



基础物理与数学科学学院
School of Fundamental Physics and Mathematical Sciences



國科大杭州高華研究院
Hangzhou Institute for Advanced Study, UCAS

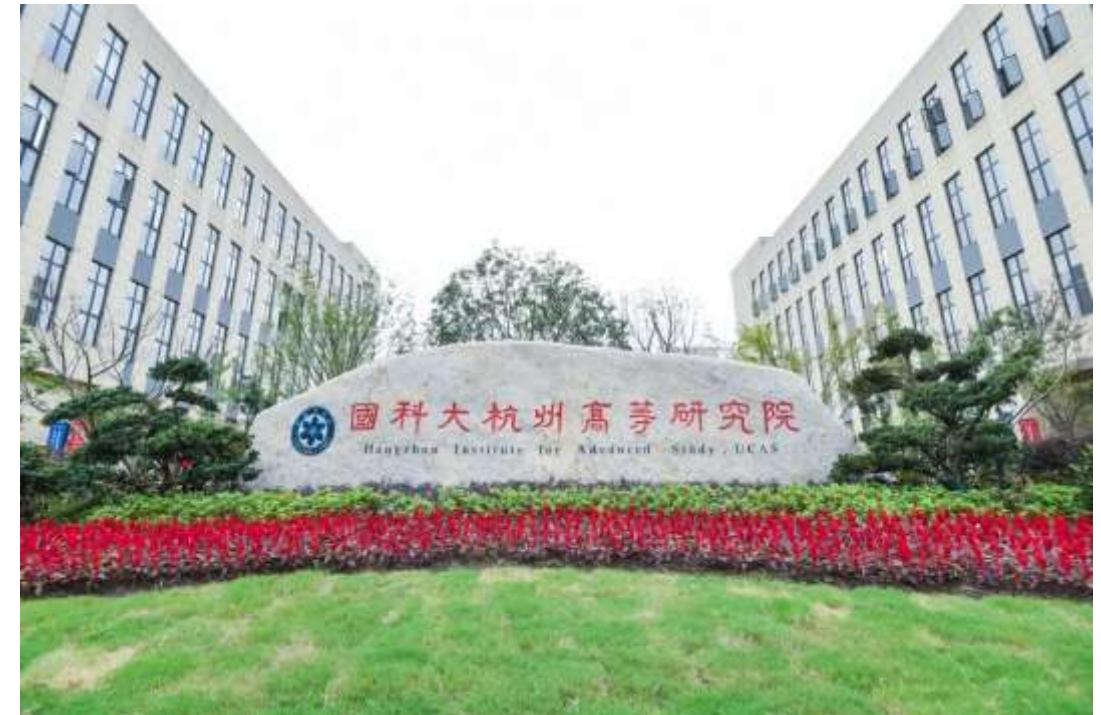
A phenomenological study of SU(5) with Type-I+III seesaw

Gao-Xiang Fang (HIAS)
2024-04-10

Based on GXF, Y.L. Zhou, arXiv:2404.xxxxxx

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GUT theory

SM: $G_{\text{SM}} = SU(3)_C \times SU(2)_L \times U(1)_Y$, neutrinos are massless

↓
Unify three fundamental forces into a single force

GUTs: $SU(5)$ GUTs, $SO(10)$ GUTs

- Unification of symmetries

$$G_{\text{GUT}} \supset G_{\text{SM}} = SU(3)_C \times SU(2)_L \times U(1)_Y$$

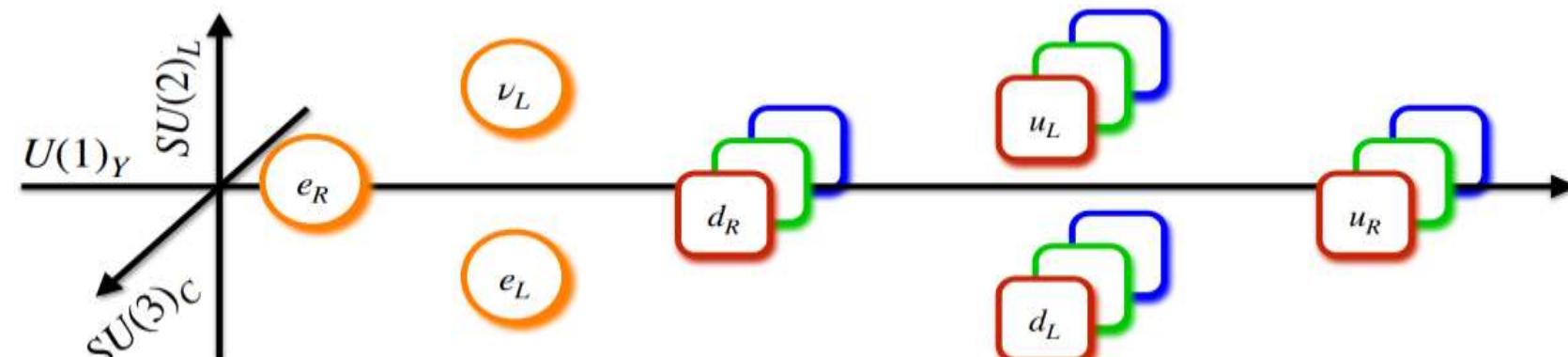
- Unification of gauge couplings

$$\downarrow \quad = \quad \downarrow \quad = \quad \downarrow$$
$$g_3 = g_2 = g_1$$

Three gauge couplings are unified at scale called M_{GUT}

- Unification of matters

Zhou' SUSY2023



Weak hypercharge: $Y=1$ $Y=-\frac{1}{2}$

$Y=\frac{1}{3}$

$Y=\frac{1}{6}$

$Y=-\frac{2}{3}$

$SU(5)$ GUTs

Georgi-Glashow model(1974) $SU(5)$

- Gauge couplings do not unify
- Neutrinos are massless and predicts wrong mass relations $m_d = m_e, m_s = m_\mu, m_t = m_\tau$

$$\bar{\mathbf{5}}_F = \begin{pmatrix} d_{rR}^c \\ d_{gR}^c \\ d_{bR}^c \\ e_L \\ -\nu_L \end{pmatrix}, \quad \mathbf{10}_F = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & u_{bR}^c & -u_{gR}^c & -u_L^r & -d_L^r \\ -u_{bR}^c & 0 & u_{rR}^c & -u_L^g & -d_L^g \\ u_{gR}^c & -u_{rR}^c & 0 & -u_L^b & -d_L^b \\ u_L^r & u_L^g & u_L^b & 0 & -e_R^c \\ d_L^r & d_L^g & d_L^b & e_R^c & 0 \end{pmatrix}$$

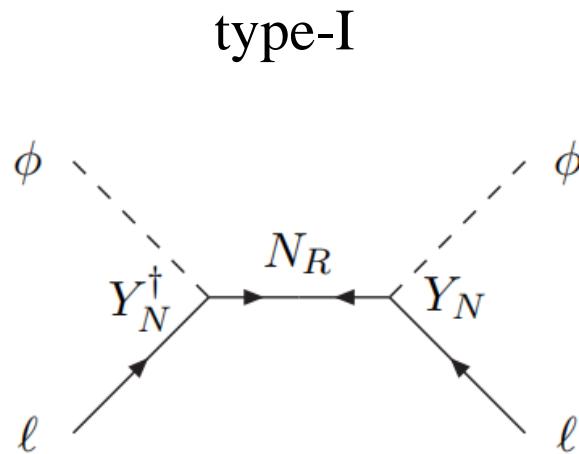
Minimally extended $SU(5)$

- Generating neutrino mass: type-I+III seesaw 24_F ; type-II seesaw 15_H
- Achieving gauge unification and predicting correct mass of quarks and leptons:
dimension=5 operators; 45_H ; vector-like fermions($5_F + \bar{5}_F, 10_F + \bar{10}_F, 15_F + \bar{15}_F$)

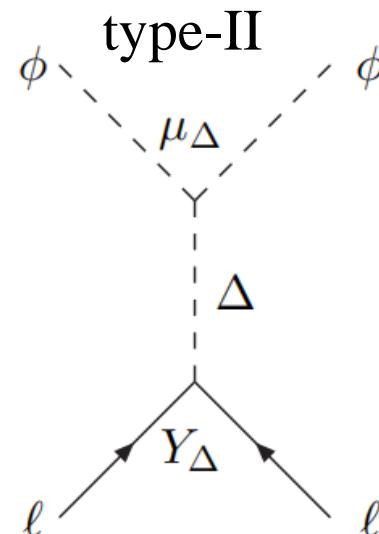
Bajc, Nemevsek, Senjanovic, 0703080
 Perez, Gross, Murgui, 1804.07831
 Calibbi, Gao, 2206.10682
 Senjanović, Zantedeschi, 2402.19224

Seesaw mechanism

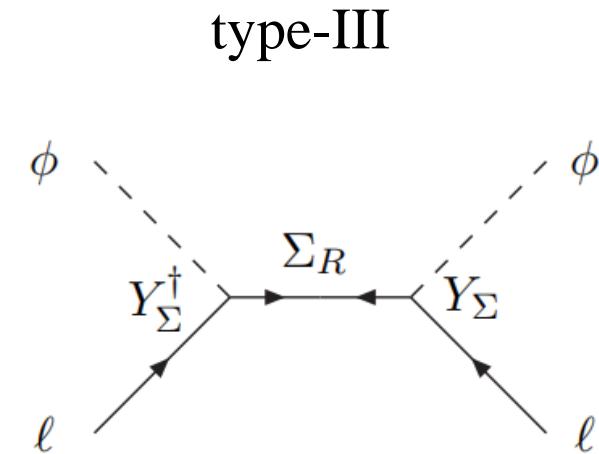
Weinberg operator: $\frac{\mathcal{L}_{d=5}}{\Lambda} = c_5 \bar{l}_L \tilde{\phi} \tilde{\phi}^T (l_L)^c + h.c.$



SM singlet fermions



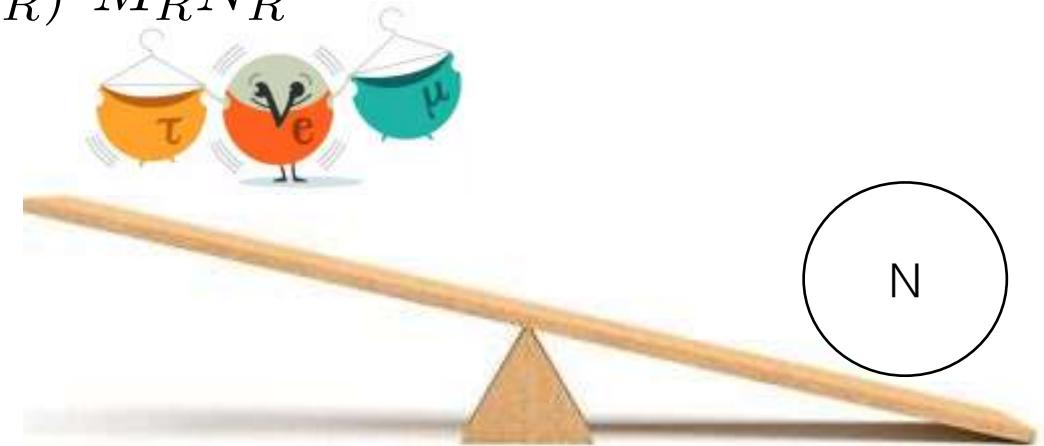
SM triplet scalars



SM triplet fermions

Majorana mass term: $-\mathcal{L} = \overline{\nu_L} M_D N_R + \frac{1}{2} \overline{(N_R)^c} M_R N_R$

Majorana neutrino: $M_\nu = -\frac{M_D M_D^T}{M_R}$



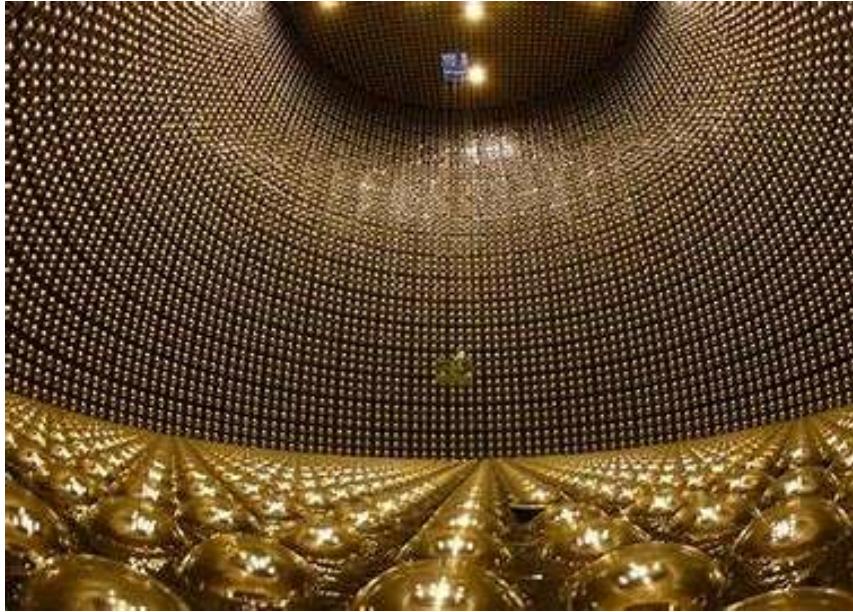
Proton decay

Theory:

Non-SUSY dimension=6 operators:

$$\frac{1}{\Lambda_1^2} \underbrace{[(\overline{u}_R^c \gamma^\mu Q)(\overline{d}_R^c \gamma_\mu L) + (\overline{u}_R^c \gamma^\mu Q)(\overline{e}_R^c \gamma_\mu Q)]}_{\text{non-SUSY SU(5)}} + \frac{1}{\Lambda_2^2} [(\overline{d}_R^c \gamma^\mu Q)(\overline{u}_R^c \gamma_\mu L) + (\overline{d}_R^c \gamma^\mu Q)(\overline{\nu}_R^c \gamma_\mu Q)]$$

Experiments:



Current Super-K:

$$\begin{aligned}\tau(p \rightarrow \pi^0 e^+) &> 2.4 \times 10^{34} \text{ yr}, \\ \tau(p \rightarrow K^+ \bar{\nu}) &> 6.6 \times 10^{33} \text{ yr}\end{aligned}$$

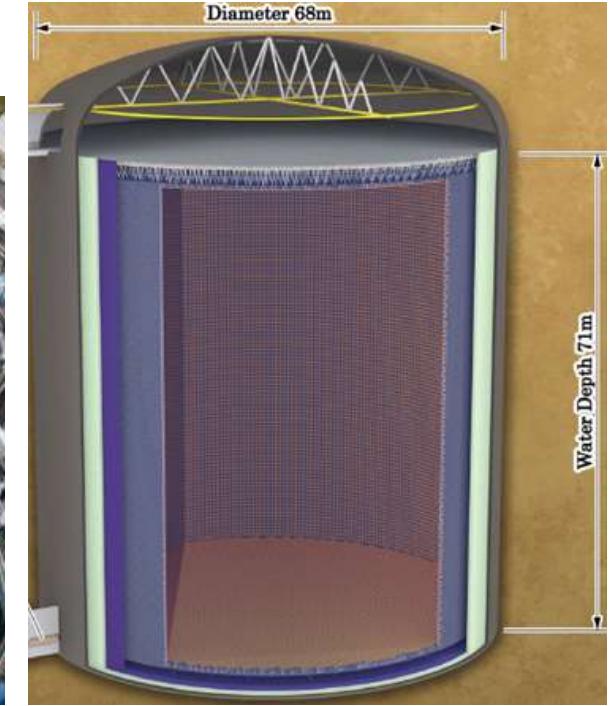
SK, 2010.16098



Future JUNO:

$$\tau(p \rightarrow K^+ \bar{\nu}) > 9.6 \times 10^{33} \text{ yr}$$

JUNO, 2212.08502



Future HK:

$$\tau(p \rightarrow K^+ \bar{\nu}) > 3.2 \times 10^{34} \text{ yr}$$

HK, 1805.04163 6

Field contents

An economical model:

Fields	$SU(5)$	$\supset SU(3)_c \times SU(2)_L \times U(1)_Y$
Fermion	$\mathbf{10}_F$	$\rightarrow (\mathbf{3}, \mathbf{2}, +\frac{1}{6})_{q_L} + (\overline{\mathbf{3}}, \mathbf{1}, -\frac{2}{3})_{u_R^c} + (\mathbf{1}, \mathbf{1}, +1)_{e_R^c}$
	$\overline{\mathbf{5}}_F$	$\rightarrow (\overline{\mathbf{3}}, \mathbf{1}, +\frac{1}{3})_{d_R^c} + (\mathbf{1}, \mathbf{2}, -\frac{1}{2})_{\ell_L}$
	$\mathbf{24}_F$	$\rightarrow (\mathbf{1}, \mathbf{1}, 0)_N + (\mathbf{1}, \mathbf{3}, 0)_\Sigma + (\mathbf{3}, \mathbf{2}, -\frac{5}{6})_{Q_L} + (\overline{\mathbf{3}}, \mathbf{2}, \frac{5}{6})_{Q_R^c} + (\mathbf{8}, \mathbf{1}, 0)_{Q_8}$
Higgs	$\mathbf{5}_H$	$\rightarrow (\mathbf{1}, \mathbf{2}, \frac{1}{2})_{h_1}$
	$\mathbf{45}_H$	$\rightarrow (\mathbf{1}, \mathbf{2}, \frac{1}{2})_{h_2}$
	$\mathbf{24}_\Phi$	$\rightarrow (\mathbf{1}, \mathbf{1}, 0)_{\phi_1}$
	$\mathbf{75}_\Phi$	$\rightarrow (\mathbf{1}, \mathbf{1}, 0)_{\phi_2}$

$$\begin{aligned}
 -\mathcal{L} \supset & \overline{\mathbf{5}}_F(Y_1 \mathbf{5}_H^\dagger + Y_2 \mathbf{45}_H^\dagger) \mathbf{10}_F + \mathbf{10}_F(Y_3 \mathbf{5}_H + Y_4 \mathbf{45}_H) \mathbf{10}_F \\
 & + \overline{\mathbf{5}}_F(Y_5 \mathbf{5}_H + Y_6 \mathbf{45}_H) \mathbf{24}_F + \mathbf{24}_F(M_1 + \kappa_1 \mathbf{24}_\Phi + \kappa_2 \mathbf{75}_\Phi) \mathbf{24}_F + \text{h.c.}
 \end{aligned}$$

$$\begin{aligned}
 M_e &= aY_1^* - 3bY_2^*, \\
 M_N &= M_1 - M_{24} + 5M_{75}, \quad M_d = aY_1^\dagger + bY_2^\dagger, \\
 M_\Sigma &= M_1 - 3M_{24} - 3M_{75}, \quad M_u = cY_3^* + dY_4^*, \quad \xrightarrow{\text{type I+III seesaw}} M_\nu = M_I M_N^{-1} M_I^T + M_{III} M_\Sigma^{-1} M_{III}^T \\
 M_Q &= M_1 - \frac{1}{2}M_{24} + M_{75}, \quad M_I = \sqrt{3}fY_5^* + \sqrt{5}gY_6^*, \\
 M_{Q_8} &= M_1 + 2M_{24} - M_{75}. \quad M_{III} = \sqrt{5}fY_5^* - \sqrt{3}gY_6^*, \\
 \text{New particles} & \qquad \qquad \qquad \text{SM particles}
 \end{aligned}$$

Renormalization group equation (RGE)

RGE at one-loop level: $\frac{d\alpha_i(t)}{dt} = \frac{a_i}{2\pi} \alpha_i^2$

↓
analytical solution

$$\alpha_i^{-1}(t_{\text{GUT}}) = \alpha_i^{-1}(0) - \frac{a_i^{\text{SM}}}{2\pi} t_{\text{GUT}} - \sum_I \frac{a_i^I}{2\pi} (t_{\text{GUT}} - t_I)$$

Fields $SU(5) \supset SU(3)_c \times SU(2)_L \times U(1)_Y$

Fermion	$\mathbf{10}_F$	$= (\mathbf{3}, \mathbf{2}, +\frac{1}{6})_{q_L} + (\bar{\mathbf{3}}, \mathbf{1}, -\frac{2}{3})_{u_R^c} + (\mathbf{1}, \mathbf{1}, +1)_{e_R^c}$
	$\bar{\mathbf{5}}_F$	$= (\bar{\mathbf{3}}, \mathbf{1}, +\frac{1}{3})_{d_R^c} + (\mathbf{1}, \mathbf{2}, -\frac{1}{2})_{\ell_L}$
	$\mathbf{24}_F$	$= (\mathbf{1}, \mathbf{1}, 0)_N + (\mathbf{1}, \mathbf{3}, 0)_\Sigma + (\mathbf{3}, \mathbf{2}, -\frac{5}{6})_{Q_L} + (\bar{\mathbf{3}}, \mathbf{2}, \frac{5}{6})_{Q_R^c} + (\mathbf{8}, \mathbf{1}, 0)_{Q_8}$

Beta coefficients at one-loop level:

$$\{a_i^{\text{SM}}\} = \begin{pmatrix} -7 \\ -\frac{19}{6} \\ \frac{41}{10} \end{pmatrix} \text{ SM contribution}$$

$$\{a_i^\Sigma\} = \begin{pmatrix} 0 \\ \frac{4}{3} \\ 0 \end{pmatrix} \quad \{a_i^{Q_8}\} = \begin{pmatrix} 2 \\ 0 \\ 0 \end{pmatrix}$$

$$\{a_i^Q\} = \begin{pmatrix} \frac{4}{3} \\ 2 \\ \frac{10}{3} \end{pmatrix}$$

Only considering fermion's contribution in RGE

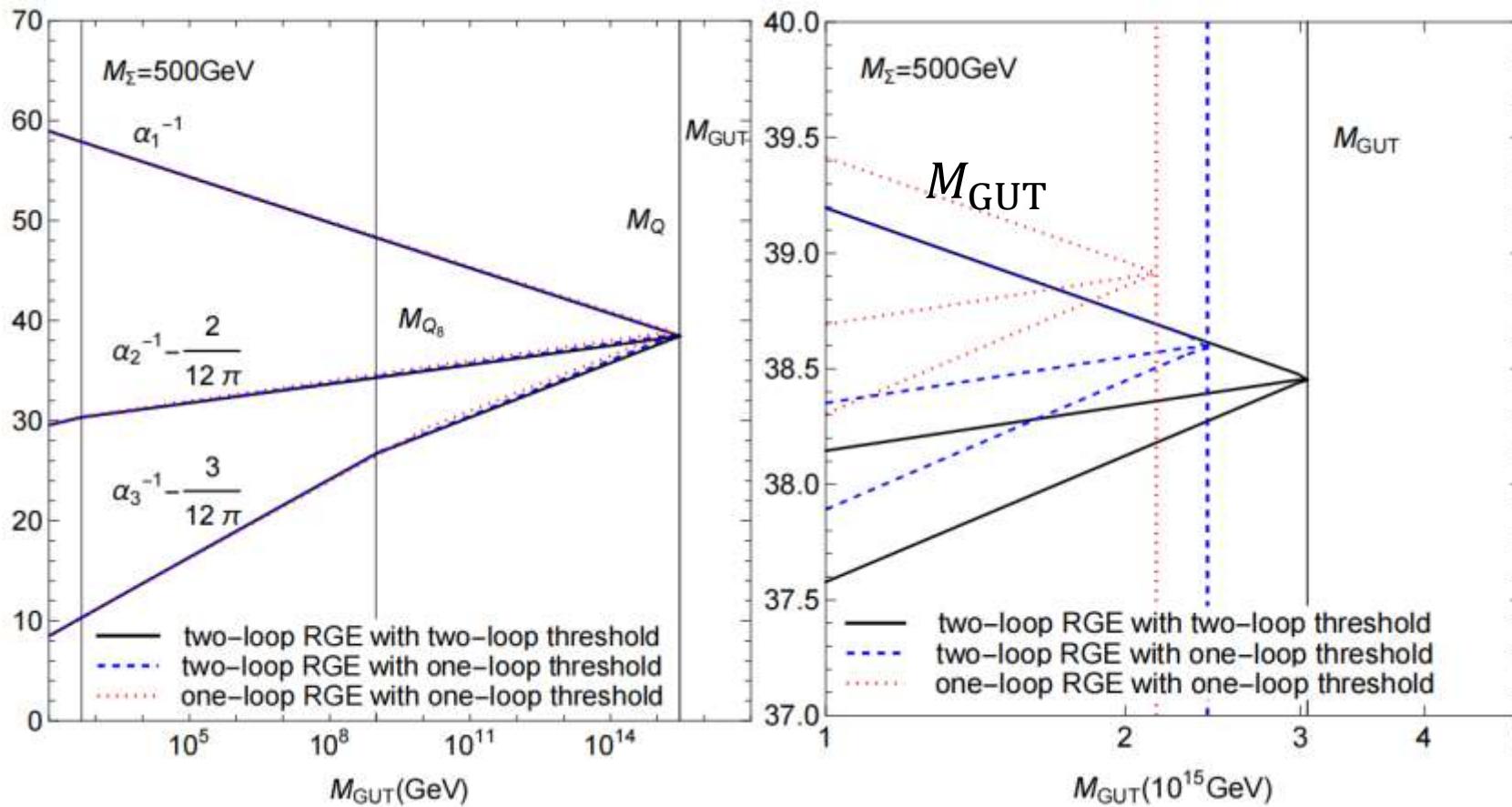
New particles contribution

Gauge coupling unification

We only focus on the situation where $\mathbf{24}_F$ has only one copy

The precise measurement of experiments requires us to calculate RGE at two-loop level

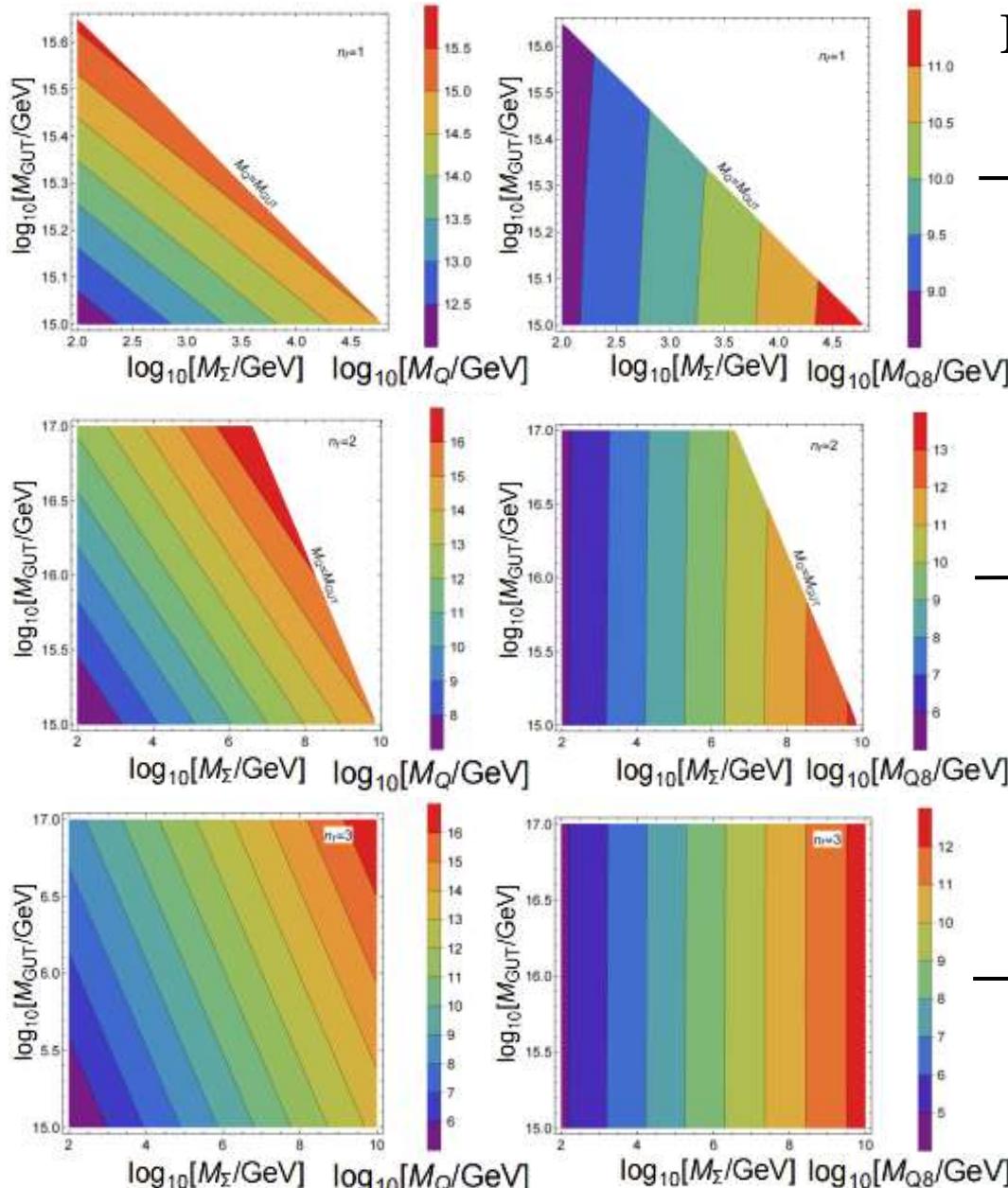
Two-loop threshold effect increase M_{GUT} $M_{\text{GUT}} = 3.05 \times 10^{15}$ GeV



one-loop matching condition:

$$\begin{aligned}\alpha_3^{-1}(t_{\text{GUT}}) - \frac{3}{12\pi} \\= \alpha_2^{-1}(t_{\text{GUT}}) - \frac{2}{12\pi} \\= \alpha_1^{-1}(t_{\text{GUT}})\end{aligned}$$

Scan Parameter Space



Mass hierarchy: $M_\Sigma < M_{Q_8} < M_Q \leq M_{\text{GUT}}$

→ $\mathbf{24}_F$ has one copy

Parameter space increases

$\mathbf{24}_F$ has one copy,
if $M_{\text{GUT}} \geq 10^{15}$ GeV,
 $M_\Sigma \in (100\text{GeV}, 63\text{TeV})$

→ $\mathbf{24}_F$ has two copy

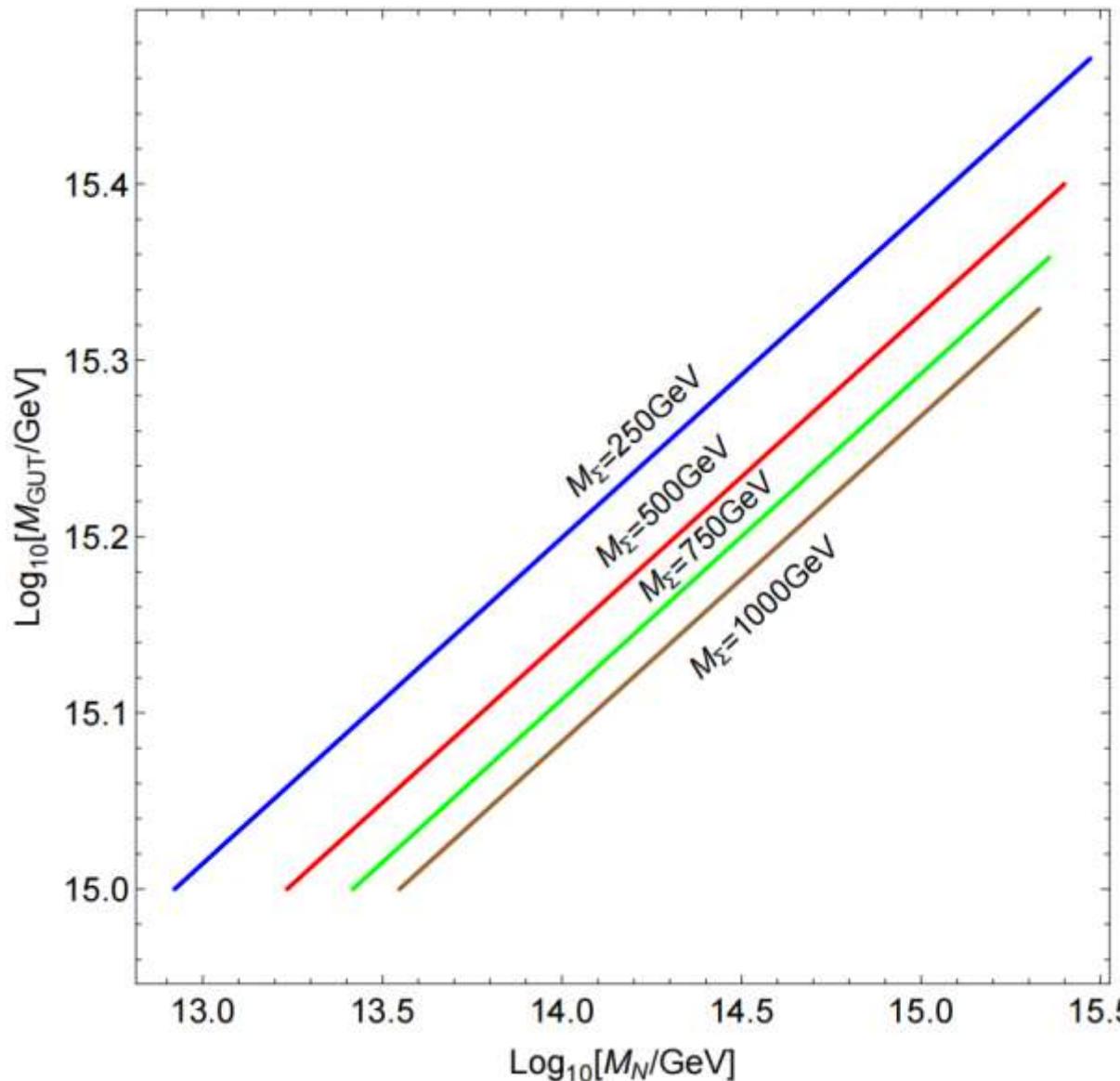
Parameter space increases

Beta coefficients at two-loop level:

$$\{b_{ij}^{\text{SM}}\} = \begin{pmatrix} -26 & \frac{9}{2} & \frac{11}{10} \\ 12 & \frac{35}{6} & \frac{9}{10} \\ \frac{44}{5} & \frac{27}{10} & \frac{199}{50} \end{pmatrix} \{b_{ij}^{Q_8}\} = \begin{pmatrix} 48 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\{b_{ij}^Q\} = \begin{pmatrix} \frac{76}{3} & 3 & \frac{5}{3} \\ 8 & \frac{49}{2} & \frac{5}{2} \\ \frac{40}{3} & \frac{15}{2} & \frac{25}{6} \end{pmatrix} \{b_{ij}^\Sigma\} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \frac{64}{3} & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Correlation between M_N and M_{GUT}



$$M_N = M_1 - M_{24} + 5M_{75},$$

$$M_\Sigma = M_1 - 3M_{24} - 3M_{75},$$

$$M_Q = M_1 - \frac{1}{2}M_{24} + M_{75},$$

$$M_{Q_8} = M_1 + 2M_{24} - M_{75}.$$

$$M_N = \frac{12}{5}M_Q - \frac{4}{5}M_{Q_8} - \frac{3}{5}M_\Sigma$$

$$M_\Sigma \uparrow M_N \downarrow M_{\text{GUT}} \downarrow$$

When $M_\Sigma = 500 \text{ GeV}$, the maximal value of M_{GUT} is $2.51 \times 10^{15} \text{ GeV}$, which is allowed by gauge unification

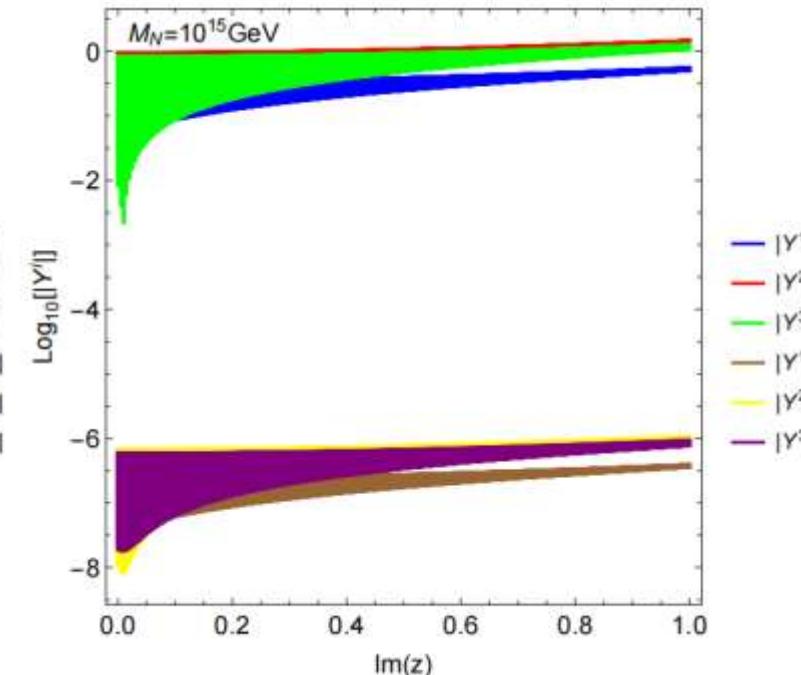
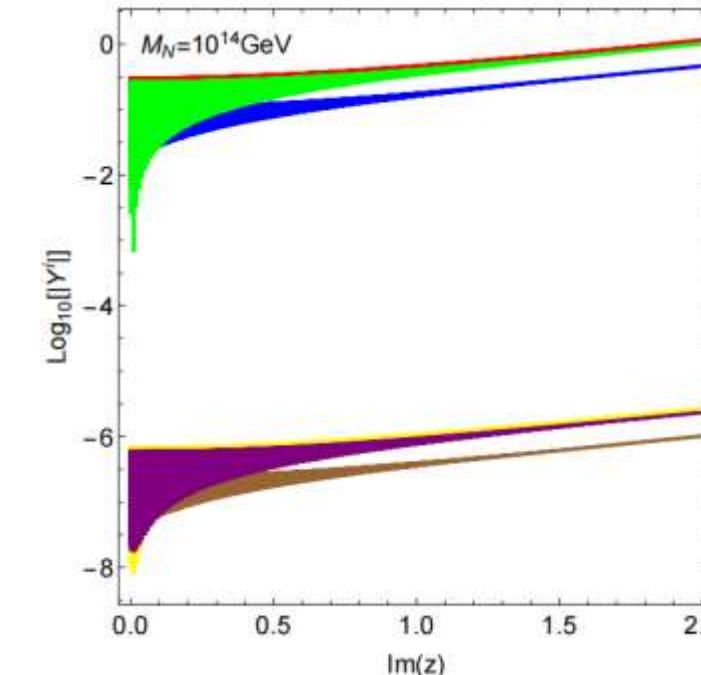
Flavor mixing

Casas-Ibarra parametrization:

$$Y_I^\alpha = \frac{\sqrt{2M_N}}{v} (\sqrt{m_2} \cos z U_{\alpha 2}^* + \sqrt{m_3} \sin z U_{\alpha 3}^*)$$

$$Y_{III}^\alpha = \frac{\sqrt{2M_\Sigma}}{v} (-\sqrt{m_2} \sin z U_{\alpha 2}^* + \sqrt{m_3} \cos z U_{\alpha 3}^*)$$

$$M_\nu = \frac{v^2}{2M_N} Y_I Y_I^T + \frac{v^2}{2M_\Sigma} Y_{III} Y_{III}^T$$



$\text{Im}(z)$ actually affects $|Y|$

Casas, Ibarra, 0103065

The smallness of three left-handed neutrinos is attributed to the smallness of $|Y_{III}|$ and the largeness of M_N

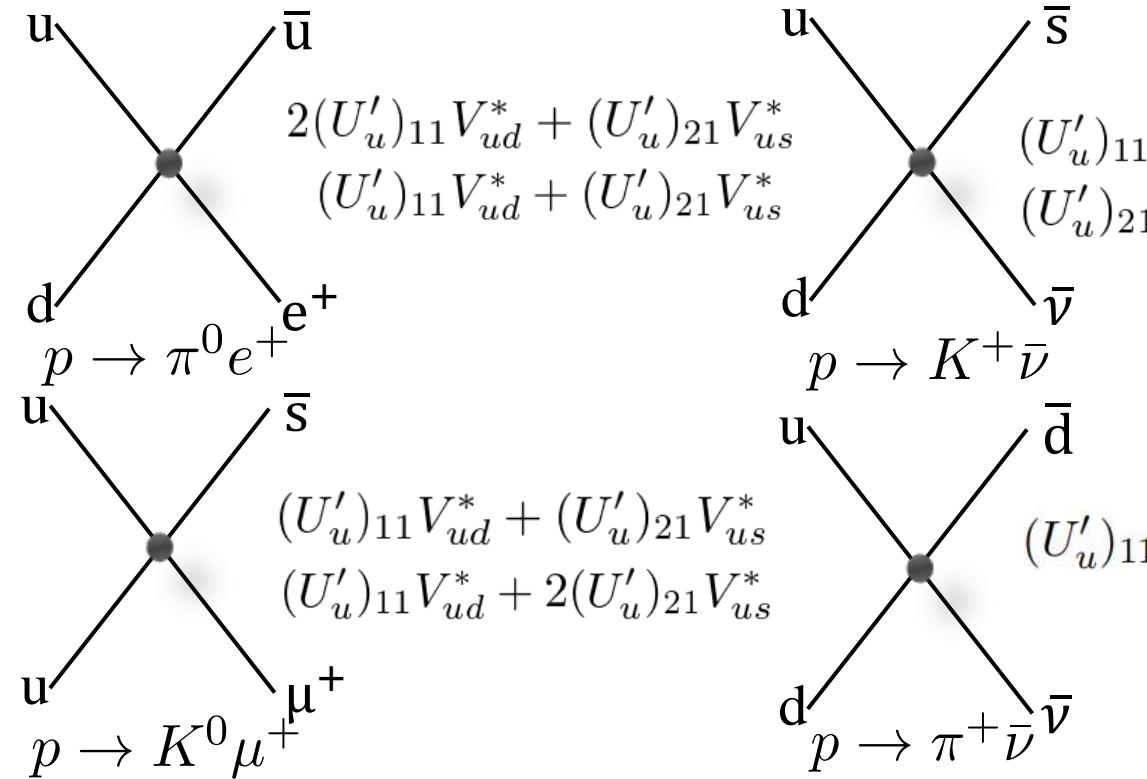
$$M_N \gg M_\Sigma$$

$$|Y_I| \gg |Y_{III}|$$

Proton decay

$$\frac{|V_{ub}| < |V_{td}| \ll |V_{ts}| < |V_{cb}| \ll |V_{cd}| < |V_{us}| \ll |V_{cs}| < |V_{ud}| < |V_{tb}|}{\lambda^3 \quad \lambda^2 \quad \lambda \quad \mathcal{O}(1)}$$

Scenario 1: $Y_d = \hat{Y}_d$



Unitarity: $| (U'_u)_{11} |^2 + | (U'_u)_{21} |^2 \leq 1$

numerical solution (V_{ud}, V_{us} are taken from PDG)

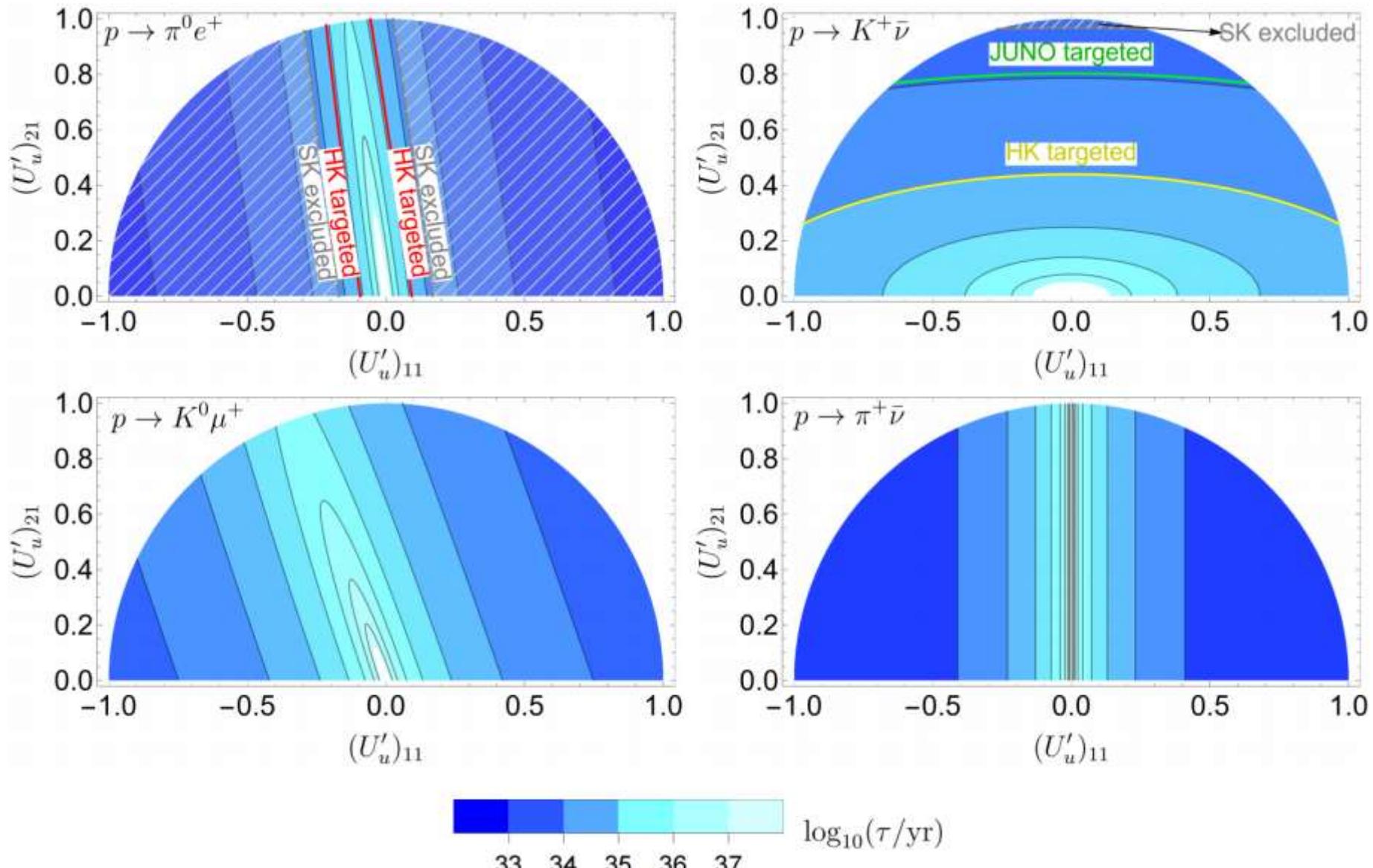
$$\begin{aligned}\tau(p \rightarrow \pi^0 e^+) &= \log_{10} \left(\frac{6.88 \times 10^{32}}{(U'_u)_{11}^2 + 0.27(U'_u)_{11}(U'_u)_{21} + 0.02(U'_u)_{21}^2} \right) \\ \tau(p \rightarrow K^+ \bar{\nu}) &= \log_{10} \left(\frac{4.62 \times 10^{34}}{(U'_u)_{11}^2 + 7.46(U'_u)_{21}^2} \right) \\ \tau(p \rightarrow K^0 \mu^+) &= \log_{10} \left(\frac{5.59 \times 10^{33}}{(U'_u)_{11}^2 + 0.69(U'_u)_{11}(U'_u)_{21} + 0.13(U'_u)_{21}^2} \right) \\ \tau(p \rightarrow \pi^+ \bar{\nu}) &= \log_{10} \left(\frac{1.68 \times 10^{33}}{(U'_u)_{11}^2} \right)\end{aligned}$$

Only two free parameters $(U'_u)_{11}, (U'_u)_{21}$

$$(U'_u)_{11} \rightarrow 0, (U'_u)_{21} \rightarrow 0, \tau \rightarrow +\infty$$

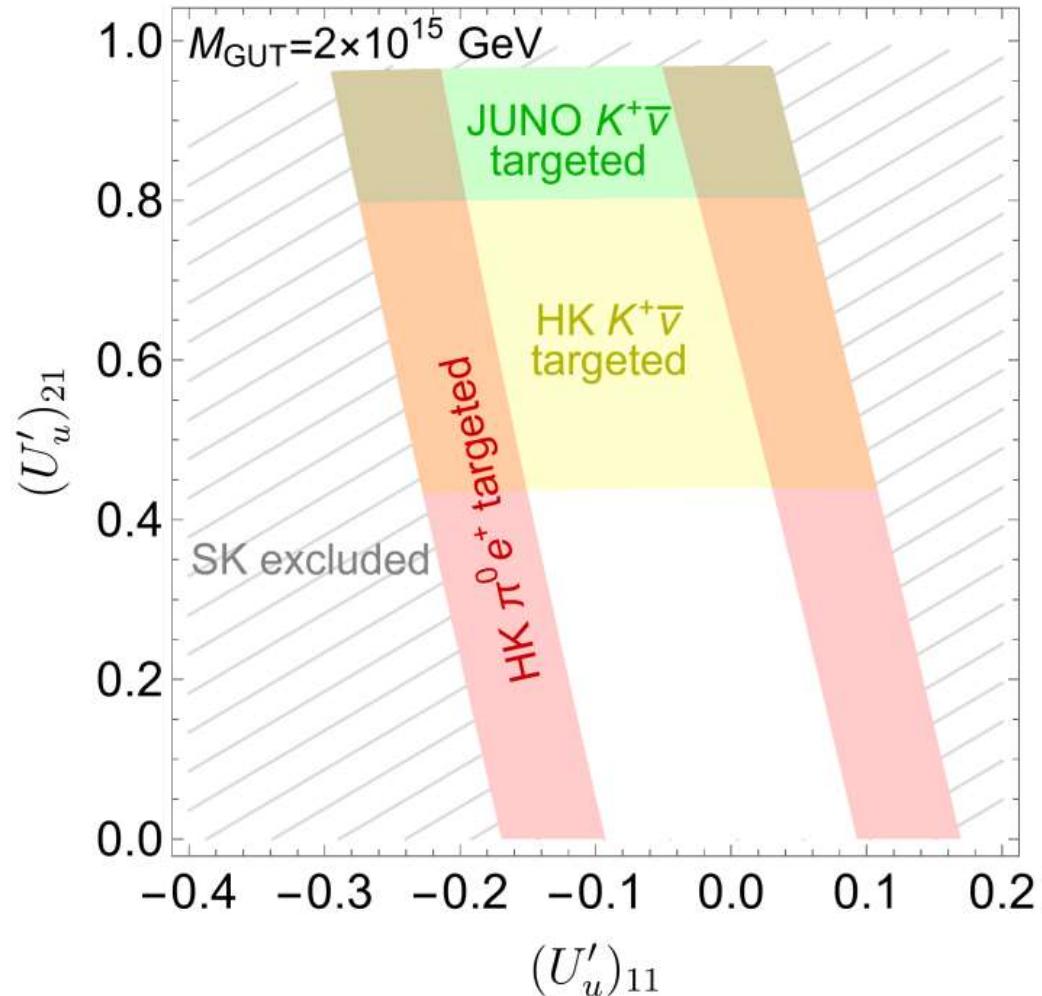
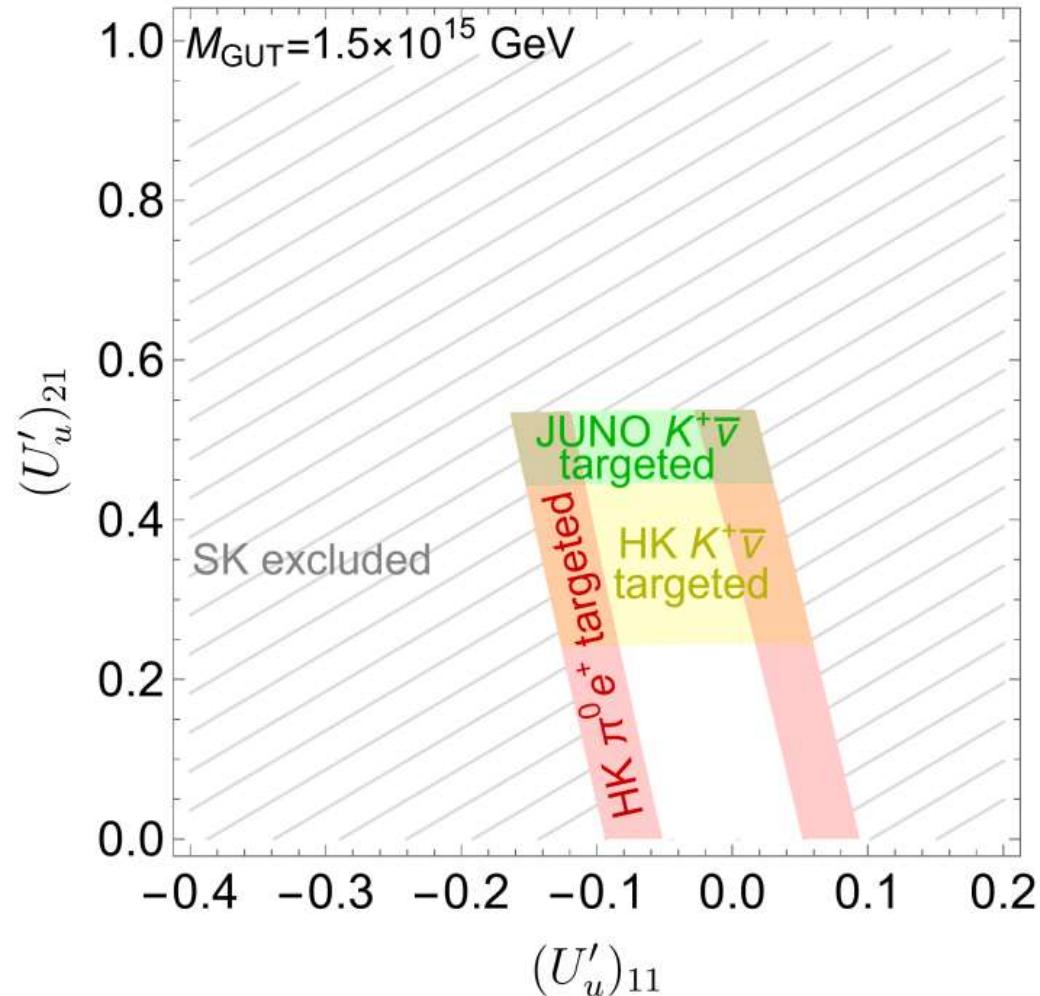
Proton decay

$M_\Sigma = 500 \text{ GeV}$, $M_{\text{GUT}} = 2 \times 10^{15} \text{ GeV}$ are determined



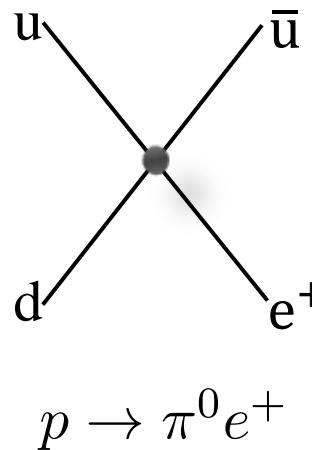
Proton decay

$M_{\text{GUT}} \uparrow$ Proton lifetime \uparrow Parameter Space \uparrow



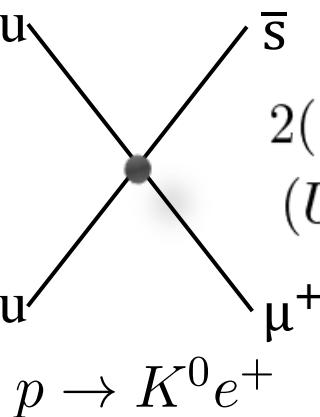
Proton decay

Scenario 2: $Y_d^\dagger = Y_d$ and $Y_u^\dagger = Y_u$



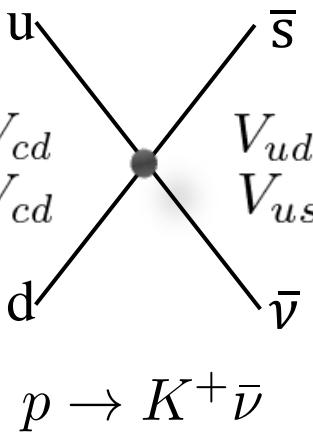
$$2(U_u)_{11}V_{ud} + (U_u)_{12}V_{cd}$$

$$(U_u)_{11}V_{ud} + (U_u)_{12}V_{cd}$$



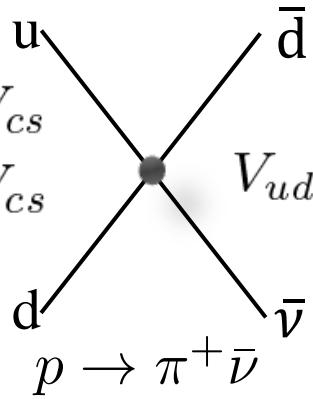
$$2(U_u)_{11}V_{us} + (U_u)_{12}V_{cs}$$

$$(U_u)_{11}V_{us} + (U_u)_{12}V_{cs}$$



$$V_{ud}$$

$$V_{us}$$



numerical solution:

$$\tau(p \rightarrow \pi^0 e^+) =$$

$$\log_{10} \left(\frac{6.88 \times 10^{32}}{(U_u)_{11}^2 + 0.27(U_u)_{11}(U_u)_{12} + 0.02(U_u)_{12}^2} \right)$$

$$\tau(p \rightarrow K^0 e^+) =$$

$$\log_{10} \left(\frac{4.22 \times 10^{34}}{(U_u)_{11}^2 + 5.22(U_u)_{11}(U_u)_{12} + 7.56(U_u)_{12}^2} \right)$$

$$\tau(p \rightarrow K^+ \bar{\nu}) = 3.5 \times 10^{34} \text{ years} > \text{HK bound}$$

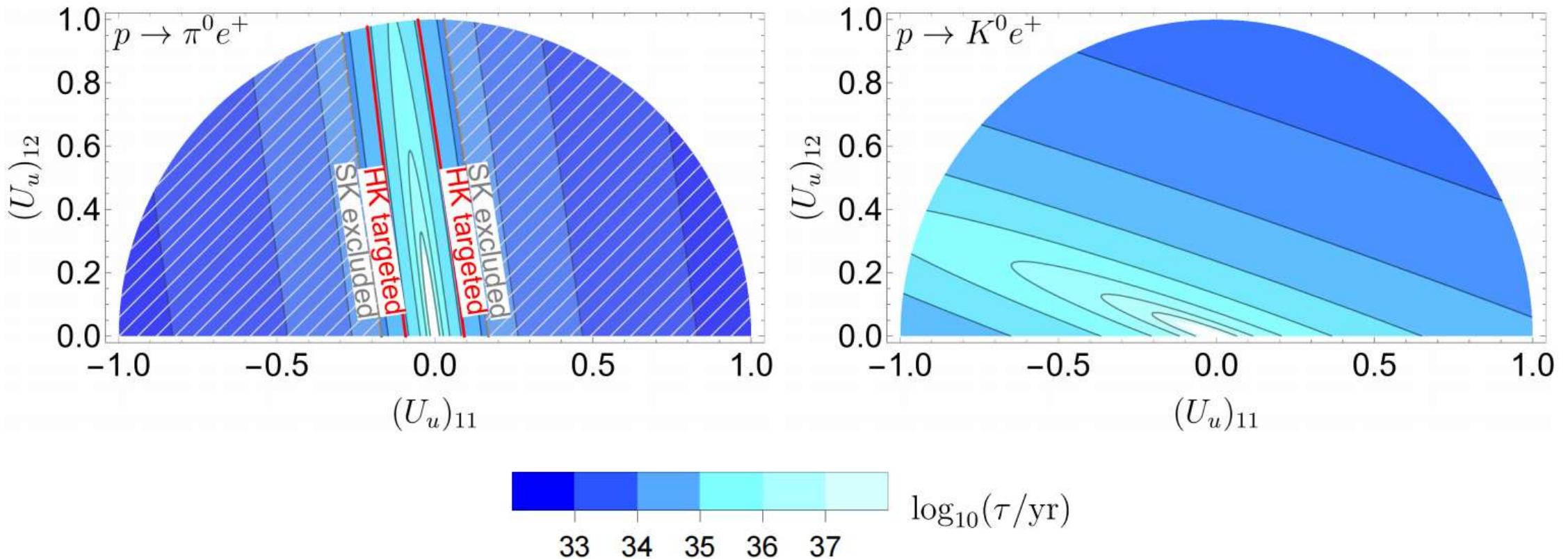
$$\tau(p \rightarrow \pi^+ \bar{\nu}) = 1.8 \times 10^{33} \text{ years} > \text{SK bound}$$

Only two free parameters: $(U_u)_{11}, (U_u)_{12}$

$$(U_u)_{11} \rightarrow 0, (U_u)_{12} \rightarrow 0, \tau \rightarrow +\infty$$

Proton decay

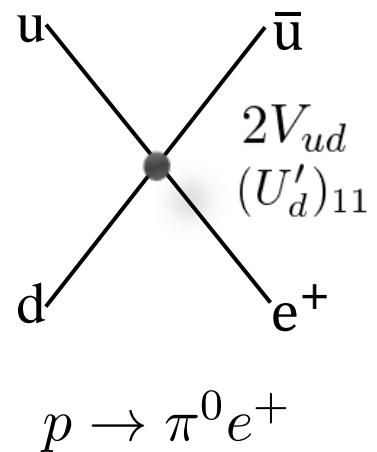
Two free parameters: $(U_u)_{11}, (U_u)_{12}$



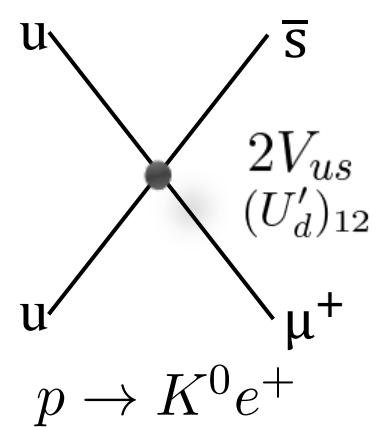
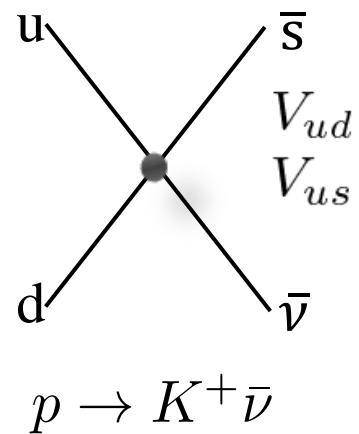
Similar to S1, there is hope to test this assumption in future proton decay experiments

Proton decay

Scenario 3: $Y_u = \hat{Y}_u$



$$\tau(p \rightarrow \pi^0 e^+)_{\max} = 8.5 \times 10^{32} \text{ years} < 2.4 \times 10^{34} \text{ years} (\text{SK bound})$$



Scenario 3 is excluded

numerical solution:

$$\tau(p \rightarrow \pi^0 e^+) = \log_{10}\left(\frac{3.3865 \times 10^{33}}{(U'_d)_{11}^2 + 3.9737}\right),$$

$$\tau(p \rightarrow K^0 e^+) = \log_{10}\left(\frac{1.0607 \times 10^{34}}{(U'_d)_{12}^2 + 0.2012}\right),$$

$$\tau(p \rightarrow K^+ \bar{\nu}) = 3.49 \times 10^{34} \text{ years} > 3.2 \times 10^{34} \text{ years} (\text{HK targeted})$$

Conclusion

1. New particles in $\mathbf{24}_F$ satisfy the following mass hierarchy $M_\Sigma < M_{Q_8} < M_Q$.
2. M_Σ should be lighter than $10^{4.8}\text{GeV}$ in order to make $M_{\text{GUT}} > 10^{15}\text{GeV}$ when $\mathbf{24}_F$ has only one copy.
3. There is hope to test this economical model in future proton decay experiments.
4. Future neutrino experiments can provide multiple tests of GUTs via multiple proton decay channels.

Thanks!