

超子的自旋结构

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Our View of the Proton

with history

- **Point-Like** 1919
- **Finite Size with Radius** 1930s-1950s
- **Quark Model** 1960s
- **QCD and Gluons** 1970s
- **Puzzles and Anomalies** 1980s-present
 - **Quark Sea of the Nucleon**
 - **Baryon-Meson Fluctuations**
 - **Statistical Features**
 -

Surprises & Unknown

about the Quark Structure of Nucleon: **Sea**

- **Spin Structure:** $\Sigma = \Delta u + \Delta d + \Delta s \approx 0.3$

“puzzle”: where is the proton’s missing spin

- **Strange Content** $\Delta s \neq 0$ $s(x) \neq \bar{s}(x)$

Brodsky & Ma, PLB381(96)317

- **Flavor Asymmetry** $\bar{u} \neq \bar{d}$

- **Isospin Symmetry Breaking** $\bar{u}_p \neq \bar{d}_n$ $\bar{d}_p \neq \bar{u}_n$

Ma, PLB 274 (92) 111

Boros, Londergan, Thomas, PRL81(98)4075

Unknown about the nucleon: **valence**

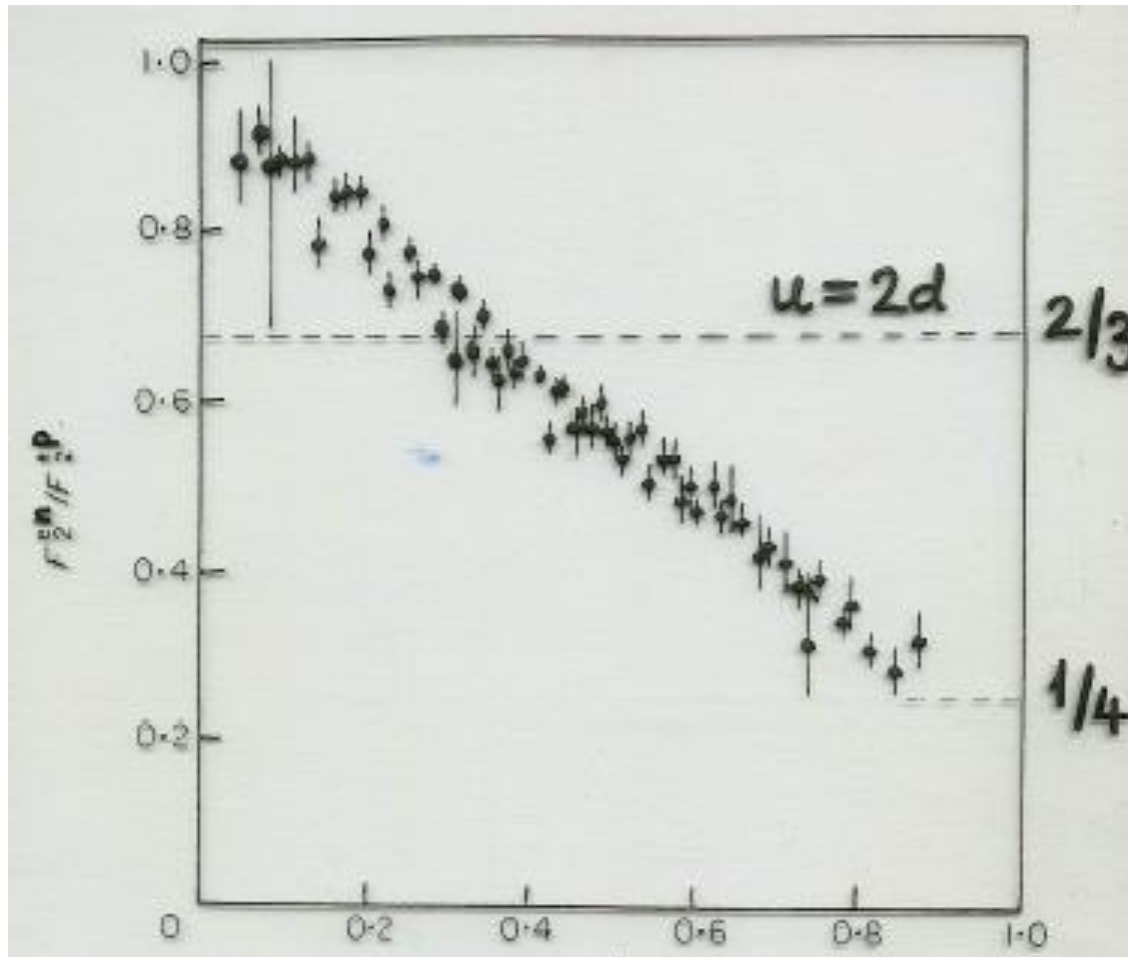
$x \rightarrow 1$ behaviors of flavor and spin

- Flavor**

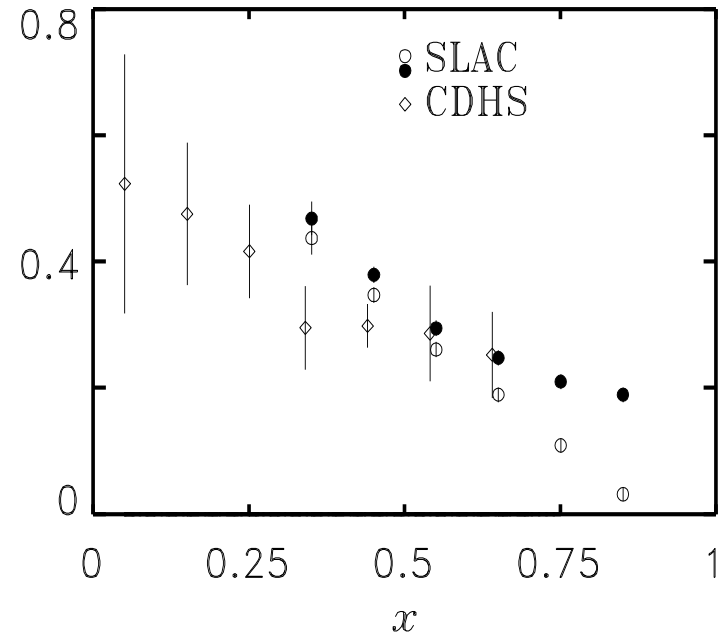
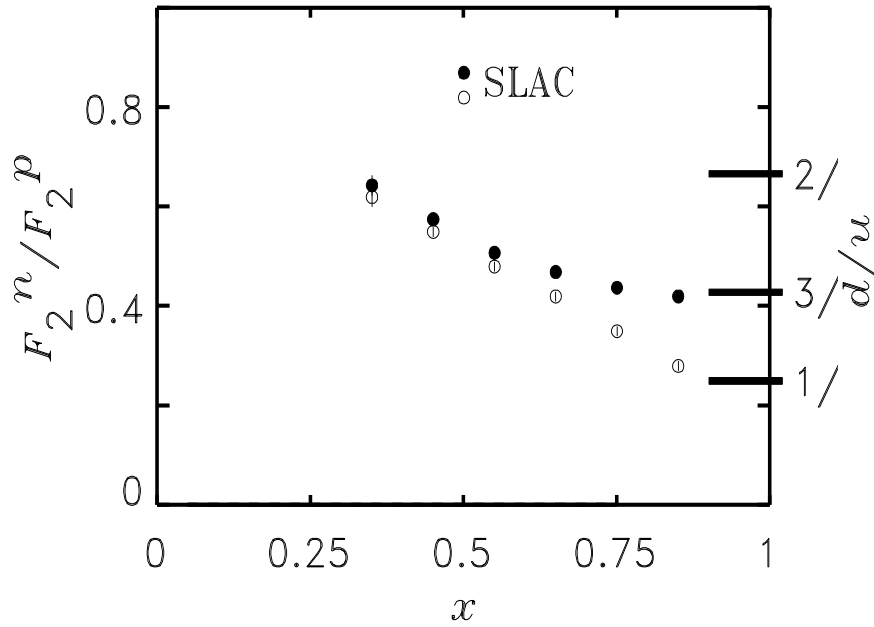
$\frac{d(x)}{u(x)}$	\nearrow	0	Diquark Model
$\frac{F_2^n(x)}{F_2^p(x)}$	\searrow	$\frac{1}{2}$	pQCD
$\frac{F_2^n(x)}{F_2^p(x)}$	\nearrow	$\frac{1}{2}$	Diquark Model
$\frac{F_2^n(x)}{F_2^p(x)}$	\searrow	$\frac{1}{2}$	pQCD
$\frac{\Delta d(x)}{d(x)}$	\rightarrow	$-\frac{1}{3}$	Diquark Model
	\searrow	1	pQCD

- Spin**

Ratio of Neutron/Proton Structure Functions

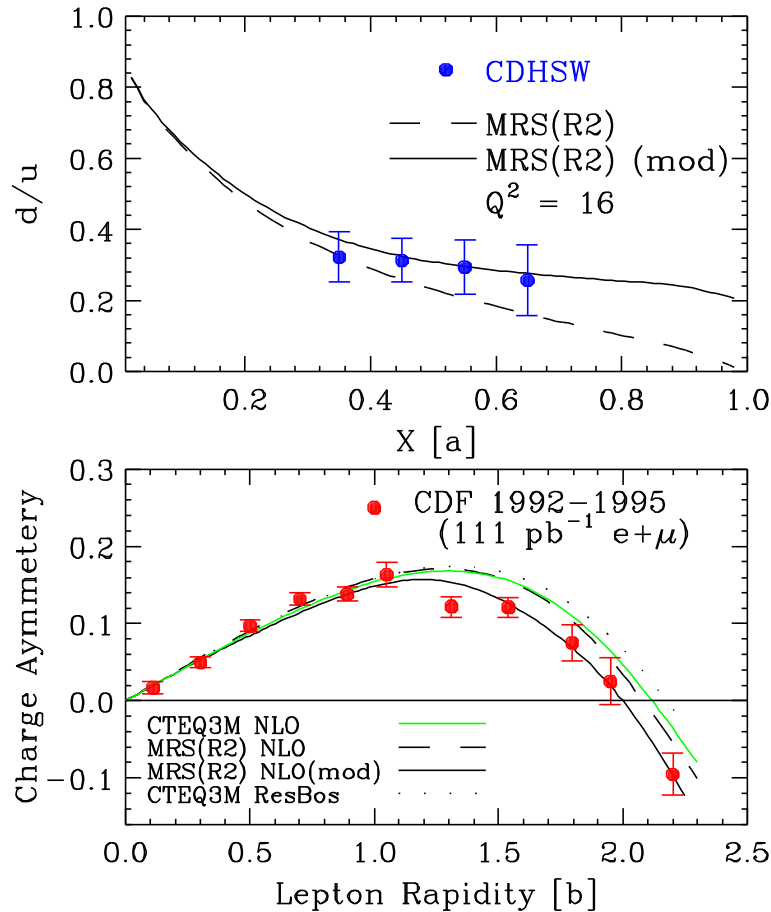


Flavor Content of the Proton with nuclear binding correction



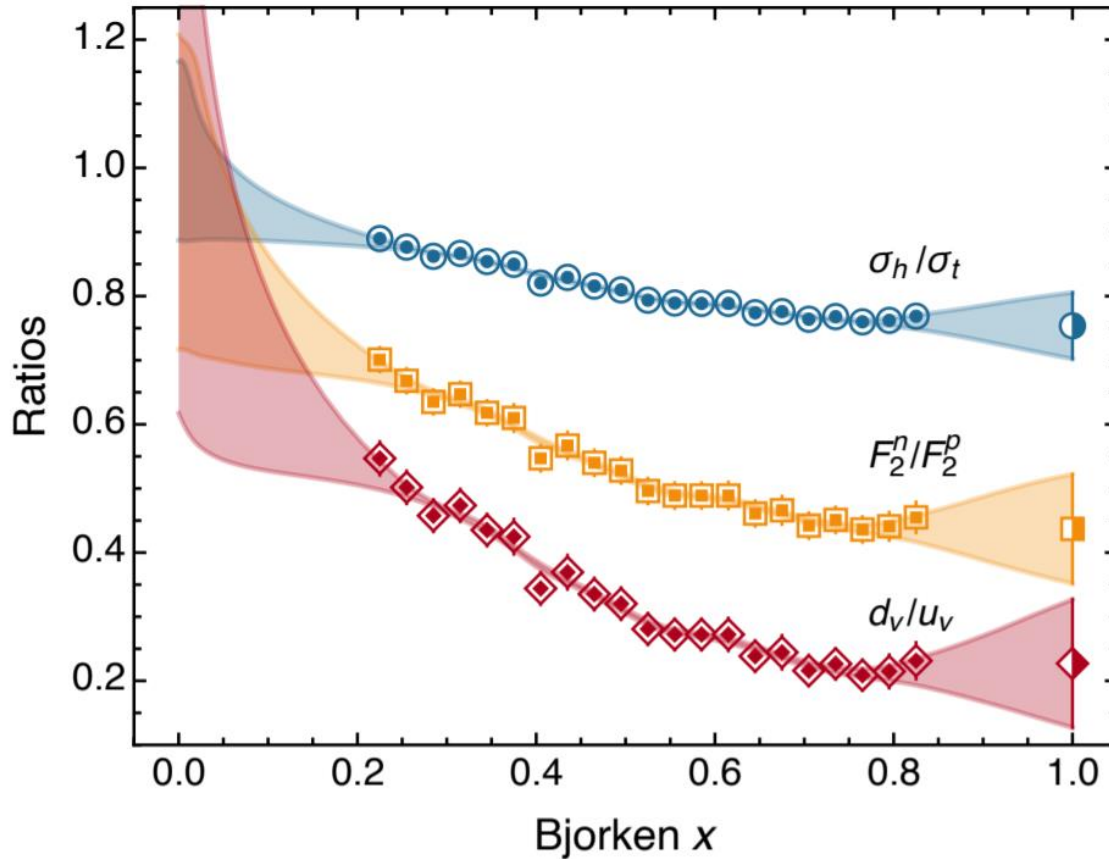
W.Melnitchouk & A.W. Thomas
PLB 377(1996) 11

Flavor Content of the Proton from DIS neutrino data analysis



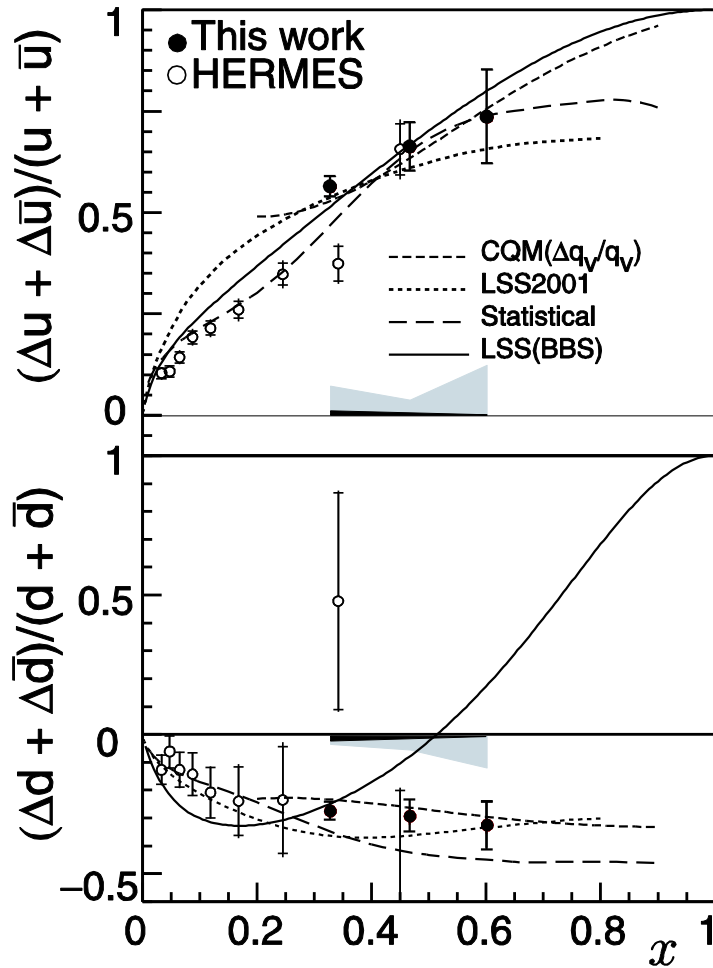
U.K. Yang & A. Bodek
PRL 82 (1999) 2467.

Flavor Content of the Proton from DIS data of ^3He and ^3H



Quark Helicity Distributions of Proton

Measurements at JLAB and HERMES



X. Zheng et al, JLAB Hall A Collaboration
nucl-ex/0308011
PRL92 (2004) 012004.

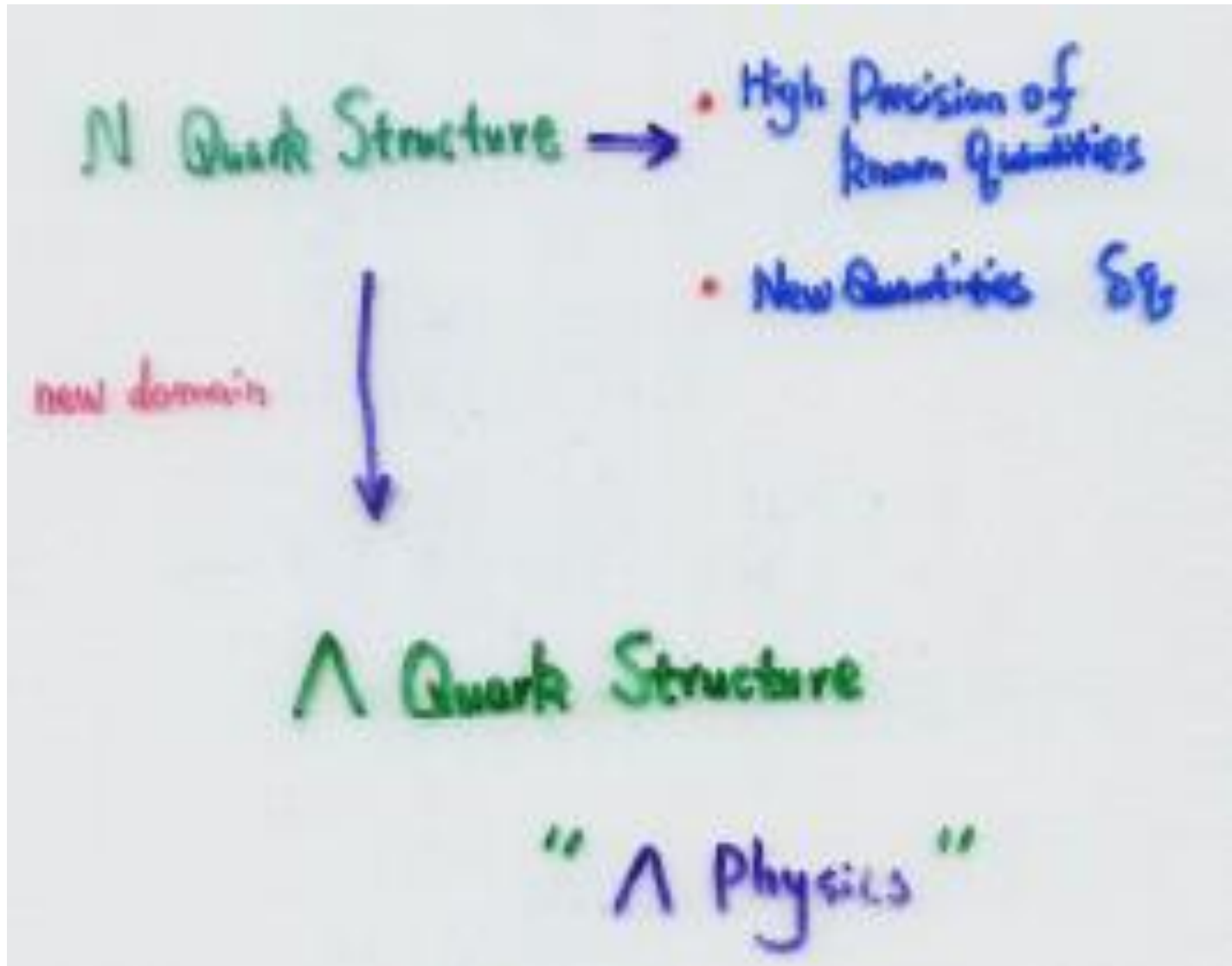
Status

of the Flavor and Spin Contents of the Proton

- Flavor favors pQCD
- Spin favors diquark model

Unclear & In Contradiction!

How to Test Various Theories?



SU(3) Symmetry together with Proton Spin Problem

PRL 70(9) 2537

Burkhardt-Luffe SU(3) Argument:

$$\int_0^1 dx g_1^{uA}(\omega) = \frac{1}{18} (2\Sigma - D)$$

$$= \int_0^1 dx g_1^{ep}(\omega) - \frac{1}{18} (2D + 3F)$$

$$= -0.042 \pm 0.019$$

$$\Delta u^A = \Delta d^A = \frac{1}{3} (\Sigma - D) = -0.23 \pm 0.06$$

$$\Delta S^A = \frac{1}{3} (\Sigma + 2D) = 0.58 \pm 0.07$$

whereas the Quark Model predicts

$$\Delta u^A = \Delta d^A = 0$$

$$\Delta S^A = 1$$

The u, d sea of Lambda versus s sea of Nucleon

Ma-Soffer

PRL 82(99) 2250

u, d polarizations in Λ

is related to s polarization in N

- $P(u, d, s) = \Lambda(u, d, s) K^+(u, s)$

- $\begin{cases} \Lambda(u, d, s, u) = P(u, d) K^-(s, u) \\ \Lambda(u, d, s, d) = P(u, d) K^-(s, d) \end{cases}$

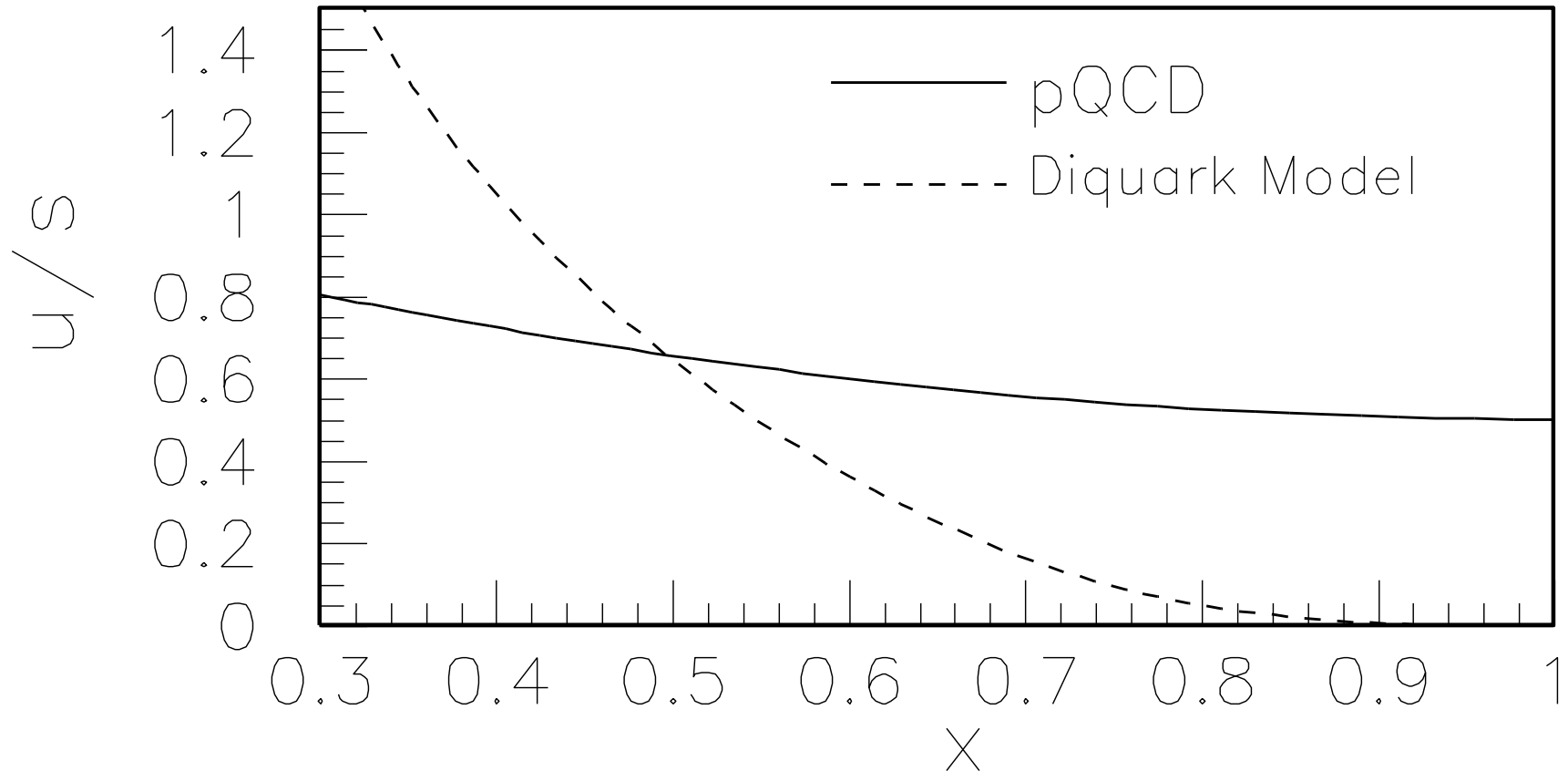
Different flavor & spin structure in different models

Phys. Lett. B 477 (2000) 107
Ma-Schmidt-Yang $x \rightarrow 1$ behaviors

- Flavor
 - $\frac{u(x)}{S(x)} \begin{cases} \rightarrow 0 & \text{Diquark Model} \\ \rightarrow \frac{1}{2} & \text{pQCD} \end{cases}$
 - $\frac{\Delta S(x)}{S(x)} \rightarrow 1$
 - $\frac{\Delta u(x)}{u(x)} \rightarrow 1$

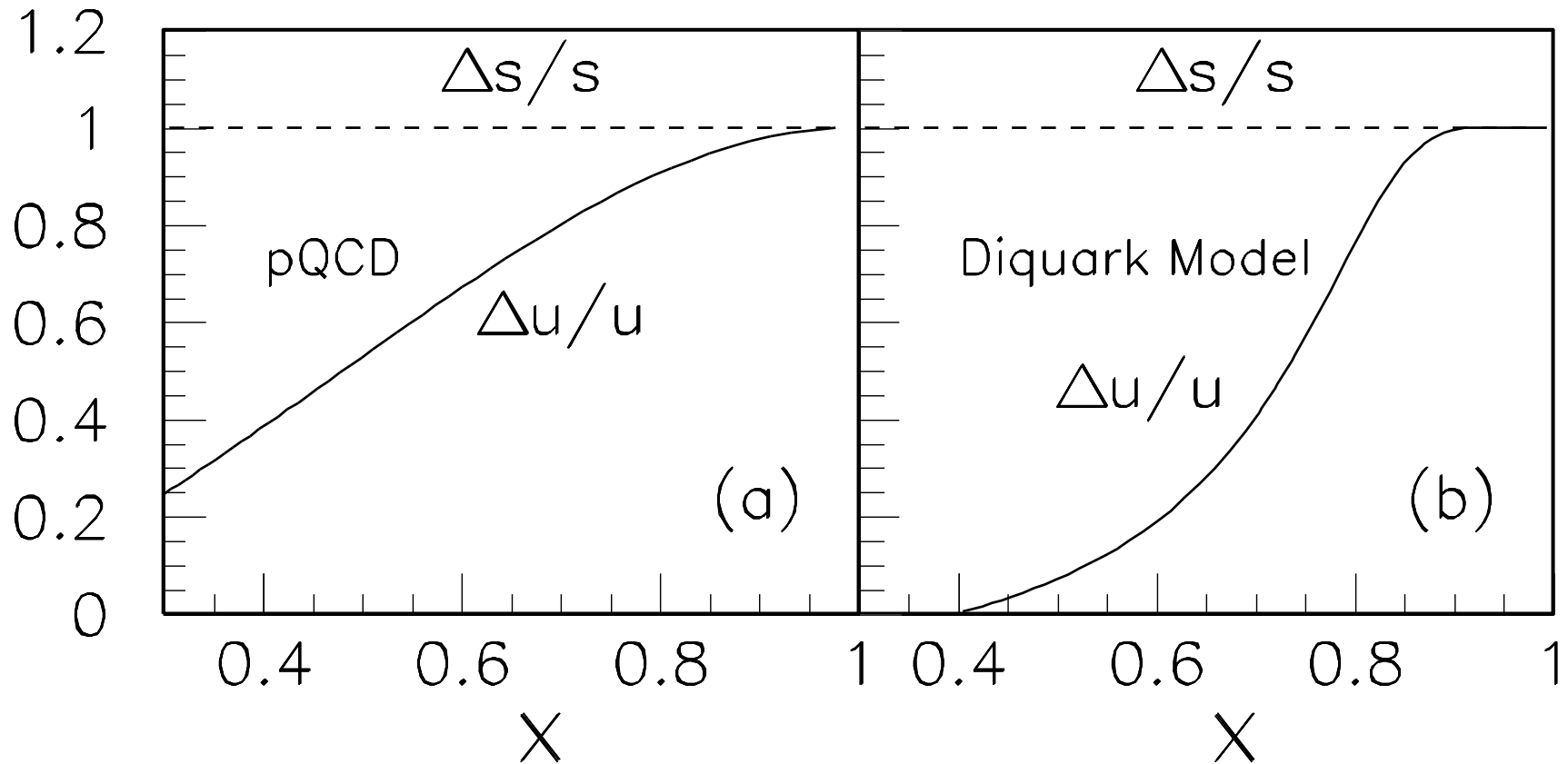
$\Delta u(x) = \Delta d(x) > 0$ at large x
 $\int_0^1 dx \Delta u(x) = \int_0^1 dx \Delta d(x) \leq 0$

Flavor structure in two different models



**B.-Q. Ma, I. Schmidt, J.-J. Yang,
Phys. Lett. B 477 (2000) 107**

Spin structure in two different models



B.-Q. Ma, I. Schmidt, J.-J. Yang,
Phys. Lett. B 477 (2000) 107

An intuitive argument

Quark-Diquark Model

$S(ud)S$

$D(ds)u$

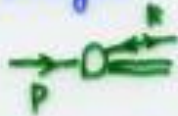
$u^*S(ds)$ $M_D = M_S$

$u^*V(ds)$ $M_D = M_U > M_S$

$$\psi(x) \approx e^{-x \left(\frac{\bar{k}_1^2 + M_U^2}{x} + \frac{\bar{k}_2^2 + M_D^2}{1-x} \right)}$$

at $x \rightarrow 1$ $\psi_U(x) \ll \psi_S(x)$

pQCD Analysis



$$q_h(x) \sim (1-x)^p \quad p = 2m-1 + 2|p_S|$$

$$dS_d = S_q - S_p$$

$$S_q = S_p \quad |dS_d| = 0$$

$$S_q \neq S_p \quad |dS_d| = 1 \quad \text{suppressed}$$

Significantly different predictions of Lambda Structure

- naive quark model predicts:
 $\Delta u = \Delta d = 0, \Delta S = 1$
- Jaffe-Burkhardt predict:
 $\Delta u = \Delta d = -0.2, \Delta S = 0.6$
- We predict:
 $\frac{\Delta u}{u} = \frac{\Delta d}{d} \rightarrow 1$ at $x \rightarrow 1$
in both quark-diquark model
and p QCD analysis

Connections between structure functions and fragmentation functions

How to Measure $f^{\wedge}(\omega)$, $\Delta f^{\wedge}(\omega)$?

$$f^{\wedge}(\omega) \propto D_q^{\wedge}(\tau)$$

$$x = \frac{\omega^2}{4q}$$
$$\tau = \frac{2q}{\omega^2}$$

- space-like time-like
- The Guber-Lipatov reciprocity relation
- parton distribution by parton fragmentation duality

•S.J. Brodsky, B.-Q. Ma, PLB 392 (1997) 452.

•V. Barone, A. Drago, B.-Q. Ma, PRC 62 (2000) 062201 (R).

•B.-Q. Ma, I. Schmidt, J. Soffer, J.-J. Yang, PLB 547 (2002) 245.

Various Processes of Polarized Fragmentation

Various Processes to Measure $D_f^{\uparrow}(\vec{e})$, $\Delta D_f^{\uparrow}(\vec{e})$

- $e^+e^- \rightarrow \vec{\Lambda} + X$

M. Ecker et al. R.L. Jaffe
PRD 70 (9) (2004)

- $\vec{e}N \rightarrow \vec{\Lambda} + X$

R.L. Jaffe, PRD 54 (8) (1996)

- $p\vec{p} \rightarrow \vec{\Lambda} + X$

D. de Florian, M. Sotomayor, A. Bac
W. Vogelsang, PRD 61 (9) (2000)

- $\nu N \rightarrow \vec{\Lambda} + X$

$$\frac{\Delta D_f^{\uparrow}(\vec{e})}{D_f^{\uparrow}(\vec{e})}$$

Kalweit-Bauer
- et al. Hamburg
EPJ C 27 (1999) 261

Flavor separation of fragmentation functions

B.-Q. Ma, J. Soffer, PRL 82 (1999) 2250

Complete flavor separation of

$$D_g^{\hat{A}}(z), \Delta D_g^{\hat{A}}(z), D_g^{\hat{B}}(z), \Delta D_g^{\hat{B}}(z)$$

- $\nu N \rightarrow \mu^- \vec{\Lambda} X$
- $\bar{\nu} N \rightarrow \mu^+ \vec{\Lambda} X$
- $\nu N \rightarrow \mu^- \vec{\bar{\Lambda}} X$
- $\bar{\nu} N \rightarrow \mu^+ \vec{\bar{\Lambda}} X$

Extension to Sigma Hyperon

Ma-Schmidt-Yang hep-ph/9907556
Nucl. Phys. B 574 (2000) 331

Σ^+ Valence: $\alpha \rightarrow 1$

$\frac{s}{u} \rightarrow 0$ Diquark Model
 $\frac{s}{u} \rightarrow \frac{1}{5}$ pQCD

$\frac{ds}{s} \rightarrow -\frac{1}{3}$ Diquark Model
 $\frac{ds}{s} \rightarrow 1$ pQCD

$\frac{du}{u} \rightarrow 1$

- Bigger difference at middle x

The advantage of using Sigma hyperons

The Advantage of Σ^\pm

charged, us beam



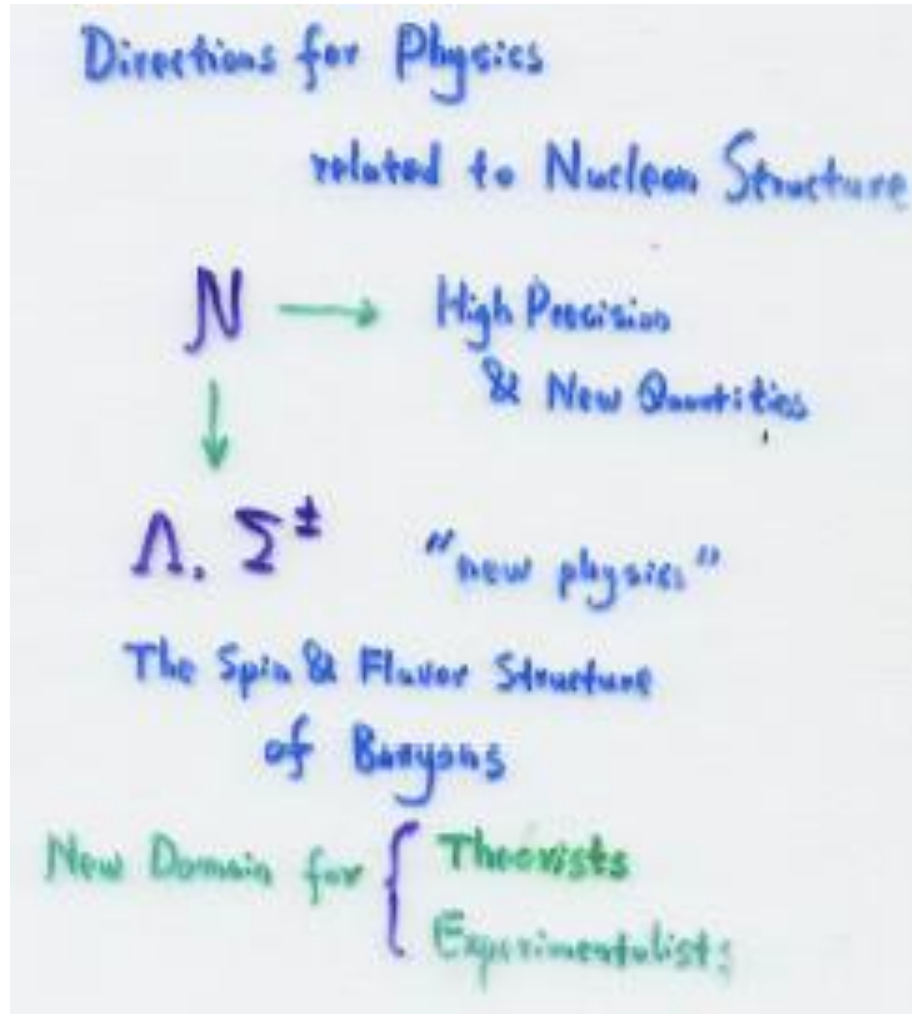
Albrey et al. PLB 379 (1967)

363

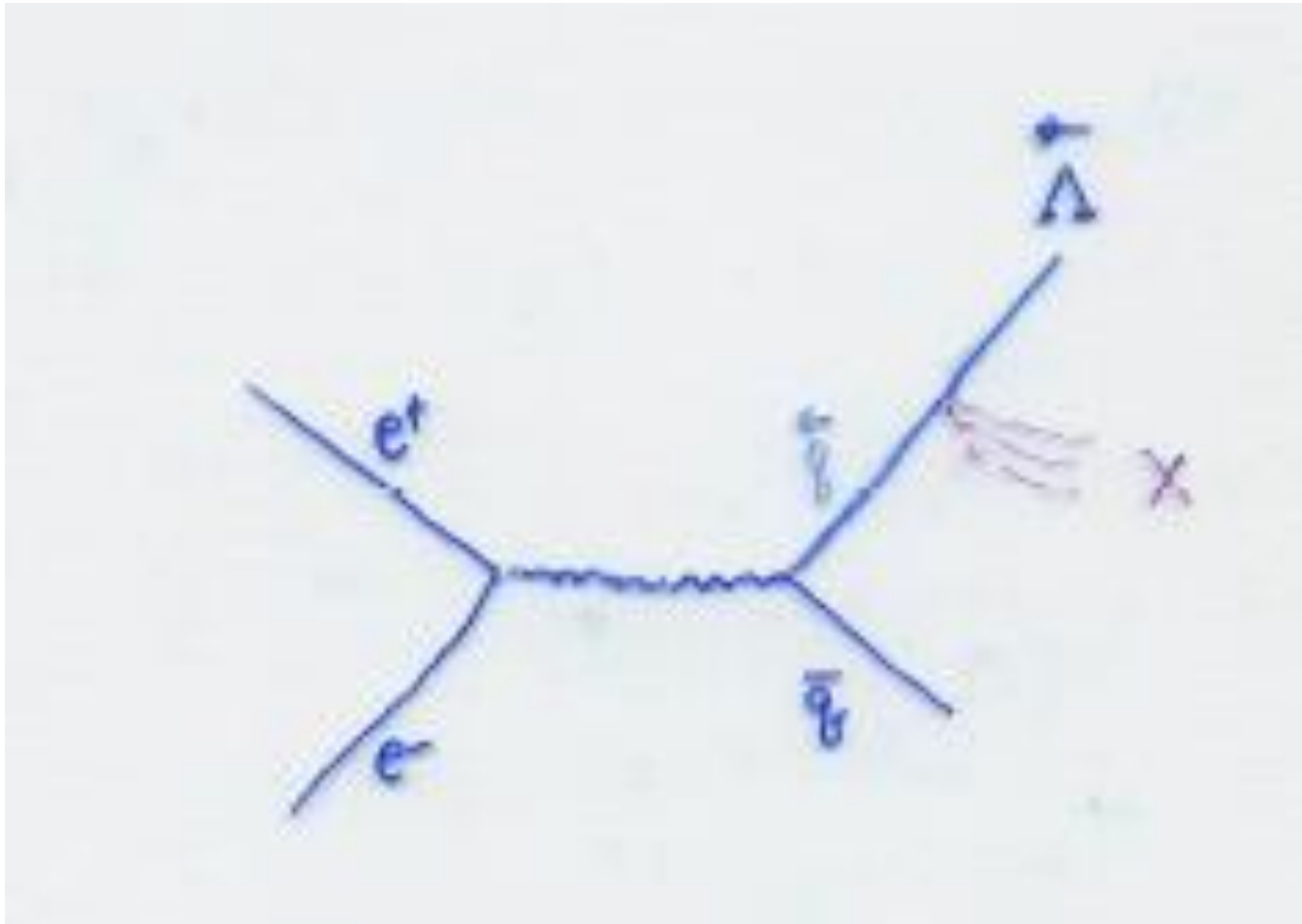
Ma, Schmidt, Yang, Topolnick

The quark distributions can be measured

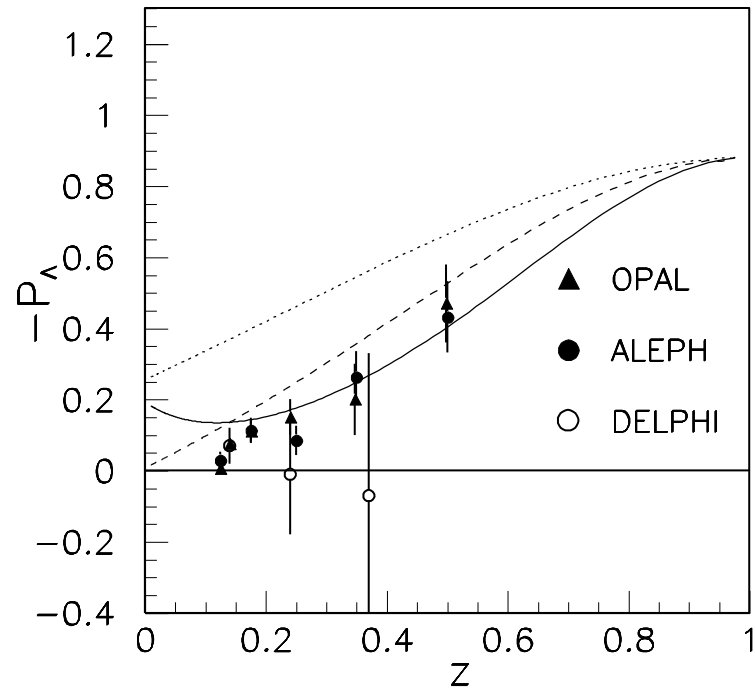
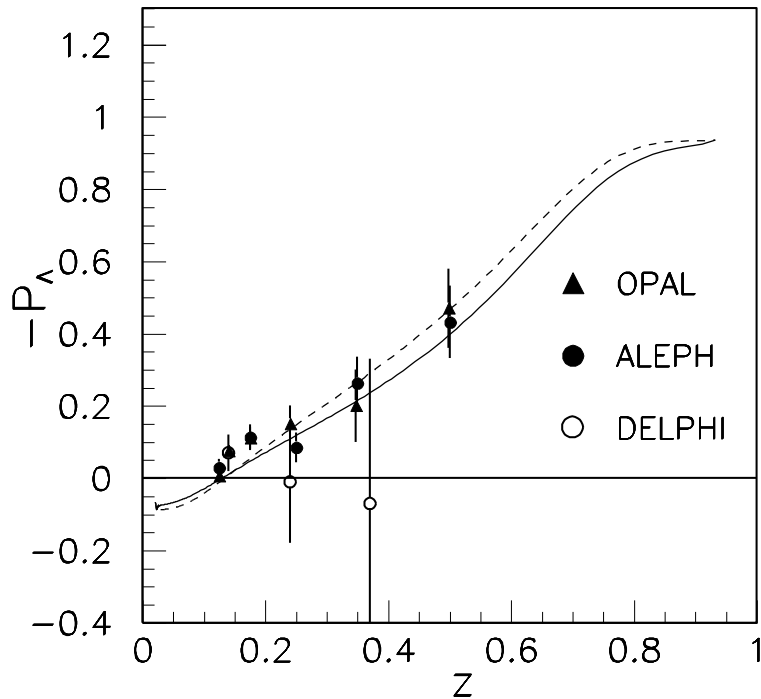
New Domain for Theorists and Experimentists



Spin structure of Lambda from Lambda polarization in Z^0 decay

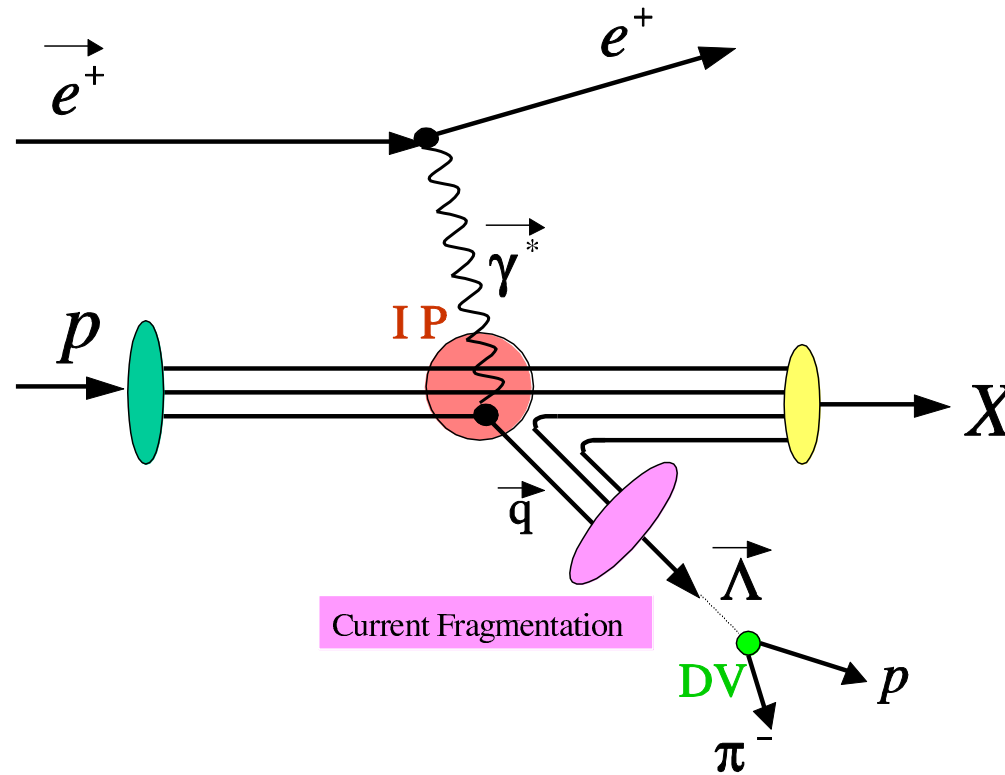


Diquark model and pQCD results

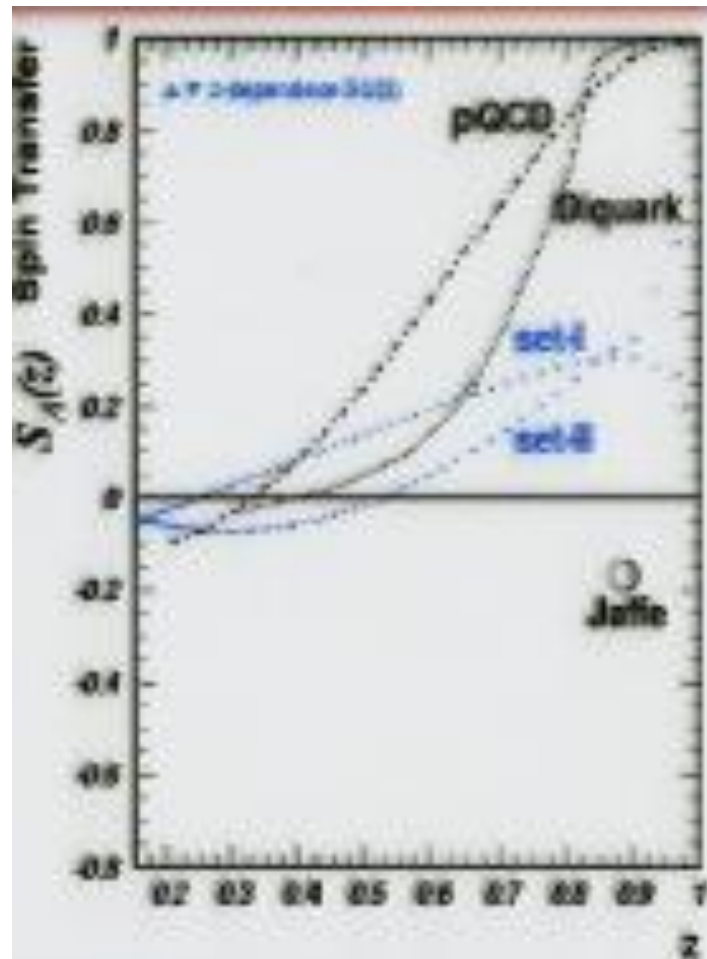


**B.-Q. Ma, I. Schmidt, J.-J. Yang,
Phys. Rev. D 61 (2000) 034017**

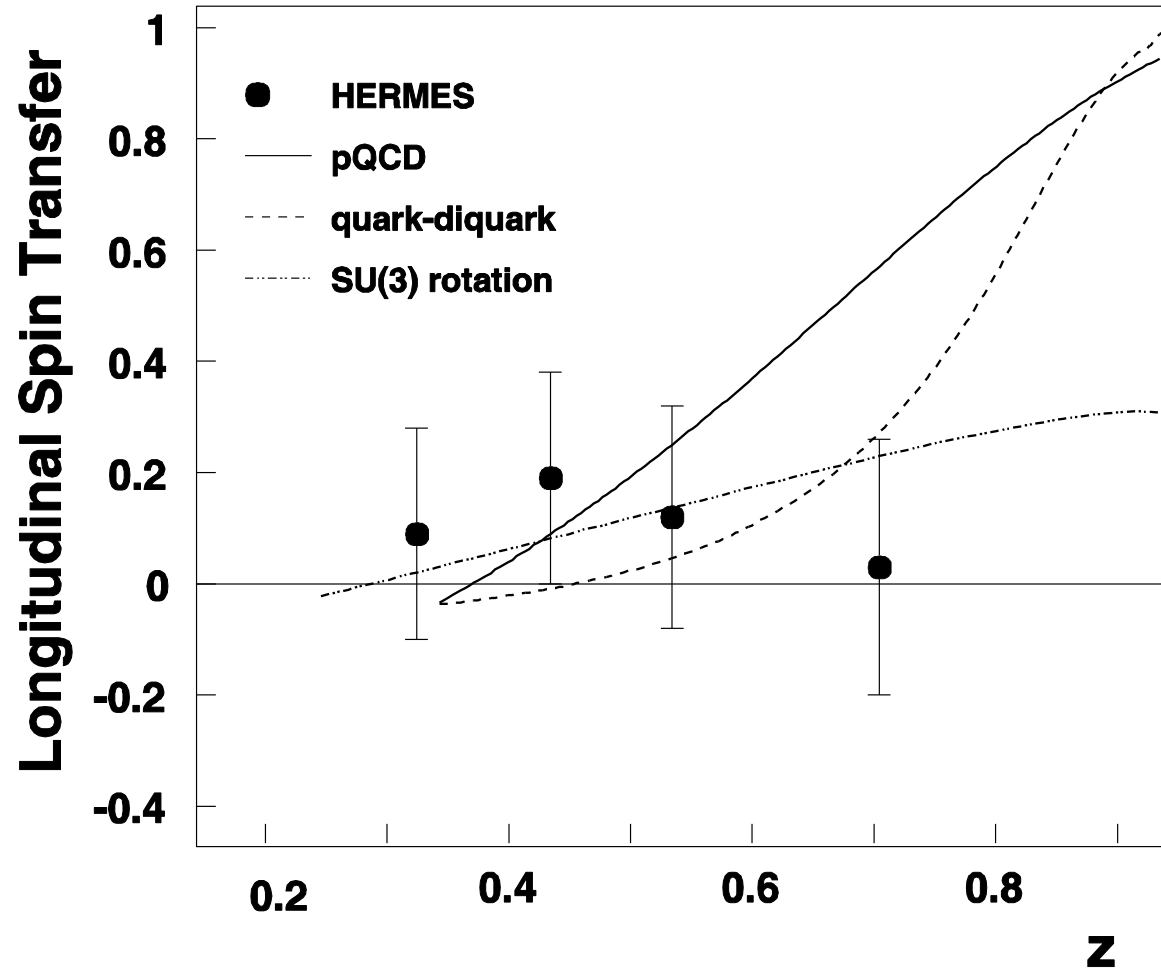
Spin Transfer to Λ in Semi-Inclusive DIS



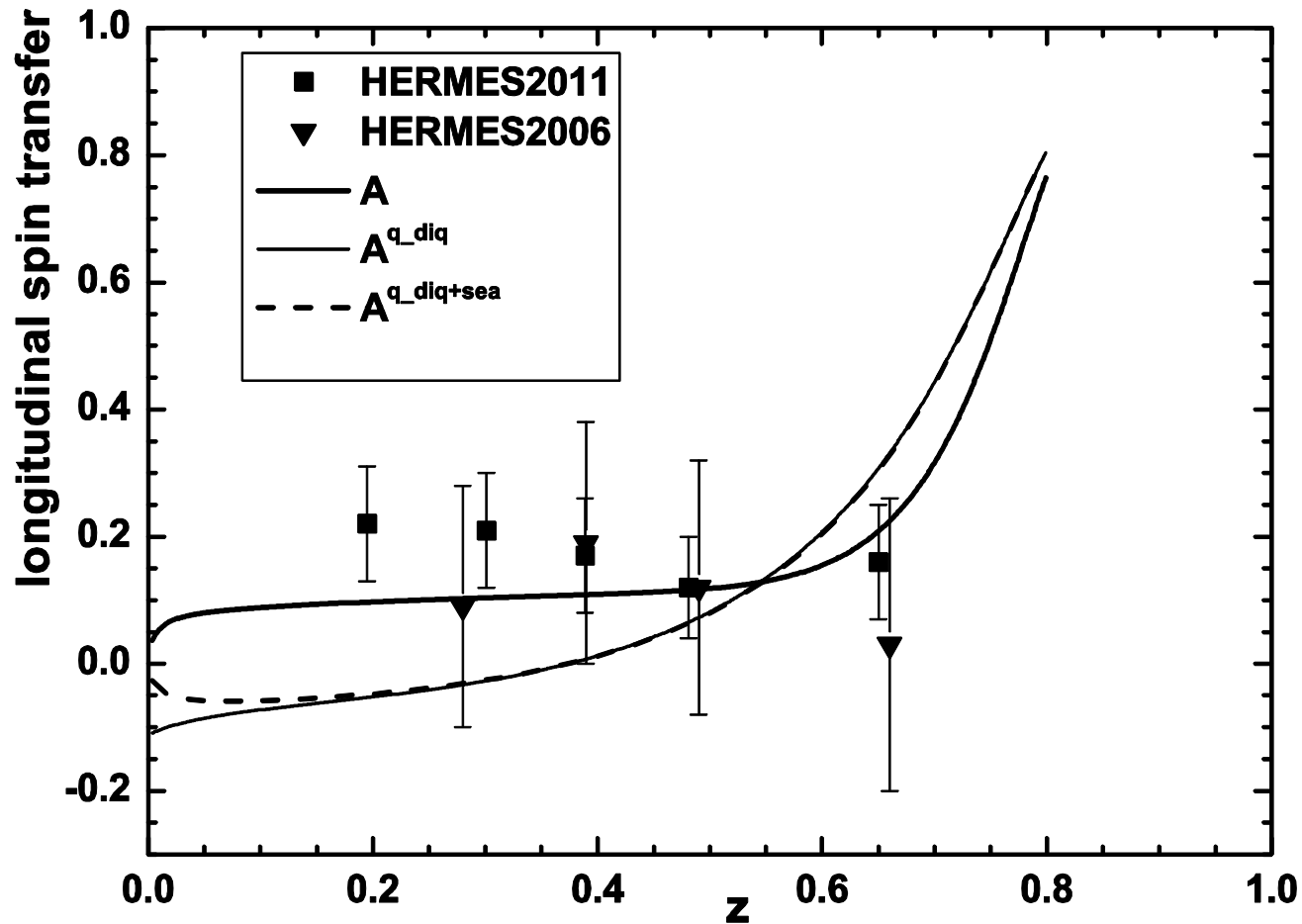
Different predictions



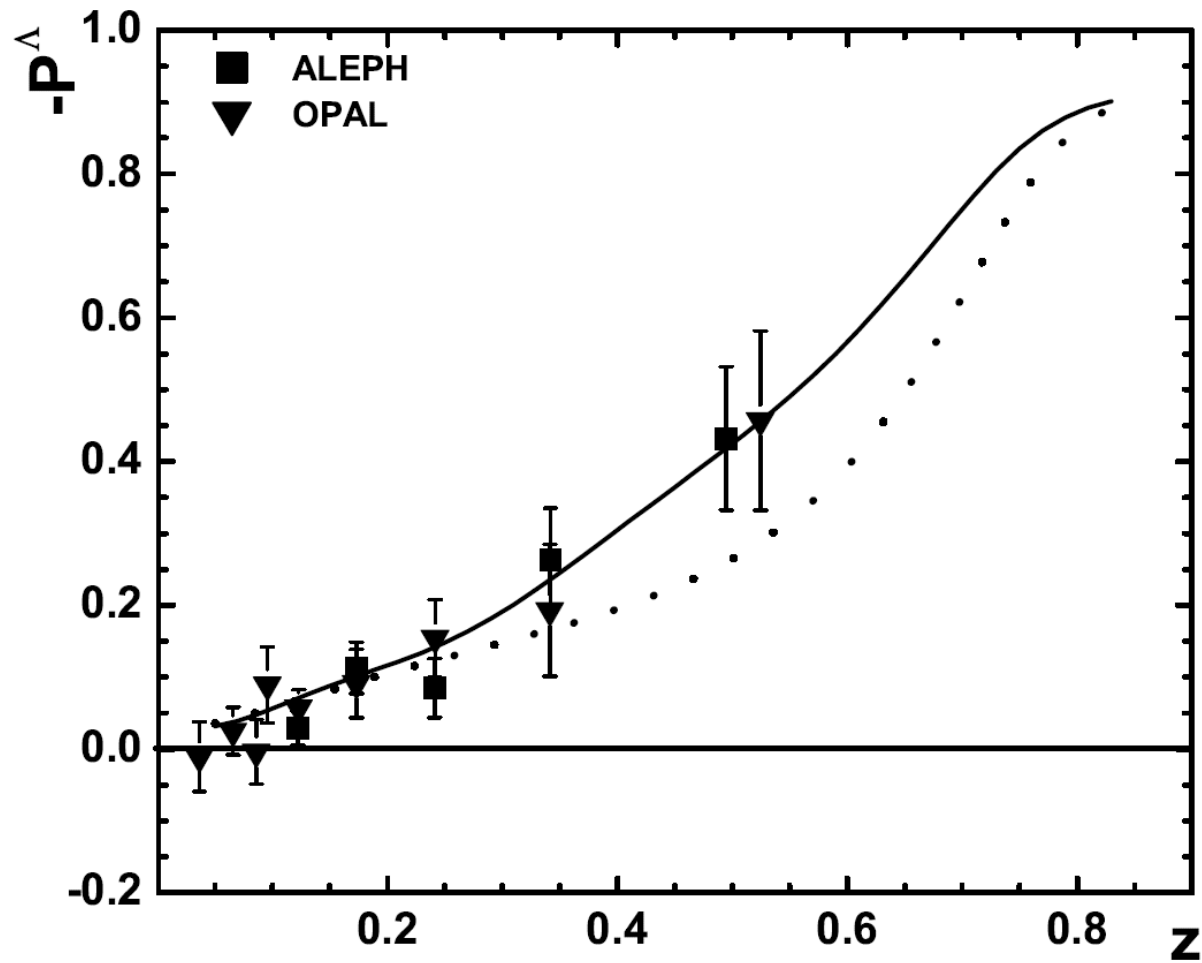
Comparison with data



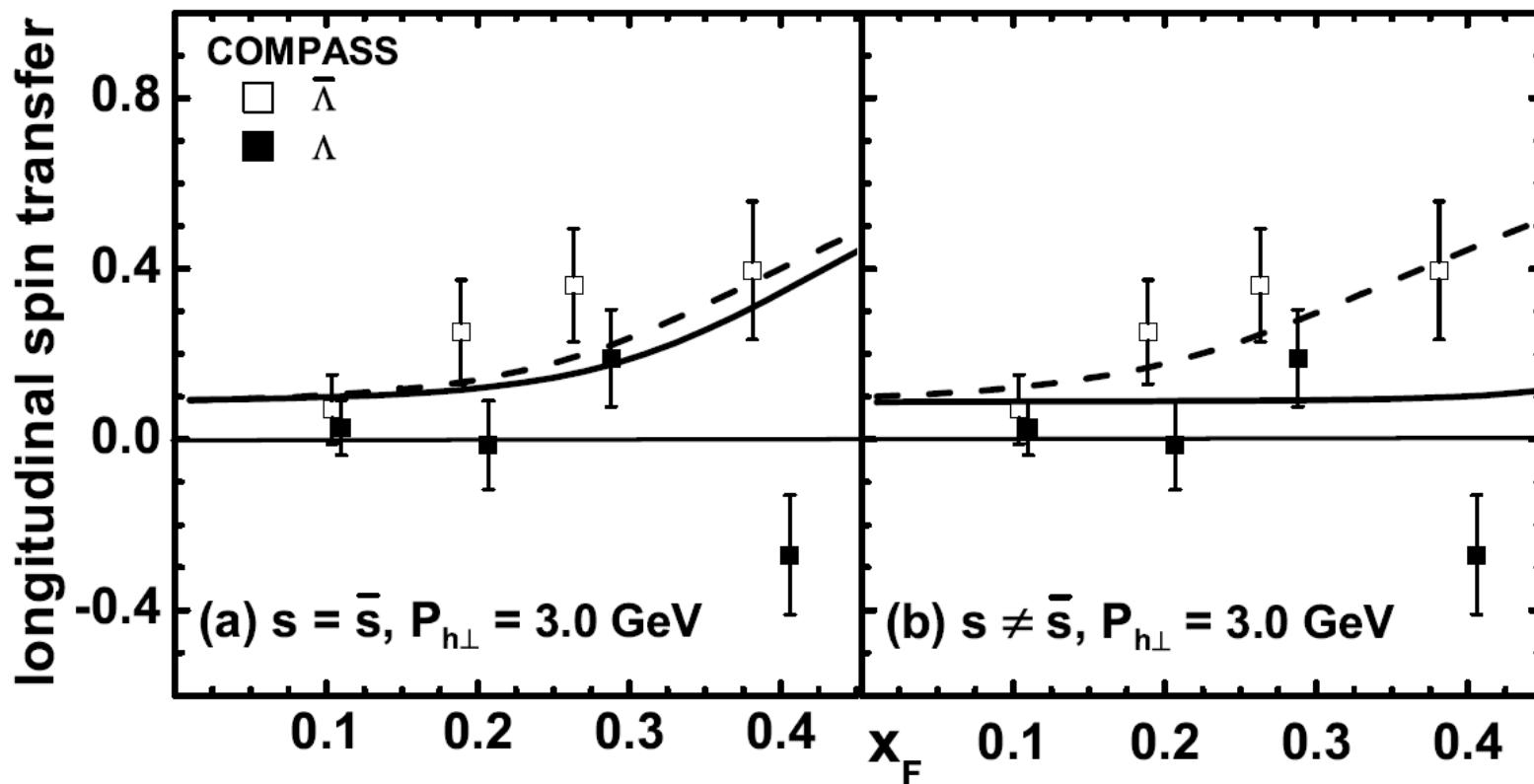
New results including both unfavored and indirect decays: SIDIS



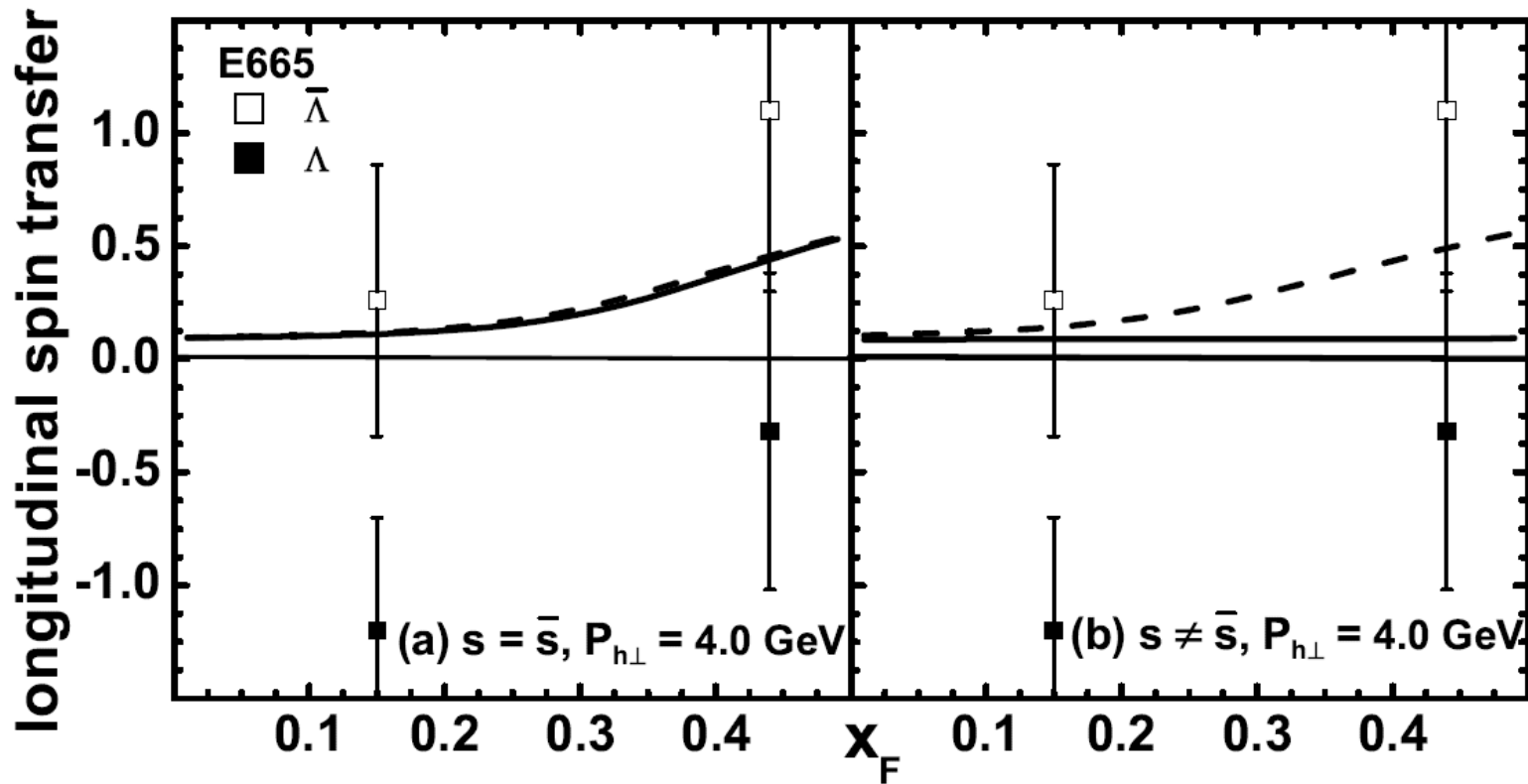
Results with new parametrization: Z-pole



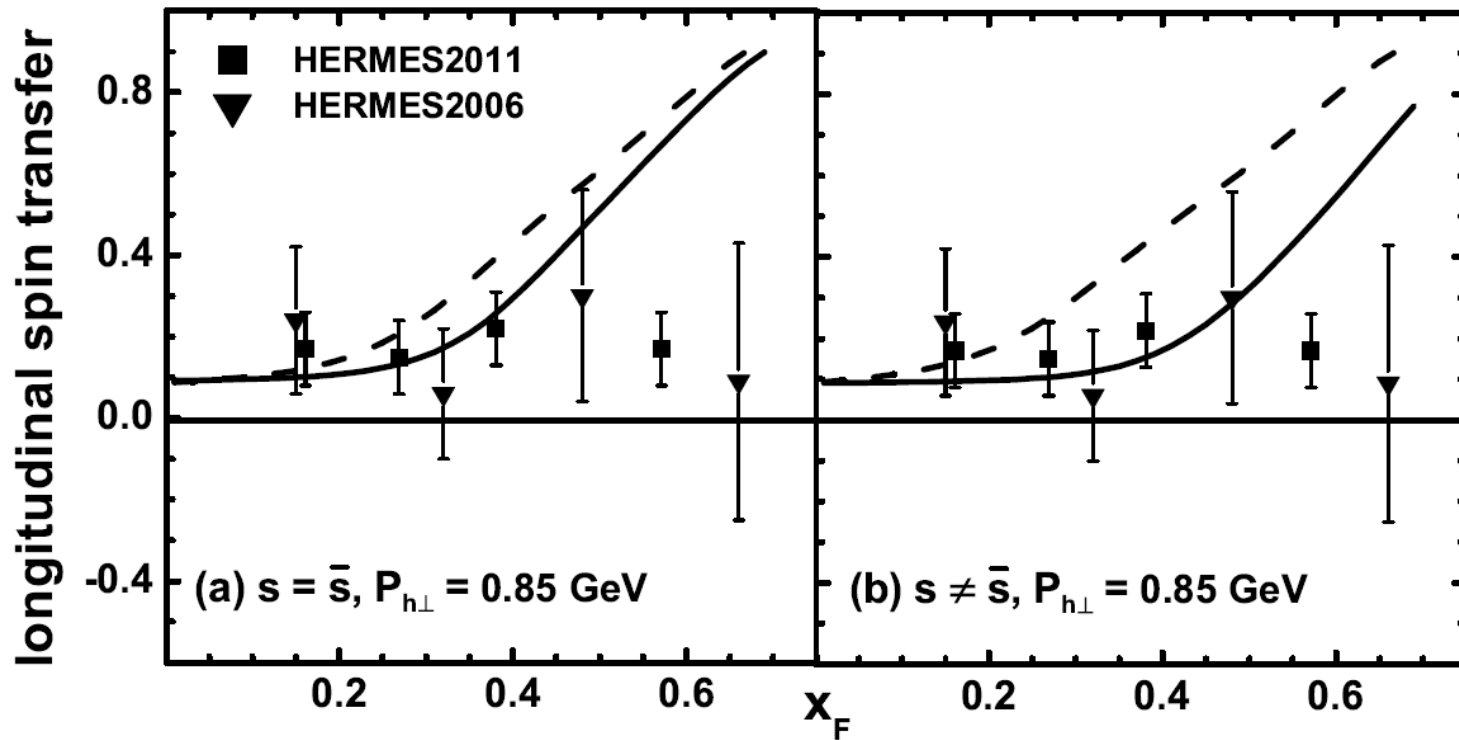
Difference between Lambda and anti-Lambda spin transfers with the COMPASS data



Difference between Lambda and anti-Lambda spin transfers with s-sbar asymmetry for E665



Difference between Lambda and anti-Lambda spin transfers with s-sbar asymmetry for HERMES

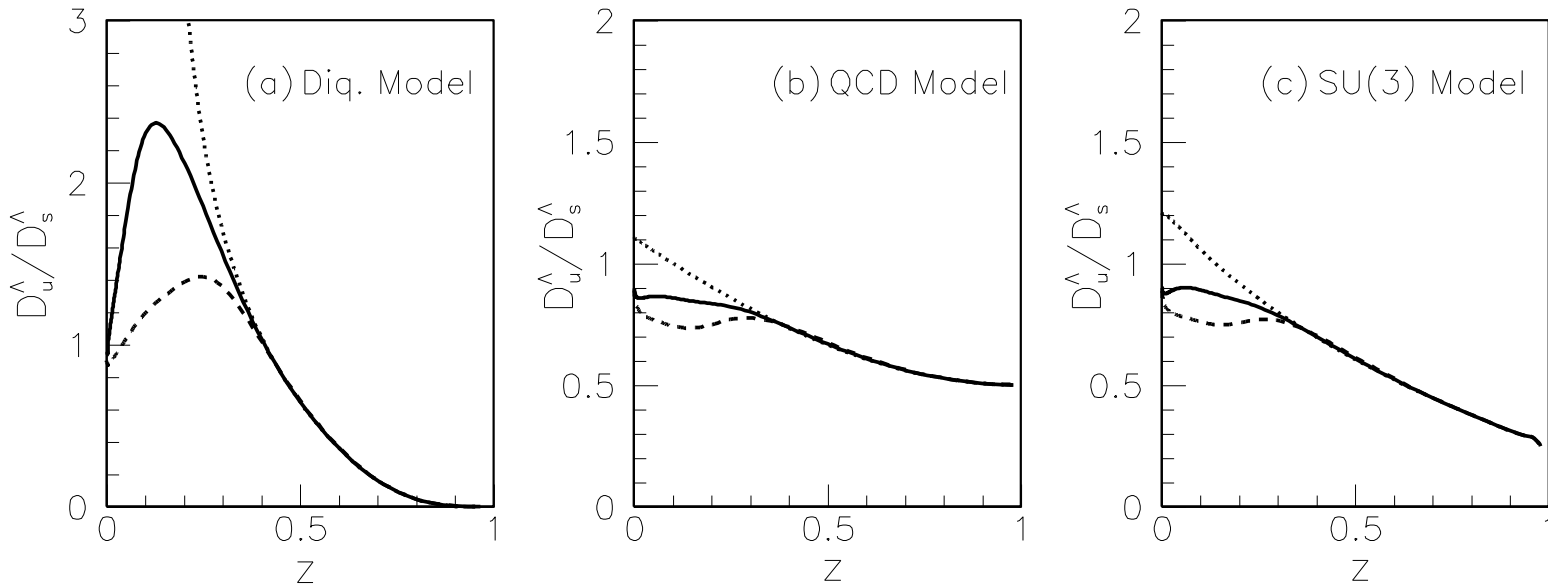


$\bar{\Lambda}/\Lambda$ Ratio in DIS Production

- A sensitive quantity that can provides information about the flavor structure of Λ hyperon.

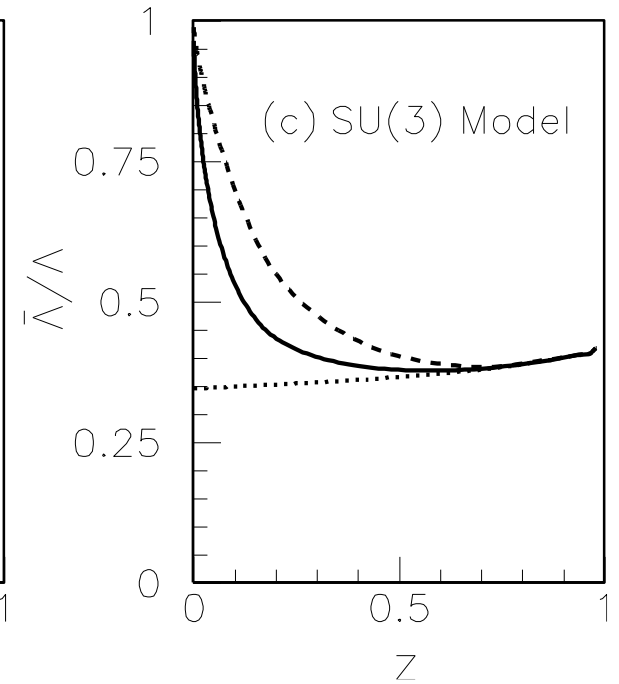
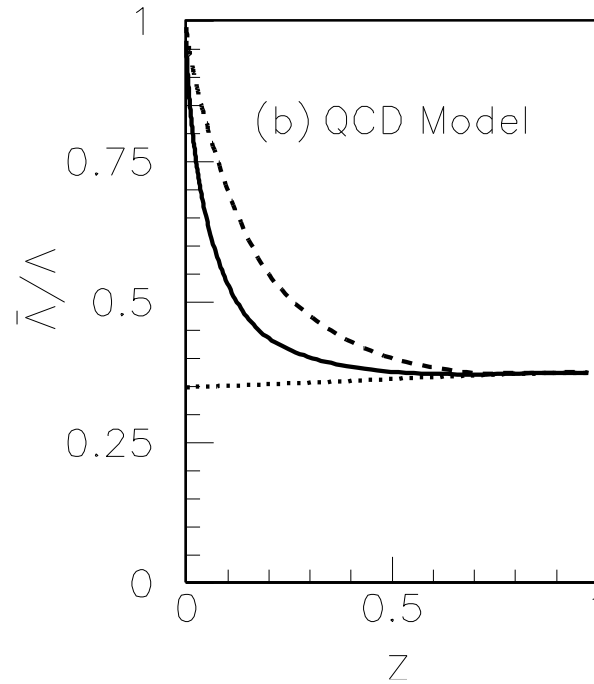
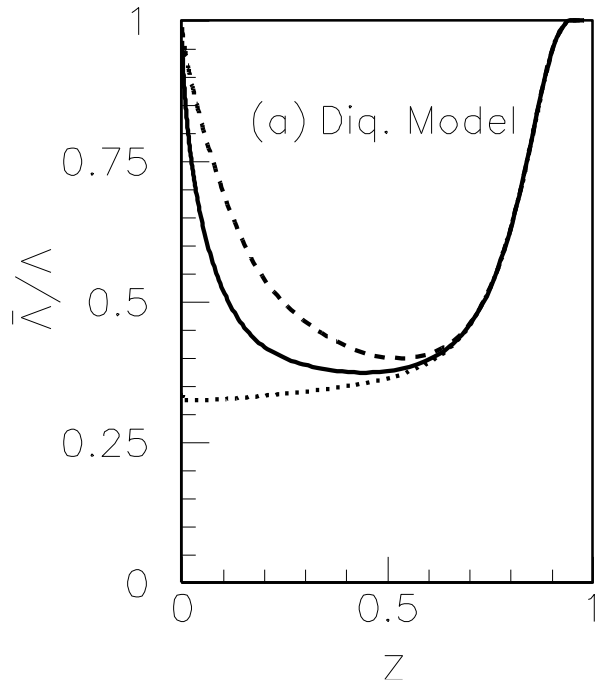
B.-Q. Ma, I. Schmidt, J.-J. Yang
Phys. Lett. B 574 (2003) 35

The flavor structure of Lambda u/s ratio with x-dependence



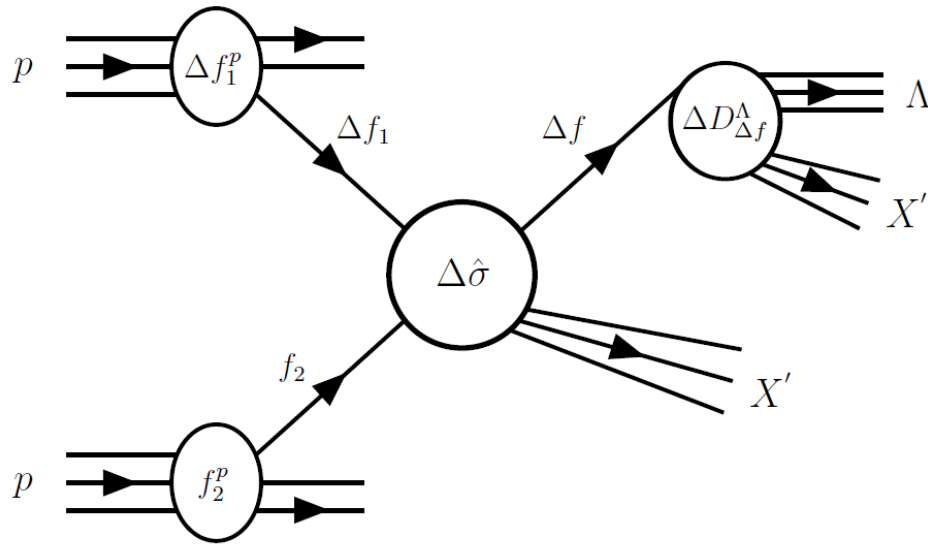
B.-Q. Ma, I. Schmidt, J.-J. Yang
Phys. Lett. B 574 (2003) 35

Different predictions



B.-Q. Ma, I. Schmidt, J.-J. Yang
Phys. Lett. B 574 (2003) 35

$$\vec{p}p \rightarrow \vec{\Lambda}X$$

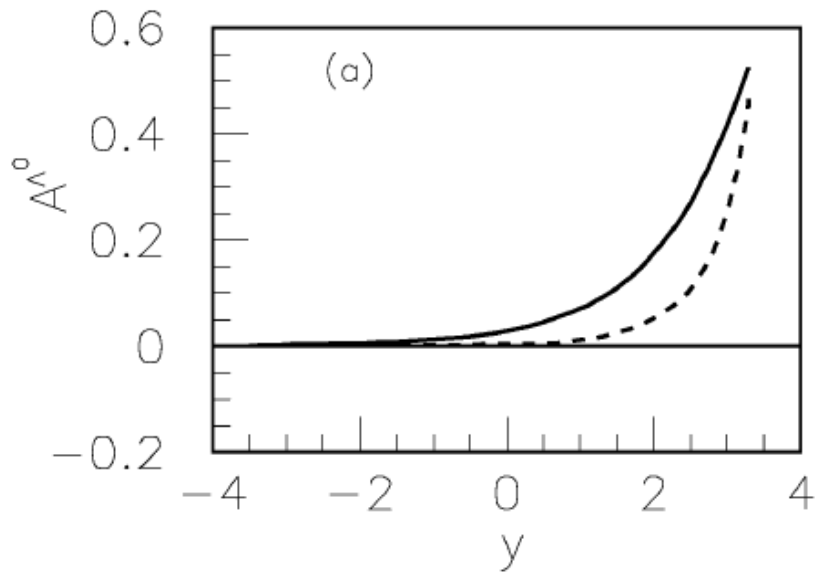


Providing information about

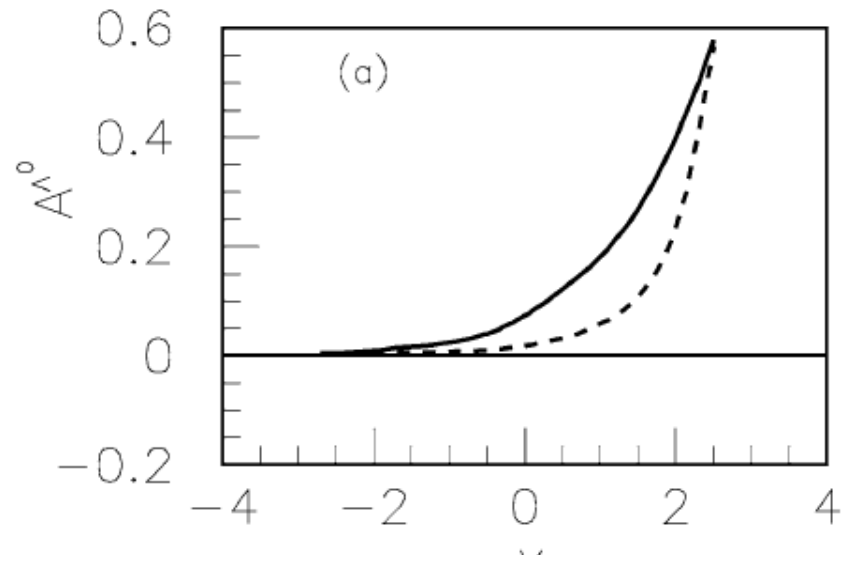
- the inclusive production of hadrons
- the strange and antistrange quark polarizations of the proton.

Spin transfer for

$$\vec{p}p \rightarrow \vec{\Lambda}X \text{ at RHIC-BNL}$$



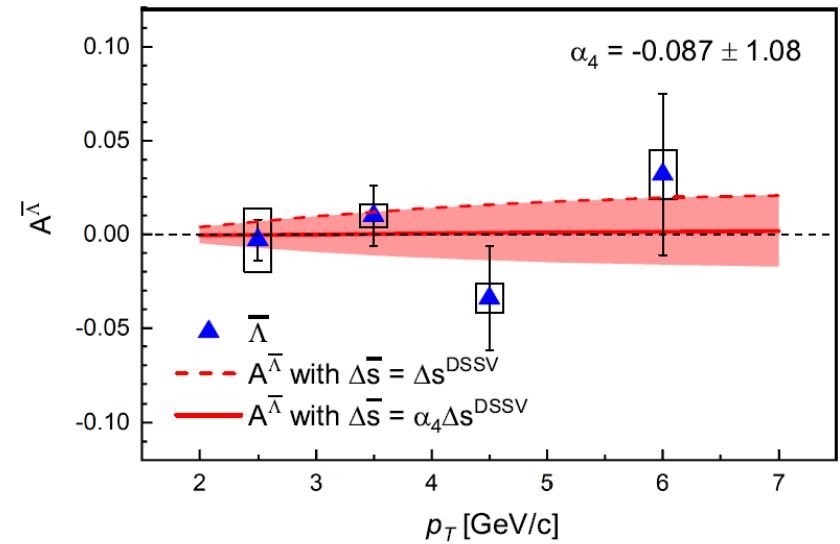
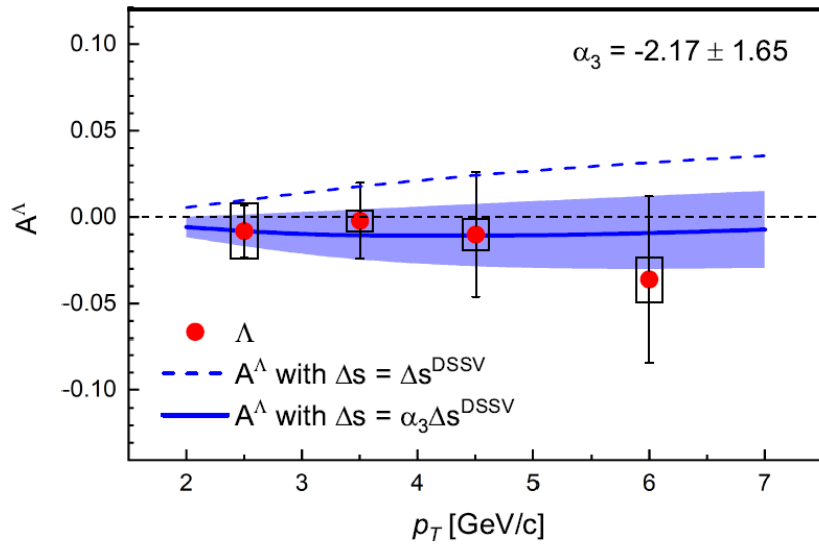
$$\sqrt{s} = 500 \text{ GeV,}$$



$$\sqrt{s} = 200 \text{ GeV,}$$

B.-Q. Ma, I. Schmidt, J.-J. Yang, J. Sofer, Nucl. Phys. A 703 (2002) 346

Fitting to STAR DATA



Xiaonan Liu, B.-Q. Ma, arXiv:1905.02360, EPJC 79 (2019) 409

Results from fitting STAR data

Table: Fitting results of α_i and calculated results of Δs and $\Delta \bar{s}$.

	value	Δs	$\Delta \bar{s}$	χ_{\min}^2
α_1	-1.20 ± 1.31	-0.014 ± 0.015		0.37
α_2	-0.24 ± 0.49		-0.003 ± 0.005	2.48
α_3	-2.17 ± 1.65	-0.025 ± 0.019		0.42
α_4	-0.087 ± 1.08		-0.001 ± 0.012	2.24

Two options: with/without gluon polarization


Comparison with Predictions & Results

The central values of the fitting results are basically compatible with

- the light-cone meson-baryon fluctuation model²⁴ prediction $\Delta s(x) \approx -0.05$ to -0.01 and $\Delta \bar{s}(x) \approx 0$.
- the recent lattice QCD determination²⁵, $\Delta s^+ = -0.02(1)$ at $Q^2 \approx 7\text{GeV}^2$.
- the results from Jefferson Lab Angular Momentum (JAM) Collaboration²⁶ $\Delta s^+(Q_0^2) = -0.03(10)$.

²⁴S. J. Brodsky and B.-Q. Ma, Phys. Lett. B 381, 317 (1996).

²⁵G. S. Bali et al. [QCDSF Collaboration], Phys. Rev. Lett. 108, 222001 (2012)

²⁶J.J.Ethier, N.Sato and W.Melnitchouk, Phys. Rev. Lett. 119, 132001 (2017) 

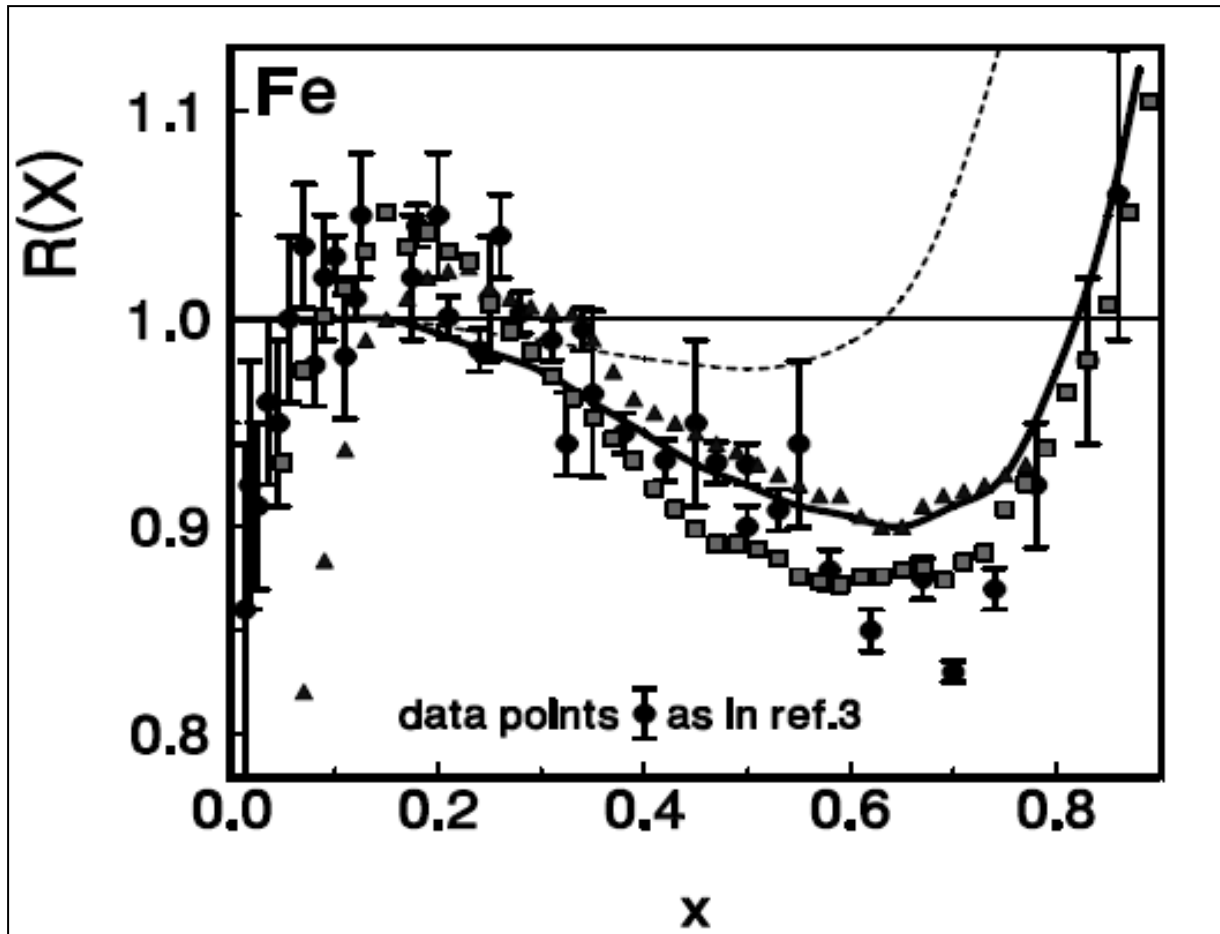
Relating Λ -production with nucleon strangeness

- The spin transfer process of $\vec{p}p \rightarrow \vec{\Lambda}X$ is feasible to study strange-antistrange polarizations of the nucleon.
- The fitting to STAR data suggests:

$$\Delta s \neq \Delta \bar{s}$$
$$\Delta s \approx -0.025 \pm 0.019$$
$$\Delta \bar{s} \approx -0.001 \pm 0.012$$

- The results are compatible with the light-cone baryon-meson fluctuation model prediction.

Nuclear EMC Effect

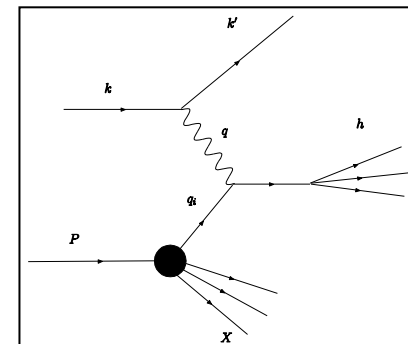


Anti-Lambda production as a probe of the nuclear sea structure?

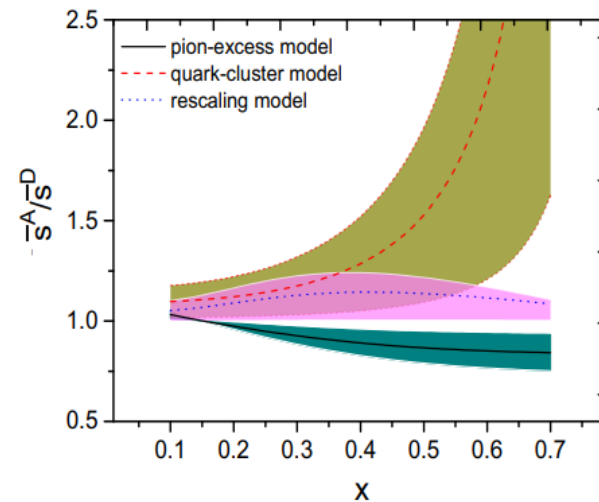
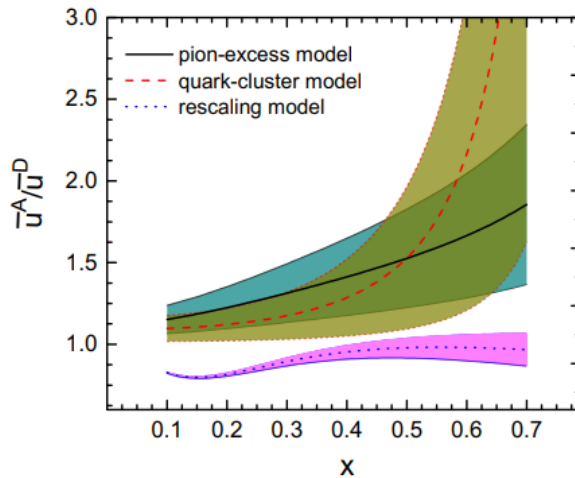
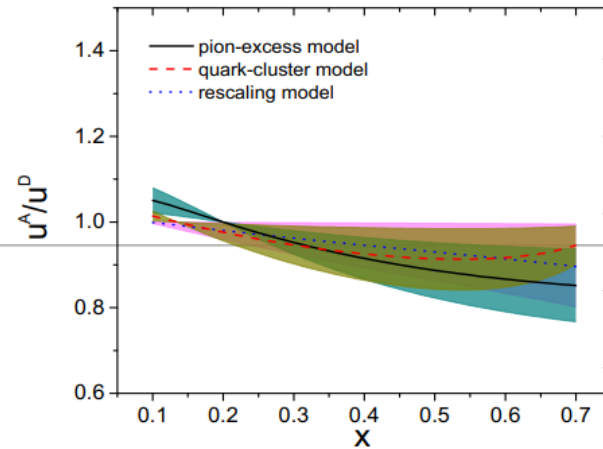
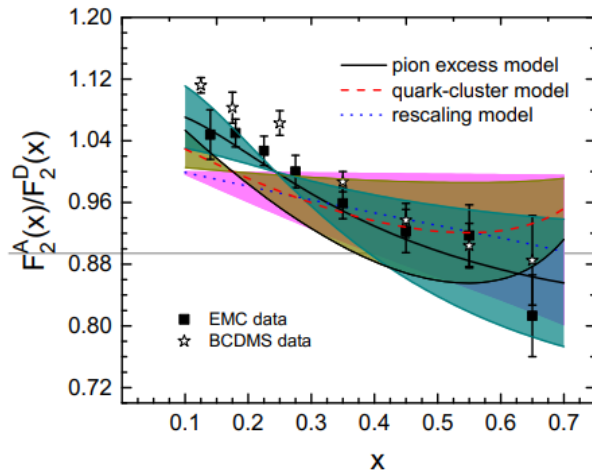
- 有三类原子核模型可以解释EMC效应: 团簇模型, Pi盈余模型, 重新标度模型.
- 通过考察原子核与核子荷电轻子半单举过程中末态强子anti-Lambda的产率比 x 的依赖性, 我们发现 anti-Lambda能够区分定性描述原子核EMC效应的三类不同原子核结构模型.

B.Lu, B.-Q. Ma, Phys.Rev.C74 (2006) 055202

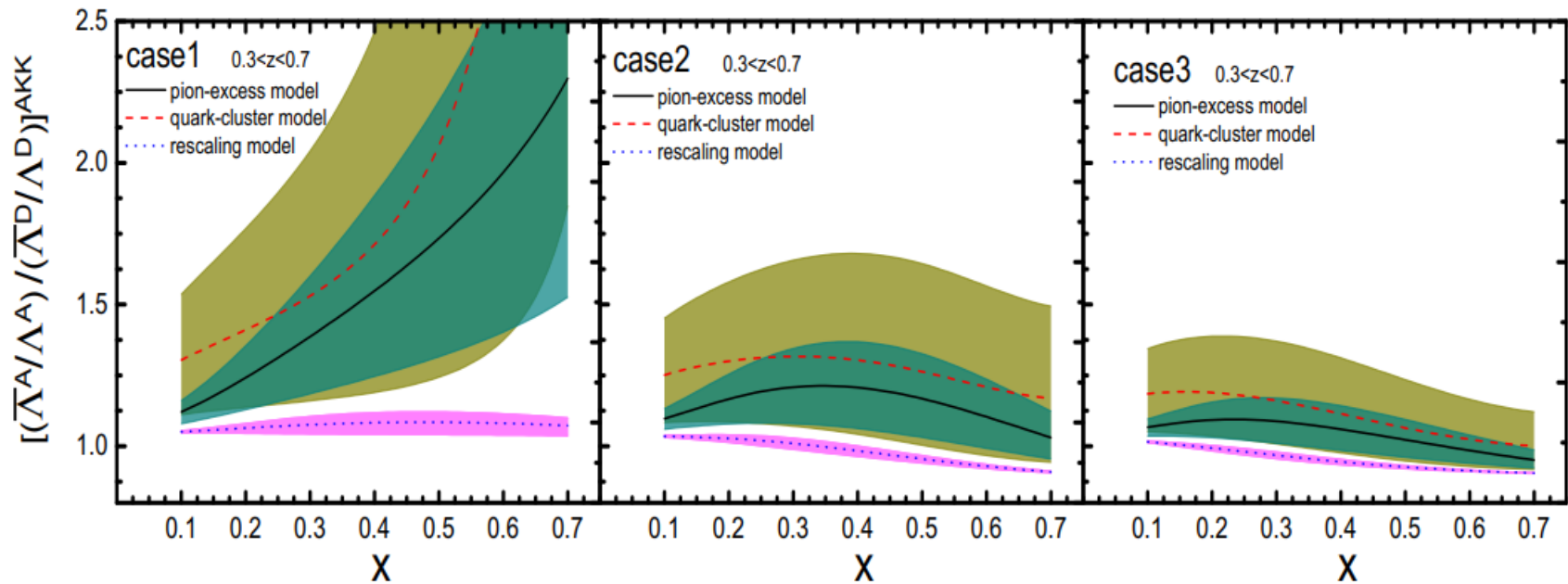
C.Gong, B.-Q. Ma, Phys.Rev.C97 (2018) 065207



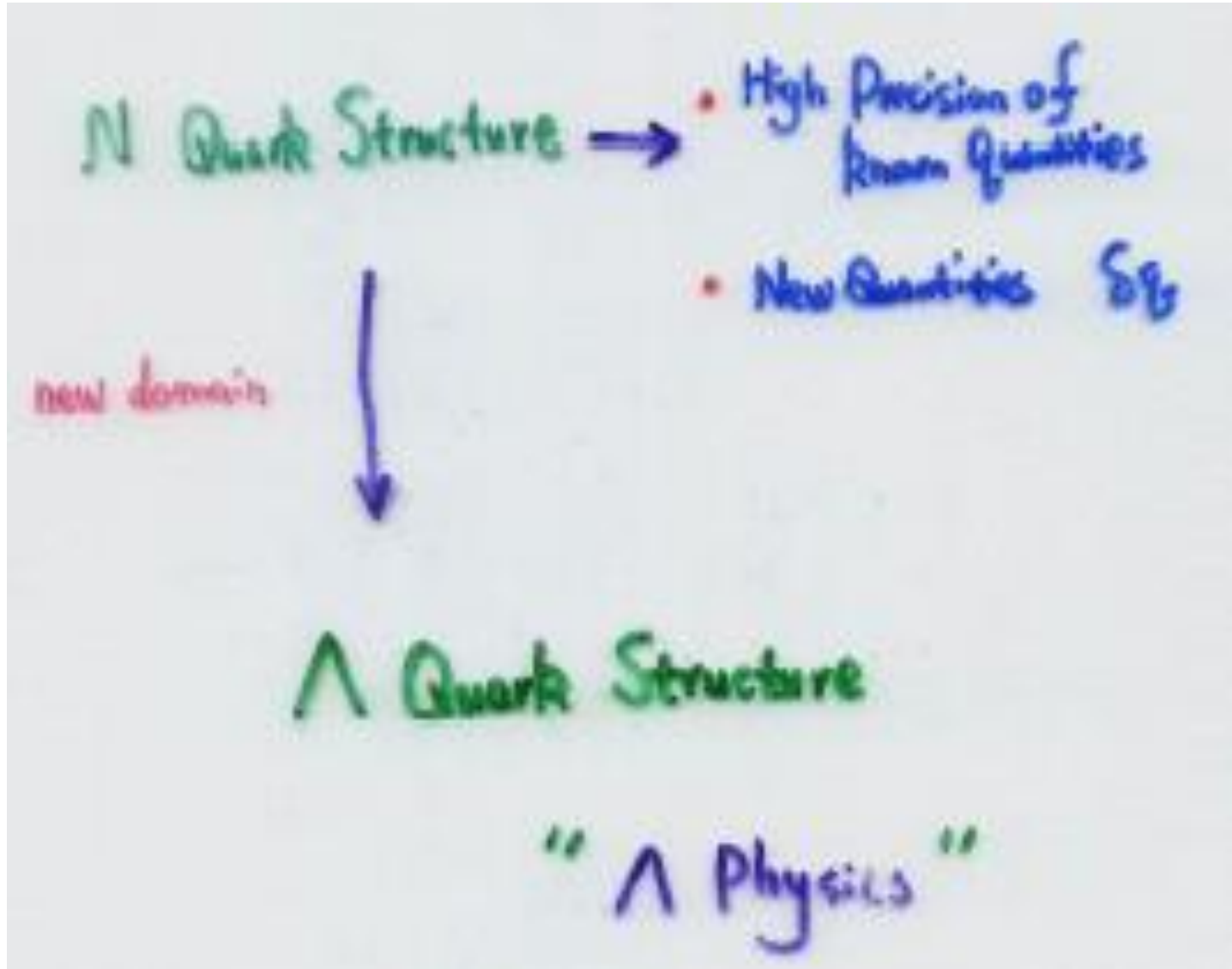
Different sea behaviors in nuclei



Different predictions of anti-lambda production



Conclusion 1: Lambda Physics



Conclusion 2: Relating Λ -production with nucleon strangeness

- The spin transfer process of $\vec{p}p \rightarrow \vec{\Lambda}X$ is feasible to study strange-antistrange polarizations of the nucleon.
- The fitting to STAR data suggests:

$$\Delta s \neq \Delta \bar{s}$$
$$\Delta s \approx -0.025 \pm 0.019$$
$$\Delta \bar{s} \approx -0.001 \pm 0.012$$

- The results are compatible with the light-cone baryon-meson fluctuation model prediction.

Conclusion 3: Anti-Lambda Production for Nuclear Physics

Anti-Lambda production charged lepton semi-inclusive deep inelastic scattering off nuclear target is ideal to figure out the nuclear sea content, which is differently predicted by different models accounting for the nuclear EMC effect.

B.Lu, B.-Q. Ma, Phys.Rev.C74 (2006) 055202

C.Gong, B.-Q. Ma, Phys.Rev.C97 (2018) 065207