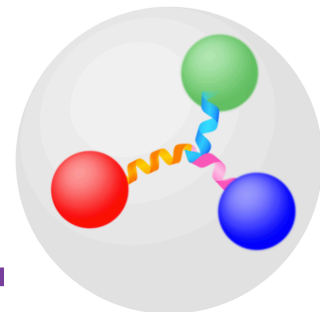




Workshop on Hyperon Physics,  
Apr. 12–15, 2024, Huizhou, IMP



# 夸克平均场对超核以及超子物质的研究

Jinniu Hu

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J.H, A. Li, H. Shen, and H. Toki, *Prog. Theor. Exp. Phys.* 2014 (2014) 013D02

J.H, A. Li, H. Toki, and W. Zuo, *Phys. Rev. C* 89 (2014) 025802

X. Xing, J.H., and H. Shen, *Phys. Rev. C* 94 (2016) 044308

X. Xing, J.H., and H. Shen, *Phys. Rev. C* 95 (2017) 054310

J.H. and H. Shen, *Phys. Rev. C* 96 (2017) 054304

L. Wu, J.H., and H. Shen, *Phys. Rev. C* 101 (2020) 024303

A. Li, Z. Zhen, E. Zhou, J. Dong, J.H., C. Xia, *JHEAp* 28(2020)19

J. H, Y. Zhang, and H. Shen, *J. Phys. G* 49 (2022)025104



# Outline

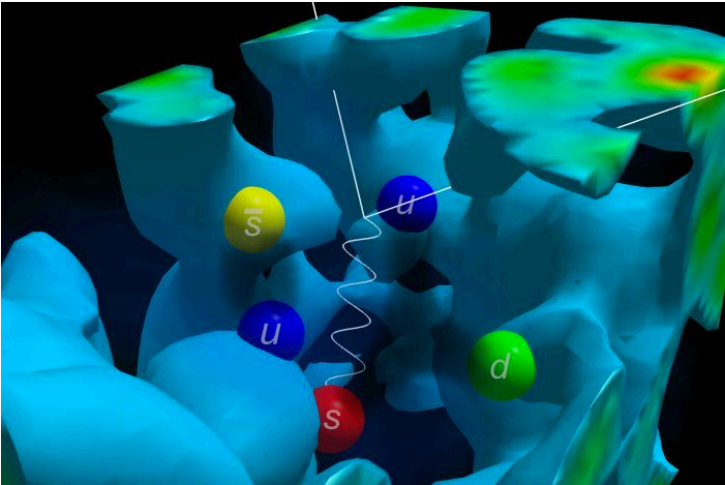
1 Introduction

2 The quark mean field model

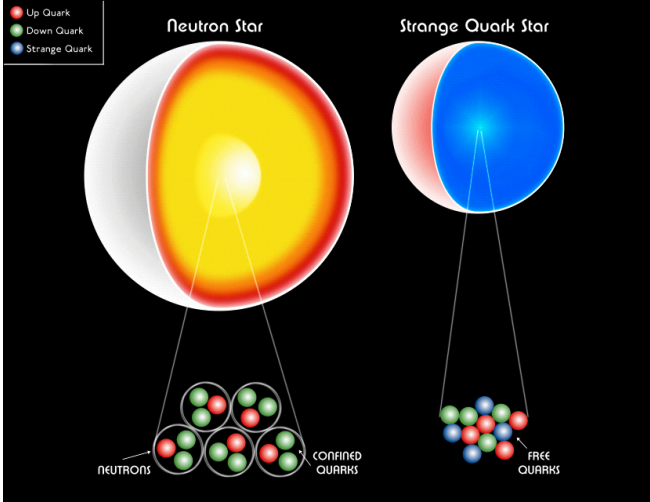
3 The strangeness with QMF model

4 The summary and perspective

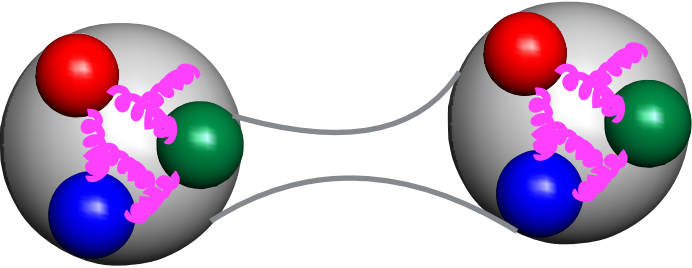
## Hadrons



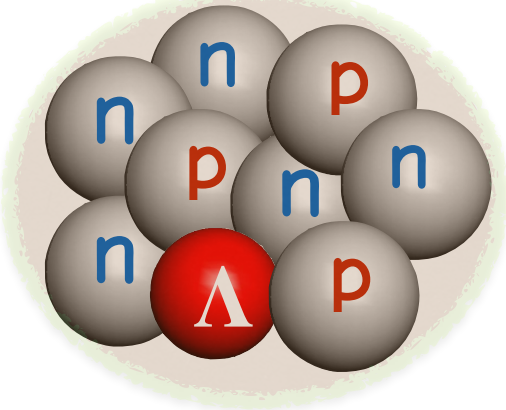
## Compact star



## Baryon-Baryon force

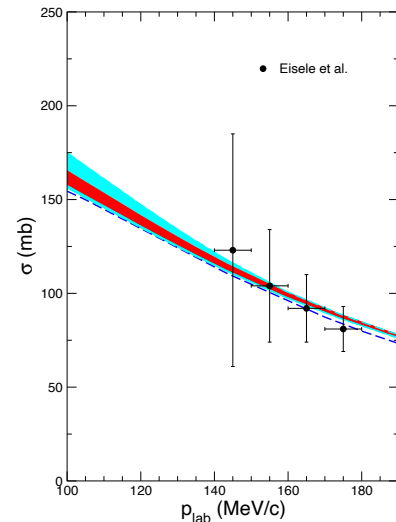
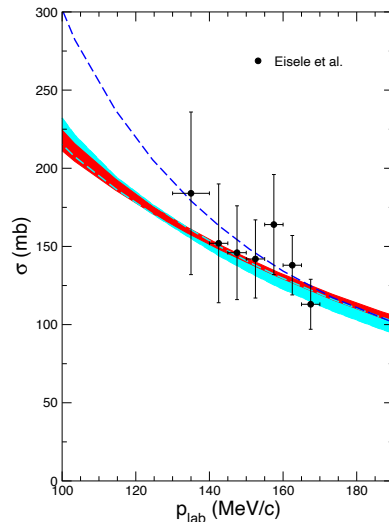
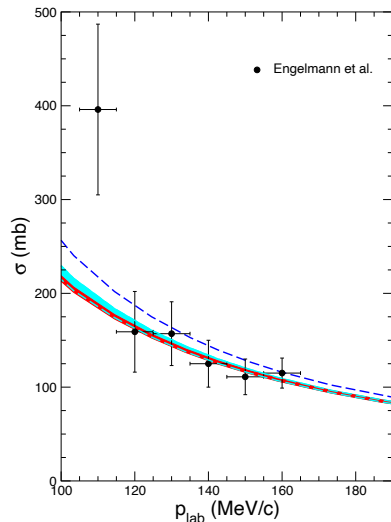
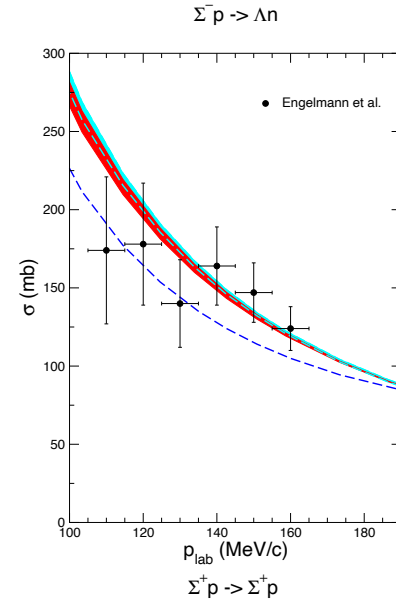
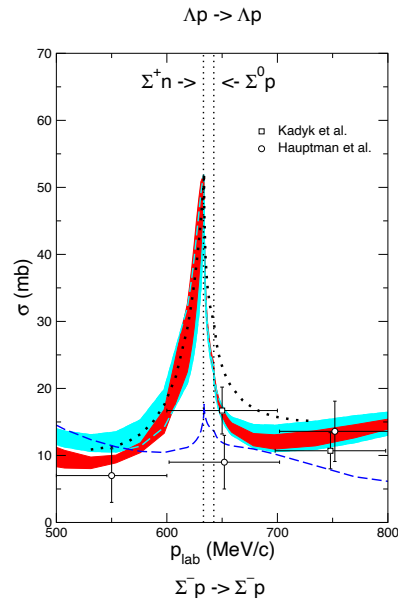
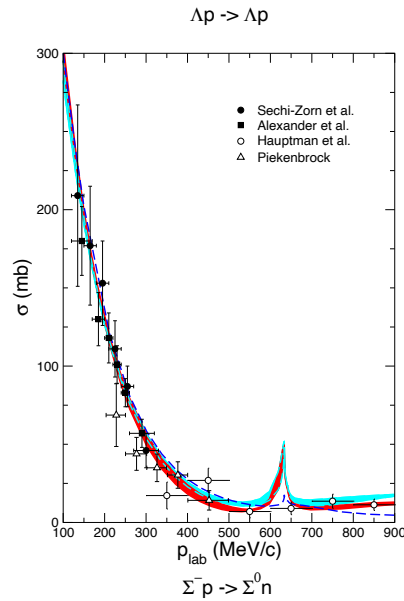


## Hypernuclei



A. Gal, E. V. Hungerford, and D. J. Millener, Rev. Mod. Phys. 88(2016)035004

# Baryon-Baryon scattering data



J. Haidenbauer, Ulf-G. Meissner, and A. Nogga, Eur. Phys J. A56, 91 (2020)

➤ **Juelich Meson exchange exchange potential (Jul94,05)**

J. Haidenbauer, Ulf-G. Meissner, Phys.Rev. C72, 044005 (2005)

➤ **Nijmegen potentials: OBE and ESC (ESC06, ESC12, ESC16)**

Th. A. Rijken and Y. Yamamoto, Phys. Rev. C 73, 044008 (2006)

➤ **Quark-Cluster models**

Y. Fujiwara, Y. Suzuki, and C. Nakamoto, Prog. Part. Nucl Phys. 58, 439 (2007)

➤ **Chiral effective potential (NLO13, NLO19)**

J. Haidenbauer, Ulf-G. Meissner, and A. Nogga, Eur. Phys J. A56, 91 (2020)

➤ **Covariant Chiral effective potential**

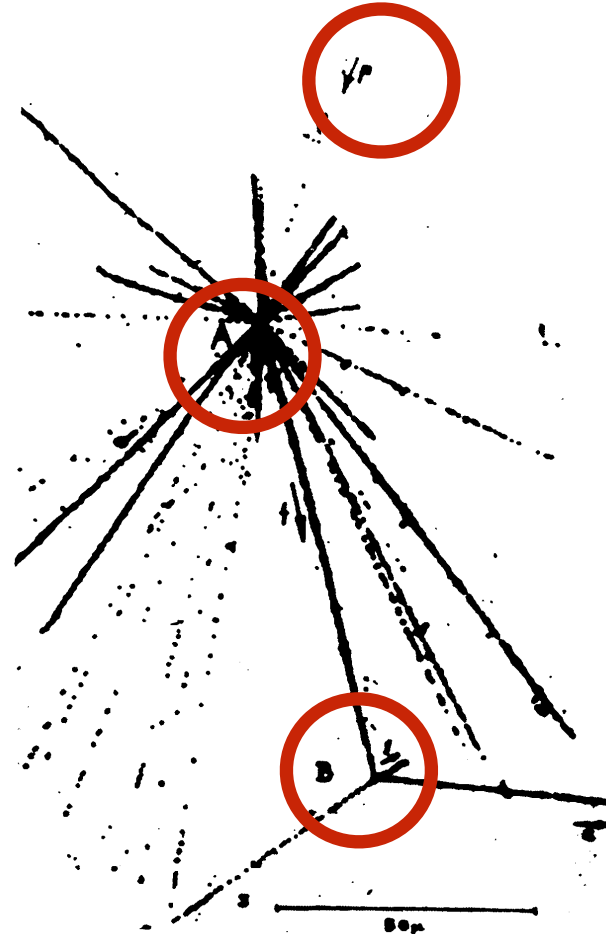
J. Song, Z. Liu, K. Li, and L. Geng, Phys. Rev. C 105, 035203 (2022)

➤ **Lattice QCD potential**

H. Nemura, arXiv: 1810.04046

# Discovery of first $\Lambda$ -hypernuclei

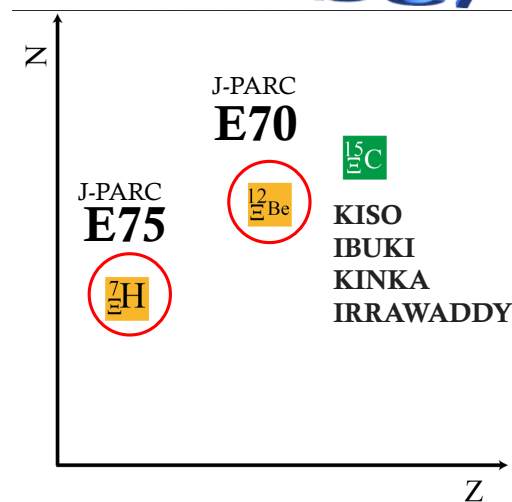
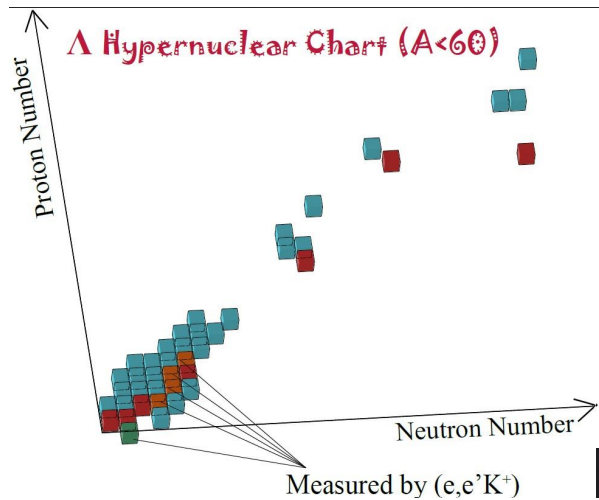
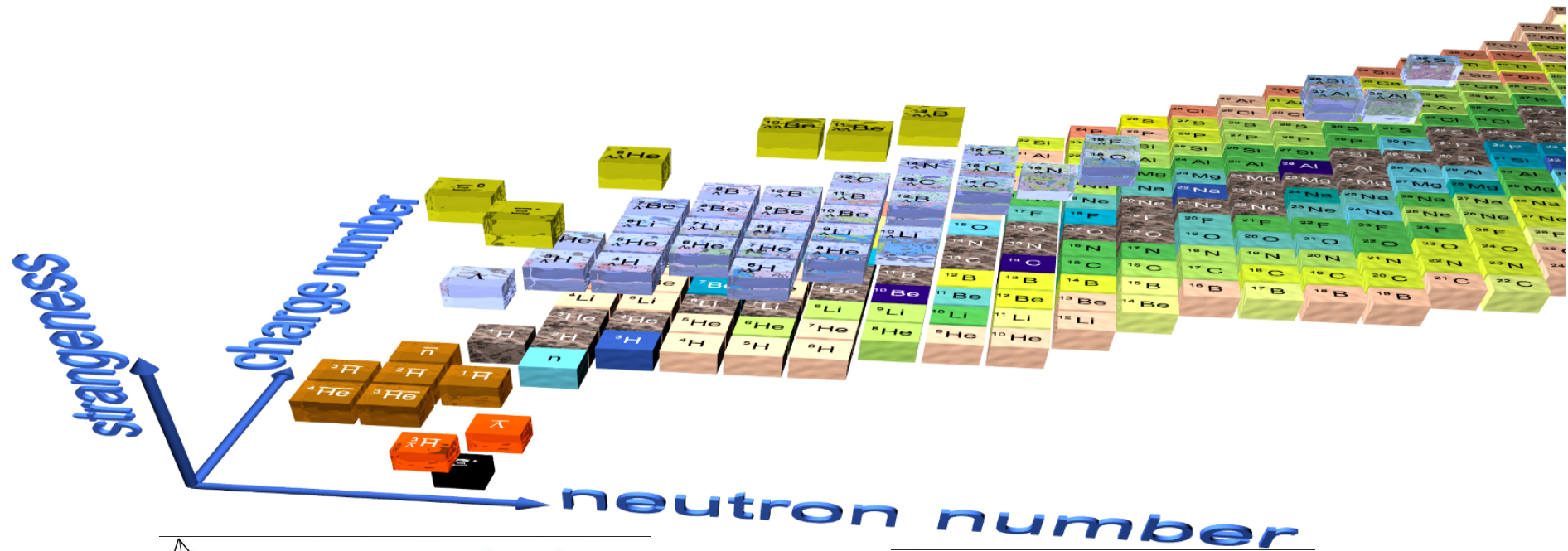
Danzysz and Pniewski discovered the first hypernucleus in Warsaw in September 1952. The hypernucleus was created when a high-energy proton interacted with a nucleus in the emulsions they were using as a detector, producing a hyperfragment.



# Hypernuclei chart



南开大学



Hypernucleus	Number of events	$B_{\Lambda} \pm \Delta B_{\Lambda}$ (MeV)
${}^3_{\Lambda}\text{H}$	204	$0.13 \pm 0.05$
${}^4_{\Lambda}\text{H}$	155	$2.04 \pm 0.04$
${}^4_{\Lambda}\text{He}$	279	$2.39 \pm 0.03$
${}^5_{\Lambda}\text{He}$	1784	$3.12 \pm 0.02$
${}^6_{\Lambda}\text{He}$	31	$4.18 \pm 0.10$
${}^7_{\Lambda}\text{He}$	16	Not averaged
${}^7_{\Lambda}\text{Li}$	226	$5.58 \pm 0.03$
${}^7_{\Lambda}\text{Be}$	35	$5.16 \pm 0.08$
${}^8_{\Lambda}\text{He}$	6	$7.16 \pm 0.70$
${}^8_{\Lambda}\text{Li}$	787	$6.80 \pm 0.03$
${}^8_{\Lambda}\text{Be}$	68	$6.84 \pm 0.05$
${}^9_{\Lambda}\text{Li}$	8	$8.50 \pm 0.12$
${}^9_{\Lambda}\text{Be}$	222	$6.71 \pm 0.04$
${}^9_{\Lambda}\text{B}$	4	$8.29 \pm 0.18$
${}^{10}_{\Lambda}\text{Be}$	3	$9.11 \pm 0.22$
${}^{10}_{\Lambda}\text{B}$	10	$8.89 \pm 0.12$
${}^{11}_{\Lambda}\text{B}$	73	$10.24 \pm 0.05$
${}^{12}_{\Lambda}\text{B}$	87	$11.37 \pm 0.06$
${}^{12}_{\Lambda}\text{C}$	6	$10.76 \pm 0.19$
${}^{13}_{\Lambda}\text{C}$	6	$11.69 \pm 0.12$
${}^{14}_{\Lambda}\text{C}$	3	$12.17 \pm 0.33$

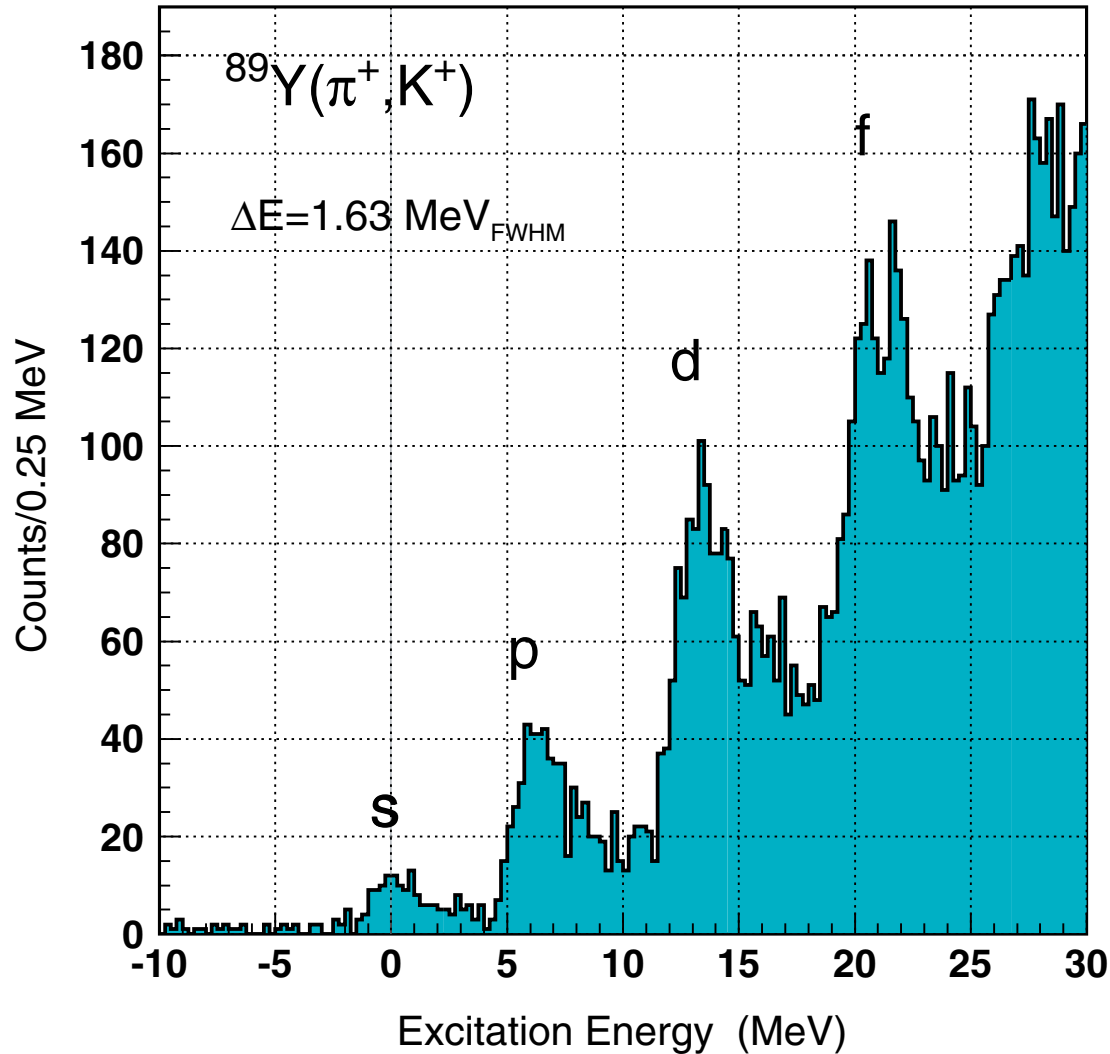
Hypernucleus	$s_{\Lambda}$	$p_{\Lambda}$	$d_{\Lambda}$	$f_{\Lambda}$	$g_{\Lambda}$
			$(\pi^+, K^+)$		
${}^{208}_{\Lambda}\text{Pb}$	26.9(8)	22.5(6)	17.4(7)	12.3(6)	7.2(6)
${}^{139}_{\Lambda}\text{La}$	25.1(12)	21.0(6)	14.9(6)	8.6(6)	2.1(6)
${}^{89}_{\Lambda}\text{Y}$	23.6(5)	17.7(6)	10.9(6)	3.7(6)	-3.8(10)
${}^{51}_{\Lambda}\text{V}$	21.5(6)	13.4(6)	5.1(6)		
${}^{28}_{\Lambda}\text{Si}$	17.2(2)	7.6(2)	-1.0(5)		
${}^{16}_{\Lambda}\text{O}$	13.0(2)	2.5(2)			
${}^{13}_{\Lambda}\text{C}$	12.0(2)	1.1(2)			
${}^{12}_{\Lambda}\text{C}$	11.36(20)	0.36(20)			
${}^{10}_{\Lambda}\text{B}$	8.7(3)				
			$(e, e'K^+)$		
${}^{52}_{\Lambda}\text{V}$	21.8(3)				
${}^{16}_{\Lambda}\text{N}$	13.76(16)	2.84(18)			
${}^{12}_{\Lambda}\text{B}$	11.52(2)	0.54(4)			
${}^{10}_{\Lambda}\text{Be}$	8.55(13)				
${}^7_{\Lambda}\text{He}$	5.55(15)				
			Emulsion		
${}^{13}_{\Lambda}\text{C}$	11.69(12)	0.8(3)			
${}^{12}_{\Lambda}\text{B}$	11.37(6)				
${}^{12}_{\Lambda}\text{C}$		0.14(5)			
${}^8_{\Lambda}\text{Li}$	6.80(3)				
${}^7_{\Lambda}\text{Be}$	5.16(8)				
			$(K^-, \pi^-)$		
${}^{40}_{\Lambda}\text{Ca}$		11.0(5)	1.0(5)		
${}^{32}_{\Lambda}\text{S}$	17.5(5)	8.2(5)	-1.0(5)		

A. Gal, E. V. Hungerford, and D. J. Millener, Rev. Mod. Phys. 88(2016)035004

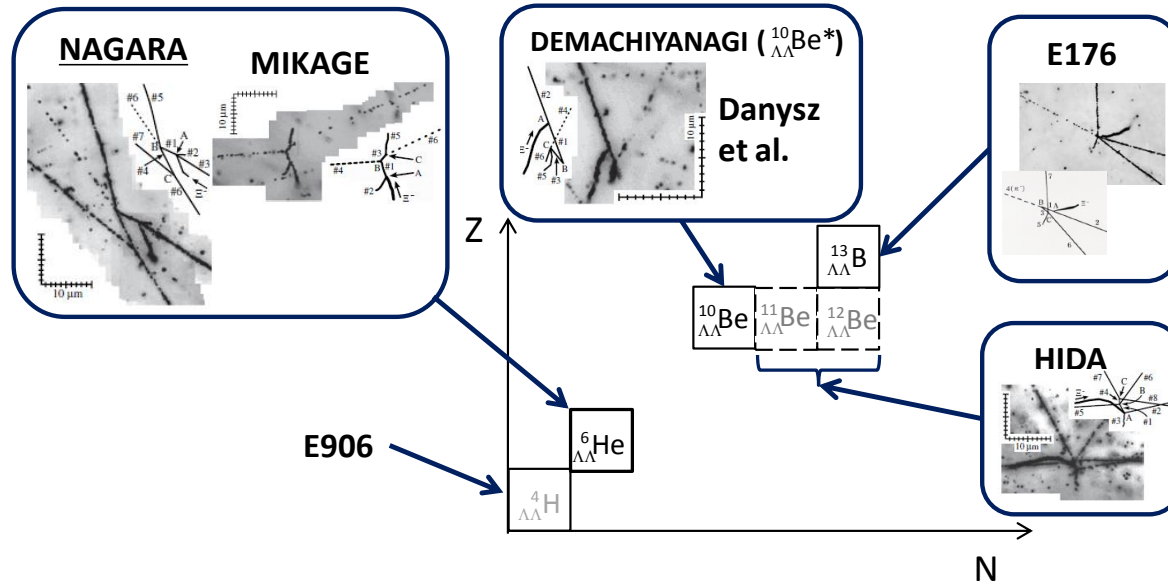


# Single $\Lambda$ potential

A. Gal, E. V. Hungerford, and D. J. Millener, Rev. Mod. Phys. 88(2016)035004



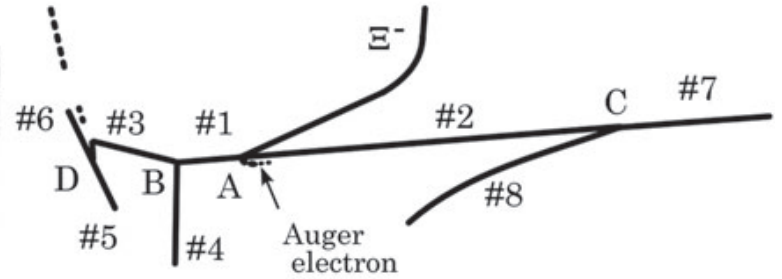
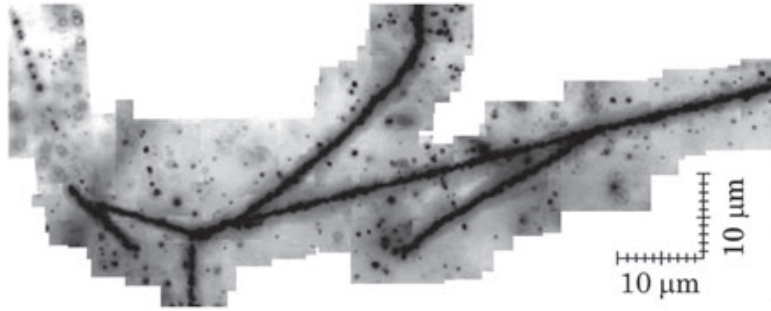
## $\Lambda\Lambda$ -Hypernuclear Chart



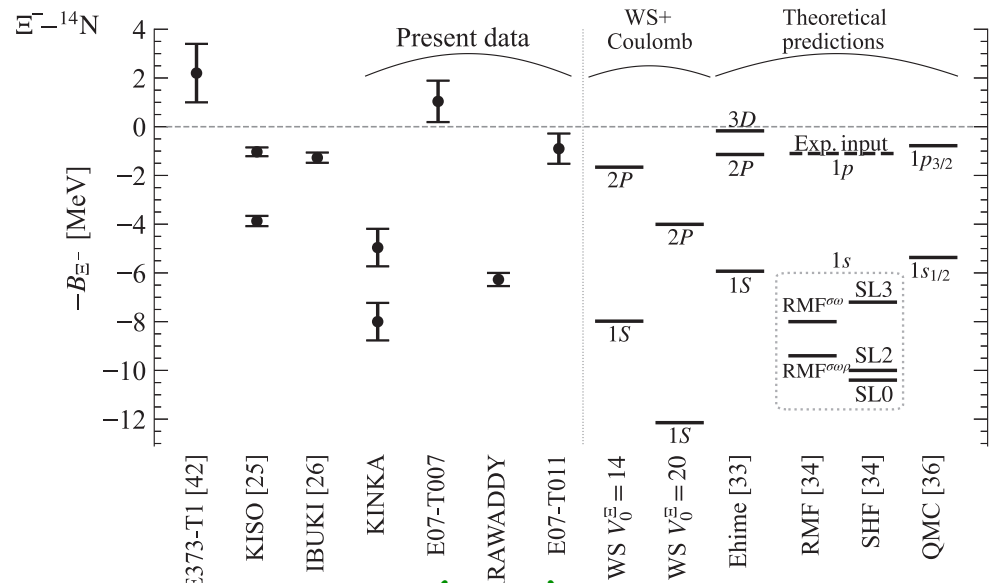
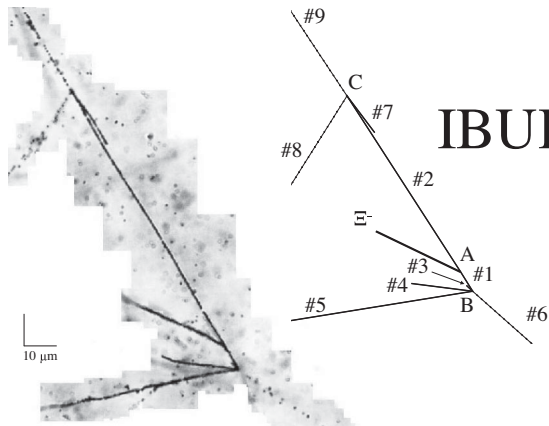
Event	${}_{\Lambda\Lambda}^AZ$	$\bar{B}_{\Lambda}(A-1Z)$	$B_{\Lambda\Lambda}^{\text{exp}}$	$B_{\Lambda\Lambda}^{\text{CM}}$	$B_{\Lambda\Lambda}^{\text{SM}}$
E373-Nagara	${}_{\Lambda\Lambda}^6\text{He}$	$3.12 \pm 0.02$	$6.91 \pm 0.16$	$6.91 \pm 0.16$	$6.91 \pm 0.16$
E373-Dem Yan	${}_{\Lambda\Lambda}^{10}\text{Be}$	$6.71 \pm 0.04$	$14.94 \pm 0.13$	$14.74 \pm 0.16$	$14.97 \pm 0.22^a$
E176-G2	${}_{\Lambda\Lambda}^{11}\text{Be}$	$8.86 \pm 0.11$	$17.53 \pm 0.71$	$18.23 \pm 0.16$	$18.40 \pm 0.28$
E373-Hida	${}_{\Lambda\Lambda}^{11}\text{Be}$	$8.86 \pm 0.11$	$20.83 \pm 1.27$	$18.23 \pm 0.16$	$18.40 \pm 0.28$
E373-Hida	${}_{\Lambda\Lambda}^{12}\text{Be}$	$10.02 \pm 0.05$	$22.48 \pm 1.21$	...	$20.72 \pm 0.20$
E176-E2	${}_{\Lambda\Lambda}^{12}\text{B}$	$10.09 \pm 0.05$	$20.02 \pm 0.78$	...	$20.85 \pm 0.20$
E176-E4	${}_{\Lambda\Lambda}^{13}\text{B}$	$11.27 \pm 0.06$	$23.4 \pm 0.7$	...	$23.21 \pm 0.21$

A. Gal, E. V. Hungerford, and D. J. Millener, Rev. Mod. Phys. 88(2016)035004

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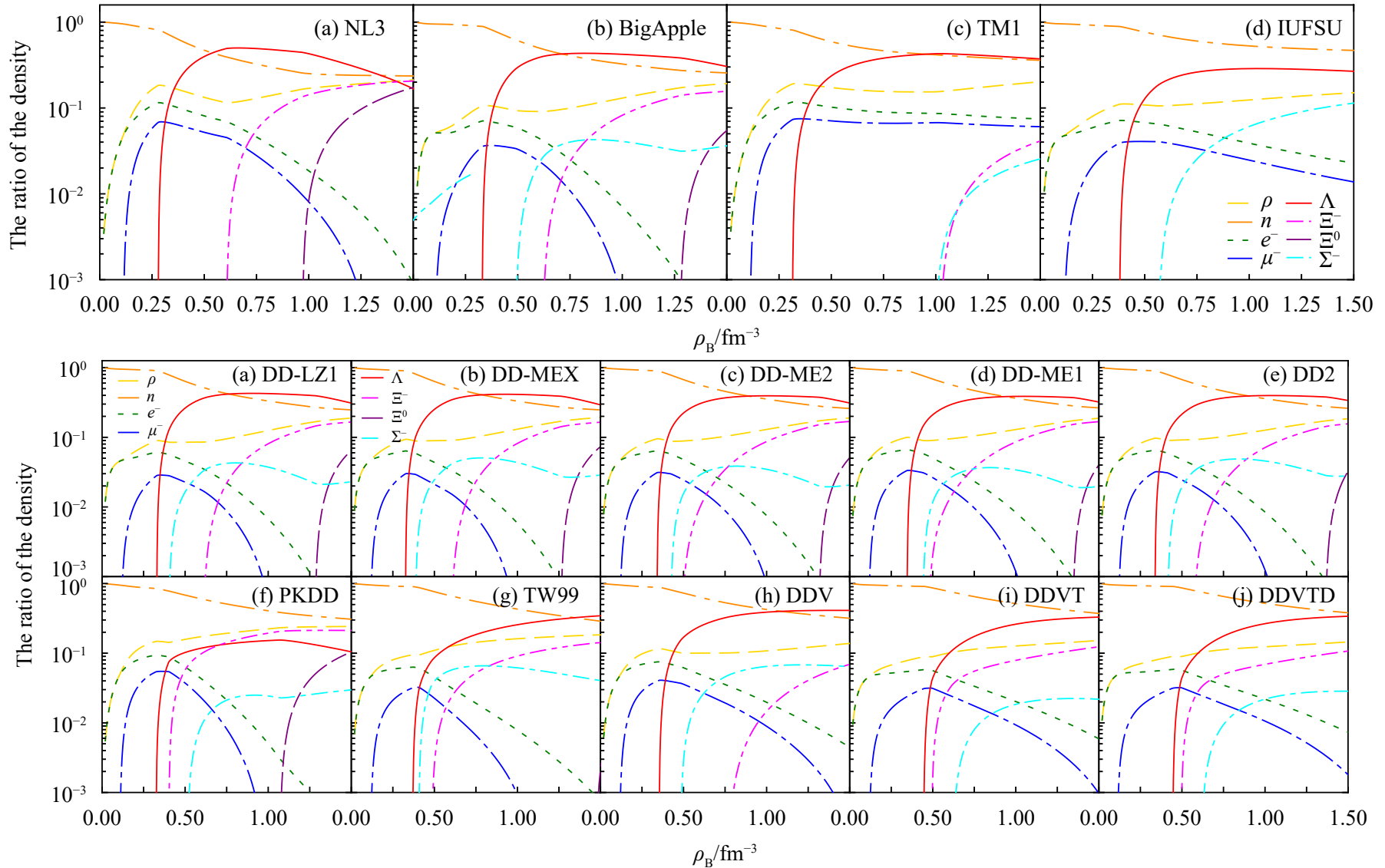


K. Nakazawa et al., Prog. Theor. Exp. Phys. 2015, 033D02 (2015)

M. Yoshimoto et al. Prog. Theor. Exp. Phys. 2021, 073D02 (2021)

S. H. Hayakawa et al. [J-PARC E07 Collaboration], Phys. Rev. Lett. 126 062501 (2021)

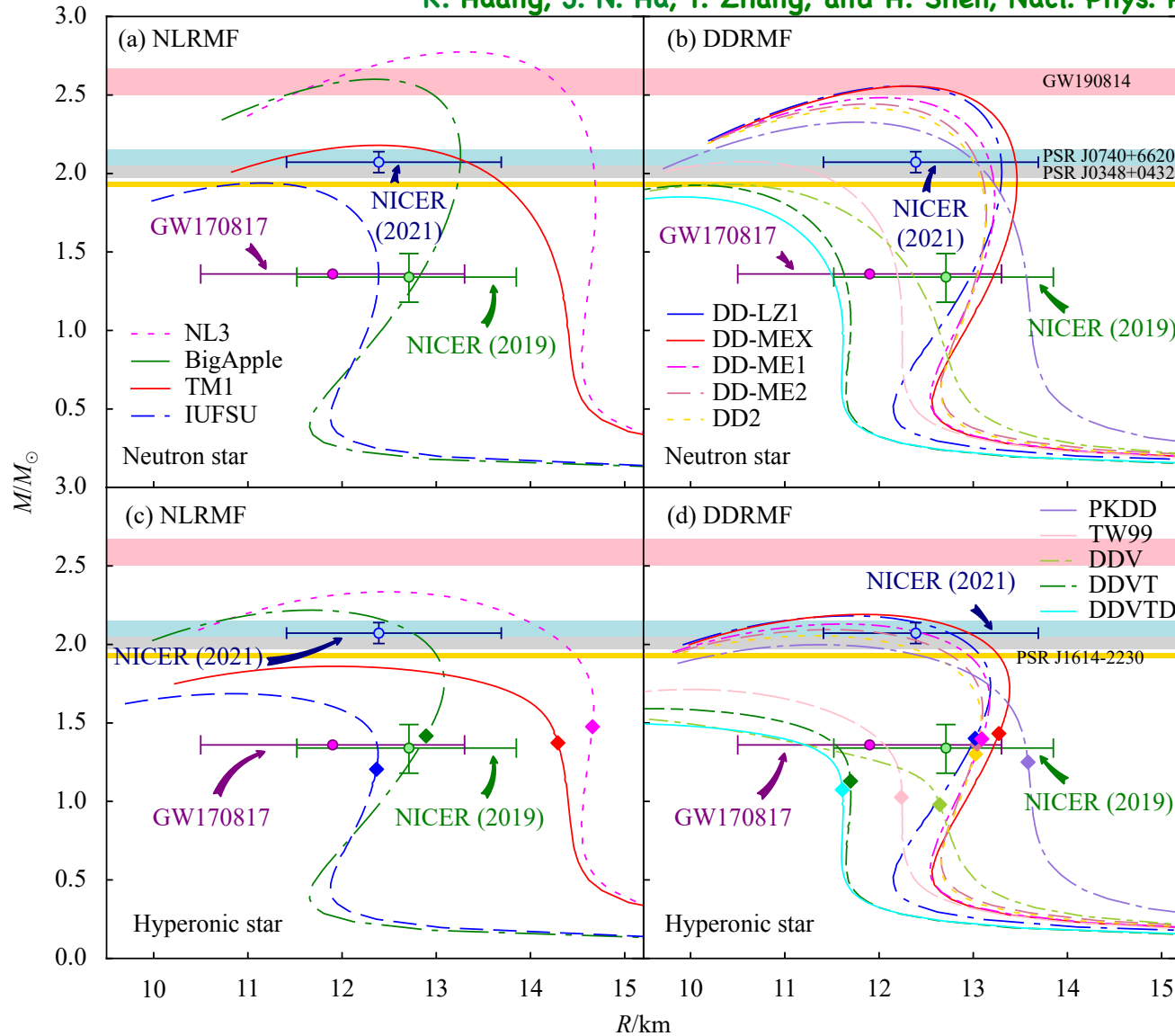
# Hyperon in neutron star

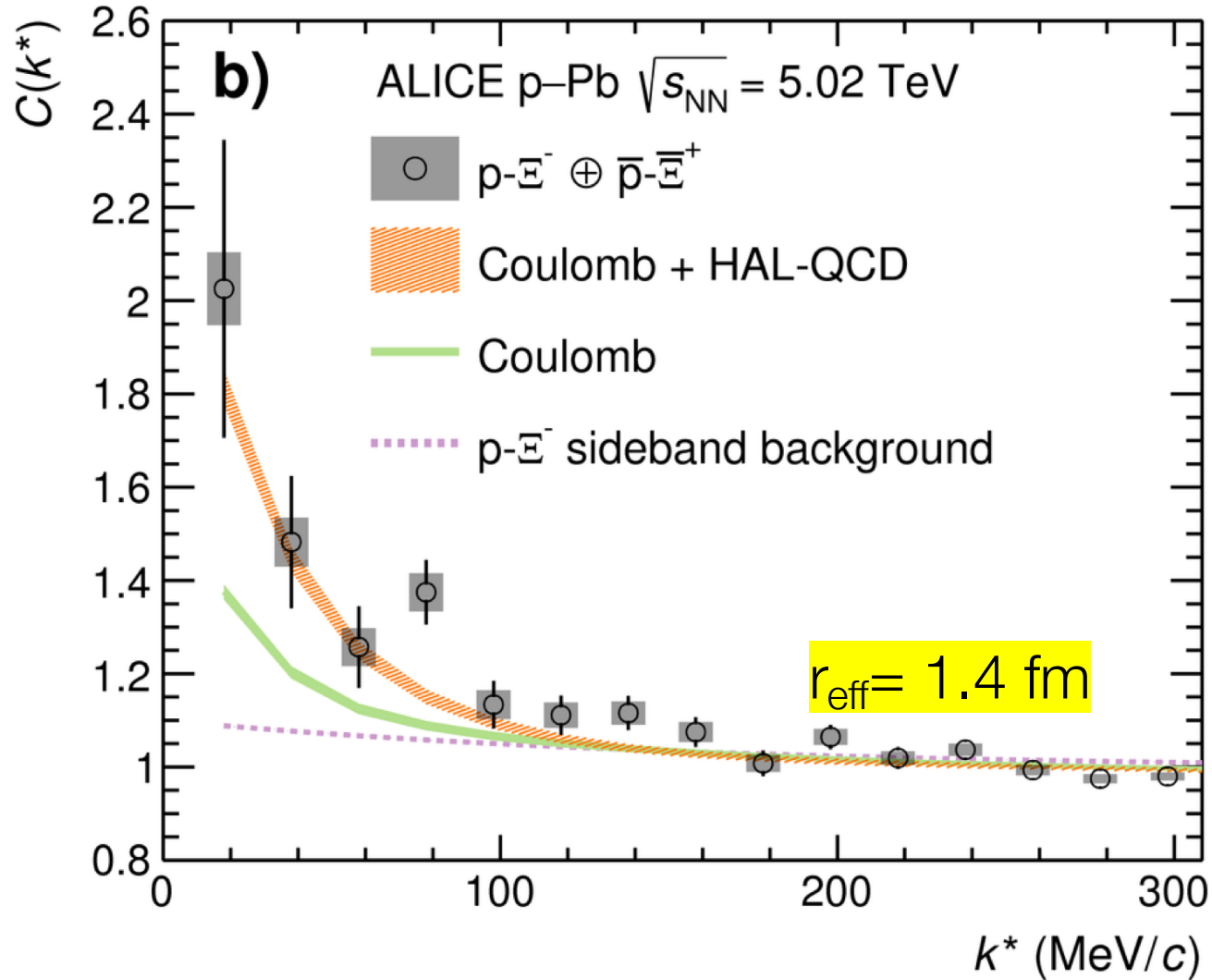


K. Huang, J. N. Hu, Y. Zhang, and H. Shen, Nucl. Phys. Rev. 39(2022)35

# Hyperon in neutron star

K. Huang, J. N. Hu, Y. Zhang, and H. Shen, Nucl. Phys. Rev. 39(2022)35







## ✓ ab initio methods

H. Nemura, Y. Akaishi, and Y. Suzuki, Phys. Rev. Lett. 89(2002)142504

E. Hiyama and T. Yamada, Prog. Part. Nucl. Phys. 63(2009)339

D. Lonardoni, S. Gandolfi, and F. Pederiva, Phys. Rev. C 87(2013)041303(R)

R. Wirth, et al. Phys. Rev. Lett. 113(2014)192502

## ✓ Shell model

D. J. Millener, Nucl. Phys. A 881(2012)298

## ✓ Skyrme Hartree-Fock model

M. Rayet, Ann. Phys. (NY) 102(1976)226

X. R. Zhou, et al. Phys. Rev. C 76(2007)034312

H.-J. Schulze and T. Rijken, Phys. Rev. C 88(2013)024322

Y. Zhang, H. Sagawa, and E. Hiyama Phys. Rev. C 103(2021)034321

## ✓ Covariant density functional theory

H. Shen, F. Yang, and H. Toki, Prog. Theor. Phys. 115(2006)325

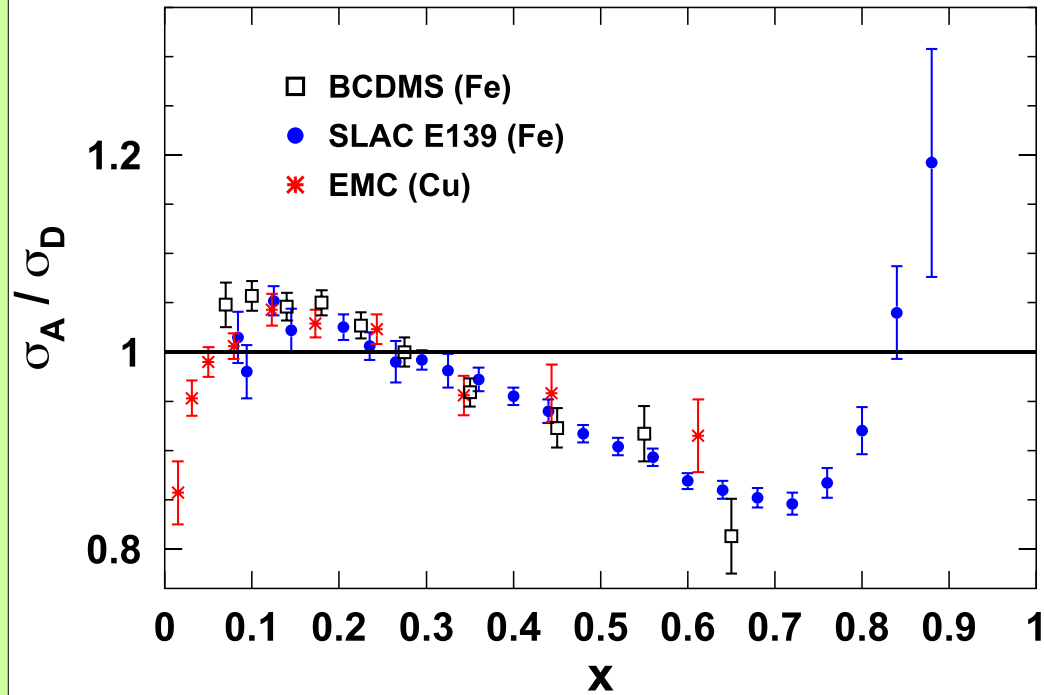
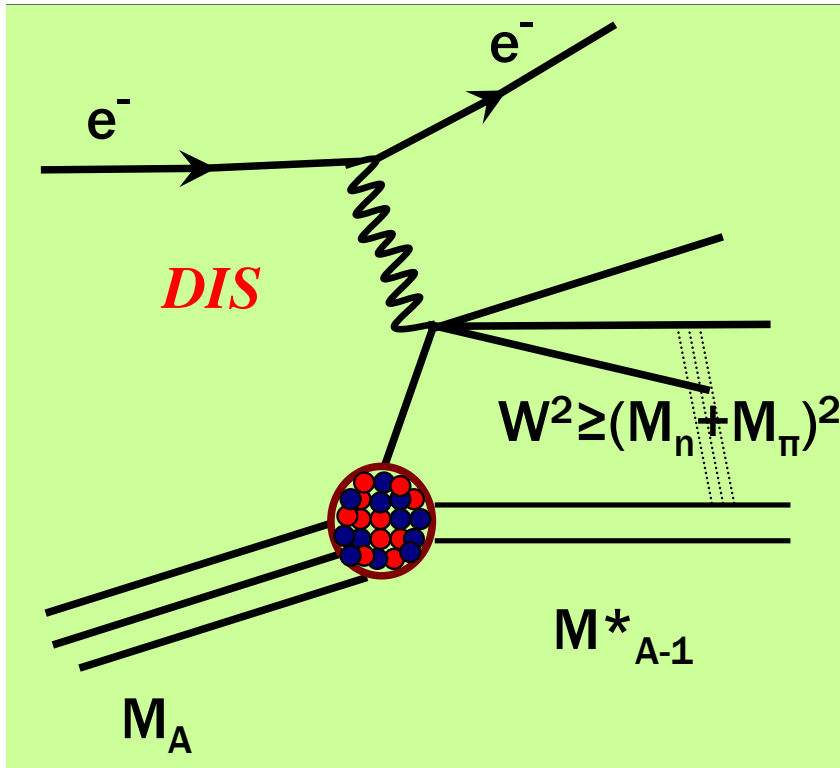
R. L. Xu, C. Wu, and Z. Z. Ren, J. Phys. G 39(2012)085107

T. T. Sun, et al., Phys. Rev. C 94(2016)064319

S. Y. Ding, Z. Qian, B. Y. Sun, and W. H. Long, Phys. Rev. C 106(2022)054311

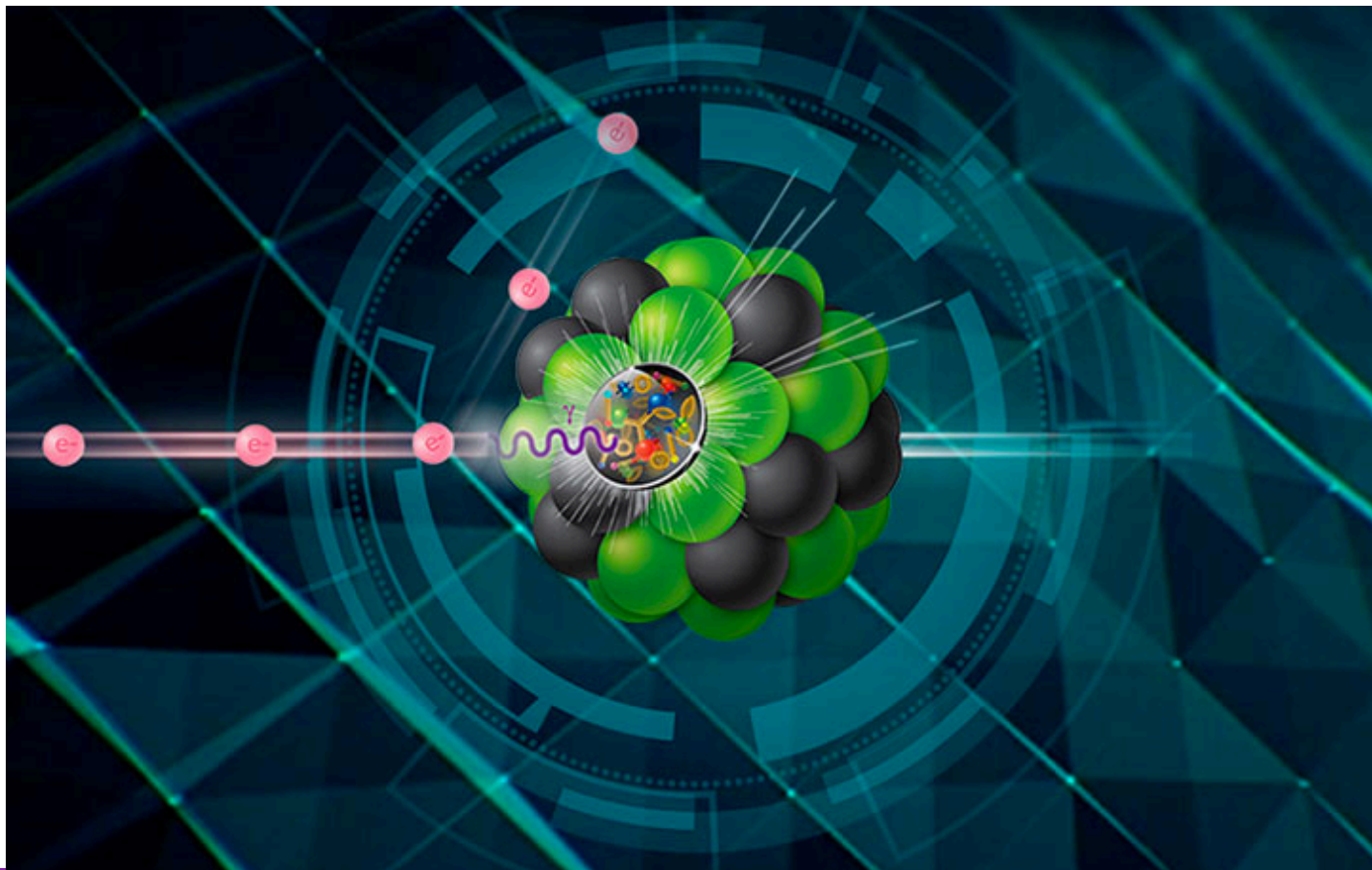
Y. T. Rong, P. W. Zhao, and S. G. Zhou, Phys. Lett, B, 807(2020)135533

I. C. Cloet, et al. J. Phys. G 46(2018)093001



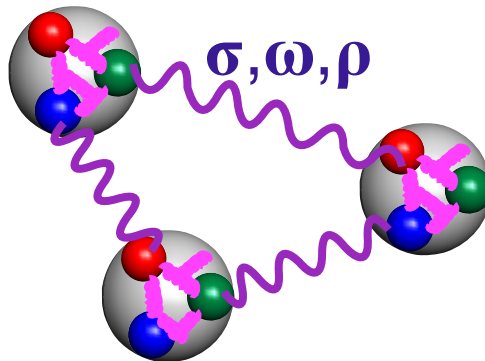


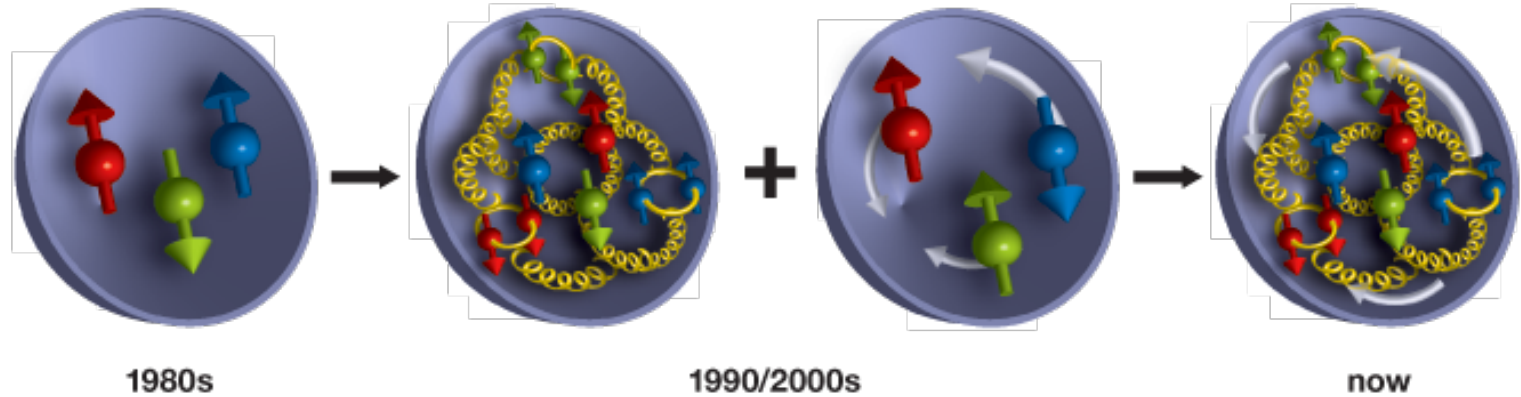
U.S. Department of Energy Selects Brookhaven National Laboratory to Host Major New Nuclear Physics Facility



# Relativistic many-body theories from quark level

- ✓ baryons are not point particles!
- ✓ baryon properties change in medium!
- ✓ quark-gluon plasma!





✓ Quark meson coupling (QMC) model

K. Tsushima, et al. Nucl. Phys. A 630(1998)691

✓ Friedberg-Lee model

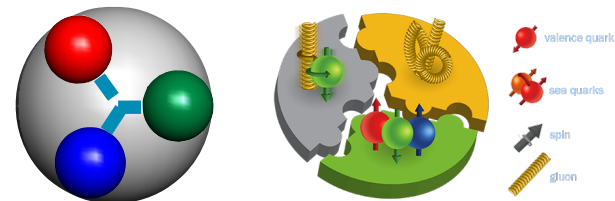
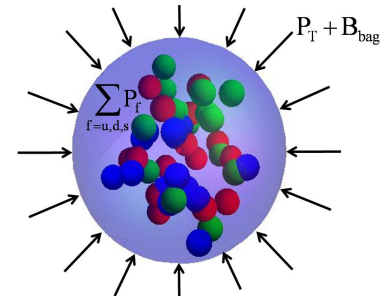
J. S. Liang and H. Shen, Phys. Rev. C88 (2013) 035208

.....

✓ Quark mean field (QMF) model

H. Toki, U. Meyer, A. Faessler, and R. Brockmann, Phys. Rev. C 58 (1998) 3749

H. Shen, H. Toki, Nucl. Phys. A 707 (2002) 469





# Outline

1 Introduction

2 The quark mean field model

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4 The summary and perspective

X. Xing, J.H., and H. Shen, Phys. Rev. C 94 (2016) 044308

- **Constituent quark in Dirac equation**

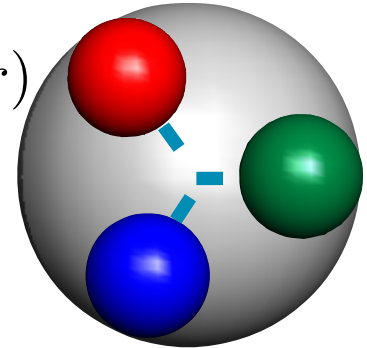
$$[-i\alpha \cdot \nabla + \beta m_i^* + \beta U(r)]q_i(r) = \varepsilon_i^* q_i(r)$$

where, the effective quark mass is

$$m_i^* = m_i + g_\sigma^i \sigma$$

and effective single particle energy is

$$\varepsilon_i^* = \varepsilon_i - g_\omega^i - g_\rho^i \rho \tau_3$$



- **Confinement potential**

$$U(r) = \frac{1}{2}(1 + \gamma^0)(ar^2 + V_0)$$

- **Center-of-mass corrections**

$$\langle B | \sum_{i=1}^3 \gamma^0(i) \left\{ \frac{1}{3} \gamma^0(i) \cdot \sum_{j=1}^3 \vec{p}_j + \frac{1}{2} (1 + \gamma^0(i)) [U(r_i) - U(\rho_i)] \right\} | B \rangle$$

- Pionic self-energy correction

$$\delta M_B^\pi = - \sum_k \sum_{B'} \frac{V_j^{\dagger BB'} V_j^{BB'}}{w_k}$$

- Gluon correction

**Color-electric**  $(\Delta E_B)_g^E = \frac{1}{8\pi} \sum_{i,j} \sum_{a=1}^8 \int \frac{d^3 r_i d^3 r_j}{|\vec{r}_i - \vec{r}_j|} \langle B | J_i^{0a}(\vec{r}_i) J_j^{0a}(\vec{r}_j) | B \rangle$

**Color-magnetic**  $(\Delta E_B)_g^M = -\frac{1}{8\pi} \sum_{i,j} \sum_{a=1}^8 \int \frac{d^3 r_i d^3 r_j}{|\vec{r}_i - \vec{r}_j|} \langle B | \vec{J}_i^a(\vec{r}_i) \cdot \vec{J}_j^a(\vec{r}_j) | B \rangle$

**Quark color current density**  $J_i^{\mu a}(x) = g_c \bar{\psi}_q(x) \gamma^\mu \lambda_i^a \psi_q(x)$

- Baryon mass

$$M_B^* = E_B^{*0} - \epsilon_{\text{c.m.}} + \delta M_B^\pi + (\Delta E_B)_g^E + (\Delta E_B)_g^M$$

## • Strangeness QMF Lagrangian

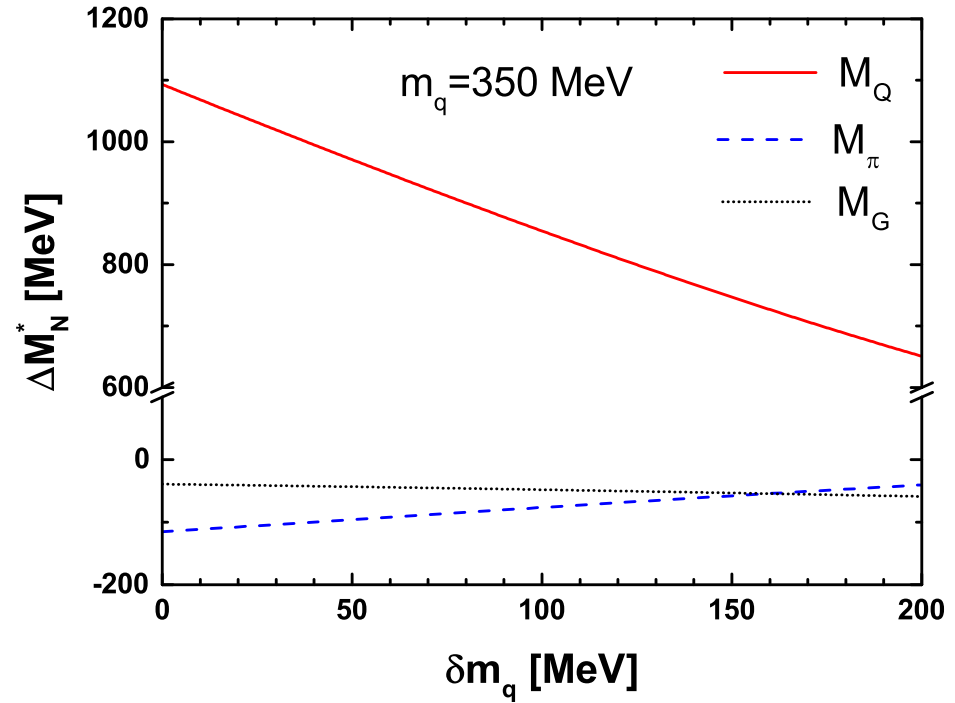
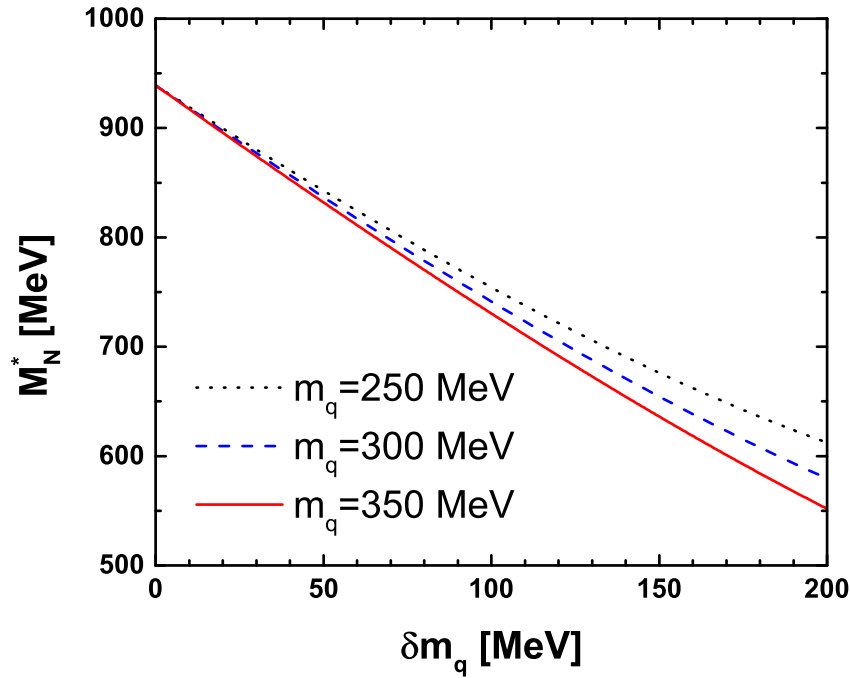
$$\begin{aligned} \mathcal{L}_{\text{QMF}} = & \bar{\psi} \left[ i\gamma_{\mu} \partial^{\mu} - M_N^* - g_{\omega} \omega \gamma^0 - g_{\rho} \rho \tau_3 \gamma^0 - e \frac{(1 - \tau_3)}{2} A \gamma^0 \right] \psi \\ & + \bar{\psi}_H \left[ i\gamma_{\mu} \partial^{\mu} - M_H^* - g_{\omega}^H \omega \gamma^0 + \frac{f_{\omega}^H}{2M_H} \sigma^{0i} \partial_i \omega \right] \psi_H \\ & - \frac{1}{2} (\nabla \sigma)^2 - \frac{1}{2} m_{\sigma}^2 \sigma^2 - \frac{1}{3} g_2 \sigma^3 - \frac{1}{4} g_3 \sigma^4 \\ & + \frac{1}{2} (\nabla \omega)^2 + \frac{1}{2} m_{\omega}^2 \omega^2 + \frac{1}{4} c_3 \omega^4 \\ & + \frac{1}{2} (\nabla \rho)^2 + \frac{1}{2} m_{\rho}^2 \rho^2 + \frac{1}{2} (\nabla A)^2, \end{aligned}$$

## • Dirac equations for baryons

$$\begin{aligned} \left[ i\gamma_{\mu} \partial^{\mu} - M_N^* - g_{\omega} \omega \gamma^0 - g_{\rho} \rho \tau_3 \gamma^0 - e \frac{(1 - \tau_3)}{2} A \gamma^0 \right] \psi &= 0, \\ \left[ i\gamma_{\mu} \partial^{\mu} - M_H^* - g_{\omega}^H \omega \gamma^0 + \frac{f_{\omega}^H}{2M_H} \sigma^{0i} \partial_i \omega \right] \psi_H &= 0. \end{aligned}$$

## • Equations of motion for mesons

$$\begin{aligned} \Delta \sigma - m_{\sigma}^2 \sigma - g_2 \sigma^2 - g_3 \sigma^3 &= \frac{\partial M_N^*}{\partial \sigma} \langle \bar{\psi} \psi \rangle + \frac{\partial M_H^*}{\partial \sigma} \langle \bar{\psi}_H \psi_H \rangle, \\ \Delta \omega - m_{\omega}^2 \omega - c_3 \omega^3 &= -g_{\omega} \langle \bar{\psi} \gamma^0 \psi \rangle - g_{\omega}^H \langle \bar{\psi}_H \gamma^0 \psi_H \rangle + \frac{f_{\omega}^H}{2M_H} \partial_i \langle \bar{\psi}_H \sigma^{0i} \psi_H \rangle, \\ \Delta \rho - m_{\rho}^2 \rho &= -g_{\rho} \langle \bar{\psi} \tau_3 \gamma^0 \psi \rangle, \\ \Delta A &= -e \langle \bar{\psi} \frac{(1 - \tau_3)}{2} \gamma^0 \psi \rangle. \end{aligned}$$



$m_q$ (MeV)	$g_\sigma^q$	$g_\omega$	$g_\rho$	$g_2$ (fm <sup>-1</sup> )	$g_3$	$c_3$	$a$ (fm <sup>-3</sup> )	$V_0$ (MeV)
250	5.15871	11.54726	3.79601	-3.52737	-78.52006	305.00240	0.57945	-24.28660
300	5.09346	12.30084	4.04190	-3.42813	-57.68387	249.05654	0.53430	-62.25719
350	5.01631	12.83898	4.10772	-3.29969	-39.87981	221.68240	0.49560	-102.04158

$$U(r) = \frac{1}{2}(1 + \gamma^0)(ar^2 + V_0)$$



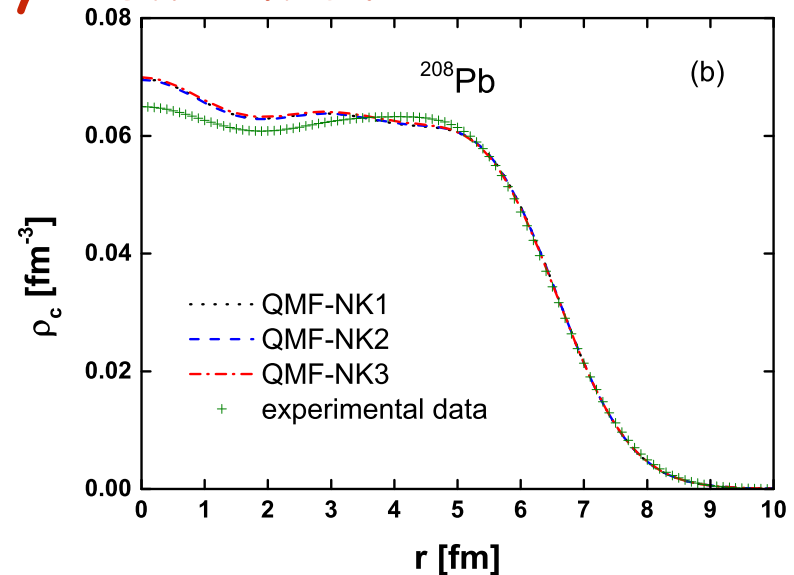
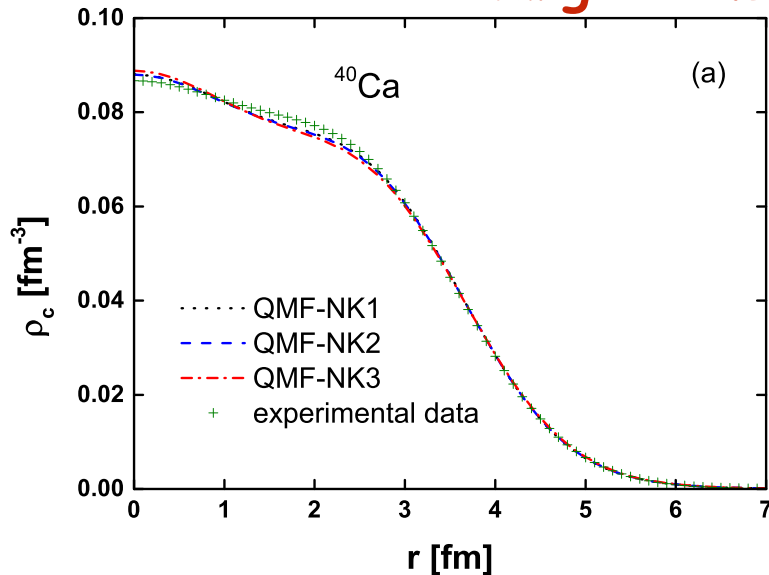
## Binding energy and charge radii

Model	$E/A$ (MeV)				$R_c$ (fm)			
	$^{40}\text{Ca}$	$^{48}\text{Ca}$	$^{90}\text{Zr}$	$^{208}\text{Pb}$	$^{40}\text{Ca}$	$^{48}\text{Ca}$	$^{90}\text{Zr}$	$^{208}\text{Pb}$
QMF-NK1	8.62	8.61	8.65	7.92	3.43	3.47	4.26	5.49
QMF-NK2	8.61	8.61	8.67	7.91	3.44	3.47	4.26	5.50
QMF-NK3	8.59	8.63	8.68	7.90	3.44	3.46	4.26	5.50
QMF[18]	8.35	8.43	8.54	7.73	3.44	3.46	4.27	5.53
Expt.	8.55	8.67	8.71	7.87	3.45	3.45	4.26	5.50

Without pion  
and gluon  
corrections



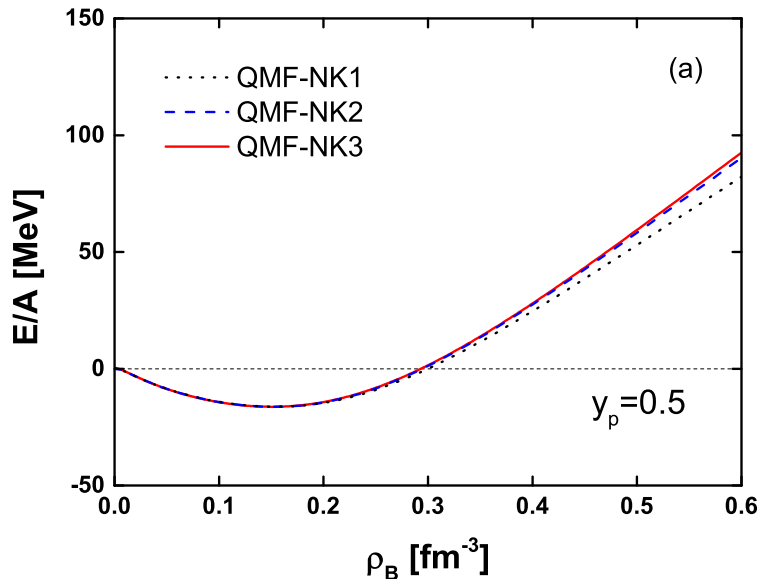
## Charge density distribution



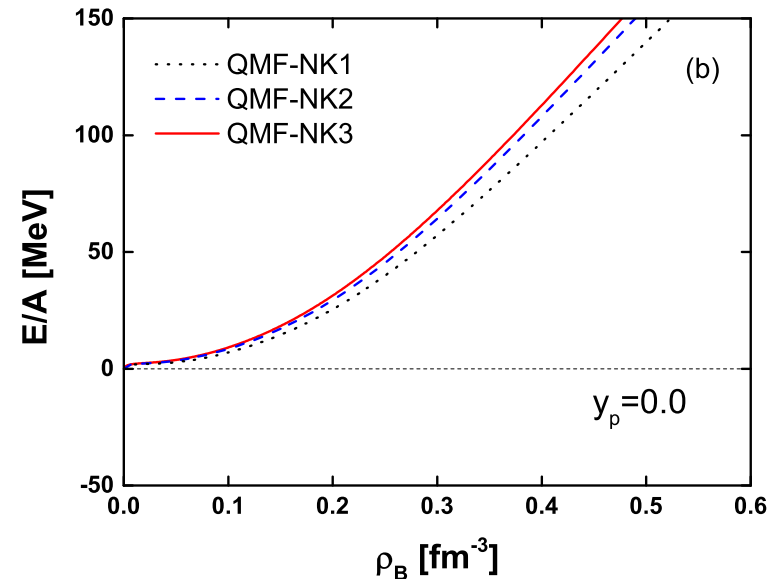
## Nuclear saturation properties

Model	$\rho_0$ ( $\text{fm}^{-3}$ )	$E/A$ (MeV)	$K_0$ (MeV)	$J$ (MeV)	$M_N^*/M_N$	$L^0$ (MeV)	$K_{\text{sym}}^0$ (MeV)	$K_{\text{asy}}$ (MeV)	$Q_0$ (MeV)	$K_\tau$ (MeV)
QMF-NK1	0.154	-16.3	323	30.6	0.70	84.8	-28.8	-537.6	495.4	-667.7
QMF-NK2	0.152	-16.3	328	32.9	0.66	93.7	-23.5	-585.7	221.0	-648.8
QMF-NK3	0.150	-16.3	322	33.6	0.64	97.3	-12.0	-595.8	263.0	-675.3

### symmetric nuclear matter



### pure neutron matter



# Outline

1 Introduction

2 The quark mean field model

3 The strangeness with QMF model


4 The summary and perspective

## The strength of quark confinement potential

	$m_u$ (MeV)	$V_u$ (MeV)	$a_u$ (fm <sup>-3</sup> )	$m_s$ (MeV)	$V_s$ (MeV)	$a_s$ (fm <sup>-3</sup> )
set A	250	-24.286601	0.579450	330	101.78180	0.097317
set B	300	-62.257187	0.534296	380	54.548210	0.087243
set C	350	-102.041575	0.495596	430	6.802695	0.079534

## The coupling constants between meson and baryons

Model	$m_u$ (MeV)	$g_\sigma^u$	$g_\omega$	$g_\omega^\Lambda$	$g_\omega^\Xi$	$g_\rho$	$g_2$	$g_3$	$c_3$
							(fm <sup>-1</sup> )		
QMF-NK1S	250	5.15871	11.54726	$0.8258g_\omega$	$0.4965g_\omega$	3.79601	-3.52737	-78.52006	305.00240
QMF-NK2S	300	5.09346	12.30084	$0.8134g_\omega$	$0.4800g_\omega$	4.04190	-3.42813	-57.68387	249.05654
QMF-NK3S	350	5.01631	12.83898	$0.8040g_\omega$	$0.4681g_\omega$	4.10772	-3.29969	-39.87981	221.68240



$$U_\Lambda^{(N)} = -30 \text{ MeV}$$

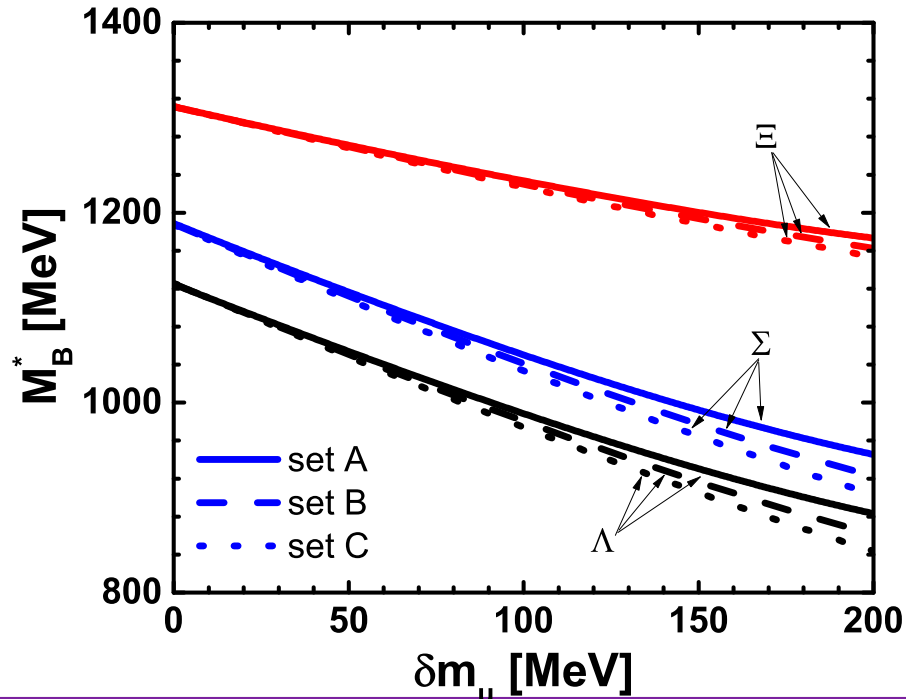
$$U_\Xi^{(N)} = -12 \text{ MeV}$$

# The effective masses of baryons



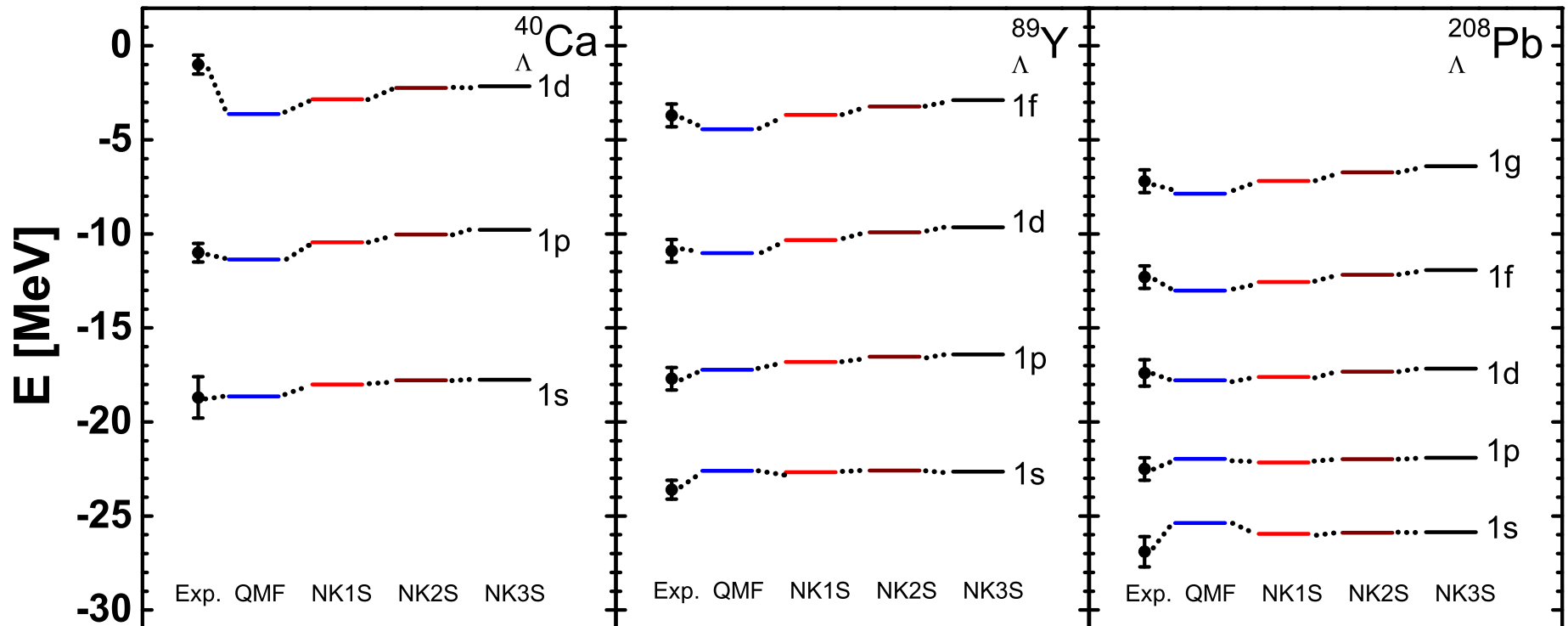
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	Baryon	$E_B^0$	$\epsilon_{c.m.}$	$\delta M_B^\pi$	$(\Delta E_B)_g$	$M_B^{\text{Theor.}}$	$M_B^{\text{Expt.}}$
set A	$\Lambda$	1446.340	231.975	-65.172	-24.390	1124.803	$1115.683 \pm 0.006$
	$\Sigma^0$	1446.340	231.975	-36.207	10.515	1188.673	$1192.642 \pm 0.024$
	$\Xi^0$	1504.254	175.047	-16.293	-1.289	1311.625	$1314.86 \pm 0.20$
set B	$\Lambda$	1433.489	220.692	-69.277	-18.313	1125.207	$1115.683 \pm 0.006$
	$\Sigma^0$	1433.489	220.692	-38.487	13.753	1188.063	$1192.642 \pm 0.024$
	$\Xi^0$	1491.611	165.564	-17.319	2.979	1311.707	$1314.86 \pm 0.20$
set C	$\Lambda$	1421.908	210.233	-72.829	-13.170	1125.676	$1115.683 \pm 0.006$
	$\Sigma^0$	1421.908	210.233	-40.461	16.203	1187.417	$1192.642 \pm 0.024$
	$\Xi^0$	1480.703	157.102	-18.207	6.377	1311.771	$1314.86 \pm 0.20$

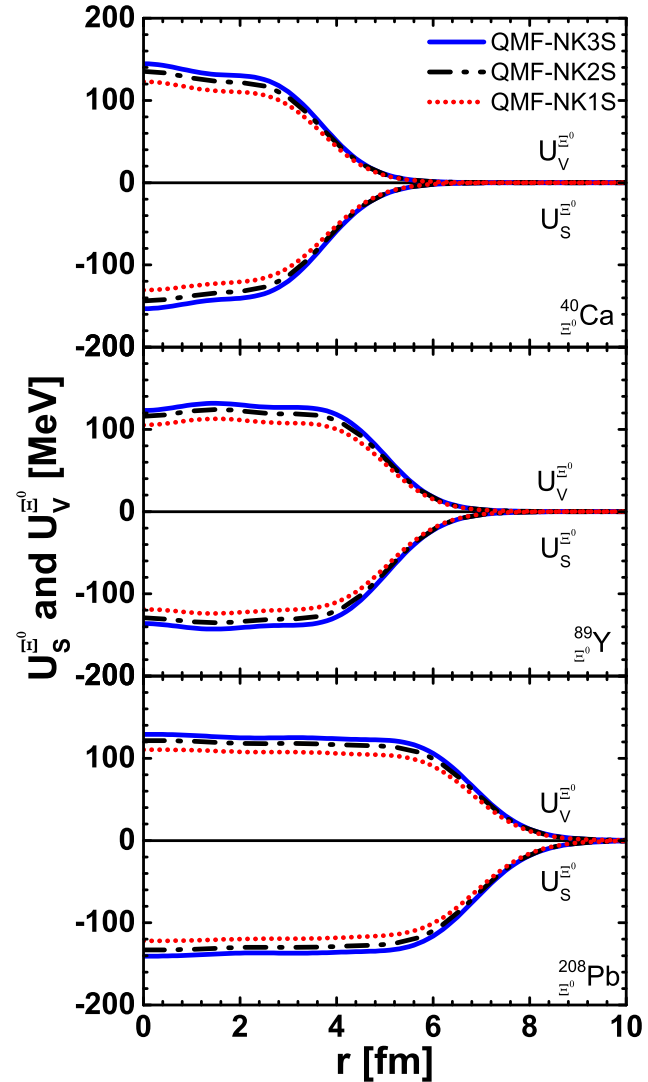
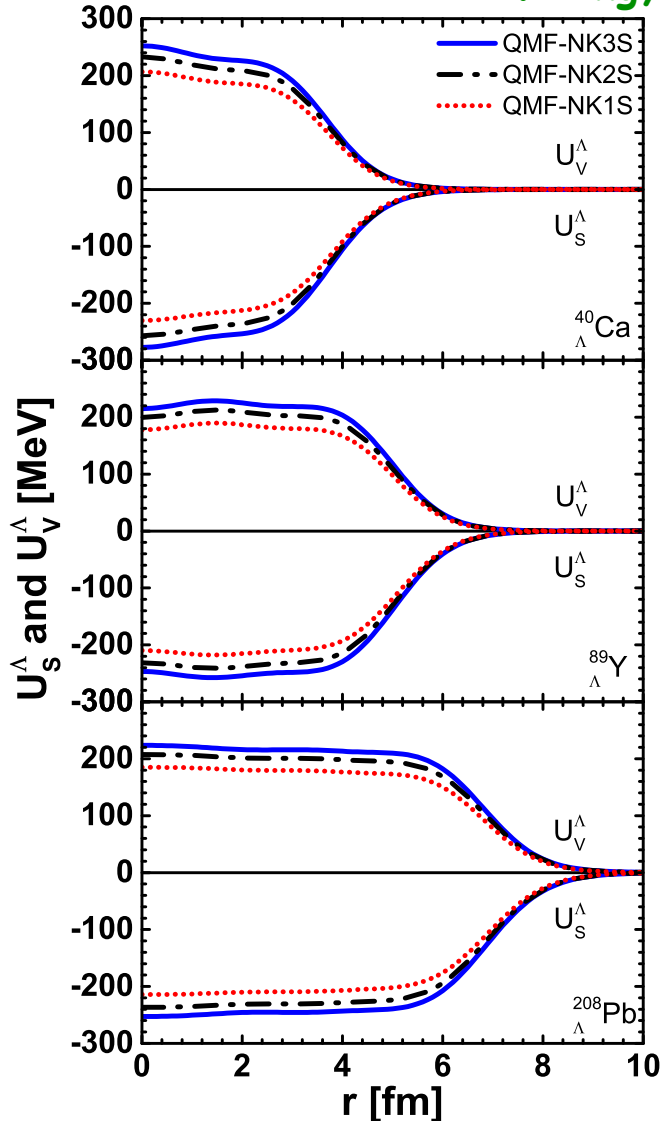


X. Xing, J.H., and H. Shen, Phys. Rev. C 95 (2017) 054310

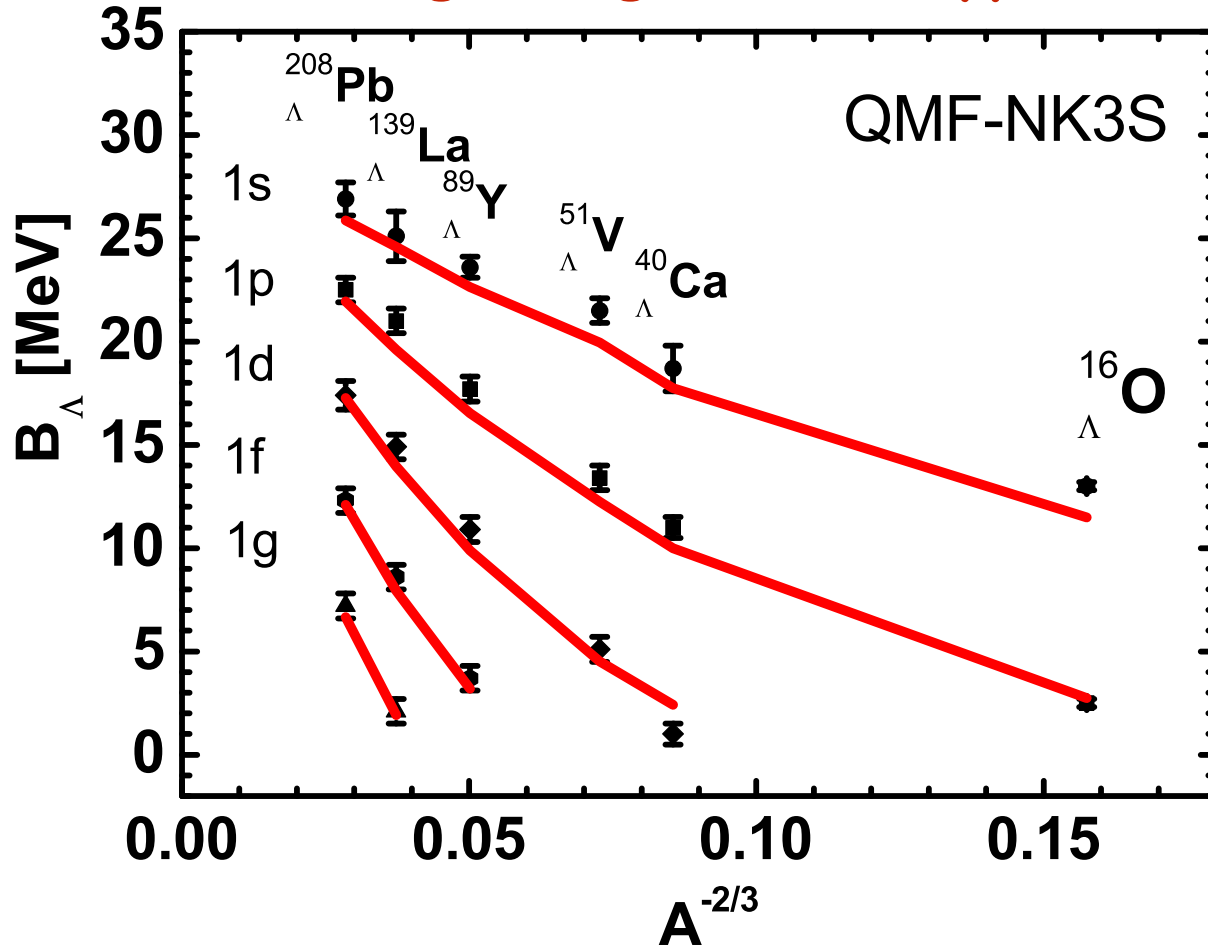
## The energy level of $\Lambda$ hypernuclei



X. Xing, J.H., and H. Shen, Phys. Rev. C 95 (2017) 054310



X. Xing, J.H., and H. Shen, Phys. Rev. C 95 (2017) 054310  
The binding energies of  $\Lambda$  hypernuclei



Exp. Data: A. Gal, E. V. Hungerford and D. J. Millener, Rev. Mod. Phys. 88(2016)035004



## The KISO event ( $\Xi^- + {}^{14}\text{N}$ ) of $\Xi^-$ hypernuclei

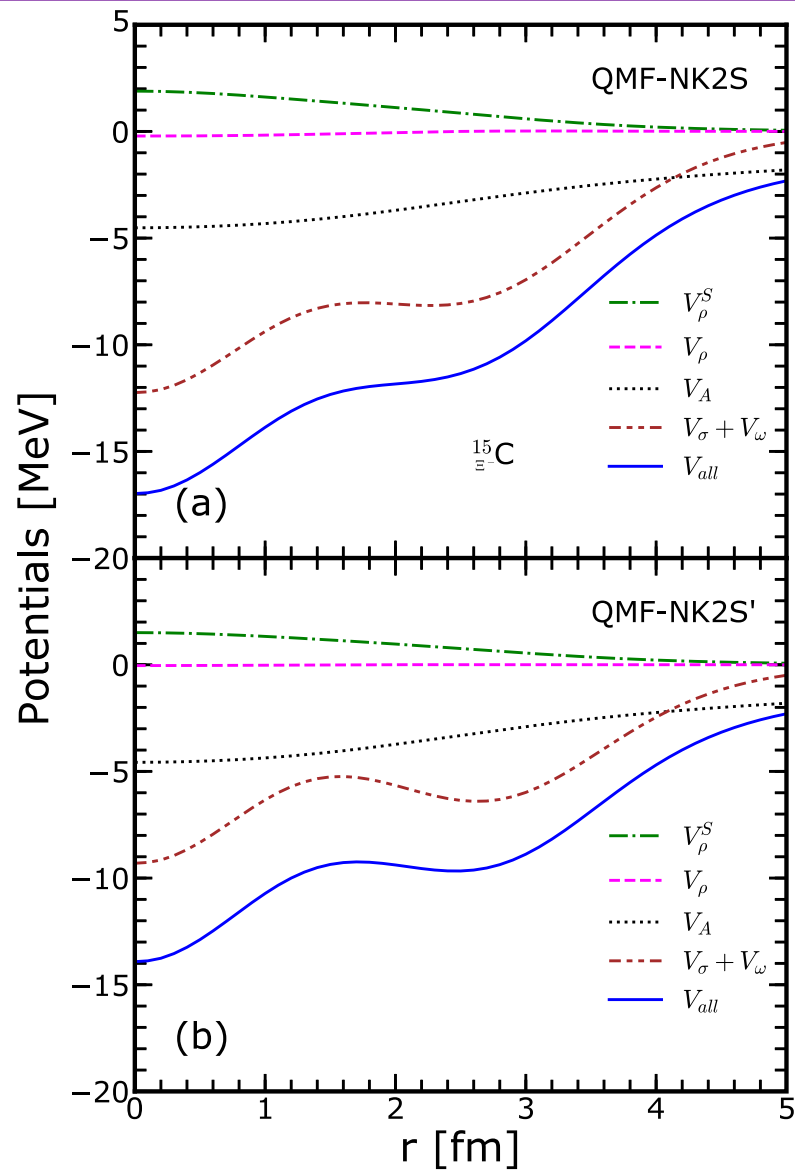
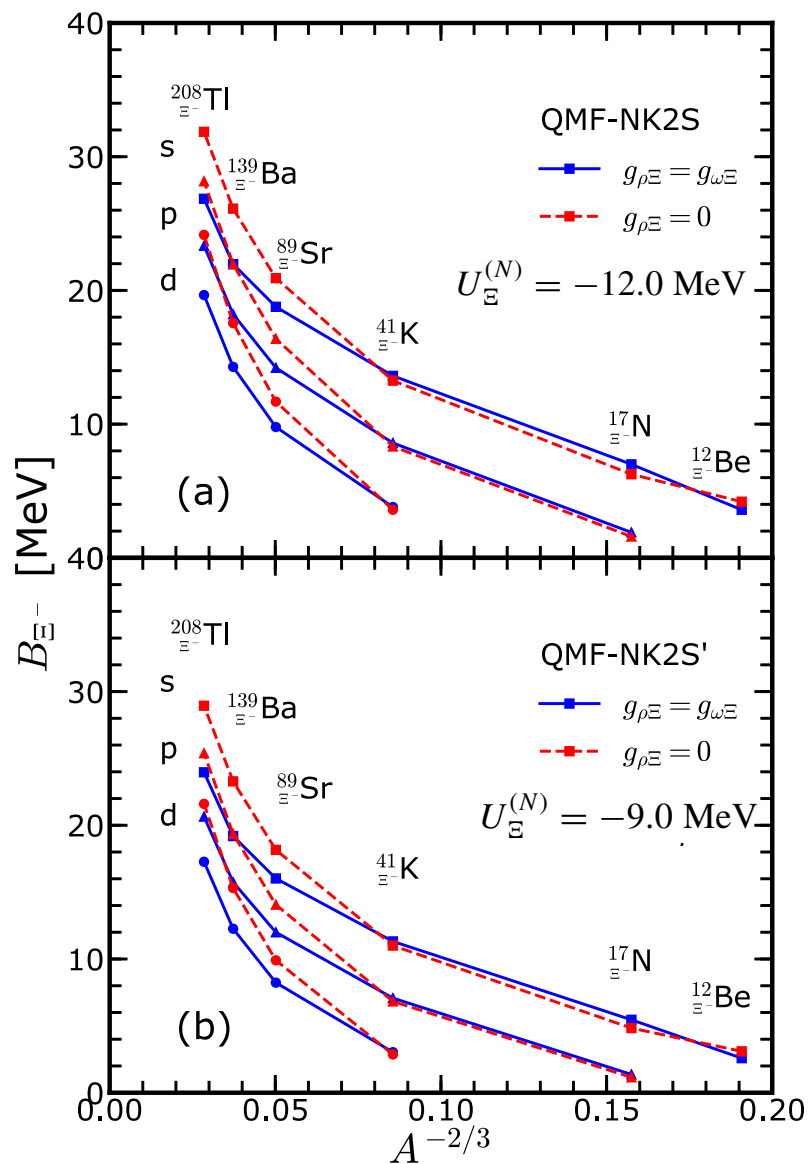
K. Nakazawa et al., Prog. Theor. Exp. Phys. 2015, 033D02 (2015)

	${}_{\Xi}^{15}\text{C}(1s)$	${}_{\Xi}^{15}\text{C}(1p)$	${}_{\Xi}^{12}\text{Be}(1s)$
QMF-NK1S( $g_{\rho\Xi} = g_{\rho N}$ )	5.82	1.21	3.78
QMF-NK2S( $g_{\rho\Xi} = g_{\rho N}$ )	5.69	1.14	3.59
QMF-NK3S( $g_{\rho\Xi} = g_{\rho N}$ )	5.61	1.08	3.49
QMF-NK1S( $g_{\rho\Xi} = 0$ )	5.80	1.21	4.35
QMF-NK2S( $g_{\rho\Xi} = 0$ )	5.65	1.14	4.20
QMF-NK3S( $g_{\rho\Xi} = 0$ )	5.58	1.08	4.11
Expt. or empirical data	$4.38 \pm 0.25$	$1.11 \pm 0.25$	3.0 – 5.5

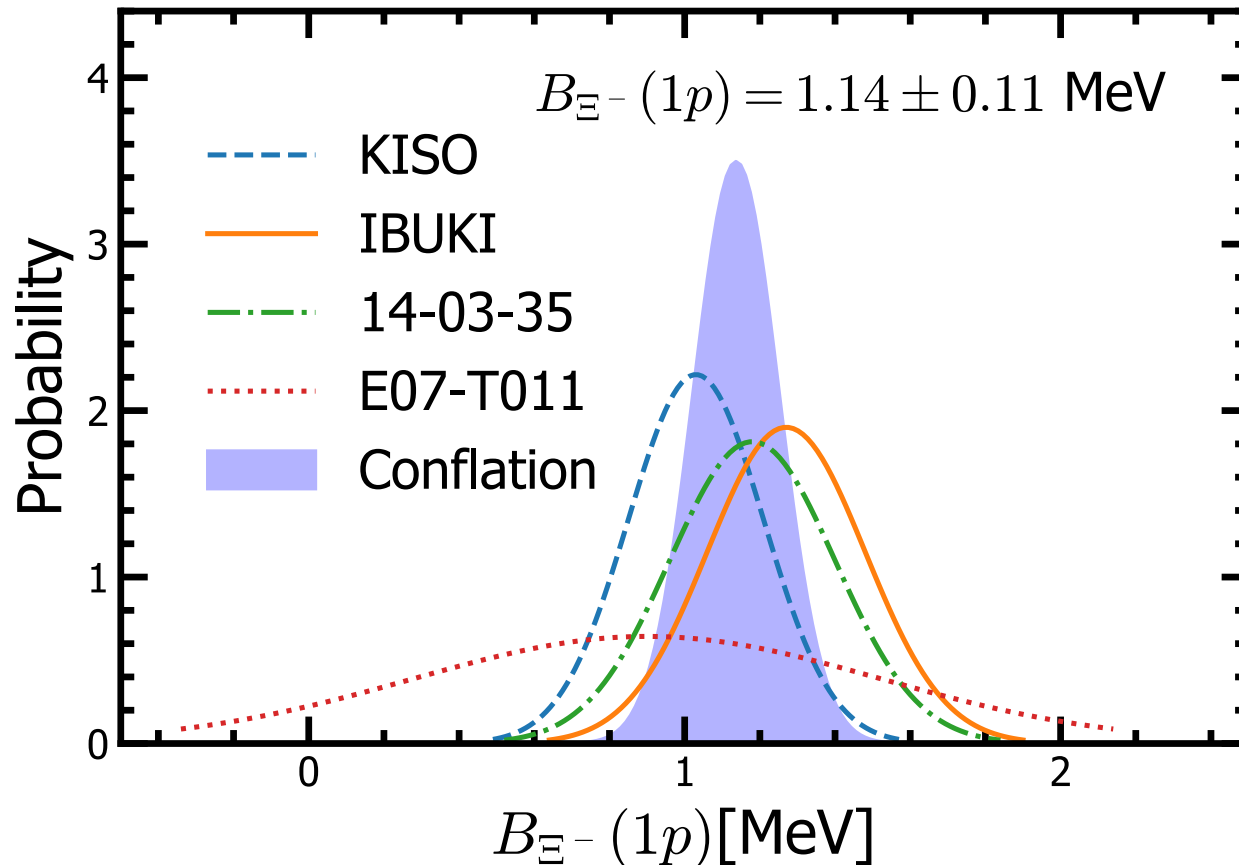
T.T.Sun, E.Hiyama, H.Sagawa, H.J.Schulze, J. Meng, Phys. Rev. C 94, 064319 (2016)

J. Margueron, E. Khan, F. Gulminelli, Phys. Rev. C 96, 054317 (2017).

**The IBUKI event (2021):**  $1.27 \pm 0.21$



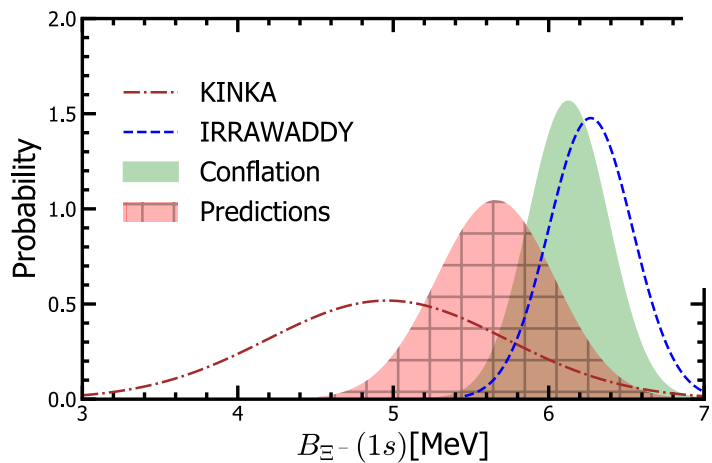
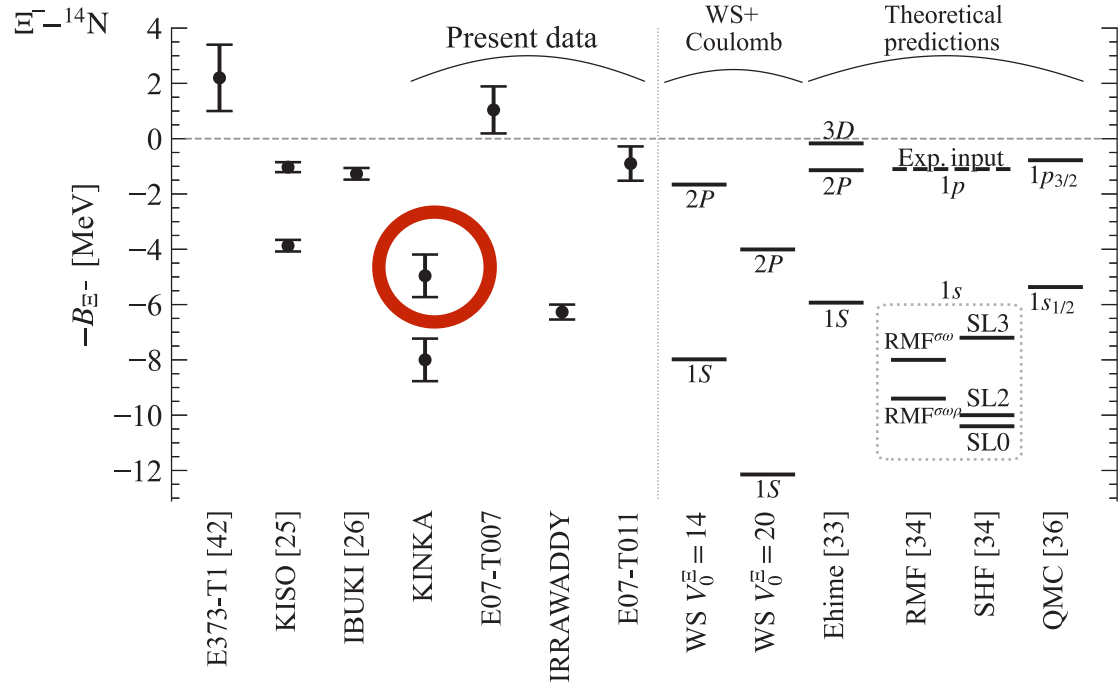
## The conflation values of $1p$ $\Xi^-$ -hyperon in $\Xi^-+^{14}\text{N}$ system



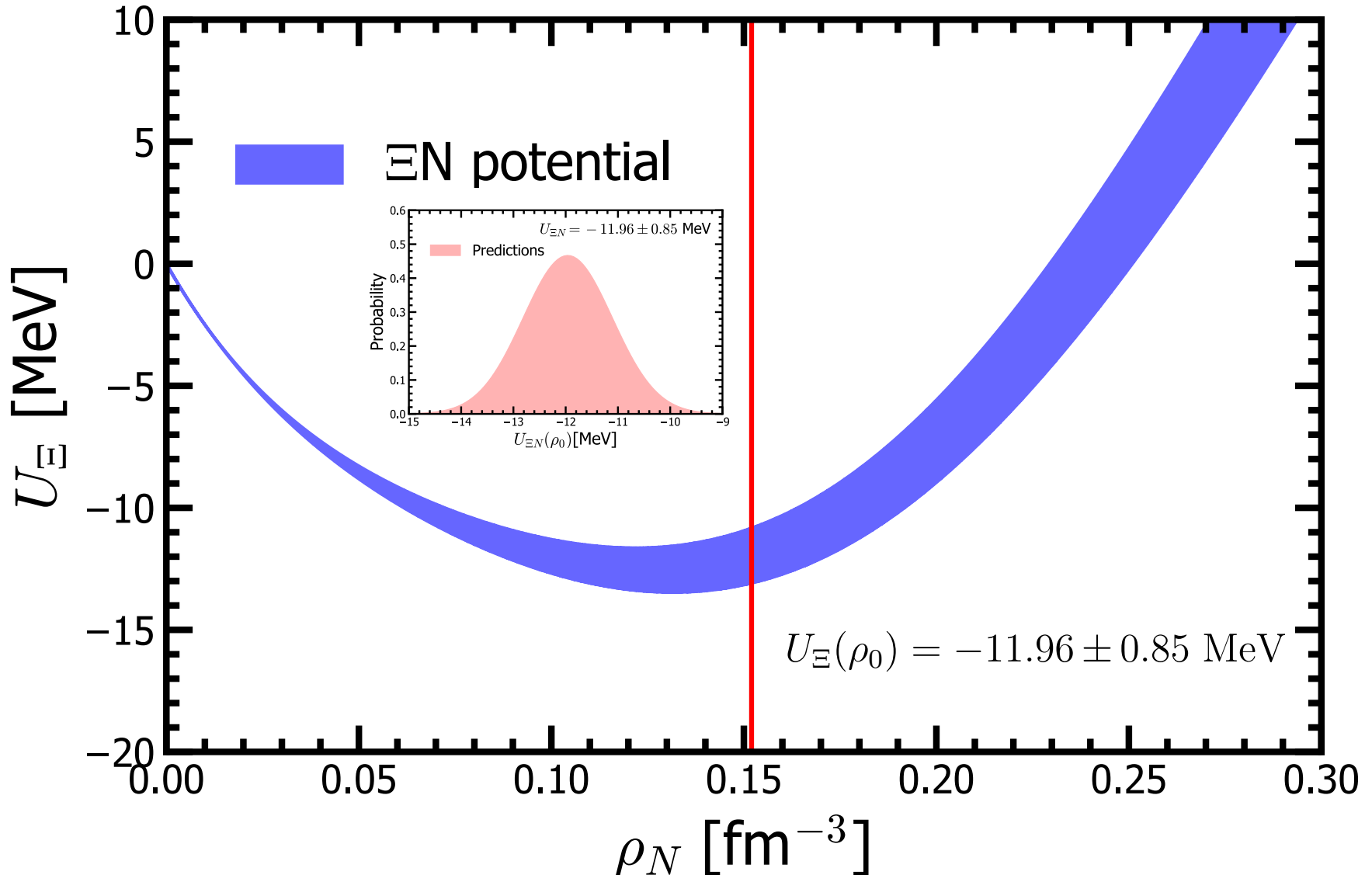
J. H., Y. Zhang, and H. Shen, J. Phys. G 49 (2022)025104

J. H. Y. Zhang, and H. Shen, J. Phys. G 49 (2022)025104

Sets	$g_{\omega\Xi}/g_{\omega N}$	$B_{\Xi^-}(1p)$	$B_{\Xi^-}(1s)$
QMF-NK1X1	0.5024	-1.03	-5.19
QMF-NK1X2	0.4954	-1.25	-5.91
QMF-NK2X1	0.4832	-1.03	-5.27
QMF-NK2X2	0.4771	-1.25	-6.01
QMF-NK3X1	0.4694	-1.03	-5.41
QMF-NK3X2	0.4638	-1.25	-6.11



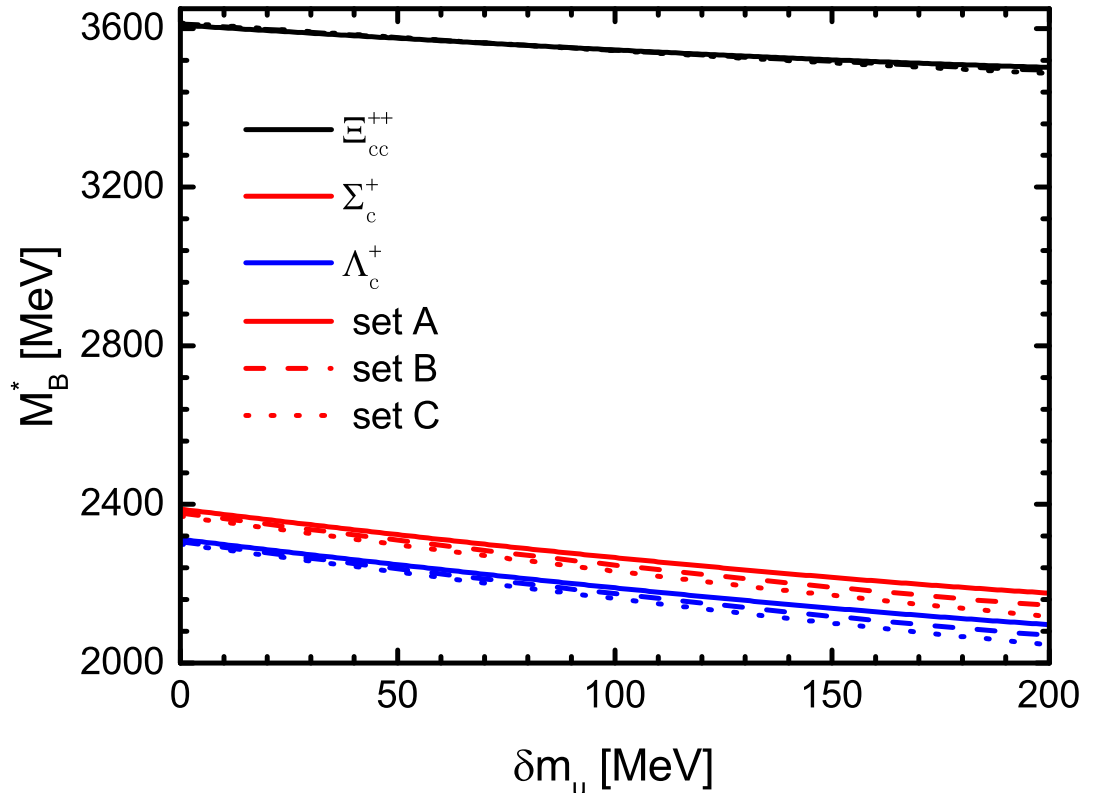
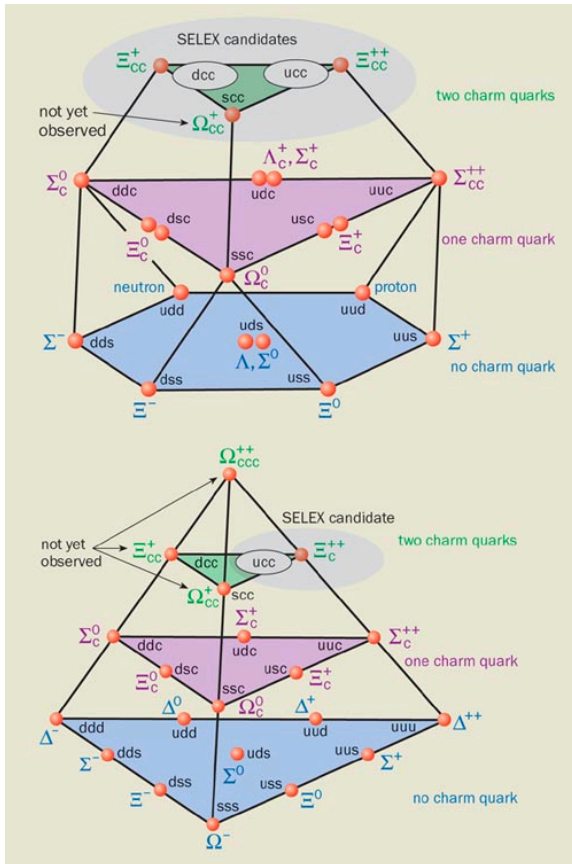
**The KINKA event should be 1s state!**



J. H., Y. Zhang, and H. Shen, *J. Phys. G* 49 (2022)025104

# The charmed baryons

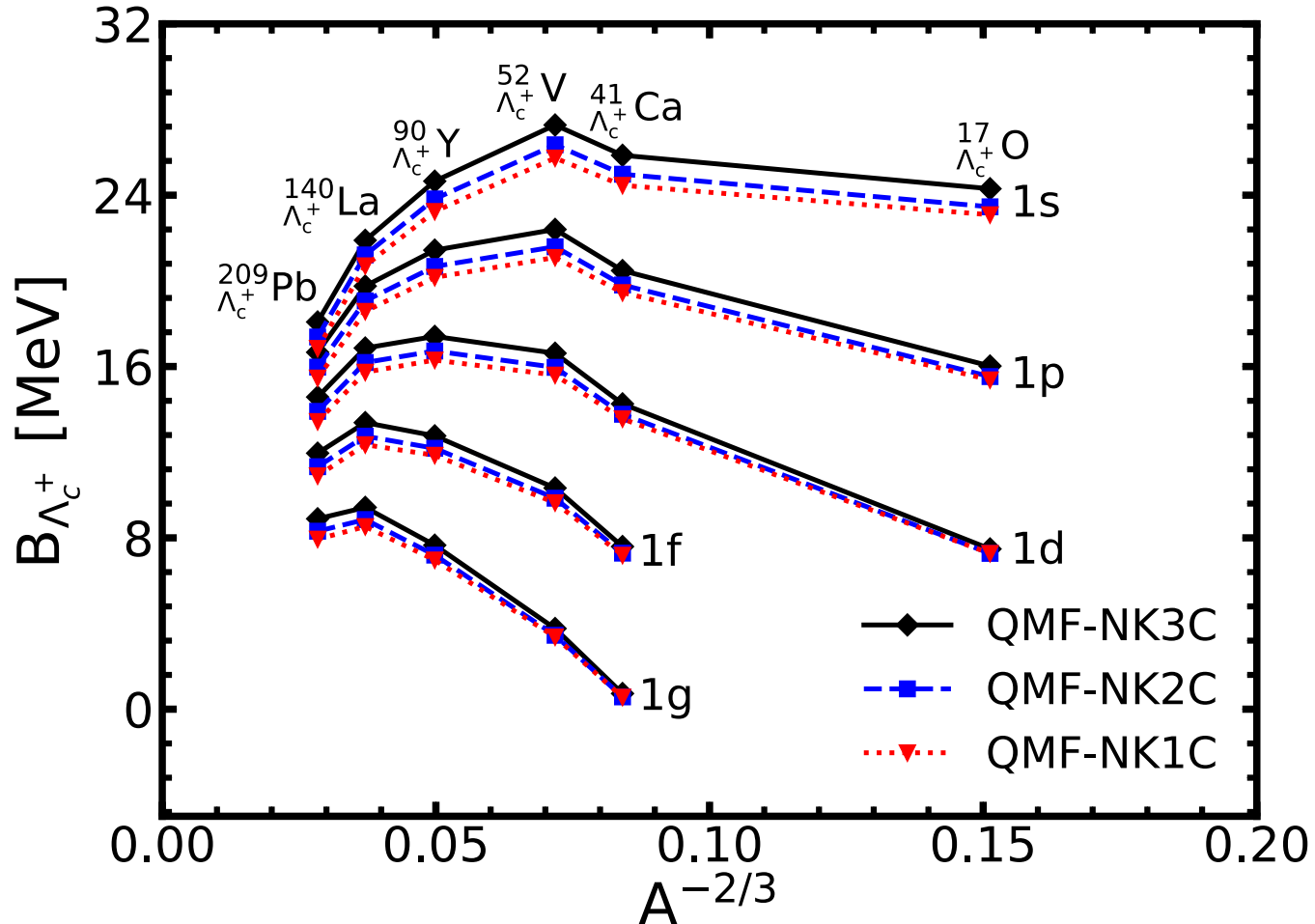
	$m_u$ (MeV)	$V_u$ (MeV)	$a_u$ (fm <sup>-3</sup> )	$m_c$ (MeV)	$V_c$ (MeV)	$a_c$ (fm <sup>-3</sup> )
Set A	250	-24.286601	0.579450	1300	284.58724	0.118172
Set B	300	-62.257187	0.534296	1350	239.53994	0.117312
Set C	350	-102.041575	0.495596	1400	193.67265	0.116036



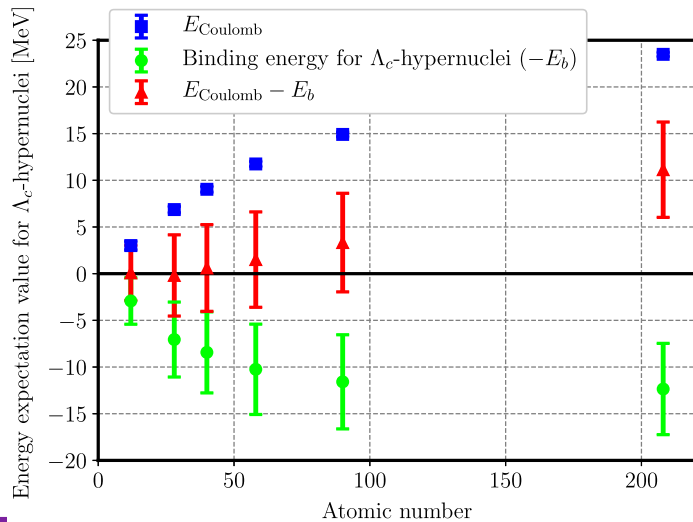
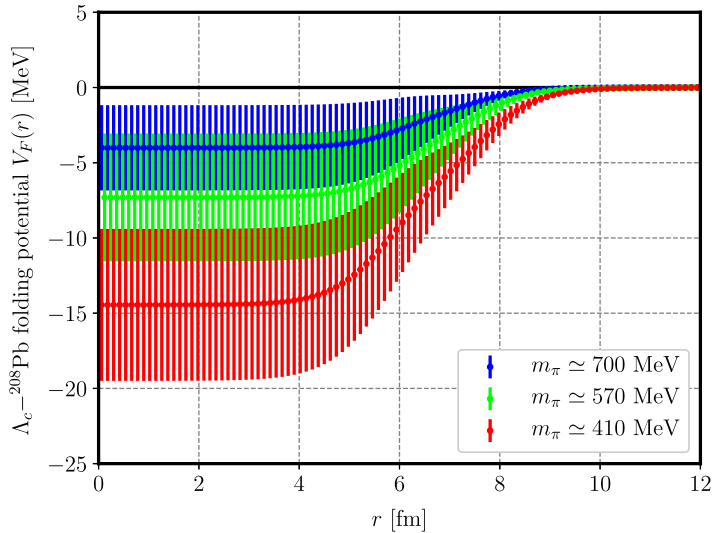
L. Wu, J.H., and H. Shen, *Phys. Rev. C* 101 (2020) 024303

L. Wu, J.H., and H. Shen, Phys. Rev. C 101 (2020) 024303

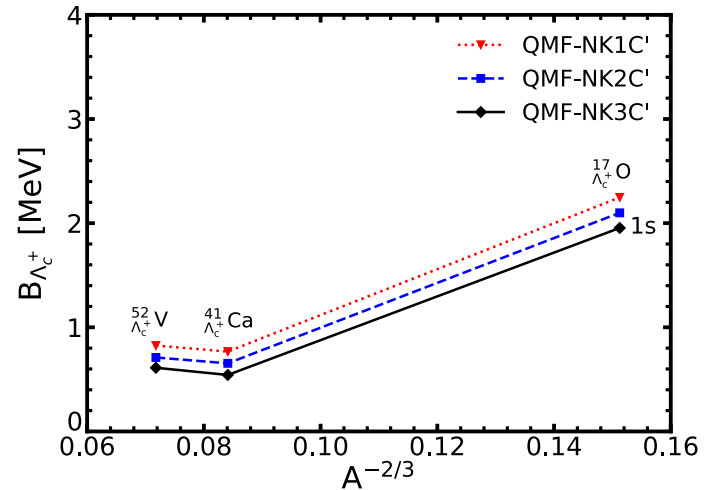
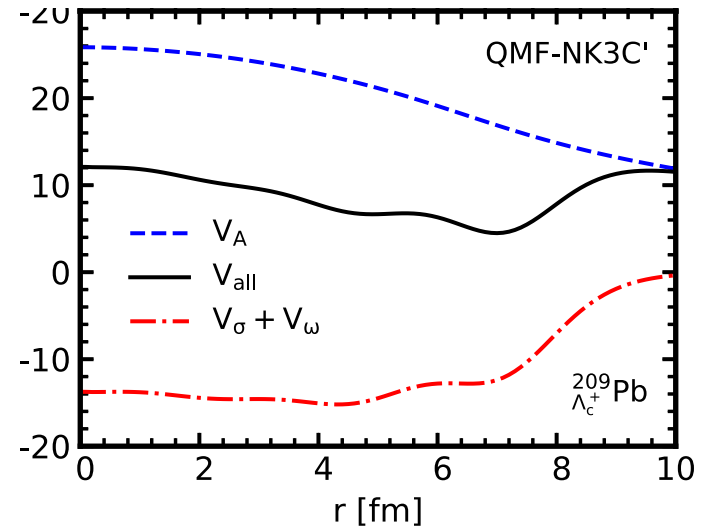
$$g_{\omega\Lambda_c^+} = 2/3g_{\omega N}$$



## HAL QCD results



## QMF results





## Binding energies and radius of charmed hypernuclei

L. Wu, J.H., and H. Shen, *Phys. Rev. C* 101 (2020) 024303

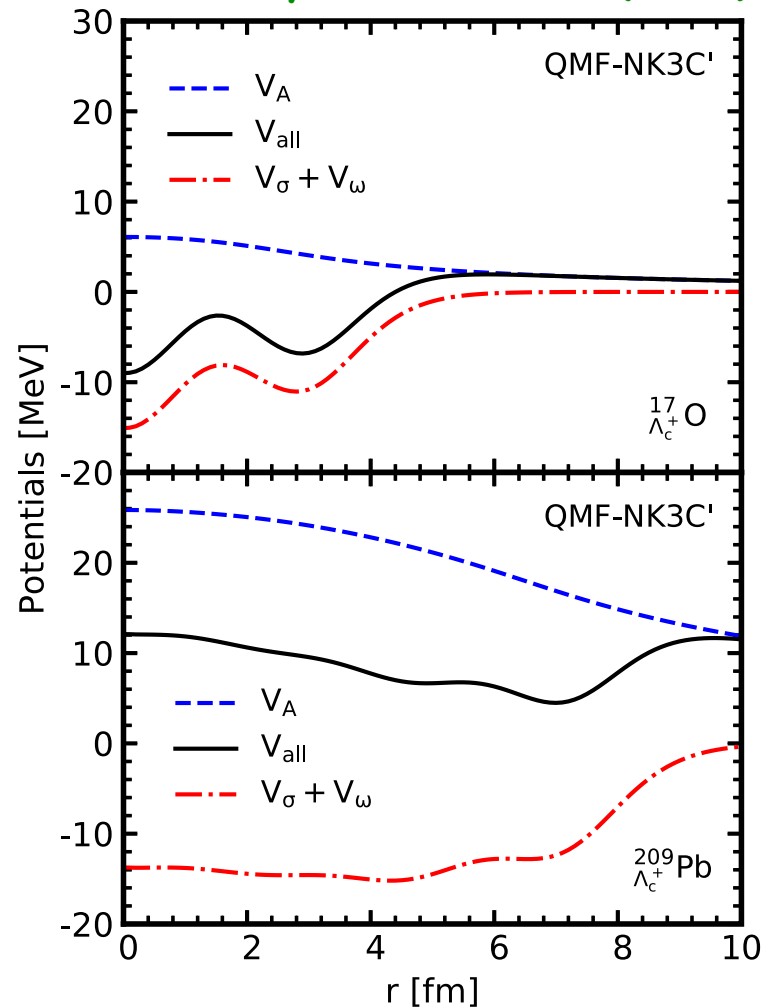
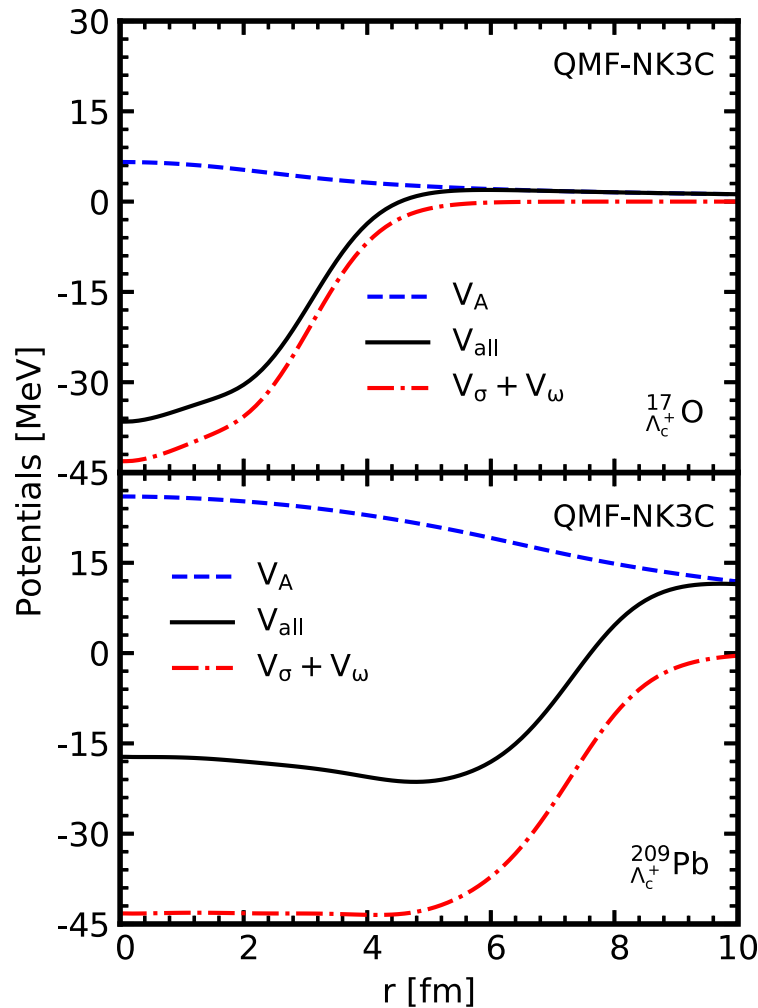
	QMF-NK3C					QMF-NK3C'				
	$-E/A$	$r_{\text{ch}}$	$r_p$	$r_n$	$r_{\Lambda_c^+}$	$-E/A$	$r_{\text{ch}}$	$r_p$	$r_n$	$r_{\Lambda_c^+}$
$^{16}\text{O}$	8.1377	2.7225	2.6042	2.5763		8.1377	2.7225	2.6042	2.5763	
$^{17}_{\Lambda_c^+}\text{O}$	9.1039	2.7298	2.6118	2.5797	1.8199	7.7937	2.7418	2.6244	2.5936	3.1746
$^{40}\text{Ca}$	8.5916	3.4562	3.3638	3.3141		8.5916	3.4562	3.3638	3.3141	
$^{41}_{\Lambda_c^+}\text{Ca}$	9.0333	3.4630	3.3708	3.3174	2.2599	8.4159	3.4692	3.3771	3.3252	3.8017
$^{51}\text{V}$	8.6403	3.6050	3.5200	3.6127		8.6403	3.6050	3.5200	3.6127	
$^{52}_{\Lambda_c^+}\text{V}$	9.0162	3.6086	3.5237	3.6123	2.3773	8.5047	3.6190	3.5343	3.6246	3.7366
$^{89}\text{Y}$	8.6990	4.2435	4.1724	4.2923		8.6990	4.2435	4.1724	4.2923	
$^{90}_{\Lambda_c^+}\text{Y}$	8.8925	4.2466	4.1755	4.2921	2.9105					
$^{139}\text{La}$	8.4276	4.8556	4.7954	4.9826		8.4276	4.8556	4.7954	4.9826	
$^{140}_{\Lambda_c^+}\text{La}$	8.5388	4.8565	4.7964	4.9812	3.5325					
$^{208}\text{Pb}$	7.8992	5.5037	5.4517	5.6898		7.8992	5.5037	5.4517	5.6898	
$^{209}_{\Lambda_c^+}\text{Pb}$	7.9623	5.5052	5.4532	5.6892	4.2618					

The  $\Lambda_c$  hyperon is inside of the nuclei with quark counting

The  $\Lambda_c$  hyperon is outside of the nuclei with lattice

## The potentials of charmed hyper nuclei

L. Wu, J.H., and H. Shen, Phys. Rev. C 101 (2020) 024303



Experiment	Event	Daughters	$B_{\Xi^-}$ [MeV]
E373	T1 [39]	${}^5_{\Lambda}\text{He} + {}^5_{\Lambda}\text{He} + {}^4\text{He} + \text{n}$	$-2.2 \pm 1.2$
E373	T2 KISO [24]	${}^{10}_{\Lambda}\text{Be} + {}^5_{\Lambda}\text{He}$	$3.87 \pm 0.21$ or $1.03 \pm 0.18$
E07	T006 IBUKI [25]	${}^{10}_{\Lambda}\text{Be} + {}^5_{\Lambda}\text{He}$	$1.27 \pm 0.21$
E373	T3 KINKA	${}^9_{\Lambda}\text{Be} + {}^5_{\Lambda}\text{He} + \text{n}$	$8.00 \pm 0.77$ or $4.96 \pm 0.77$
E07	T007	${}^9_{\Lambda}\text{Be} + {}^5_{\Lambda}\text{He} + \text{n}$	$-1.04 \pm 0.85$
E07	T010 IRRAWADDY	${}^5_{\Lambda}\text{He} + {}^5_{\Lambda}\text{He} + {}^4\text{He} + \text{n}$	$6.27 \pm 0.27$
E07	T011	${}^5_{\Lambda}\text{He} + {}^5_{\Lambda}\text{He} + {}^4\text{He} + \text{n}$	$0.90 \pm 0.62$

## The conflation method

$$f(x) = \frac{1}{\sigma_i \sqrt{2\pi}} \exp \left[ -\frac{(x - m_i)^2}{2\sigma_i^2} \right]$$

$m$  is the mean value  
 $\sigma$  is the standard deviation



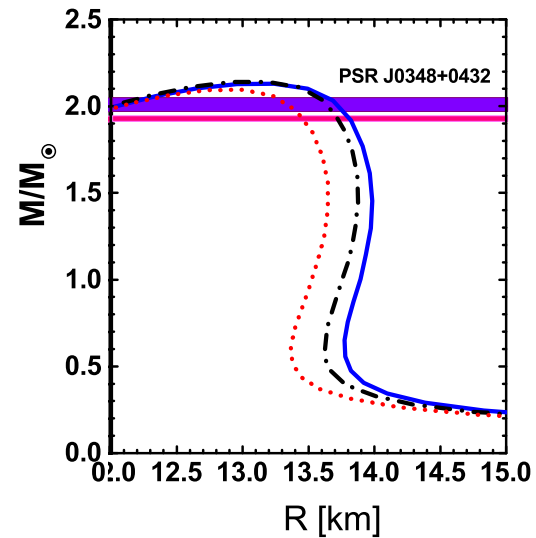
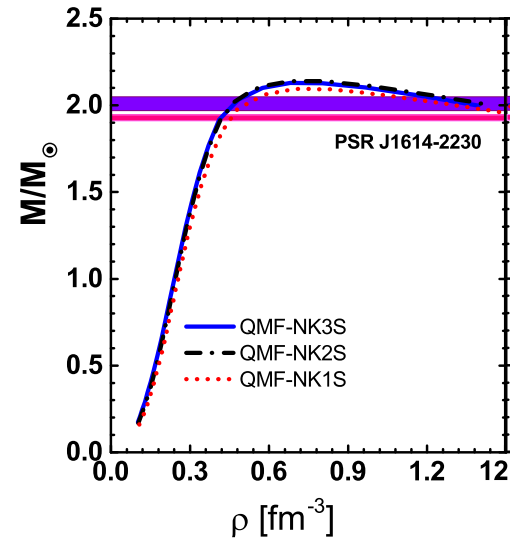
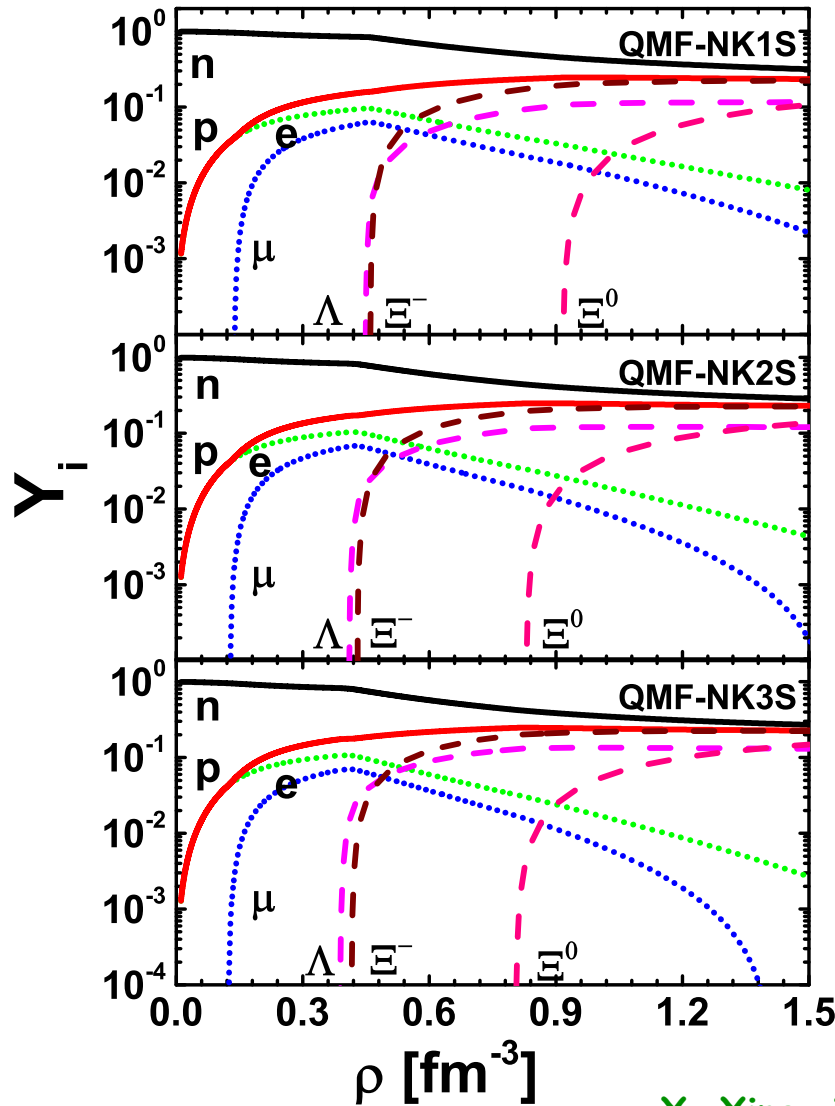
$$m_t = \frac{\sum_i^n m_i \sigma_i^{-2}}{\sum_i^n \sigma_i^{-2}}^{-1},$$

$$\sigma_t = \left( \sum_i^n \sigma_i^{-2} \right)^{-1/2}.$$

# The strangeness in neutron star



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X. Xing, J.H., and H. Shen, Phys. Rev. C 95 (2017) 054310



## Outline

1 Introduction

2 The quark mean field model

3 The strangeness with QMF model

4 The summary and perspective

The effects of chiral dynamics and gluons are introduced into the QMF model.

The improved QMF model was applied to the finite nuclei,  $\Lambda$ ,  $\Xi^0$ ,  $\Xi^-$ ,  $\Lambda_c$  hypernuclei and neutron star, which can describe the many-body strangest system very well.

The EMC effect will be studied within QMF model.