Fragmentation functions for doubly heavy mesons

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Outline

1. Background

2. Fragmentation functions at NLO

3. Applications

4. Summary

1. Background

Top quark cannot form a hadron

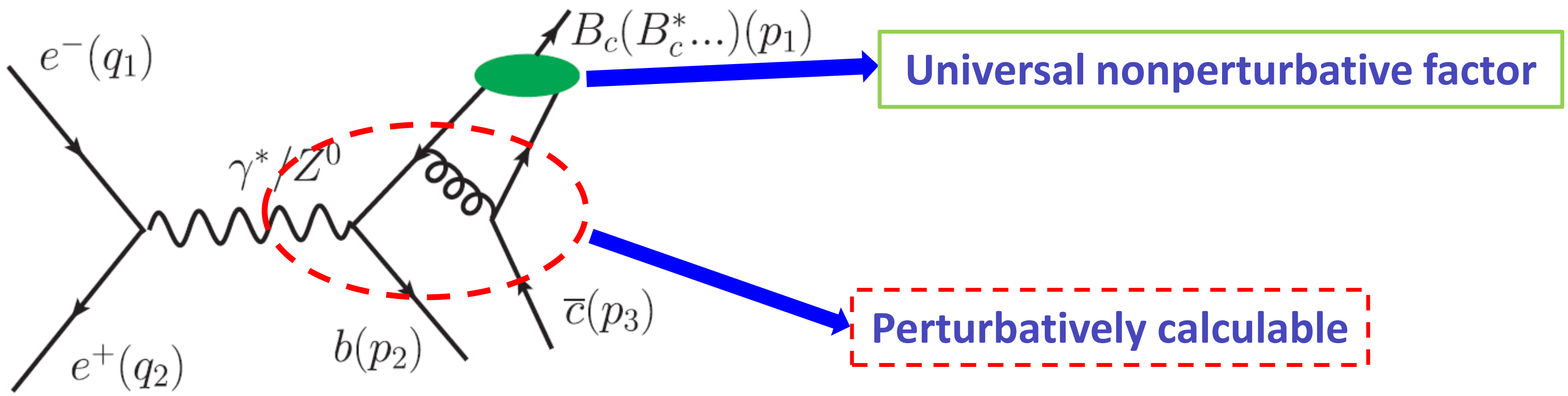
Doubly heavy mesons：
$c\bar{c}$, $b\bar{b}$ (J/ψ, Υ, etc.), $c\bar{b}$(B_c, B_c^*, etc.)

- **Production： perturbative, non-perturbative QCD**

  To test pQCD, NRQCD

- **Decay： weak interaction**

- NRQCD factorization



$$
\begin{aligned}
& d\sigma(e^+ + e^- \to Bc + b + \bar{c}) \\
& = \sum_n d\hat{\sigma}(e^+ + e^- \to c\bar{b}[n] + b + \bar{c}) \left\langle O^{Bc}(n) \right\rangle
\end{aligned}
$$

NRQCD factorization

Short-distance coefficients

Long-distance matrix elements

- NRQCD factorization

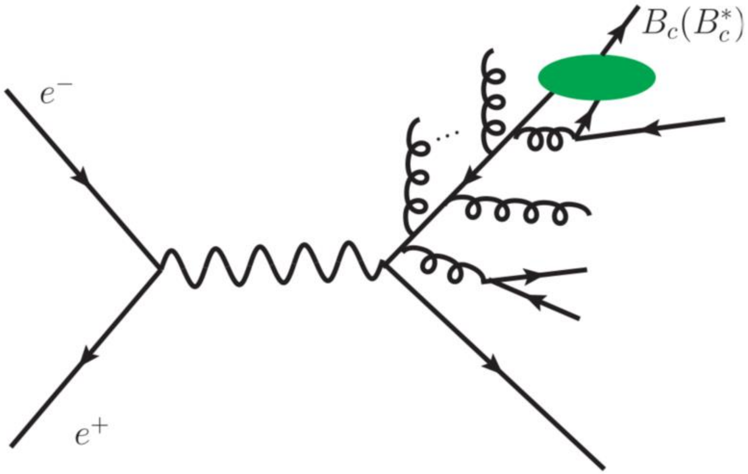
$$d\sigma(e^+ + e^- \to Bc + b + \bar{c})$$
$$= \sum_n d\sigma(e^+ + e^- \to (c\bar{b})[n] + b + \bar{c})\langle O^{\mathrm{Bc}}(n)\rangle$$

Energy scales:
$\sqrt{s}, m_Q$

Log-terms appear in short-distance coefficients:

$$\boxed{\alpha_s^m \sum_{n=0}^{\infty} \alpha_s^n \ln^n(s/m_Q^2)}$$

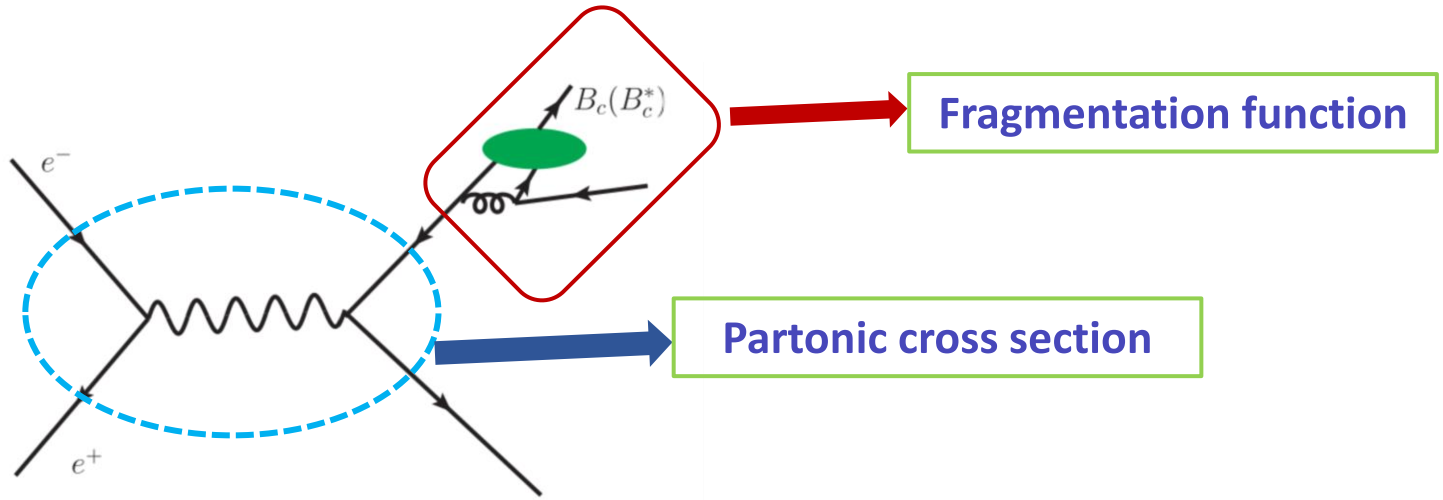
Collinear emission



Spoil or weak the convergence of the series

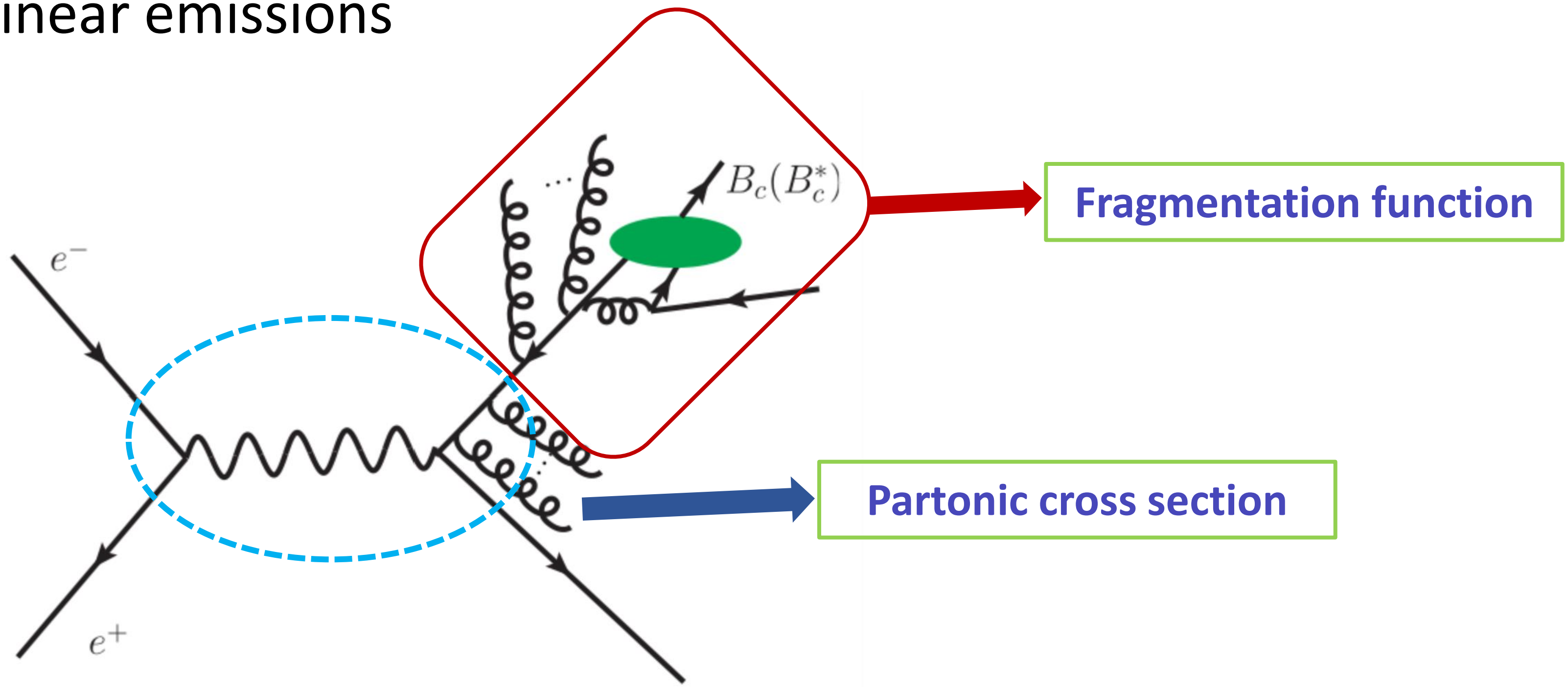
$\ln(p_t^2/m_Q^2)$ appearing in the production at a hadron collider

- Fragmentation-function approach



$$d\sigma(e^+ + e^- \to Bc(p) + b + \bar{c})$$
$$= \sum_i d\hat{\sigma}(e^+ + e^- \to i + X)(p/z, \mu_F) \otimes D_{i\to Bc}(z, \mu_F) + \mathrm{O}(m_Q^2 / s)$$

Collinear emissions



Large Logrithms of s/m_Q^2?

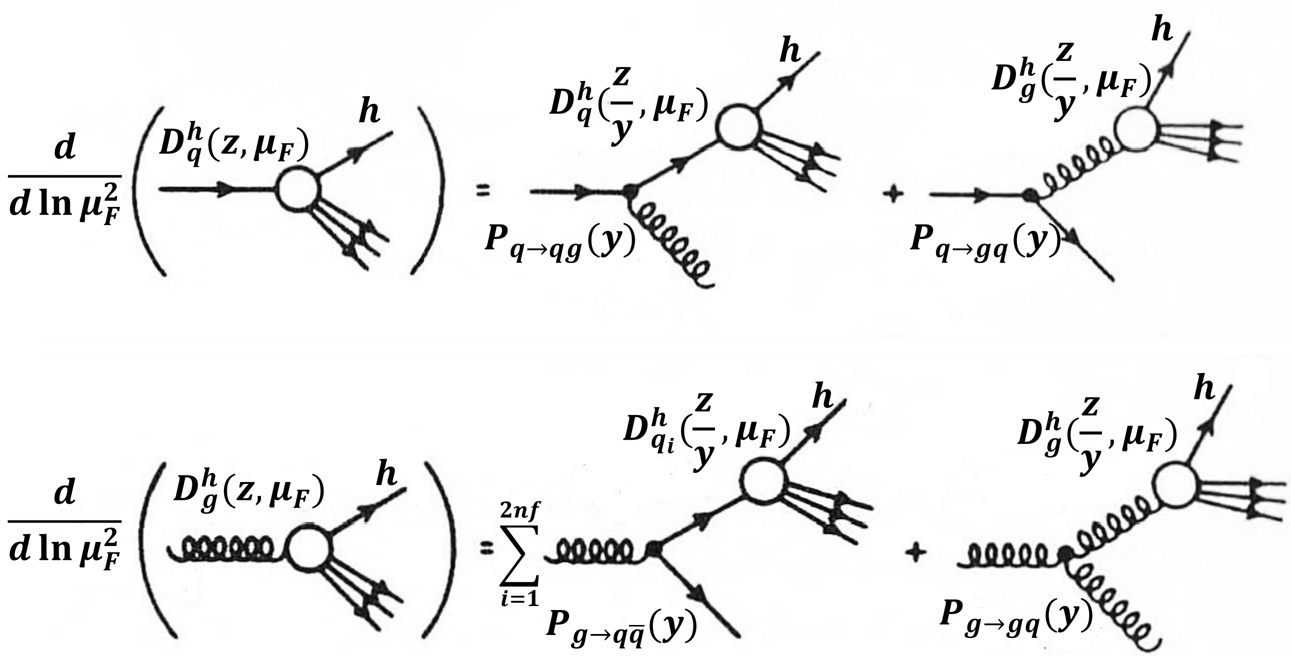
Involving $\ln(s/\mu_F^2)$

$\mu_F = O(\sqrt{s})$

$$d\sigma(e^+ + e^- \to Bc(p) + b + \bar{c})$$
$$= \sum_i \boxed{d\hat{\sigma}(e^+ + e^- \to i + X)(p/z, \mu_F)} \otimes \boxed{D_{i\to Bc}(z, \mu_F)} + O(m_Q^2/s)$$

Involving $\ln(\mu_F^2/m_Q^2)$

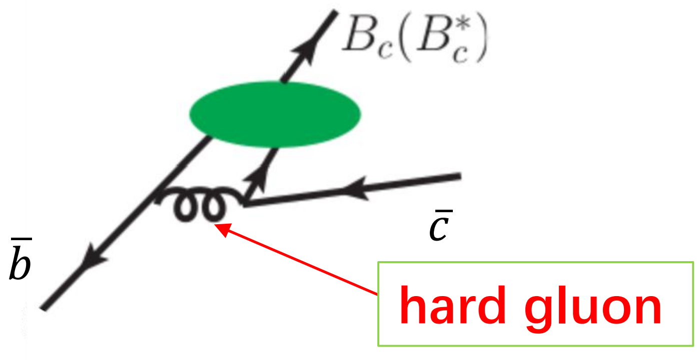
DGLAP evolution



Collinear log-terms can be resummed through the DGLAP evolution.

Boundary condition: $D_i^h(z,\mu_{F0})$ with $\mu_{F0} = O(m_Q)$

NRQCD factorization for FFs:



The fragmentation process contains perturbatively calculable information

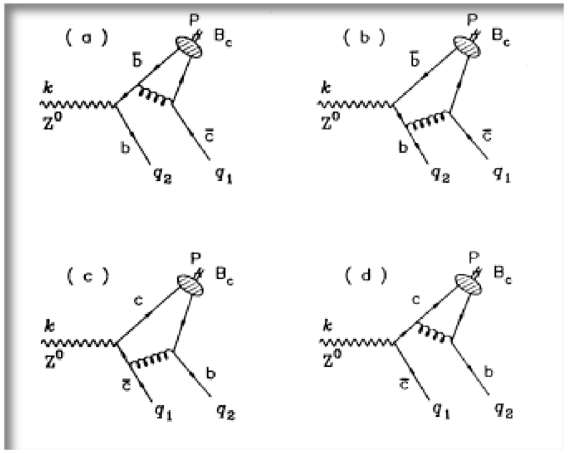
$$D_{i\to Bc}(z,\mu_{F0}) = \sum_n d_{i\to c\bar{b}[n]}(z,\mu_{F0})\left\langle O^{\mathrm{Bc}}(n)\right\rangle$$

Involving $\ln(\mu_{F0}^2/m_Q^2)$ $\mu_{F0} = O(m_Q)$

The FFs at a higher scale can be obtained by solving the DGLAP equations

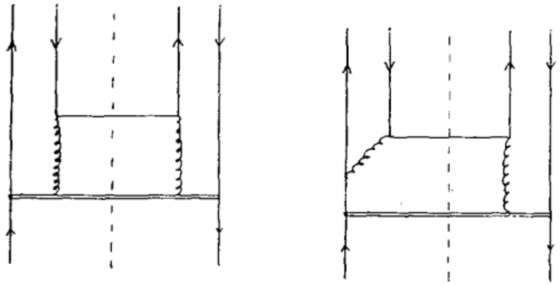
➢ LO fragmentation functions for the Bc production

- Extracting from the LO calculation of process $Z^0 \to Bc + b + \bar{c}$



C.-H. Chang, Y.-Q. Chen, Phys. Rev. D 46, 3845, (1992);
E. Braaten, K. Cheung, T.C. Yuan, D 48, R5049, (1993).

- Calculating from the definition:



J.-P. Ma, Phys. Lett. B 332, 398, (1994).

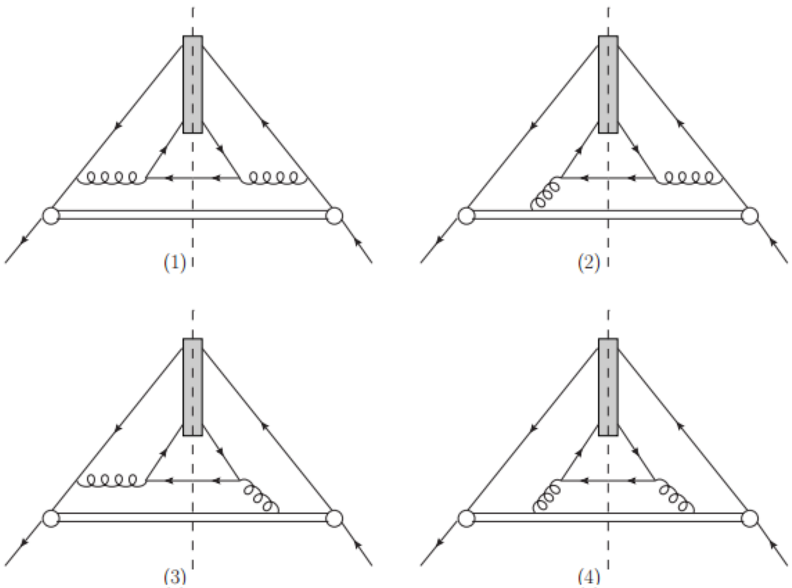
2. FFs at NLO

Based on the definition of FFs by **Collins and Soper.**

Nucl. Phys. B 194, 445, (1982).

Process independent approach

LO cut diagrams:



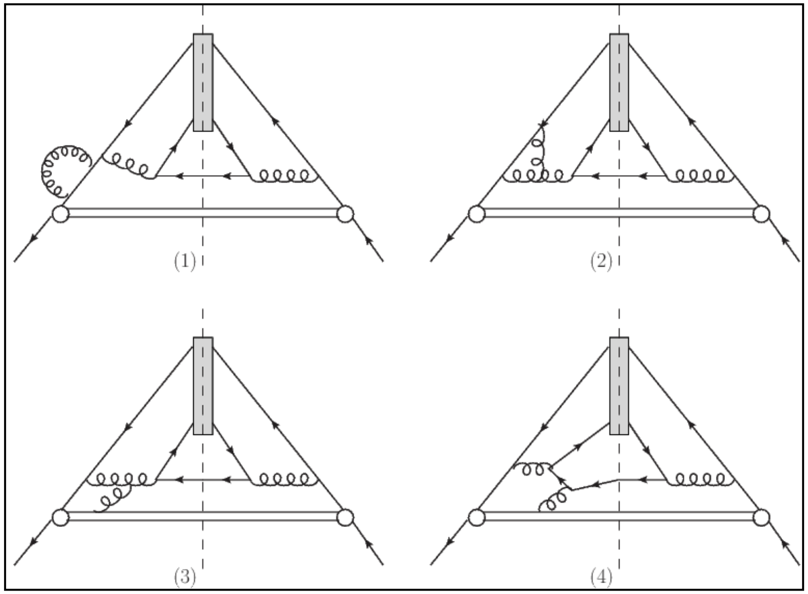
LO fragmentation functions:

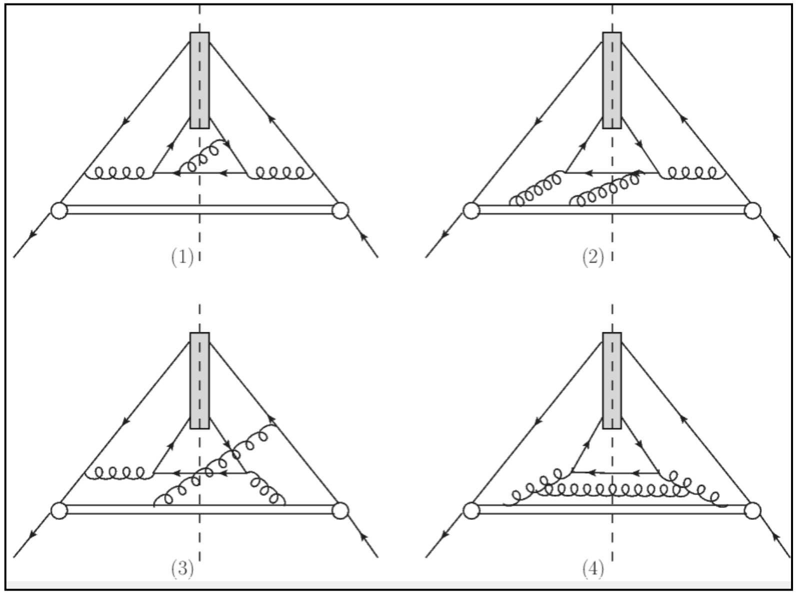
$$D^{\mathrm{LO}}_{\bar{b}\to B_c}(z) = \frac{2\alpha_s^2 z(1-z)^2|R_S(0)|^2}{81\pi r_c^2(1-r_b z)^6 M^3}[6-18(1-2r_c)z+(21-74r_c+68r_c^2)z^2 -2r_b(6-19r_c+18r_c^2)z^3+3r_b^2(1-2r_c+2r_c^2)z^4],$$

$$D^{\mathrm{LO}}_{\bar{b}\to B_c^*}(z) = \frac{2\alpha_s^2 z(1-z)^2|R_S(0)|^2}{27\pi r_c^2(1-r_b z)^6 M^3}[2-2(3-2r_c)z+3(3-2r_c+4r_c^2)z^2 -2r_b(4-r_c+2r_c^2)z^3+r_b^2(3-2r_c+2r_c^2)z^4].$$

NLO corrections

Sample NLO cut diagrams





54 virtual cut diagrams， 72 real cut diagrams.

➢ Virtual corrections

Loop-integral reduction

Many integrals containing an eikonal line, e.g,

$$\int \frac{d^D l}{[(l-p_1)^2 - m_1^2 + i\varepsilon][(l-p_2)^2 - m_2^2 + i\varepsilon][(l-p_3)^2 - m_3^2 + i\varepsilon](l\cdot n + i\varepsilon)}$$

➢ Real corrections

UV and IR divergences!

$$\mathrm{D}_{\bar{b}\to Bc}^{real}(z) = \int N_{CS} d\phi_{real} (A_{real} - A_S) + \int N_{CS} d\phi_{real} A_S$$

Calculated in 4 dimensions

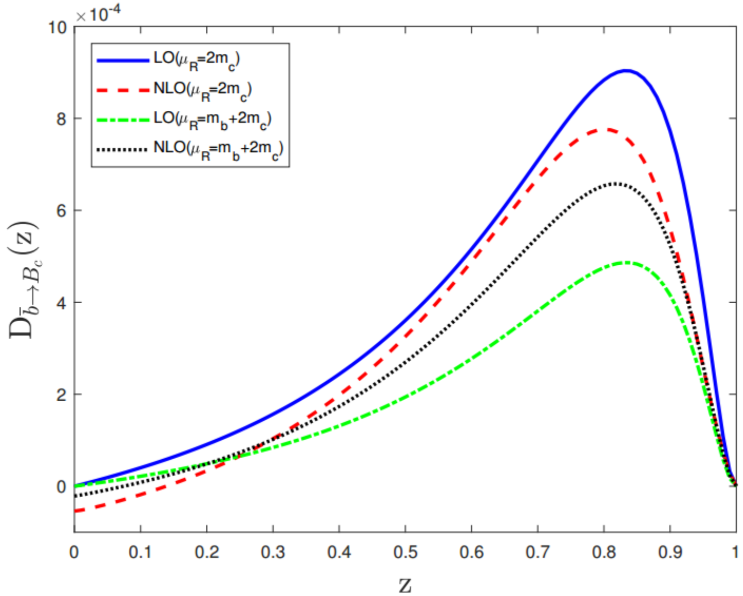
Calculated in d dimensions

Various types of subtraction terms need to be integrated!

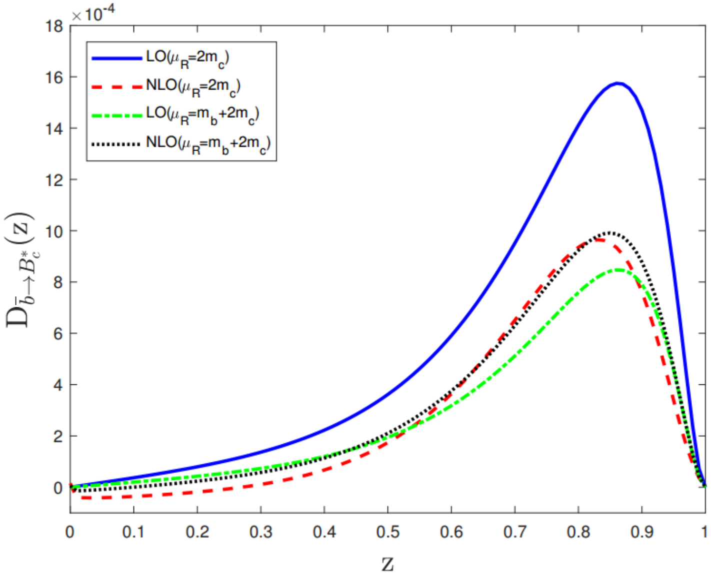
NLO results

Phys. Rev. D 100, 034004, (2019),
X.-C. Zheng, C.-H. Chang, T.-F Feng, X.-G. Wu.

NLO fragmentation functions for $\bar{b} \to B_c$ and $\bar{b} \to B_c^*$

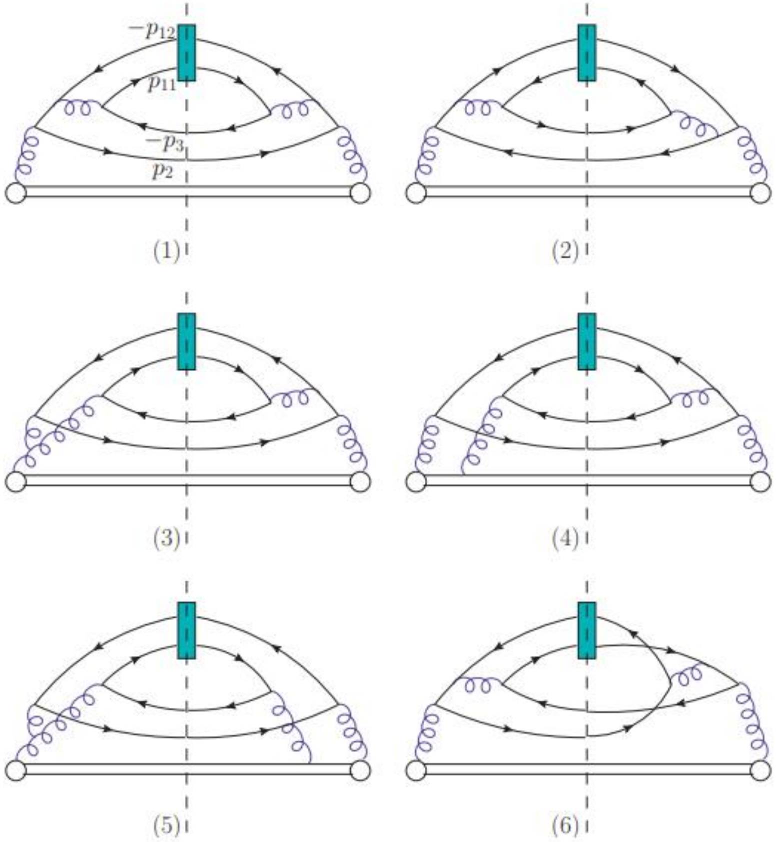


$$D_{\bar{b}\to Bc}(z, \mu_F = m_b + 2m_c)$$



$$D_{\bar{b}\to B_c^*}(z, \mu_F = m_b + 2m_c)$$

g->Bc(Bc*) FFs

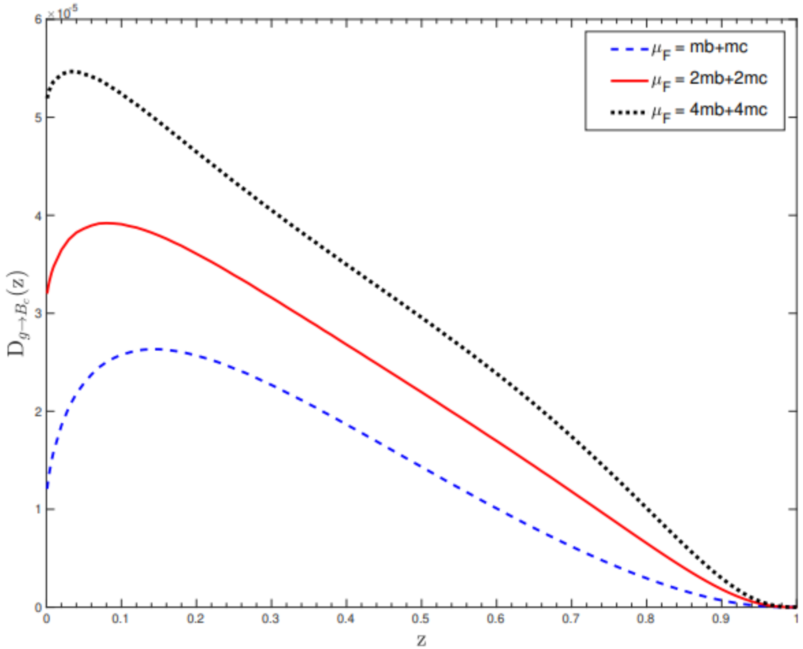


末态有三个粒子,且相空间积分存在紫外发散。

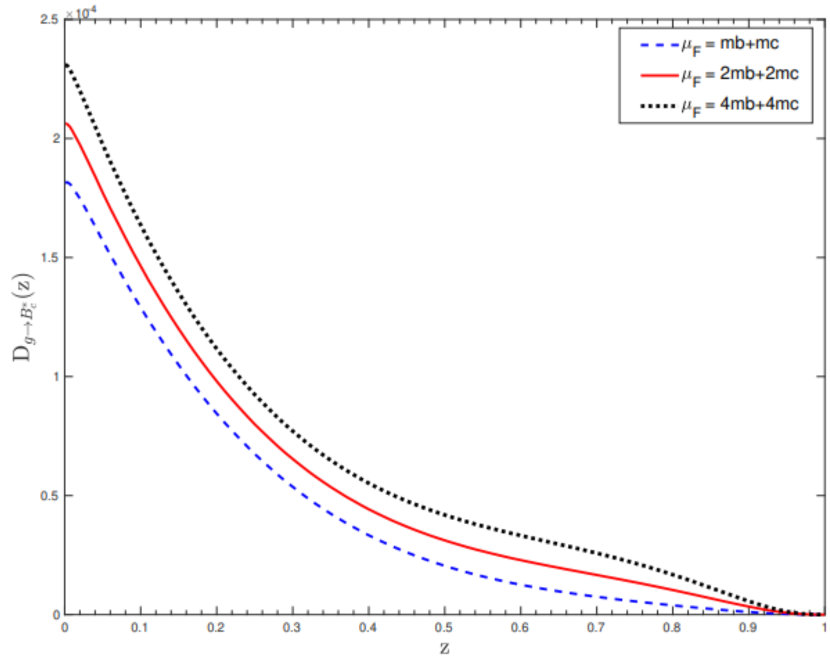
Six of the 49 cut diagrams

g->Bc(Bc*) FFs

JHEP 05, 036, (2022),
X.-C. Zheng, C.-H. Chang, X.-G. Wu.



g->Bc FFs



g->Bc* FFs

FFs at NLO

PHYSICAL REVIEW D **106,** 054030 (2022)

Gluon fragmentation into $B_c^{(*)}$ in NRQCD factorization

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[1]*China University of Mining and Technology, Beijing 100083, China*
[2]*Institute of High Energy Physics, Chinese Academy of Sciences, Beijing 100049, China*
[3]*School of Physical Sciences, University of Chinese Academy of Sciences, Beijing 100049, China*

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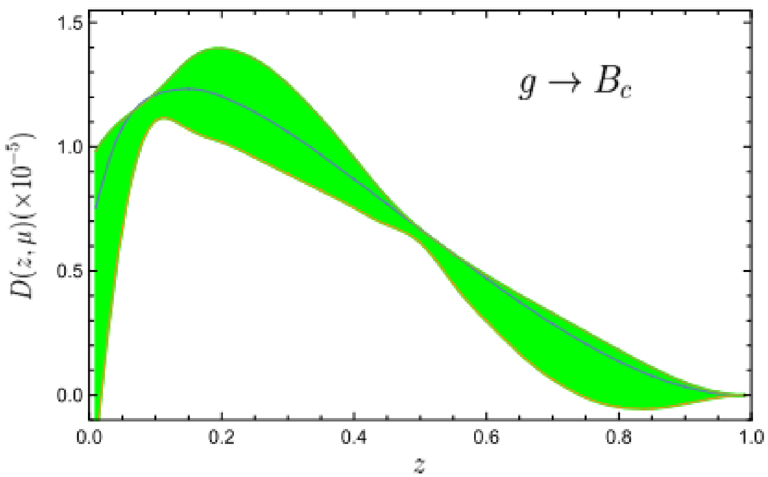
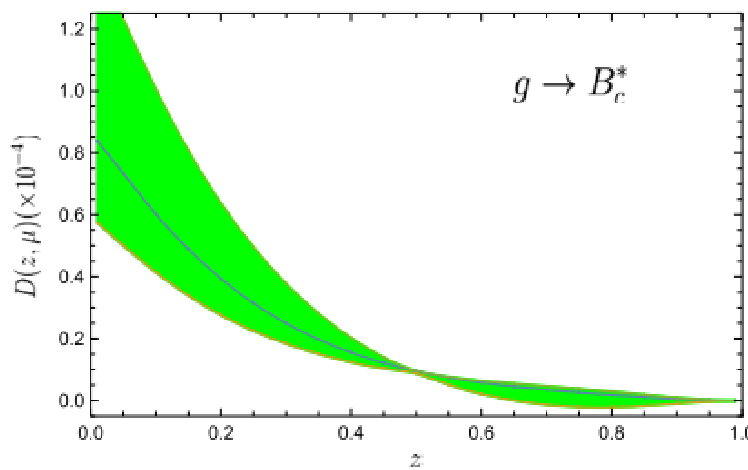
The universal fragmentation functions of gluon into the flavored quarkonia B_c and (polarized) B_c^* are computed within NRQCD factorization framework at the lowest order in velocity expansion and strong coupling constant. It is mandatory to invoke the Dokshitzer-Gribov-Lipatov-Altarelli-Parisi renormalization program to render the NRQCD short-distance coefficients UV finite in a pointwise manner. The calculation is facilitated with the sector decomposition method, with the final results presented with high numerical accuracy. This knowledge is useful to enrich our understanding toward the large-p_T behavior of $B_c^{(*)}$ production at LHC experiment.

DOI: 10.1103/PhysRevD.106.054030

$c_1^{Bc}(z)$	z=0.1	Z=0.2	Z=0.5	Z=0.8	Z=0.9
Zheng et al	0.2323	0.2311	0.1282	0.02612	0.006144
Feng et al	0.2324	0.2311	0.1282	0.02612	0.006143

$c_1^{Bc^*}(z)$	z=0.1	Z=0.2	Z=0.5	Z=0.8	Z=0.9
Zheng et al	1.155	0.7554	0.1822	0.03412	0.009589
Feng et al	1.155	0.7550	0.1822	0.03411	0.009586





➢ Quarkonium FFs at NLO

E. Braaten et al, NPB 586, 427, (2000) , $g \to Q\bar{Q}(\,^3S_1^{[8]})$

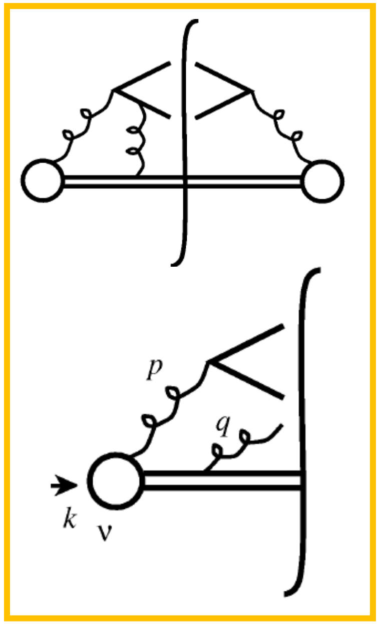
E. Braaten et al, JHEP 04, 121, (2015) , $g \to Q\bar{Q}(\,^1S_0^{[1]})$

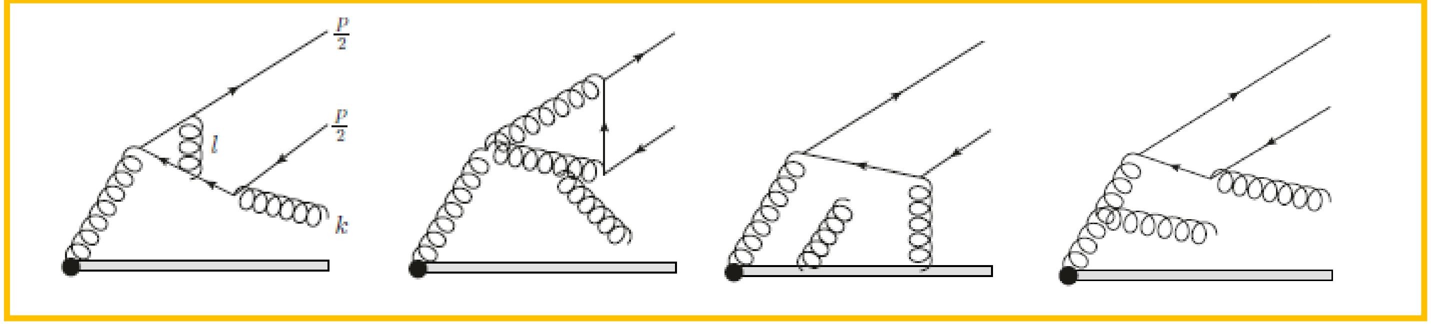
E. Braaten et al, JHEP 01, 227, (2019) , $g \to Q\bar{Q}(\,^1S_0^{[8]})$

Y. Jia et al, arXiv:1810.04138(2018) , $g \to Q\bar{Q}(\,^1S_0^{[1,8]})$

Y. –Q. Ma et al, JHEP 04, 116, (2019) , $g \to Q\bar{Q}(\,^1S_0^{[1,8]})$

Y. –Q. Ma et al, JHEP 08, 111, (2021) , $g \to Q\bar{Q}(\,^3P_J^{[1,8]})$



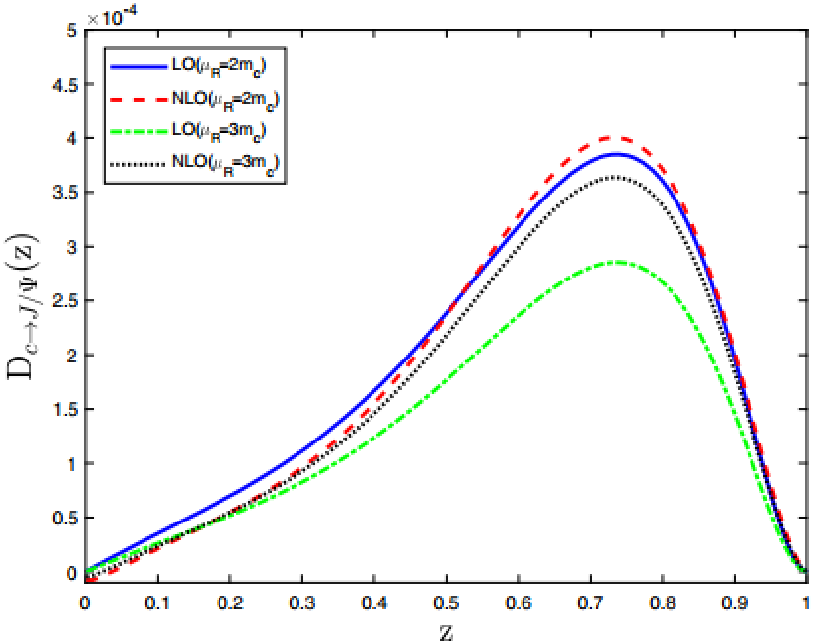


Several typical diagrams

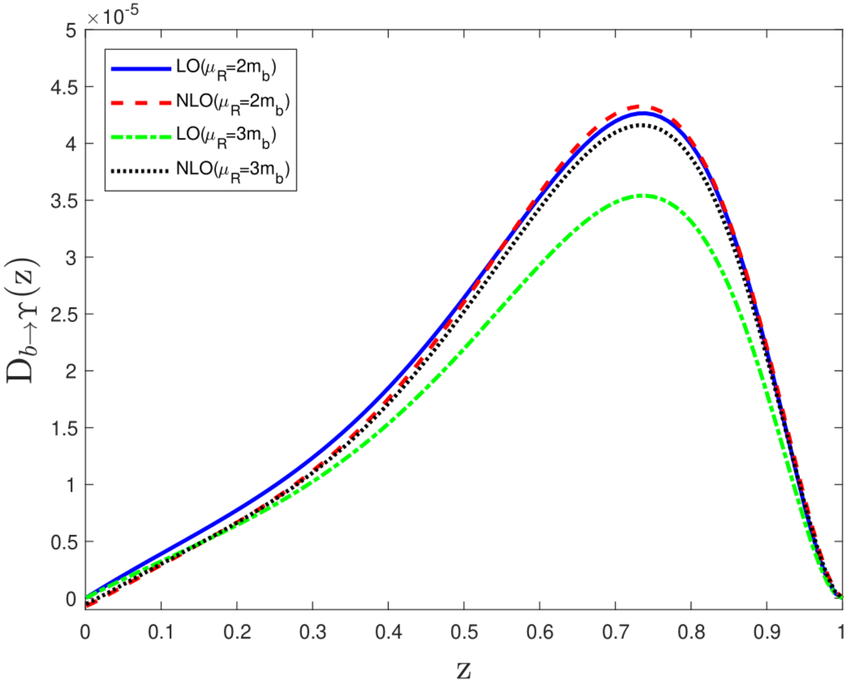
Quarkonium FFs

Phys. Rev. D 100, 014005, (2019),
X.-C. Zheng, C.-H. Chang, X.-G. Wu.

NLO fragmentation functions for $c \to J/\psi$ and $b \to \Upsilon$



$$D_{c\to J/\psi}(z, \mu_F = 3m_c)$$

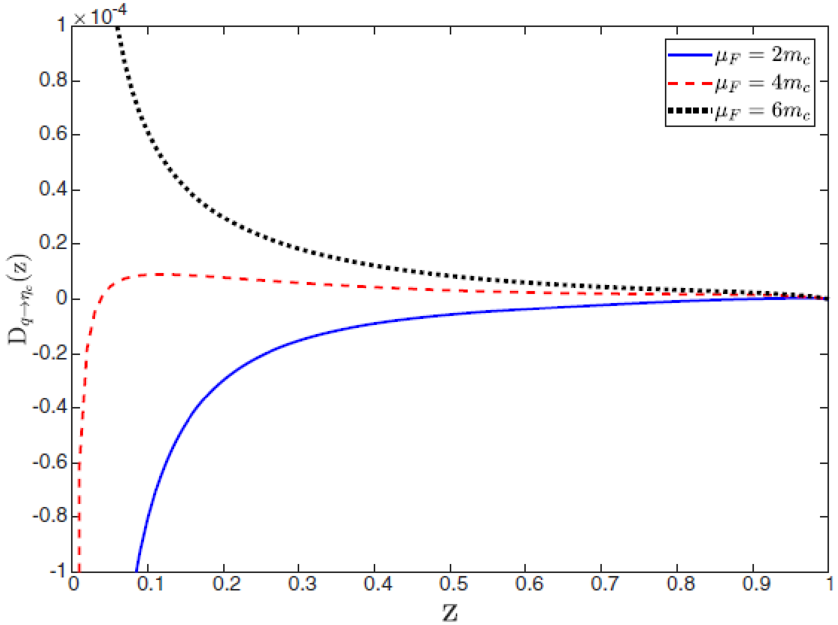


$$D_{b\to \Upsilon}(z, \mu_F = 3m_b)$$

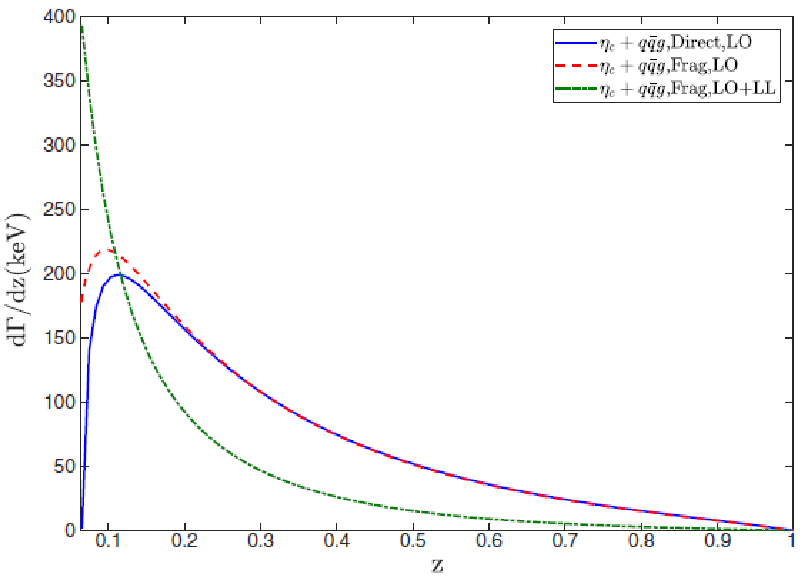
Quarkonium FFs

Phys. Rev. D 103, 074004, (2021),
X.-C. Zheng, Z.-Y. Zhang, X.-G. Wu.

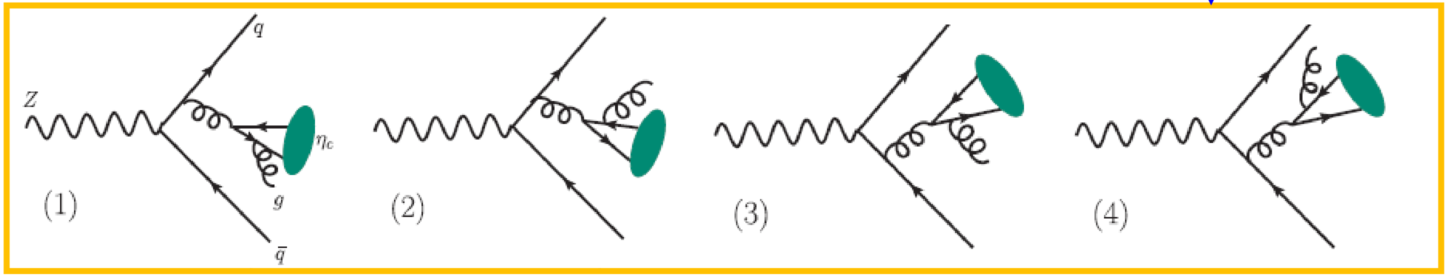
Fragmentation functions for $q \to \eta_Q (q \neq Q)$



$D_{q\to\eta_c}(z, \mu_F)$



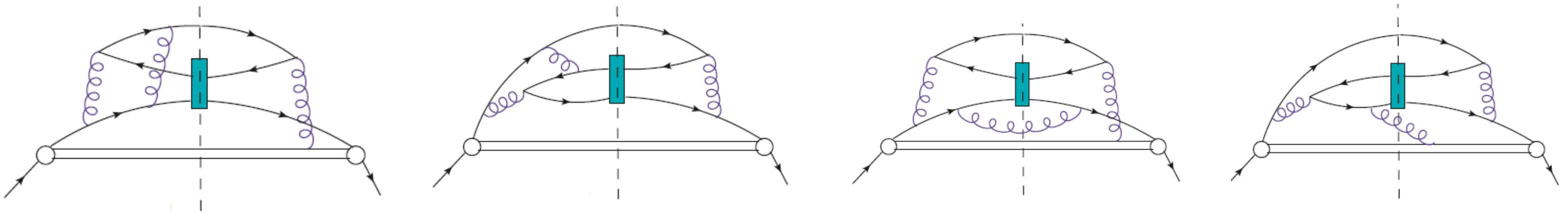
$d\Gamma/dz$ for $Z \to \eta_c + q\bar{q}g$

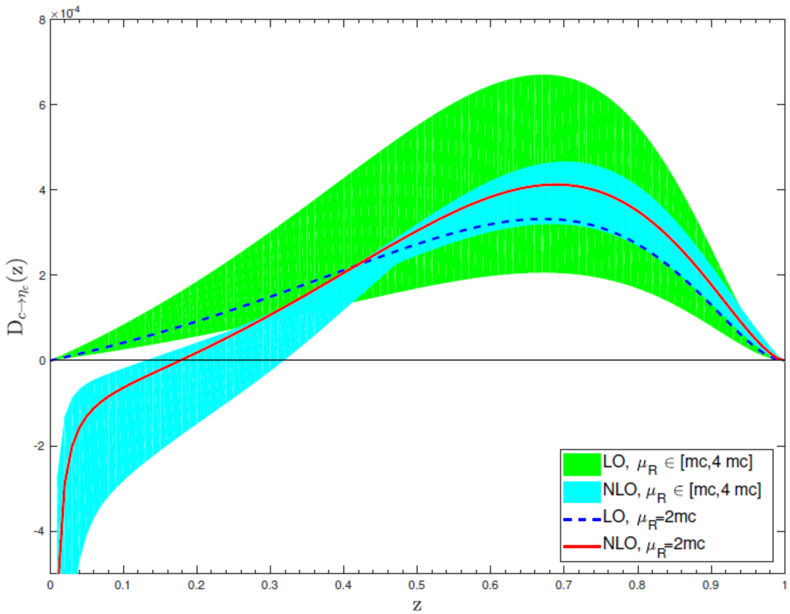


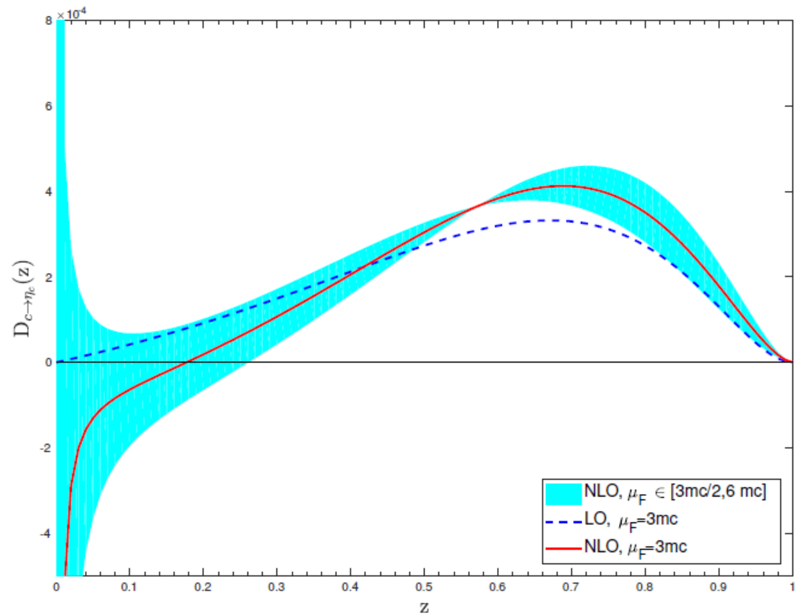
Quarkonium FFs

JHEP 07, 014, (2021),
X.-C. Zheng, X.-G. Wu, X.-D. Huang.

NLO fragmentation functions for $Q \to \eta_Q$







$D_{c\to\eta_c}(z,\mu_F)$ $D_{c\to\eta_c}(z,\mu_F)$

Quarkonium FFs

Eur. Phys. J. C (2021) 81:597
https://doi.org/10.1140/epjc/s10052-021-09390-4

THE EUROPEAN PHYSICAL JOURNAL C

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Regular Article - Theoretical Physics

Next-to-leading-order QCD corrections to heavy quark fragmentation into ${}^1S_0^{(1,8)}$ quarkonia

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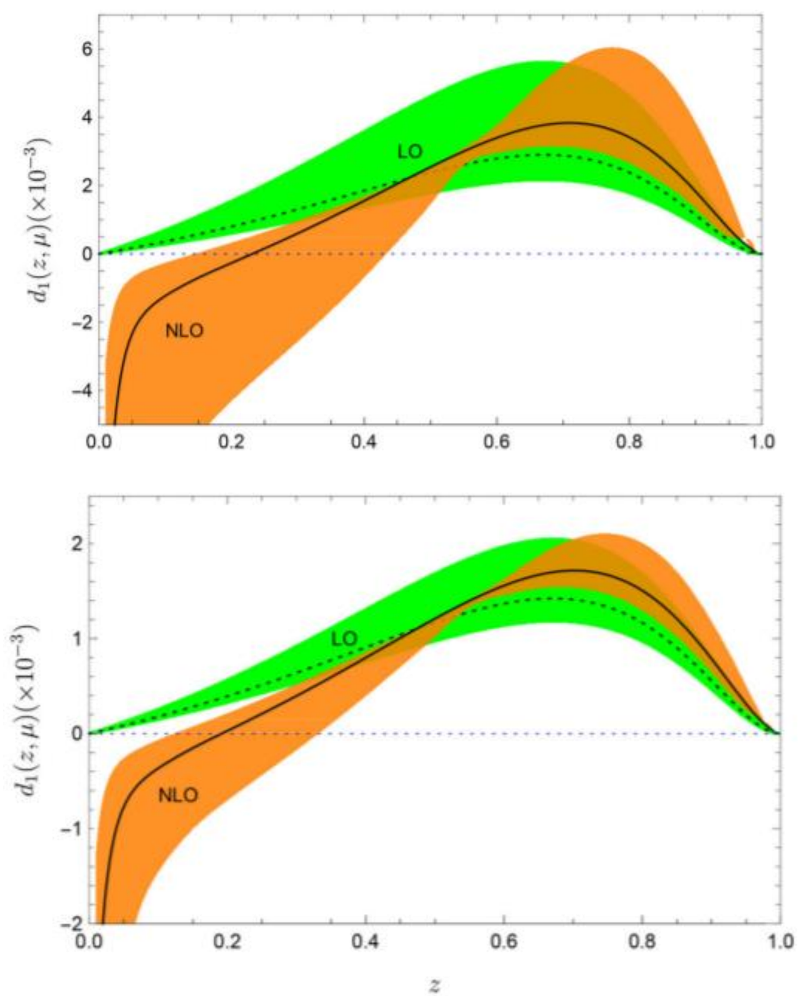
Abstract Within NRQCD factorization framework, in this work we compute, at the lowest order in velocity expansion, the next-to-leading-order (NLO) perturbative corrections to the short-distance coefficients associated with heavy quark fragmentation into the ${}^1S_0^{(1,8)}$ components of a heavy quarkonium. Starting from the Collins and Soper's operator definition of the quark fragmentation function, we apply the sector decomposition method to facilitate the numerical manipulation. It is found that the NLO QCD corrections have a significant impact.

hadronization mechanism. Similar to parton distribution functions, the scale dependence of FFs is governed by the celebrated Dokshitzer–Gribov–Lipatov–Altarelli–Parisi (DGLAP) equation:

$$\frac{d}{d\ln\mu^2}D_{i\to H}(z,\mu) = \sum_i \int_z^1 \frac{dy}{y} P_{ji}(y,\alpha_s(\mu)) D_{j\to H}\left(\frac{z}{y},\mu\right), \tag{2}$$

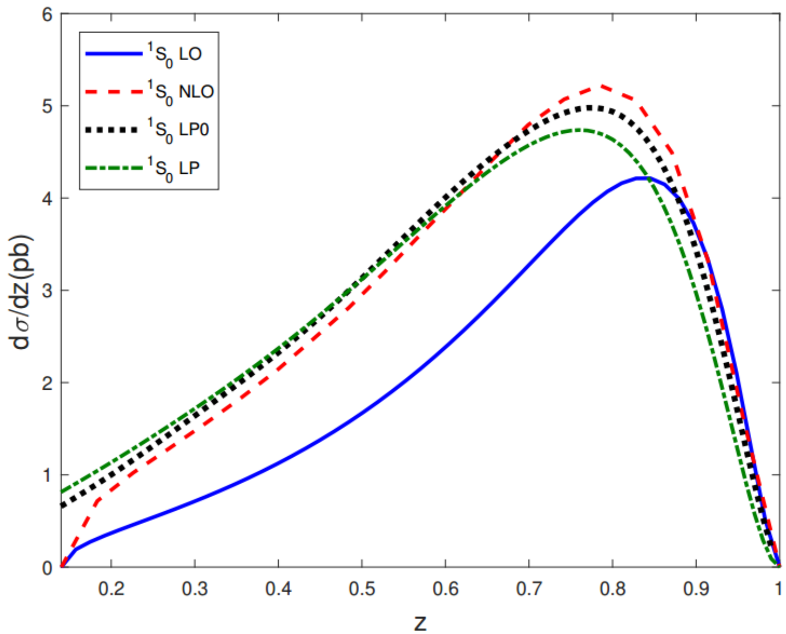
Note added After this work was completed and while we were preparing the manuscript, very recently a preprint [52] has appeared, which also computes the NLO perturbative corrections to the heavy quark fragmentation into a ${}^1S_0^{(1)}$ quarkonium. Their numerical results appear to be compatible with ours in this color-singlet channel.



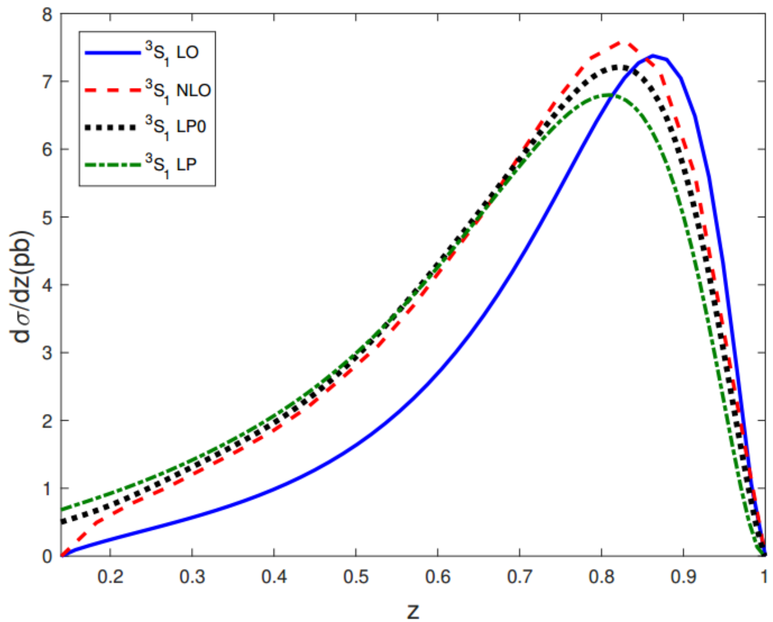
3. Applications

Phys. Rev. D 100, 034004, (2019),
X.-C. Zheng, C.-H. Chang, X.-G. Wu.

- Bc production at the Z pole



$d\sigma/dz(Bc)$



$d\sigma/dz(Bc^*)$

LO,NLO: direct NRQCD approach

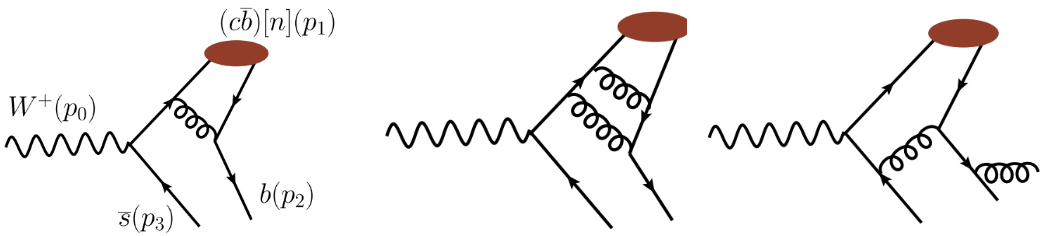
LP0： fragmentation approach, no DGLAP evolution.

LP： fragmentation approach, evolved with DGLAP equation.

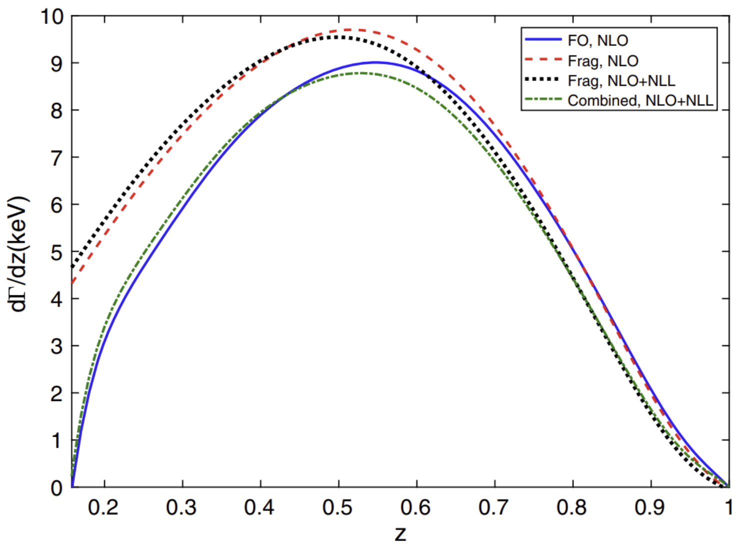
Applications

- ➢ Bc production via W⁺-boson decay

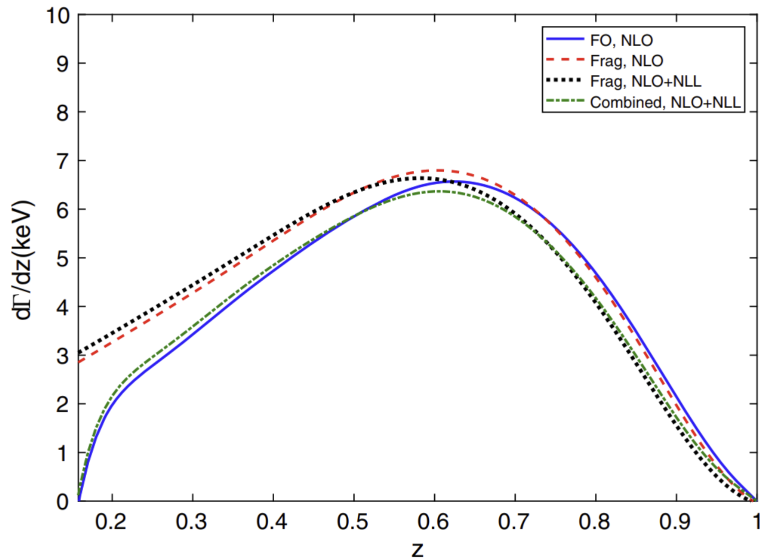
Phys. Rev. D 101, 034029, (2020),
X.-C. Zheng, C.-H. Chang, X.-G. Wu, et al.



Typical Feynman diagrams under the Fixed-order approach



$d\Gamma / dz(Bc)$



$d\Gamma / dz(Bc^*)$

- Bc production via Higgs-boson decay

Phys. Rev. D 107, 074005, (2023), X.-C. Zheng, X.-G. Wu et al.

Two sources of large logarithms:

- Renormalization of the Yukawa couplings;
- Collinear emission of gluons and quarks.

Running quark masses

DGLAP evolution

Contributions	Direct NRQCD	FF approach
$\bar{b}$-fragmentation	1.20	1.22
c-fragmentation	4.13×10^{-3}	4.26×10^{-3}
Interference	1.25×10^{-2}	
Total	1.22	1.22

Decay width under the fixed-order calculation

Contributions	B_c
$\bar{b}$-fragmentation	0.673
c-fragmentation	1.47×10^{-3}
g-fragmentation	-1.80×10^{-3}
Triangle top-loop	4.59×10^{-2}
Total	0.719

Decay width after resuming the large log terms

Summary

- Fragmentation-function approach can be used to **resum the large logarithms** arising from the collinear emissions;
- The **NLO fragmentation functions** for a quark or gluon into a doubly heavy meson **(Bc, J/ψ, Υ, η_c, η_b)** have been obtained;
- These **fragmentation functions** can be studied at the high energy colliders, such as CEPC, FCC-ee, etc.

Thank you！