Deuteron Tensor Structure

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Deuteron Polarization

$$
E_m = -h\nu_D m + h\nu_Q(\cos^2\theta - 1)(3m^2 - 2)
$$

Deuteron NMR Line-shap

Keller, D. Eur. Phys. J. A53 (2017)

DNP Technique

Requirements

- High magnetic field (at Jlab typically 5T)
- Low temperature $(\sim 1K)$
- Microwaves (induce spin transitions)
- CW NMR
- Irradiated material $ND₃$

Dynamic Nuclear Polarization: technique used to enhance vector polarization

DNP enhancement carries tensor polarization enhancement.

 \star Typical average vector polarization in Jefferson lab $P \sim 45\%$ which corresponds to $Q \sim 16\%$

Enhancing tensor polarization

- Low tensor polarization has limited physics experiments.
- New target developments are ongoing, with an enhancement of up to 30%.
- Two experiments to measure tensor observables have been approved.
- Several new experiments are underway.

Quasi-elastic Measurements

Scattering from an unpolarized deuteron target $p_{unp}(p_m) = u(p_m)^2 + w(p_m)^2$ $u(p_m)$: S-partial wave of the deuteron $w(p_m)$: D-partial wave of the deuteron

Scattering from tensor polarized target

$$
\rho_{20}(p_m, \theta_N) = \frac{2\cos^2(\theta_N) - 1}{2} \left[2\sqrt{2}u(p_m)w(p_m) - w^2(p_m) \right]
$$

 θ_N : direction of internal momenta with respect to the polarization axis of the deuteron

Measured asymmetry

$$
A_d^T = A_{20}(p_m, \theta_N) = \frac{\rho_{20}(p_m, \theta_N)}{\rho_{unp}(p_m)}
$$

Experimental Measurement

: target vector polarization Q: target tensor polarization

 A_D^V : vector analyzing power A_D^T : vector analyzing power

$$
\sigma_{pol}(P,Q) = \sigma_{unpol}\left[1 + PA_D^V + \frac{1}{2}QA_D^T\right]
$$

$$
\sigma_{pol}(-P,Q) = \sigma_{unpol}\left[1 - PA_z + \frac{1}{2}QA_{zz}\right]
$$

$$
\sigma_{pol}(P,0) = \sigma_{unpol}[1 + PA_z]
$$

$$
\sigma_{pol}(-P,0) = \sigma_{unpol}[1 - PA_z]
$$

$$
A_d^T = \frac{2}{Q} \left(\frac{\sigma_{pol}(P,Q) + \sigma_{pol}(-P,Q)}{\sigma_{pol}(P,0) + \sigma_{pol}(-P,0)} - 1 \right)
$$

New Ideas: Exclusive measurement

$$
\rho_{node}(p_m) = \rho_{unp}(p_m) + \frac{2\rho_{20}}{3\cos^2(\theta_N) - 1} = u^2(p_m) + 2\sqrt{2}u(p_m)w(p_m)
$$

We could measure
\n
$$
A_{node}(p_m) \equiv \frac{\rho_{node}}{\rho_{unp}} = 1 + \frac{2A_{20}(p_m, \theta_N)}{3\cos^2(\theta_N) - 1} = \frac{u^2(p_m) + 2\sqrt{2}u(p_m)w(p_m)}{u(p_m)^2 + w(p_m)^2}
$$
\n
$$
A_{node} = \begin{cases} u(p_m) = -2\sqrt{2}w(p_m) & \implies p_m \sim 180MeV = 0\\ u(p_m) = 0 & \implies p_m > 400MeV \end{cases}
$$

Most direct measurement of the repulsive strength of the nuclear core ever done in electro-nuclear processes.

•has large sensitivity to the position of these nodes

•has sensitivity to the choice of the potential used in calculating the deuteron wave function at $p_m > 300 \text{ MeV}$

•measures the contribution of the S -partial wave of the deuteron with respect to the missing momentum

Our observable

PWIA vs FSI

 θ_{na} : angle between the virtual photon and the recoiling neutron

Possible Experimental setup at Hall C

Semi-Inclusive Deep Inelastic Scattering

Leading twist distribution functions

I

After integrating over the transverse momentum:

Phys. Rev. D 62 (2000)

Leading twist distribution functions

After integrating over the transverse momentum:

•Only b_1 has been measured by H Phys.Rev.Lett. 95 (2005).

•A new measurement of b_1 will be d JLab (E12-13-011).

SIDIS spin-1 measurements open door to a complete new set observables that can tell us about degrees of freedom and beyond sta hadron physics.

Hermes Experiment: First Measurement of b₁

 $\cdot 0.5 \text{ GeV}^2 < Q < 5 \text{ GeV}^2$

 $\cdot 0.01 \le x \le 0.45$

•Positrons in the momentum range of GeV to 27 GeV

•The average target vector P and ten Q polarizations are typically more th •Polarized gas target (integrated lum 42 pb-1)

•The rise of for decreasing values of *be interpreted to originate from the same mechanism that leads to nuclear shadowing in unpolarized scattering.*

Phys.Rev.Lett. 95 (20

Theory predictions of b1

We found that a significant antiquark tensor polarization exists if the overall tensor polarization vanishes for the valence quarks although such a result could depend on the assumed functional form. Further experimental measurements are needed for $b₁$ such as at JLab as *well as Drell-Yan measurements with tensor-polarized deuteron at hadron facilities, J-PARC and GSI-FAIR.*

Hidden-color model: six-quark cordinations $(with \sim 0.15\%$ probability to ex*deuteron) proposed* and *found substantial contributions for values of x > 0.2.*

Phys. Rev. D 82, 017501 (2010) Phys. Rev. C 89, 045203 (2

E12-13-011Approved Experiment at Jefferson Lab

0.012 0.03 Projected ♦ Projected 0.01 **HERMES HERMES** 0.02 Miller b16q Kumano (With δ_{τ} qbar) 0.008 Sargsian (lc) Kumano (No δ_τqbar) 0.006 Sargsian (vn) 0.01 0.004 μ Kumano (With δ, qbar) A \blacktriangleleft $\bf{0}$ Kumano (No δ. qbar) 0.002 s Miller (One at Exch θ -0.01 -0.002 -0.004 -0.02 -0.006 -0.03 0.1 0.2 0.3 0.4 0.5 0.6 0.1 0.2 0.3 0.4 0.5 0.6 $\bf{0}$ θ x_{Bjorken} x_{Bjorken} $0.16 < x < 0.49$ $0.8 < Q^2 < 5.0$ GeV² Incident beam 11 GeV

Inclusive Measurement

Slifer, Chen, Kalantarians, Keller, Long, Rondon, Santiesteban, Solvignon

SI[DIS processes with](https://inspirehep.net/literature?sort=mostrecent&size=25&page=1&q=find%20eprint%202011.08583) a Spin-1 target

Theory developments

•Leading twist: A. Bacchetta (thesis) arXiv:hep-ph/0212025

•Leading twist: Phys. Rev. D 62 (2000)

•Phys. Rev. C 102, 065204 (2020)

•Up to twist 4: Phys. Rev. D 103 (2021)

Explicit cross-sections weren't co estimated for all processes. A theory currently being done.

Longitudinally polarized target

$$
\frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x}\right)
$$
\n
$$
\left\{\frac{F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos \phi_h \frac{F_{UU}^{\cos \phi_h}}{F_{UU}^{\cos \phi_h}} + \varepsilon \cos(2\phi_h) \frac{F_{UU}^{\cos 2\phi_h}}{F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin \phi_h \frac{F_{UU}^{\sin \phi_h}}{F_{UL}^{\sin 2\phi_h}}\right\}
$$
\nvector\n
$$
+ S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_h \frac{F_{UL}^{\sin \phi_h}}{F_{UL}^{\cos \phi_h}} + \varepsilon \sin(2\phi_h) \frac{F_{UL}^{\sin 2\phi_h}}{F_{UL}^{\cos \phi_h}}\right]
$$
\ntensor\n
$$
+ T_{\parallel\parallel} \left[F_{U(LL),T} + \varepsilon F_{U(LL),L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos \phi_h \frac{F_{U(LL)}^{\cos \phi_h}}{F_{U(LL)}^{\cos \phi_h}}\right].
$$
\n
$$
+ \varepsilon \cos(2\phi_h) \frac{F_{U(LL)}^{\cos 2\phi_h}}{F_{U(LL)}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin \phi_h \frac{F_{\perp\parallel}^{\sin \phi_h}}{F_{L(LL)}^{\sin \phi_h}}\right].
$$

Courtesy of A. Bacchetta (private communication) 2023.

Tensor-polarized structure functions

$$
F_{U(LL),T} = \mathcal{C}\left[f_{1LL}D_1\right], \qquad \text{To be}
$$
\n
$$
F_{U(LL),L} = 0,
$$
\n
$$
F_{U(LL)}^{cos \phi_h} = \frac{2M}{Q} \mathcal{C}\left[-\frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M_h} \left(xh_{LL}H_1^\perp + \frac{M_h}{M} \frac{f_{1LL} \tilde{D}^\perp}{f_{1LL}}\right) - \frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M} \left(xf_{LL}^\perp D_1 + \frac{M_h}{M} \frac{h_{1LL}^\perp \tilde{H}}{h_{1LL}^\perp}\right)\right],
$$
\n
$$
F_{U(LL)}^{cos 2\phi_h} = \mathcal{C}\left[-\frac{2\left(\hat{\mathbf{h}} \cdot \mathbf{k}_T\right)\left(\hat{\mathbf{h}} \cdot \mathbf{p}_T\right) - \mathbf{k}_T \cdot \mathbf{p}_T}{MM_h}\frac{h_{1LL}^\perp}{h_{1LL}^\perp}H_1^\perp\right],
$$
\n
$$
F_{L(LL)}^{sin \phi_h} = \frac{2M}{Q} \mathcal{C}\left[-\frac{\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M_h} \left(xe_{LL}H_1^\perp + \frac{M_h}{M} \frac{f_{1LL} \tilde{G}^\perp}{f_{1LL}^\perp}\right) + \frac{\hat{\mathbf{h}} \cdot \mathbf{p}_T}{M} \left(xg_{LL}^\perp D_1 + \frac{M_h}{M} \frac{h_{1LL}^\perp \tilde{E}}{h_{1LL}^\perp \tilde{E}}\right)\right].
$$

Spin-1 leading twist

Courtesy of A. Bacchetta (private communication) 2023.

Simplified version:

$$
\sigma_{meas}^{total} = \sigma_u^D + P\sigma_v + Q\sigma_T + \sum \sigma_i
$$

Summing over positive and negative vector polarization:

$$
\frac{\sigma_T}{\sigma_u^D} = \frac{1}{f} \frac{\sigma_{meas}^{total} - \sigma_{meas}^u}{\sigma_{meas}^u}
$$

 f : Dilution factor due to all other nuclei in the target sample σ_i

CAA: Spin 1 Transverse Momentum Dependent Tensor Structure Functions in CLAS12 Data: Run Group C

- This measurement will help to understand the tensor contribution.
- Currently assuming 10% of the unpolarized contribution as the inclusive measurement.
- Our predictions imply a 60% uncertainty.
- Crucial to propose new experiments.

Poudel, Santiesteban, Chen, Slifer, Ruth, Fernando, Keller, Long, Bacchetta

LOI: Spin-1 TMDs and Structure Functions of the Deuteron

The kinematic ranges assumed for the chosen momentum setting in SHMS (electron) and SBS (hadron)

Ruth, Santiesteban, Chen, Slifer, Poudel, Fernando, Keller, Long, Bacchetta

Step 3: Future program at SoLID

 $0.3 < z < 0.7$ $Q^2 > 1.0$ GeV W > 2.3 GeV W' > 1.6 GeV

 $•$ Pure D $-$ 2 1n + 1p

New Ideas: Spin-1 SIDIS program

•No predictions: Use Hall B data (Run group $C \sim 12\%$ tensor polarization) to estimate the rates and possible sensitivity to structure functions shape/structure.

•Exploratory measurement: Propose a run in the short term (probably around the time of the already approved tensor experiments) to map the longitudinal distributions with better precision.

•Continue target development and plan for all possible configurations of polarization and higher polarizations.

•Formalize a plan to measure the distributions with the SoLID detector.

Thank you!