# Possible evidences from  $H(z)$  parameter data for physics beyond ΛCDM

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# **Motivation**

- The Universe is experiencing an accelerated expansion.
- Hubble tension may also provide evidences for physics beyond ΛCDM.
- Statistical methods, such as the maximum likelihood, are generally used to analyze the observational data to fit the parameters.
- These statistical method yields the best statistical results, but it is easy to eliminate some interesting (possibly important) data.
- Here we propose a model-independent method by using the Lagrange mean value theorem to analyze  $H(z)$  parameter data.

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consider a spatially flat FRWL spacetime here, the Friedmann equations take the form

$$
H2 = \frac{8\pi G}{3}\rho,
$$
  
\n
$$
\frac{\ddot{a}}{a} = -\frac{4\pi G}{3}(\rho + 3\rho),
$$
\n(1)

or equivalently

<span id="page-4-0"></span>
$$
\dot{H} = -4\pi G(\rho + \rho),\tag{3}
$$

where the  $H = \frac{a}{a}$  is the Hubble parameter with the dot denoting the derivative with respect to the cosmic time  $t$ . The total energy density  $\rho$  and pressure  $\rho$  contain contributions coming from the radiation, nonrelativistic matter, and other components.

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Because  $dz = -(1+z)Hdt$ , we have

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$$
\dot{H} = -(1+z)H\frac{dH}{dz}.
$$
\n(4)

Combining Eqs. [\(3\)](#page-4-0) and [\(4\)](#page-5-0), yields

<span id="page-5-1"></span>
$$
\frac{dH}{dz} = \frac{4\pi G}{(1+z)H}(\rho + \rho) = \frac{4\pi G\rho(1 + w_t)}{(1+z)H},
$$
(5)

where  $w_{\rm t}$  is the total EoS. From this equation, we can judge whether the total EoS is greater than, equal to, or less than  $-1$ : see for example, if  $dH/dz < 0$ , we have  $w_x < w_t < -1$  because of the positive of H and  $\rho$ . In an era dominated by dark energy, we can also determine with Eq. [\(5\)](#page-5-1) wether the EoS of dark energy is equal to  $-1$ : if  $dH/dz = 0$ , then one have  $w_x \simeq w_t = -1$ . If  $dH/dz \le 0$ , we know the Universe is experiencing an accelerated expansion.

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But if  $dH/dz > 0$ , we can't judge whether the Universe speeds up. At this point, we need another important physical quantity, the deceleration parameter, which is defined as

<span id="page-6-0"></span>
$$
q=-1+(1+z)\frac{1}{H}\frac{dH}{dz}.
$$
\n(6)

Now a question naturally rise: if we have some  $H(z)$  parameter data, how can we use them to directly determine  $dH/dz$  or q?

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Think of Lagrange mean value theorem in Calculus, which states: for a continuous and differentiable function  $f(x)$ , there exists  $x_1 < x_1$ <sub>2</sub>  $x_2$  satisfying

<span id="page-7-0"></span>
$$
\frac{df}{dx}\big|_{x=x_{12}} = \frac{f(x_2) - f(x_1)}{x_2 - x_1}.\tag{7}
$$

Applying this theorem to Hubble parameter which we assume is a continuous and differentiable function of z, and taking function  $H(z)$  as  $f(x)$  in [\(7\)](#page-7-0), we have

$$
H'(z_{ij}) \equiv \frac{dH}{dz}\Big|_{z=z_{ij}} = \frac{H(z_i) - H(z_j)}{z_i - z_j},
$$
 (8)

where  $z_j < z_{ij} < z_i$ . Combining this equation and Hubble parameter data, namely, taking the  $H(z)$  data from the table [14](#page-12-0) for  $H(z_i)$  and  $H(z_j)$  and the corresponding z data for  $z_i$  and  $z_j$ , we can obtain a lot of data for  $H'(z_{ij})$ . (ロ) (母) (ヨ) (ヨ) ニヨ

Since the difference between different data of Hubble parameters  $H(z)$  in table [14](#page-12-0) in general is large,  $H'(z_{ij})$  will be larger if  $z_i - z_j \ll 1$ , which will make relevant results less credible. In order to make the results credible, we will impose restrictions on  $H(z_i) - H(z_i)$  and  $z_i - z_i$  during the process of data analysis. When applying equation [\(6\)](#page-6-0) to analyze the data in table [14,](#page-12-0)  $z_{ii}$  and  $H(z_{ii})$ are unknown, we take approximatively:  $z_{ii} = (z_i + z_i)/2$  and  $H(z_{ii}) \simeq [H(z_i) + H(z_i)]/2$ , which can be called as mid-value approximate method.

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if the difference between  $H(z_i)$  and  $H(z_i)$  and the difference between  $z_i$  and  $z_i$  are reasonable, this approximate method in general is credible. Then we have

$$
q(z_{ij}) \simeq -1 + \frac{(2+z_i+z_j) H'(z_{ij})}{H(z_i) + H(z_j)}.
$$
\n(9)

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The uncertainties associated to the  $H'(z)$  and  $q(z)$  are given by, respectively

$$
\sigma_{H'} = \frac{\sqrt{\sigma_{H_i}^2 + \sigma_{H_j}^2}}{z_i - z_j},
$$
\n(10)

and

$$
\sigma_q = \frac{2(2+z_i+z_j)H_i}{(H_i+H_j)^2} \frac{\sqrt{\sigma_{H_i}^2 + \sigma_{H_j}^2}}{z_i - z_j}.
$$
 (11)

With  $\sigma_{H'}$  and  $\sigma_q$ , we can determine whether the results are credible at 1  $\sigma$  confidence level.

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TABLE I: Hubble parameter compilation from cosmic chronometers (DA) or from the radial BAO surveys (clustering).

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TABLE II:  $H'(z)$  and  $q(z)$  data obtained from  $H(z)$  parameter data.

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- At redshifts z61, z71, z91, z131, z141, z171, z181, z201, z251, z271, z311, z139, z179, z189, z209, z259, z279, z319, z339, z2014, z2714, z359, z409, z419, z4324, z5150, z5550, z5650, z5552, z<sup>5652</sup> z5554,  $z_{5654}$ , and  $z_{5854}$ , the Universe experiences an accelerated expansion at  $1 \sigma$  confidence level.
- At redshifts  $z_{3418}$ ,  $z_{3420}$ ,  $z_{3718}$ ,  $z_{3431}$ ,  $z_{3727}$ , and  $z_{3731}$ , the Universe experiences an decelerated expansion at 1  $\sigma$ confidence level.
- At redshifts  $z_{139}$ ,  $z_{2014}$ ,  $z_{5150}$ ,  $z_{5550}$ ,  $z_{5650}$ ,  $z_{5552}$ , and  $z_{5854}$ , since  $H'(z) < 0$ , implying  $w_{\rm x} \leq w_{\rm t} < -1$ , but not at at 1  $\sigma$ confidence level. However, we can infer that  $w_x \leq w_t \leq -1$  at redshifts  $z_{5652}$ ,  $z_{5554}$ , and  $z_{5654}$  at at 1  $\sigma$  confidence level.

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- There exists decelerated phase at 1  $\sigma$  confidence level in the redshift range (0.38,0.59);
- The EoS of dark energy may be less than  $-1$  at 1  $\sigma$  confidence level at some redshifts in the redshift range (1.3,1.53);
- There exist accelerated phase at 1  $\sigma$  confidence level in the redshift range (1.037,1.944).
- These results suggest that the dark energy maybe dynamic with EoS crossing −1 and the Universe may be accelerated first, then decelerated, and accelerated again recently.

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# Thanks!

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