m_{B_c^*}^2}\right)^2 \times m_{B_c^*}^3,$$

$$\Gamma(B_c^*(\lambda = 0) \rightarrow \ell \nu_\ell) = \frac{m_\ell^2 \Gamma(B_c^{*+}(\lambda = \pm 1) \rightarrow \ell \nu_\ell)}{2m_{B_c^*}^2},$$

### 3、 Polarization analysis of other $B_c^*$ decays



LHCb, arXiv:1204.0079



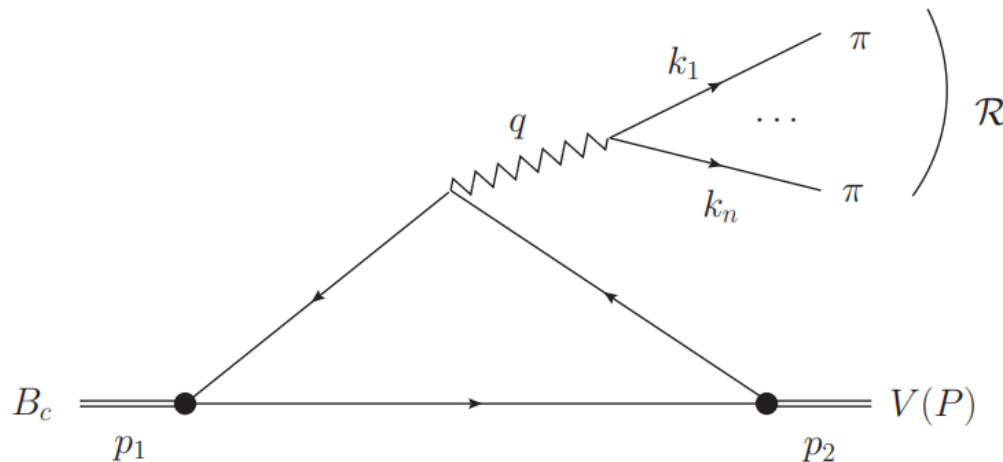
LHCb, arXiv:2111.03001  
Around  $10^5$   $B_c$  to  $J/\psi + X$  events

# Invariant mass distribution in $B_c^*$ decays to $J/\psi + n h$

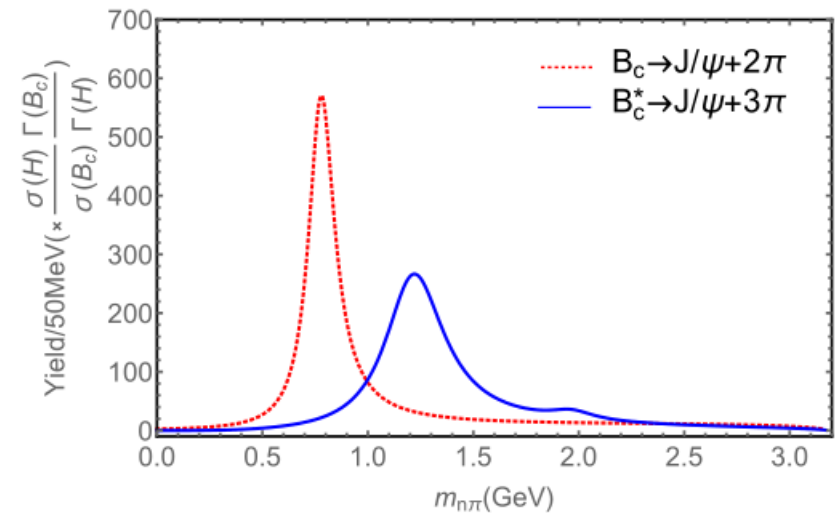
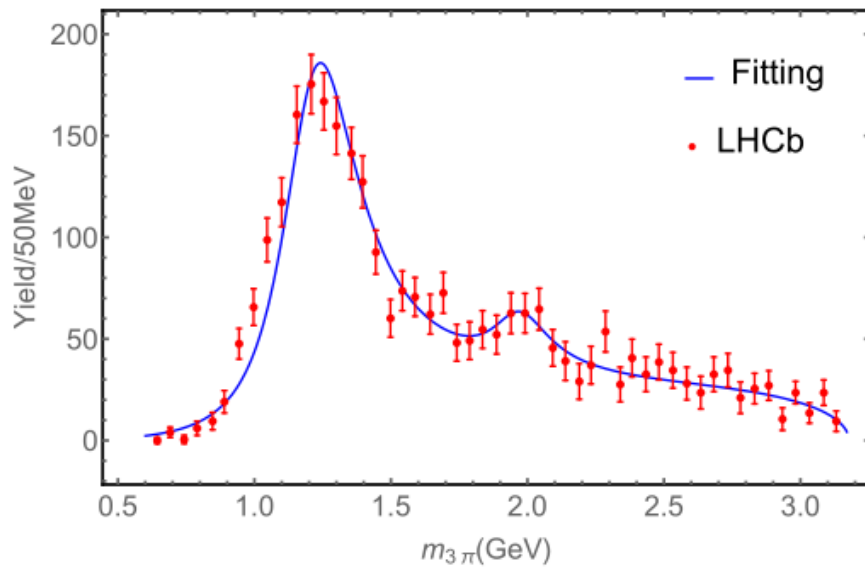
## ➤ Helicity decomposition of weak decay width

$$\frac{d\Gamma(B_c^{(*)} \rightarrow J/\psi + nh)}{dq^2} = \sum_{\lambda_i} \frac{|V_{cb}|^2 G_F^2 a_1^2 |\mathbf{P}'|}{32\pi M^2} \Gamma_{J_1 \lambda_1 J_2 \lambda_2 \lambda_{nh}},$$

$$\Gamma_{111110} = 2 \left[ V_1^2 \left( (M - M')^2 - q^2 \right) \left( (M' + M)^2 - q^2 \right) + (A_1 (M^2 - M'^2) + A_2 q^2)^2 \right] \rho_T^{nh}(q^2),$$



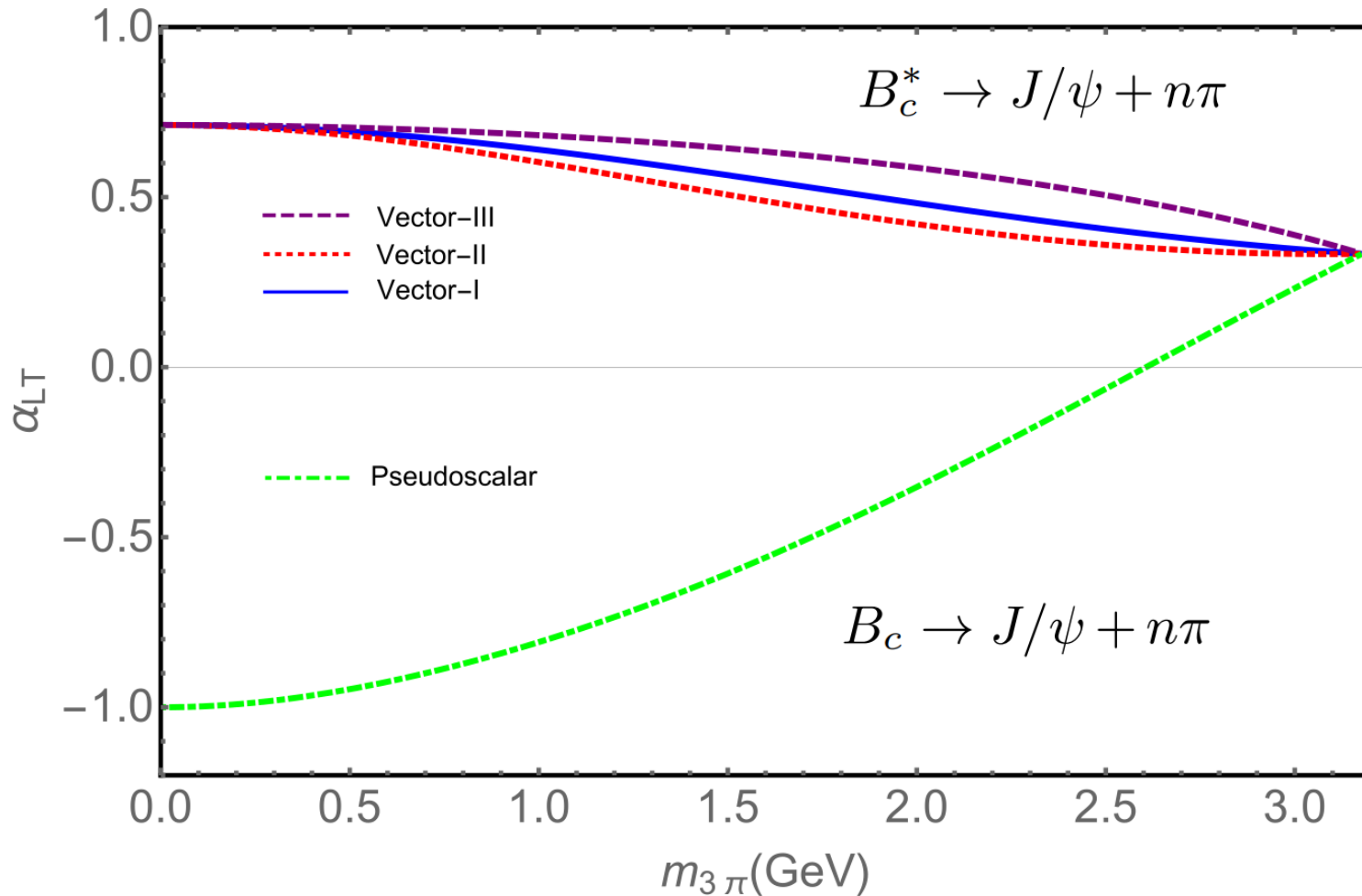
# Results of invariant mass distribution



LHCb, arXiv:2111.03001

# Polarization Asymmetry (A general law in V(P) to V transitions)

$$\alpha_{LT} = \sum_{\lambda_1, \lambda_{nh}} \frac{\Gamma_{J_1 \lambda_1 11 \lambda_{nh}} - \Gamma_{J_1 \lambda_1 10 \lambda_{nh}}}{\Gamma_{J_1 \lambda_1 11 \lambda_{nh}} + \Gamma_{J_1 \lambda_1 10 \lambda_{nh}}},$$



# Summary and Outlook

## Summary

- ✓  $B_c^*$  decay width is studied in QCD effective theory
- ✓ Convergent  $B_c^*$  decay constant up to three-loop accuracy is obtained
- ✓ Distinguishing vector  $B_c^*$  meson at LHC is possible by helicity decomposition

## Outlook

- Nontrivial high-order calculation of wave function in un-equal mass cases
- Experimental analysis of  $B_c$  family at LHC/CEPC/Super-Z/EicC

**Thank you a lot!**

# Hyperfine splitting

## ➤ Hyperfine splitting relation

$$(\Delta M)_{i\bar{j}} = 32\pi\alpha_s(2\mu_{i\bar{j}})|\Psi_{i\bar{j}}(0)|^2/9m_i m_j,$$

$$\Delta M_{c\bar{b}} = \alpha_s(2m_r)x^{1-2q} \left( \frac{\Delta M_{c\bar{c}}}{\alpha_s(m_c)} \right)^{1-q} \left( \frac{\Delta M_{b\bar{b}}}{\alpha_s(m_b)} \right)^q.$$

$$\Delta M_{c\bar{b}(1S)} = 63.8_{-8.4}^{+5.5}(q)_{-1.2}^{+1.2}(exp) \text{ MeV},$$

$$\Delta M_{c\bar{b}(2S)} = 26.4_{-3.3}^{+2.1}(q)_{-1.7}^{+1.5}(exp) \text{ MeV},$$

$$\Delta M_{b\bar{b}(1S)} = 62.3 \pm 3.2 \text{ MeV}$$

$$\Delta M_{b\bar{b}(2S)} = 24 \pm 4 \text{ MeV}$$

CMS 2019

$$\Delta M_{c\bar{b}(2S)} = 29.1 \pm 1.5(stat) \pm 0.7(syst) \text{ MeV}$$

LHCb 2019

$$\Delta M_{c\bar{b}(2S)} = 31.0 \pm 1.4(stat) \pm 0.0(syst) \text{ MeV}$$

HPQCD lattice results

$$\Delta M_{c\bar{b}(1S)} = 54 \pm 4 \text{ MeV}$$

# Calculation procedure

---

- Feynman Diagrams & Amplitudes  
(Packages: FeynRules/FeynArts / QGraf)
- Feynman Amplitudes Simplification: Trace & Contraction  
(Packages: FeynCalc / FormCalc / FormLink)
- Feynman Integrals Reduction  
(Packages: Apart(Feng) / FIRE/Kira /...)
- Feynman Master Integrals Calculation:  
(Packages: AMFlow(Ma et al) / FIESTA /...)