Deuteron Tensor Structure

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Deuteron Polarization



DNP Technique



Requirements

- High magnetic field (at Jlab typically 5T)
- Low temperature (~1K)
- Microwaves (induce spin transitions)
- CW NMR
- Irradiated material ND₃

Dynamic Nuclear Polarization: technique used to enhance vector polarization



DNP enhancement carries tensor polarization enhancement.

★ Typical average vector polarization in Jefferson lab $P \sim 45\%$ which corresponds to $Q \sim 16\%$

Enhancing tensor polarization



- Low tensor polarization has limited physics experiments.
- New target developments are ongoing, with an enhancement of up to 30%.
- Two experiments to measure tensor observables have been approved.
- Several new experiments are underway.

Quasi-elastic Measurements

Scattering from an unpolarized deuteron target $\rho_{unp}(p_m) = u(p_m)^2 + w(p_m)^2$ $u(p_m): S$ -partial wave of the deuteron $w(p_m): D$ -partial wave of the deuteron

Scattering from tensor polarized target

$$\rho_{20}(p_m, \theta_N) = \frac{2\cos^2(\theta_N) - 1}{2} \left[2\sqrt{2}u(p_m)w(p_m) - w^2(p_m) \right]$$

 θ_N : direction of internal momenta with respect to the polarization axis of the deuteron

Measured asymmetry

$$A_d^T = A_{20}(p_m, \theta_N) = \frac{\rho_{20}(p_m, \theta_N)}{\rho_{unp}(p_m)}$$

Experimental Measurement

P: target vector polarization *Q*: target tensor polarization

 A_D^V : vector analyzing power A_D^T : vector analyzing power

$$\sigma_{pol}(P,Q) = \sigma_{unpol} \left[1 + PA_D^V + \frac{1}{2}QA_D^T \right]$$

$$\sigma_{pol}(-P,Q) = \sigma_{unpol} \left[1 - PA_z + \frac{1}{2}QA_{zz} \right]$$

$$\sigma_{pol}(P,0) = \sigma_{unpol} [1 + PA_z]$$

$$\sigma_{pol}(-P,0) = \sigma_{unpol} [1 - PA_z]$$

$$\sigma_{pol}(-P,0) = \sigma_{unpol}[1-PA_z]$$

$$A_{d}^{T} = \frac{2}{Q} \left(\frac{\sigma_{pol}(P,Q) + \sigma_{pol}(-P,Q)}{\sigma_{pol}(P,0) + \sigma_{pol}(-P,0)} - 1 \right)$$



New Ideas: Exclusive measurement

$$\rho_{node}(p_m) = \rho_{unp}(p_m) + \frac{2\rho_{20}}{3\cos^2(\theta_N) - 1} = u^2(p_m) + 2\sqrt{2}u(p_m)w(p_m)$$

$$We could measure Courtesy of Sargsian (2024)$$

$$A_{node}(p_m) \equiv \frac{\rho_{node}}{\rho_{unp}} = 1 + \frac{2A_{20}(p_m, \theta_N)}{3cos^2(\theta_N) - 1} = \frac{u^2(p_m) + 2\sqrt{2}u(p_m)w(p_m)}{u(p_m)^2 + w(p_m)^2}$$

$$A_{node} = \begin{cases} u(p_m) = -2\sqrt{2}w(p_m) & \mapsto p_m \sim 180MeV \\ u(p_m) = 0 & \mapsto p_m > 400MeV \end{cases} = 0$$

Most direct measurement of the repulsive strength of the nuclear core ever done in electro-nuclear processes.

•has large sensitivity to the position of these nodes

•has sensitivity to the choice of the potential used in calculating the deuteron wave function at $p_m > 300 \, MeV$

•measures the contribution of the *S*-partial wave of the deuteron with respect to the missing momentum

Our observable



PWIA vs FSI



 θ_{nq} : angle between the virtual photon and the recoiling neutron

Possible Experimental setup at Hall C



Semi-Inclusive Deep Inelastic Scattering

Leading twist distribution functions

Quark	$\mathbf{U}\left(\boldsymbol{\gamma}^{+}\right)$		$L(\gamma^+\gamma_5)$		T $(i\sigma^{i+}\gamma_5/\sigma^{i+})$	
Hadron	T-even	T-odd	T-even	T-odd	T-even	T-odd
U	f_1					$[h_1^{\perp}]$
L			g_{1L}		$[h_{1L}^{\perp}]$	
Т		$f_{1\mathrm{T}}^{\scriptscriptstyle \perp}$	g _{1T}		$[h_1], [h_{1\mathrm{T}}^{\perp}]$	
LL	f_{1LL}					$[h_{1\mathrm{LL}}^{\perp}]$
LT	f_{1LT}			g _{1LT}		$[h_{1LT}], [h_{1LT}^{\perp}]$
ТТ	f _{1TT}			g _{1TT}		$[h_{1\mathrm{TT}}], [h_{1\mathrm{TT}}^{\perp}]$

After integrating over the transverse momentum:

Quark	$U(\gamma^+)$		$L(\gamma^+\gamma_5)$		T $(i\sigma^{i+}\gamma_5 / \sigma^{i+})$	
Hadron	T-even	T-odd	T-even	T-odd	T-even	T-odd
U	f_1					
L			$g_{1L}(g_1)$			
Т					$[h_1]$	
LL	$f_{1LL}(b_1)$					
LT						*1
ТТ						



Phys. Rev. D 62 (2000)

Leading twist distribution functions

Quark	U (γ ⁺)		$L(\gamma^+\gamma_5)$		T $(i\sigma^{i+}\gamma_5/\sigma^{i+})$	
Hadron	T-even	T-odd	T-even	T-odd	T-even	T-odd
U	f_1					$[h_1^{\perp}]$
L			g _{1L}		$[h_{1\mathrm{L}}^{\perp}]$	
Т		$f_{1\mathrm{T}}^{\scriptscriptstyle \perp}$	g _{1T}		$[h_1], [h_{1\mathrm{T}}^{\perp}]$	
LL	f_{1LL}					$[h_{1\mathrm{LL}}^{\perp}]$
LT	f _{1LT}			g _{1LT}		$[h_{1LT}], [h_{1LT}^{\perp}]$
TT	f _{1TT}			g _{1TT}		$[h_{1\mathrm{TT}}], [h_{1\mathrm{TT}}^{\perp}]$

After integrating over the transverse momentum:

Quark	U (γ ⁺)		L (γ	$(\gamma^{+}\gamma_{5})$	T $(i\sigma^{i+}\gamma_5/\sigma^{i+})$	
Hadron	T-even	T-odd	T-even	T-odd	T-even	T-odd
U	f_1					
L			$g_{1L}(g_1)$			
Т					[<i>h</i> ₁]	
LL	$f_{1LL}(b_1)$					
LT						*1
ТТ						

Only b₁ has been measured by Hermes <u>Phys.Rev.Lett. 95 (2005)</u>.
A new measurement of b₁ will be done at JLab (<u>E12-13-011)</u>.

SIDIS spin-1 measurements open the door to a complete new set of observables that can tell us about color degrees of freedom and beyond standard hadron physics.

Hermes Experiment: First Measurement of b₁



- •0.5 GeV² < Q < 5 GeV²
- •0.01 < x < 0.45

•Positrons in the momentum range of 2.5 GeV to 27 GeV

The average target vector P and tensor Q polarizations are typically more than 80%
Polarized gas target (integrated luminosity 42 pb⁻¹)

•The rise of for decreasing values of x can be interpreted to originate from the same mechanism that leads to nuclear shadowing in unpolarized scattering.

Phys.Rev.Lett. 95 (2005)

Theory predictions of b₁



We found that a significant antiquark tensor polarization exists if the overall tensor polarization vanishes for the valence quarks although such a result could depend on the assumed functional form. Further experimental measurements are needed for b_1 such as at JLab as well as Drell-Yan measurements with tensor-polarized deuteron at hadron facilities, J-PARC and GSI-FAIR.

Hidden-color model: six-quark configurations (with ~ 0.15% probability to exist in the deuteron) proposed and found to give substantial contributions for values of x > 0.2.

Phys. Rev. D 82, 017501 (2010)

Phys. Rev. C 89, 045203 (2014)

E12-13-011Approved Experiment at Jefferson Lab

0.012 0.03 Projected Projected 0.01 HERMES HERMES 0.02 Miller b16q Kumano (With δ_rqbar) 0.008 Sargsian (Ic) Kumano (No δ_{τ} qbar) 0.006 0.01 Sargsian (vn) A_d^T 0.004 Kumano (With δ_πqbar) e Kumano (No δ, gbar) 0 0.002 -2 Miller (One at Exch 0 -0.01 -0.002 -0.004 -0.02 -0.006 -0.03 0.6 0.2 0.1 0.2 0.3 0.4 0.5 0.1 0.3 0.4 0.5 0.6 0 0 x_{Bjorken} x_{Bjorken} 0.16 < x < 0.49 $0.8 < Q^2 < 5.0 \text{ GeV}^2$ Incident beam 11 GeV

Inclusive Measurement

Slifer, Chen, Kalantarians, Keller, Long, Rondon, Santiesteban, Solvignon

SIDIS processes with a Spin-1 target

Theory developments

•Leading twist: A. Bacchetta (thesis) arXiv:hep-ph/0212025

- •Leading twist: Phys. Rev. D 62 (2000)
- •Phys. Rev. C 102, 065204 (2020)

•Up to twist 4: Phys. Rev. D 103 (2021)

Explicit cross-sections weren't completely estimated for all processes. A theory effort is currently being done.

Longitudinally polarized target

$$\begin{aligned} \frac{d\sigma}{dx\,dy\,d\psi\,dz\,d\phi_h\,dP_{h\perp}^2} &= \frac{\alpha^2}{xyQ^2} \frac{y^2}{2\left(1-\varepsilon\right)} \left(1+\frac{\gamma^2}{2x}\right) & \text{published} \\ &\left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)}\cos\phi_h \frac{F_{UU}^{\cos\phi_h}}{F_{UU}^{\cos\phi_h}} + \varepsilon\cos(2\phi_h) \frac{F_{UU}^{\cos2\phi_h}}{F_{UU}^{\cos2\phi_h}} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)}\sin\phi_h \frac{F_{LU}^{\sin\phi_h}}{F_{UL}^{\sin\phi_h}} \right. \\ &\left. + \varepsilon \cos(2\phi_h) \frac{F_{UU}^{\cos2\phi_h}}{F_{UU}^{\cos\phi_h}} + \varepsilon\sin(2\phi_h) \frac{F_{UL}^{\sin2\phi_h}}{F_{UL}^{\sin\phi_h}} \right] \\ &\left. + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} \frac{F_{LL}}{F_{LL}} + \sqrt{2\varepsilon(1-\varepsilon)}\cos\phi_h \frac{F_{UL}^{\cos\phi_h}}{F_{UL}^{\cos\phi_h}} \right] \\ &\left. + \varepsilon\cos(2\phi_h) \frac{F_{U(LL),T}}{F_{U(LL),L}} + \varepsilon \frac{F_{U(LL),L}}{F_{U(LL)}} + \sqrt{2\varepsilon(1-\varepsilon)}\cos\phi_h \frac{F_{U(LL)}^{\cos\phi_h}}{F_{U(LL)}^{\cos\phi_h}} \right] \right\}. \end{aligned}$$

Courtesy of A. Bacchetta (private communication) 2023.

Tensor-polarized structure functions

$$\begin{split} F_{U(LL),T} &= \mathcal{C}[f_{1LL}D_{1}], & \text{To be}\\ \text{published} \\ F_{U(LL),L} &= 0, \end{split}$$

$$\begin{split} F_{U(LL),L}^{\cos\phi_{h}} &= \frac{2M}{Q} \mathcal{C}\left[-\frac{\hat{h} \cdot k_{T}}{M_{h}}\left(xh_{LL}H_{1}^{\perp} + \frac{M_{h}}{M}f_{1LL}\frac{\tilde{D}^{\perp}}{z}\right) - \frac{\hat{h} \cdot p_{T}}{M}\left(xf_{LL}^{\perp}D_{1} + \frac{M_{h}}{M}h_{1LL}\frac{\tilde{H}}{z}\right)\right], \\ F_{U(LL)}^{\cos2\phi_{h}} &= \mathcal{C}\left[-\frac{2\left(\hat{h} \cdot k_{T}\right)\left(\hat{h} \cdot p_{T}\right) - k_{T} \cdot p_{T}}{MM_{h}}h_{1LL}^{\perp}H_{1}^{\perp}\right], \\ F_{L(LL)}^{\sin\phi_{h}} &= \frac{2M}{Q} \mathcal{C}\left[-\frac{\hat{h} \cdot k_{T}}{M_{h}}\left(xe_{LL}H_{1}^{\perp} + \frac{M_{h}}{M}f_{1LL}\frac{\tilde{G}^{\perp}}{z}\right) + \frac{\hat{h} \cdot p_{T}}{M}\left(xg_{LL}^{\perp}D_{1} + \frac{M_{h}}{M}h_{1LL}\frac{\tilde{E}}{z}\right)\right]. \end{split}$$

Spin-1 leading twist

Courtesy of A. Bacchetta (private communication) 2023.



Simplified version:

$$\sigma_{meas}^{total} = \sigma_u^D + \frac{P\sigma_v}{P\sigma_v} + \frac{Q\sigma_T}{Q\sigma_T} + \sum \sigma_i$$

Summing over positive and negative vector polarization:

$$\frac{\sigma_T}{\sigma_u^D} = \frac{1}{f} \frac{\sigma_{meas}^{total} - \sigma_{meas}^u}{\sigma_{meas}^u}$$

f: Dilution factor due to all other nuclei in the target sample σ_i

CAA: Spin 1 Transverse Momentum Dependent Tensor Structure Functions in CLAS12 Data: Run Group C



- This measurement will help to understand the tensor contribution.
- Currently assuming 10% of the unpolarized contribution as the inclusive measurement.
- Our predictions imply a 60% uncertainty.
- Crucial to propose new experiments.

Poudel, Santiesteban, Chen, Slifer, Ruth, Fernando, Keller, Long, Bacchetta





LOI: Spin-1 TMDs and Structure Functions of the Deuteron



	heta (deg.)	$\phi~({ m deg.})$	P (GeV)
Electron	10.3 - 12.4	-2.87 - 2.87	4.0 - 5.4
Hadron	5.0 - 15.0	167 - 193	2.0 - 4.0

The kinematic ranges assumed for the chosen momentum setting in SHMS (electron) and SBS (hadron)

Ruth, Santiesteban, Chen, Slifer, Poudel, Fernando, Keller, Long, Bacchetta

Step 3: Future program at SoLID



0.3 < z < 0.7 Q² > 1.0 GeV W > 2.3 GeV W' > 1.6 GeV



•Pure D-> 1n + 1p

New Ideas: Spin-1 SIDIS program



•No predictions: Use Hall B data (Run group C ~ 12% tensor polarization) to estimate the rates and possible sensitivity to structure functions shape/structure.

•Exploratory measurement: Propose a run in the short term (probably around the time of the already approved tensor experiments) to map the longitudinal distributions with better precision.

•Continue target development and plan for all possible configurations of polarization and higher polarizations.

•Formalize a plan to measure the distributions with the SoLID detector.

Thank you!