



Hyperon-nucleon scattering at BESIII

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Henan University

Workshop on Hyperon Physics

2024.04 in Huizhou

Outline

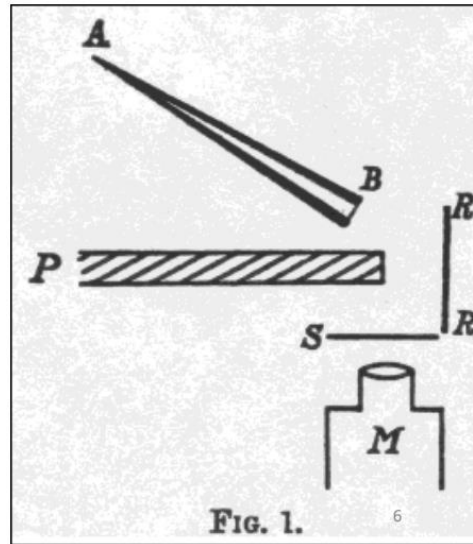
- Introduction
- Recent experimental measurements on hyperon-nucleon scattering
CLAS, JPARC-E40, **BESIII**
- Summary

Scattering experiments of particle beams bombarding target materials

1911



$\alpha + Au$



➔ Nuclear structure
model of atom

1919

$\alpha + N$



Observation of proton

1932

$\alpha + Be$



Observation of neutron



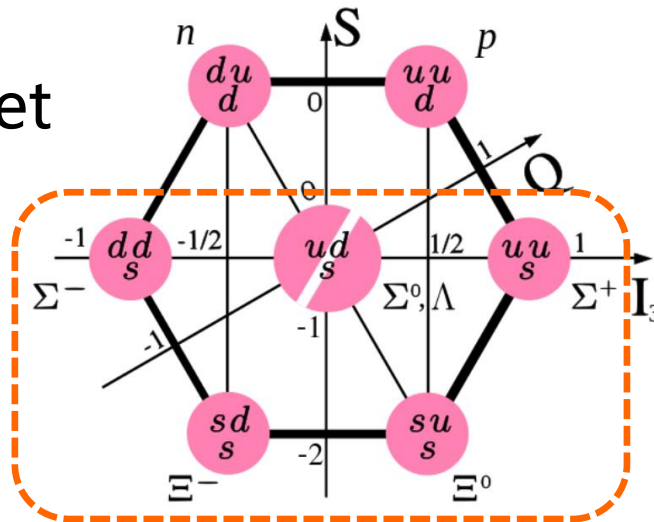
James Chadwick

Scattering experiment must have **particle source**,
target material, and detector.

Hyperon source

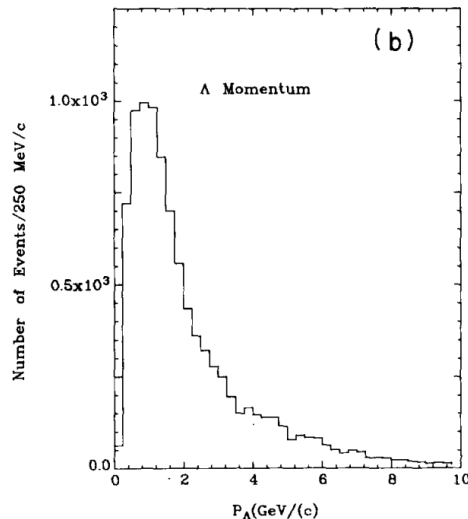
Baryon octet

One of main goals of nuclear physics is to understand baryon-baryon interaction in a unified perspective

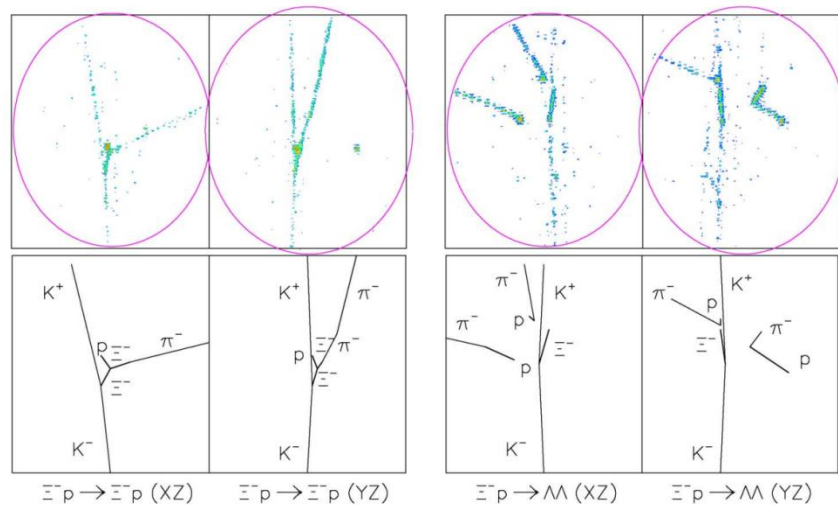


Limited by availability and short-lifetime of hyperon beams

- Hyperons are obtained by bombarding hydrogen bubble chamber or scintillating fiber target with K^- .



Nucl. Phys. B 125, 29 (1977)



Phys. Lett. B 633, 214 (2006)

Hyperon source

- Hyperons are obtained by bombarding hydrogen bubble chamber or scintillating fiber target with K^- .
- Intensity of hyperon beams is low, experimental measurements are scarce and have large uncertainty.
- No anti-hyperon source.

Reaction	Number of events	
$\Lambda p \rightarrow \Lambda p$ (elastic)	584	(1)
$\Lambda p \rightarrow \Sigma^- p \pi^+$	132	(2)
$\Lambda p \rightarrow \Sigma^+ p \pi^-$	60	(3)
$\Lambda p \rightarrow \Lambda p \pi^+ \pi^-$	181	(4)
$\Lambda p \rightarrow \Sigma^0 p$	35	(5)
various $\Xi^0 p$ interactions	25	

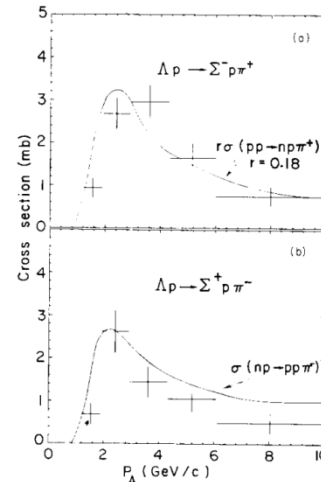
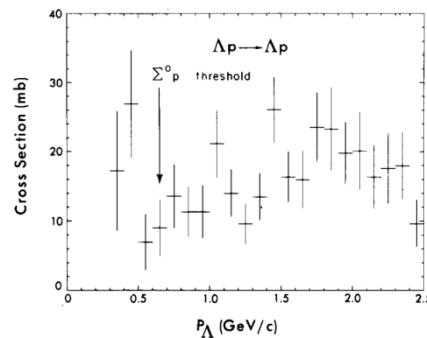
Phys. Lett. B 38, 123 (1972)

reaction	events *	signature	cross-section events **	cross-section (mb)
$\Xi^0 + p \rightarrow \Xi^0 + p$	2	K, Λ	1	8
$\Xi^0 + p \rightarrow \Lambda + \Sigma^+$	6	Λ	4	24
$\Xi^0 + p \rightarrow \Sigma^0 + \Sigma^+$	1	Λ	1	6
$\Xi^0 + p \rightarrow \pi^+ + \Lambda + \Lambda$	1	K, Λ	1	6
$\Xi^0 + p \rightarrow \pi^0 + \Lambda + \Sigma^+$	1	Λ	1	6
$\Xi^0 + p \rightarrow \pi^+ + \Xi^- + p$	1	K or Λ	1	5
$\Xi^0 + p \rightarrow \pi^+ + \pi^+ + \Xi^- + n$	1	K, Λ	1	6
$\Xi^0 + p \rightarrow \Xi^- + p$	2	Λ	2	8
$\Xi^0 + p \rightarrow \Sigma^- + \Sigma^+$	1	K	1	4
$\Xi^0 + p \rightarrow \Sigma^- + K^0 + p$	1	K	1	4

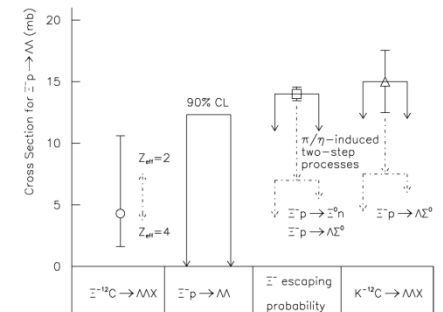
Phys. Lett. B 32, 720 (1970)

Reaction	Momentum interval (GeV/c)	Number of events	σ (mb)
$\Lambda p \rightarrow \text{all}$	0.5 - 1.0	25.8 ± 6.2	
	1.0 - 1.5	31.3 ± 6.5	
	1.5 - 2.0	42.8 ± 7.1	
	2.0 - 2.5	37.5 ± 7.2	
	2.5 - 3.0	34.1 ± 8.3	
	3.0 - 4.0	41.8 ± 10.0	
$\Lambda p \rightarrow \Lambda p$	0.5 - 1.0	20	22.2 ± 5.0
	1.0 - 1.5	21	12.9 ± 2.8
	1.5 - 2.0	37	22.0 ± 3.6
	2.0 - 2.5	28	16.1 ± 3.1
	2.5 - 3.0	12	11.0 ± 3.2
	3.0 - 4.0	13	12.5 ± 3.4
$\Lambda p \rightarrow \Sigma^0$	0.66 - 4.0	11	1.5 ± 0.5
	0.88 - 4.0	29	4.1 ± 0.8
	1.36 - 4.0	12	1.9 ± 0.6
$\Sigma^+ p \rightarrow \Sigma^+ p$	0.5 - 1.5	10	31.2 ± 10.1
	1.5 - 2.5	8	18.7 ± 6.6
	2.5 - 4.0	4	15.3 ± 7.8
$\Sigma^- p \rightarrow \Sigma^- p$	0.5 - 1.5	6	13.2 ± 4.7
	1.5 - 2.5	11	13.9 ± 4.1
	2.5 - 4.0	4	7.5 ± 3.8
$\Xi^- p \rightarrow \Xi^- p$	1.0 - 4.0	6	13 ± 6
	1.0 - 4.0	4	19 ± 10

Nucl. Phys. B 125, 29 (1977)

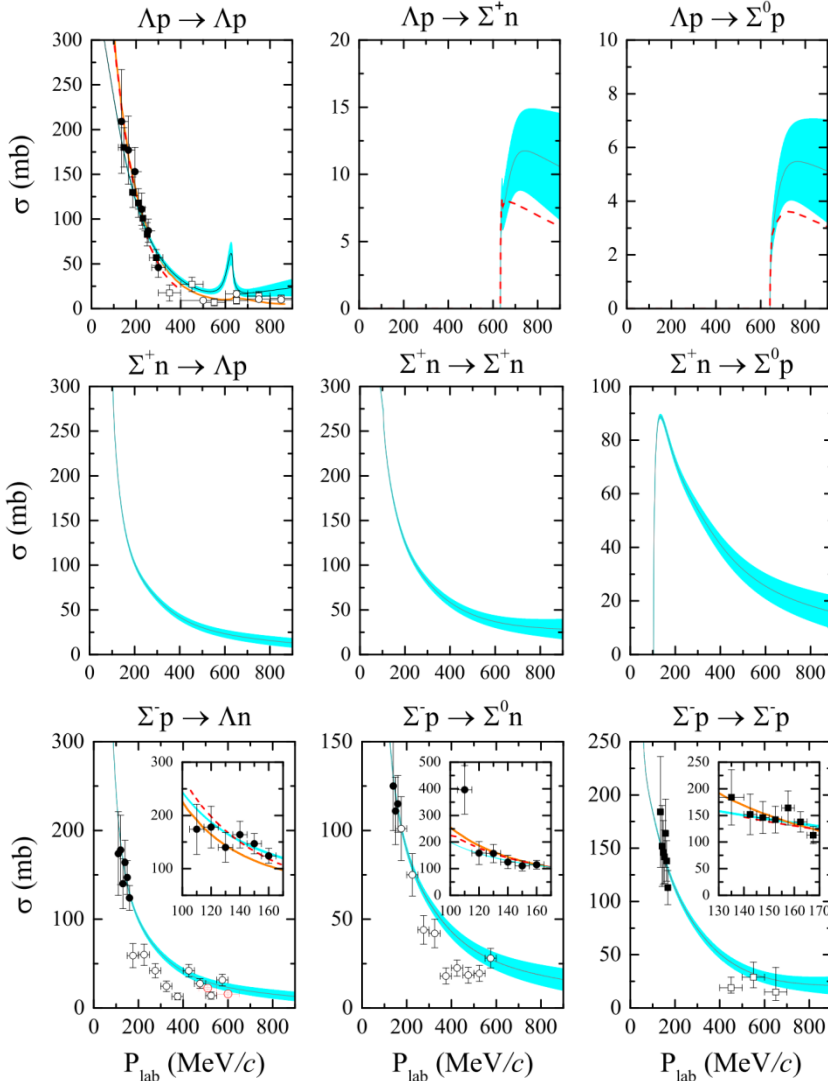


Phys. Lett. B 633, 214 (2006)

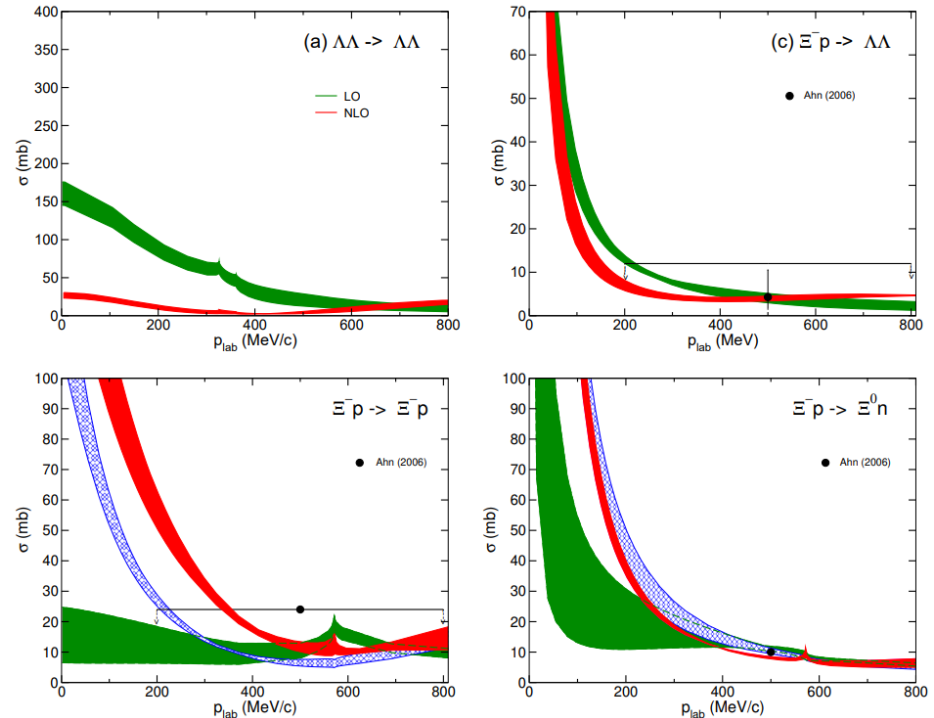


Theory of hyperon-nucleon (YN) interaction has large uncertainty due to lack of relevant measurements

Phys. Rev. C 105, 035203 (2022)



LO : H. Polinder, J.H., U.-G. Meißner, PLB 653 (2007) 29
 NLO16: J.H., U.-G. Meißner, S. Petschauer, NPA 954 (2016) 273
 NLO19: J.H., U.-G. Meißner, EPJA 55 (2019) 23



"Hyperon puzzle" of neutron stars

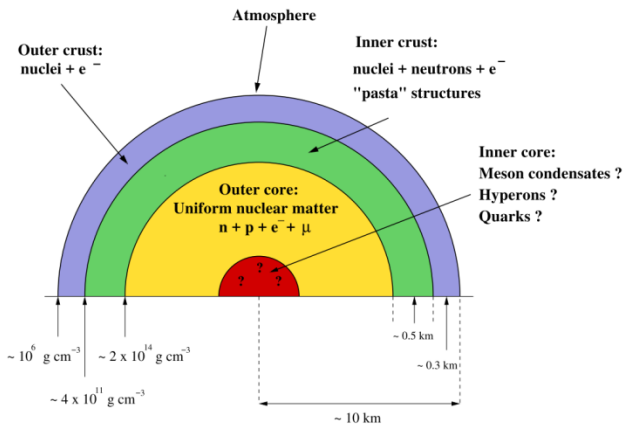
- Hyperons are believed to be appeared in inner core of neutron stars.

$$B_1 \rightarrow B_2 + l + \bar{\nu}_l, B_2 + l \rightarrow B_1 + \nu_l$$

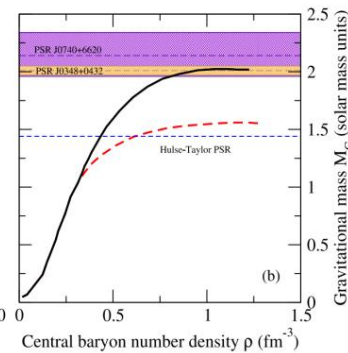
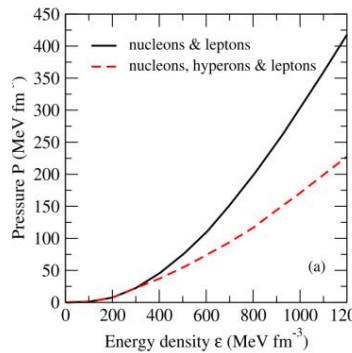
$$n \rightarrow p + e^- + \bar{\nu}_e, p + e^- \rightarrow n + \nu_e \quad \Sigma^- \rightarrow n + e^- + \bar{\nu}_e, n + e^- \rightarrow \Sigma^- + \nu_e$$

$$\Lambda \rightarrow p + e^- + \bar{\nu}_e, p + e^- \rightarrow \Lambda + \nu_e \quad \Xi^- \rightarrow \Lambda + e^- + \bar{\nu}_e, \Lambda + e^- \rightarrow \Xi^- + \nu_e$$

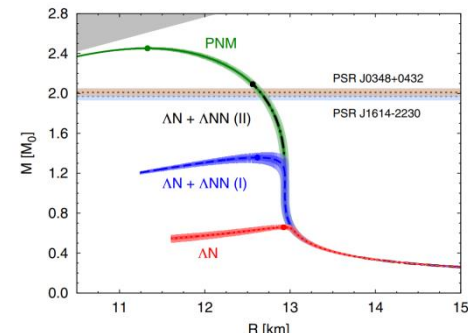
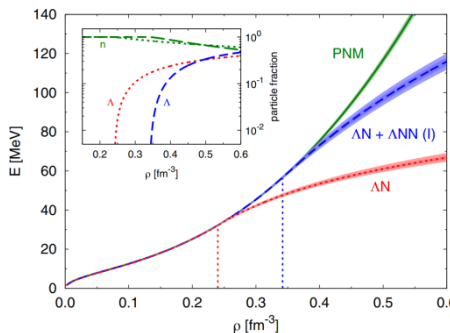
- Appearance of hyperons softens equation of state, lead to maximum mass that neutron stars can sustain is less than mass of already-observed neutron stars.
- A repulsive force is introduced to stiffen equation of state in theory, such as a combination of ΛN and ΛNN interactions. Study of hyperon-nucleon interaction is crucial to solve "hyperon puzzle" of neutron stars.



Phys. Rev. Lett. 114, 092301 (2015)



Prog. Part. Nucl. Phys. 112, 103770 (2020)

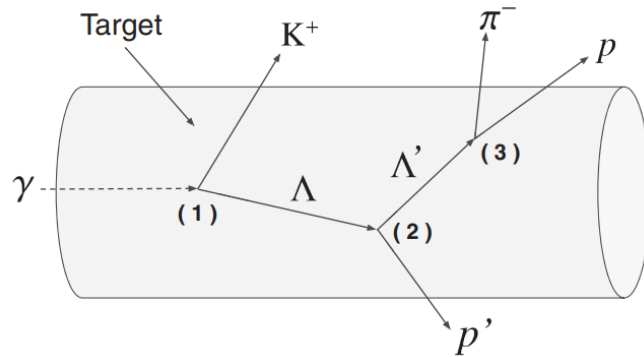


Some recent experimental results on hyperon-nucleon scattering

PHYSICAL REVIEW LETTERS **127**, 272303 (2021)

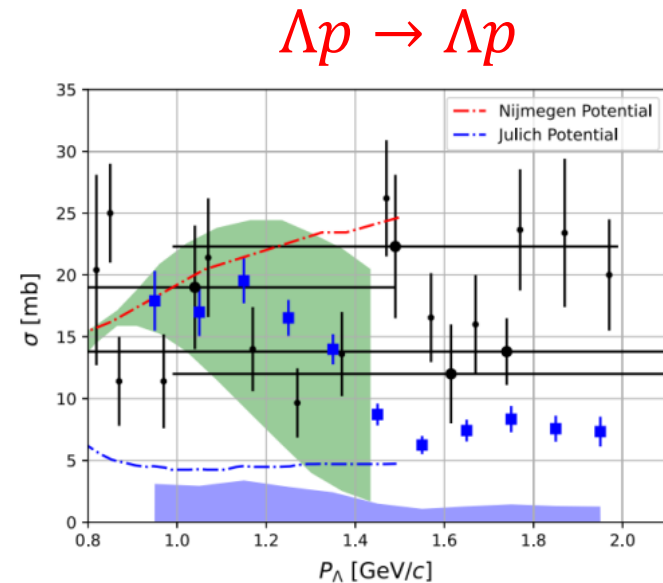
(CLAS Collaboration)

Improved Λp Elastic Scattering Cross Sections between 0.9 and 2.0 GeV/c as a Main Ingredient of the Neutron Star Equation of State



$$\sigma(p_\Lambda) = \frac{Y(p_\Lambda)}{A(p_\Lambda) \times \mathcal{L}(p_\Lambda) \times \Gamma}$$

$$\mathcal{L}(p_\Lambda) = \frac{N_A \times \rho_T \times l}{M} N_\Lambda(p_\Lambda)$$

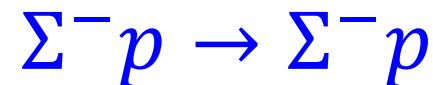
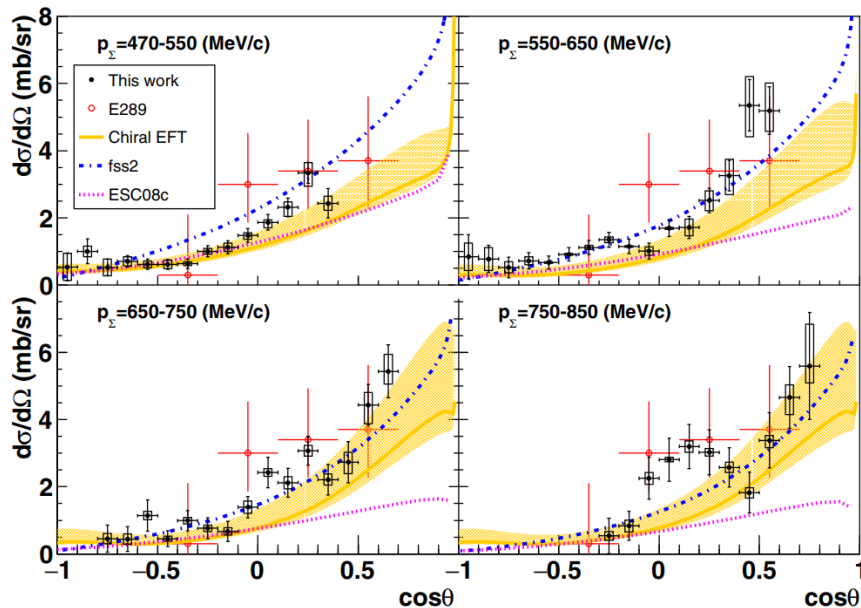
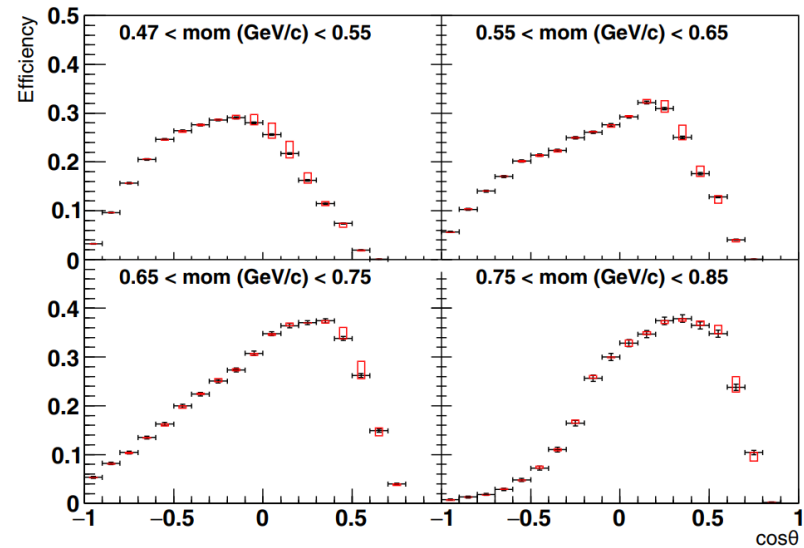
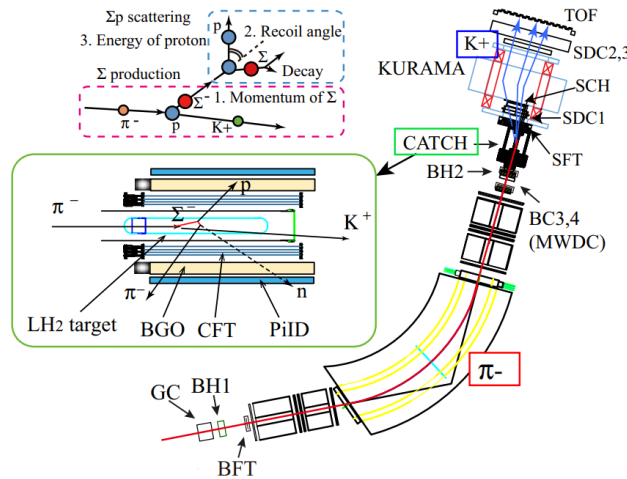


This is the first data on this reaction since the 1970s.

Some recent experimental results on hyperon-nucleon scattering

J-PARC E40 Collaboration

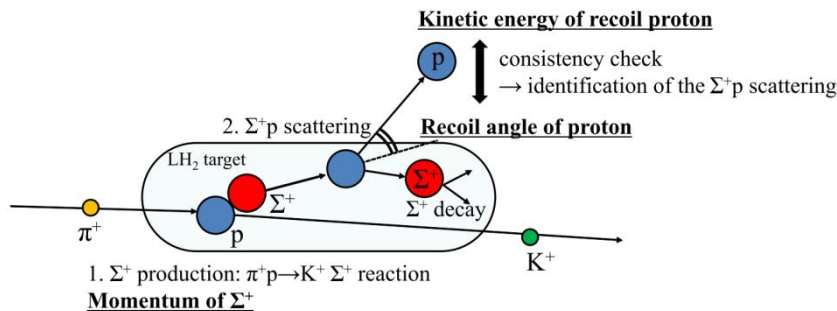
Phys. Rev. C 104, 045204 (2021)



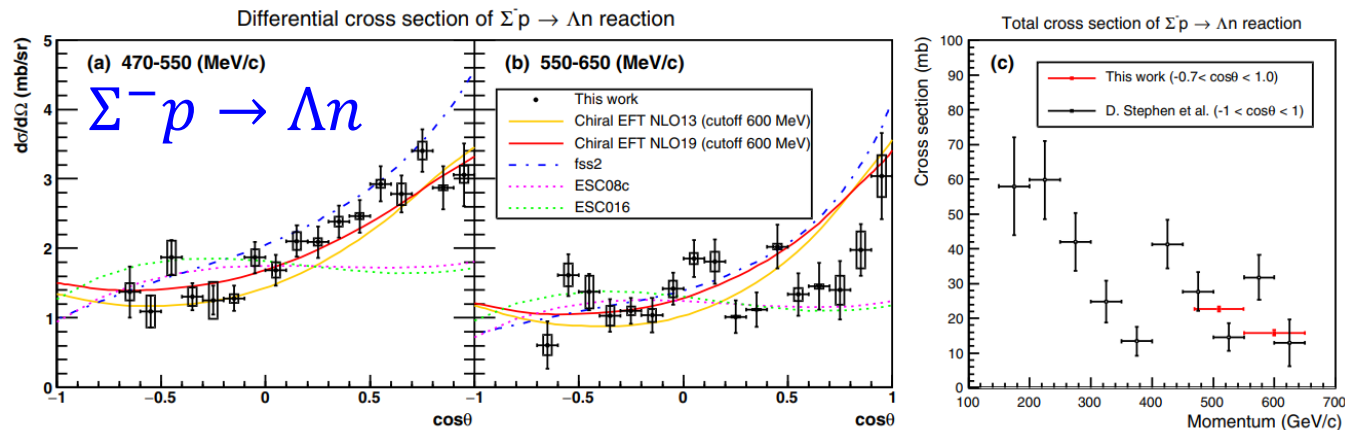
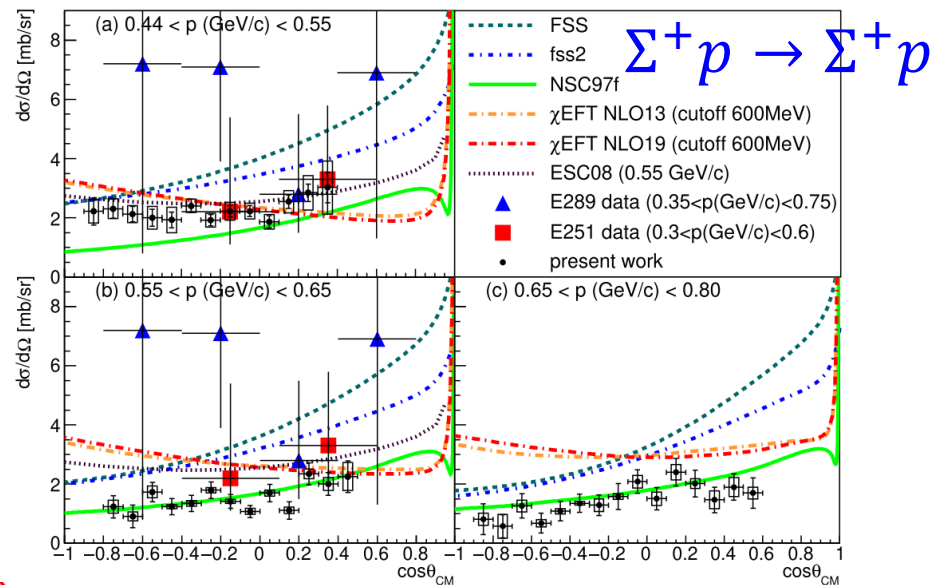
Some recent experimental results on hyperon-nucleon scattering

J-PARC E40 Collaboration

PTEP 2022, 093D01 (2022)

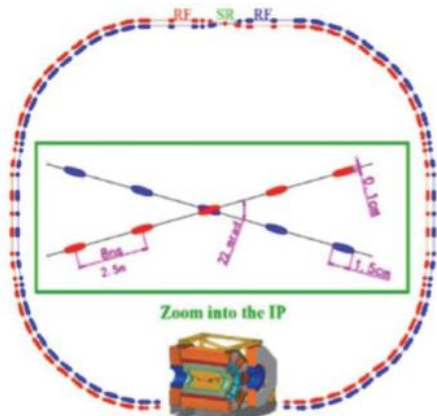


Phys. Rev. Lett. 128, 072501 (2022)



Beijing Electron Positron Collider II (BEPCII) and Beijing Spectrometer III (BESIII)

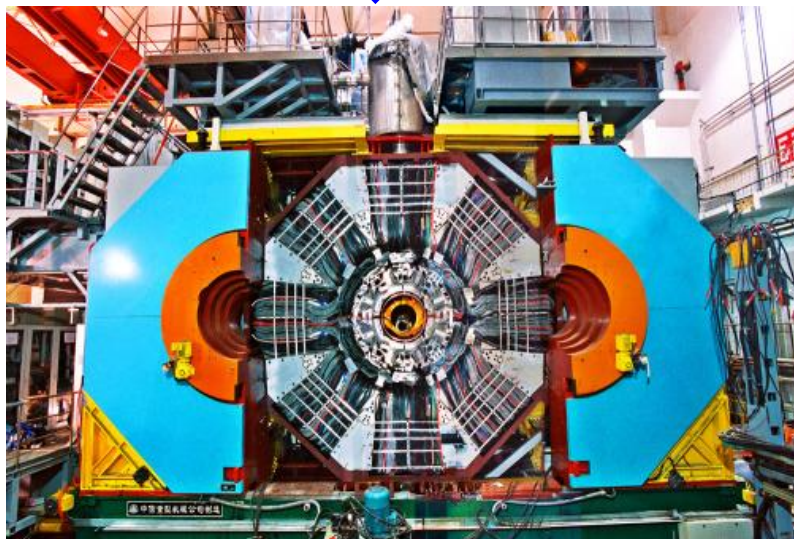
tau-charm energy region



Storage ring



Linear accelerator

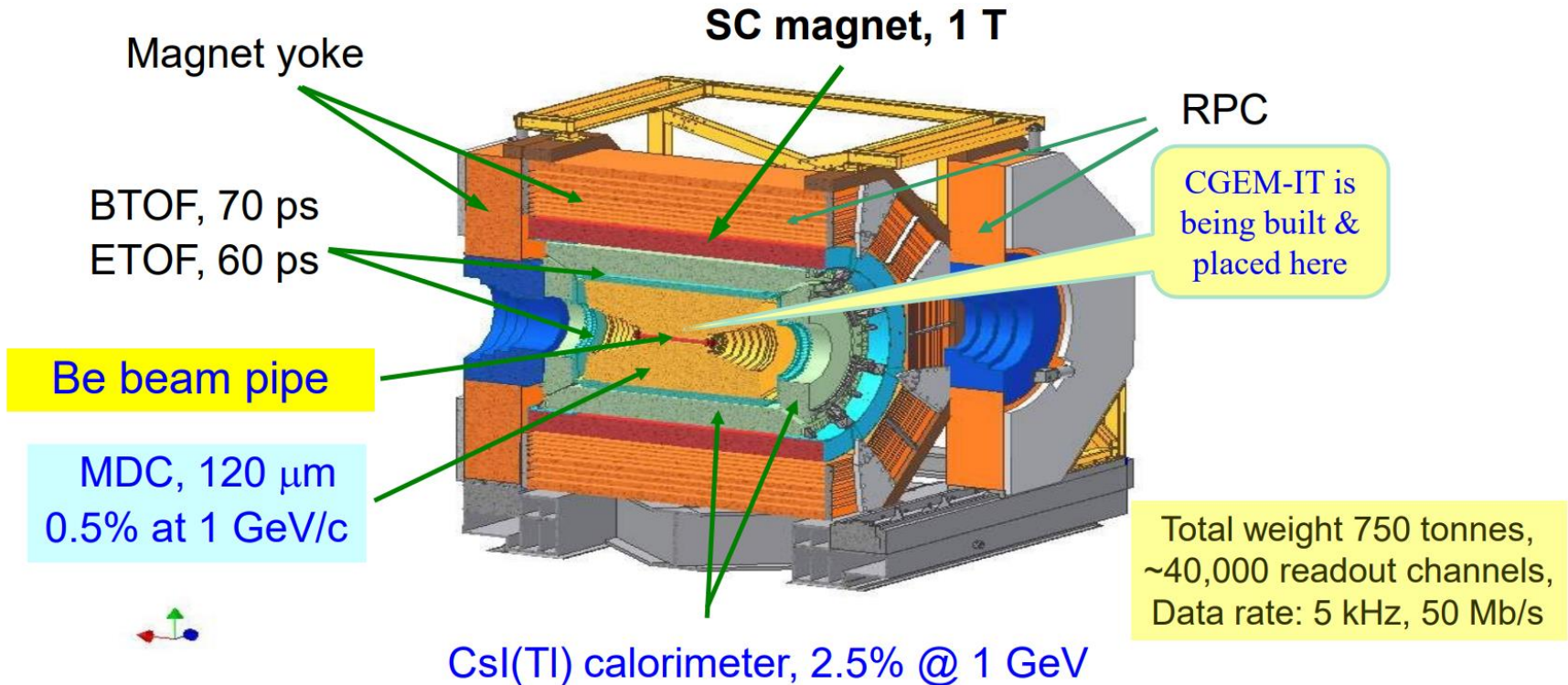


BESIII detector



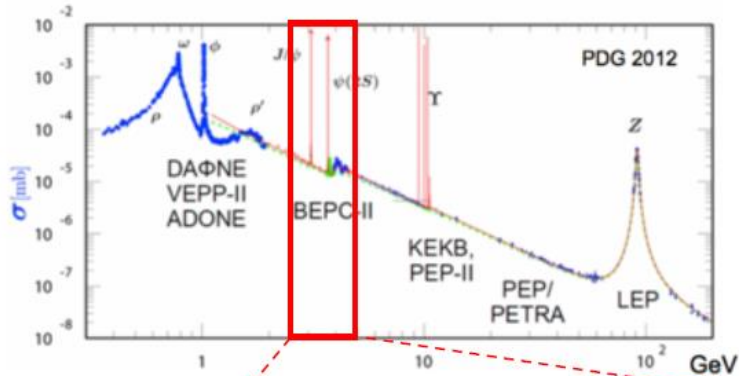
BEPCII

BESIII detector



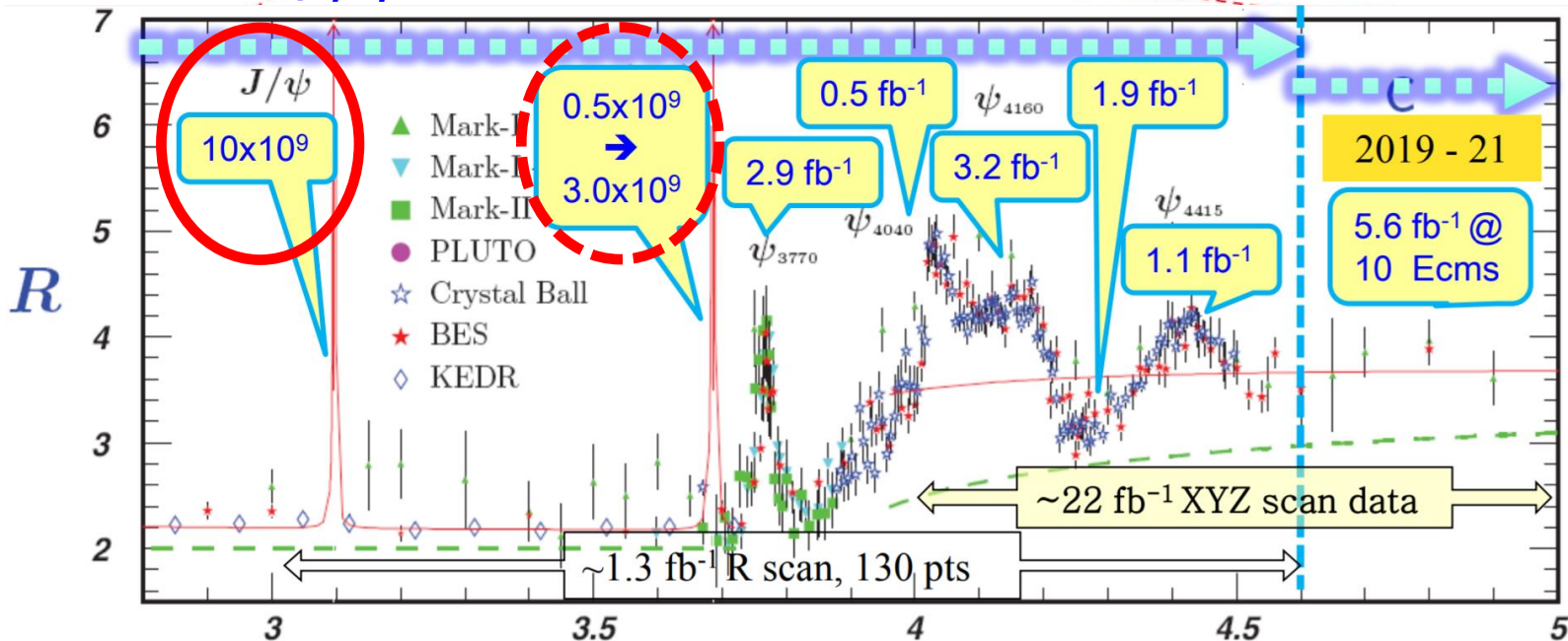
Has been in full operation since 2008,
all subdetectors are in very good status!

BESIII data samples

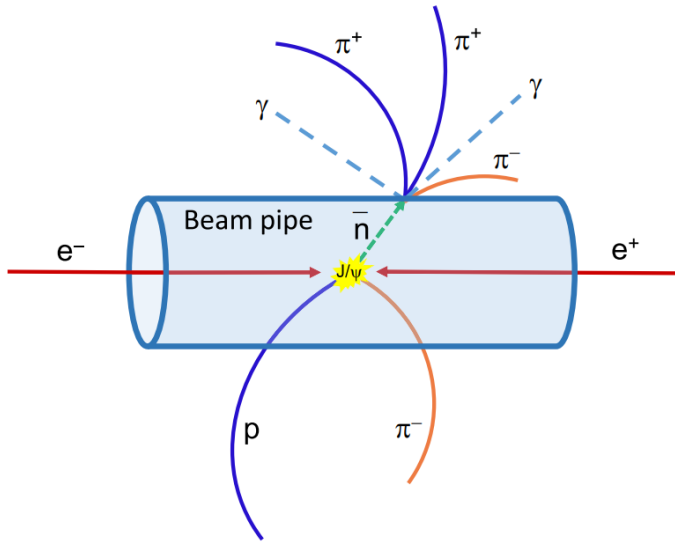


BESIII has collected the largest data samples of the J/ψ and $\psi(3686)$ in the world, and $> 20 \text{ fb}^{-1}$ above 4.0 GeV in total.

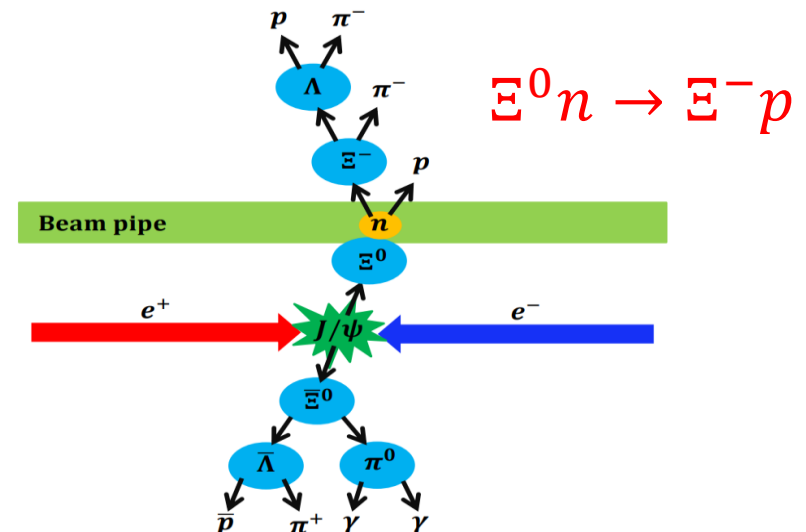
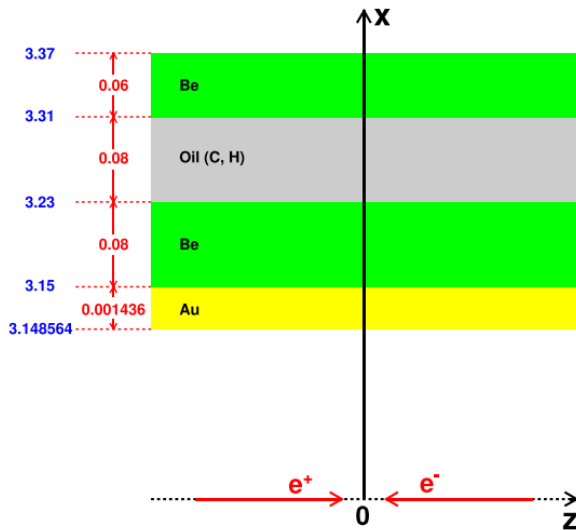
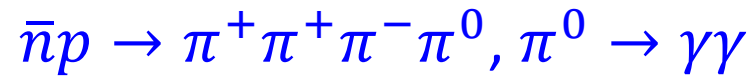
10 billion J/ψ events



Experimental study on particle targeting at BESIII



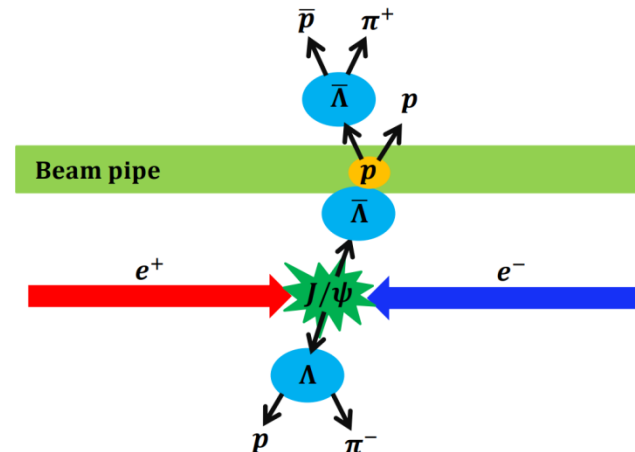
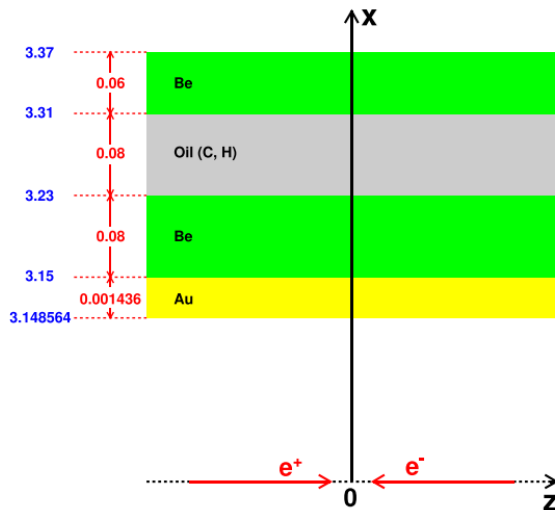
Phys. Rev. Lett. 127, 012003 (2021)
arXiv: 2209.12601



particle source: hyperon from J/ψ decays
target material: beam pipe
detector: BESIII detector

Experimental study on particle targeting at BESIII

Hyperon/Antihyperon	τ ($\times 10^{-10}$ s)	Decay mode	\mathcal{B} ($\times 10^{-3}$)	P (GeV/c)	E_{cm} (GeV)	N ($\times 10^5$)	N^{bp} ($\times 10^5$)
$\Lambda/\bar{\Lambda}$	2.63	$J/\psi \rightarrow \Lambda\bar{\Lambda}$	1.89	1.074	2.24	189	103
$\Sigma^+/\bar{\Sigma}^-$	0.80	$J/\psi \rightarrow \Sigma^+\bar{\Sigma}^-$	1.07	0.992	2.28	107	15
$\Xi^0/\bar{\Xi}^0$	2.90	$J/\psi \rightarrow \Xi^0\bar{\Xi}^0$	1.17	0.818	2.35	117	51
$\Xi^-/\bar{\Xi}^+$	1.64	$J/\psi \rightarrow \Xi^-\bar{\Xi}^+$	0.97	0.807	2.35	97	23
$\Lambda/\bar{\Lambda}$	2.63	$\psi(2S) \rightarrow \Lambda\bar{\Lambda}$	0.38	1.467	2.36	11	7
$\Sigma^+/\bar{\Sigma}^-$	0.80	$\psi(2S) \rightarrow \Sigma^+\bar{\Sigma}^-$	0.24	1.408	2.40	7	2
$\Xi^0/\bar{\Xi}^0$	2.90	$\psi(2S) \rightarrow \Xi^0\bar{\Xi}^0$	0.23	1.291	2.47	7	4
$\Xi^-/\bar{\Xi}^+$	1.64	$\psi(2S) \rightarrow \Xi^-\bar{\Xi}^+$	0.29	1.284	2.47	9	3
$\Omega^-/\bar{\Omega}^+$	0.82	$\psi(2S) \rightarrow \Omega^-\bar{\Omega}^+$	0.06	0.774	2.67	2	0.1



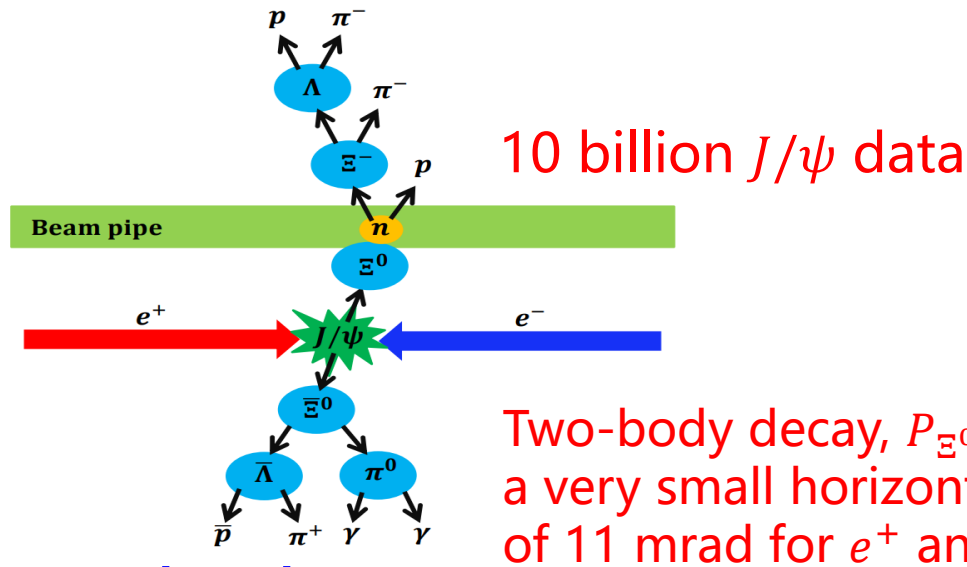
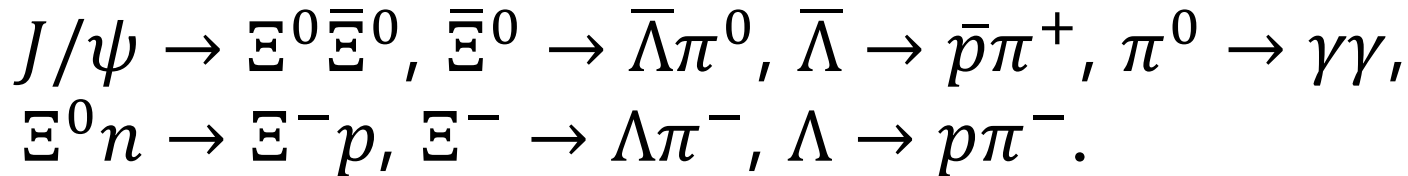
New results on hyperon-nucleon scattering at BESIII

- **First Study of Reaction $\Xi^0 n \rightarrow \Xi^- p$ Using Ξ^0 -Nucleus Scattering at an Electron-Positron Collider**
Phys. Rev. Lett. 130, 251902 (2023)
- **First measurement of ΛN inelastic scattering with Λ from $e^+ e^- \rightarrow J/\psi \rightarrow \Lambda \bar{\Lambda}$**
arXiv:2310.00720
- **First study of antihyperon-nucleon scattering $\bar{\Lambda} p \rightarrow \bar{\Lambda} p$ and measurement of $\Lambda p \rightarrow \Lambda p$**
arXiv:2401.09012

Study of $\Xi^0 n \rightarrow \Xi^- p$

Phys. Rev. Lett. 130, 251902 (2023)

Reaction chain :

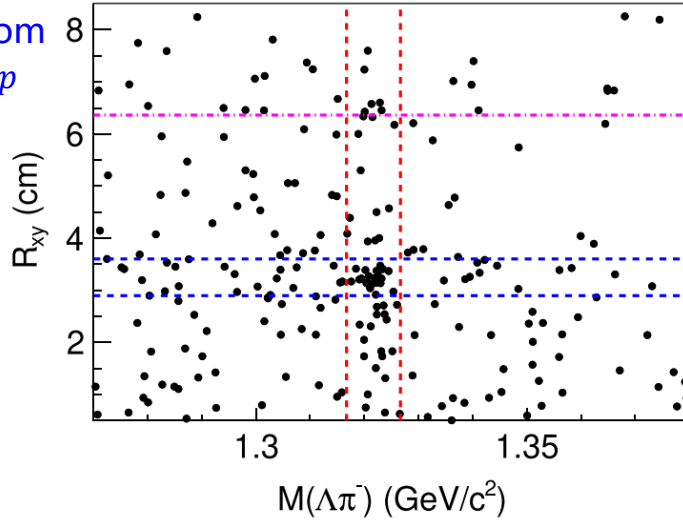


Analysis method :

Using $\bar{\Xi}^0$ to tag the event and requiring the recoiling mass in Ξ^0 region. Then reconstructing Ξ^- and p in the signal side.

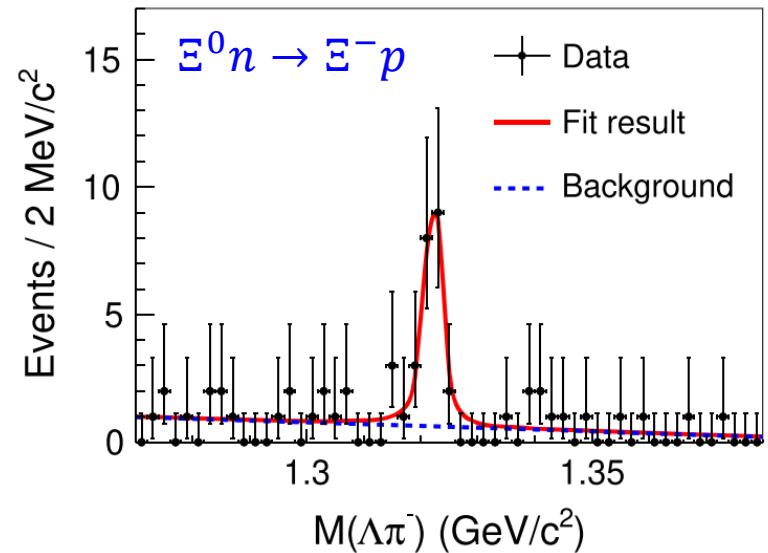
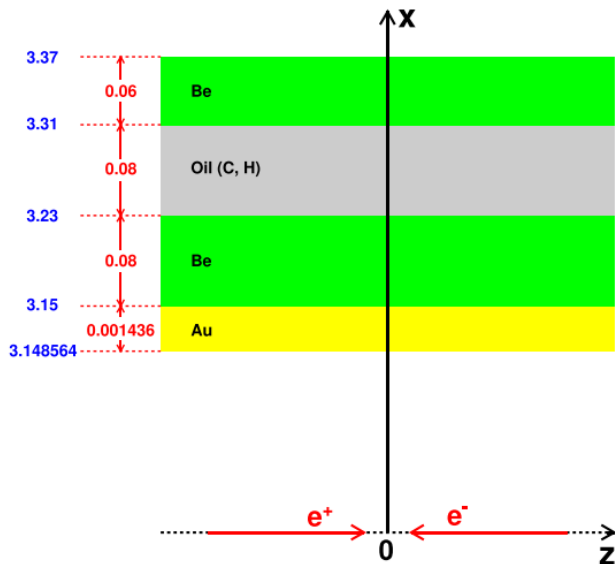
Study of $\Xi^0 n \rightarrow \Xi^- p$

R_{xy} is distance from reconstructed $\Xi^- p$ vertex to z axis



Inner wall of MDC

Beam pipe



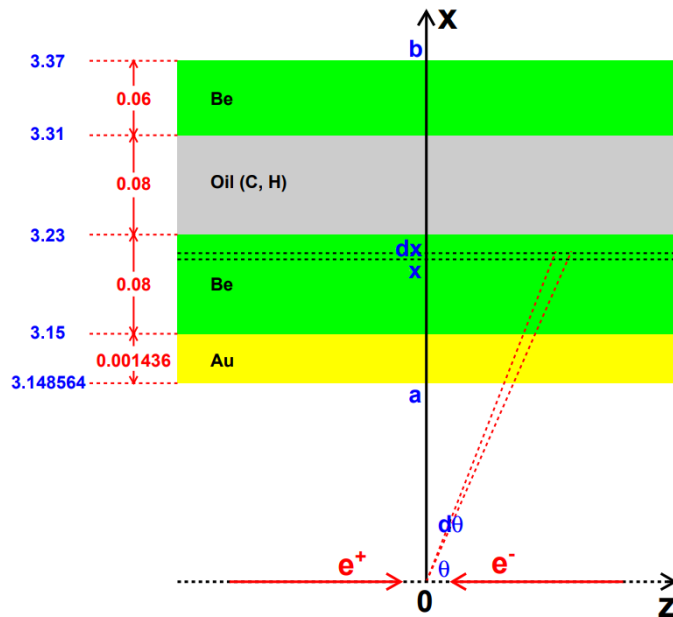
$$N = 22.9 \pm 5.5$$

$$S = 7.1\sigma$$

Cross section of $\Xi^0 + {}^9\text{Be} \rightarrow \Xi^- + p + {}^8\text{Be}$

$$\sigma(\Xi^0 + {}^9\text{Be} \rightarrow \Xi^- + p + {}^8\text{Be}) = \frac{N^{\text{sig}}}{\epsilon \mathcal{B} \mathcal{L}_{\text{eff}}}$$

$$\mathcal{L}_{\text{eff}} = \frac{N_{J/\psi} \mathcal{B}_{J/\psi}}{2 + \frac{2}{3}\alpha} \int_a^b \int_0^\pi (1 + \alpha \cos^2 \theta) e^{-\frac{x}{\sin \theta \beta \gamma L}} N(x) C(x) d\theta dx$$



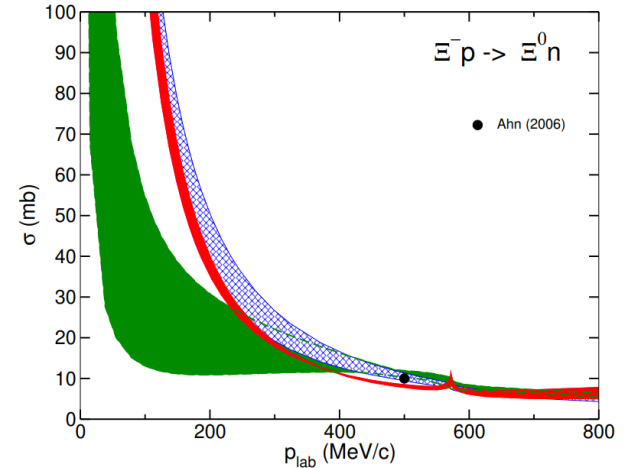
pure surface process assumption
(proportional to number of neutrons)

Parameter	Result
N^{sig}	22.9 ± 5.5
ϵ	1.873%
\mathcal{B}	$(40.114 \pm 0.444)\%$ [53]
$N_{J/\psi}$	$(1.0087 \pm 0.0044) \times 10^{10}$ [46]
$\mathcal{B}_{J/\psi}$	$(0.117 \pm 0.004)\%$ [53]
α	0.514 ± 0.016 [56]
L	(8.69 ± 0.27) cm [53]
E_{beam}	1.5485 GeV
m_{Ξ^0}	(1.31486 ± 0.00020) GeV/ c^2 [53]
a	3.148564 cm [45]
b	3.37 cm [45]
$N(x)$	$\begin{cases} 5.91 \times 10^{22} \text{ cm}^{-3}, & 3.148564 \leq x \leq 3.15 \text{ cm} \\ 1.24 \times 10^{23} \text{ cm}^{-3}, & 3.15 < x \leq 3.23 \text{ cm} \\ 3.45 \times 10^{22} \text{ cm}^{-3}, & 3.23 < x \leq 3.31 \text{ cm} \\ 1.24 \times 10^{23} \text{ cm}^{-3}, & 3.31 < x \leq 3.37 \text{ cm} \end{cases}$
$C(x)$	$\begin{cases} 8.437(23.6), & 3.148564 \leq x \leq 3.15 \text{ cm} \\ 1.000(1.00), & 3.15 < x \leq 3.23 \text{ cm} \\ 1.090(1.20), & 3.23 < x \leq 3.31 \text{ cm} \\ 1.000(1.00), & 3.31 < x \leq 3.37 \text{ cm} \end{cases}$

Study of $\Xi^0 n \rightarrow \Xi^- p$

The measured cross section of the reaction process $\Xi^0 + {}^9\text{Be} \rightarrow \Xi^- + p + {}^8\text{Be}$ is $\sigma(\Xi^0 + {}^9\text{Be} \rightarrow \Xi^- + p + {}^8\text{Be}) = (22.1 \pm 5.3_{\text{stat}} \pm 4.5_{\text{sys}})$ mb at $P_{\Xi^0} \approx 0.818$ GeV/c.

If we take the effective number of reaction neutrons in ${}^9\text{Be}$ nucleus as 3, the cross section of $\Xi^0 n \rightarrow \Xi^- p$ for single neutron is determined to be $\sigma(\Xi^0 n \rightarrow \Xi^- p) = (7.4 \pm 1.8_{\text{stat}} \pm 1.5_{\text{sys}})$ mb, consistent with theoretical predictions.

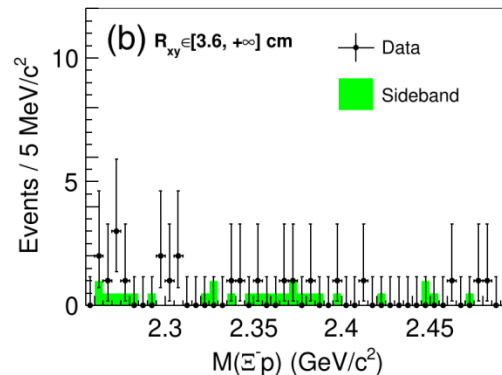
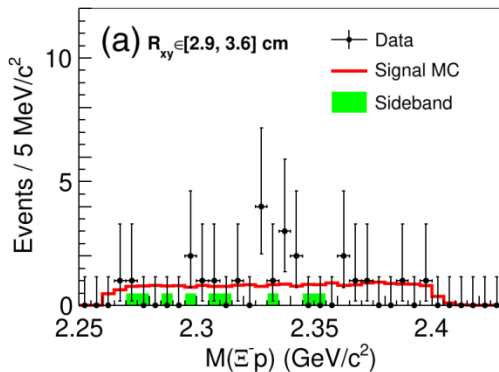


LO : H. Polinder, J.H., U.-G. Meißner, PLB 653 (2007) 29

NLO16: J.H., U.-G. Meißner, S. Petschauer, NPA 954 (2016) 273

NLO19: J.H., U.-G. Meißner, EPJA 55 (2019) 23

No significant H-dibaryon signals are seen



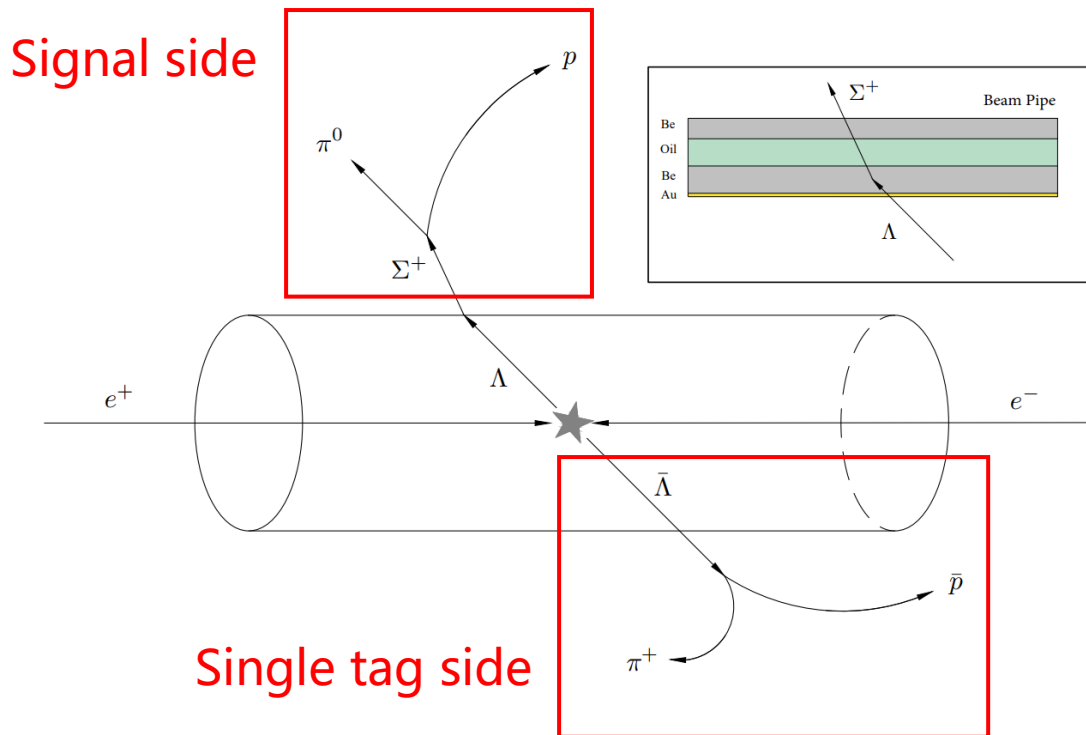
This work is the first study of hyperon-nucleon interaction in electron-positron collisions, and opens up a new direction for such research.

Study of $\Lambda N \rightarrow \Sigma^+ X$

arXiv:2310.00720

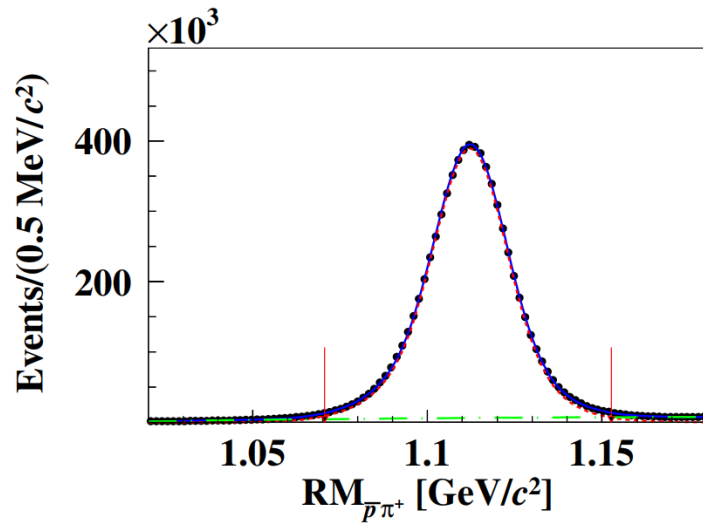
Reaction chain :

$$J/\psi \rightarrow \Lambda \bar{\Lambda}, \quad \bar{\Lambda} \rightarrow \bar{p} \pi^+, \quad \Lambda + N(\text{nucleus}) \rightarrow \Sigma^+ + X(\text{anything}), \\ \Sigma^+ \rightarrow p \pi^0, \quad \pi^0 \rightarrow \gamma \gamma.$$

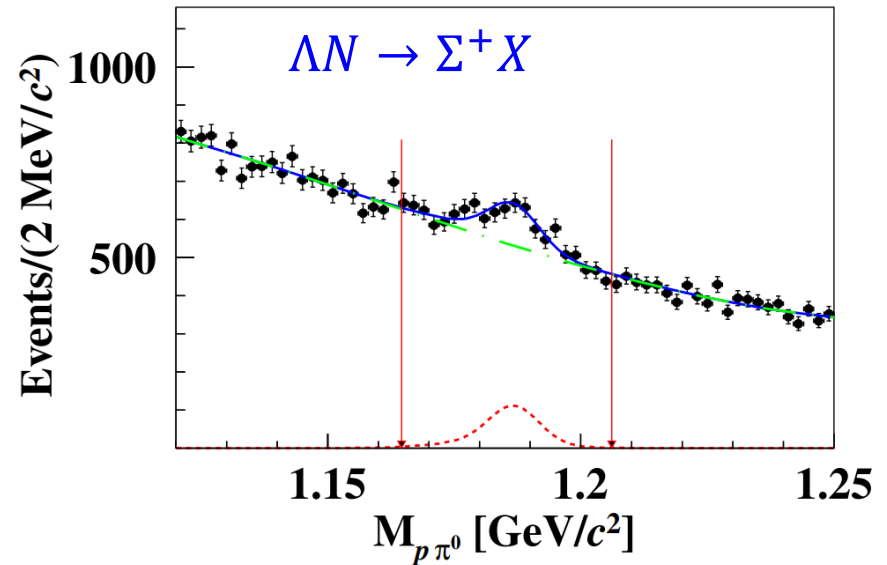


Two-body decay,
 $P_\Lambda \approx 1.074 \text{ GeV}/c$,
a very small horizontal
crossing angle of 11 mrad
for e^+ and e^- beams,
resulting in a small range of
0.017 GeV/c above and
below 1.074 GeV/c for P_Λ .

Study of $\Lambda N \rightarrow \Sigma^+ X$

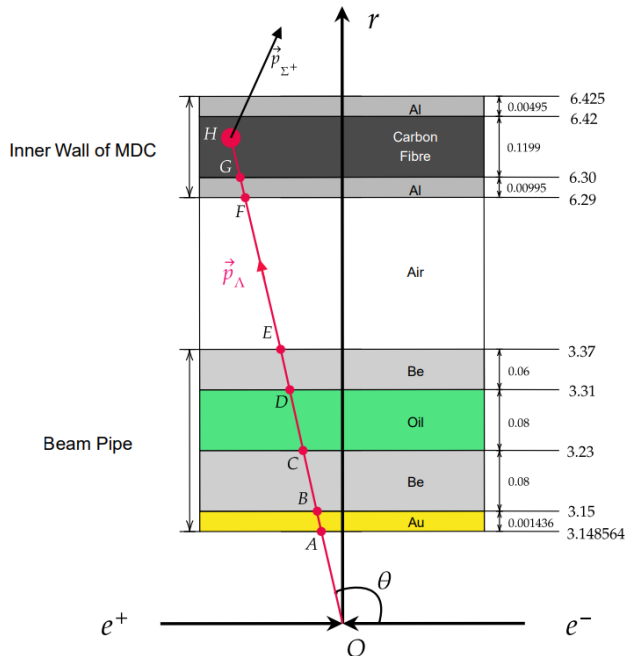


$$N_{ST} = 7207565 \pm 3741$$



$$N_{DT} = 795 \pm 101$$

The reaction position can not be determined. These signal events mainly come from the reaction with beam pipe and inner wall of MDC.



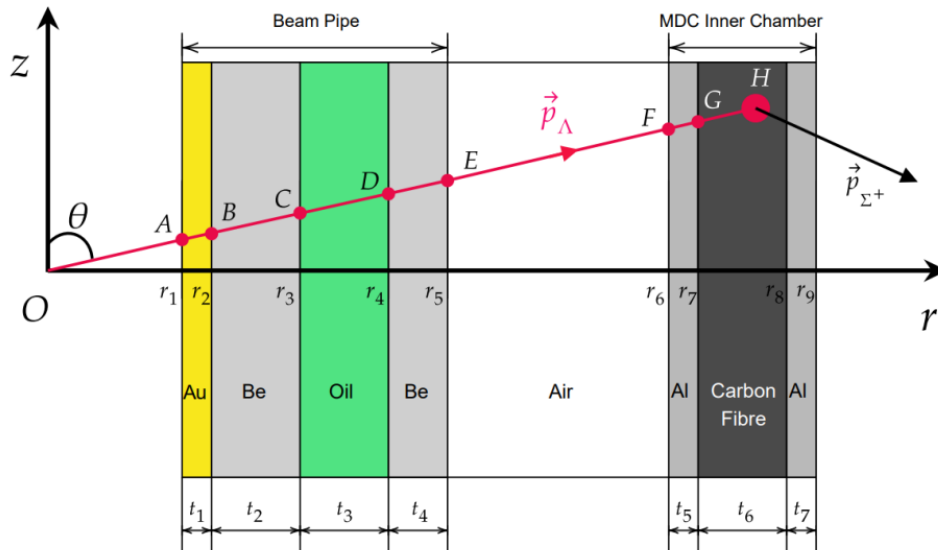
Cross section of $\Lambda + {}^9\text{Be} \rightarrow \Sigma^+ + X$

$$\sigma(\Lambda + {}^9\text{Be} \rightarrow \Sigma^+ + X) = \frac{N_{\text{DT}}}{\epsilon_{\text{sig}} \cdot \mathcal{L}_{\Lambda}} \cdot \frac{1}{\mathcal{B}(\Sigma^+ \rightarrow p\pi^0)}$$

$$\mathcal{L}_{\Lambda} = N_{\text{ST}} \cdot \frac{N_A}{N_{\text{ST}}^{\text{MC}}} \cdot \sum_j^7 \sum_i^{N_{\text{ST}}^{\text{MC}}} \frac{\rho_T^j \cdot l^{ij}}{M^j} \cdot \mathcal{R}_{\sigma}^j$$

path length of incident Λ of i^{th} event inside j^{th} layer

pure surface process assumption (proportional to number of protons)



Parameter	Value
N_{DT}	795 ± 101
ϵ_{sig}	24.32%
\mathcal{L}_{Λ}	$(17.00 \pm 0.01) \times 10^{28} \text{ cm}^{-2}$
$\mathcal{B}(\Sigma^+ \rightarrow p\pi^0)$	$(51.57 \pm 0.30)\%$

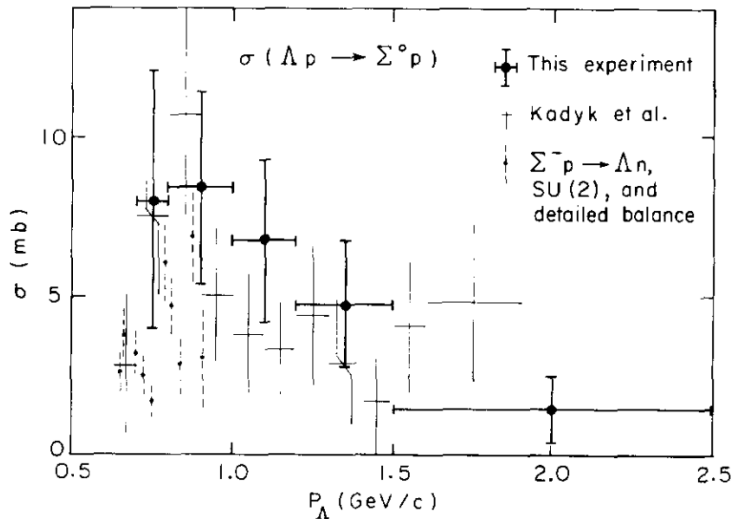
Study of $\Lambda N \rightarrow \Sigma^+ X$

The measured cross section of the reaction process $\Lambda + {}^9\text{Be} \rightarrow \Sigma^+ + X$ is $\sigma(\Lambda + {}^9\text{Be} \rightarrow \Sigma^+ + X) = (37.3 \pm 4.7_{\text{stat}} \pm 3.5_{\text{sys}})$ mb at $P_\Lambda \approx 1.074$ GeV/c. This work represents the first attempt to investigate Λ -nucleus interaction at an e^+e^- collider.

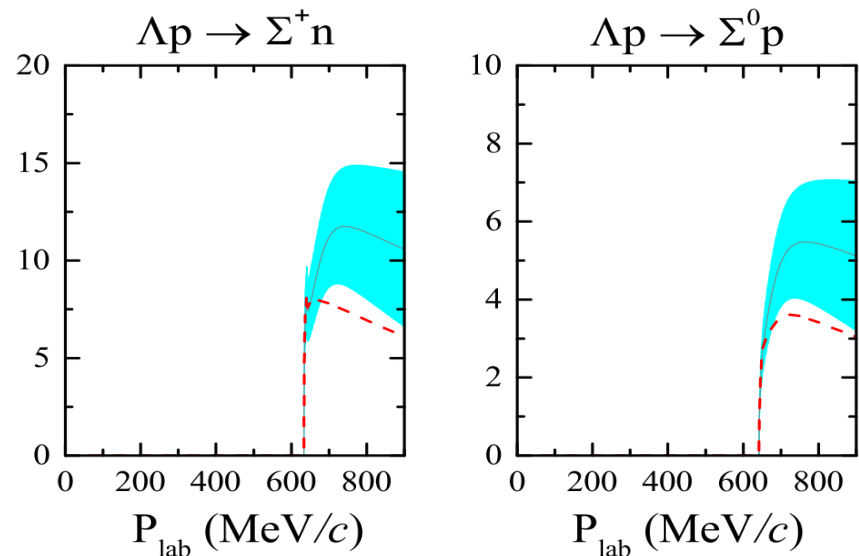
If taking the effective number of reaction protons in ${}^9\text{Be}$ nucleus as 1.93, the cross section of $\Lambda p \rightarrow \Sigma^+ X$ for single proton is determined to be $\sigma(\Lambda p \rightarrow \Sigma^+ X) = (19.3 \pm 2.4_{\text{stat}} \pm 1.8_{\text{sys}})$ mb.

$\sigma(\Lambda p \rightarrow \Sigma^+ n)$ is twice of $\sigma(\Lambda p \rightarrow \Sigma^0 p)$

Nucl. Phys. B 125, 29 (1977)

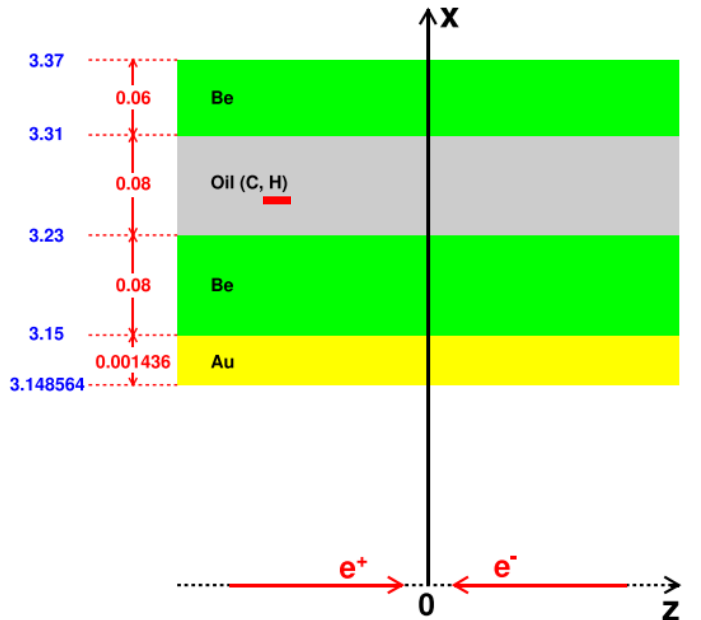


Phys. Rev. C 105, 035203 (2022)



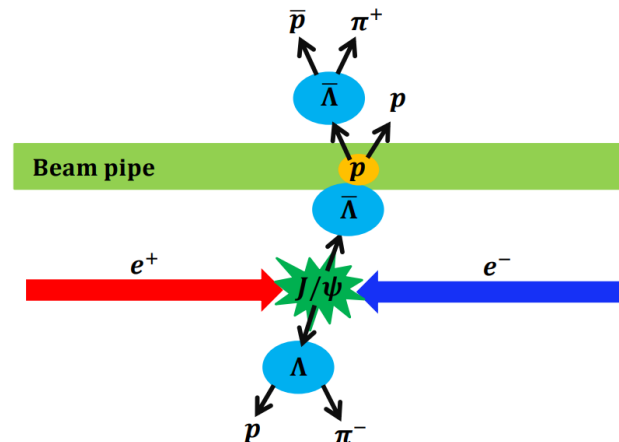
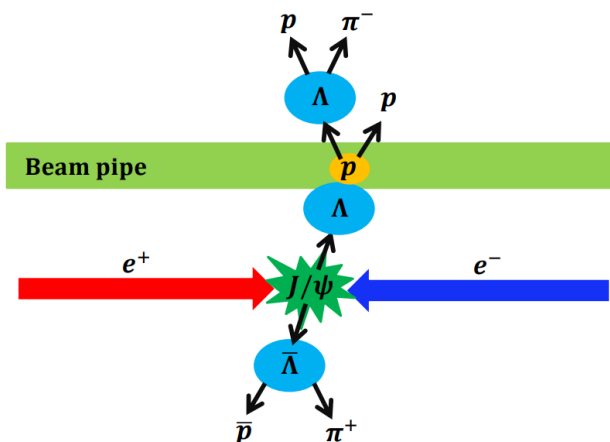
Study of $\Lambda p \rightarrow \Lambda p$ and $\bar{\Lambda} p \rightarrow \bar{\Lambda} p$

arXiv:2401.09012



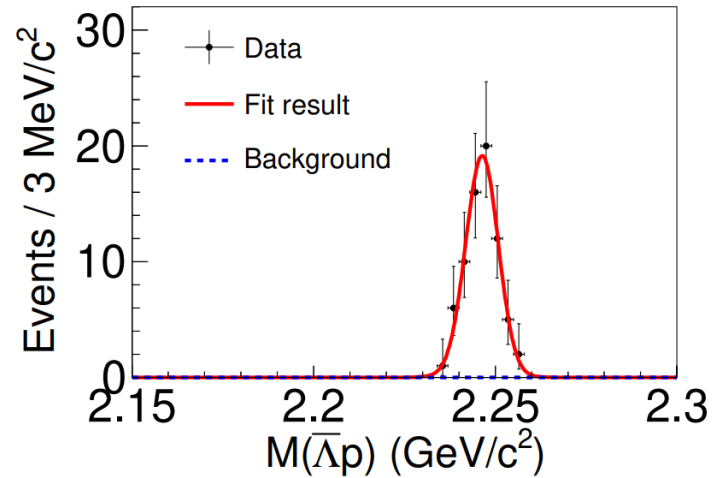
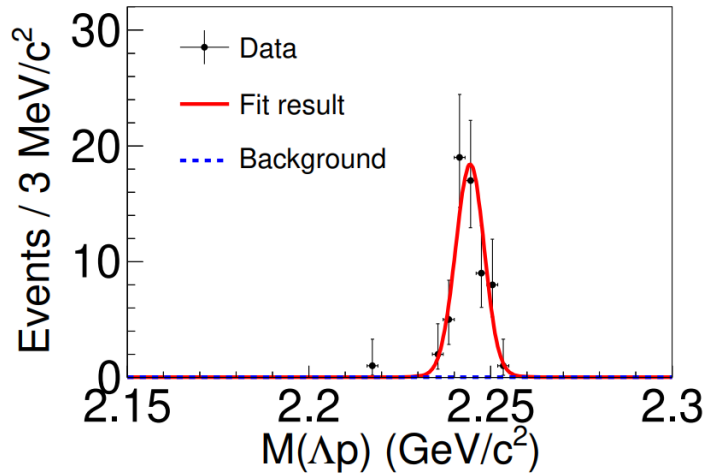
Taking the hydrogen in the cooling oil of the beam pipe as target material, the information on the hyperon-proton scattering can be directly extracted.

Two-body decay, $P_{\Lambda/\bar{\Lambda}} \approx 1.074 \text{ GeV}/c$



Study of $\Lambda p \rightarrow \Lambda p$ and $\bar{\Lambda} p \rightarrow \bar{\Lambda} p$

The center-of-mass energy for the incident $\Lambda/\bar{\Lambda}$ and a static p is about $2.243 \text{ GeV}/c^2$.

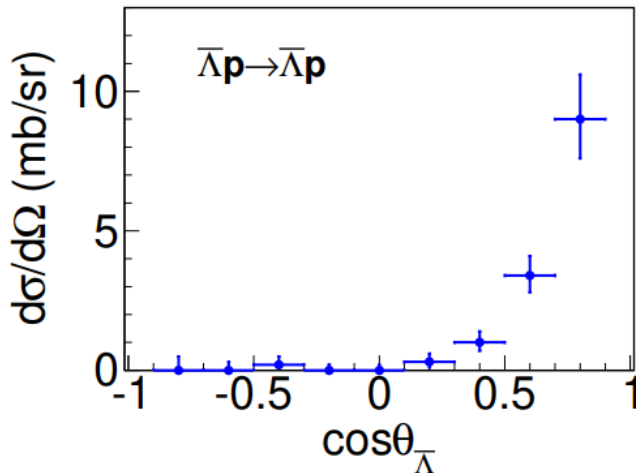
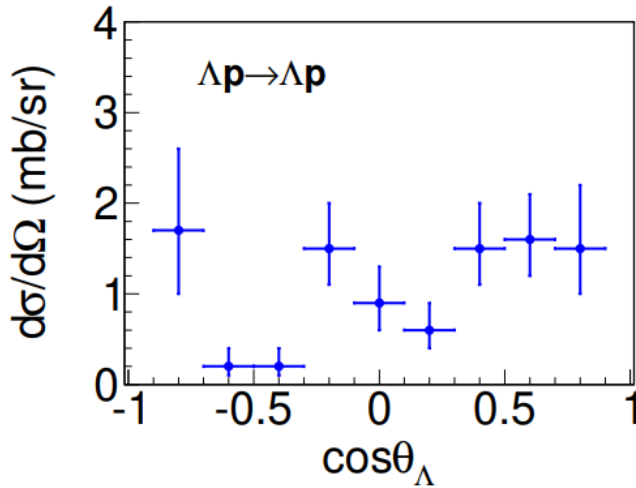


$$\sigma(\Lambda p \rightarrow \Lambda p / \bar{\Lambda} p \rightarrow \bar{\Lambda} p) = \frac{N_{\Lambda p / \bar{\Lambda} p}^{\text{sig}}}{\epsilon_{\Lambda p / \bar{\Lambda} p} \mathcal{B} \mathcal{L}_{\text{eff}}}$$

$$\mathcal{L}_{\text{eff}} = \frac{N_{J/\psi} \mathcal{B}_{J/\psi}}{2 + \frac{2}{3} \alpha} \int_a^b \int_0^\pi (1 + \alpha \cos^2 \theta) e^{-\frac{x}{\sin \theta \beta \gamma L}} N_H d\theta dx$$

$$\left(\frac{d\sigma}{d\Omega} \right)_i = \frac{N_i^{\text{sig}}}{\epsilon_i \mathcal{B} \mathcal{L}_{\text{eff}} \Delta\Omega}$$

Study of $\Lambda p \rightarrow \Lambda p$ and $\bar{\Lambda} p \rightarrow \bar{\Lambda} p$



$\cos\theta_{\Lambda/\bar{\Lambda}}$	N_i^{sig}	ϵ_i (%)	$\frac{d\sigma}{d\Omega}$ (mb/sr)
$[-0.9, -0.7]$	$(5.0^{+2.6}_{-1.9}, 0.0^{+1.1}_{-0.0})$	(6.94, 4.93)	$(1.7^{+0.9}_{-0.7}, 0.0^{+0.5}_{-0.0})$
$(-0.7, -0.5]$	$(1.0^{+1.4}_{-0.7}, 0.0^{+1.1}_{-0.0})$	(14.13, 10.44)	$(0.2^{+0.2}_{-0.1}, 0.0^{+0.3}_{-0.0})$
$(-0.5, -0.3]$	$(1.0^{+1.4}_{-0.7}, 1.0^{+1.4}_{-0.7})$	(17.32, 13.27)	$(0.2^{+0.2}_{-0.1}, 0.2^{+0.3}_{-0.1})$
$(-0.3, -0.1]$	$(11.0^{+3.7}_{-3.0}, 0.0^{+1.1}_{-0.0})$	(17.74, 14.66)	$(1.5^{+0.5}_{-0.4}, 0.0^{+0.2}_{-0.0})$
$(-0.1, 0.1]$	$(6.9^{+3.0}_{-2.3}, 0.0^{+1.1}_{-0.0})$	(19.11, 15.79)	$(0.9^{+0.4}_{-0.3}, 0.0^{+0.2}_{-0.0})$
$(0.1, 0.3]$	$(5.0^{+2.6}_{-1.9}, 2.0^{+1.8}_{-1.1})$	(19.53, 16.82)	$(0.6^{+0.3}_{-0.2}, 0.3^{+0.3}_{-0.2})$
$(0.3, 0.5]$	$(12.0^{+3.8}_{-3.1}, 7.0^{+3.0}_{-2.3})$	(19.21, 17.68)	$(1.5^{+0.5}_{-0.4}, 1.0^{+0.4}_{-0.3})$
$(0.5, 0.7]$	$(13.0^{+3.9}_{-3.3}, 25.0^{+5.3}_{-4.7})$	(19.71, 17.60)	$(1.6^{+0.5}_{-0.4}, 3.4^{+0.7}_{-0.6})$
$(0.7, 0.9]$	$(6.0^{+2.8}_{-2.1}, 37.0^{+6.4}_{-5.8})$	(9.80, 9.93)	$(1.5^{+0.7}_{-0.5}, 9.0^{+1.6}_{-1.4})$

Cross sections in $-0.9 \leq \cos\theta_{\Lambda/\bar{\Lambda}} \leq 0.9$

are measured to be

$$\sigma(\Lambda p \rightarrow \Lambda p) = (12.2 \pm 1.6 \pm 1.1) \text{ mb and}$$

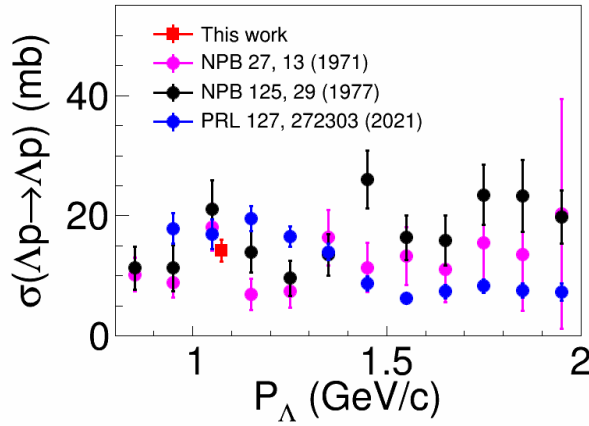
$$\sigma(\bar{\Lambda} p \rightarrow \bar{\Lambda} p) = (17.5 \pm 2.1 \pm 1.6) \text{ mb}$$

Total cross sections are determined to be

$$\sigma_t(\Lambda p \rightarrow \Lambda p) = (14.2 \pm 1.8 \pm 1.3) \text{ mb and}$$

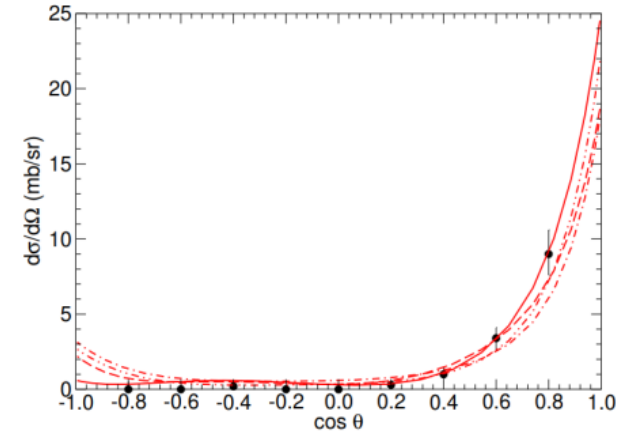
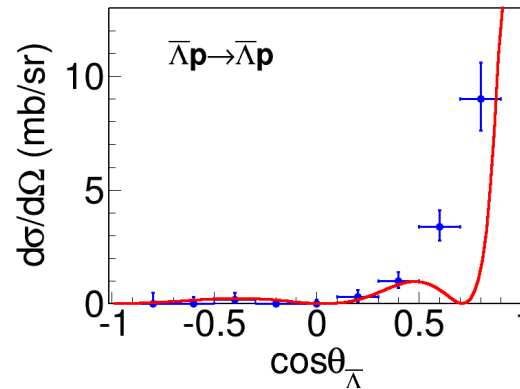
$$\sigma_t(\bar{\Lambda} p \rightarrow \bar{\Lambda} p) = (27.4 \pm 3.2 \pm 2.5) \text{ mb}$$

Study of $\Lambda p \rightarrow \Lambda p$ and $\bar{\Lambda} p \rightarrow \bar{\Lambda} p$



strong absorption/annihilation

J. Haidenbauer et al,
arXiv:2403.18706



Phys. Rev. 112, 1303 (1958)

“black sphere” scattering

$$\frac{d\sigma}{d\Omega} = k^2 R^4 \left[\frac{J_1(2kR \sin(\theta/2))}{2kR \sin(\theta/2)} \right]^2,$$

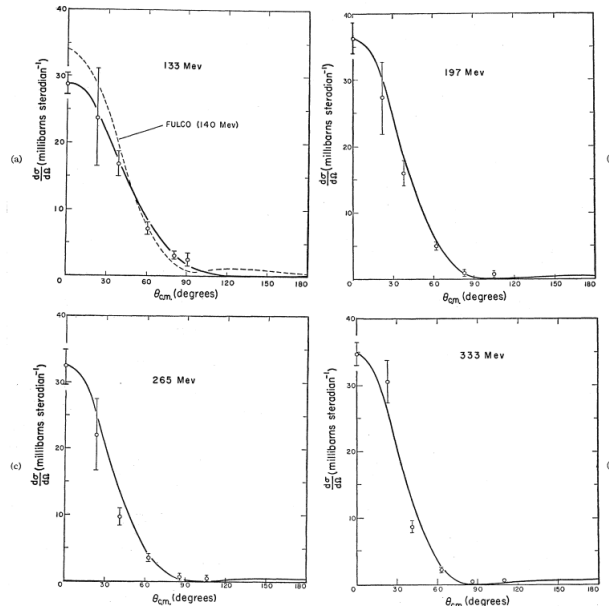


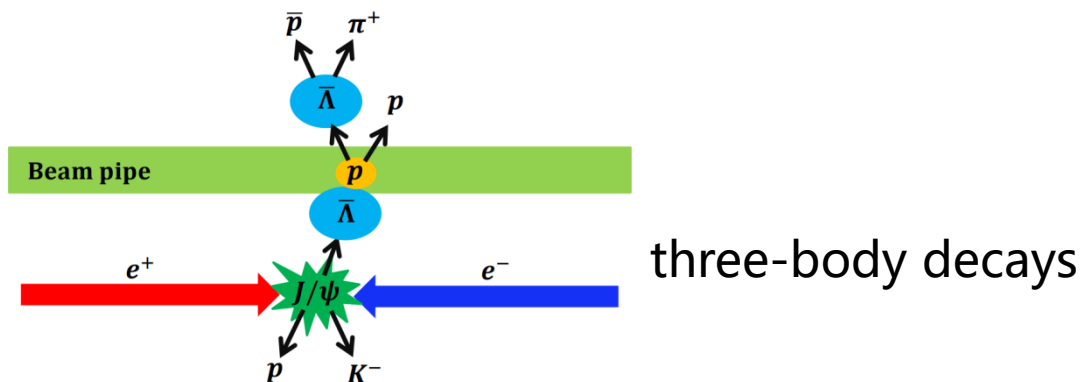
Fig. 1 Differential cross section for $p\bar{\Lambda}$ scattering at $p_{lab} = 1.074 \pm 0.017$ GeV [59]. The curves are predictions by the $\Lambda\bar{\Lambda}$ interactions I-IV, see Ref. [35], at 1.05 GeV/c.

Table 2 $\Lambda\bar{\Lambda}$ scattering lengths (in fm) in the 1S_0 and 3S_1 partial waves of the employed $\Lambda\bar{\Lambda}$ potentials [41,43]. The spin-averaged value by the ALICE Collaboration is from an analysis of the $\Lambda\bar{\Lambda}$ correlation function measured in Pb-Pb collisions [21]

potential	$a(^1S_0)$	$a(^3S_1)$
I	0.32 – i0.52	0.74 – i0.56
II	0.67 – i1.14	0.66 – i0.37
III	1.42 – i1.15	1.00 – i0.44
IV	1.56 – i1.40	0.98 – i0.65
ALICE	$(0.90 \pm 0.16) - i(0.40 \pm 0.18)$	

More (anti)hyperon-nucleon scattering can be studied at BESIII

Elastic scattering	Inelastic scattering
$\Lambda p \rightarrow \Lambda p$	$\Lambda p \rightarrow \Sigma^0 p, \Lambda p \pi^0, \Sigma^+ p \pi^-$
$\bar{\Lambda} p \rightarrow \bar{\Lambda} p$	$\bar{\Lambda} p \rightarrow \bar{\Sigma}^0 p, \bar{\Lambda} p \pi^0, \bar{\Sigma}^- p \pi^+$
$\Sigma^+ p \rightarrow \Sigma^+ p$	$\Sigma^+ p \rightarrow \Lambda p \pi^+, \Sigma^+ p \pi^0, \Sigma^0 p \pi^+$
$\bar{\Sigma}^- p \rightarrow \bar{\Sigma}^- p$	$\bar{\Sigma}^- p \rightarrow \bar{\Lambda} p \pi^-, \bar{\Sigma}^- p \pi^0, \bar{\Sigma}^0 p \pi^-$
$\Xi^0 p \rightarrow \Xi^0 p$	$\Xi^0 p \rightarrow \Lambda \Sigma^+, \Sigma^0 \Sigma^+, \Xi^0 p \pi^0, \Xi^- p \pi^+$
$\bar{\Xi}^0 p \rightarrow \bar{\Xi}^0 p$	$\bar{\Xi}^0 p \rightarrow \bar{\Xi}^0 p \pi^0, \bar{\Xi}^+ p \pi^-$
$\Xi^- p \rightarrow \Xi^- p$	$\Xi^- p \rightarrow \Lambda \Lambda, \Lambda \Sigma^0, \Sigma^0 \Sigma^0, \Xi^- p \pi^0, \Xi^0 p \pi^-$
$\bar{\Xi}^+ p \rightarrow \bar{\Xi}^+ p$	$\bar{\Xi}^+ p \rightarrow \bar{\Xi}^+ p \pi^0, \bar{\Xi}^0 p \pi^+$
$\Omega^- p \rightarrow \Omega^- p$	$\Omega^- p \rightarrow \Lambda \Xi^0, \Sigma^0 \Xi^0, \Sigma^+ \Xi^-$
$\bar{\Omega}^+ p \rightarrow \bar{\Omega}^+ p$...





Summary

1. Recently, some hyperon-nucleon scattering processes have been measured with unprecedented precision by CLAS and J-PARC E40 Collaborations. Hyperon-nucleon scattering can also be measured at BESIII now.

- $\Xi^0 n \rightarrow \Xi^- p$
- $\Lambda N \rightarrow \Sigma^+ X$
- $\Lambda p \rightarrow \Lambda p$
- $\bar{\Lambda} p \rightarrow \bar{\Lambda} p$

2. This is the first study of hyperon-nucleon scattering in electron-positron collisions, and opens up a new direction for such research. Especially, antihyperon-nucleon scattering is studied for the first time.

3. More (anti)hyperon-nucleon scattering processes can be studied at BESIII.

Thanks for your attention!