# Highlights of hadron physics@ ₩5III

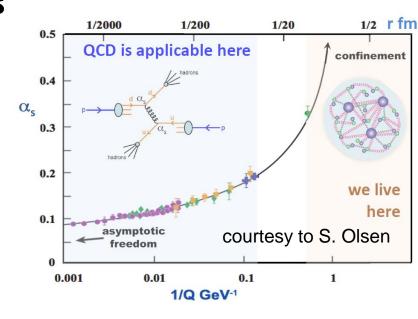
Beijiang Liu (on behalf of BESIII)

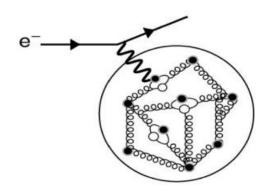
Institute of High Energy Physics, Chinese Academy of Sciences

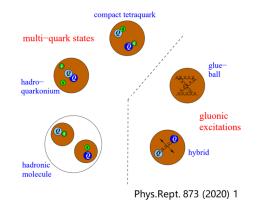
The 12th Workshop on Hadron Physics and Opportunities Worldwide, Dalian, 2024

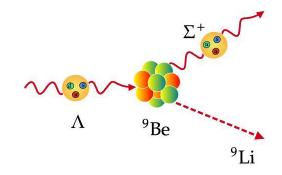
### Hadrons

- Dominant part of visible matter in the universe
- To fully understand the strong interaction
  - Understanding the rich and complex features of its bound states, hadrons
  - ➤ How are hadrons formed from quarks and gluons?
  - ➤ What is the origin of confinement?
  - ➤ How is the mass of hadron generated in QCD?
  - ➤ What is the dynamics of effective DoF in hadrons?









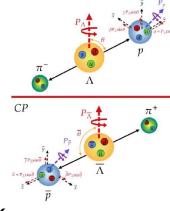
**Structure** 

**Spectroscopy** 

**Interactions** 

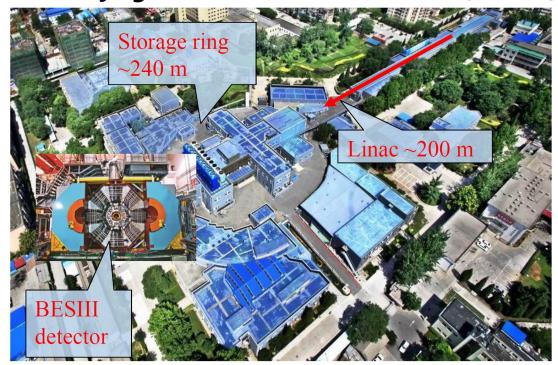
**Hadron physics** 





### BESIII@BECPII

#### **Beijing Electron Positron Collider(BEPCII)**



Double-ring, symmetry, multi-bunch e<sup>+</sup> e<sup>-</sup> collider

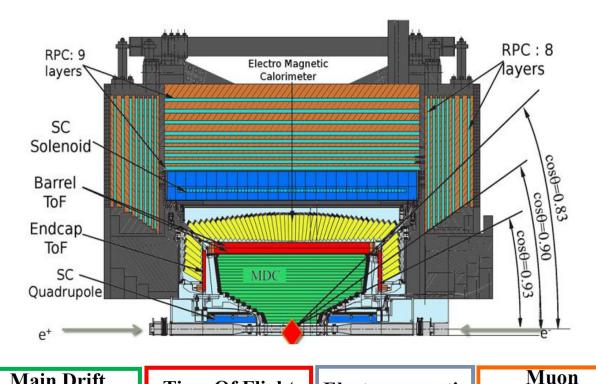
 $E_{cm} = 1.84 \text{ to } 4.95 \text{ GeV}$ 

Energy spread:  $\Delta E \approx 5 \times 10^{-4}$ 

Peak luminosity in continuously operation @E<sub>cm</sub>=

3.77 GeV:  $1.1 \times 10^{33} \text{ cm}^{-2} \text{s}^{-1}$ 

#### **Beijing Spectrometer(BESIII)**



Main Drift Chamber Small cell, 43 layer  $\sigma_{xy}$ =130  $\mu$ m dE/dx~6%  $\sigma_p/p$ = 0.5% at 1

GeV

Time Of Flight
Plastic scintillator  $\sigma_T$ (barrel): 65 ps  $\sigma_T$ (endcap): 110 ps
(update to 60 ps
with MRPC)

Calorimeter
CsI(Tl): L=28 cm
Barrel  $\sigma_E$ =2.5%
Endcap  $\sigma_E$ =5.0%

Electromagnetic

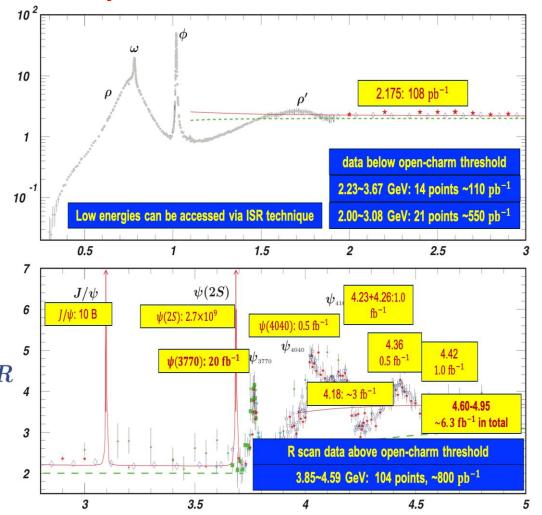
Counter RPC
Barrel: 9 layers
Endcap: 8
layers  $\sigma_{spatial}$ : 1.48
cm

### World's largest $\tau$ – charm data sets in $e^+e^-$ annihilation

#### Data sets collected so far include

- $> 10 \times 10^9$  J/ $\psi$  events
- $\geq 2.7 \times 10^9 \ \psi(2S)$  events
- $\gt$  20 fb<sup>-1</sup>  $\psi(3770)$
- Scan data [1.84, 3.08] GeV; [3.735, 4.600]GeV, 143
   energy points, ~2.0 fb<sup>-1</sup>
- $\triangleright$  Large data sets for XYZ study  $\sim$ 22 fb<sup>-1</sup>
- $\triangleright$  Entangled hadron pair-productions near thresholds  $^R$

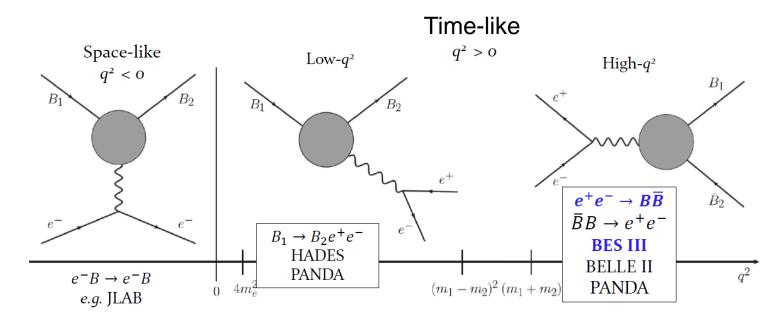
#### Totally about 50 fb<sup>-1</sup> from 2.0-4.95 GeV



#### **Rich physics program:**

Spectroscopy & decays of light hadrons and charmonium, charm physics, precision measurements of QCD parameters, tests of fundamental symmetry, .....

#### Electromagnetic Form Factors (EMFFs)



### Hadron structure with BESIII

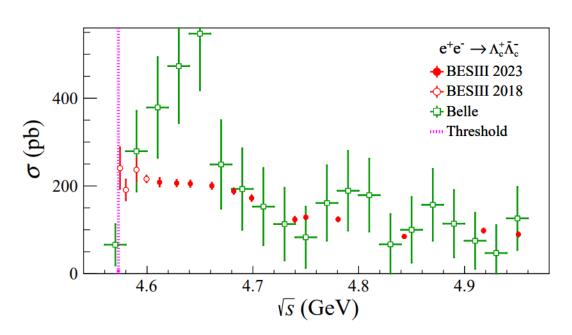
$$e^+e^- \rightarrow \Lambda_c^+ \overline{\Lambda_c}^-$$

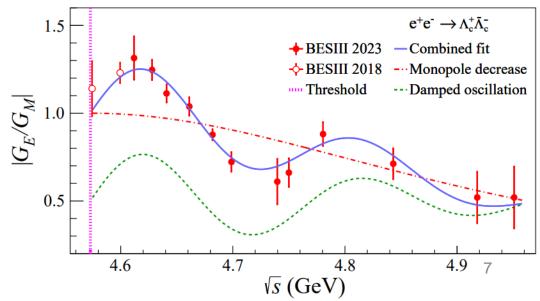
#### Energy scans from 4.61 to 4.95 GeV

- Sharp rise in cross section near threshold
- Disagreement with Belle data near 4.6 GeV
- No discernible oscillations of the effective form factors  $G_{\text{eff}}$

With the polar-angle distribution of  $\Lambda_c^+$ 

- |G<sub>E</sub>| and |G<sub>M</sub>| are extracted
- Energy dependence of R =  $|\frac{G_E}{G_M}|$ :
  - → Damped oscillations with frequency
  - ~3.5 times larger than for the proton

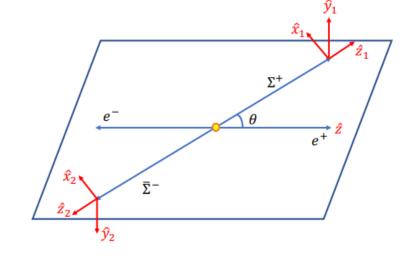




### Complete decomposition of $\Sigma^+$ EMFFs

• Using a fully differential angular description of the final state particles  $e^+e^- \to \Sigma^+ (\to p\pi^0)\bar{\Sigma}^- (\to \bar{p}\pi^0)$ , the relative magnitude and phase of  $\Sigma^+$  EMFFs can be extracted:

$$W(\xi) \propto \frac{\mathcal{F}_0(\xi) + \alpha \mathcal{F}_5(\xi)}{\mathcal{F}_0(\xi) + \alpha \mathcal{F}_5(\xi)}$$
 Unpolarized part 
$$+ \alpha_1 \alpha_2 (\mathcal{F}_1(\xi) + \sqrt{1 - \alpha^2} \cos(\Delta \Phi) \mathcal{F}_2(\xi) + \alpha \mathcal{F}_6(\xi))$$
 Correlated part 
$$+ \sqrt{1 - \alpha^2} \sin(\Delta \Phi) (-\alpha_1 \mathcal{F}_3(\xi) + \alpha_2 \mathcal{F}_4(\xi)),$$
 Polarized part



$$\begin{split} \mathcal{F}_0(\xi) &= 1 \\ \mathcal{F}_1(\xi) &= \sin^2\theta \sin\theta_1 \sin\theta_2 \cos\phi_1 \cos\phi_2 - \cos^2\theta \cos\theta_1 \cos\theta_2 \\ \mathcal{F}_2(\xi) &= \sin\theta \cos\theta (\sin\theta_1 \cos\theta_2 \cos\phi_1 - \cos\theta_1 \sin\theta_2 \cos\phi_2) \\ \mathcal{F}_3(\xi) &= \sin\theta \cos\theta \sin\theta_1 \sin\phi_1 \\ \mathcal{F}_4(\xi) &= \sin\theta \cos\theta \sin\theta_2 \sin\phi_2 \\ \mathcal{F}_5(\xi) &= \cos^2\theta \\ \mathcal{F}_6(\xi) &= \sin^2\theta \sin\theta_1 \sin\theta_2 \sin\phi_1 - \cos\theta_1 \cos\theta_2. \end{split}$$

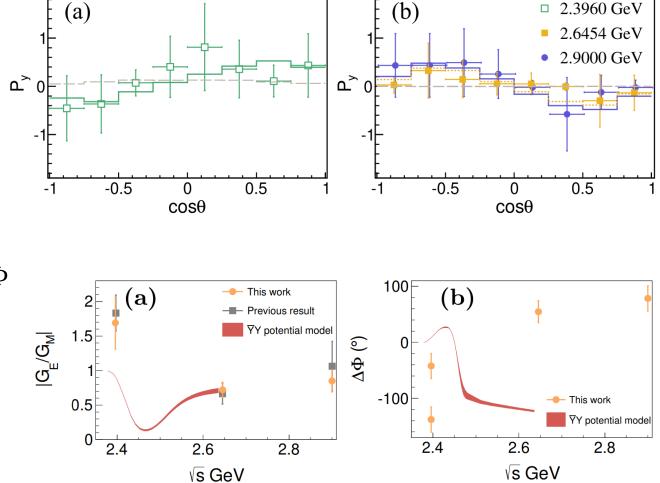
• A nonzero relative phase leads to polarization  $P_y$  of the out going baryons:

$$P_{y} = \frac{\sqrt{1 - \alpha^{2}} \sin\theta \cos\theta}{1 + \alpha \cos^{2}\theta} \sin(\Delta\Phi)$$

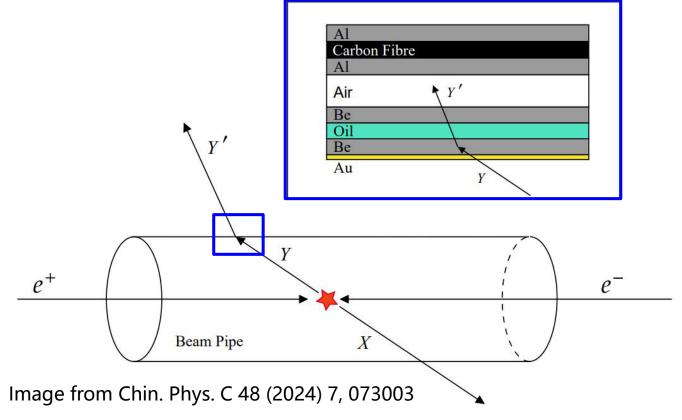
### Complete decomposition of $\Sigma^+$ EMFFs

BESIII PRL 132, 081904 (2024)

- Polarization is observed at  $\sqrt{s}$ =2.396, 2.644 and 2.90 GeV with a significance of 2.2 $\sigma$ , 3.6 $\sigma$  and 4.1 $\sigma$
- Relative phase is determined for the first time in a wide  $q^2$  range
  - $|G_E/G_M|$  and  $\Delta\Phi$  line-shape is compared with  $\overline{Y}Y$  model [PRD 103, 014028 (2021)], different tendency in  $\Delta\Phi$
  - ΔΦ evolution is an important input for understanding its asymptotic behavior and the dynamics of baryons



\*  $\Delta\Phi$  /  $180^{\circ}$  –  $\Delta\Phi$  ambiguity



# Hadron (YN) interactions with BESIII

- Crucial component to predict hypernuclei properties
- Key to understand the hyperon puzzle of neutron stars

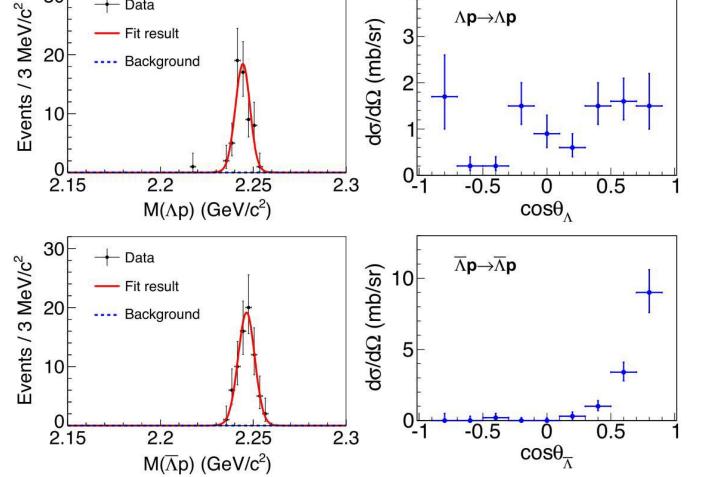
BESIII PRL 130, 251902 (2023) BESIII PRC 109, L052201 (2024) BESIII PRL 132, 231902 (2024)

### First study of antihyperon-nucleon scattering

BESIII PRL 132, 231902 (2024)

Beam:  $\Lambda/\overline{\Lambda}$  from  $J/\psi \to \Lambda\overline{\Lambda}$ , using 10B  $J/\psi$   $p_{\Lambda} = 1.074 \pm 0.017 \text{ GeV/c}^2$ ,  $\left|\cos\theta_{\Lambda(\overline{\Lambda})}\right| < 0.9$ 

Target: proton, the hydrogen nuclei in the cooling oil of the beam pipe



$$\sigma(\Lambda p \to \Lambda p) = (12.2 \pm 1.6(\text{stat.}) \pm 1.1(\text{syst.})) \text{ mb}$$
  
$$\sigma(\overline{\Lambda} p \to \overline{\Lambda} p) = (17.5 \pm 2.1(\text{stat.}) \pm 1.6(\text{syst.})) \text{ mb}$$

- Slight tendency of forward scattering for  $\Lambda p \to \Lambda p$
- Strong forward peak for  $\overline{\Lambda}p 
  ightarrow \overline{\Lambda}p$

# Atomic Spectrum: Bohr model → QED

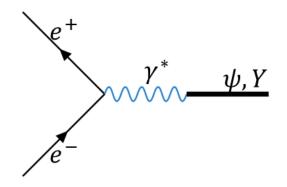


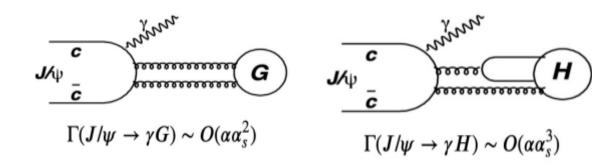
#### Hadron spectrum:

Quark model → QCD



# Hadron spectroscopy with BESIII





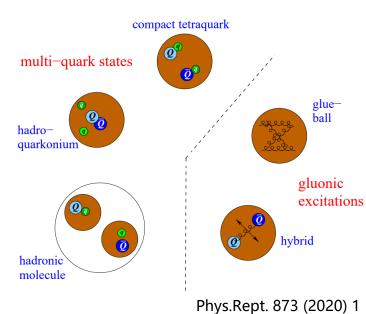
#### **QCD** exotics

## Hadron spectroscopy

- How does QCD give rise to hadrons?
  - Quark model seems to work really well. Why?
- Key to access the effective degree of freedom of QCD
  - Strong evidences for multi-quark in heavy quark sector



Evidence for gluonic excitations remains sparse



#### Manifestly exotic: with forbidden QN Physical meson

A linear superposition of all allowed color-singlet configurations

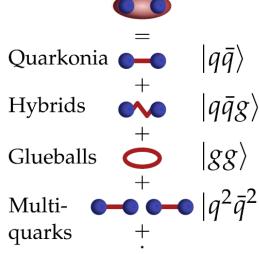
> **Identification of exotics** is challenging

Flavor exotic:  $Z_c$ ,  $T_{cc}$ ,  $T_{\psi\psi}$  ... ... Spin exotic:  $I^{PC} = 0^{--}$ ,  $even^{+-}$ ,  $odd^{-+}$ 

#### Crypto exotic: with QN as $q\bar{q}$

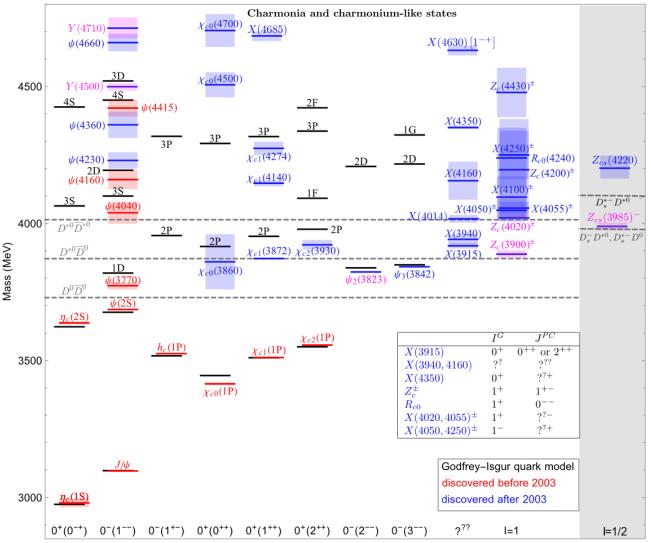
Supernumerary states - -> glueball Abnormal properties

+ Kinematic effects





### Charmonium-like states

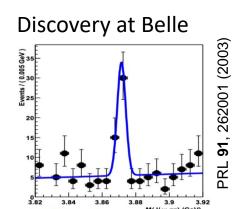


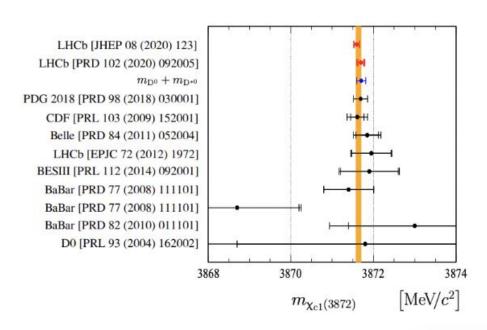
arXiv:2203.08290

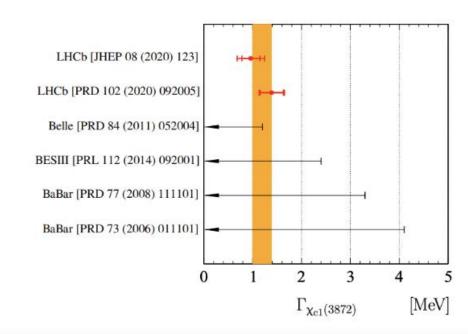
- Conventional  $c\overline{c}$  meson fit well with potential model
- Abundance of new states with various probes
  - *b*-hadron decays
  - hadron/heavy-ion collisions
  - γγ processes
  - $e^+e^-$  collisions
    - BESIII: dominant for vectors and states produced from vector decays

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### X(3872)







#### Many experiments contribute to it

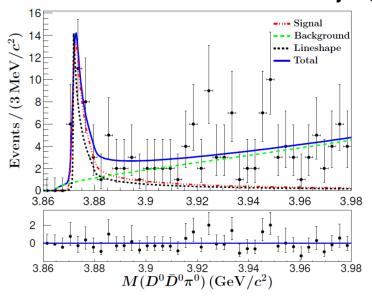
- Spin assignment: **J**<sup>PC</sup> = **1**<sup>++</sup>
- Mass is consistent with m(D<sup>0</sup>) + m(D<sup>\*0</sup>)
- Width is **surprisingly narrow**
- **Prompt production:**  $X(3872)-\psi(2S)$  yield ratio from p-p with increasing multiplicities toward p-Pb and Pb-Pb collisions
- Decay properties:  $\rightarrow \omega J/\psi$ ,  $\rho J/\psi$ ;  $\rightarrow \gamma J/\psi$ ,  $\gamma \psi(2S)$

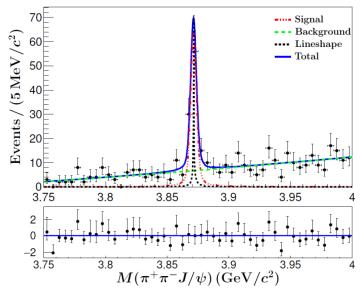
#### Its nature is still under debate!

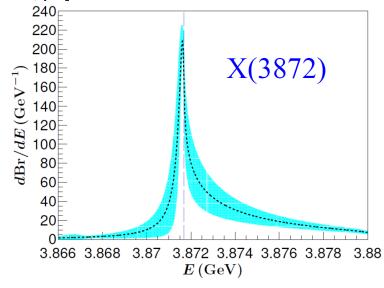
 $\rightarrow$  conventional  $\chi_{c1}(2^3P_1)$ , DD\* molecular state, tetraquark, hybrid, vector glueball, or mixed?

# X(3872) line shape @BESIII

 $e^+e^- o \gamma X(3872)$ ,  $X(3872) o D^0 \overline{D}{}^0 \pi^0$  and  $\pi^+\pi^- J/\psi$  BESIII PRL 132, 151903 (2024)







#### Pole positions

Two sheets with respect to  $D^{*0}\overline{D}{}^{0}$  branch cut

• Sheet I: 
$$E - E_X - g\sqrt{-2\mu(E - E_R + i\Gamma/2)}$$

• Sheet II: 
$$E - E_X + g\sqrt{-2\mu(E - E_R + i\Gamma/2)}$$

$$E_{\rm I} = (7.04 \pm 0.15^{+0.07}_{-0.08}) + (-0.19 \pm 0.08^{+0.14}_{-0.19})i \text{ MeV}$$

$$E_{\text{II}} = (0.26 \pm 5.74^{+5.14}_{-38.32}) + (-1.71 \pm 0.90^{+0.60}_{-1.96})i \text{ MeV}$$

	LHCb	Belle	BESIII	
g	$0.108 \pm 0.003^{+0.005}_{-0.006}$	$0.29^{+2.69}_{-0.15}$	$0.16 \pm 0.10^{+1.12}_{-0.11}$	
$Re[E_I]$ [MeV]	7.10	7.12	$7.04 \pm 0.15^{+0.07}_{-0.08}$	
$Im[E_I]$ [MeV]	-0.13	-0.12	$-0.19 \pm 0.08^{+0.14}_{-0.19}$	
$Re[k^+]$ [MeV]	-13.9	-15.3	$-12.6 \pm 5.5^{+6.6}_{-6.2}$	
$Im[k^+]$ [MeV]	8.8	7.7	$12.3 \pm 6.8^{+6.0}_{-6.4}$	
a (fm)	-27.1	-31.2	$-16.5^{+7.0}_{-27.6}{}^{+5.6}_{-27.7}$	
$r_e$ (fm)	-5.3	$-3.0^{+1.3}_{-1.5}$	$-4.1^{+0.9}_{-3.3}{}^{+2.8}_{-4.4}$	
$ar{Z}_A$	0.15 (0.33)	$0.08^{+0.04}_{-0.03}$	$0.18^{+0.06}_{-0.17}~^{+0.19}_{-0.16}$	

Weinberg's compositeness: Z=1: pure elementary state; Z=0: pure bound (composite) state

### Vector states: $Y(4260) \rightarrow Y(4230)$

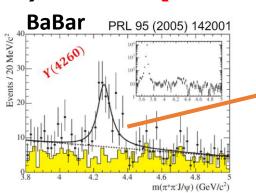
E<sub>CM</sub><sup>4.3</sup>(GeV)

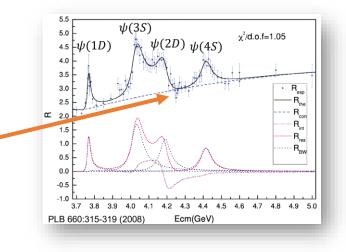
- Y(4260) firstly seen by BaBar
  - Inconsistent with simple  $c\overline{c}$  scenario
  - Candidates for exotics:

4.0 4.1

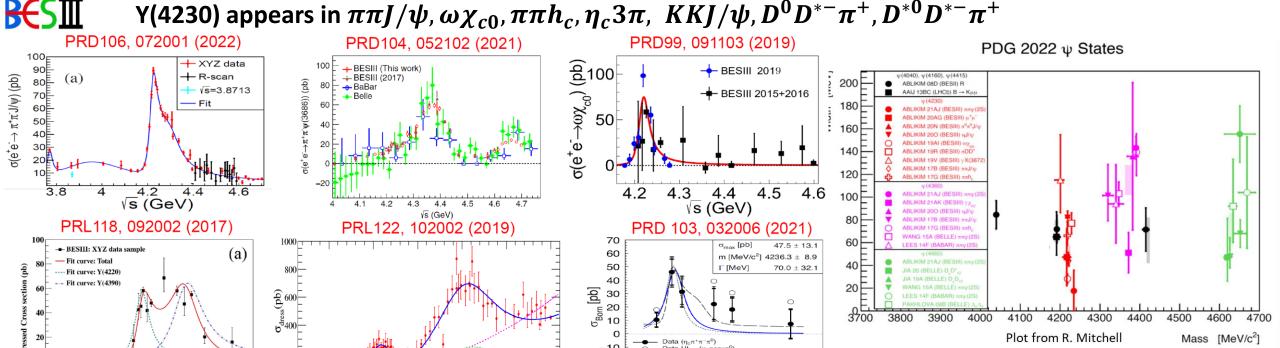
4.2 4.3 4.4

Hybrid /molecule /Tetraquark ?



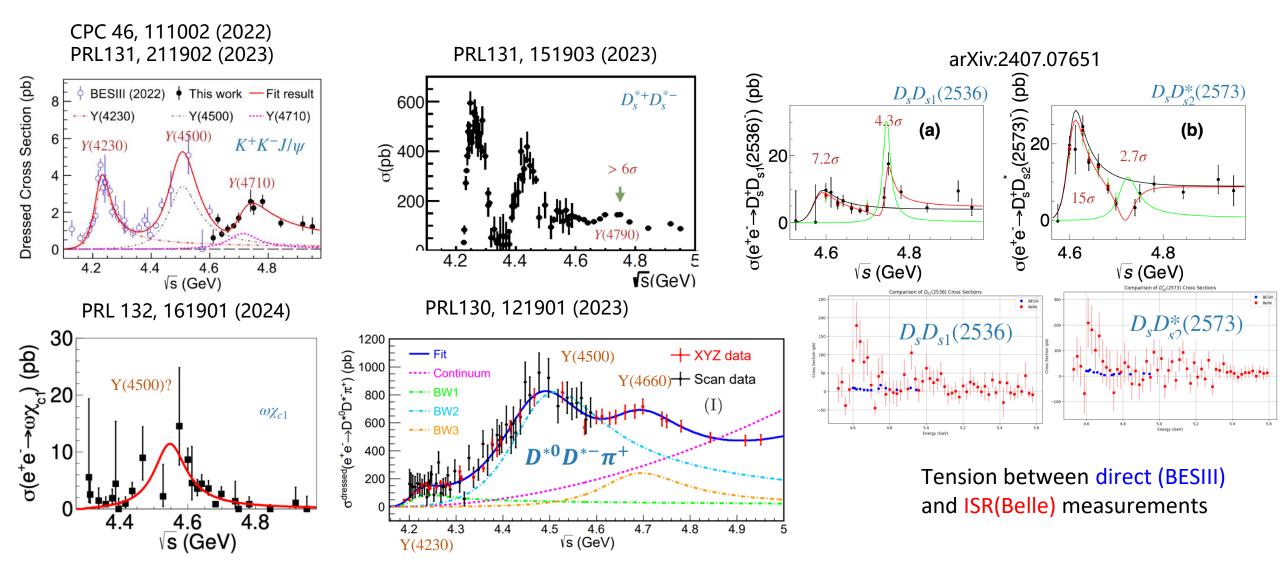


Mass~4220 MeV, width~ 50 MeV

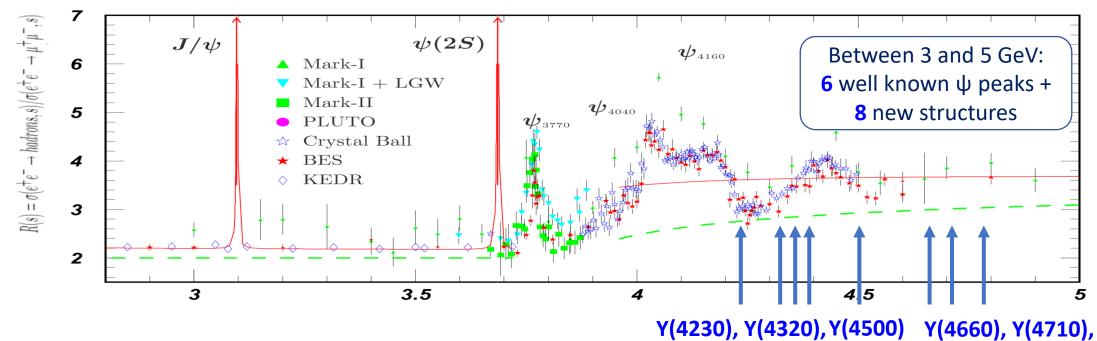


√s [GeV]

### Observations of Y(4500), Y(4710) and Y(4790)

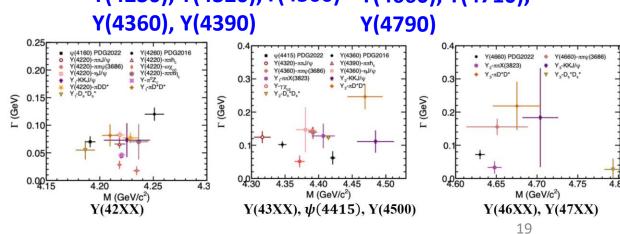


### How many vectors in charmonium energy region?



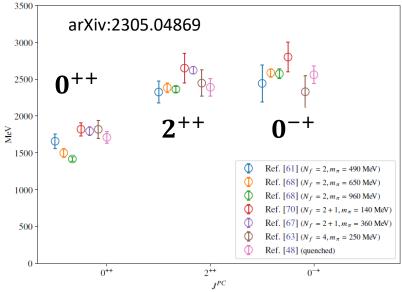
Besides  $c\overline{c}$  states, we also expect  $gc\overline{c}$  hybrids, and  $c\overline{c}q\overline{q}$  tetraquark states. Have they already been observed?

→ More theoretical/experimental efforts necessary

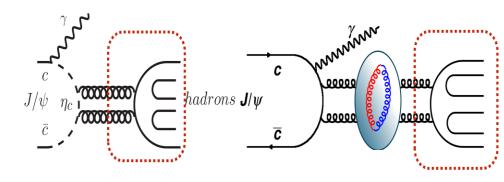


# Glueball hunting for over 40 years

- Glueballs: the most direct prediction of QCD
  - Gluon self-interactions
  - Can massless gluons form massive, exotic matter?
- Theoretical predictions from LQCD and QCDinspired models mostly consistent
- Supernumerary states that do not fit into  $q \overline{q}$  multiplets
- Production: Strongly produced in gluon-rich processes
- Decay: gluon is flavor-blind
  - No rigorous predictions
    - Could be analogy to OZI suppressed decays of charmonium, as they all decay via gluons [PLB 380 189(1996), Commu. Theor. Phys. 24.373(1995)]



Light Yang-Mills glueballs on lattice (quenched and unquenched results)



### Scalar glueball candidate

- Supernumerary scalars suggest additional degrees of freedom
  - However, mixing scenarios are controversial
- Measured B(J/ $\psi \rightarrow \gamma f_0(1710)$ ) is **x10 larger** than  $f_0(1500)$  BESIII [PRD 87 092009, PRD 92 052003, PRD 98 072003]
  - LQCD:  $\Gamma(J/\psi \to \gamma G_{0+})/\Gamma_{total} = 3.8(9) \times 10^{-3}$  [PRL 110, 091601(2013)]
  - **BESIII:**  $f_0(1710)$  largely overlays with the scalar glueball
  - ➤ Identification of scalar glueball with coupled-channel analyses based on BESIII data

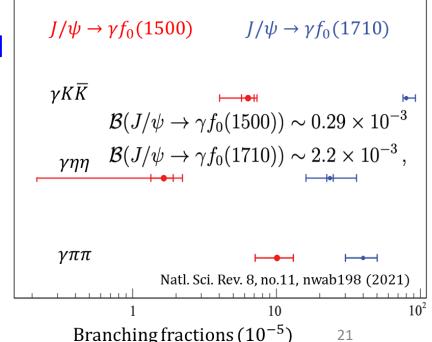
[PLB 816, 136227 (2021), EPJC 82, 80 (2022), PLB 826, 136906 (2022)]

• Further more, suppression of  $f_0(1710)\to\eta\eta'$  supports  $f_0(1710)$  has a large overlap with glueball

 $f_0(1370)$   $f_0(1500)$   $f_0(1710)$   $u\bar{u} + d\bar{d}$   $s\bar{s}$  ggCheng et al, Phys. Rev. D74 (2006) 094005  $f_0(1370)$   $f_0(1500)$   $f_0(1710)$ 

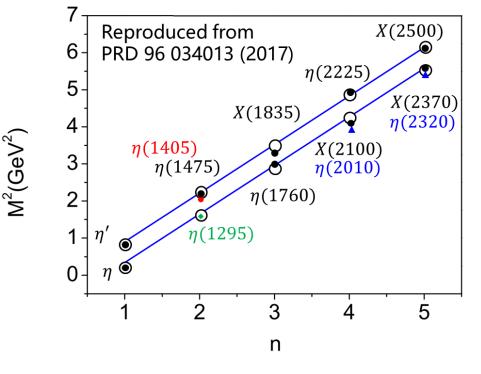
 $\blacksquare gg$ 

Close and Kirk, PLB483 (2000) 345



# Where is the 0<sup>-+</sup> glueball

- Pseudoscalar sector, a promising window
  - Only  $\eta$ ,  $\eta'$  (& radial excitations) from quark model
- Mass
  - LQCD: 0<sup>-+</sup> glueball (2.3~2.6 GeV)
  - The first glueball candidate:  $\iota(1440)$  (Split into  $\eta(1405)$  and  $\eta(1475)$ )
    - Mass incompatible with LQCD
  - Little experimental information above 2 GeV
- Production
  - LQCD:  $\Gamma(J/\psi \to \gamma G_{0-})/\Gamma_{total} = 2.31(80) \times 10^{-4}$ , at the same level as 0<sup>-+</sup> mesons [PRD.100.054511(2019)]
- Decays
  - Possible guidance: OZI suppressed decays of  $\eta_c$
  - 3 pseudoscalar final state is a good place to look for (0<sup>-+</sup> → 2P is forbidden)



#### $\eta_c \rightarrow 3 \text{ P in PDG}$

Decays involving hadronic resonances

 $(1.7 \pm 0.5)\%$ 

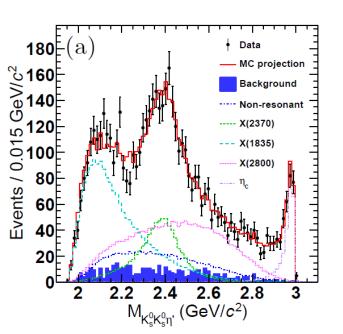
```
\eta'(958)\pi\pi ( 1.87\pm0.26) % ( 1.61\pm0.25) % 

\eta'(958)K\overline{K} Decays into stable hadrons ( 7.0\pm0.4) % ( 1.32\pm0.15) %
```

- No dominant decay
  - Flavor symmetric

# A glueball-like state X(2370)

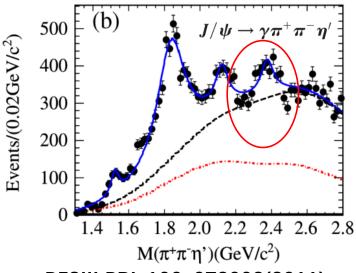
- Discovered by BESIII in  $J/\psi \rightarrow \gamma \eta' \pi \pi$  in 2011
- Confirmed by BESIII in  $J/\psi \rightarrow \gamma \eta' \pi \pi$ ,  $\gamma \eta' KK$ 
  - Not seen in J/ $\psi \to \gamma \eta' \eta \eta$  [BESIII PRD 103 012009 (2021)], J/ $\psi \to \gamma \gamma \varphi$  [BESIII arXiv: 2401.00918]. Upper limits of BF are well consistent with predictions of  $0^{-+}$  glueball
- Mass consistent with LQCD prediction for  $0^{-+}$  glueball
- Spin-parity determined to be  $0^{-+}$  BESIII PRL 132, 181901(2024)



 $J/\psi \to \gamma \eta' K^0_S K^0_S$ 

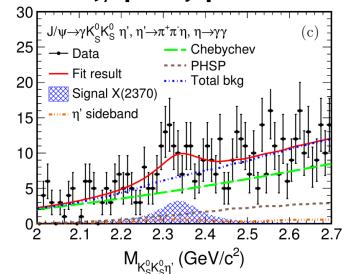
$$\begin{split} J^{pc} &= 0\text{-+ with significance} > 9.8\sigma \\ M &= 2395 \pm 11^{+26}\text{-}94 \text{ MeV} \\ \Gamma &= 188^{+18}\text{-}17^{+124}\text{-}33 \text{ MeV} \\ B(J/\psi \rightarrow \gamma X(2370))B(X(2370) \rightarrow f_0(980)\eta')B(f_0(980) \rightarrow \text{K}^0\text{s}\text{K}^0\text{s}) \\ &= 1.31 \pm 0.22^{+2.85}\text{-}0.84 \times 10^{-5} \end{split}$$

#### $J/\psi \rightarrow \gamma \eta' \pi \pi$



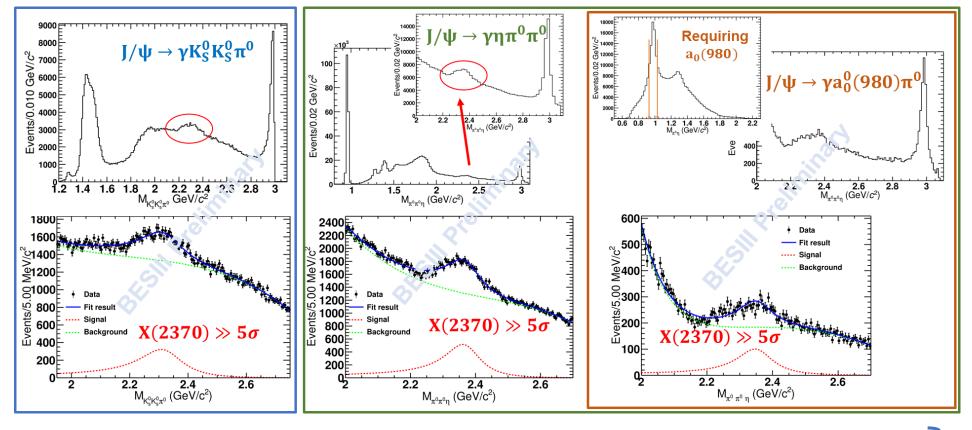
BESIII PRL 106, 072002(2011), PRL 117, 042002 (2016)





BESIII EPJC 80 746(2020)

Events/(0.01GeV/c<sup>2</sup>



# New decay modes

X(2370) observed in the gluon-rich  $J/\psi$  radiative decays

- J<sup>PC</sup> determined to be 0<sup>-+</sup>
- Mass and production rate consistent with LQCD
- Decay modes  $X(2370)\to \eta'\pi\pi, \eta'KK, K^0_SK^0_S\eta, K^0_SK^0_S\pi^0, \eta\pi^0\pi^0, a^0_0(980)\pi^0$  observed, in analog to  $\eta_c$

Consistent with  $\mathbf{0}^{-+}$  glueball

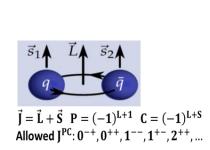
### Light hadrons with exotic quantum numbers

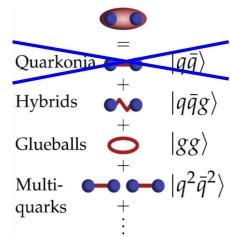
- Unambiguous signature: exotic quantum numbers forbidden for  $q\bar{q}$ :  $J^{PC}=0^{--}$ , even<sup>+-</sup>, odd<sup>-+</sup>
- Only 3 candidates over 30 yrs:

All 1<sup>-+</sup> isovectors 
$$\pi_1(1400), \pi_1(1600), \pi_1(2015)$$

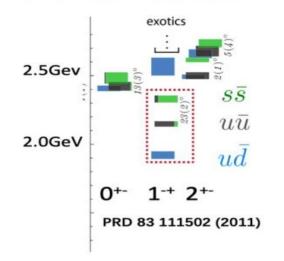
\*  $\pi_1(1400)$  and  $\pi_1(1600)$  can be explained as one resonance with recent coupled channel analyses

- Lightest spin-exotic state in LQCD: 1<sup>-+</sup> hybrid
- Isoscalar 1<sup>-+</sup> is critical to establish the nonet
  - Can be produced in the gluon-rich charmonium decays
  - Can decay to  $\eta \eta'$  in P-wave





#### **Lattice QCD Predictions:**



$$\pi_{1} I^{G}(J^{PC}) = \mathbf{1}^{-}(\mathbf{1}^{-+})$$

$$K_{1} I^{G}(J^{P}) = \frac{1}{2}^{-}(1^{-})$$

$$\eta_{1} I^{G}(J^{PC}) = \mathbf{0}^{+}(\mathbf{1}^{-+})$$

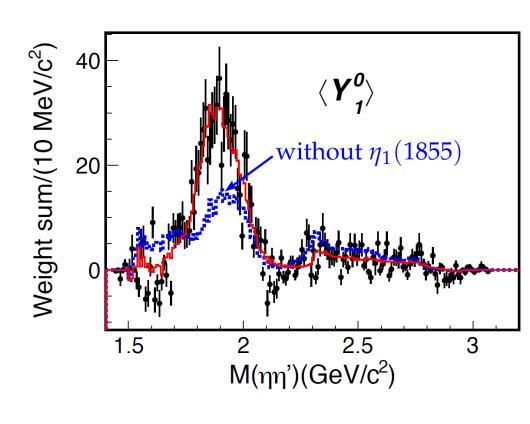
### Observation of An Exotic 1<sup>-+</sup> Isoscalar State $\eta_1(1855)$

PRL 129 192002(2022), PRD 106 072012(2022)

• An isoscalar 1<sup>-+</sup> ,  $\eta_1(1855)$ , has been observed in  $J/\psi \rightarrow \gamma \eta \eta'$  (>19 $\sigma$ )

$$\begin{split} M &= \left(1855 \pm 9^{+6}_{-1}\right) \text{MeV/c}^2, \ \Gamma = \left(188 \pm 18^{+3}_{-8}\right) \text{MeV/c}^2 \\ B(J/\psi \to \gamma \eta_1 (1855) \to \gamma \eta \eta') &= \left(2.70 \pm 0.41^{+0.16}_{-0.35}\right) \times 10^{-6} \end{split}$$

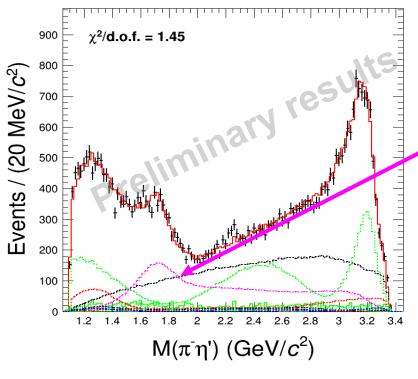
- Mass consistent with hybrid on LQCD
- Inspired many interpretations: Hybrid/ $K\overline{K}_1$ Molecule/Tetraquark?
- Opens a new direction to completing the picture of spin-exotics



"Here, the result by the BESIII experiment of a possible observation of an  $\eta_1(1855)$  state could be a breakthrough."

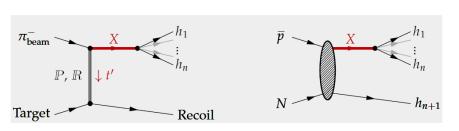
## Observation of $\pi_1(1600)$ in $\chi_{c_1} \to \eta' \pi^+ \pi^-$

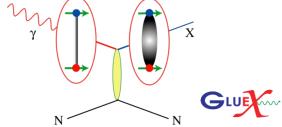
#### $2.7 \times 10^9 \psi(3686)$ @BESIII [preliminary]

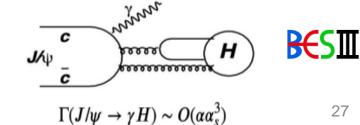


- Amplitude analysis of  $\chi_{c_1} \rightarrow \eta' \pi^+ \pi^-$  is performed
- $\pi_1(1600)$  observed> $10\sigma$
- with a significant BW phase motion
- I<sup>PC</sup>= 1<sup>-+</sup>, better than other assignments well over  $10\sigma$ 
  - Evidence of  $\pi_1 \to \eta' \pi$  at CLEO-c is confirmed [ PR D84 112009 (2011)]

#### Observations of $\pi_1$ and $\eta_1$ in charmonium decays provide a new path to study $1^{-+}$



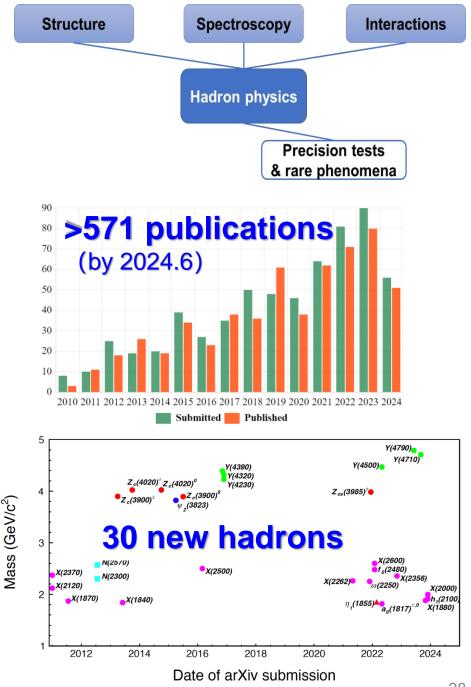




27

# Summary

- BESIII has a rich program of hadron physics
- Lots of progress
- Great potential to be fully explored
  - 50 fb  $^{-1}$  data on disk, including  $10 \times 10^9$  J/ $\psi$  and  $2.7 \times 10^9$   $\psi'$
  - Running until ~2030
  - Upgrade in this summer
    - $\mathcal{L} \times 3 @\sqrt{s} = 4.7 \text{ GeV}$
    - $\sqrt{s} \rightarrow 5.6 \ GeV$ , starting from 2028
    - CGEM inner tracker



# Thank you for your attention

### Indications of tensor glueball

$$\Gamma(J/\psi o\gamma G_{2^+})=1.01(22)keV$$
  $\Gamma(J/\psi o\gamma G_{2^+})/\Gamma_{tot}=1.1 imes10^{-2}$  CLQCD, Phys. Rev. Lett. 111, 091601 (2013)

#### Experimental results

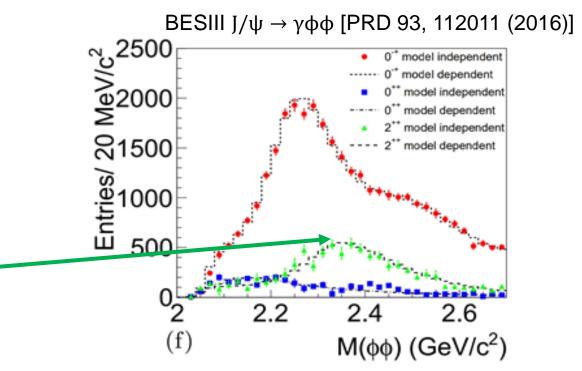
$$\begin{array}{c} Br(J/\psi \to \gamma f_2(2340) \to \gamma \eta \eta) = \left(3.8^{+0.62}_{-0.65}{}^{+2.37}_{-2.07}\right) \times 10^{-5} \\ \text{BESIII PRD 87,092009 (2013)} \end{array}$$

$$Br(J/\psi \to \gamma f_2(2340) \to \gamma \varphi \varphi) = \left(1.91 \pm 0.14^{+0.72}_{-0.73}\right) \times 10^{-4}$$
 BESIII PRD 93, 112011 (2016)

$$\begin{split} Br(J/\psi \to \gamma f_2(2340) \to \gamma K_s K_s) &= \left(5.\, 54^{+0.34+3.82}_{-0.40-1.49}\right) \times 10^{-5} \\ \text{BESIII PRD 98,072003 (2018)} \end{split}$$

$$Br(J/\psi \to \gamma f_2(2340) \to \gamma \eta' \eta') = \left(8.67 \pm 0.70^{+0.16}_{-1.67}\right) \times 10^{-6}$$
 BESIII PRD 105,072002 (2022)

still desired to study more decay modes

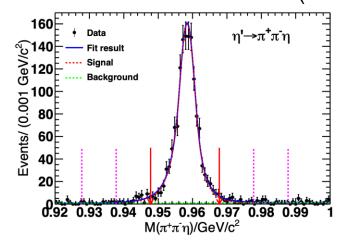


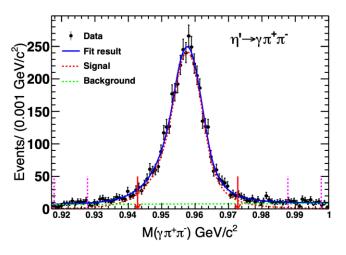
- f<sub>2</sub>(2010), f<sub>2</sub>(2300) and f<sub>2</sub>(2340) in πp reactions are all observed in J/ψ → γφφ with a strong production of f<sub>2</sub>(2340)
- Consistent with double-Pomeron exchange from WA102@CERN

## Spin-parity Determination of X(2370) in $J/\psi \rightarrow \gamma \eta' K_S^0 K_S^0$

- $\eta'$  reconstructed with  $\eta \pi^+ \pi^-$  and  $\gamma \pi^+ \pi^-$
- $K_S^0$  reconstructed with  $\pi^+\pi^-$
- Almost background free
  - Negligible mis-combination for  $K_S^0$  ( <0.1%)
  - No background from  $J/\psi \to \pi^0 \eta' K_S^0 K_S^0$  or  $\eta' K_S^0 K_S^0$ 
    - Forbidden by exchange symmetry and CP conservation
  - No peaking background
  - Little Non-  $\eta'$  backgrounds estimated from  $\eta'$  sidebands
    - 1.8% for  $\eta^\prime \to \eta \pi^+ \pi^-,$  6.8% for  $\eta^\prime \to \gamma \pi^+ \pi^-$

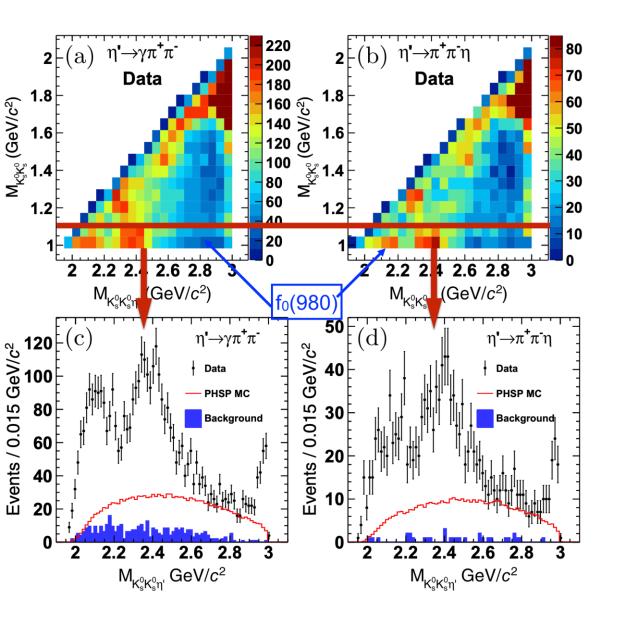
#### BESIII PRL 132 181901(2024)





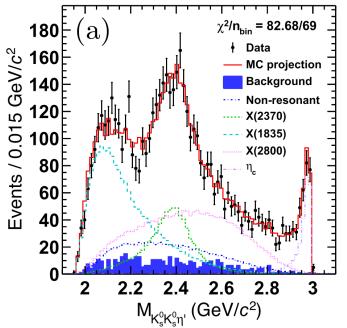
### Spin-parity Determination of X(2370) in $J/\psi \rightarrow \gamma \eta' K_S^0 K_S^0$

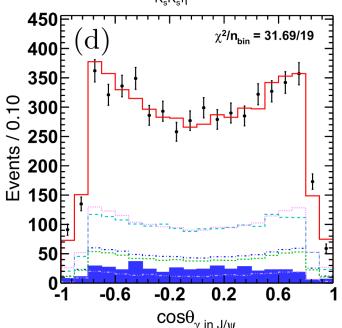
BESIII PRL 132 181901(2024)



- A clear connection between the  $f_0(980)$  and  $X(2370)/\eta_c$ 
  - $f_0(980)$  selection with  $M(K_S^0K_S^0) < 1.1$ GeV/ $c^2$
  - Clear signals of the X(2370) and  $\eta_c$
- Amplitude analysis
  - Quasi two-body decay amplitudes in the sequential decay processes J/ψ → γX, X → Yη', Y → K<sub>S</sub><sup>0</sup>K<sub>S</sub><sup>0</sup> and J/ψ → γX, X → ZK<sub>S</sub><sup>0</sup>, Z → K<sub>S</sub><sup>0</sup>η' are constructed using the covariant tensor formalism[Eur. Phys. J. A 16, 537]

## Spin-parity Determination of X(2370) in $J/\psi \rightarrow \gamma \eta' K_S^0 K_S^0$





BESIII PRL 132 181901(2024)

#### Nominal fit solution

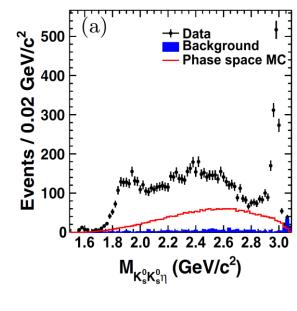
state	$J^{PC}$	Decay mode	Mass $(MeV/c^2)$	Width $(MeV/c^2)$	Significance
X(2370)	0-+	$f_0(980)\eta'$	$2395^{+11}_{-11}$	$188^{+18}_{-17}$	$14.9\sigma$
X(1835)	0-+	$f_0(980)\eta'$	1844	192	$22.0\sigma$
X(2800)	0-+	$f_0(980)\eta'$	2799 <sup>+52</sup> <sub>-48</sub>	660 <sup>+180</sup> <sub>-116</sub>	$16.4\sigma$
$\eta_c$	0-+	$f_0(980)\eta'$	2983.9	32.0	> 20.0 <i>\sigma</i>
PHSP	0-+	$\eta'(K_S^0K_S^0)_{S-wave}$			$9.0\sigma$
		$\eta'(K_S^0K_S^0)_{D-wave}$			$16.3\sigma$

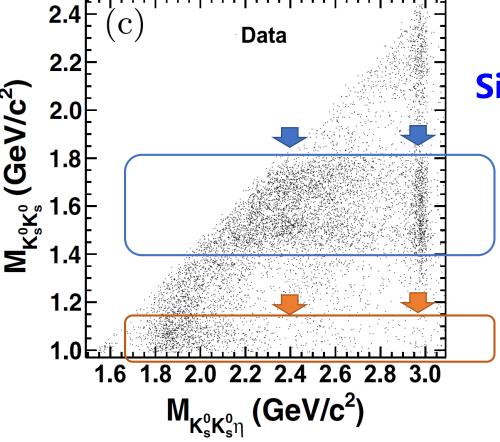
- $X(2370)'s J^{PC} = 0^{-+}$  with 9.8  $\sigma$
- Product branching fraction:

$$\begin{split} &B(J/\psi \to \gamma X(2370) B\big(X(2370) \to \eta' K_S^0 K_S^0\big) B\big(f_0(980) \to K_S^0 K_S^0\big) \\ &= \big(1.31 \pm 0.22^{+2.85}_{-0.84}\big) \times 10^{-5} \end{split}$$

# X(2370) seen in $J/\psi \rightarrow \gamma K_S^0 K_S^0 \eta$

Observation and Spin-Parity Determination of the X(1835) in  $J/\psi \rightarrow \gamma K_S^0 K_S^0 \eta$ BESIII PRL 115 091803(2015)



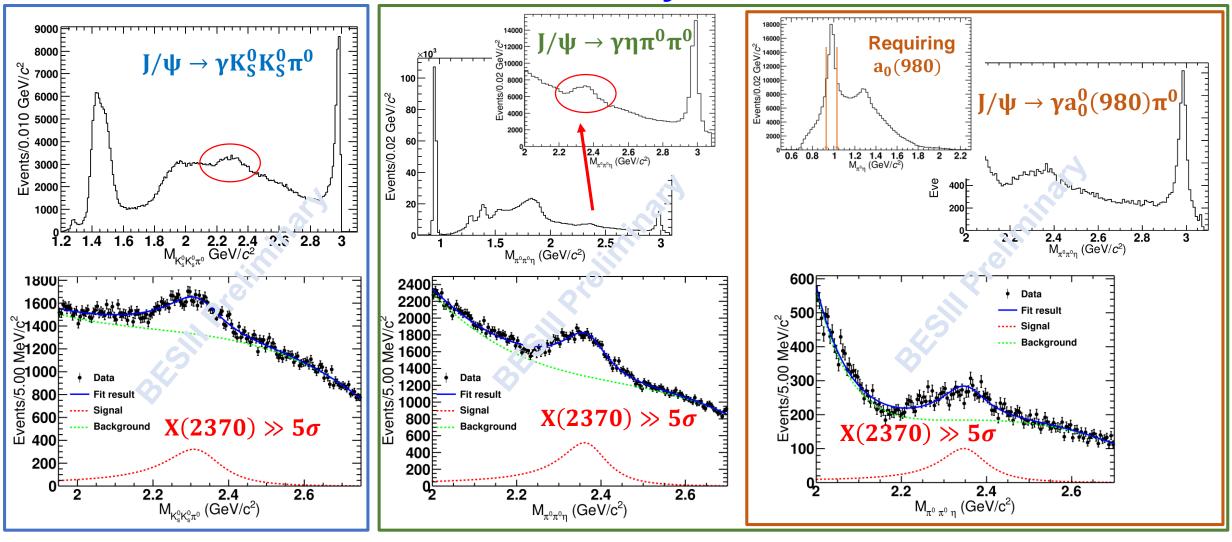


Similar decay patterns of the X(2370) and  $\eta_c$ 

clear X(2370) AND  $\eta_c$  signals

no X(2370) OR  $\eta_c$  signal

# Observation of new decay modes of X(2370)



•  $X(2370) \rightarrow K_S^0 K_S^0 \pi^0$ , \*  $\eta \pi^0 \pi^0$ ,  $a_0^0 (980) \pi^0$  firstly observed, all accompanied with  $\eta_c$ 

<sup>\*</sup>  $\eta(2320) \rightarrow \eta \eta \eta, \eta \pi \pi$  [PL B496 145(2000)] could be the current X(2370) at BESIII

### What we have learned before

-- from MarkIII, BES, Crystal barrel, OBELIX, WA102, GAMS, E852, ...

### Scalar: 1 nonet in quark model, f<sub>0</sub> & f<sub>0</sub>'

Exp: overpopulation

LQCD: ground state 0+ glueball ~1.7 GeV;

$$\Gamma(J/\psi \rightarrow \gamma G_{0+})/\Gamma_{total} = 3.8(9) \times 10^{-3}$$

#### Tensor: 2 nonets(3P2,3F2), complicated

Exp: large uncertainty

LQCD: 2++(2.3~2.4 GeV);

 $\Gamma(J/\psi \rightarrow \gamma G_{2+})/\Gamma_{total} = 1.1(2) \times 10^{-2}$ 

#### Pseudoscalar: η & η', "simple"

Exp: lacking of info. above 2 GeV; puzzles  $\eta(1295)$ ?

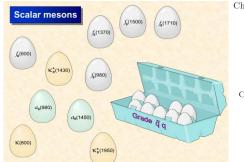
 $\eta(1405/1475)$ ?

LQCD:  $0^{-+}(2.3\sim2.6 \text{ GeV})$ 

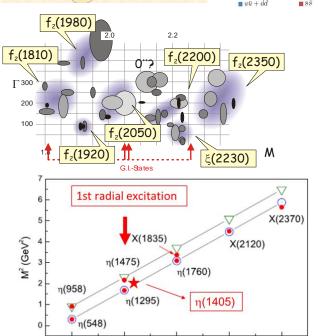
$$\Gamma(J/\psi \to \gamma G_{0-})/\Gamma_{total} = 2.31(80) \times 10^{-4}$$

#### e<sup>+</sup>e<sup>-</sup> annihilation

pp annihilation central exclusive production charge-exchange reactions







# Landscape of glueballs has been updated with BESIII's inputs

## Scalar: 1 nonet in quark model, f<sub>0</sub> & f<sub>0</sub>

Exp: overpopulation

LQCD : ground state 0+ glueball ~1.7 GeV;

$$\Gamma(J/\psi \rightarrow \gamma G_{0+})/\Gamma_{total} = 3.8(9) \times 10^{-3}$$

### Tensor: 2 nonets(3P2,3F2), complicated

Exp: large uncertainty

LQCD:  $2^{++}(2.3\sim2.4 \text{ GeV})$ ;  $\Gamma(J/\psi \to \gamma G_{2+})/\Gamma_{total} = 1.1(2) \times 10^{-2}$ 

#### Pseudoscalar: η & η', "simple"

Exp: lacking of info. above 2 GeV; puzzles  $\eta(1295)$ ?  $\eta(1405/1475)$ ?

LQCD:  $0^{-+}(2.3 \sim 2.6 \text{ GeV})$  $\Gamma(J/\psi \rightarrow \gamma G_{0-})/\Gamma_{total} = 2.31(80) \times 10^{-4}$ 

- √f<sub>0</sub>(1710) is largely overlapped with the scalar glueball, according to its production and decay properties
- **✓ Large production rate of** f<sub>2</sub>(2340) in J/ψ radiative decays
- **✓ Non-observation of**  $\eta(1295)$
- ✓ Insights of  $\eta(1405/1475)$
- ✓X(2370): a good candidate with analogy decay pattern as  $η_c$

## Scalar glueball candidate: decay properties

#### Flavor-blindness of glueball decays

$$\frac{1}{P.S.}\Gamma(G \to \pi\pi: K\overline{K}: \eta\eta: \eta\eta': \eta'\eta') = 3:4:1:0:1$$

\*with chiral suppression

PRL 95 172001, PRL 98 149103

Expectation:

$$\Gamma(G \to \pi\pi)/\Gamma(G \to K\bar{K}) \approx \frac{f_{\pi}^4}{f_K^4} \approx 0.48$$

Measured:



$$\frac{1}{P \cdot S} \Gamma(G \to \pi \pi: K\bar{K}: \eta \eta) \approx \underline{1.3: 3.16: 1}$$

New inputs from  $J/\psi \rightarrow \gamma \eta \eta'$ 

[BESIII PRL 129 192002(2022), PRD 106 072012(2022]

consistent with PDG

• Significant  $f_0(1500)$ 

$$\frac{B(f_0(1500) \to \eta \eta')}{B(f_0(1500) \to \pi \pi)} = (1.66^{+0.42}_{-0.40}) \times 10^{-1}$$

• Absence of  $f_0(1710)$ 

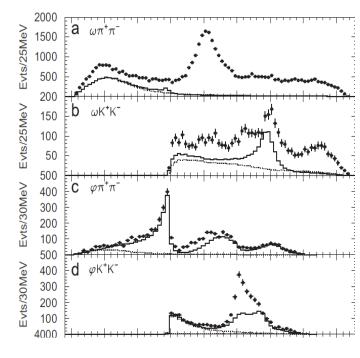
$$\frac{B(f_0(1710) \to \eta \eta')}{B(f_0(1710) \to \pi \pi)} < 2.87 \times 10^{-3} @90\% \text{ C. L.}$$

- Supports to the hypothesis that  $f_0(1710)$  overlaps with the ground state scalar glueball
  - Scalar glueball expected to be suppressed  $B(G \to \eta \eta')/B(G \to \pi \pi) < 0.04$

[PR D 92, 121902; PR D 92, 114035]

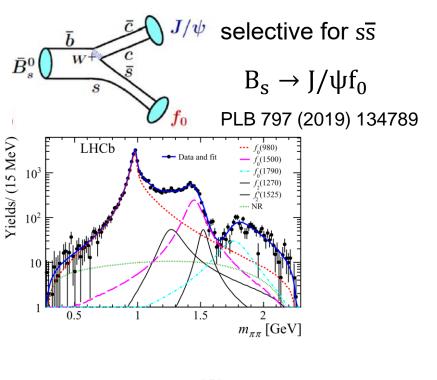
## More scalars

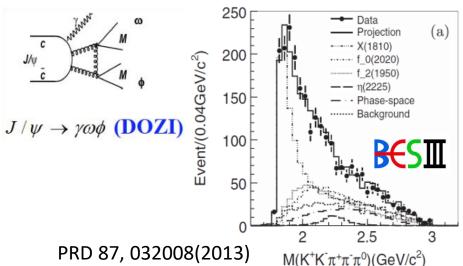
 $f_0(1710)/f_0(1790)$ ?



- ωK+K<sup>-</sup> → Peak around 1700 MeV/c²
   (OZI rule: n̄n structure)
- $\phi \pi^+ \pi^-$  Enhancement at 1790 MeV/ $c^2$
- $\phi K^+ K^ \rightarrow$  No peak around 1700 MeV/ $c^2$

 $f_0(1800)$ 

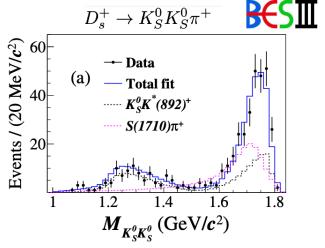




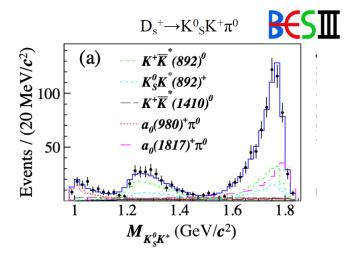
 $a_0(1817)$ 

Isovector partner of  $f_0(1800)$ ?

[Shulei 's talk] PRD105, L051103 (2022)



PRL129, 182001 (2022)



# Two photon couplings

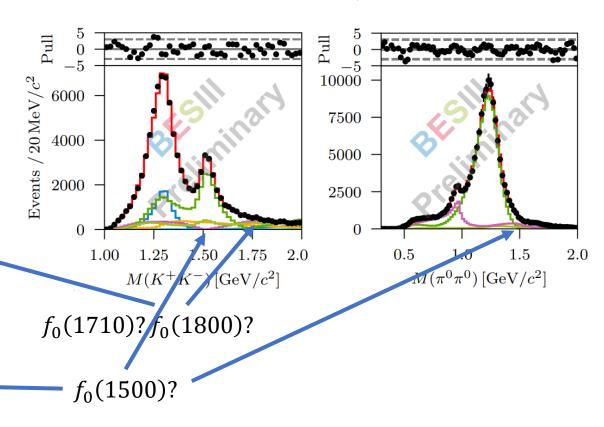
 $\gamma\gamma \rightarrow K_SK_S$ Belle PTEP 2013 (2013) 12, 123C01

Parameter	$f_0(1710) \text{ fit}$				
	fit-H	fit-L	$_{\mathrm{H,L}}$ combined	PDG	
$\chi^2/ndf$	694.2/585	701.6/585	_	_	
$\overline{\mathrm{Mass}(f_J) \; (\mathrm{MeV}/c^2)}$	$1750^{+5+29}_{-6-18}$	$1749^{+5+31}_{-6-42}$	$1750^{+6+29}_{-7-18}$	$1720 \pm 6$	
$\Gamma_{\mathrm{tot}}(f_J) \; (\mathrm{MeV})$	$138^{+12}_{-11}^{+96}$	$145^{+11+31}_{-10-54}$	$139^{+11+96}_{-12-50}$	$135 \pm 6$	
$\Gamma_{\gamma\gamma}\mathcal{B}(K\bar{K})_{f_J} \text{ (eV)}$	$12^{+3+227}_{-2-8}$	$21^{+6}_{-4-26}$	$12^{+3+227}_{-2-8}$	unknown	

$$\gamma\gamma \rightarrow \pi^0\pi^0$$
 Belle PRD 78 (2008) 052004

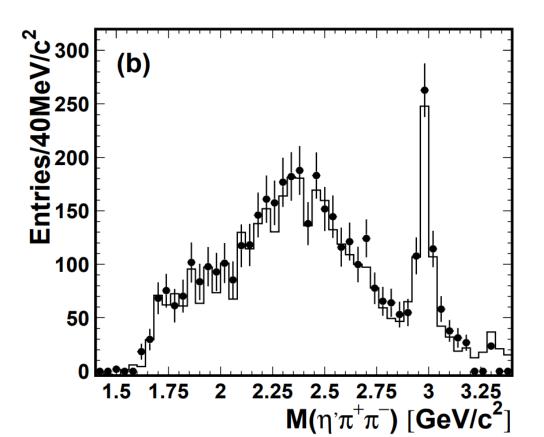
Parameter	Nominal	$r_{02} = 0$	No $f_0(Y)$	Unit
$Mass(f_0(980))$	$982.2 \pm 1.0$	$980.2 \pm 1.0$	$983.7^{+1.5}_{-1.0}$	$MeV/c^2$
$\Gamma_{\gamma\gamma}(f_0(980))$	$285.5^{+17.2}_{-17.1}$	$297.0^{+14.2}_{-13.7}$	$370.5^{+20.2}_{-18.7}$	$\mathrm{eV}$
$g_{f_0(980)\pi\pi}$	$1.82 \pm 0.03$	$1.79 \pm 0.03$	$1.89 \pm 0.03$	${ m GeV}$
$\operatorname{Mass}(f_0(Y))$	$1469.7 \pm 4.7$	$1466.8 \pm 0.6$	_	$MeV/c^2$
$\Gamma(f_0(Y))$	$89.7^{+8.1}_{-6.6}$	$422.4^{+18.4}_{-19.8}$	_	MeV
$\Gamma_{\gamma\gamma}\mathcal{B}(f_0(Y)\to\pi^0\pi^0)$	$11.2^{+5.0}_{-4.0}$	$6780.2^{+626.5}_{-574.7}$	0 (fixed)	eV

#### **BESIII** preliminary



• 
$$\gamma\gamma \to \eta'\pi^+\pi^-$$

#### Belle PRD 86 052002(2012)



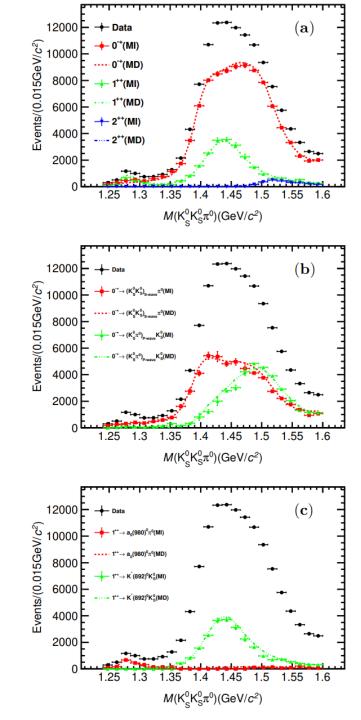
## Shed new lights on the $\eta(1405)/\eta(1475)$ puzzle

 $J/\psi \rightarrow \gamma K_S K_S \pi^0$ 

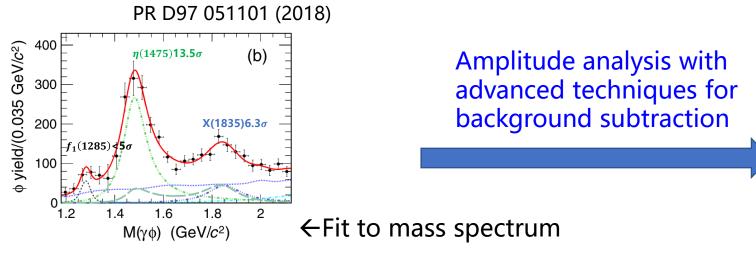
**BESIII JHEP 03 121(2023)** 

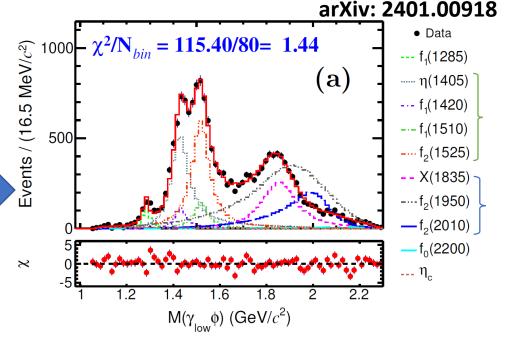
- Mass Independent PWA in bins of  $M(K_SK_S\pi^0)$  to detangle  $J^{PC}$  components
  - Valuable inputs to develop models
- Mass Dependent PWA with BW to extract resonances
- Consistency between MI and MD results
- Dominated by 0<sup>-+</sup>
  - Two BWs around 1.4 GeV is needed

- $\eta(1405)/\eta(1475)$  poles in coupled-channel analysis
  - PRD 107, L091505 (2023); PRD 109, 014021 (2024)



# $J/\psi \rightarrow \gamma \gamma \phi$ , a s̄s flavor filter





#### From the amplitude analysis,

- $\eta(1405)$  is observed, while  $\eta(1475)$  can not be excluded
- $X(1835) \rightarrow \gamma \phi$  suggests its assignment of  $\eta'$  excitation
- $\eta_c o \gamma \phi$  are observed. The very first radiative decay mode of  $\eta_c$
- Observation of  $f_2(1950)$  and  $f_0(2200) \rightarrow \gamma \phi$  unfavored their glueball interpretations[PRD 108, 014023, arXiv: 2404.01564]
- No evidence of  $X(2370)/\eta_1(1855)$ , well consistent with the predictions for glueball/hybrid [PRD 107, 114020, NPA 1037, 122683]

# Amplitude analysis

Amplitude analysis is a key tool of hadron spectroscopy to disentangle contributions from individual resonances and to extract the resonance's spin-parity, mass, width and decay properties

$$Prob(\xi:\alpha) = \frac{\omega(\xi,\alpha)\epsilon(\xi)}{\int d\xi\omega(\xi,\alpha)\epsilon(\xi)} \qquad \xi \text{ (the four-momenta of the final-state particles),} \\ \omega(\xi,\alpha) = \frac{d\sigma}{d\Phi} = |\sum_i A_i|^2 \text{ differential cross section,} \\ \ln L = \sum_{i=1}^{N_{data}} \ln(Prob(\xi,\alpha)) \qquad \epsilon(\xi) \text{ efficiency}$$

For  $J/\psi$  radiative decays

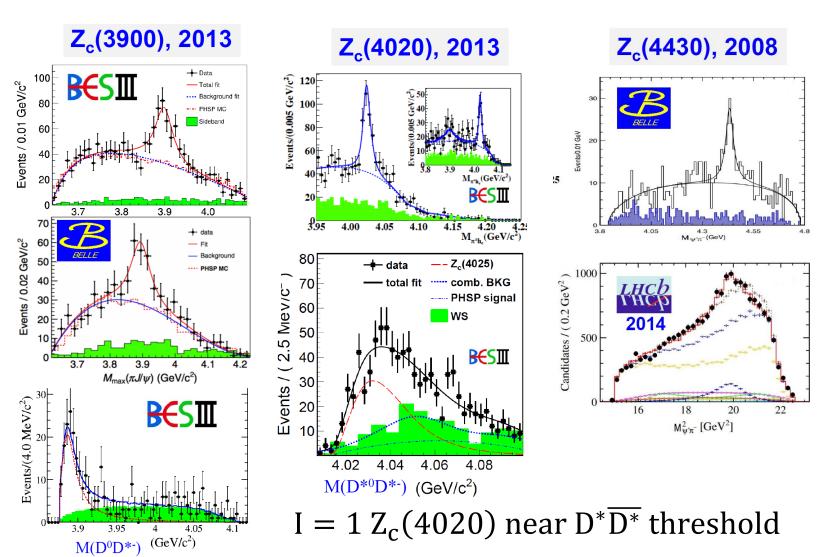
[Eur. Phys. J. A 16, 537]

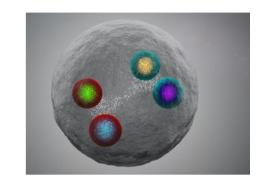
$$\begin{split} A &= \psi_{\mu}(m_1) e_{\nu}^*(m_2) A^{\mu\nu} = \psi_{\mu}(m_1) e_{\nu}^*(m_2) \sum_i \varLambda_i U_i^{\mu\nu} \\ \text{e.g. J/$\psi$} &\to \gamma 0^{-+}, 0^{-+} \to f_0 \eta, f_0 \pi \pi \\ &\langle \gamma 0^{-+} | (f_0 \eta) 1 \rangle = S_{\mu\nu} B_1(Q_{\psi\gamma X}) f_{(12)}^{(f_0)} \\ &S_{\mu\nu} = \epsilon_{\mu\nu\alpha\beta} p_{\psi}^{\alpha} q^{\beta} \end{split}$$

 $B_1(Q_{\psi \nu X})$  is Blatt-Weisskopf centrifugal barrier for  $J/\psi \rightarrow \gamma X$ 

Perform an un-binned loglikelihood fit (fit the data event-wise to high-dimensional distributions using complex weights) to make our model for  $\omega$  agree with the experimental distribution by varying the  $\alpha$ 

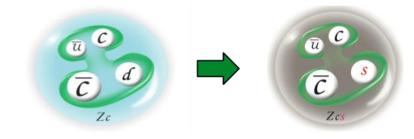
# $Z_c$ states





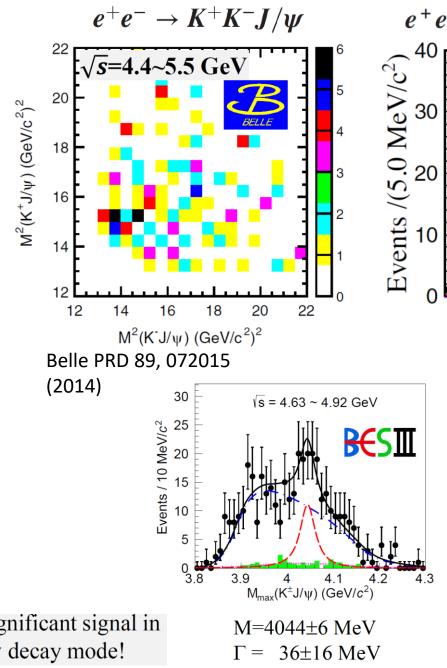
All are observed in  $\pi$ +charmonium  $c\overline{c}u\overline{d}$ 

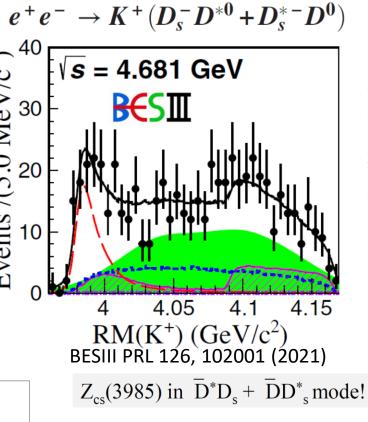
Existence of states with  $d \rightarrow s$ ?

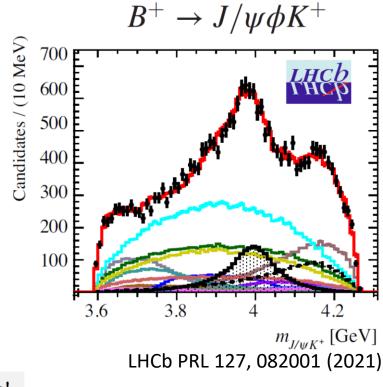


Search for states decay into  $KJ/\psi$ ,  $D_s^*\overline{D}$ ,  $D_s$ 

 $I = 1 Z_c(3900)$  near  $D\overline{D^*}$  threshold







 $Z_{cs}(4000)$  and  $Z_{cs}(4220)$ in  $K^{\pm}J/\psi$  decay mode!

State	Signif.	JP	Mass (MeV)	Width (MeV)
$Z_{cs}(3985)$	5.3σ	??	$3982.5^{+1.8}_{-2.6} \pm 2.1$	$12.8^{+5.3}_{-4.4} \pm 3.0$
$Z_{cs}(4000)$	15σ	1+	$4003 \pm 6^{+4}_{-14}$	131±15±26
Z. (4220)	5.90	1+	$4216 \pm 24^{+43}$	$233 \pm 52^{+97}$

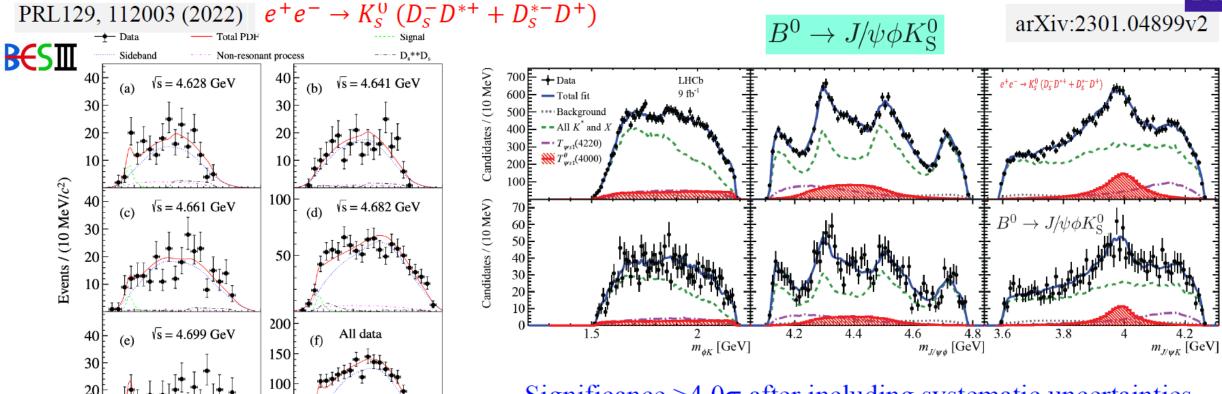
Same state?

No significant signal in  $K^{\pm}J/\psi$  decay mode!

Significance: 2.3σ

BESIII PRL131, 211902 (2023)

# Evidence for the neutral $Z_{cs}(3985)$ , $Z_{cs}(4000)$



Significance >4.0 $\sigma$  after including systematic uncertainties Significance 5.4 $\sigma$  with isospin symmetry imposed

State	Mass (MeV/ $c^2$ )	Width (MeV)	Significance
$Z_{cs}(3985)^+$	$3985.2^{+2.1}_{-2.0} \pm 1.7$	$13.8^{+8.1}_{-5.2} \pm 4.9$	$5.3\sigma$
$Z_{cs}(3985)^0$	$3992.2 \pm 1.7 \pm 1.6$	$7.7^{+4.1}_{-3.8} \pm 4.3$	$4.6\sigma$

4.05 4.1 4.15 4.2 3.95 4 4.05 4.1 4.15 4.2 RM(K<sub>c</sub>)(GeV/c<sup>2</sup>)

Mass (MeV)	${\rm Width}~({\rm MeV})$	Fit fraction (%)	$\Delta M \; ({ m MeV})$
$3991^{+12}_{-10}{}^{+9}_{-17}$	$105^{+29}_{-25}^{+17}_{-23}$	$7.9 \pm 2.5  {}^{+3.0}_{-2.8}$	$-12{}^{+11}_{-10}{}^{+6}_{-4}$

$$I = \frac{1}{2} Z_{cs}(3985)$$

## Time-like EMFFs: theoretic review

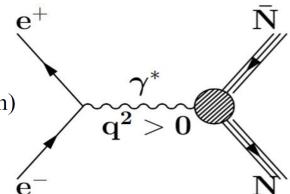
#### **1961**, first paper by N. Cabibbo and R. Gatto *Phys. Rev.* 124 (1961) 1577-1595

• The production cross section of  $e^+e^- \to B\overline{B}$  (1/2 baryon) is given:

$$\frac{\mathrm{d}\sigma_{B\overline{B}}}{\mathrm{d}\cos\theta} = \frac{\pi\alpha^2C\beta}{2q^2} \left[ \left(1 + \cos^2\theta\right) |G_M|^2 + \frac{1}{\tau} |G_E|^2 \sin^2\theta \right], \tau = \frac{q^2}{4m_B^2}$$

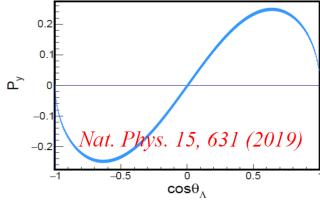
Integrated version:  $\sigma_{B\bar{B}} = \frac{4\pi\alpha^2 C\beta}{3q^2} \left[ |G_M|^2 + \frac{1}{2\tau} |G_E|^2 \right]$  (Born cross section)

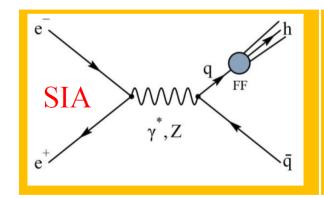
$$\xrightarrow{|G_E|=|G_M|} \sigma_{B\bar{B}} = \frac{2\pi\alpha^2 C\beta}{q^2} |G_{\text{eff}}|^2$$



• The complex feature of TLFF leads to transversely polarized baryon even the beams are unpolarized. Nuov Cim A 109, 241–256 (1996)

$$P_{y} = -\frac{\sin 2\theta \operatorname{Im}[G_{E}G_{M}^{*}]/\sqrt{\tau}}{\frac{|G_{E}|^{2}\sin^{2}\theta}{\tau} + |G_{M}|^{2}(1 + \cos^{2}\theta)}$$





$$e^+e^-$$
:  $\sigma = \sum_q \sigma(e^+e^- \to q\bar{q}) \otimes FF$ 

- No PDFs necessary
- Calculations know at NNLO
- Flavor structure not directly accessible

#### Two types of fragmentation functions can be studied at BEPCII/BESIII

- Unpolarized fragmentation function
  - Unique Q<10 GeV data
  - More results from charged  $\pi/K$  and heavy flavor

Normalized differential cross section

$$\frac{1}{\sigma_{\text{had}}} \frac{d\sigma_{\pi^0}}{dp_{\pi^0}} = \frac{N_{\pi^0}}{N_{\text{had}}} \frac{1}{\Delta p_{\pi^0}}$$

 $\pi^0/K_S$ , PRL 130 231901(2023)  $\eta$ , arXiv:2401.17873

Collins fragmentation function

 $\pi\pi$ , PRL 116, 042001 (2016)

- Essential input in the 3D imaging era of the nucleon structure study
- More results from  $K\pi + X$  and KK + X