

# Jet Origin Identification & Quantum-based Jet Clustering

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2024.08.15

第十四届全国粒子物理学术会议@青岛

# Motivation:

- 1, Quarks and gluons carry color charge and can not travel freely. Once generated in high-energy collisions, quarks, and gluon would fragment into numerous particles.
- 2, Deep learning can promote the development of high-energy physics.
- 3, Exploring the application of quantum technologies to jet clustering is key to fostering innovation for quantum computing and high-energy physics.

# Contents:

1. Jet Origin Identification (JOI)
2. Application of Quantum Approximate Optimization Algorithm (QAOA) to Jet Clustering

# Jet Origin Identification

categorizes jets into 5 quarks ( $b, c, s, u, d$ ), 5 anti-quark ( $\bar{b}, \bar{c}, \bar{s}, \bar{u}, \bar{d}$ ), and gluon  
= jet flavor tagging + jet charge measurement + s-quark tagging + gluon finding.

PHYSICAL REVIEW LETTERS 132, 221802 (2024)

Eur. Phys. J. C (2024) 84:152  
<https://doi.org/10.1140/epjc/s10052-024-12475-5>

THE EUROPEAN  
PHYSICAL JOURNAL C



Regular Article - Experimental Physics

## Jet-Origin Identification and Its Application at an Electron-Positron Higgs Factory

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(Received 16 October 2023; revised 26 April 2024; accepted 1 May 2024; published 31 May 2024)

To enhance the scientific discovery power of high-energy collider experiments, we propose and realize the concept of jet-origin identification that categorizes jets into five quark species ( $b, c, s, u, d$ ), five antiquarks ( $\bar{b}, \bar{c}, \bar{s}, \bar{u}, \bar{d}$ ), and the gluon. Using state-of-the-art algorithms and simulated  $\nu\bar{\nu}H, H \rightarrow jj$  events at 240 GeV center-of-mass energy at the electron-positron Higgs factory, the jet-origin identification simultaneously reaches jet flavor tagging efficiencies ranging from 67% to 92% for bottom, charm, and strange quarks and jet charge flip rates of 7%–24% for all quark species. We apply the jet-origin identification to Higgs rare and exotic decay measurements at the nominal luminosity of the Circular Electron Positron Collider and conclude that the upper limits on the branching ratios of  $H \rightarrow s\bar{s}, u\bar{u}, d\bar{d}$  and  $H \rightarrow sb, db, uc, ds$  can be determined to  $2 \times 10^{-4}$  to  $1 \times 10^{-3}$  at 95% confidence level. The derived upper limit for  $H \rightarrow s\bar{s}$  decay is approximately 3 times the prediction of the standard model.

DOI: [10.1103/PhysRevLett.132.221802](https://doi.org/10.1103/PhysRevLett.132.221802)

## ParticleNet and its application on CEPC jet flavor tagging

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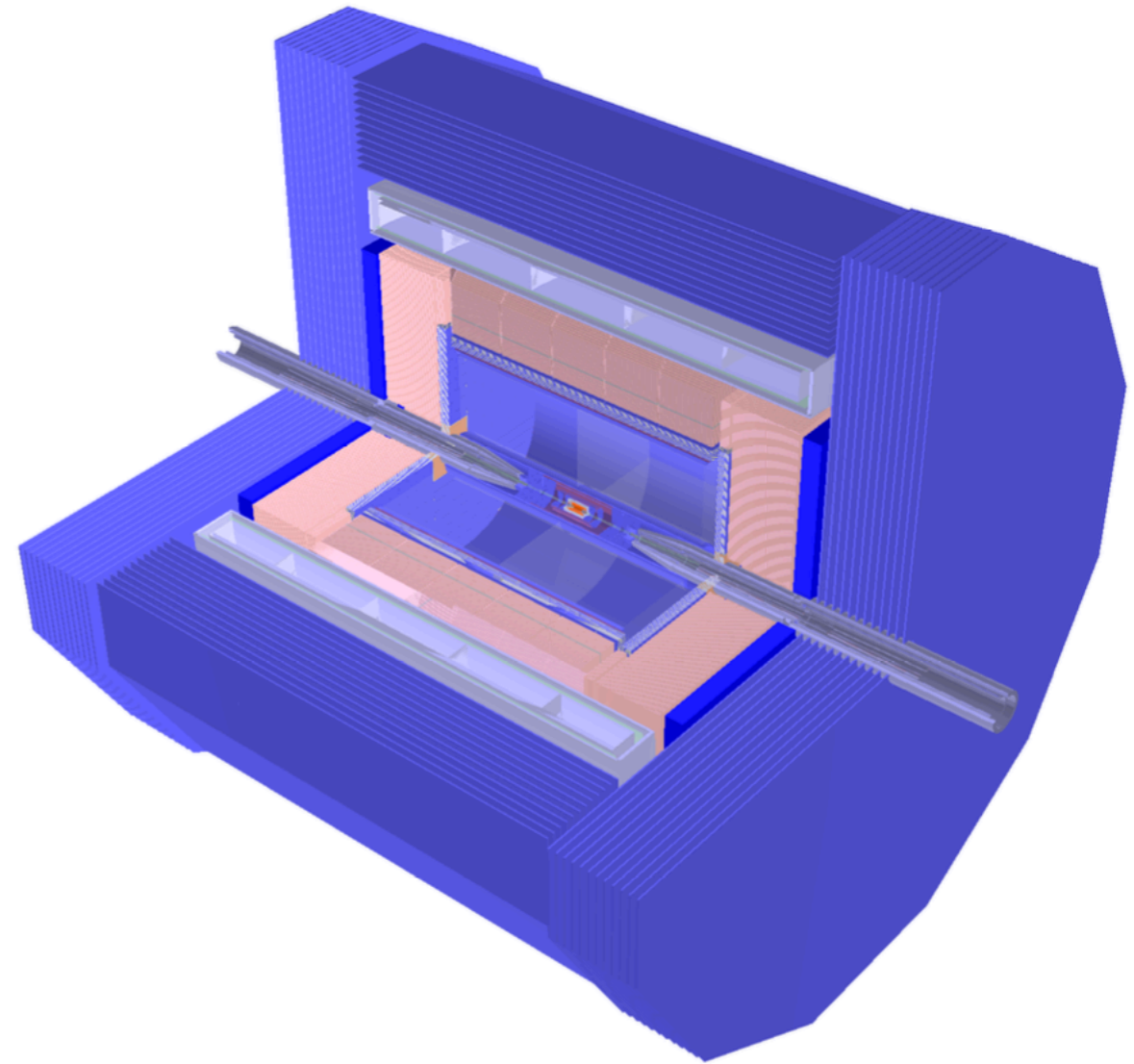
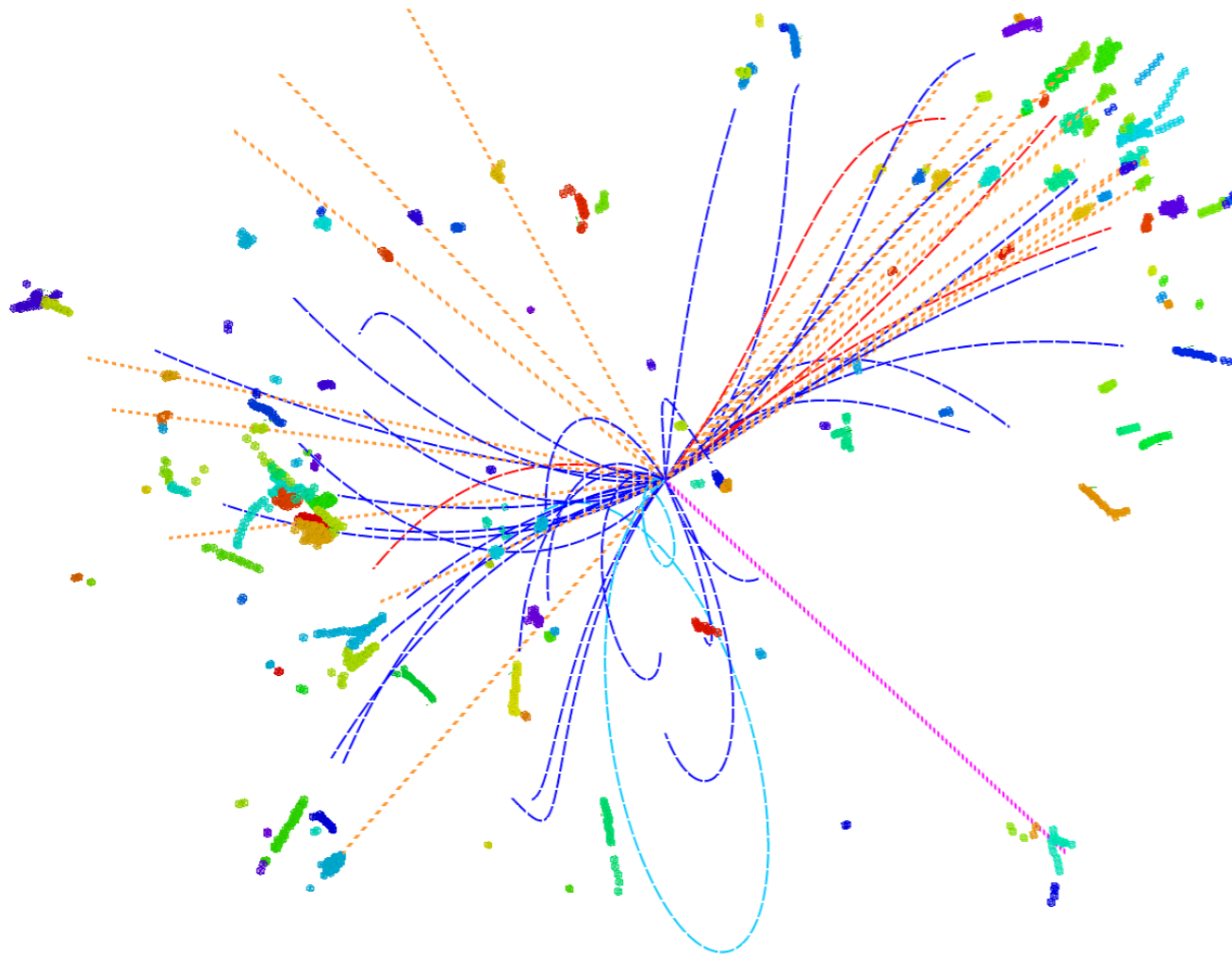
Received: 15 November 2023 / Accepted: 23 January 2024  
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**Abstract** Quarks (except top quarks) and gluons produced in collider experiments hadronize and fragment into sprays of stable particles, called jets. Identification of quark flavor is desired for collider experiments in high-energy physics, relying on flavor tagging algorithms. In this study, using a full simulation of the Circular Electron Positron Collider (CEPC), we investigate the flavor tagging performance of two different algorithms: ParticleNet, based on a Graph Neural Network, and LCFIPlus, based on the Gradient Boosted Decision Tree. Compared to LCFIPlus, ParticleNet significantly enhances flavor tagging performance, resulting in a significant improvement in benchmark measurement accuracy, i.e., a 36% improvement for  $\sigma(ZH) \cdot Br(Z \rightarrow \nu\bar{\nu}, H \rightarrow c\bar{c})$  measurement and a 75% improvement for  $|V_{cb}|$  measurement via  $W$  boson decay, respectively, when the CEPC operates as a Higgs factory at the center-of-mass energy of 240 GeV and collects an integrated luminosity of  $5.6 \text{ ab}^{-1}$ . We compare the performance of ParticleNet and LCFIPlus at different vertex detector configurations, observing that the inner radius is the most sensitive parameter, followed by material budget and spatial resolution.

light on the properties of massive SM particles and is critical for experimental exploration at the high-energy frontier. Flavor tagging is used to distinguish jets which hadronize from quarks of different flavors or from gluons. To promote the development of future electron-positron Higgs factories, which is regarded as a high priority future collider [4], accurate performance analysis and optimization of both detectors and algorithms are essential. Jet flavor tagging and relevant benchmark analyses serve as good objectives.

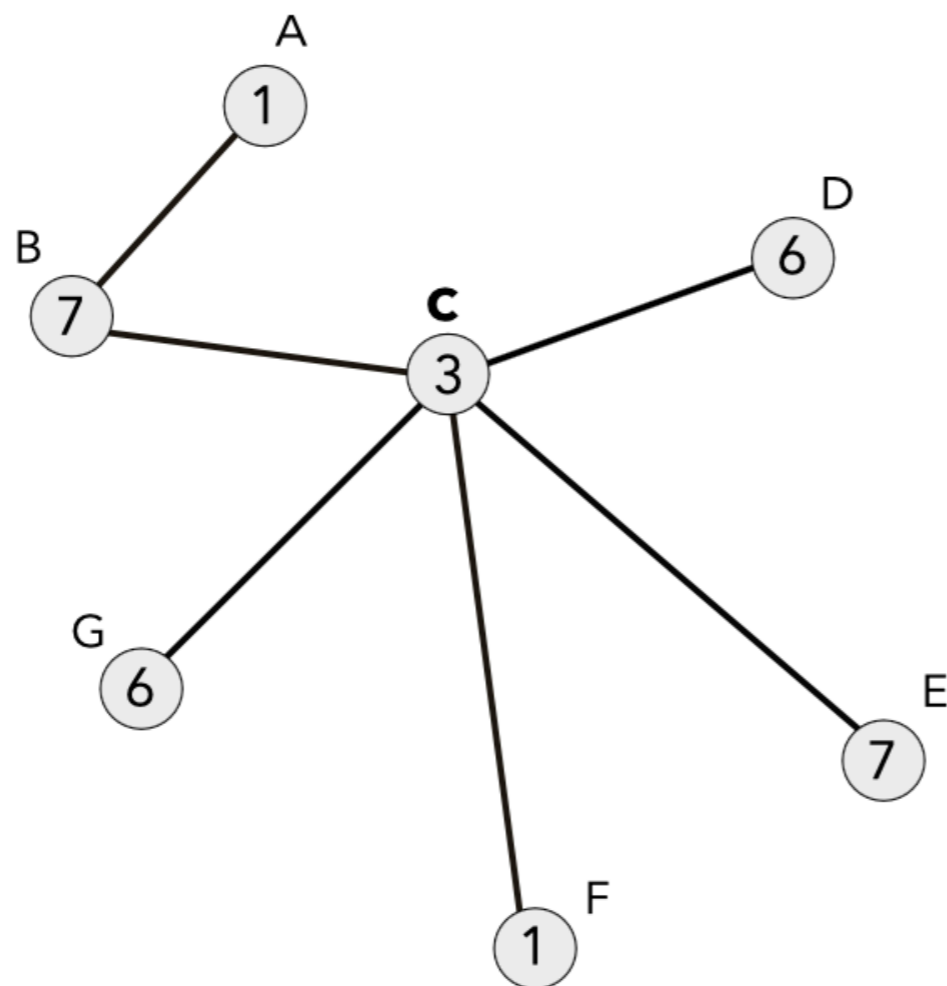
The Circular Electron Positron Collider (CEPC) [5] is a large-scale collider facility that was proposed after the discovery of the Higgs boson in 2012. It is designed to have a circumference of 100 km with two interaction points. It will be able to operate at multiple center-of-mass energies, including 240 GeV as a Higgs factory, 160 GeV for a  $W^+W^-$  threshold scan, and 91 GeV as a  $Z$  factory. It also can be upgraded to 360 GeV for a  $t\bar{t}$  threshold scan. Table 1 summarizes its baseline operating scheme and the corresponding boson yield predictions [6]. One of the main scientific objectives of the CEPC is the precise measurement of properties of the Higgs boson. Additionally, trillions of  $Z \rightarrow q\bar{q}$  events can provide

# Samples



Full Simulated  $vvH$ , Higgs to two jets sample at CEPC baseline configuration: CEPC-v4 detector, 1 Million samples each, 60/20/20% for training, validation & test

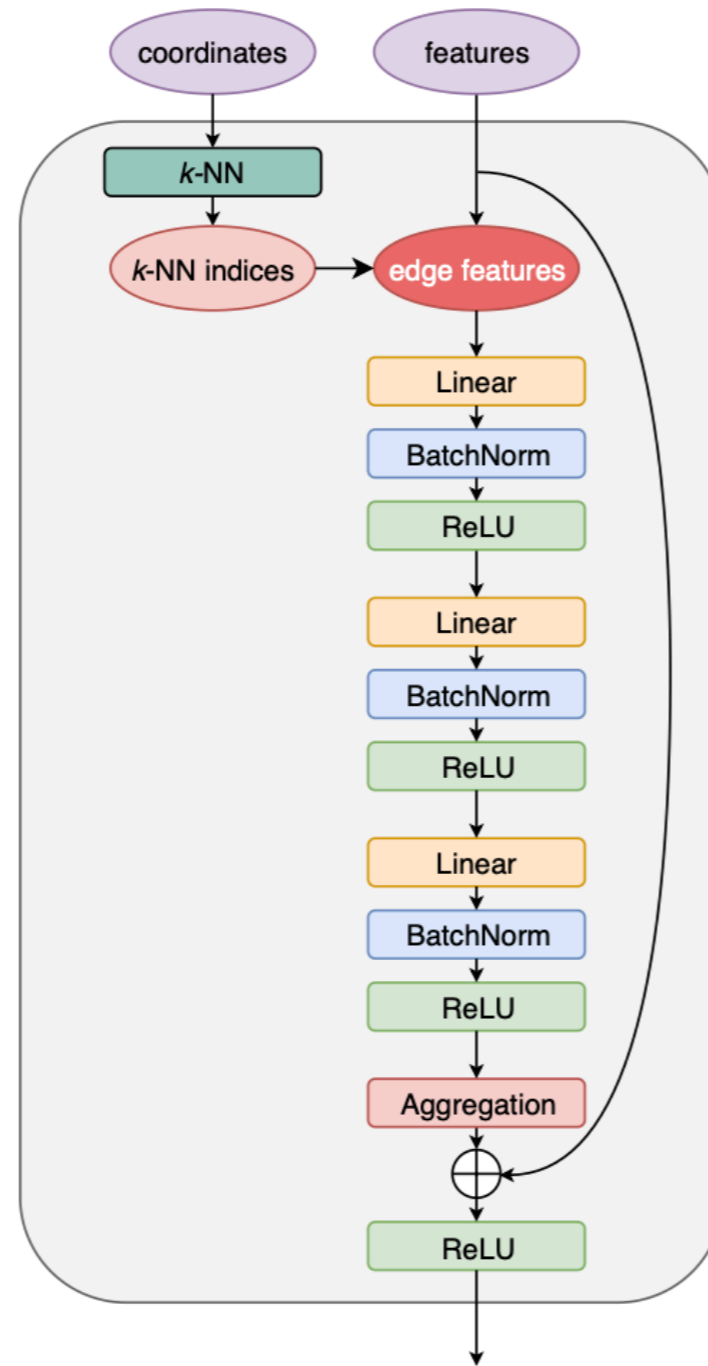
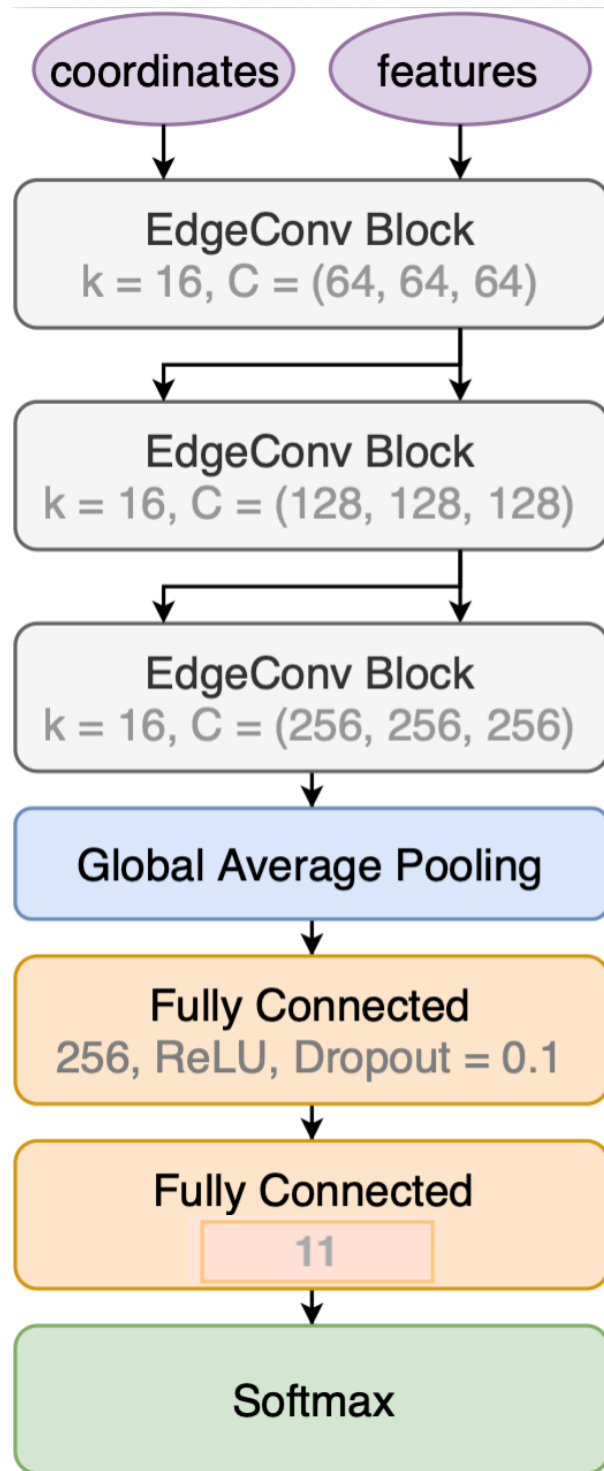
jet represented as :



Graph = (Vertices, Edges)

Construct GNN  
with "message passing"

# the ParticleNet and input features



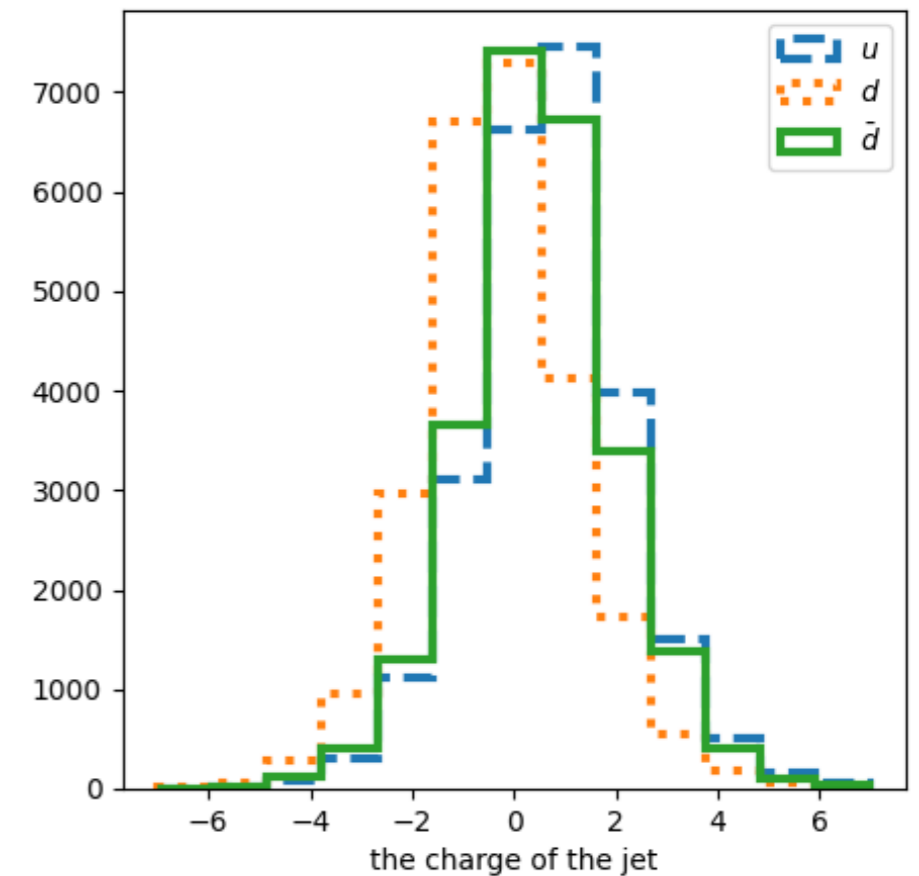
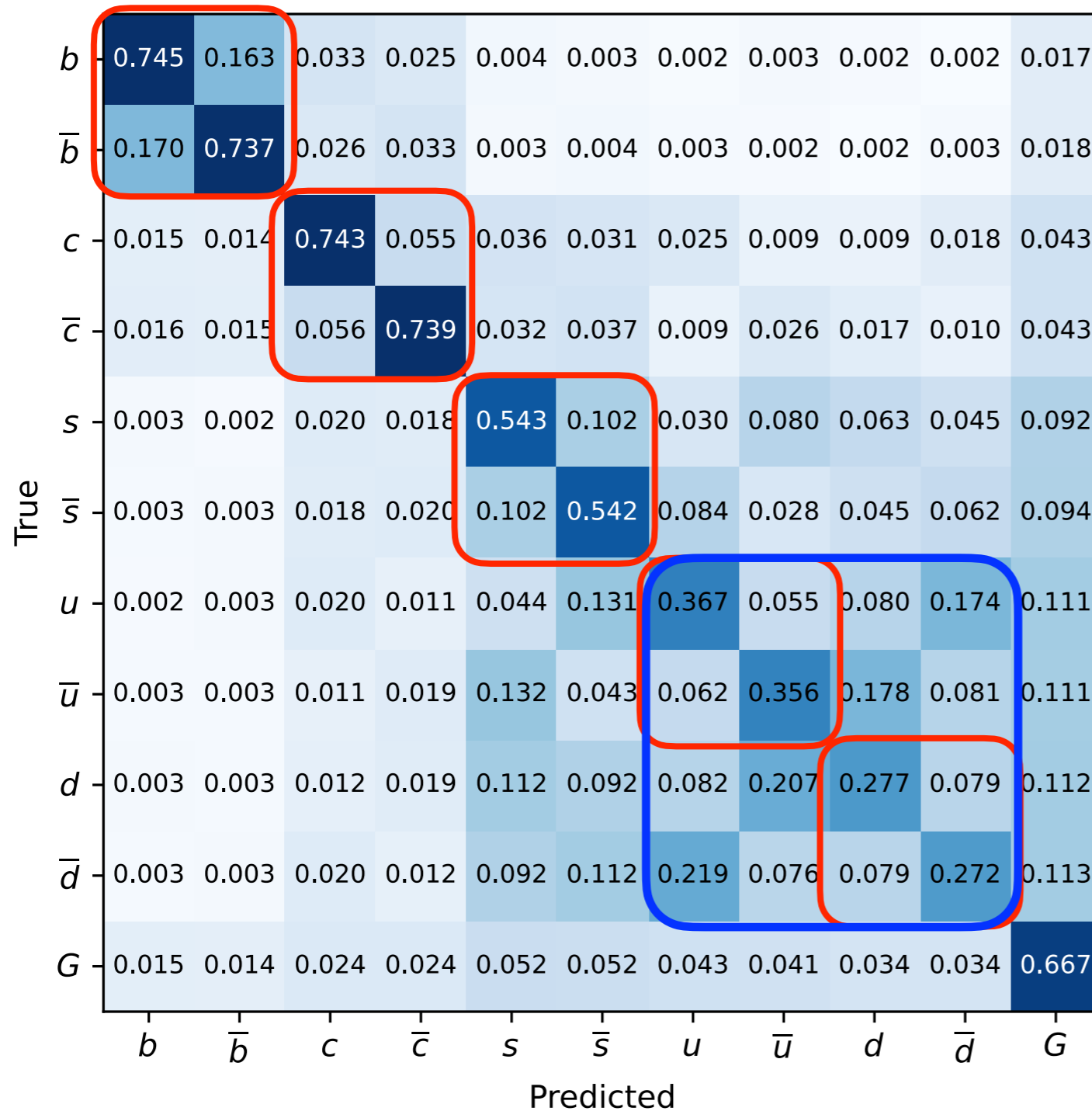
## input features :

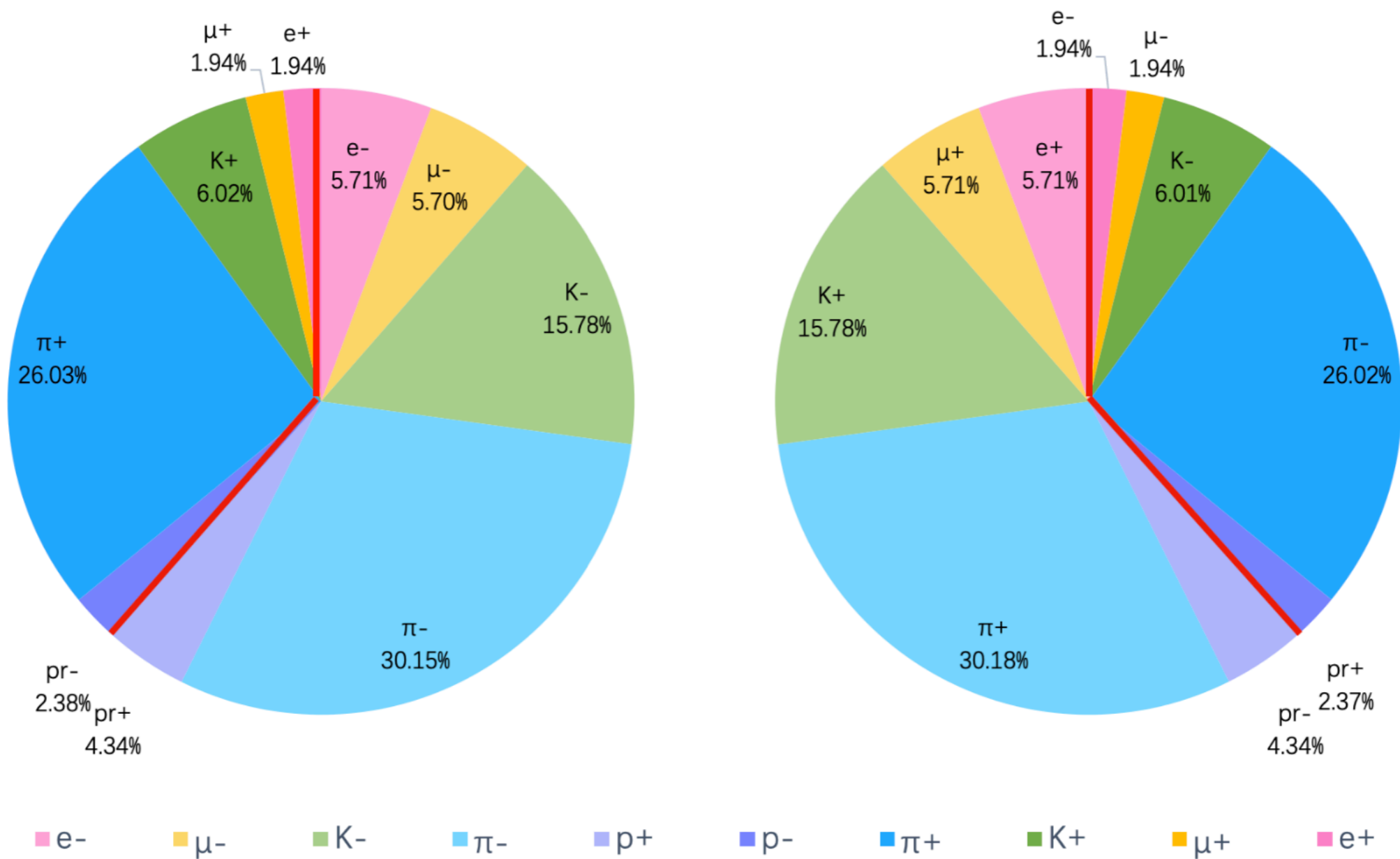
**Table 3** The input variables used in ParticleNet for jet flavor tagging at the CEPC

Variable	Definition
$\Delta \eta$	Difference in pseudorapidity between the particle and the jet axis
$\Delta \phi$	Difference in azimuthal angle between the particle and the jet axis
$\log P_t$	Logarithm of the particle's $P_t$
$\log E$	Logarithm of the particle's energy
$\log \frac{P_t}{P_t(jet)}$	Logarithm of the particle's $P_t$ relative to the jet $P_t$
$\log \frac{E}{E(jet)}$	Logarithm of the particle's energy relative to the jet energy
$\Delta R$	Angular separation between the particle and the jet axis
$d_0$	Transverse impact parameter of the track
$d_{0err}$	Uncertainty associated with the measurement of the $d_0$
$z_0$	Longitudinal impact parameter of the track
$z_{0err}$	Uncertainty associated with the measurement of the $z_0$
Charge	Electric charge of the particle
isElectron	Whether the particle is an electron
isMuon	Whether the particle is a muon
isChargedKaon	Whether the particle is a charged Kaon
isChargedPion	Whether the particle is a charged Pion
isProton	Whether the particle is a proton
isNeutralHadron	Whether the particle is a neutral hadron
isPhoton	Whether the particle is a photon

# the performance of jet origin identification

ParticleNet algorithm attaches each jet with 11 likelihoods corresponding to 11 types of jets. Then the jet type is determined according to the maximum likelihood.





**Figure 5.** The percentages of species of final state leading charged particles within the  $b$  jet (left) and the  $\bar{b}$  jet (right) by WHIZARD 1.95.

[arXiv:2306.14089](https://arxiv.org/abs/2306.14089)



# Benchmark physics analyses

Begin with the existing analyses of  $\nu\bar{\nu}H, H \rightarrow b\bar{b}/c\bar{c}/gg$ , (arXiv:2203.01469) and combining the jet origin identification, we obtain the upper limits on branching ratios of seven Higgs rare and FCNC hadronic decay modes.

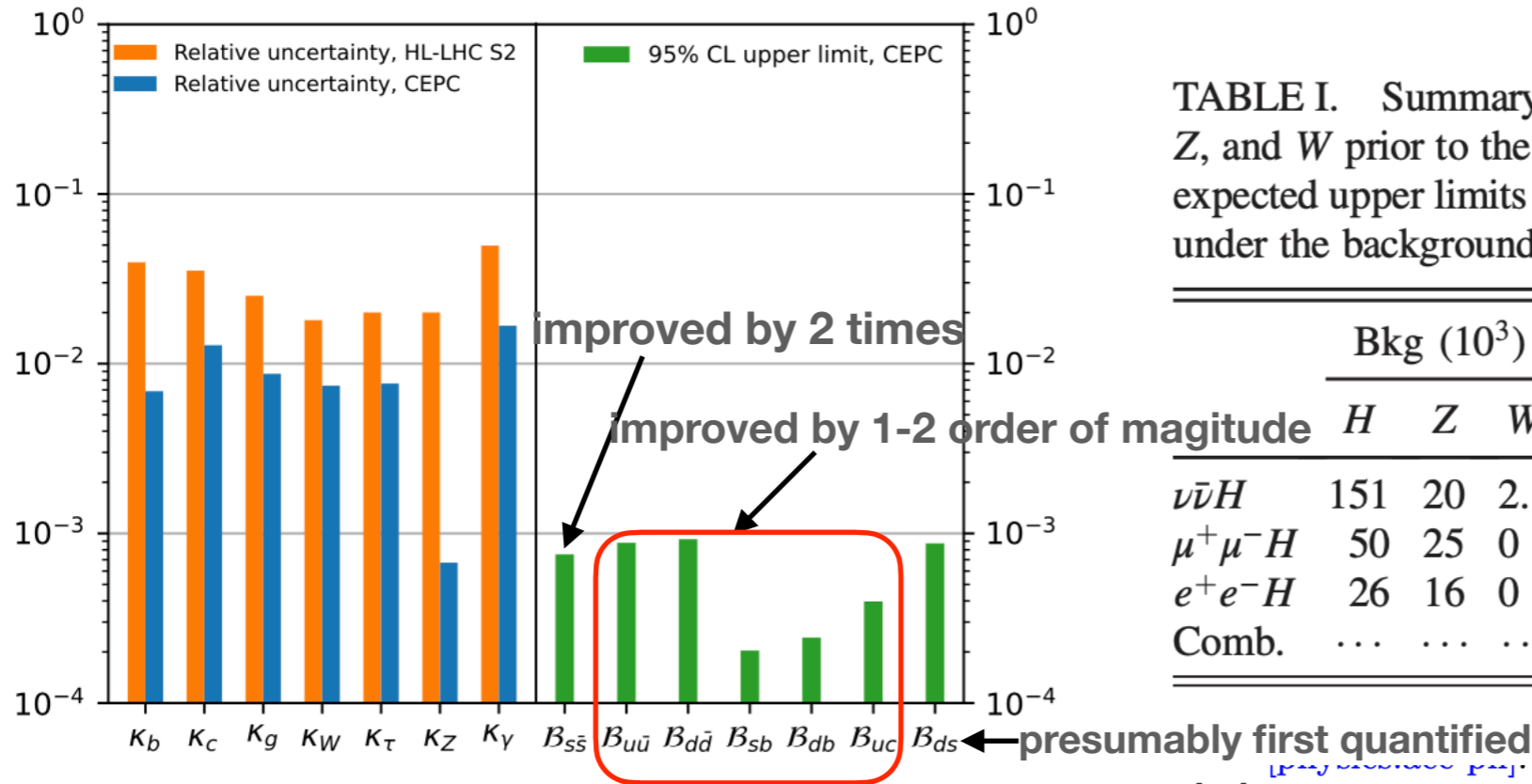


FIG. 5. Expected upper limits on the branching ratios of rare Higgs boson decays from this Letter (green) and the relative uncertainties of Higgs couplings anticipated at CEPC [19] (blue) and HL-LHC [43] (orange) under the kappa-0 fit scenario [54] and scenario S2 of systematics [55], as cited in Ref. [19]. The limit on  $B_{s\bar{s}}$  corresponds to an upper limit of 1.7 on the Higgs-strange coupling modifier  $\kappa_s$  (not shown).

TABLE I. Summary of background yields from  $H \rightarrow b\bar{b}/c\bar{c}/gg$ ,  $Z$ , and  $W$  prior to the flavor-based event selection, along with the expected upper limits on Higgs decay branching ratios at 95% CL under the background-only hypothesis.

	Bkg ( $10^3$ )			Upper limits on Br ( $10^{-3}$ )						
	$H$	$Z$	$W$	$s\bar{s}$	$u\bar{u}$	$d\bar{d}$	$sb$	$db$	$uc$	$ds$
$\nu\bar{\nu}H$	151	20	2.1	0.81	0.95	0.99	0.26	0.27	0.46	0.93
$\mu^+\mu^-H$	50	25	0	2.6	3.0	3.2	0.5	0.6	1.0	3.0
$e^+e^-H$	26	16	0	4.1	4.6	4.8	0.7	0.9	1.6	4.3
Comb.	...	...	...	0.75	0.91	0.95	0.22	0.23	0.39	0.86

[23] J. Duarte-Campderros, G. Perez, M. Schlaffer, and A. Soffer, *Phys. Rev. D* **101**, 115005 (2020), arXiv:1811.09636 [hep-ph].

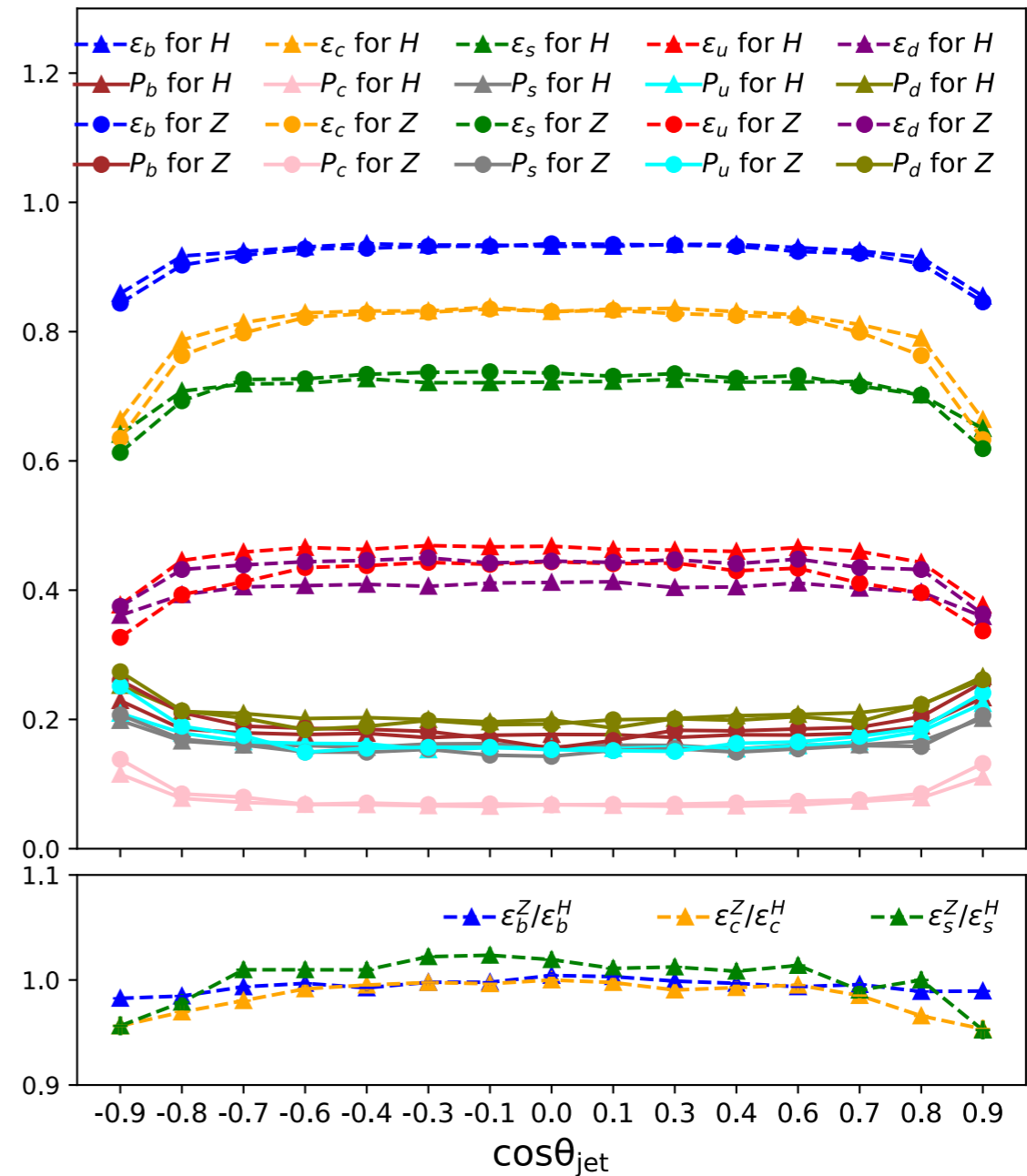
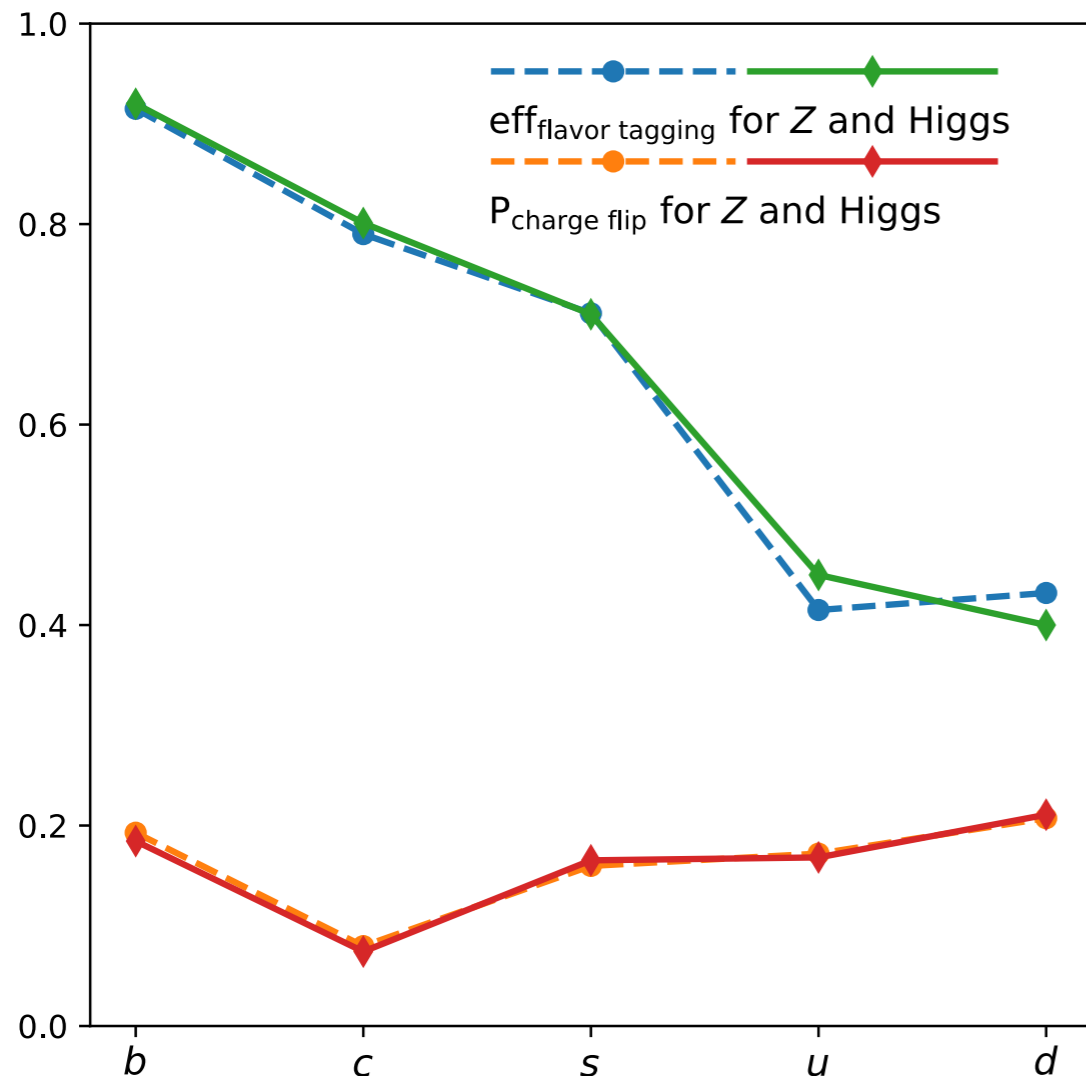
[44] A. Albert *et al.*, "Strange quark as a probe for new physics in the higgs sector," (2022), arXiv:2203.07535 [hep-ex].

[53] J. de Blas *et al.*, *JHEP* **01**, 139 (2020), arXiv:1905.03764 [hep-ph].

[54] J. De Blas, G. Durieux, C. Grojean, J. Gu, and A. Paul, *JHEP* **12**, 117 (2019), arXiv:1907.04311 [hep-ph].

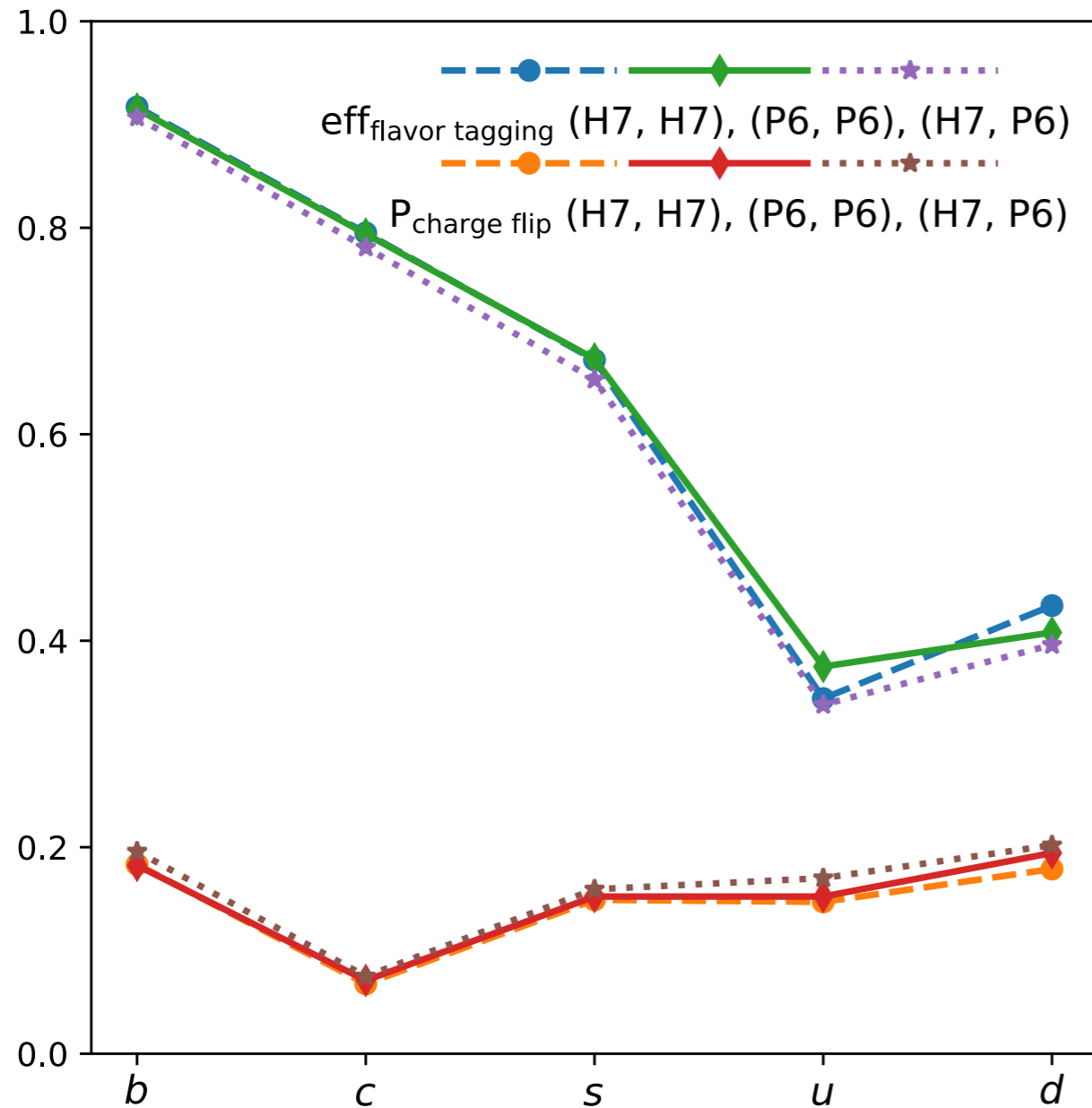
# Comparison between different physics processes

$Z \rightarrow q\bar{q}$  at  $\sqrt{s} = 91.2 \text{ GeV}$   
 $\nu\bar{\nu}H, H \rightarrow q\bar{q}$  at  $\sqrt{s} = 240 \text{ GeV}$



The jet origin identification performance agrees with each other, especially in the fiducial barrel region of the detector for the flavor tagging performance of b, c, and s.

# Comparison between different hadronization models



Pythia-6.4  
Herwig-7.2.2

$\nu\bar{\nu}H, H \rightarrow jj$  at  $\sqrt{s} = 240 \text{ GeV}$

(A, B) means: training on A, test on B

The jet origin identification performance agrees with each other, especially for b, c, and s jets, while exhibits small but visible differences for u and d jets.

# A Novel Quantum Realization of Jet Clustering in High-Energy Physics Experiments

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arXiv:2407.09056

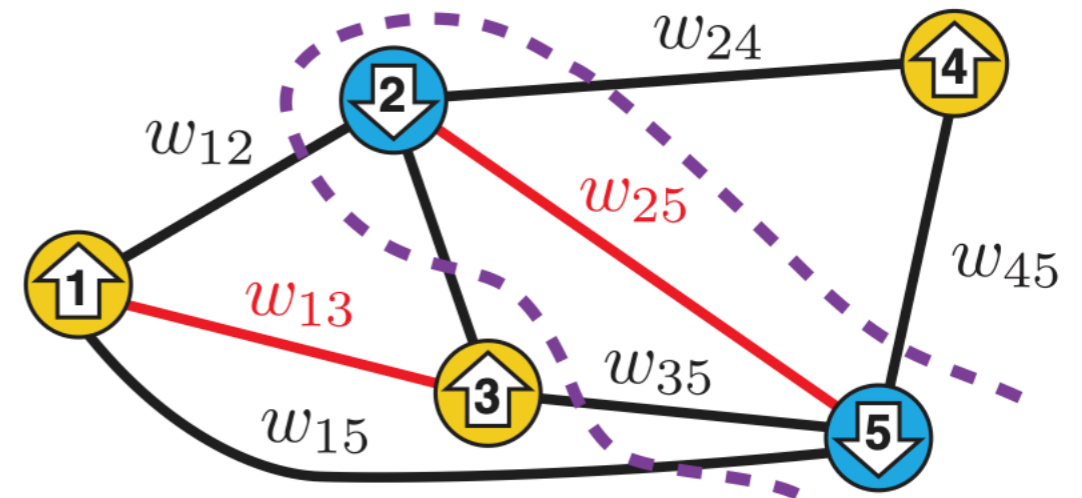
# MaxCut problem

undirected graph  $G = (V, E)$

$V$ : set of vertices

$E$ : set of edges

$w_{ij} > 0$ : the weight of the edge  $(ij) \in E$



Phys. Rev. X 10, 021067

goal: partition the graph vertices into two complementary subsets to maximize the total weight of edges with two vertices belong to two

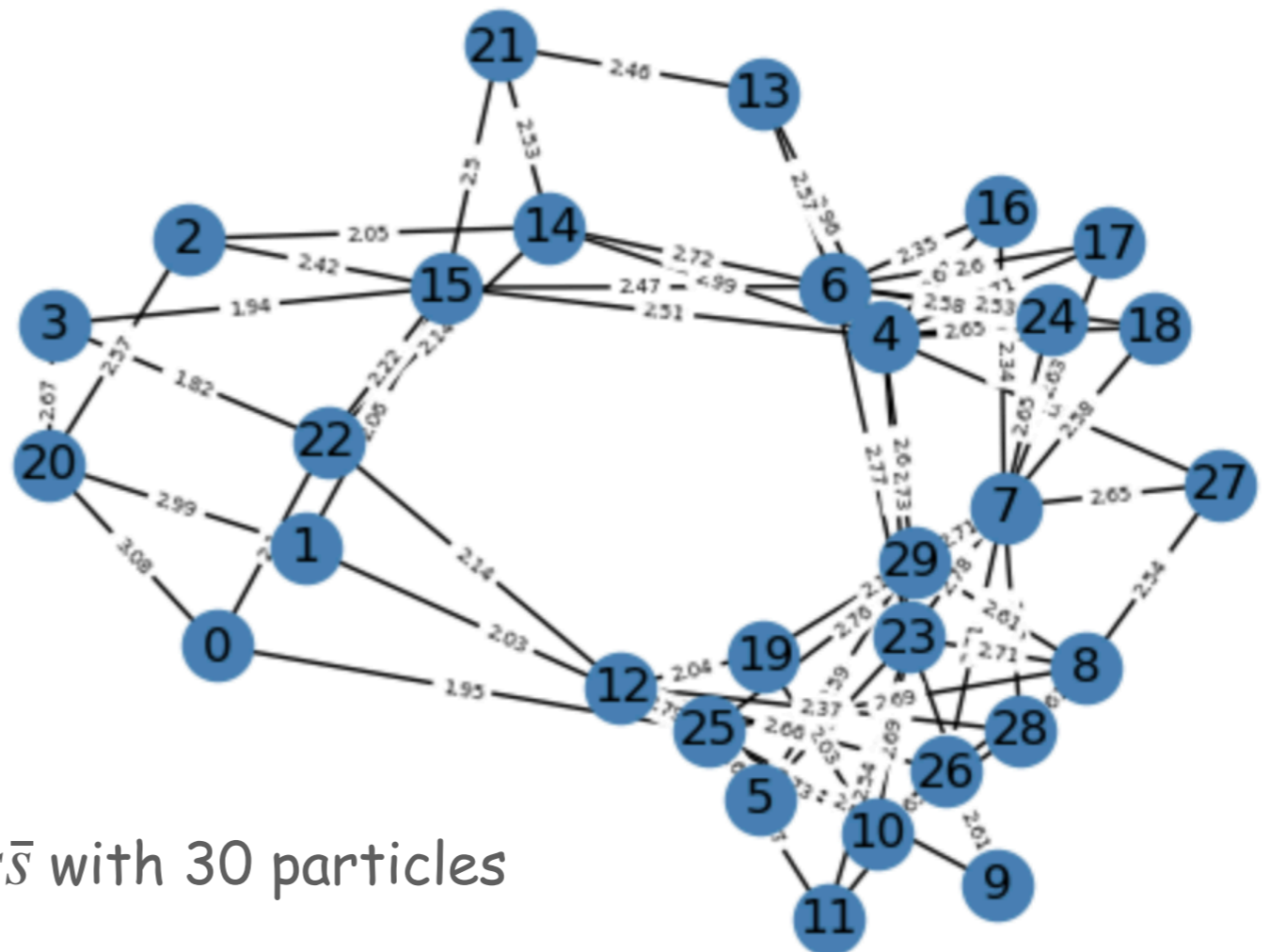
subsets 
$$C(x) = \sum_{i,j=1}^{|V|} w_{ij} x_i (1 - x_j)$$

A collision event can be represented as a graph

particles as vertices

the angle of two particles as the edge weight

only the k leading large edges are retained (k-regular graph)



$e^+e^- \rightarrow ZH \rightarrow \nu\bar{\nu}s\bar{s}$  with 30 particles

the process of AQC :

1, define the Hamiltonian:  $H(t) = (1 - s(t))H_D + s(t)H_C$  and let our quantum system evolve under it,  $U(t) = \tau e^{\frac{-i}{\hbar} \int_0^t H(T) dT}$ .

2, We discretize  $U(T)$  into intervals  $\Delta t$  small enough that the Hamiltonian is approximately constant over each interval.

3, Let  $U(b, a)$  represent time evolution from time a to time b

$$U(T,0) = U(T, T - \Delta t)U(T - \Delta t, T - 2\Delta t) \dots U(\Delta t, 0) = \prod_{j=1}^P U(j\Delta t, (j-1)\Delta t) \approx \prod_{j=1}^P e^{-iH(j\Delta t)\Delta t}$$

$$U(T,0) \approx \prod_{j=1}^P e^{-i(1-s(j\Delta t))H_D\Delta t} e^{-is(j\Delta t)H_C\Delta t} = \prod_{j=1}^P e^{-i\beta_j H_D} e^{-i\gamma_j H_C} = \prod_{j=1}^P \hat{U}_D(\beta_j) \hat{U}_C(\gamma_j)$$

Thus we can approximate AQC by repeatedly letting the system evolve under  $H_C$  for some time  $\gamma_j$  and then  $H_D$  for some time  $\beta_j$ .

# QAOA for MaxCut problem

$$1. U(T,0) \approx \prod_{j=1}^P e^{-i\beta_p H_D} e^{-i\gamma_p H_C} = \prod_{j=1}^P \hat{U}_D(\beta_j) \hat{U}_C(\gamma_j)$$

$$2. H_D : B = \sum_{j=1}^n \sigma_j^x$$

$$3. H_C : C = \frac{1}{2} \sum_{(i,j) \in E} W_{ij} (I - \sigma_i^z \sigma_j^z)$$

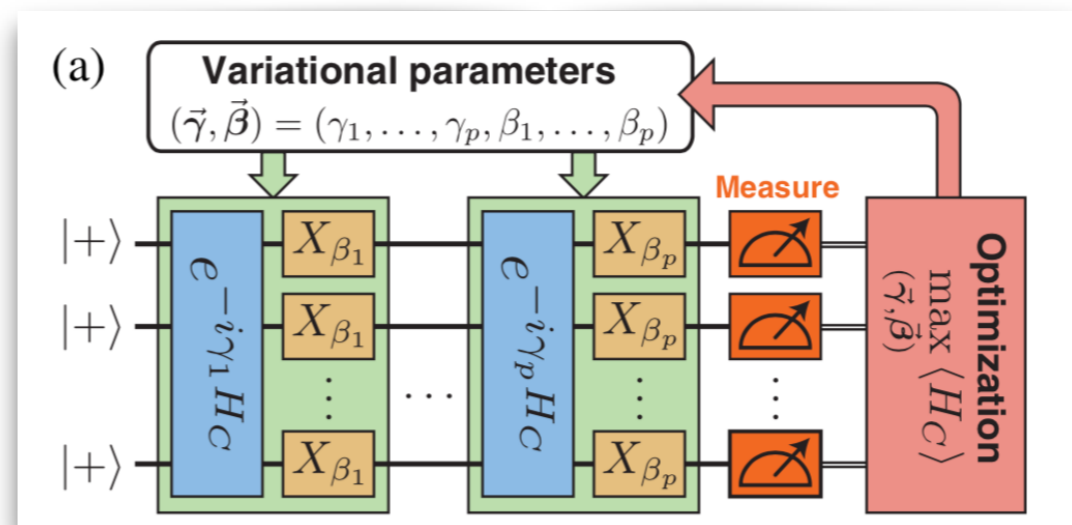
$$4. \text{ Initialize the system in the state } |s\rangle = |+\rangle^{\otimes n} = \frac{1}{\sqrt{2^n}} \sum_{x \in \{0,1\}^n} |x\rangle$$

5. Construct the circuit (ansatz) by applying the unitaries  $\hat{U}_C(\gamma_j)$  and  $\hat{U}_D(\beta_j)$  repeatedly P times

$$6. \text{ The final state output by the circuit is } |\psi_P(\gamma, \beta)\rangle = \hat{U}_D(\beta_P) \hat{U}_C(\gamma_P) \dots \hat{U}_D(\beta_1) \hat{U}_C(\gamma_1) |s\rangle$$

7. The expectation value  $\hat{H}_C$  with respect to the state  $|\psi_P(\gamma, \beta)\rangle$  is calculated through repeated measurements (1024 times in this analysis) of the final state on the computational basis,  $F_P(\gamma, \beta) = \langle \psi_P(\gamma, \beta) | \hat{H}_C | \psi_P(\gamma, \beta) \rangle$

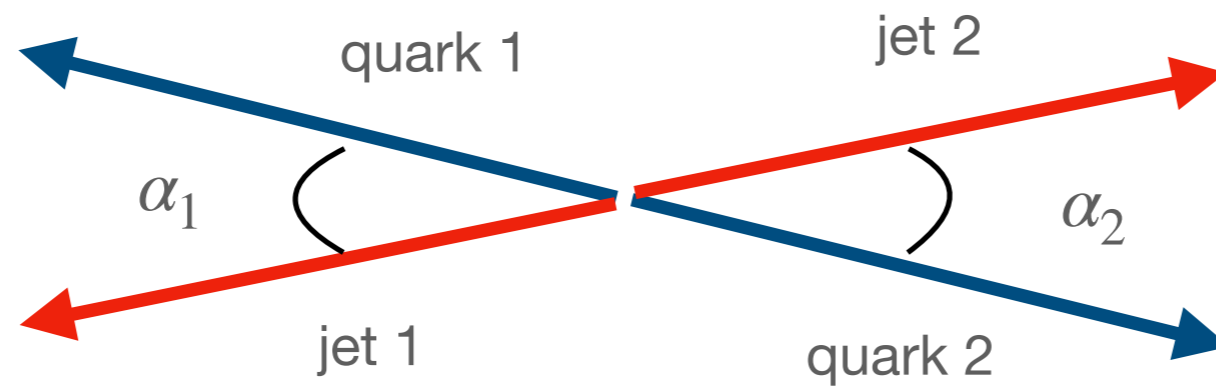
8. A classical optimization algorithm is employed to iteratively update the parameters  $\gamma$  and  $\beta$  to find the optimal set of parameters  $(\gamma^*, \beta^*)$  such that the expectation value  $F_P(\gamma, \beta)$  is maximized.





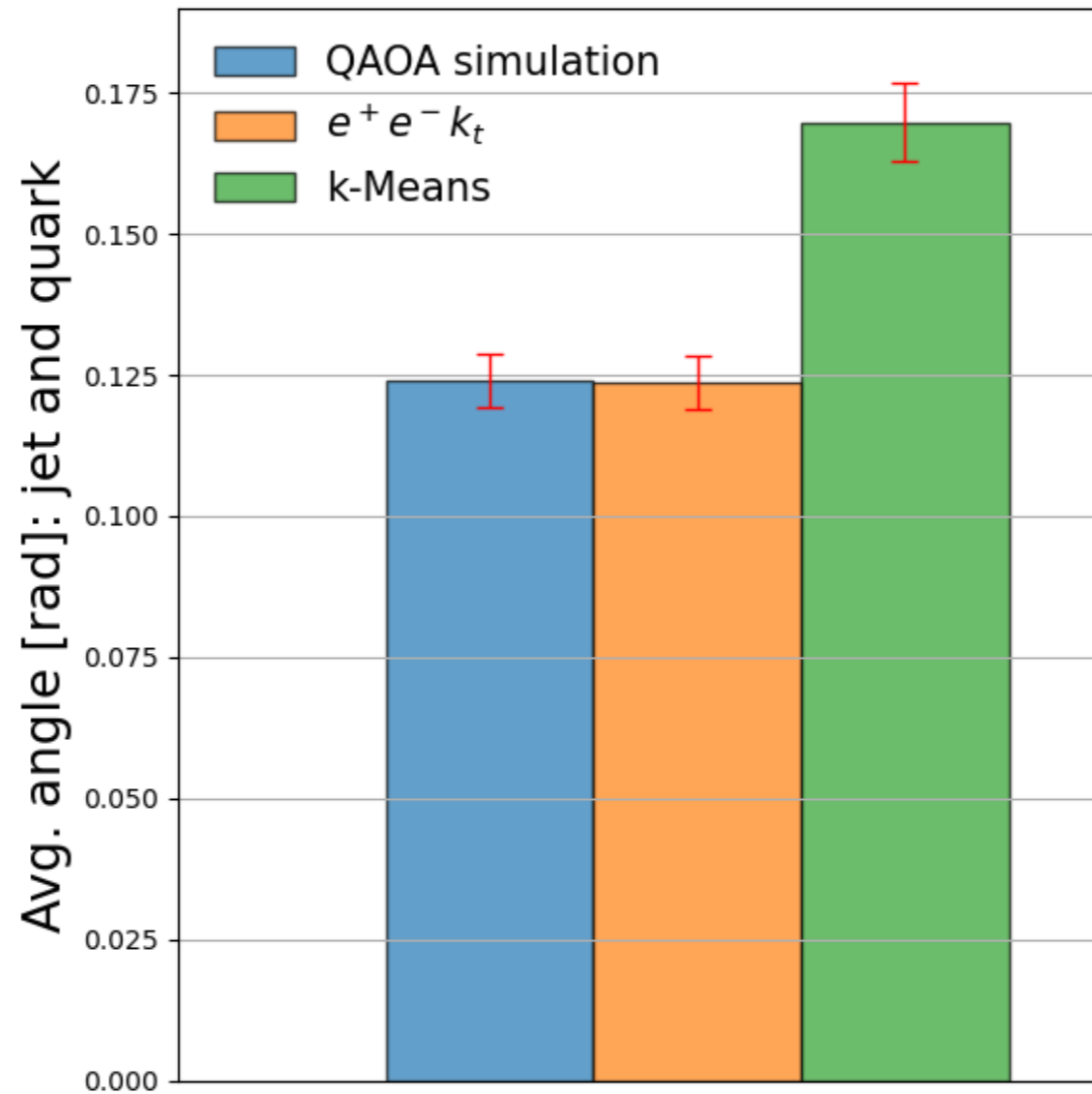
# samples and the criteria of jet clustering performance

4000  $e^+e^- \rightarrow ZH \rightarrow v\bar{v}s\bar{s}$  with 30 particles



criteria is  $\alpha = \alpha_1 + \alpha_2$

compare jet clustering performance obtained by QAOA,  
 $e^+e^-k_t$  and k-Means

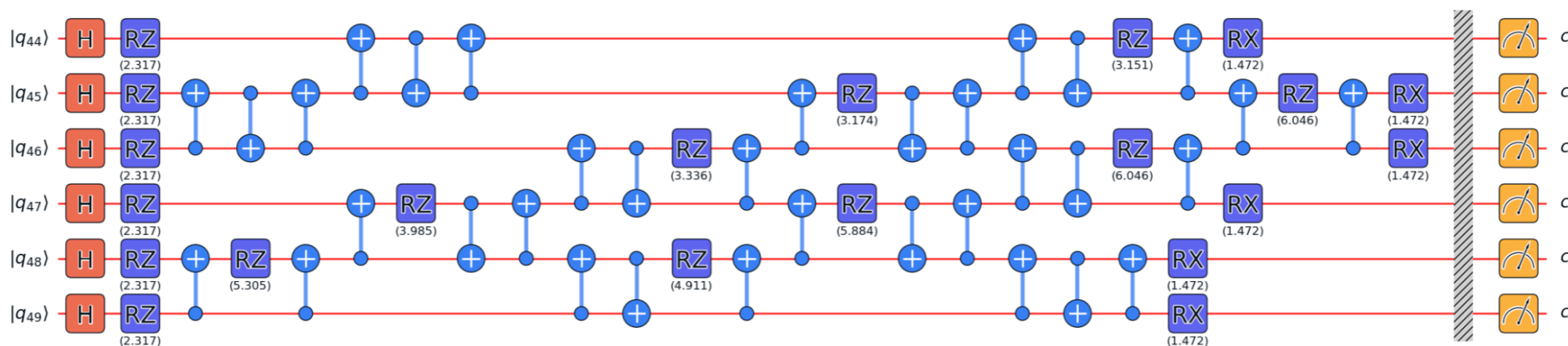
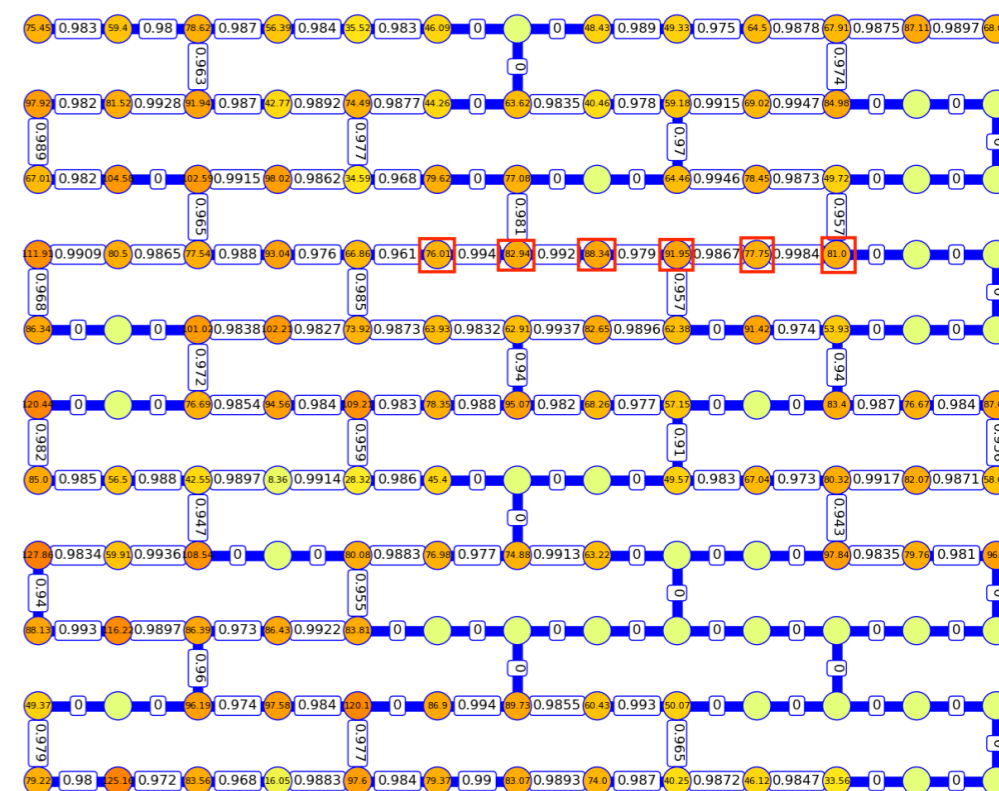


depth=3  
k=7

This comparison highlights the potential of the QAOA in the jet clustering problem.

# Conduct on the Baihua processor

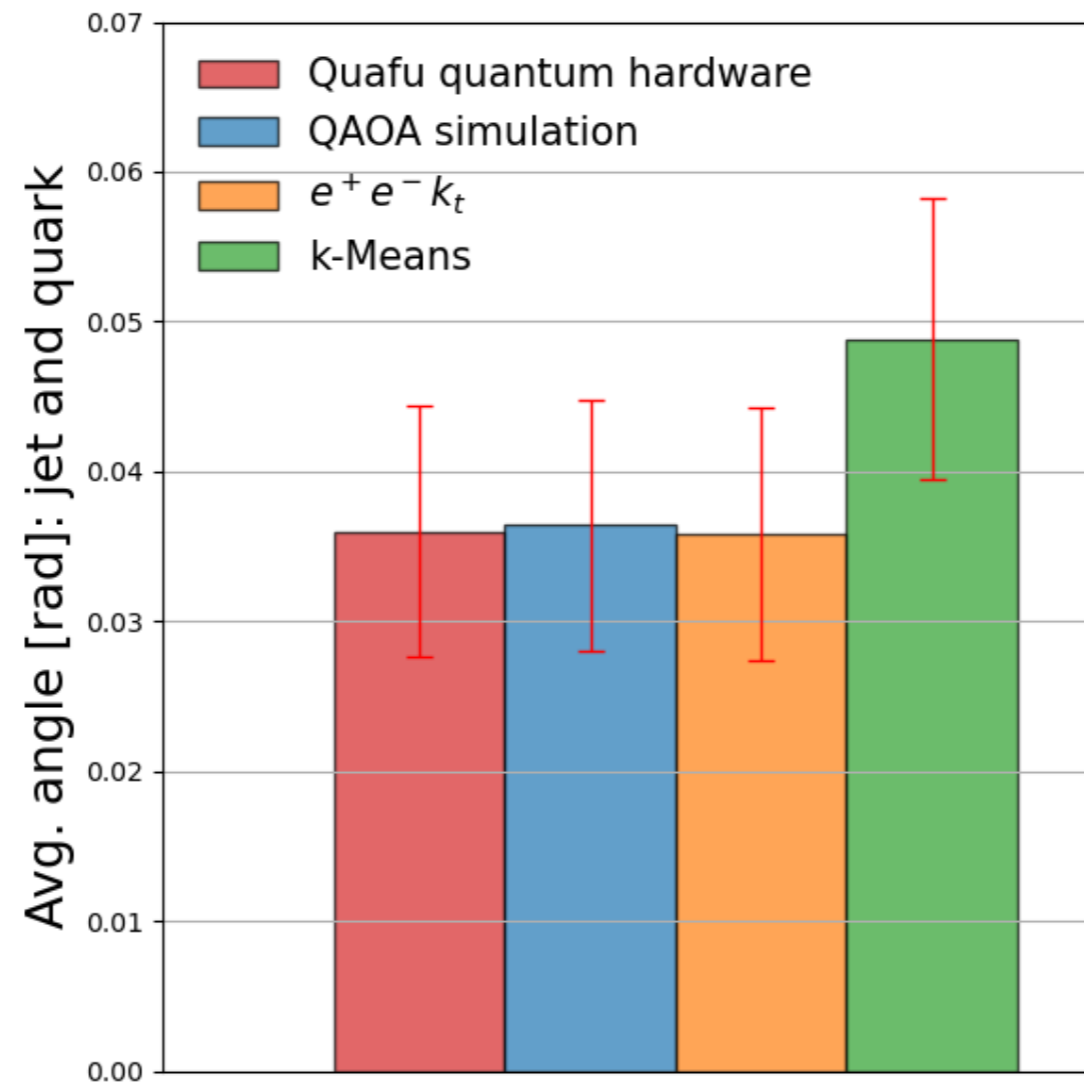
- 123 operational qubits
- relaxation time  $T_1$  of 73.994  $\mu s$
- dephasing time  $T_2^*$  of 29.02  $\mu s$
- fidelity of single-qubit gate 99.9%
- fidelity of two-qubit gates (CZ) 98.9%



$e^+e^- \rightarrow ZH \rightarrow \nu\bar{\nu}s\bar{s}$  with 6 particles

compiled QAOA circuit on Baihua processor reaches a depth of 26 with 34 CNOT gates and 27 single-qubit gates.

$180 e^+ e^- \rightarrow ZH \rightarrow \nu\bar{\nu} s\bar{s}$  with 6 particles



For this small-sized problem, the quantum hardware can achieve similar performance to a noiseless quantum computer simulator.

# Summary

- We proposed and developed a jet origin identification algorithm for jet flavor tagging and jet charge measurement, and achieved significant improvement in the measurement of Higgs rare and FCNC decay.
- The rapid development of quantum algorithms and hardware devices enables the execution of small-scale but representative applications on quantum computers for fundamental sciences. By mapping collision events into graphs, we have obtained promising results in applying QAOA to jet clustering.

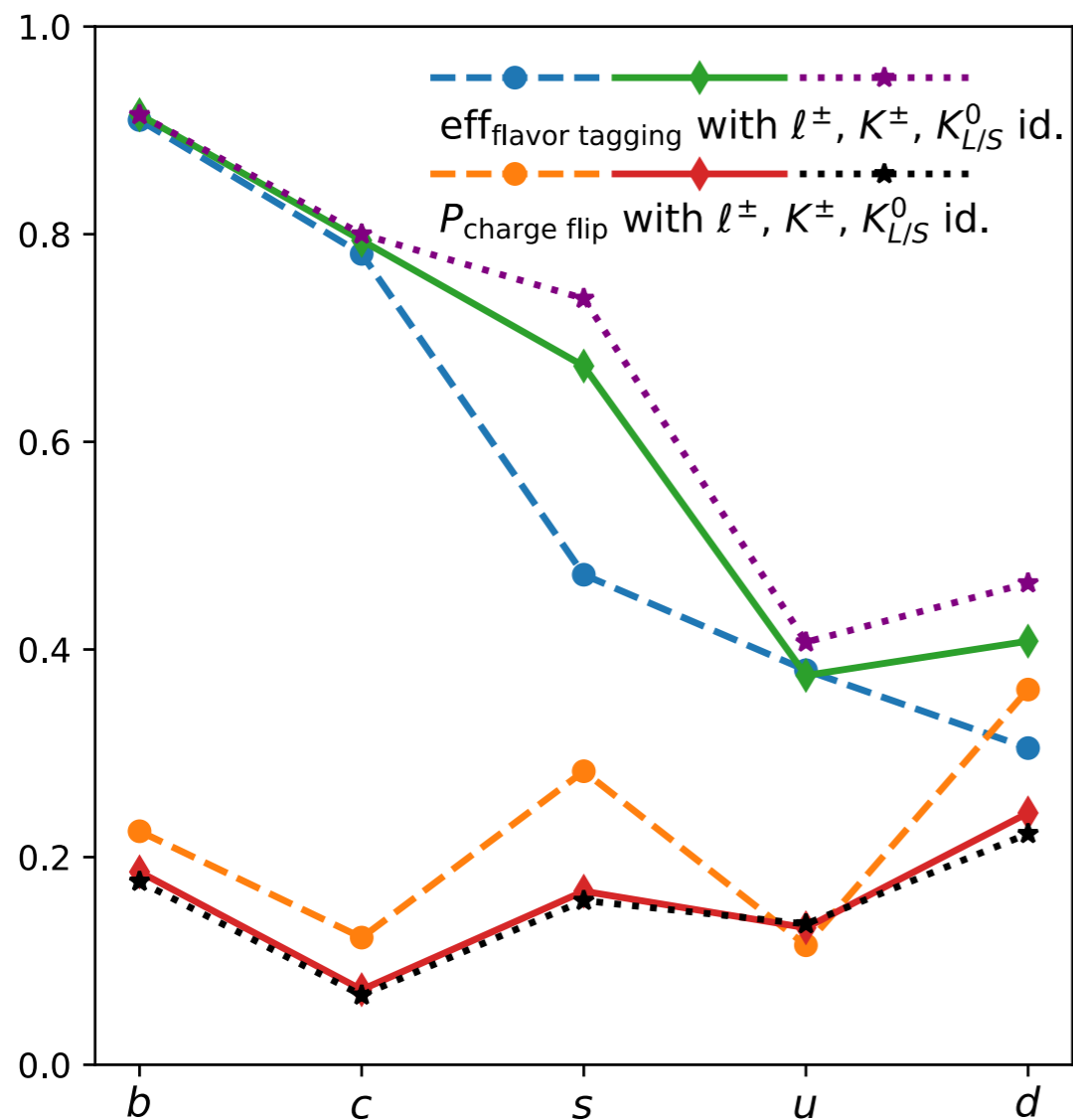
Many thanks !

Backup

# the performance of jet origin identification

**jet flavor** is defined as  $\max(b + \bar{b}, c + \bar{c}, s + \bar{s}, u + \bar{u}, d + \bar{d}, g)$

**jet charge** is assigned by comparing the quark and anti-quark likelihoods of the corresponding flavor



To understand the impact of PID, three scenarios are compared.

1, assumes perfect identification of **charged leptons** ( $\ell^\pm$ )

2, further assumes perfect identification of the **charged hadrons** ( $K^\pm$ )

3, on top of the second scenario, assumes perfect identification of  $K_L$  and  $K_S$ .

default scenario: 2 scenario, based on:

[Eur. Phys. J. C 80, 7 \(2020\)](#)

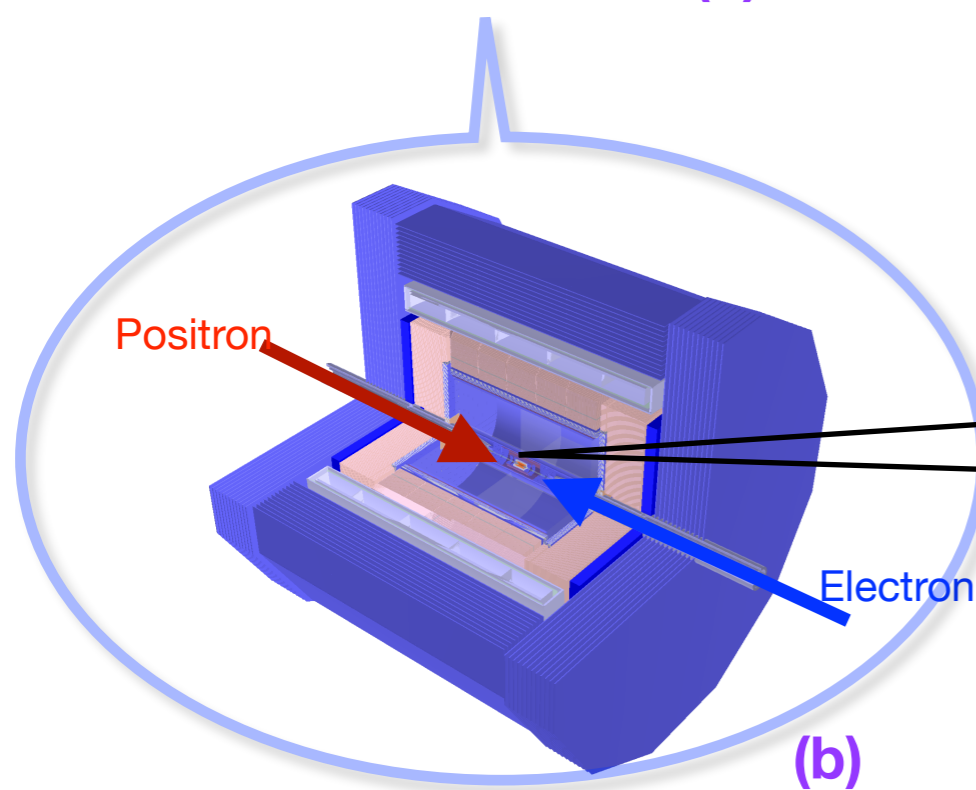
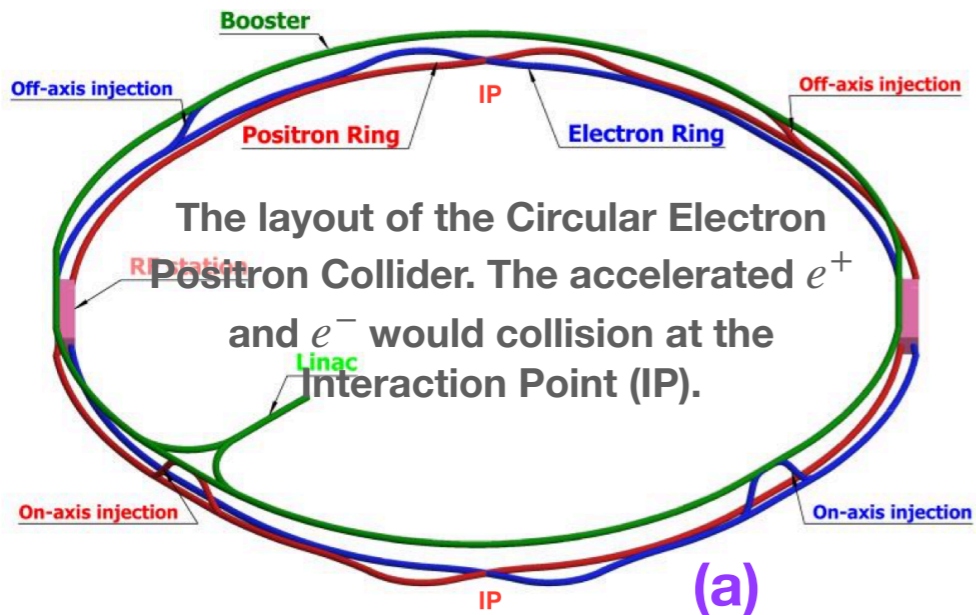
[Journal of Instrumentation 16, P06013 \(2021\)](#)

[Eur. Phys. J. C 78, 464 \(2018\)](#)

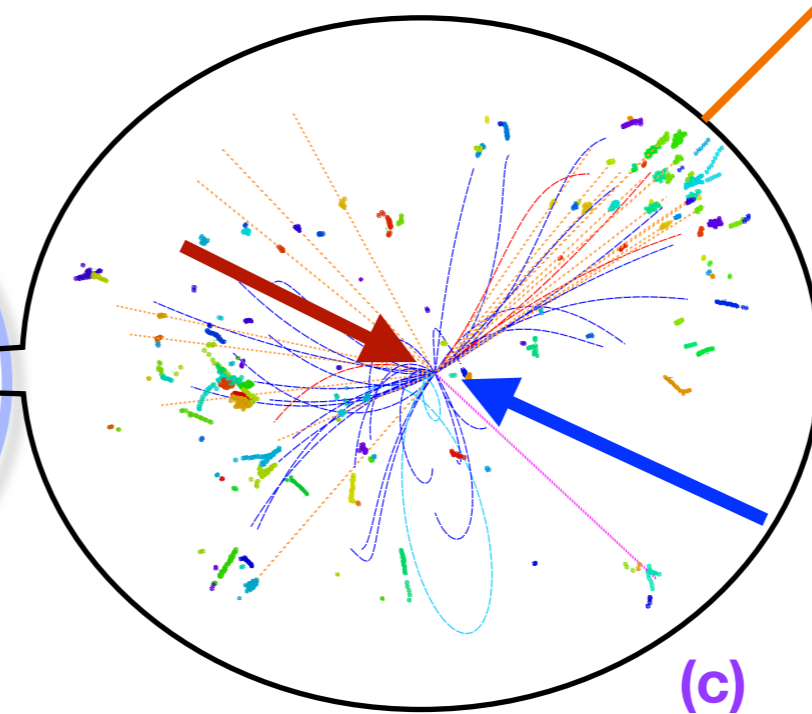
[Eur. Phys. J. C 83, 93 \(2023\)](#)

[Nucl. Instrum. Meth. A 1047, 167835 \(2023\)](#)

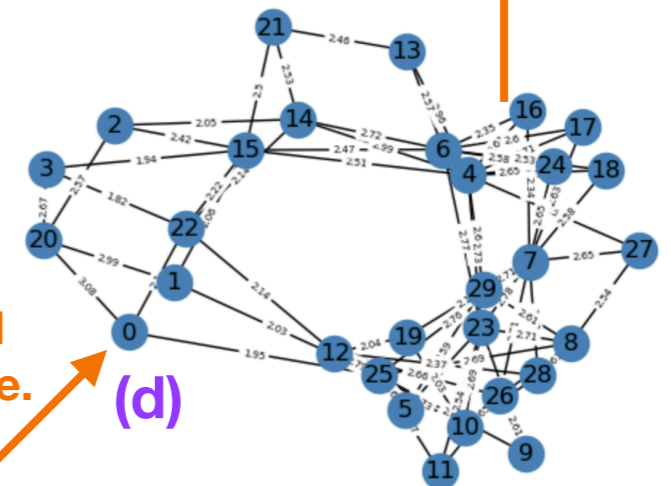




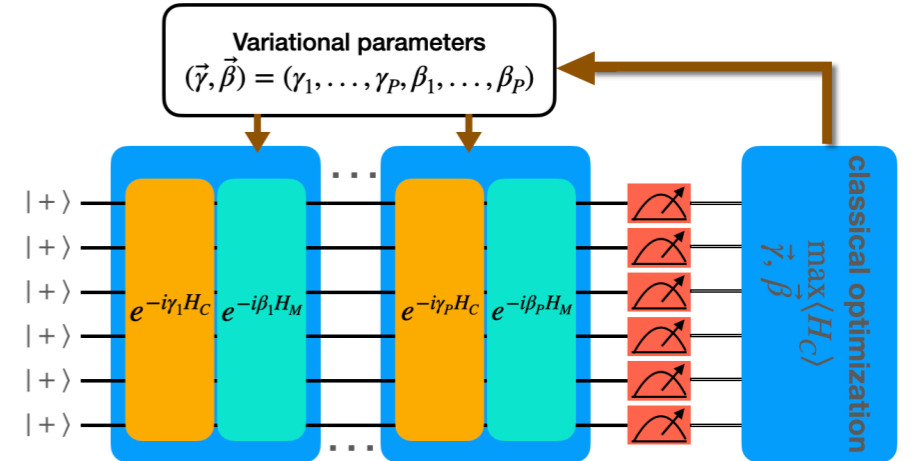
The collision of  $e^+$  and  $e^-$  can generate quarks, gluons, and leptons.



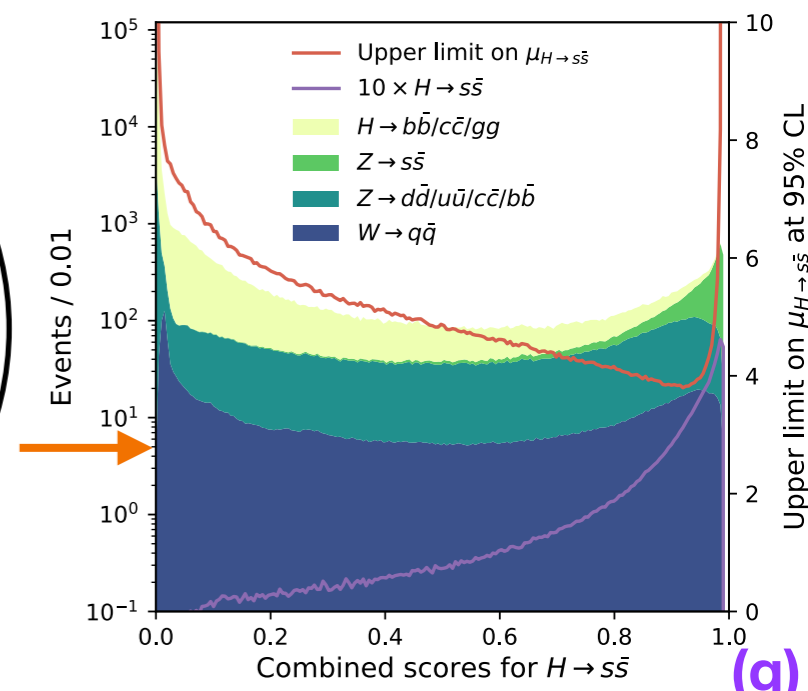
The quarks and gluons would immediately transform into collimated particle sprays known as jets.



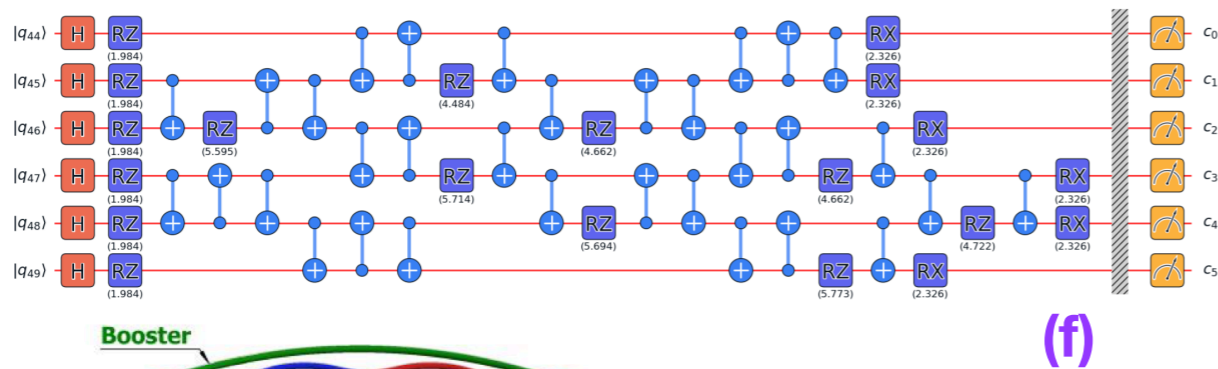
See the event as a graph, where particles as nodes and angle of two particles as edge.



jet clustering with QAOA



With jet clustering and other techniques, the related physics analyses can be performed.



compiled quantum circuit

# jet clustering performance v.s. QAOA depth (P) and k-value

