

# Generalized Polarizabilities from VCS Experiments

Hamza Atac-Nikos Sparveris

The 12<sup>th</sup> Workshop on Hadron Physics & Opportunities  
Worldwide

August 2024

# Outline

Introduction to the VCS and GPs

Status

Results from recent experiments / Jlab & MAMI

Spatial information & polarizability radii

Prospects

# Proton

Despite being the only stable composite particle, several fundamental characteristics of the proton remain not fully understood, including its **mass**, **spin**, **charge radius**, and **polarizabilities**.

If we want to understand the characteristics of the proton as a building block of the universe... **we need to understand the dynamics of the proton's constituents and how they contribute to those emergent characteristics.**

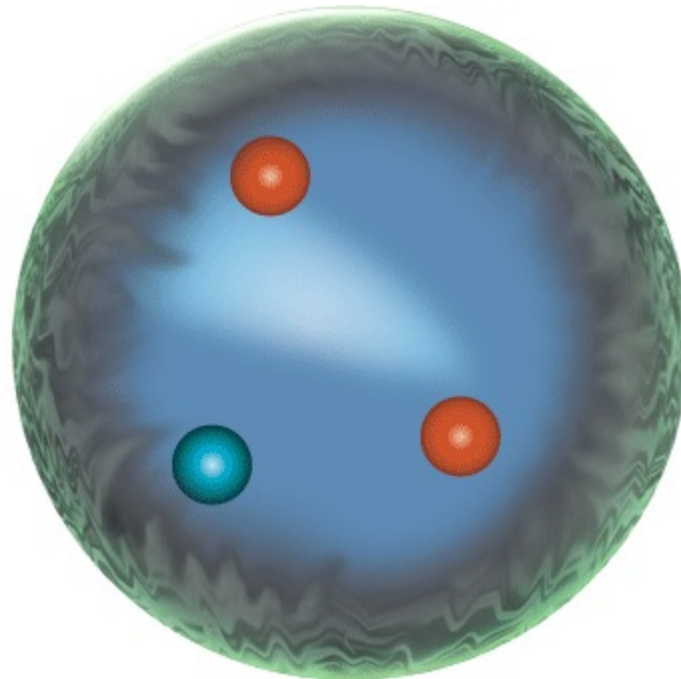
Mass

Spin

Size

Shape

Polarizabilities



.....

# Proton Polarizabilities

Fundamental structure constants  
(such as mass, size, shape, ...)

Response of the nucleon to external EM field

Sensitive to the full excitation spectrum

Accessed experimentally through Compton Scattering

RCS: static polarizabilities

Virtual Compton Scattering:

Virtuality of photon gives access to the GPs :  $\alpha_E(Q^2)$  &  $\beta_M(Q^2)$  + spin GPs

- spatial distribution of the polarization densities
- electric & magnetic polarizability radii

Fourier transform of densities of electric charges and magnetization of a nucleon deformed by an applied EM field

PDG

150 Baryon Summary Table

<b>N BARYONS</b> <b>(S = 0, I = 1/2)</b> <i>p, N<sup>+</sup> = uud; n, N<sup>0</sup> = udd</i>
--

**p**

$$I(J^P) = \frac{1}{2}(\frac{1}{2}^+)$$

Mass  $m = 1.00727646681 \pm 0.00000000009$  u

Mass  $m = 938.272046 \pm 0.000021$  MeV [a]

$|m_p - m_{\bar{p}}|/m_p < 7 \times 10^{-10}$ , CL = 90% [b]

$|\frac{q_p}{m_p}|/(\frac{q_e}{m_e}) = 0.99999999991 \pm 0.00000000009$

$|q_p + q_{\bar{p}}|/e < 7 \times 10^{-10}$ , CL = 90% [b]

$|q_p + q_e|/e < 1 \times 10^{-21}$  [c]

Magnetic moment  $\mu = 2.792847356 \pm 0.000000023 \mu_N$

$(\mu_p + \mu_{\bar{p}}) / \mu_p = (0 \pm 5) \times 10^{-6}$

Electric dipole moment  $d < 0.54 \times 10^{-23}$  e cm

Electric polarizability  $\alpha = (11.2 \pm 0.4) \times 10^{-4} \text{ fm}^3$

Magnetic polarizability  $\beta = (2.5 \pm 0.4) \times 10^{-4} \text{ fm}^3$  (S = 1.2)

Charge radius,  $\mu p$  Lamb shift =  $0.84087 \pm 0.00039$  fm [d]

Charge radius,  $e p$  CODATA value =  $0.8775 \pm 0.0051$  fm [d]

Magnetic radius =  $0.777 \pm 0.016$  fm

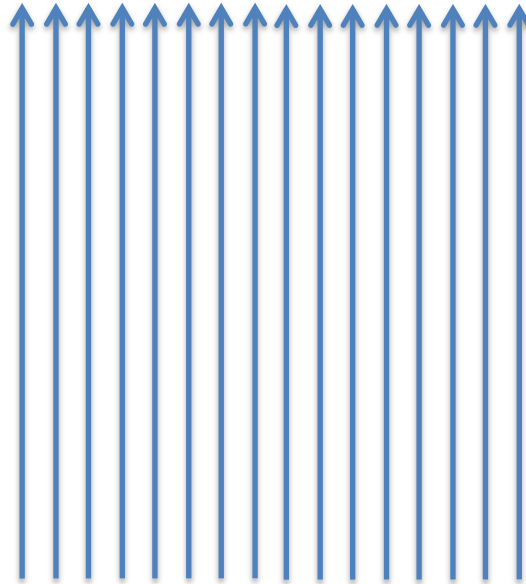
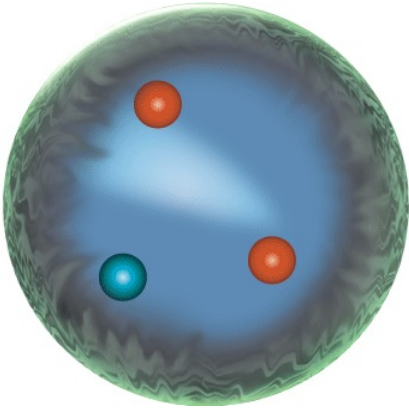
Mean life  $\tau > 2.1 \times 10^{29}$  years, CL = 90% [e] ( $p \rightarrow$  invisible mode)

Mean life  $\tau > 10^{31}$  to  $10^{33}$  years [e] (mode dependent)

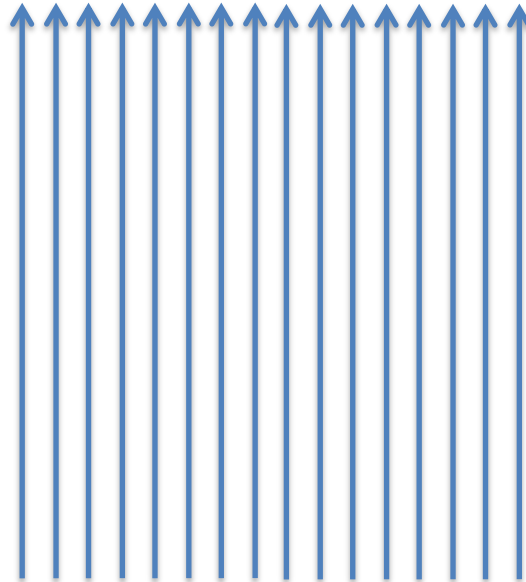
# Scalar Polarizabilities

Response of internal structure to an applied EM field

Interaction of the EM field with the internal structure of the nucleon



$\vec{E}$

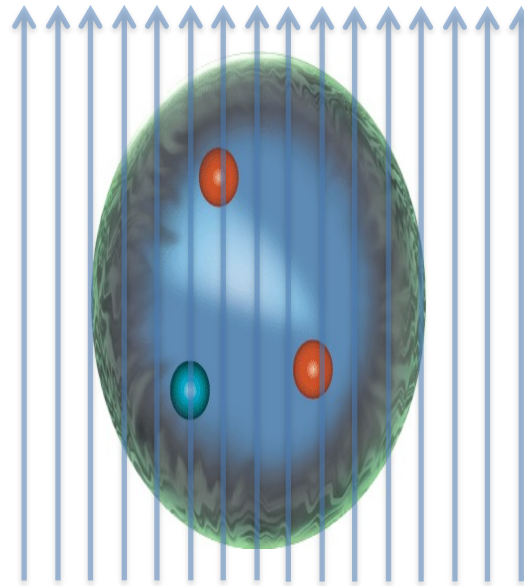
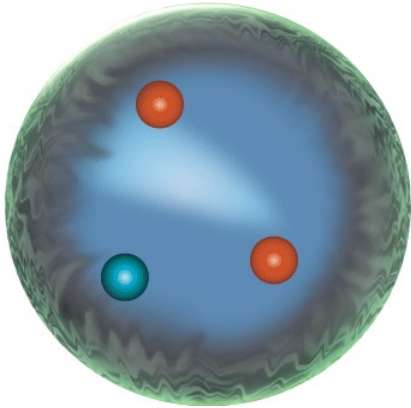


$\vec{B}$

# Scalar Polarizabilities

Response of internal structure to an applied EM field

Interaction of the EM field with the internal structure of the nucleon

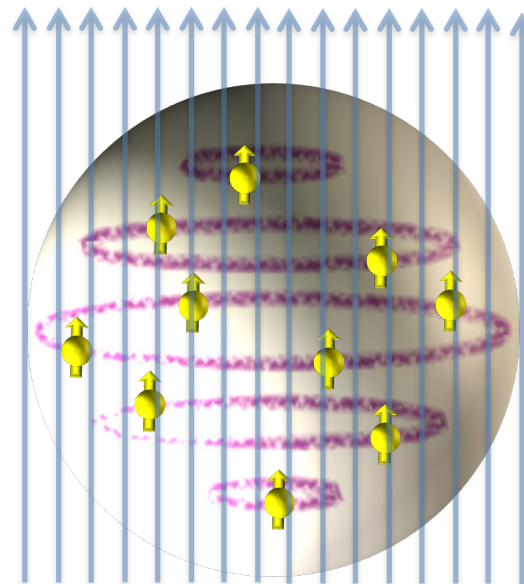


$\vec{E}$

“stretchability”

$$\vec{d}_{E \text{ induced}} \sim \alpha \vec{E}$$

External field deforms the charge distribution



$\vec{B}$

“alignability”

$$\vec{d}_{M \text{ induced}} \sim \beta \vec{B}$$

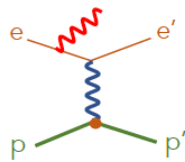
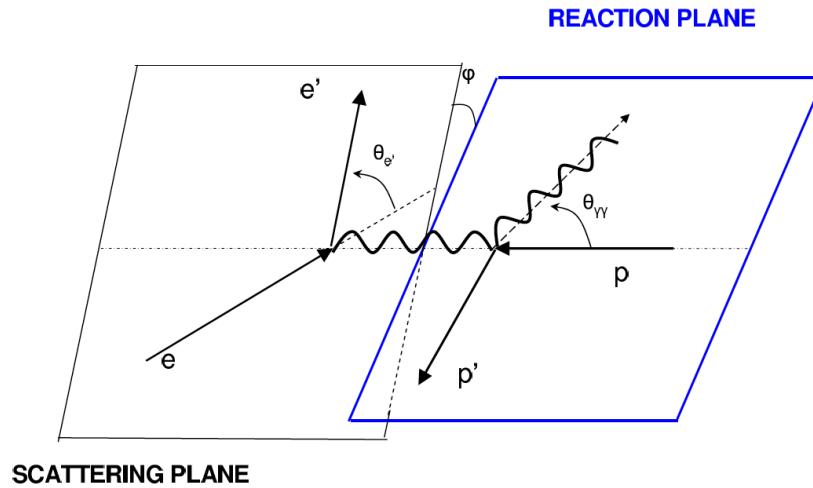
$$\beta_{\text{para}} > 0$$

$$\beta_{\text{diam}} < 0$$

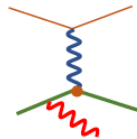
Paramagnetic: proton spin aligns with the external magnetic field

Diamagnetic:  $\pi$ -cloud induction produces field counter to the external perturbation

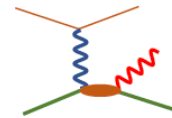
# Virtual Compton Scattering



Bethe-Heitler



Born VCS



non-Born VCS

Elastic FFs

GPs

# Virtual Compton Scattering

DR

valid below & above  
Pion threshold

Dispersive integrals  
for Non Born amplitudes

Spin GPs are fixed

Scalar GPs have  
an unconstrained part

Fit to the experimental  
cross sections at each  $Q^2$

LEX

valid only below  
Pion threshold

Response functions

$$d^5\sigma = d^5\sigma^{BH+Born} + q'_{cm} \cdot \phi \cdot \Psi_0 + \mathcal{O}(q'^2_{cm})$$

$$\Psi_0 = v_1 \cdot \left( P_{LL} - \frac{1}{\epsilon} P_{TT} \right) + v_2 \cdot P_{LT}$$

Subtract the spin part

$$P_{TT} = [P_{TT} \text{ spin}]$$

$$P_{LT} = -\frac{2M}{\alpha_{em}} \sqrt{\frac{q'^2_{cm}}{Q^2}} \cdot G_E^p(Q^2) \cdot \beta_M(Q^2) + [P_{LT} \text{ spin}]$$

utilize DR

scalar GPs  $\alpha_E$  and  $\beta_M$

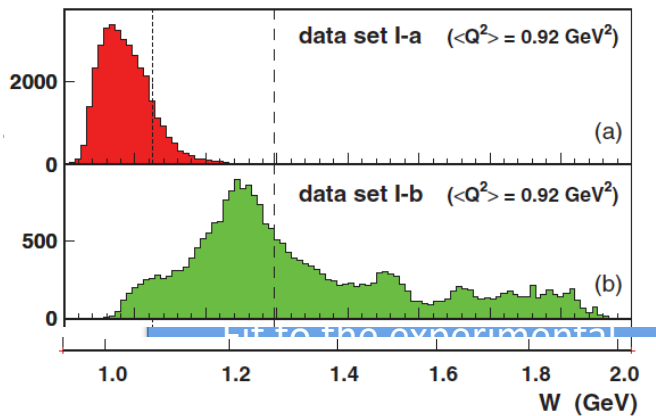


# Virtual Compton Scattering

DR

valid below & above  
Pion threshold

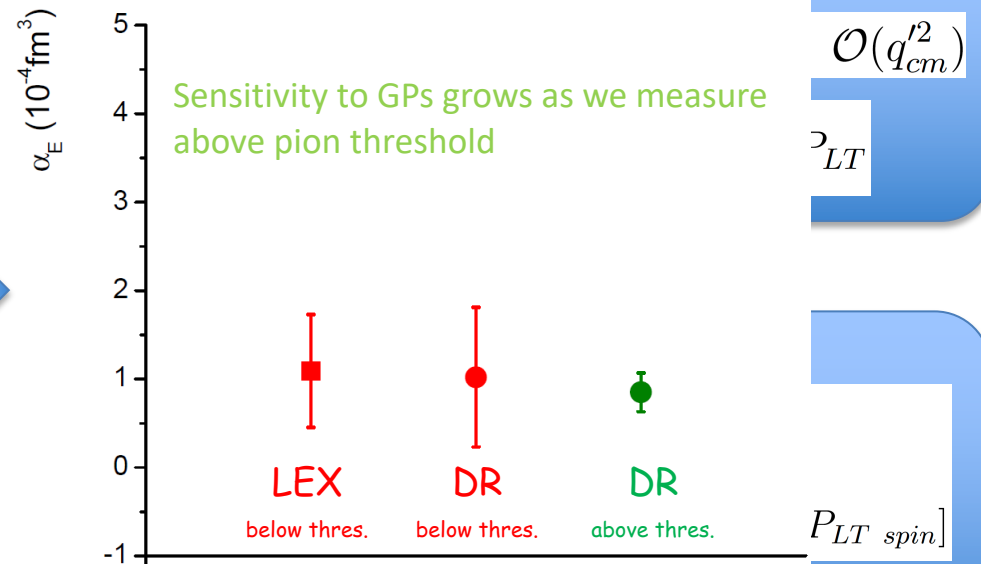
Dispersive integrals  
for Non Born amplitudes



LEX

valid only below  
Pion threshold

Response functions

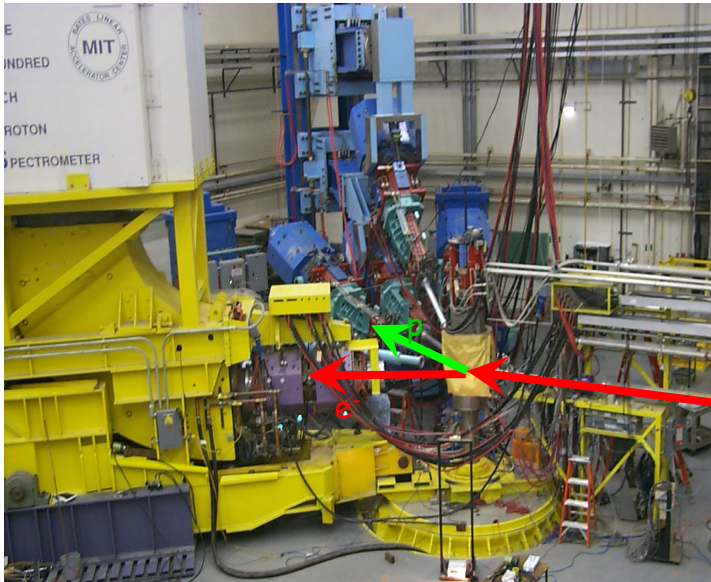


utilize DR

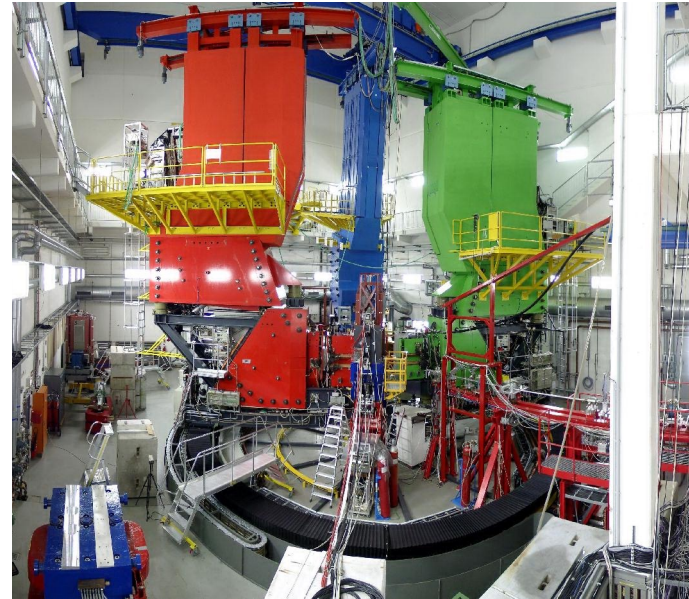
scalar GPs  $\alpha_E$  and  $\beta_M$

# Early Experiments

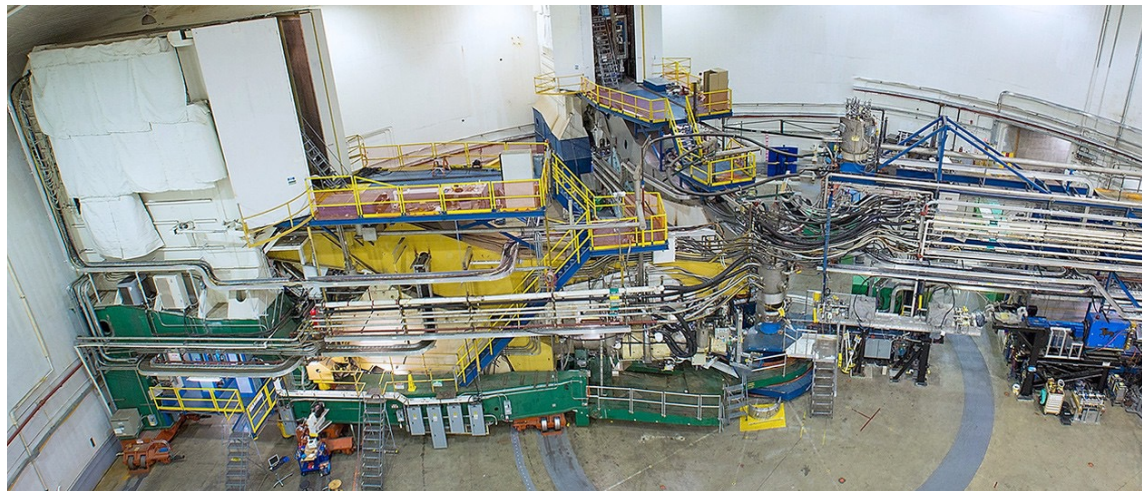
MIT-Bates @  $Q^2=0.06 \text{ GeV}^2$



MAMI-A1 @  $Q^2=0.33 \text{ GeV}^2$



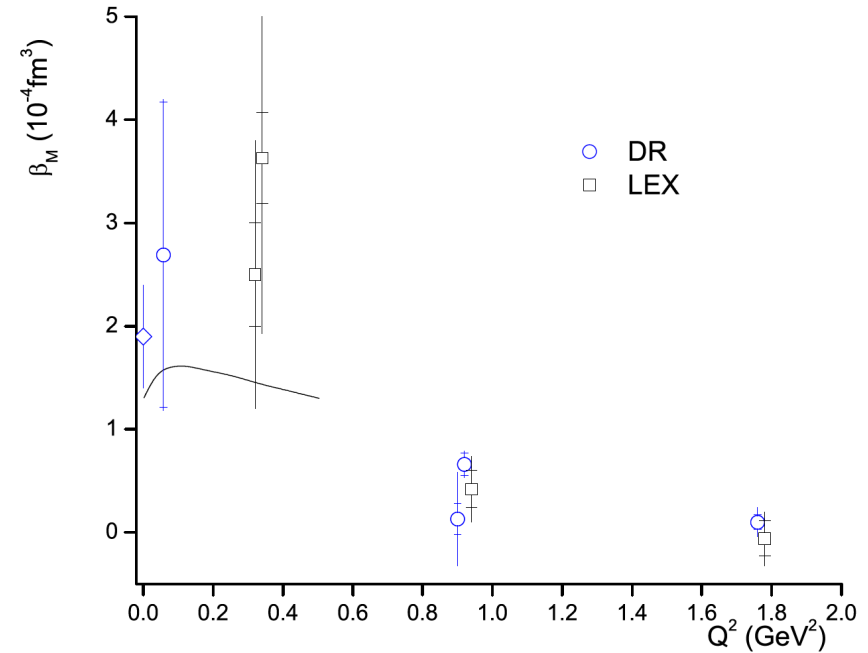
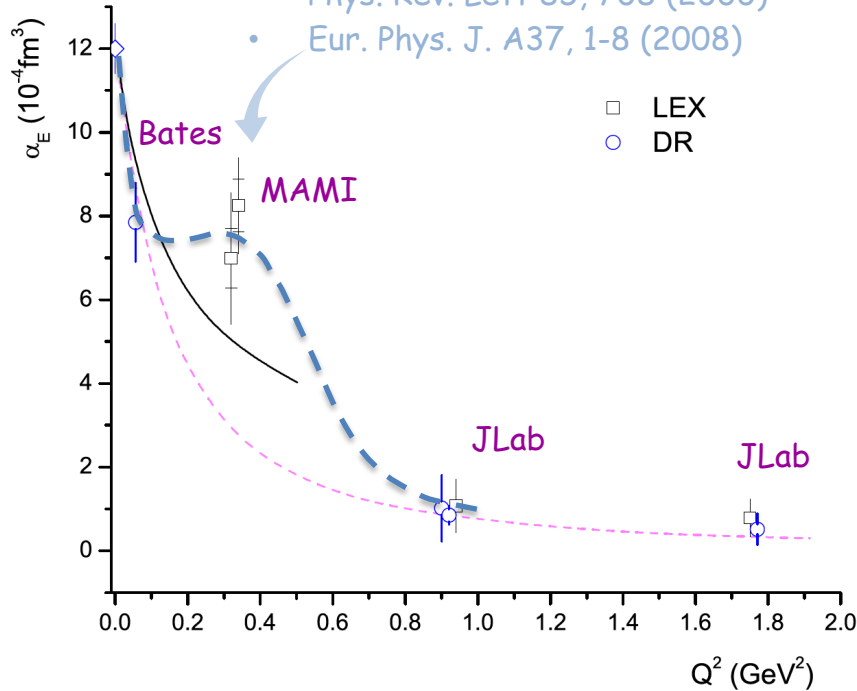
Jlab-Hall A @  $Q^2=0.9 \text{ \& } 1.8 \text{ GeV}^2$



# Early Experiments

$Q^2 = 0.33 \text{ (GeV/c)}^2$  measured twice at MAMI:

- Phys. Rev. Lett 85, 708 (2000)
- Eur. Phys. J. A37, 1-8 (2008)



$a_E \approx 10^{-3} V_N$  (stiffness / relativistic character)

Data: non-trivial  $Q^2$  dependence of  $a_E$  (?)

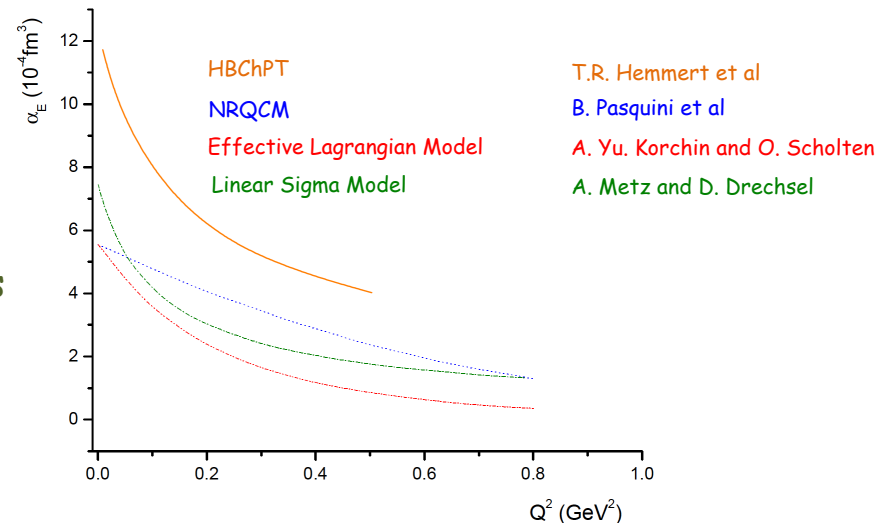
Theory: monotonic fall-off

$\beta_M$  small  $\leftrightarrow$  cancellation of competing mechanisms

Large uncertainties

Higher precision measurements needed

$\rightarrow$  Quantify balance between dia/para-magnetism



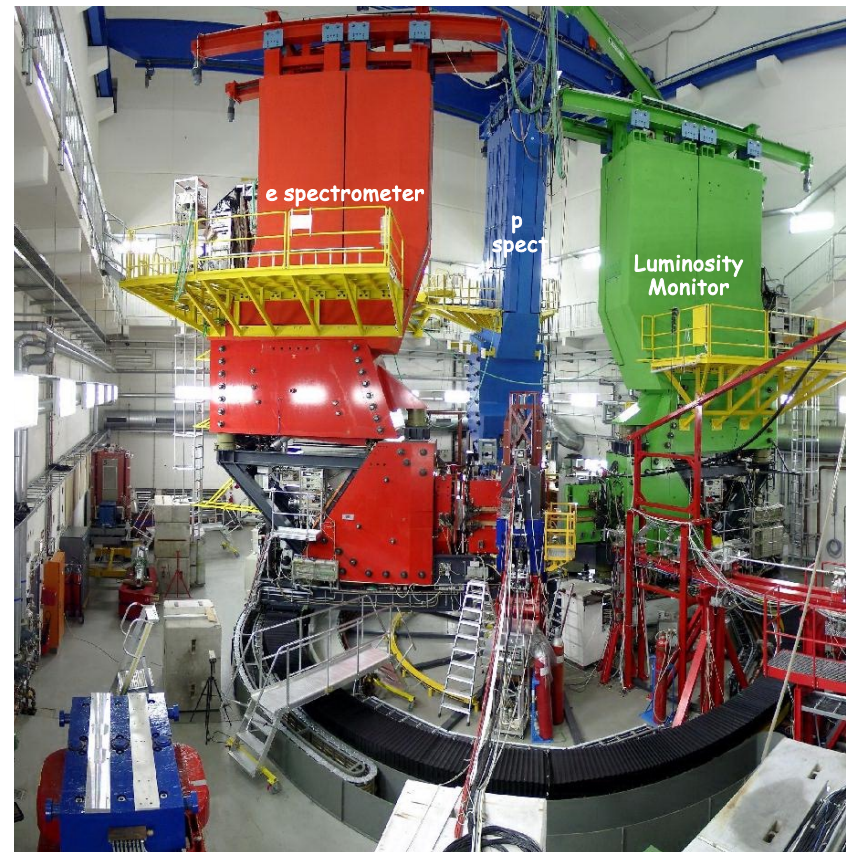
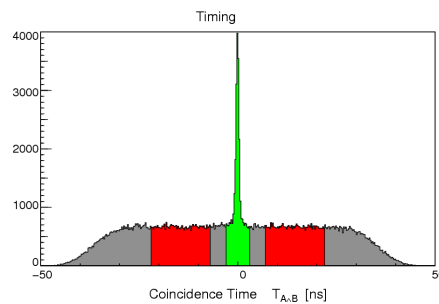
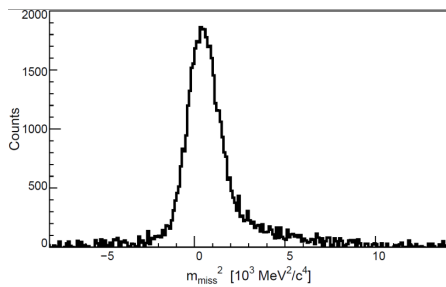
# Recent Experiments

# Recent Measurements: MAMI

MAMI A1/1-09 (vcsq2) below threshold

MAMI A1/3-12 (vcsdelta) above threshold

Both experiments utilized the A1 setup at MAMI



# A1/1-09 @ MAMI

Several improvements were implemented compared to the early MAMI experiments.

e.g. for LEX the higher order terms have to be kept small / under control

$$d^5\sigma = d^5\sigma^{BH+Born} + q'_{cm} \cdot \phi \cdot \Psi_0 + \mathcal{O}(q'^2_{cm})$$

Refined analysis procedure / phase space masking to keep these terms smaller than  $\sim 2\%$ - $3\%$  level

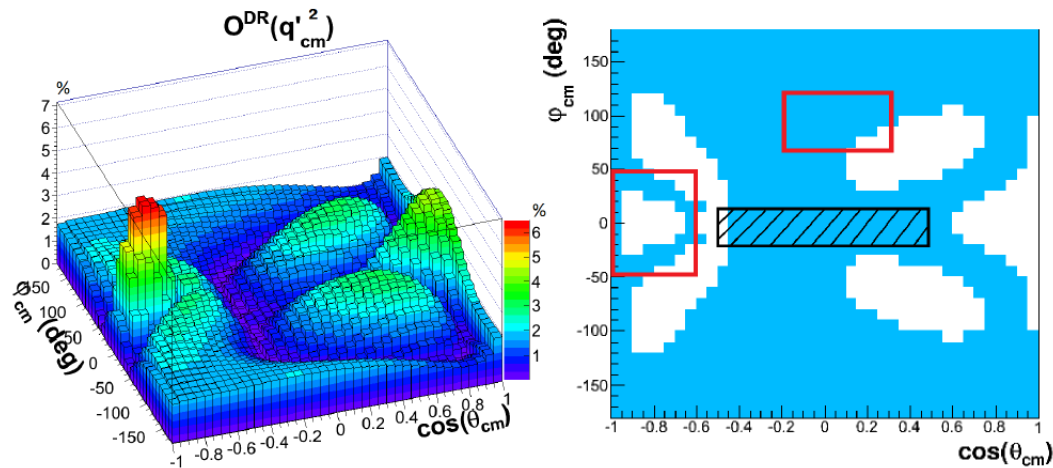
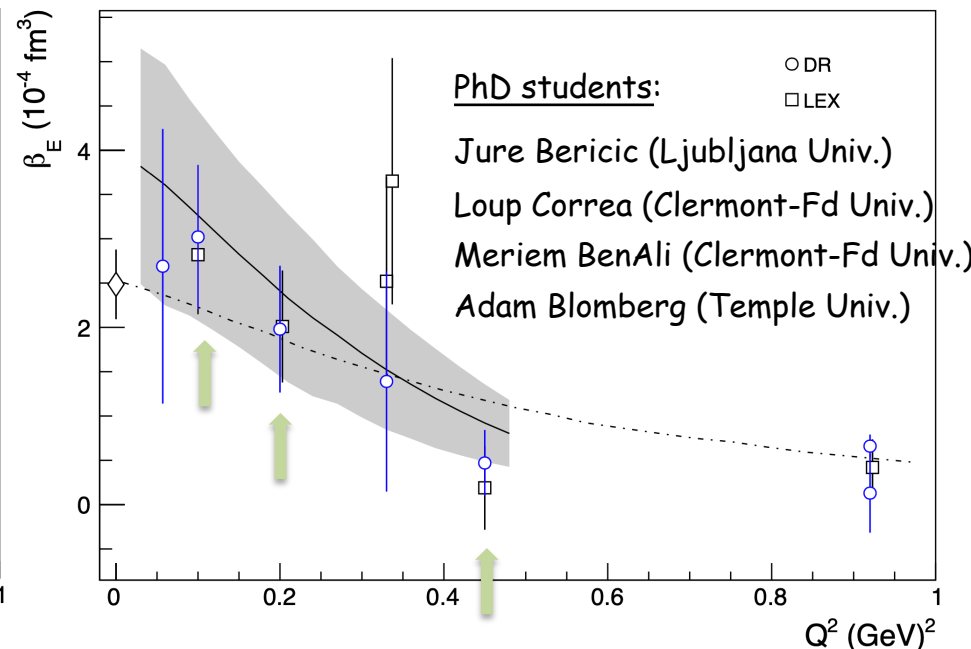
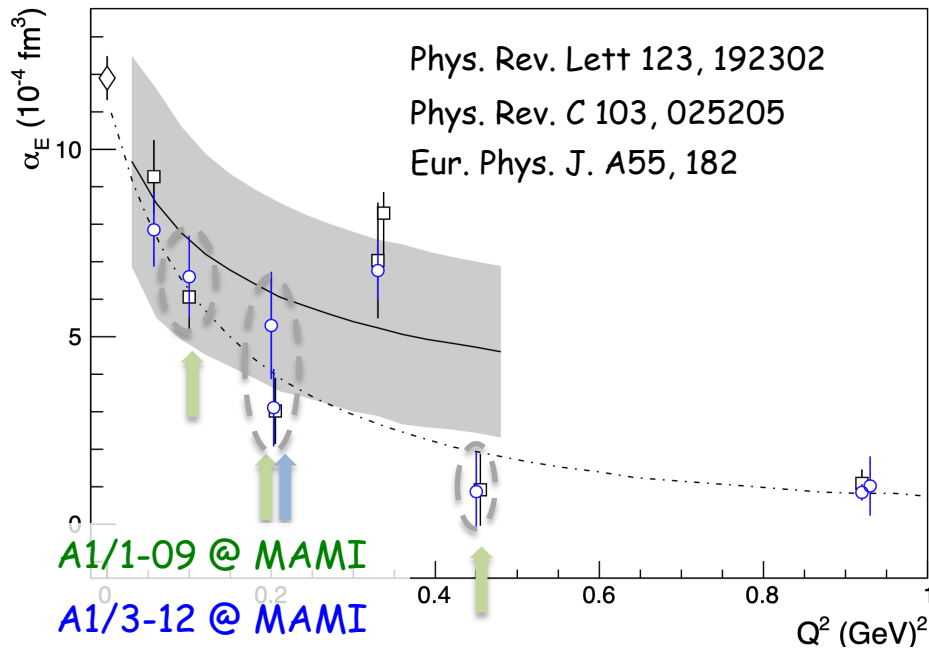


Figure 3.13: (Left) behavior of  $\mathcal{O}^{DR}(q'^2_{cm})$  in the  $(\cos(\theta_{cm}), \varphi_{cm})$ -plane at  $q'_{cm} = 87.5 \text{ MeV}/c$  and (right) two-dimensional representation of the angular region where  $\mathcal{O}^{DR}(q'^2_{cm}) < 2\%$  (blue), the red squares correspond to the two areas of interest to perform the GP extraction.

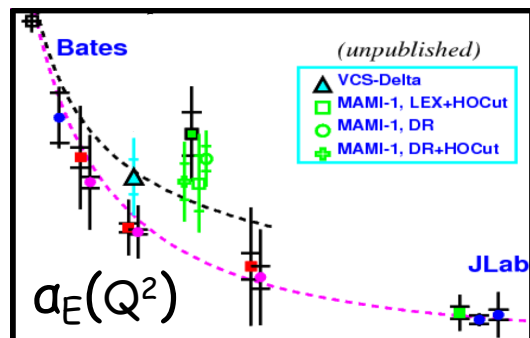
# MAMI Results



## Revisiting the $Q^2=0.33 \text{ GeV}^2$ data

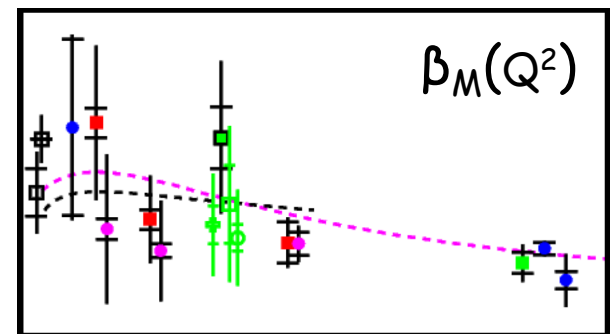
Analysis revisited (unpublished):

The  $\alpha_E$  puzzle still holds



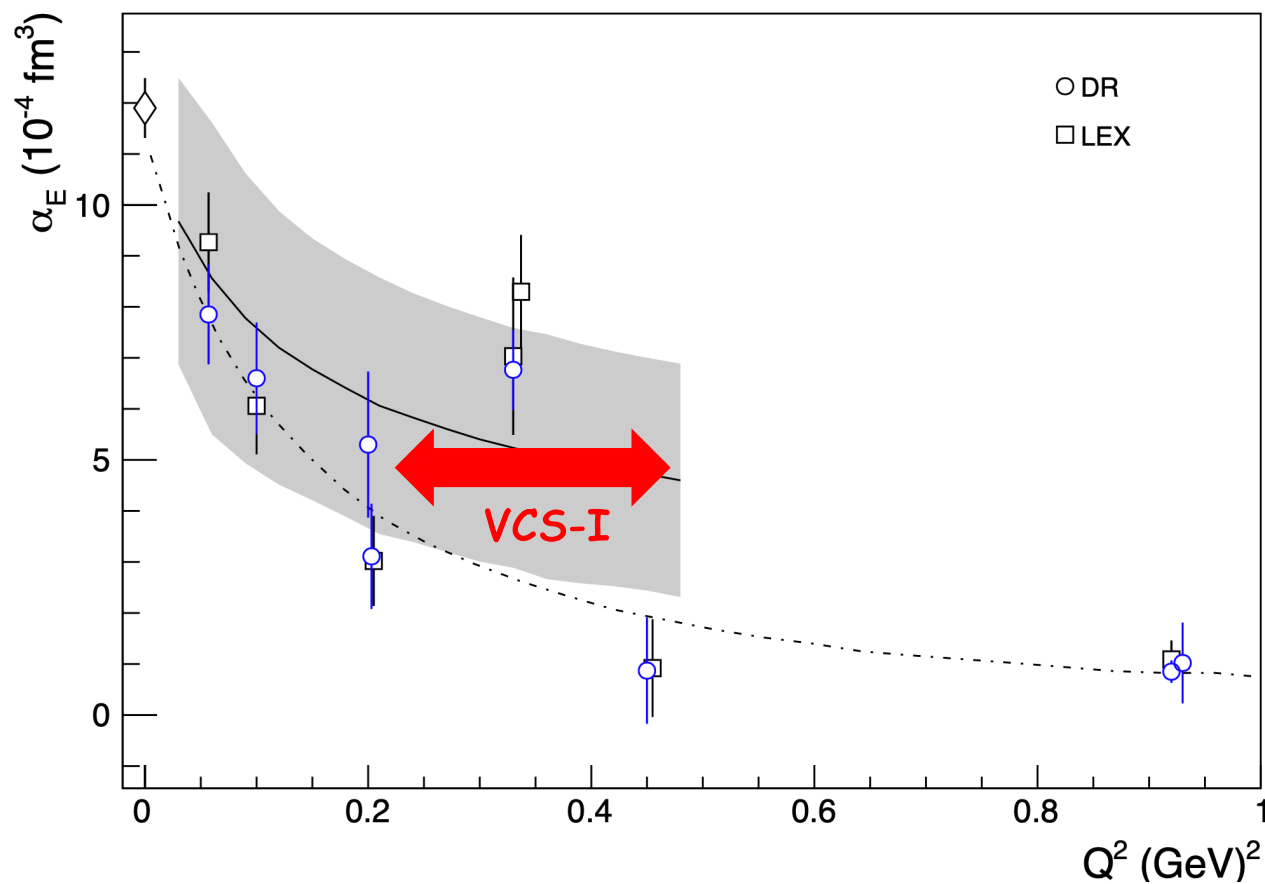
Re-fits at  
 $Q^2=0.33$   
 $\text{GeV}^2$   
 (H.F.)

LEX and DR  
 Updated HO-cut



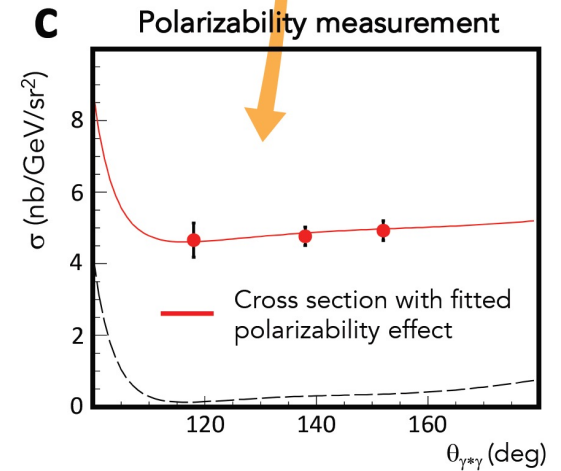
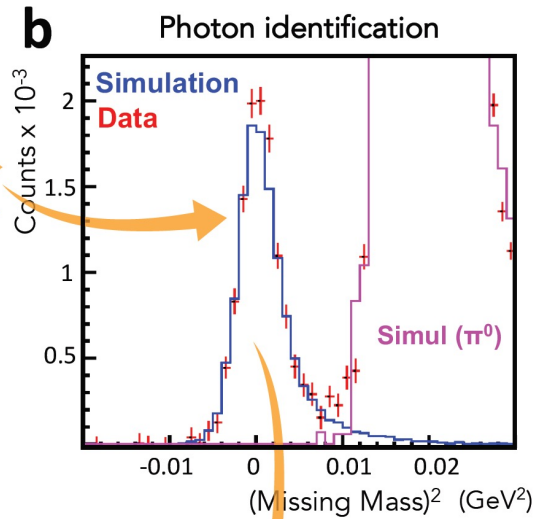
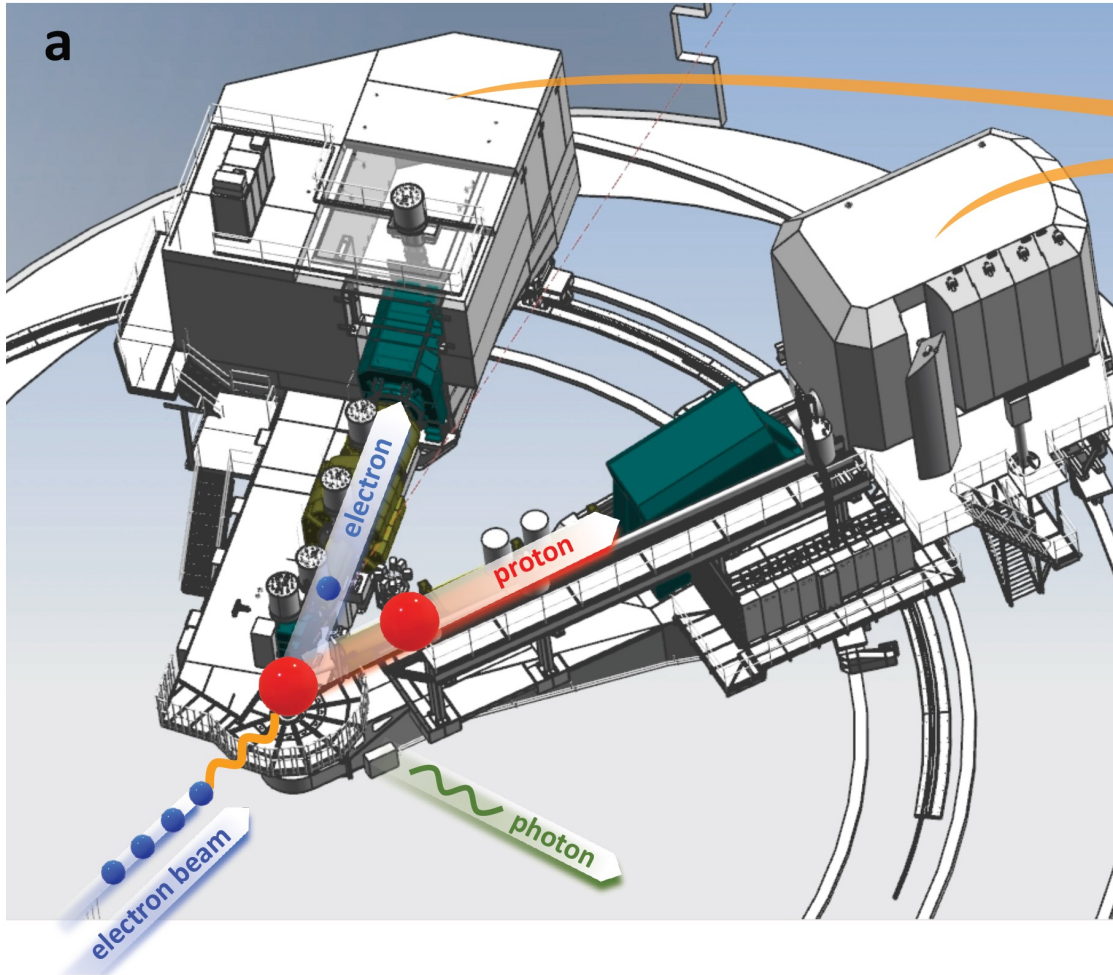
# Jlab : VCS-I Experiment (E12-15-001) in Hall C

High precision measurements targeting explicitly the kinematics of interest for  $a_E$





# The experiment



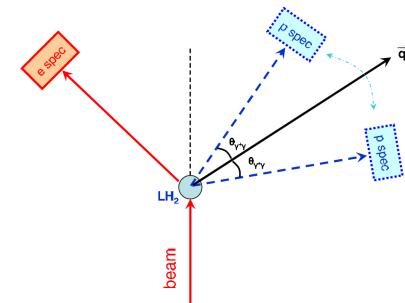
Hall C: SHMS, HMS  
 4.56 GeV  
 20 μA  
 Liquid hydrogen 10 cm

cross sections & azimuthal asymmetries

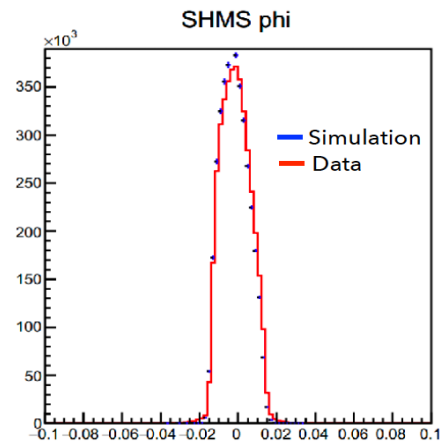
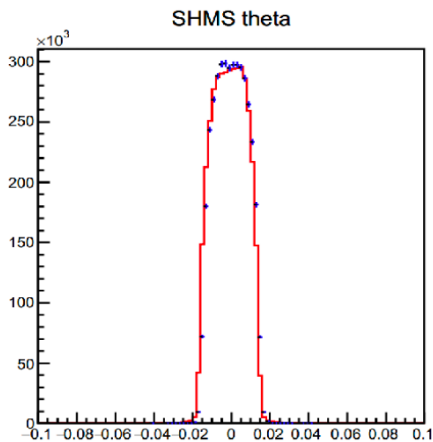
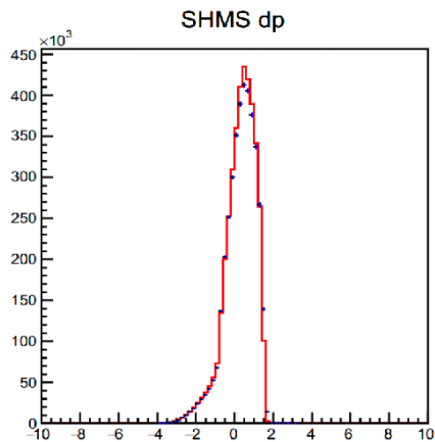
$$A_{(\phi_{\gamma^*\gamma}=0,\pi)} = \frac{\sigma_{\phi_{\gamma^*\gamma}=0} - \sigma_{\phi_{\gamma^*\gamma}=180}}{\sigma_{\phi_{\gamma^*\gamma}=0} + \sigma_{\phi_{\gamma^*\gamma}=180}}$$

sensitivity to GPs

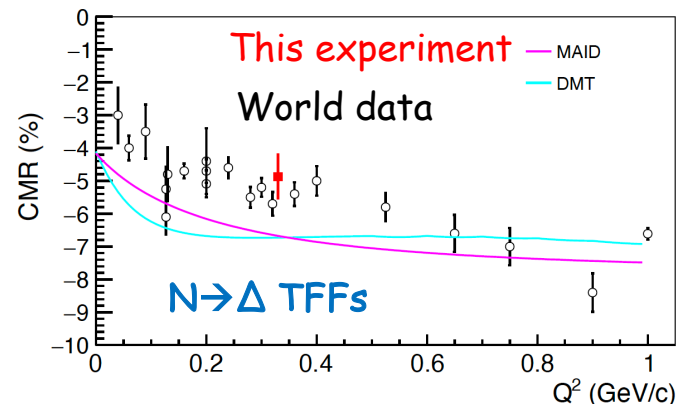
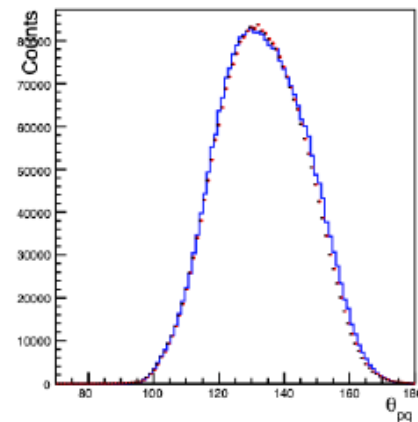
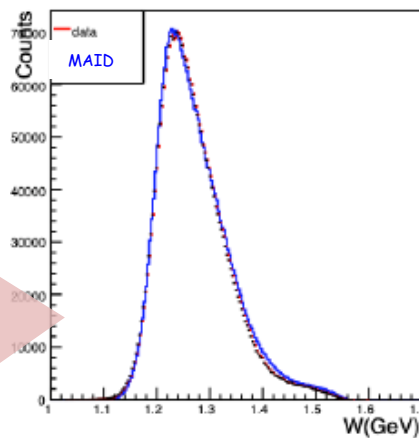
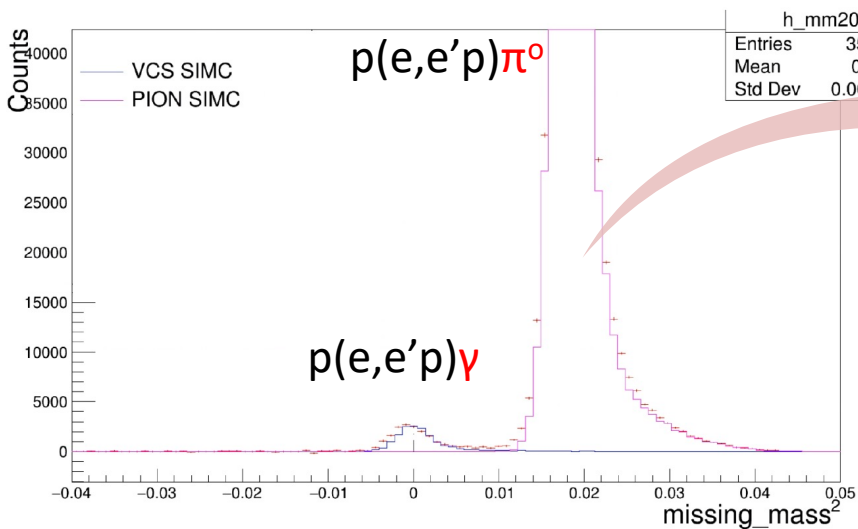
suppression of systematic asymmetries



# Elastic data

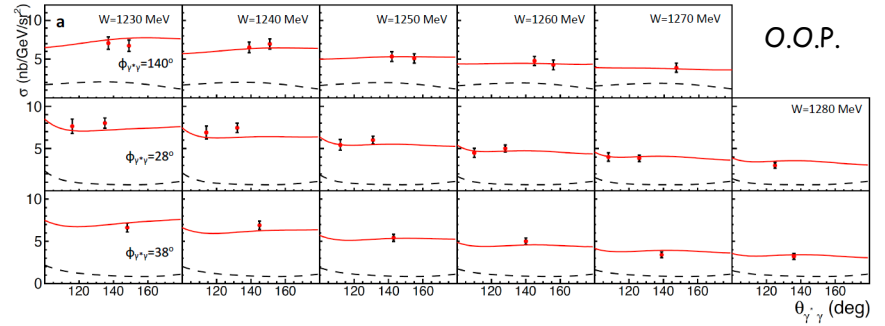
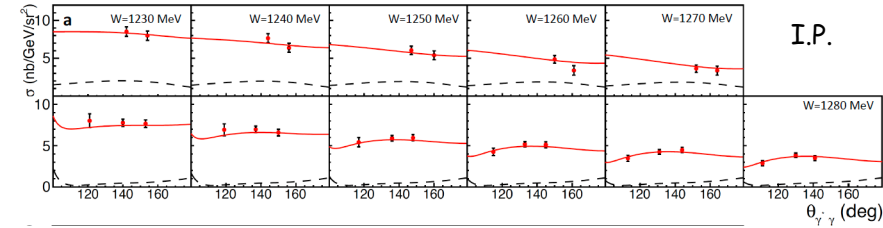


# $p(e,e'p)\pi^0$

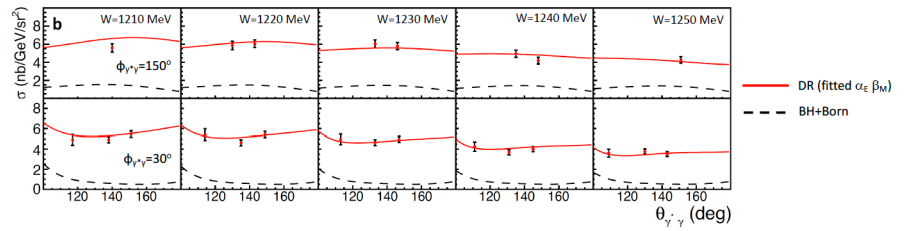
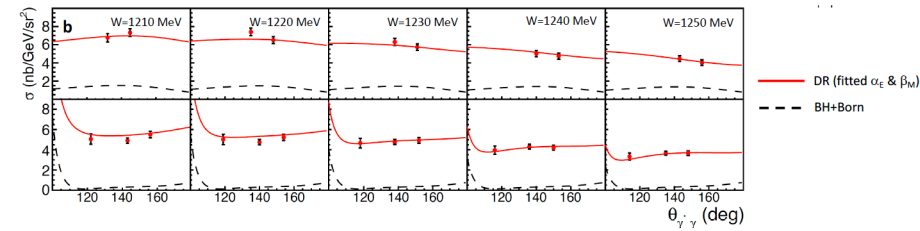


# VCS-I results: cross sections

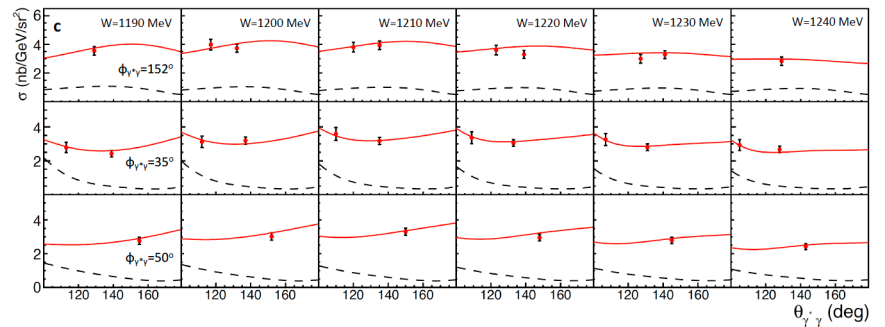
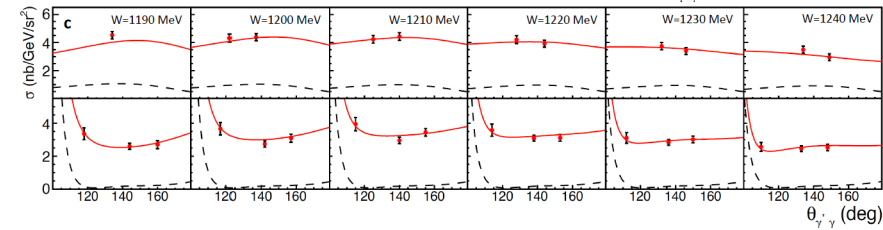
$Q^2=0.27 \text{ GeV}^2$



$Q^2=0.33 \text{ GeV}^2$

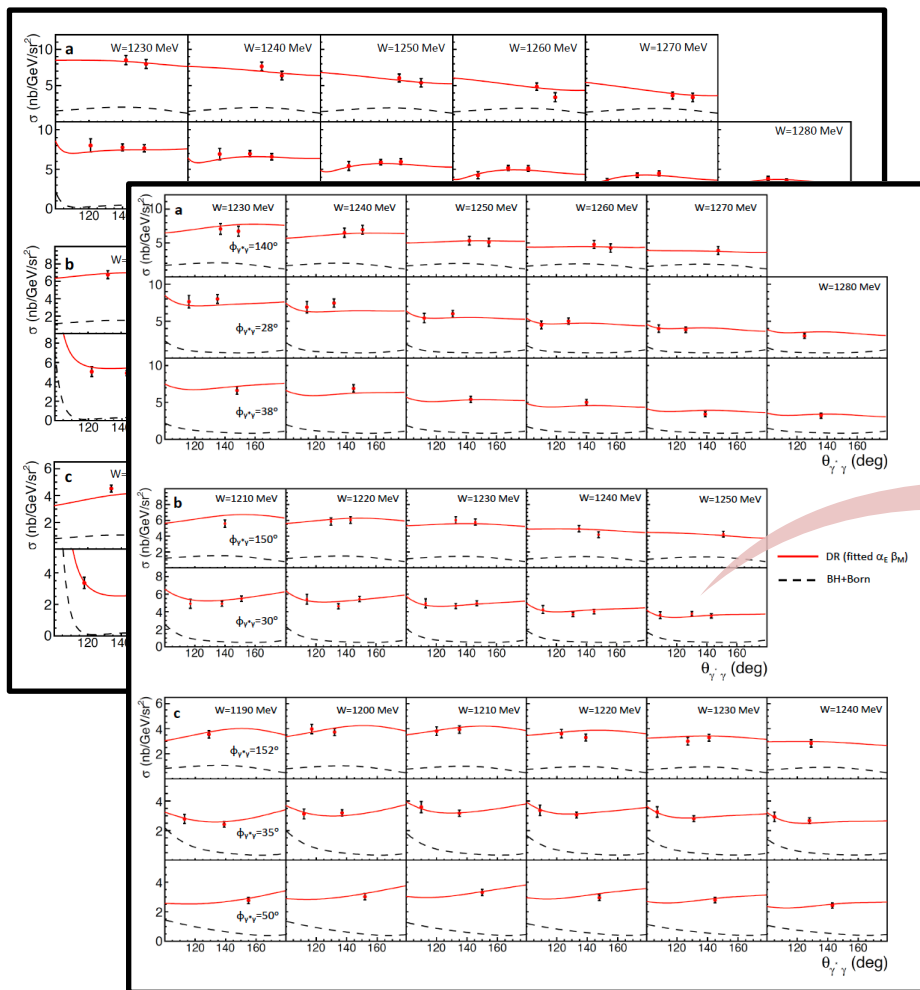


$Q^2=0.40 \text{ GeV}^2$

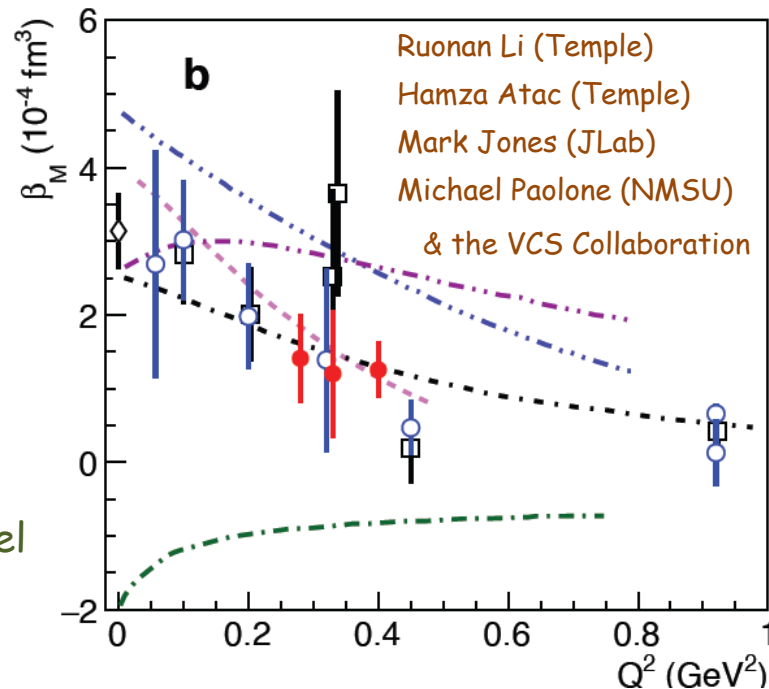
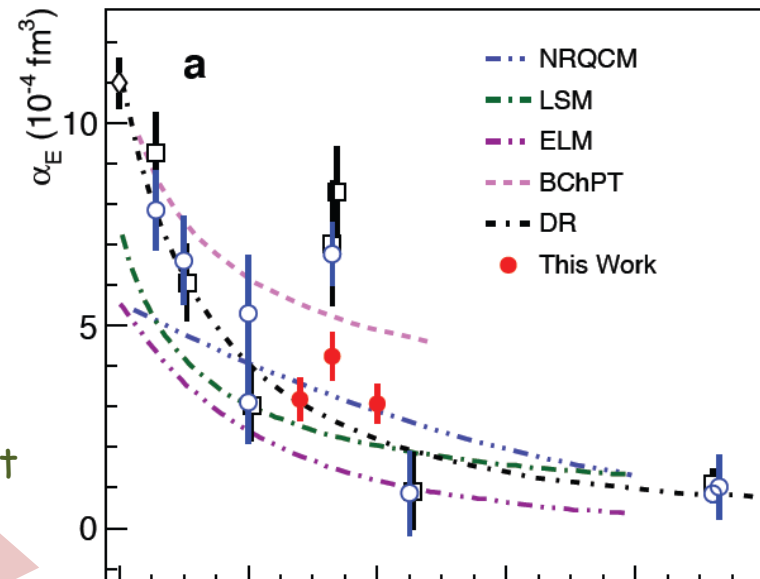


# VCS-I results: GPs

Nature 611, 265 (2022)



DR fit

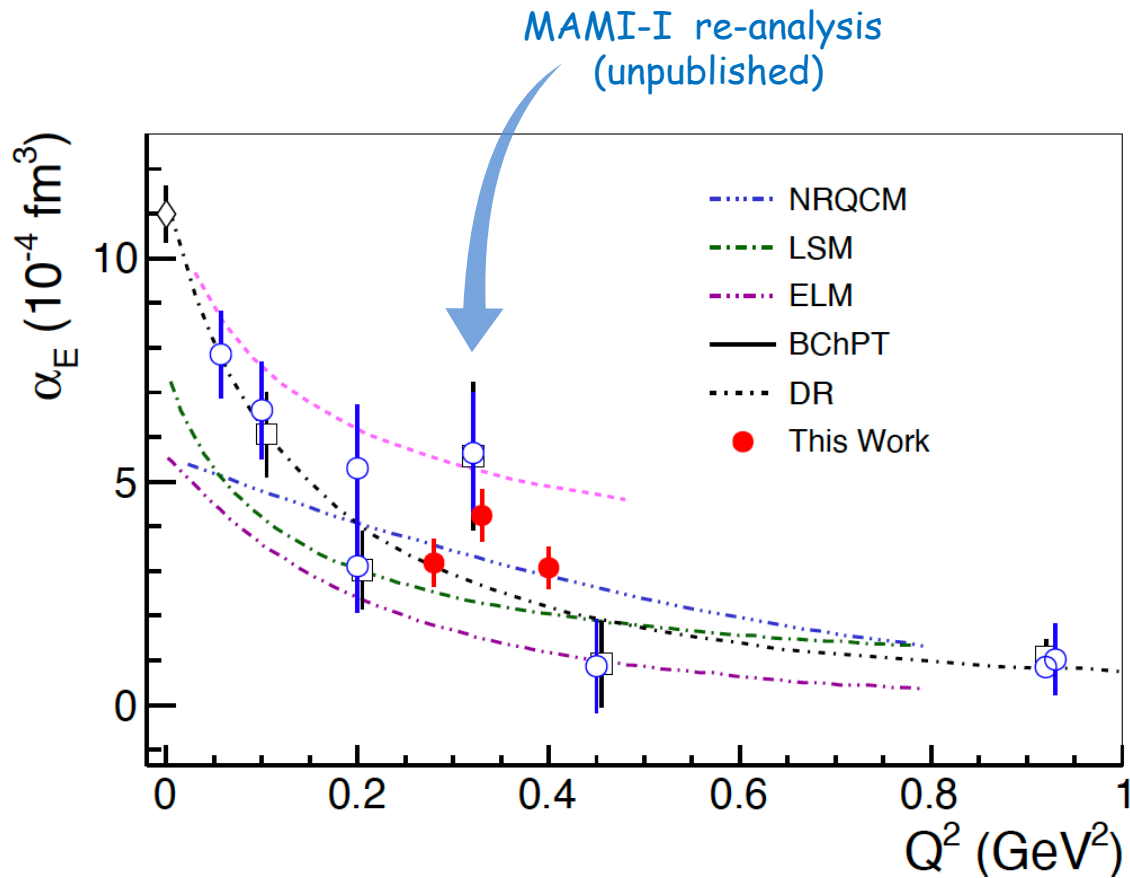


Experimental cross sections are compared to the DR model predictions for all possible values for the GPs

→  $\alpha_E(Q^2)$  and  $\beta_M(Q^2)$  are fitted by a  $\chi^2$  minimization

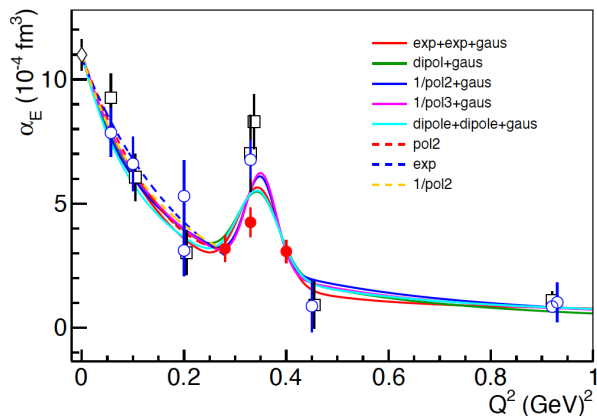
# Electric GP

Is there a non-trivial structure vs  $Q^2$  ?

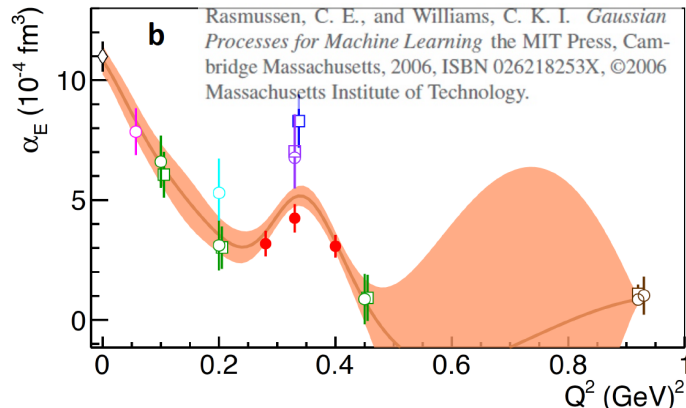


# Electric GP

Traditional fits



Data-driven techniques:  
no underlying functional  
form is assumed



Is the observed  $a_E$  structure coincidental or not?

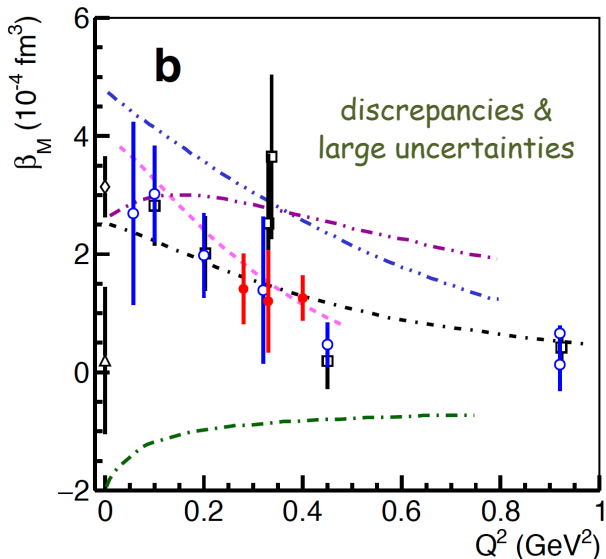
If true: Measure the shape precisely  $\rightarrow$  input to theory

If not: We are able to show it with more measurements

Strong tension between world data (?)

Something we do not yet understand well?  
Underestimated uncertainties? ...

# Magnetic GP



Magnetic GP: Large uncertainties & discrepancies

Precision and consistent systematics are needed to disentangle para/dia-magnetism in the proton

# Spatial information & polarizability radii

# Spatial dependence of induced polarizations

Nucleon form factor data → light-front quark charge densities

Formalism extended to the deformation of these quark densities when applying an external e.m. field:

GPs → spatial deformation of charge & magnetization densities under an applied e.m. field

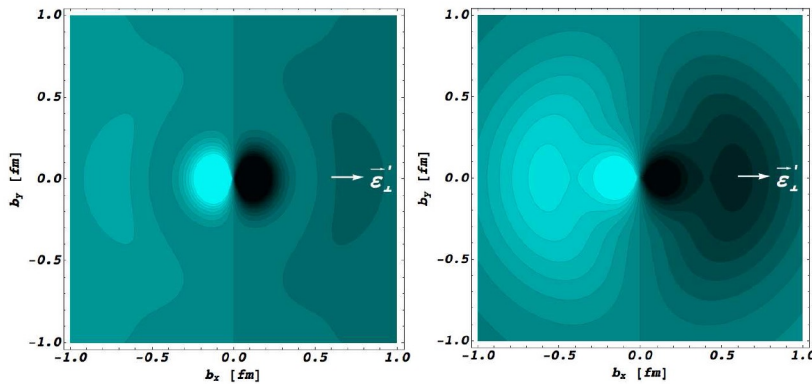
## Induced polarization in a proton when submitted to an e.m. field

Phys. Rev. Lett. 104, 112001 (2010)

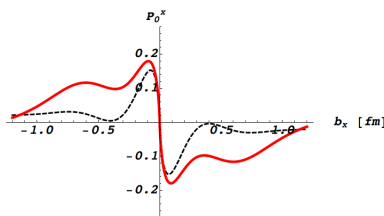
M. Gorchtein, C. Lorce, B. Pasquini, M. Vanderhaeghen

GP I

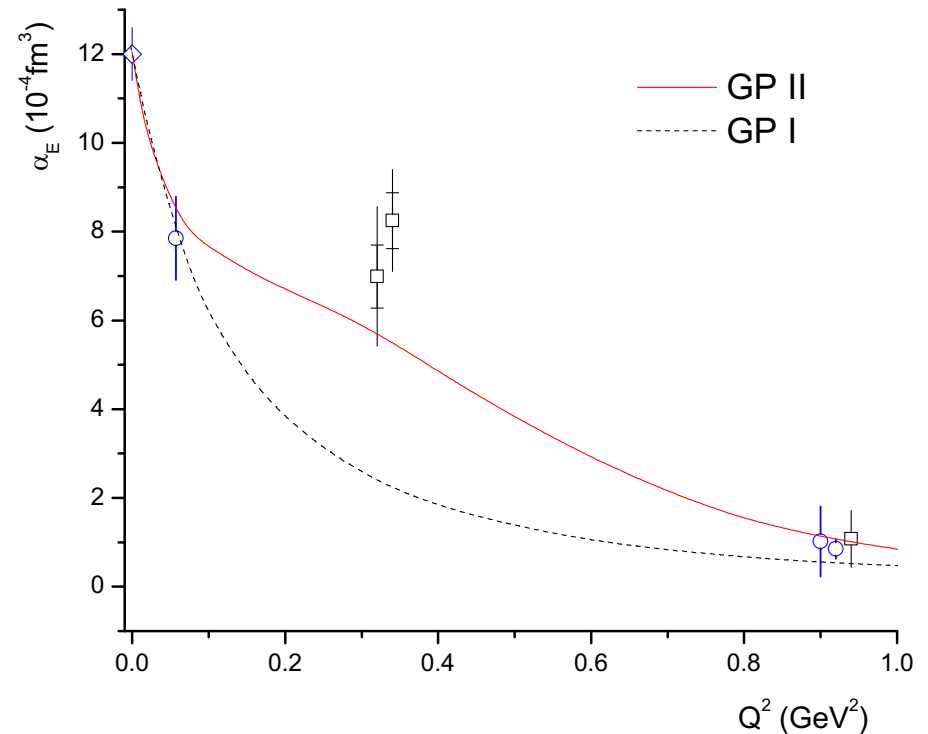
GP II



Light (dark) regions → largest (smaller) values  
(photon polarization along x-axis, as indicated)



Induced polarization along  $b_y=0$



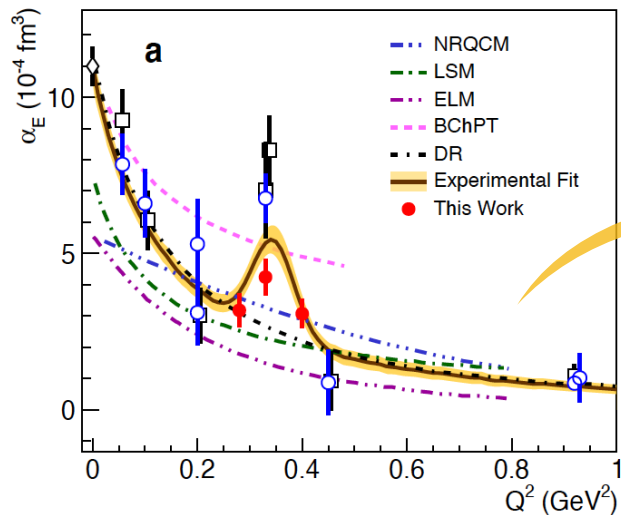


# Spatial dependence of induced polarizations

Nucleon form factor data → light-front quark charge densities

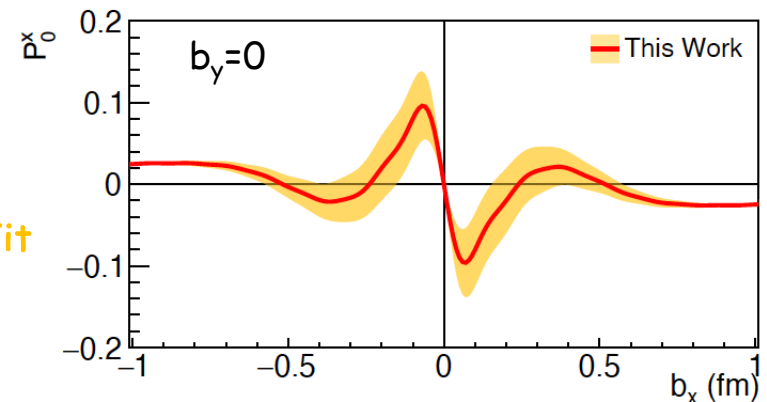
Formalism extended to the deformation of these quark densities when applying an external e.m. field:

GPs → spatial deformation of charge & magnetization densities under an applied e.m. field



Experimental Fit

## Induced polarization in a proton when submitted to an e.m. field



$x$ - $y$  defines the transverse plane with the  $z$ -axis being the direction of the fast-moving proton

## Polarizability radii

$$\langle r_{\alpha_E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2=0}$$

$$\langle r_{\beta_M}^2 \rangle = \frac{-6}{\beta_M(0)} \cdot \frac{d}{dQ^2} \beta_M(Q^2) \Big|_{Q^2=0}$$

# Polarizability radii

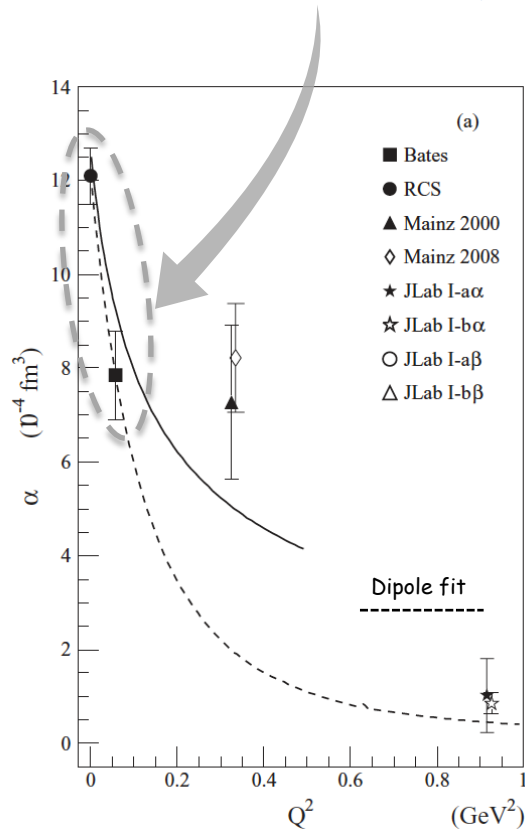
$$\langle r_{\alpha E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2=0}$$

$$\langle r_{\beta M}^2 \rangle = \frac{-6}{\beta_M(0)} \cdot \frac{d}{dQ^2} \beta_M(Q^2) \Big|_{Q^2=0}$$

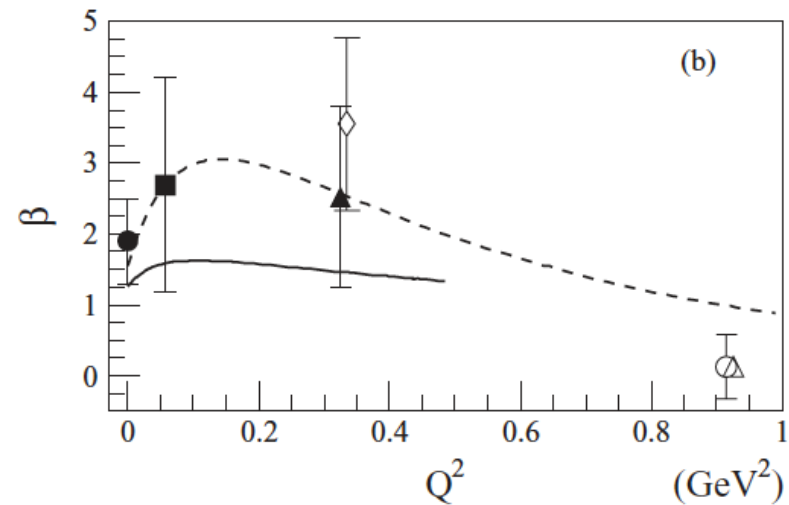
First extraction made possible with the MIT-Bates measurement ( $Q^2=0.06 \text{ GeV}^2$ )

PRL **97**, 212001 (2006)

PHYSICAL REVIEW C **84**, 035206 (2011)



$$\langle r_{\alpha}^2 \rangle = 2.16 \pm 0.31 \text{ fm}^2$$

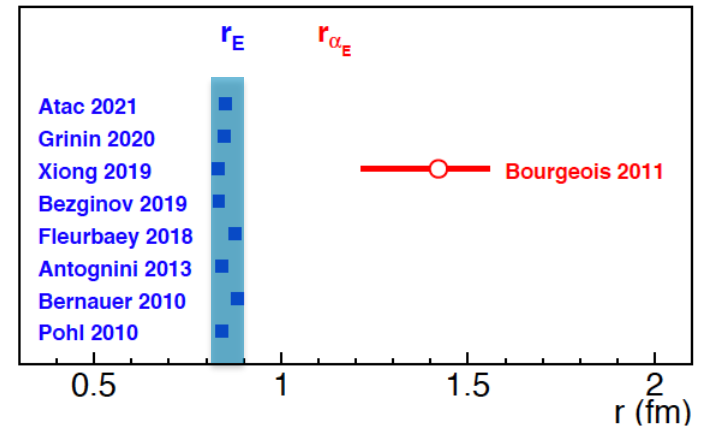
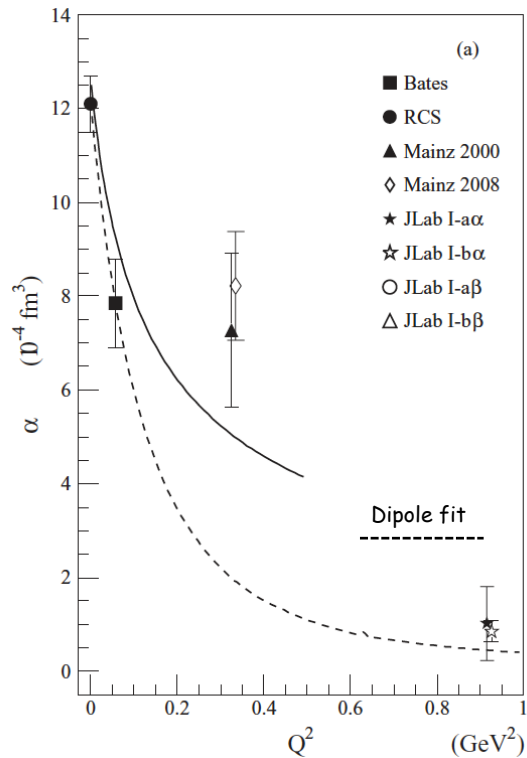


$$\frac{\langle r_{\beta}^2 \rangle}{-4.67^{+5.36}_{-13.04}}$$

# Polarizability radii

$$\langle r_{\alpha_E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2=0}$$

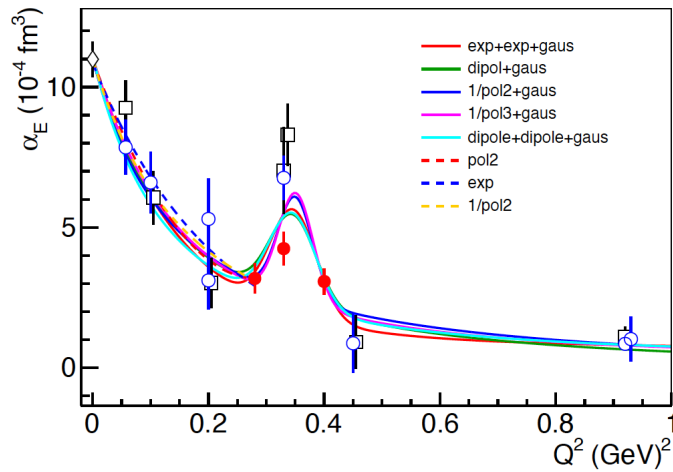
First extraction made possible with the MIT-Bates measurement ( $Q^2=0.06 \text{ GeV}^2$ )



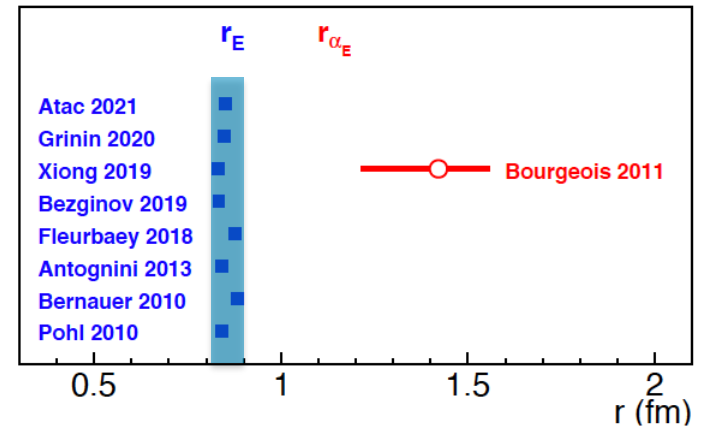
$$\langle r_{\alpha}^2 \rangle = 2.16 \pm 0.31 \text{ fm}^2$$

# Polarizability radii

$$\langle r_{\alpha_E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2=0}$$



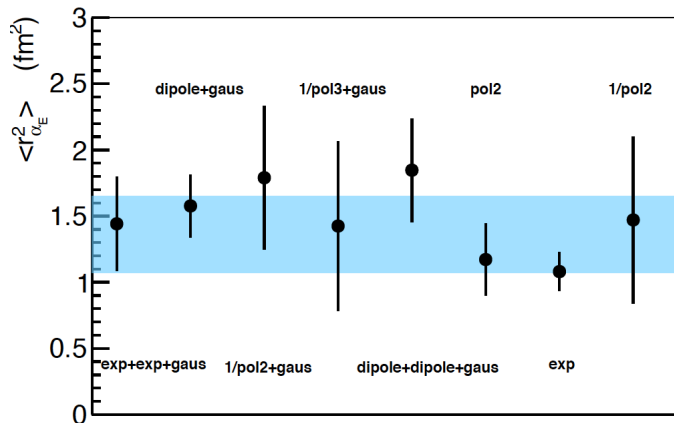
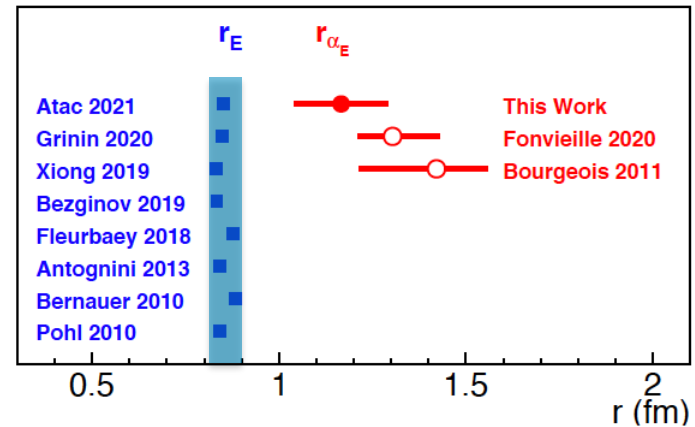
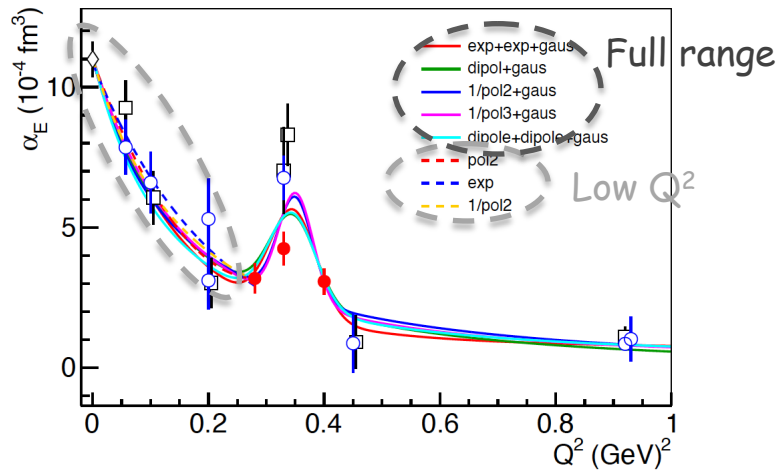
Since then: more data and more comprehensive treatment of the radius extraction



# Polarizability radii

$$\langle r_{\alpha_E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2=0}$$

Since then: more data and more comprehensive treatment of the radius extraction

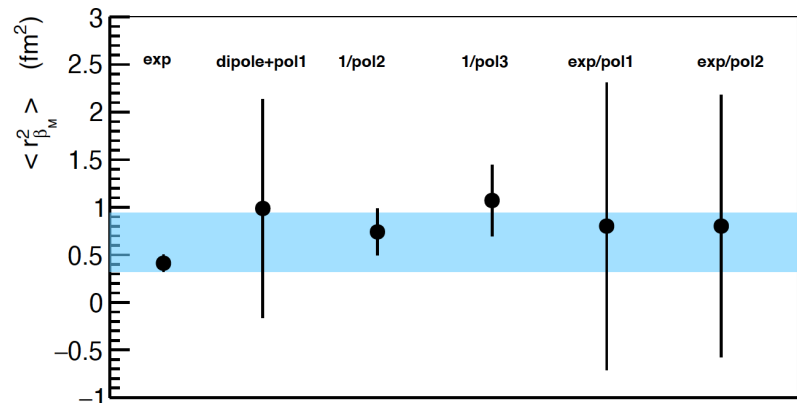
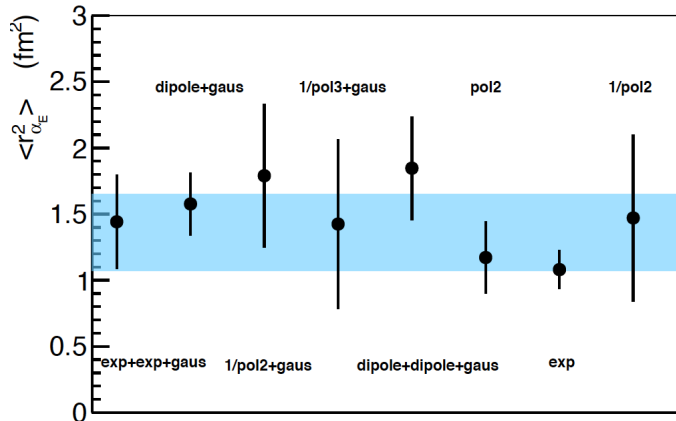
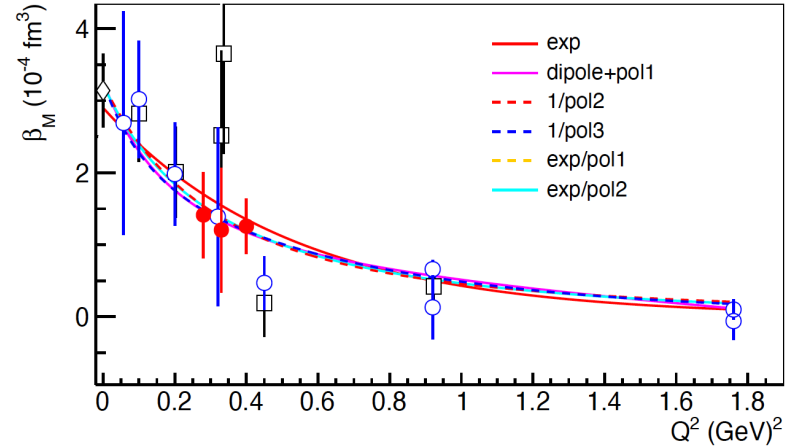
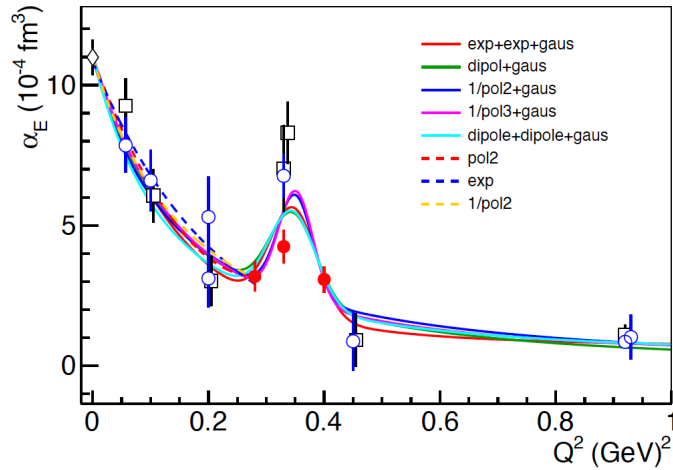


$$\langle r_{\alpha_E}^2 \rangle = 1.36 \pm 0.29 \text{ fm}^2$$

# Polarizability radii

$$\langle r_{\alpha_E}^2 \rangle = \frac{-6}{\alpha_E(0)} \cdot \frac{d}{dQ^2} \alpha_E(Q^2) \Big|_{Q^2=0}$$

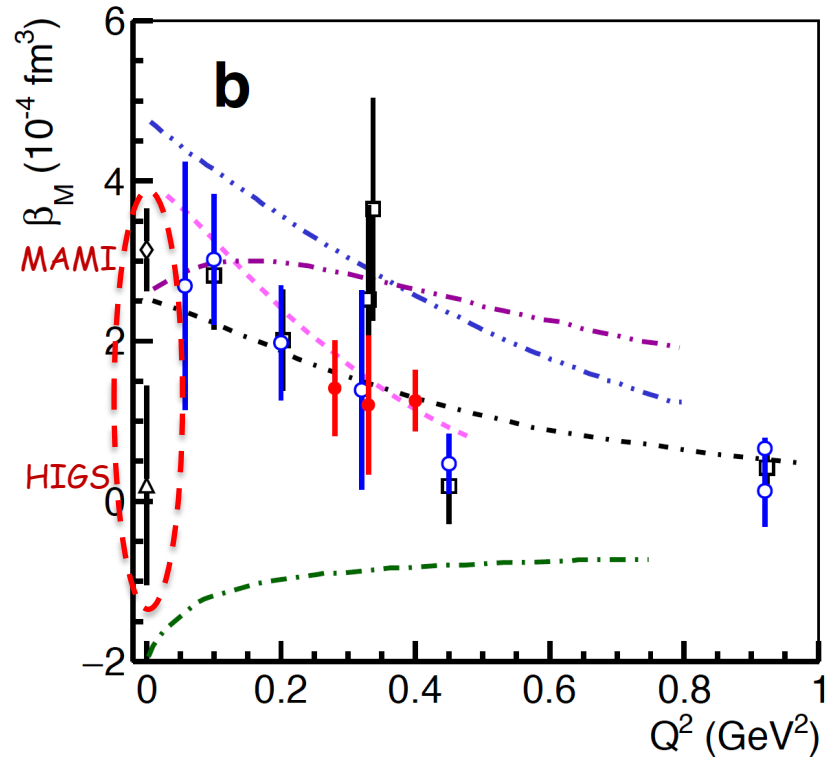
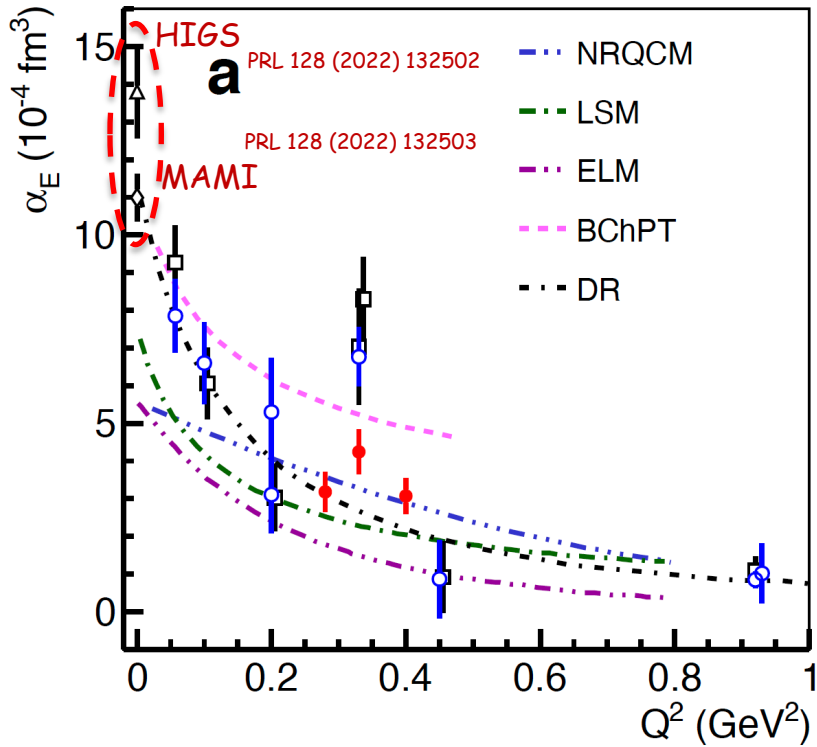
$$\langle r_{\beta_M}^2 \rangle = \frac{-6}{\beta_M(0)} \cdot \frac{d}{dQ^2} \beta_M(Q^2) \Big|_{Q^2=0}$$



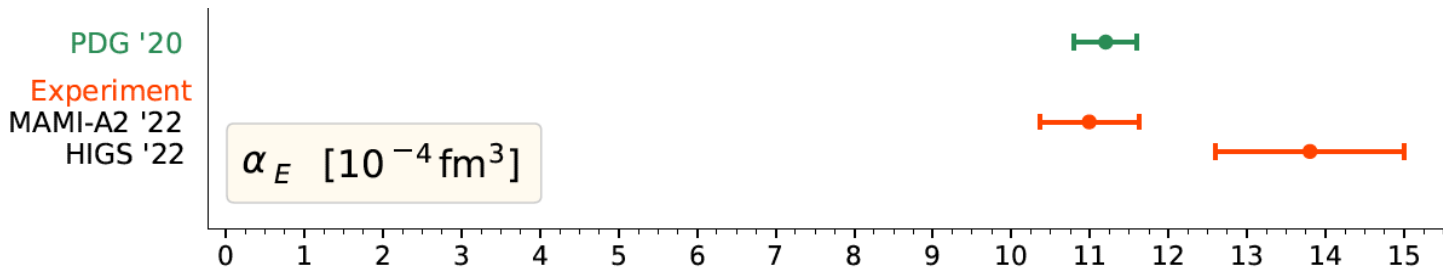
$$\langle r_{\alpha_E}^2 \rangle = 1.36 \pm 0.29 \text{ fm}^2$$

$$\langle r_{\beta_M}^2 \rangle = 0.63 \pm 0.31 \text{ fm}^2$$

# Static Polarizabilities

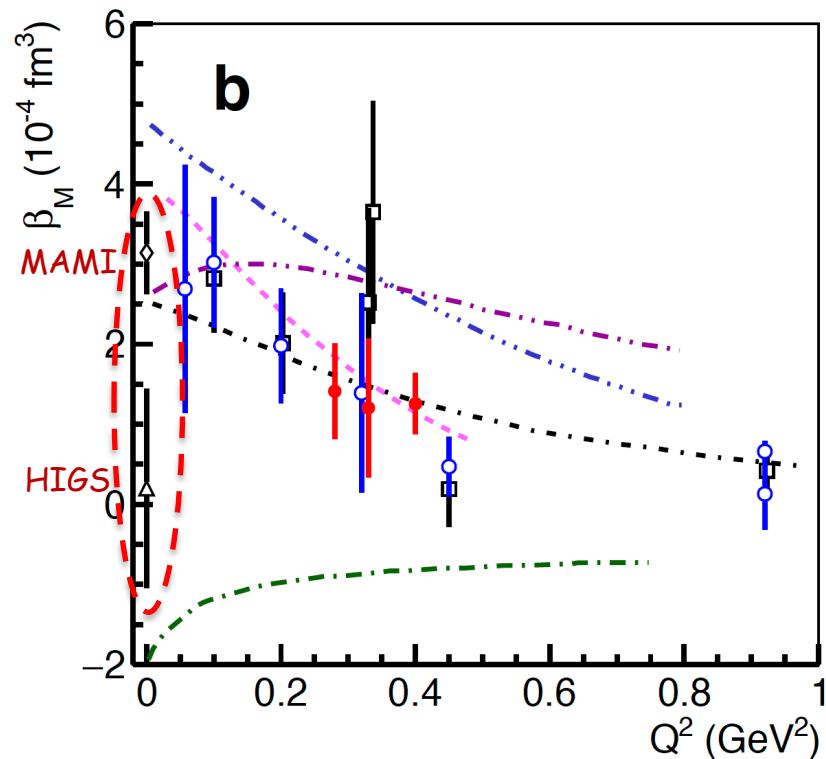
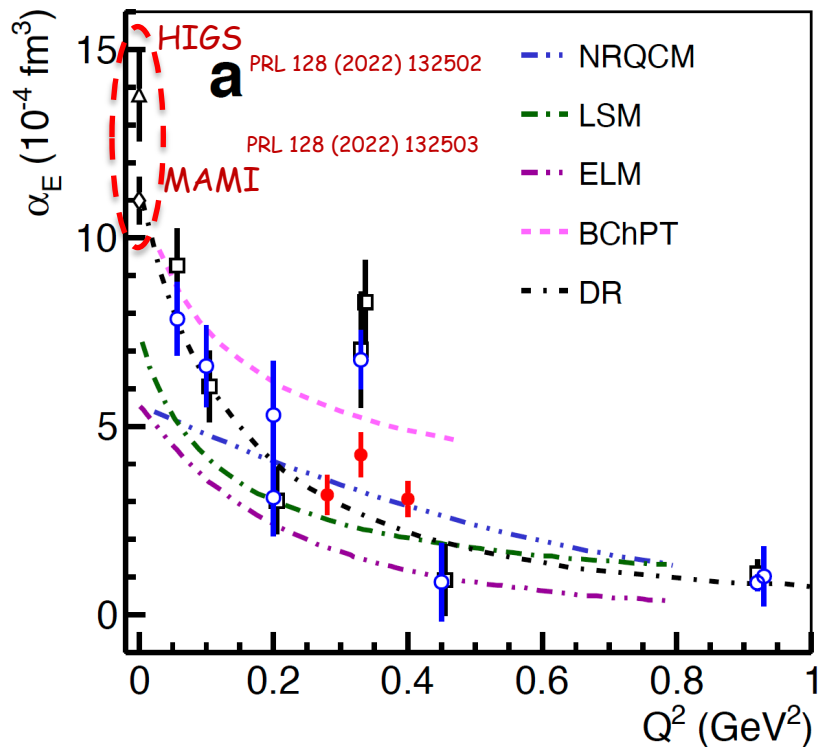


Recent measurements exhibit tension  $\rightarrow$  affects the pol. radius extraction



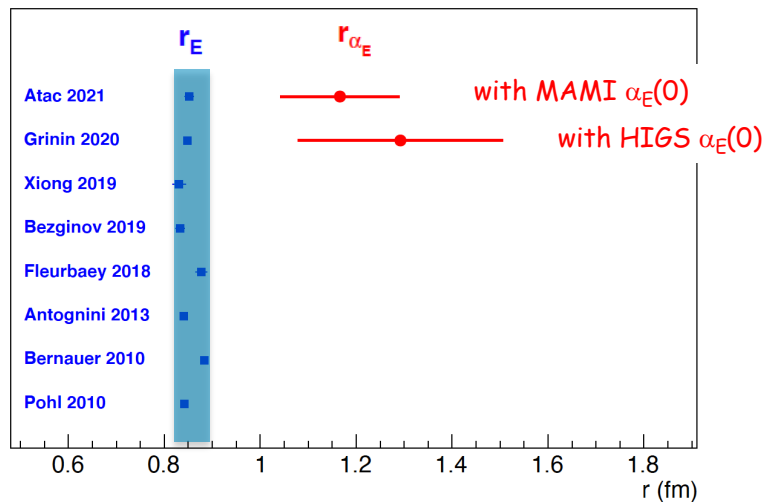


# Polarizability radii - Static Polarizabilities



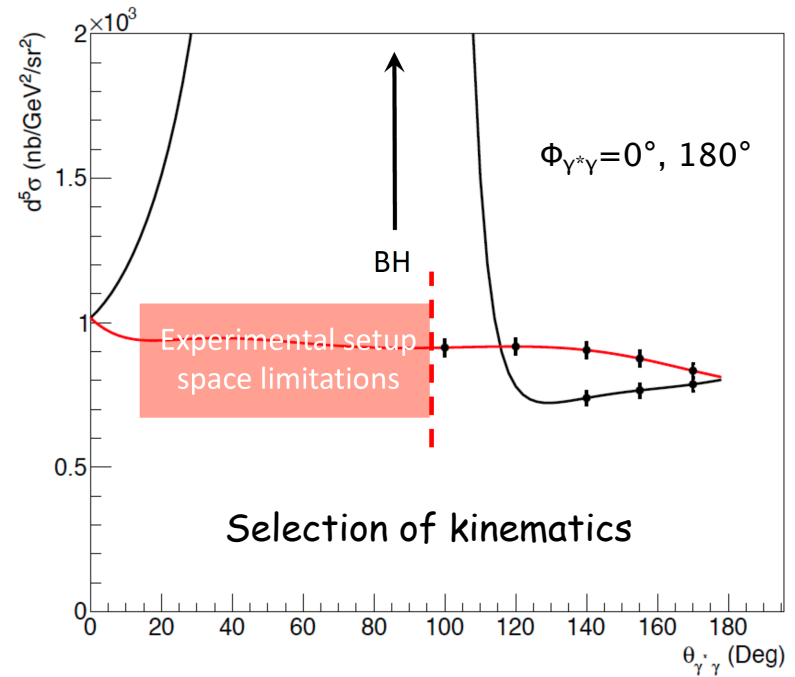
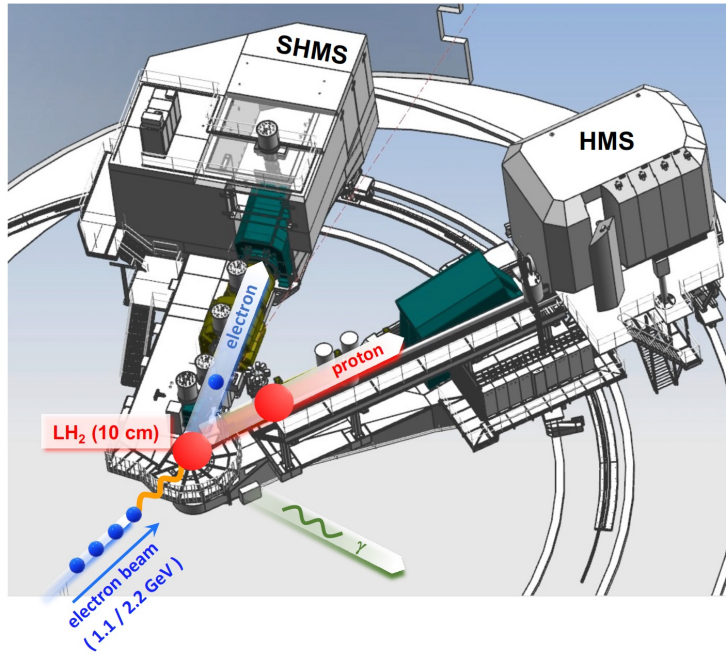
$$\langle r_{\alpha_E}^2 \rangle = 1.36 \pm 0.29 \text{ fm}^2$$

$$\langle r_{\alpha_E}^2 \rangle = 1.67 \pm 0.50 \text{ fm}^2$$



Moving Forward

# VCS-II (E12-23-001) @ JLab

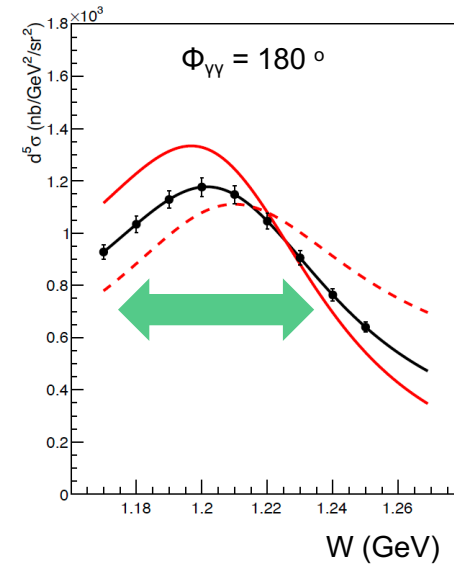
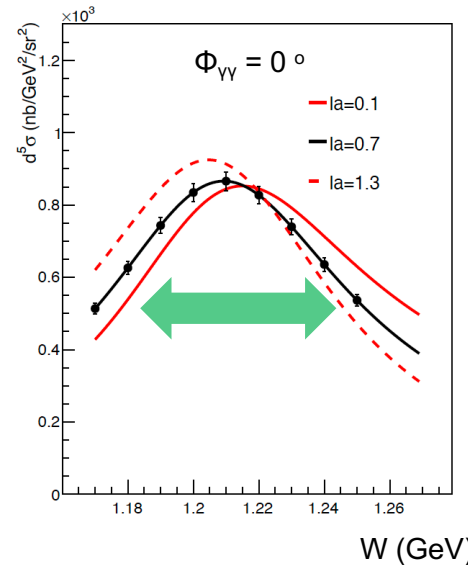


Extend  $Q^2$  range & targeted measurements to fully exploit the sensitivity to the EM GPs

**APPROVED  
PAC 51**

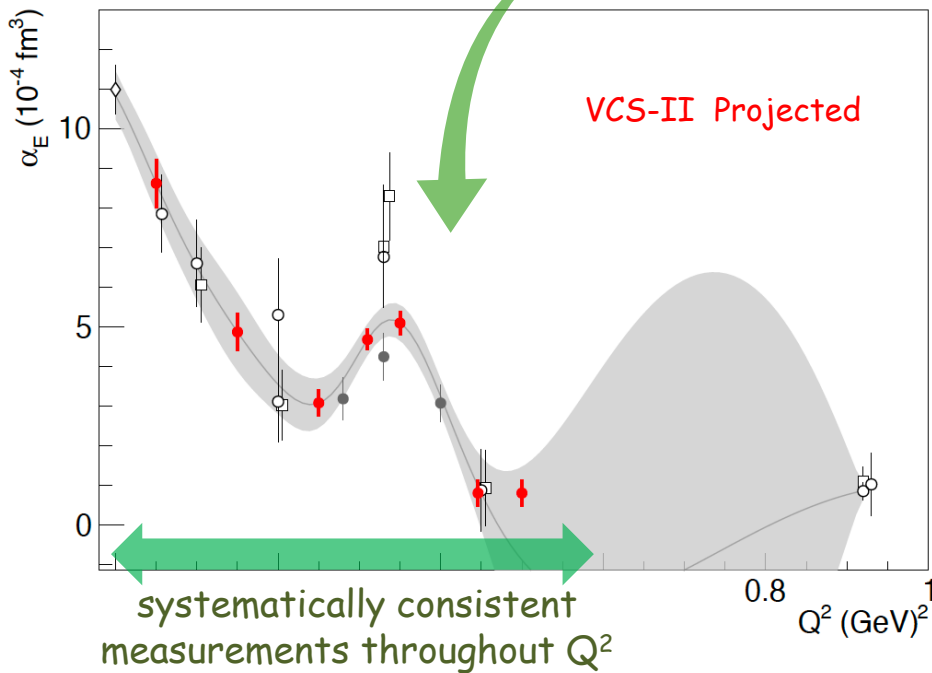
Production ( $E_0 = 1.1 \text{ GeV}$ ): 6 days  
 Production ( $E_0 = 2.2 \text{ GeV}$ ): 53 days  
 Studies (optics/dummy/calibrations): 3 days

**Total: 62 days**

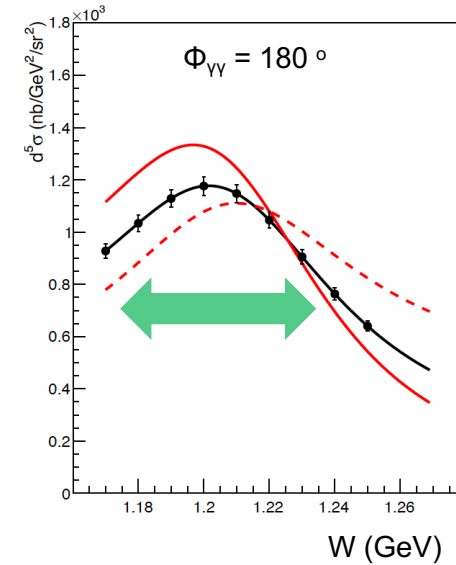
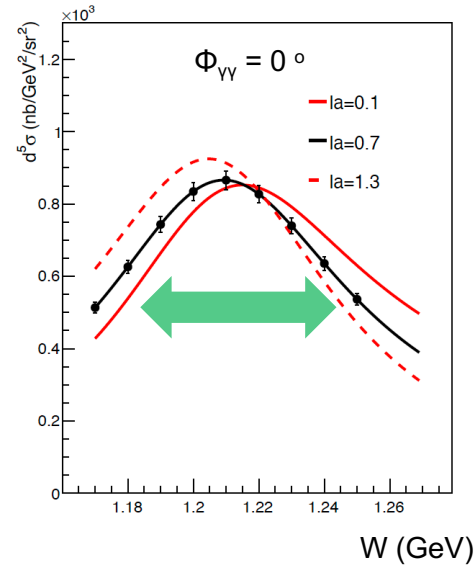


# VCS-II Projected Measurements

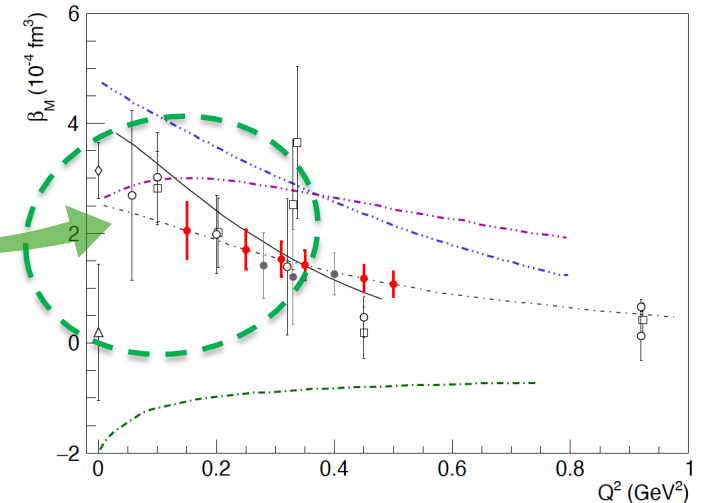
High precision measurements combined with a fine mapping in  $Q^2$



Targeted measurements to fully exploit the sensitivity to the GPs



Improve upon  $\beta_M$  :  
Pin down the competing para/dia-magnetic contributions in the nucleon



# Can we measure with a different method ?

Yes: positrons and/or beam spin asymmetries

Positrons allow for an independent path to access experimentally the GPs

Eur. Phys. J. A 57 (2021) 11, 316

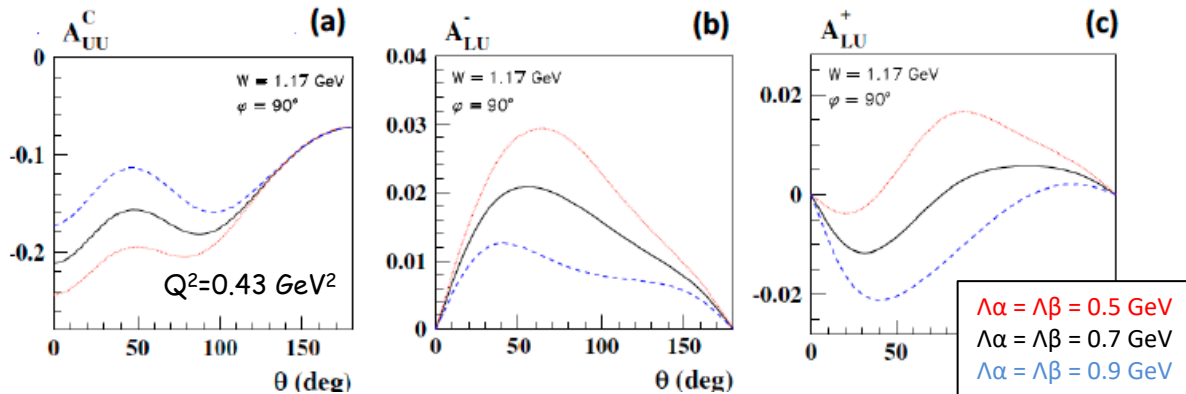
Virtual Compton scattering at low energies with a positron beam

Barbara Pasquini<sup>a,1,2</sup>, Marc Vanderhaeghen<sup>b,3</sup>

<sup>1</sup>Dipartimento di Fisica, Università degli Studi di Pavia, 27100 Pavia, Italy

<sup>2</sup>Istituto Nazionale di Fisica Nucleare, Sezione di Pavia, 27100 Pavia, Italy

<sup>3</sup>Institut für Kernphysik and PRISMA<sup>+</sup> Cluster of Excellence, Johannes Gutenberg Universität, D-55099 Mainz, Germany



(a): The beam-charge asymmetry as a function of the photon scattering angle at  $Q^2 = 0.43 \text{ GeV}^2$ .

(b) & (c): The electron and positron beam-spin asymmetry as a function of the photon scattering angle for out-of-plane kinematics.

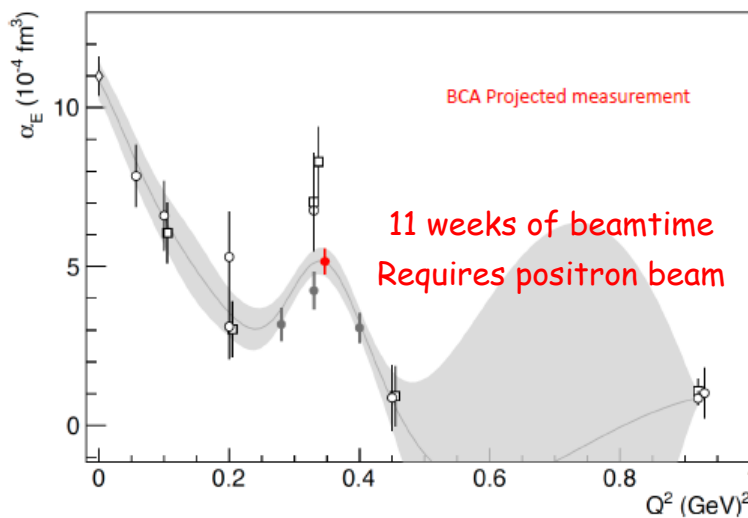
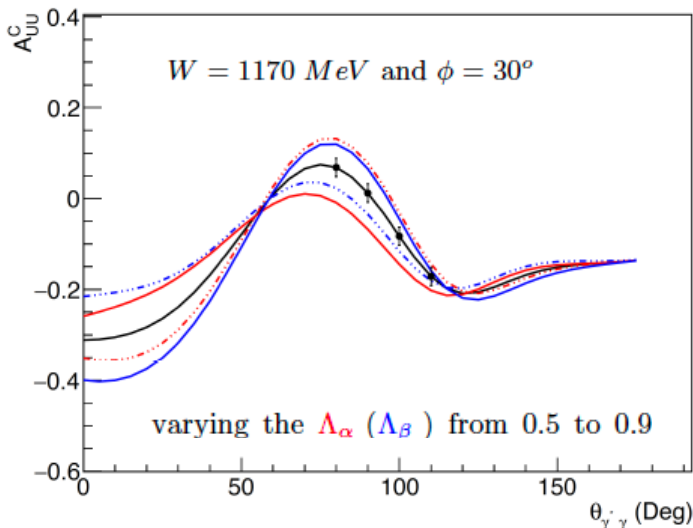
Unpolarized beam charge asymmetry (BCA):

$$A_{UU}^C = \frac{(d\sigma_+^+ + d\sigma_-^+) - (d\sigma_+^- + d\sigma_-^-)}{d\sigma_+^+ + d\sigma_-^+ + d\sigma_+^- + d\sigma_-^-}$$

Lepton beam spin asymmetry (BSA):

$$A_{LU}^e = \frac{d\sigma_+^e - d\sigma_-^e}{d\sigma_+^e + d\sigma_-^e}$$

## BCA (electrons & positrons)



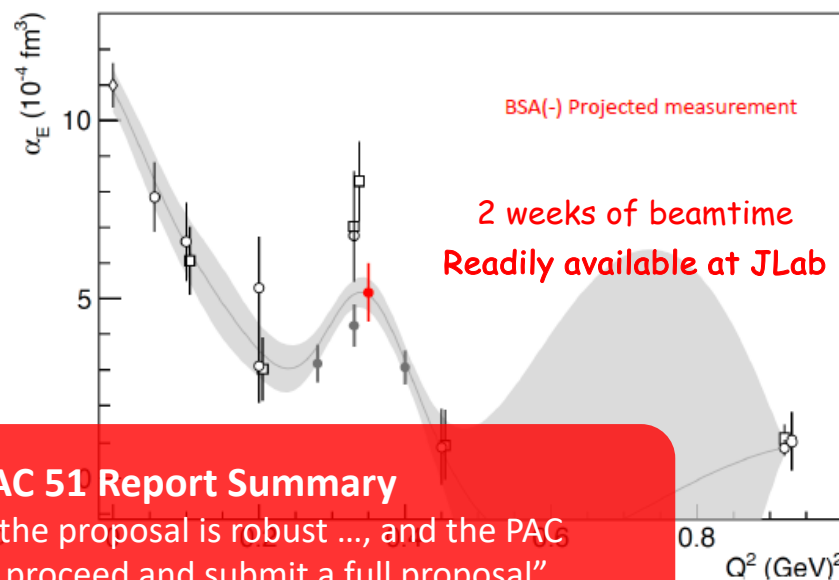
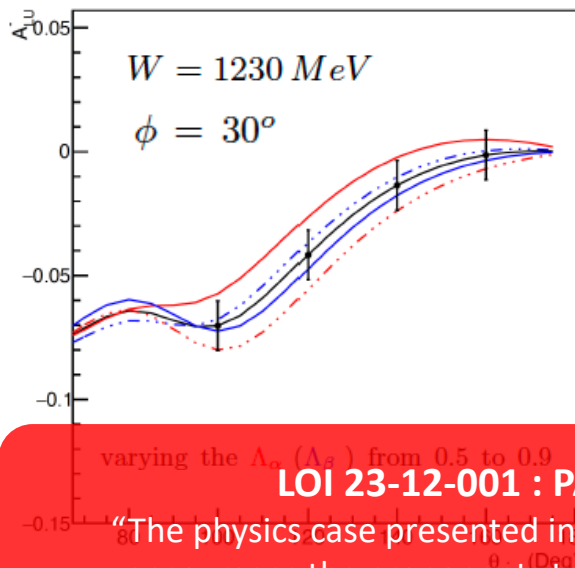
Hall C (SHMS / HMS)

$e^-$  : ~ 1 week @ 50  $\mu\text{A}$

and

$e^+$  : ~ 10 weeks @ 5  $\mu\text{A}$

## BSA (electrons or positrons)



$e^-$  (pol. 85% @ 70  $\mu\text{A}$ )

~ 2 weeks of beamtime

or

$e^+$  (pol. 60% @ 50 nA)

~ 3 orders of magnitude  
more beamtime

LOI 23-12-001 : PAC 51 Report Summary

"The physics case presented in the proposal is robust ..., and the PAC encourages the proponents to proceed and submit a full proposal"

# Summary

JLab: leading the efforts of the VCS program, past/ present / future

## Fundamental proton properties

Insight to spatial deformation of the nucleon densities under an applied EM field, interplay of para/dia-magnetism in the proton, polarizability radii, ...

Electric GP: { possibility for a non-trivial (non-monotonic) behavior in  $a_E(Q^2)$   
(albeit with a smaller magnitude than originally suggested)  
or  
at minimum: strong tension between world data

## Experiment is ahead of theory

Stringent constraints to theoretical predictions

High precision benchmark data for upcoming LQCD calculations

## Future experimental goals

Improve  $\beta_M$

Identify the shape of the  $a_E$  structure (if it exists)

pin it down with precision - important input for the theory

Conduct independent cross-check

Measure via a different channel (BS asymmetries & positrons)

Thank you!