



第十四届全国粒子物理学学术会议

Semi-inclusive DIS in the Target Fragmentation Region

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Based on:

K.B. Chen, J.P. Ma and X.B. Tong, JHEP 05 (2024) 298

K.B. Chen, J.P. Ma and X.B. Tong, PRD 108, 094015 (2023)



Contents

- Introduction
- Twist-3 contributions for TFR SIDIS
- One-loop contributions for TFR SIDIS
- Summary



Contents

➤ Introduction

➤ Twist-3 contributions for TFR SIDIS

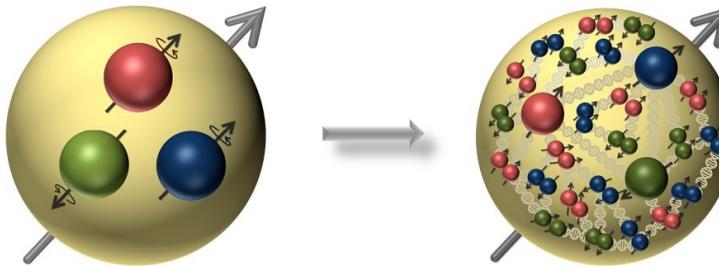
➤ One-loop contributions for TFR SIDIS

➤ Summary

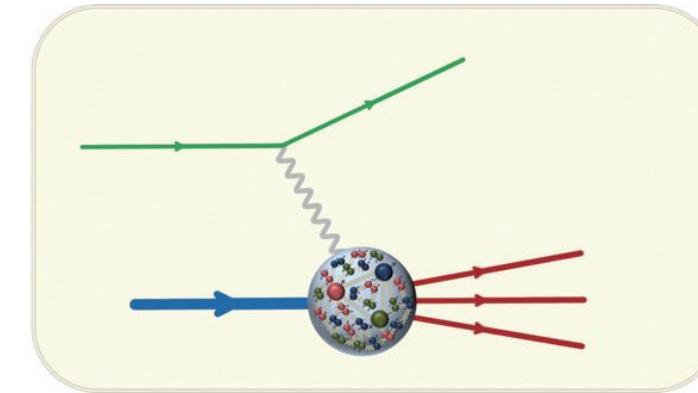


Introduction

■ Probing nucleon structure



Nucleon structure



Lepton-nucleon deep inelastic scattering (DIS)

Images taken from
Front. Phys. 16(6), 64701

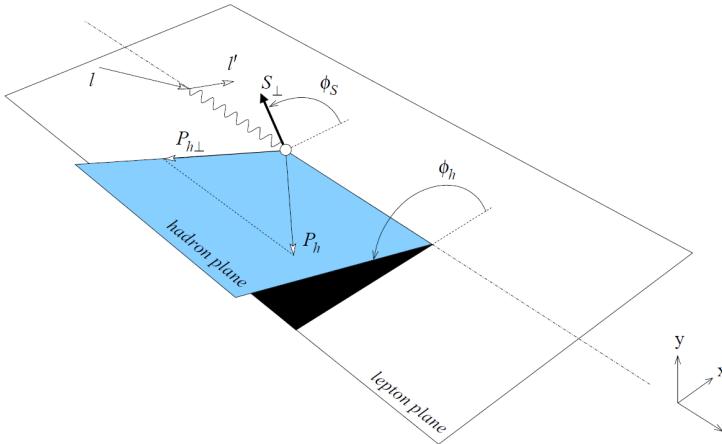
Why nucleon structure important?

- Fundamental components of the structure of matter
- Insights into the properties of strong interaction
- Inputs for describing high energy reactions
- Prerequisite for doing new physics researches
-



Introduction

■ Semi-inclusive DIS (SIDIS)



A. Bacchetta et al., JHEP 02, 093 (2007)

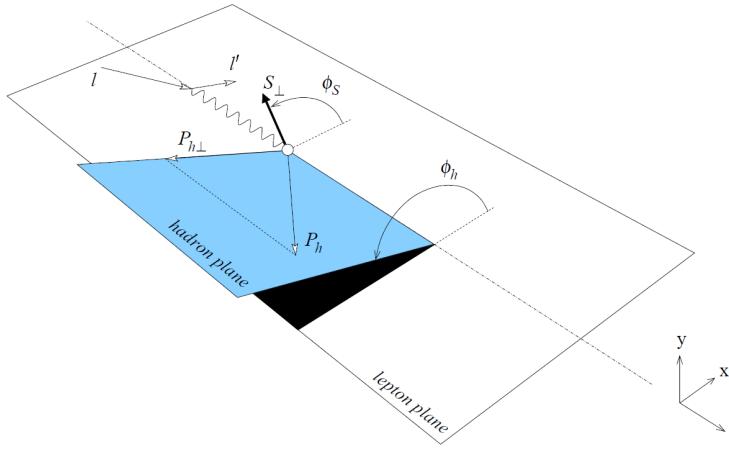
18 structure functions for SIDIS
with polarized lepton and nucleon

$$\begin{aligned} \frac{d\sigma}{dx dy d\psi dz d\phi_h dP_{h\perp}^2} = & \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\varepsilon)} \left(1 + \frac{\gamma^2}{2x} \right) \\ & \left\{ F_{UU,T} + \varepsilon F_{UU,L} + \sqrt{2\varepsilon(1+\varepsilon)} \cos \phi_h F_{UU}^{\cos \phi_h} \right. \\ & + \varepsilon \cos(2\phi_h) F_{UU}^{\cos 2\phi_h} + \lambda_e \sqrt{2\varepsilon(1-\varepsilon)} \sin \phi_h F_{LU}^{\sin \phi_h} \\ & + S_{\parallel} \left[\sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_h F_{UL}^{\sin \phi_h} + \varepsilon \sin(2\phi_h) F_{UL}^{\sin 2\phi_h} \right] \\ & + S_{\parallel} \lambda_e \left[\sqrt{1-\varepsilon^2} F_{LL} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_h F_{LL}^{\cos \phi_h} \right] \\ & + |S_{\perp}| \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \varepsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right. \\ & + \varepsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \varepsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)} \\ & + \sqrt{2\varepsilon(1+\varepsilon)} \sin \phi_S F_{UT}^{\sin \phi_S} + \sqrt{2\varepsilon(1+\varepsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \left. \right] \\ & + |S_{\perp}| \lambda_e \left[\sqrt{1-\varepsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\varepsilon(1-\varepsilon)} \cos \phi_S F_{LT}^{\cos \phi_S} \right. \\ & \left. \left. + \sqrt{2\varepsilon(1-\varepsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\} \end{aligned}$$



Introduction

■ Semi-inclusive DIS (SIDIS)



A. Bacchetta et al., JHEP 02, 093 (2007)

$$P_{h\perp} \sim M \ll Q$$

TMD factorization applies.

Structure functions are expressed by the convolution of TMD PDFs and FFs.

$$F_{UU,T} \sim f_1 \otimes D_1$$

Number density

$$F_{LL} \sim g_{1L} \otimes D_1$$

Helicity

$$F_{UT,T}^{\sin(\phi_h - \phi_S)} \sim f_{1T}^\perp \otimes D_1$$

Sivers

$$F_{LT}^{\cos(\phi_h - \phi_S)} \sim g_{1T} \otimes D_1$$

Worm-gear L-T

$$F_{UT}^{\sin(\phi_h + \phi_S)} \sim h_1 \otimes H_1^\perp$$

Transversity

$$F_{UU}^{\cos 2\phi_h} \sim h_1^\perp \otimes H_1^\perp$$

Boer-Mulders

$$F_{UL}^{\sin 2\phi_h} \sim h_{1L}^\perp \otimes H_1^\perp$$

Worm-gear T-L

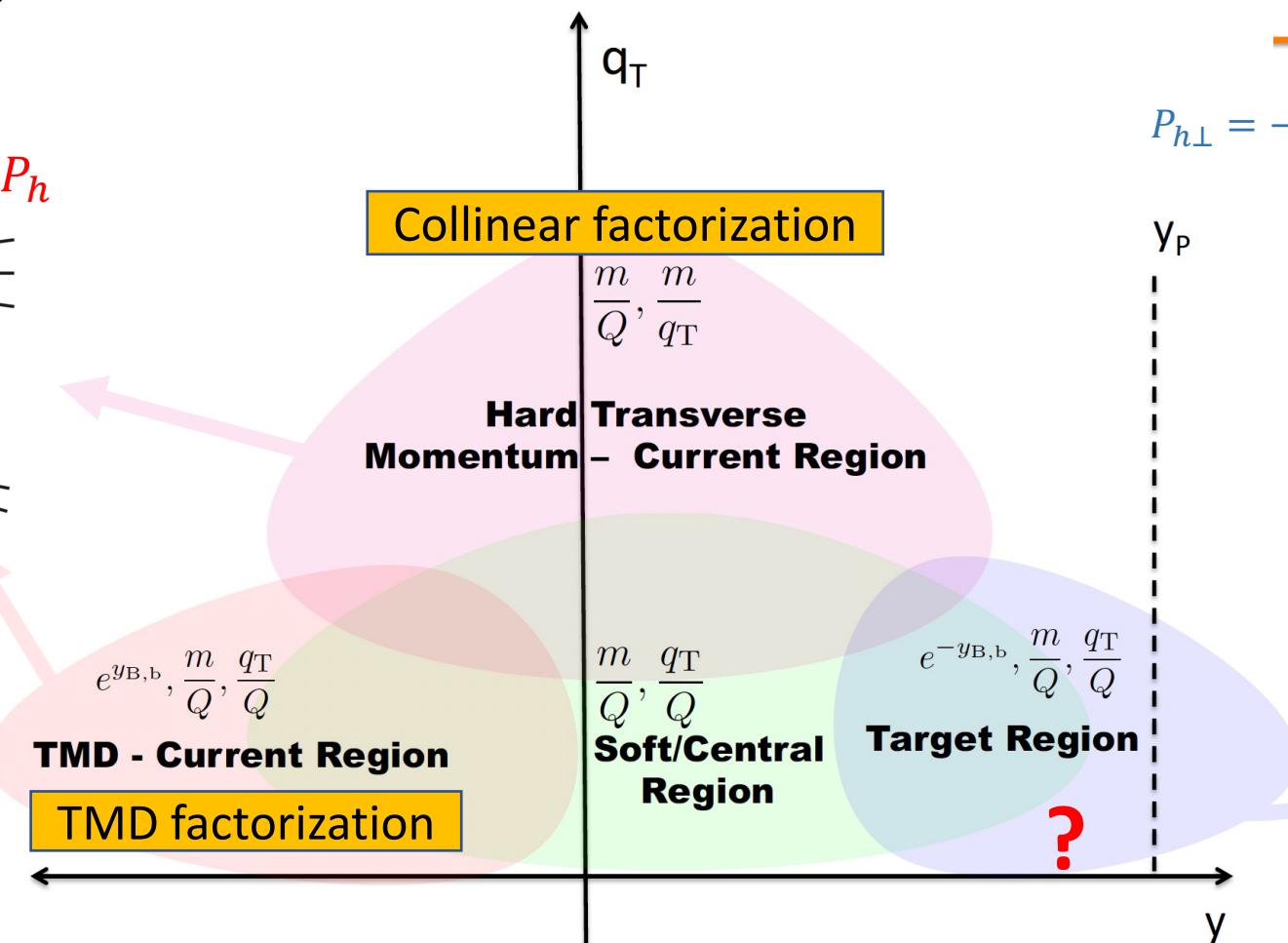
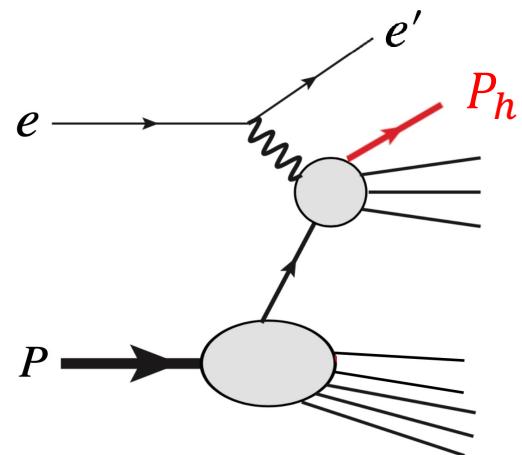
$$F_{UT}^{\sin(3\phi_h - \phi_S)} \sim h_{1T}^\perp \otimes H_1^\perp$$

Pretzelosity

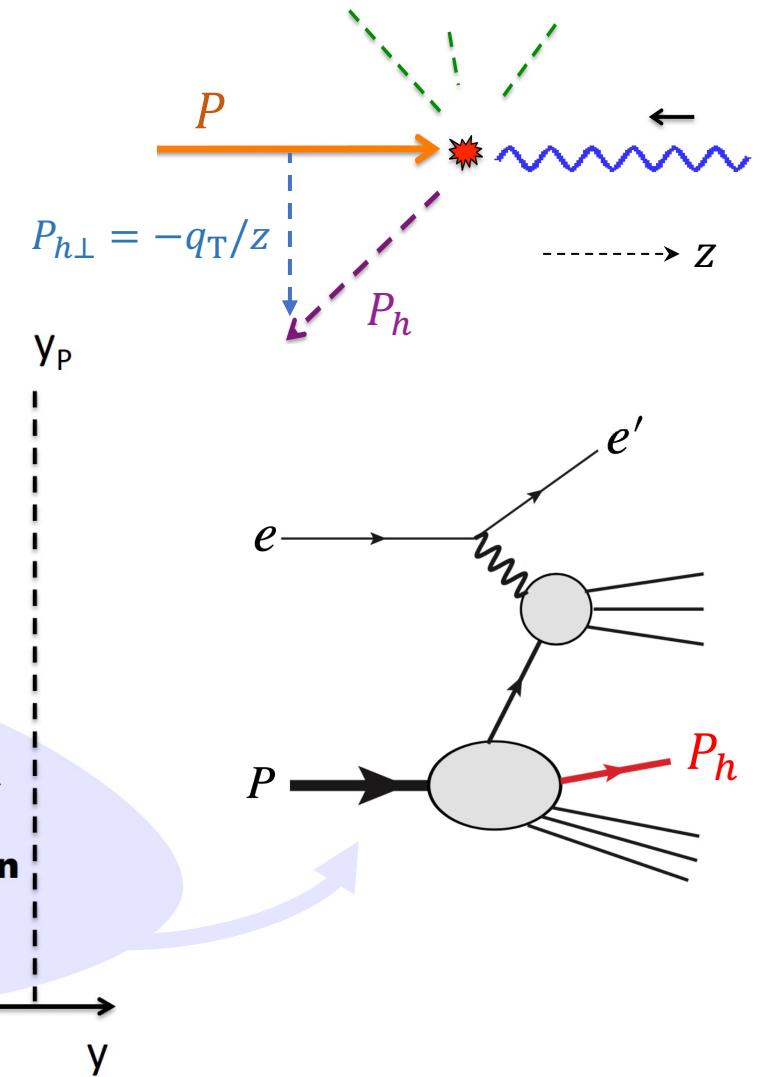
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Introduction

■ Kinematic regions of SIDIS



Mapping kinematical regions of SIDIS in terms of the produced hadron's **rapidity** and **transverse momentum** in the Breit frame



M. Boglione et al. JHEP 10 (2019) 122



Introduction

■ SIDIS in the target fragmentation region (TFR)

- TFR SIDIS events were found at HERA.

M. Derrick et al. (ZEUS Collaboration), Phys. Lett. B346, 399 (1995).

T. Ahmed et al. (H1 Collaboration), Phys. Lett. B348, 681 (1995).

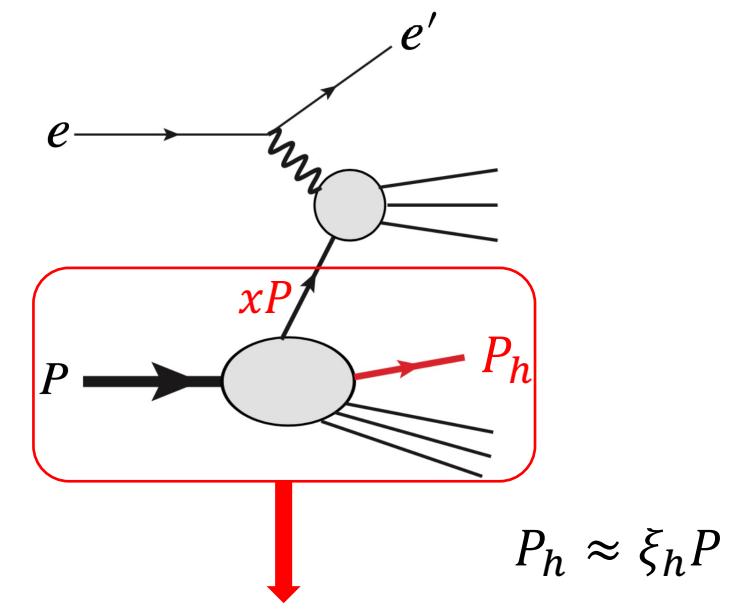
- Theoretical description: the concept of fracture function

$F(x, \xi_h)$: Fracture function

Parton distributions in the presence of an almost collinear particle observed in the final state.

(conditional probability, hybrid of PDFs and FFs)

L. Trentadue and G. Veneziano, Phys. Lett. B323, 201 (1994)





Introduction

■ SIDIS in the target fragmentation region (TFR)

- Structure functions in the TFR

$$\frac{d\sigma^{\text{TFR}}}{dx_B dy d\xi d^2 \mathbf{P}_{h\perp} d\phi_S} = \frac{2\alpha_{\text{em}}^2}{Q^2 y} \left\{ \left(1 - y + \frac{y^2}{2} \right) \sum_a e_a^2 \left[M(x_B, \xi, \mathbf{P}_{h\perp}^2) - |\mathbf{S}_\perp| \frac{|\mathbf{P}_{h\perp}|}{m_h} M_T^h(x_B, \xi, \mathbf{P}_{h\perp}^2) \sin(\phi_h - \phi_S) \right] \right.$$
$$\left. + \lambda_L y \left(1 - \frac{y}{2} \right) \sum_a e_a^2 \left[S_\parallel \Delta M_L(x_B, \xi, \mathbf{P}_{h\perp}^2) + |\mathbf{S}_\perp| \frac{|\mathbf{P}_{h\perp}|}{m_h} \Delta M_T^h(x_B, \xi, \mathbf{P}_{h\perp}^2) \cos(\phi_h - \phi_S) \right] \right\}.$$

$F_{UU,T}$ $F_{UT,T}^{\sin(\phi_h - \phi_S)}$
 F_{LL} $F_{LT}^{\cos(\phi_h - \phi_S)}$

Can we go beyond twist-2 or tree level? For:

1. More accurate description of the process
2. Novel azimuthal/spin asymmetries

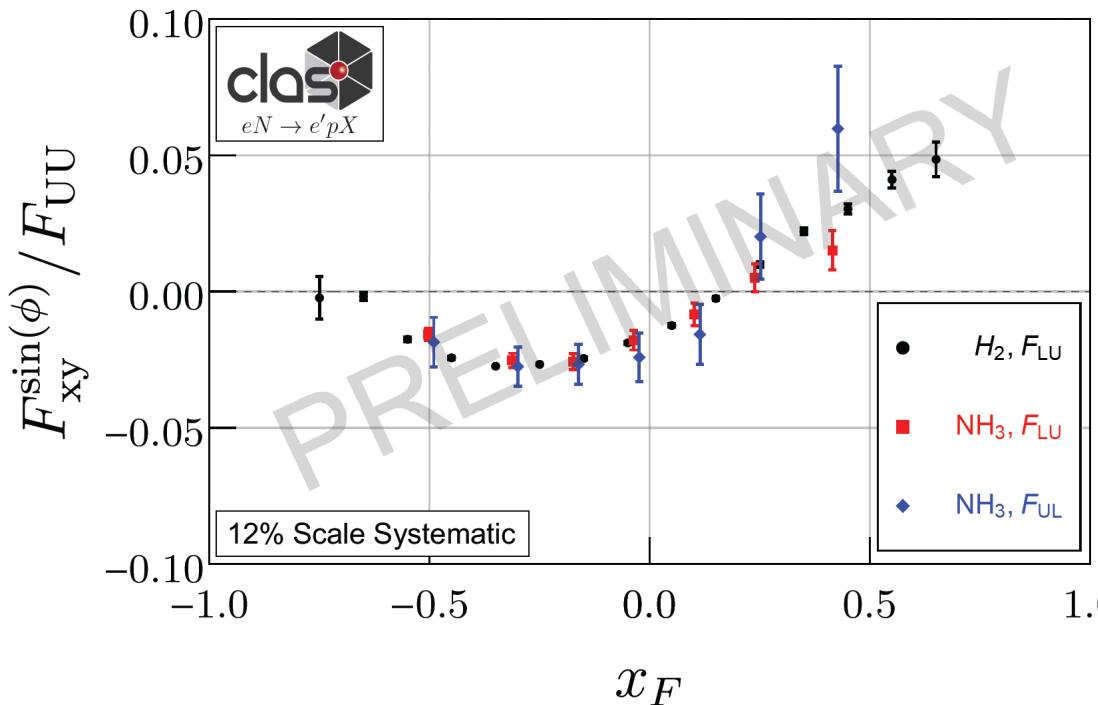
Only four structure functions at twist-2 and the tree level.

*Anselmino, Barone, and Kotzinian,
Phys. Lett. B699, 108 (2011);
Phys. Lett. B706, 46 (2011).*

Introduction

■ SIDIS in the target fragmentation region (TFR)

- Preliminary results from JLab



A. Accardi et al., arXiv:2306.09360.

CLAS12 beam-spin asymmetry results
show different behavior in CFR and TFR !

Sensitive to $F_{LU}^{\sin \phi}, F_{UL}^{\sin \phi}$.

Structure functions beyond twist-2 !



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Twist-3 contributions for TFR SIDIS

■ Kinematics

$$e(l, \lambda_e) + h_A(P, S) \rightarrow e(l') + h(P_h) + X$$

$$\frac{d\sigma}{dx_B dy d\xi_h d\psi d^2 P_{h\perp}} = \frac{\alpha^2 y}{4\xi_h Q^4} L_{\mu\nu}(l, \lambda_e, l') W^{\mu\nu}(q, P, S, P_h)$$

$$L^{\mu\nu}(l, \lambda_e, l') = 2(l^\mu l'^\nu + l^\nu l'^\mu - l \cdot l' g^{\mu\nu}) + 2i\lambda_e \epsilon^{\mu\nu\rho\sigma} l_\rho l'_\sigma$$

$$W^{\mu\nu}(q, P, S, P_h) = \sum_X \int \frac{d^4 x}{(2\pi)^4} e^{iq \cdot x} \langle S; h_A | J^\mu(x) | hX \rangle \langle Xh | J^\nu(0) | h_A; S \rangle$$

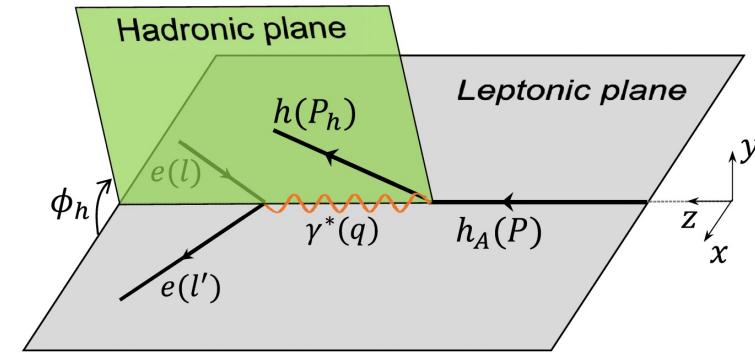
$$P^\mu \approx (P^+, 0, 0, 0)$$

$$l^\mu = \left(\frac{1-y}{y} x_B P^+, \frac{Q^2}{2x_B y P^+}, \frac{Q\sqrt{1-y}}{y}, 0 \right)$$

$$P_h^\mu = (P_h^+, P_h^-, \vec{P}_{h\perp})$$

$$q^\mu = \left(-x_B P^+, \frac{Q^2}{2x_B P^+}, 0, 0 \right)$$

$$S^\mu = \left(\frac{S_L P^+}{M}, -\frac{S_L M}{2P^+}, \vec{S}_\perp \right)$$



$$Q^2 = -q^2,$$

$$x_B = \frac{Q^2}{2P \cdot q},$$

$$y = \frac{P \cdot q}{P \cdot l},$$

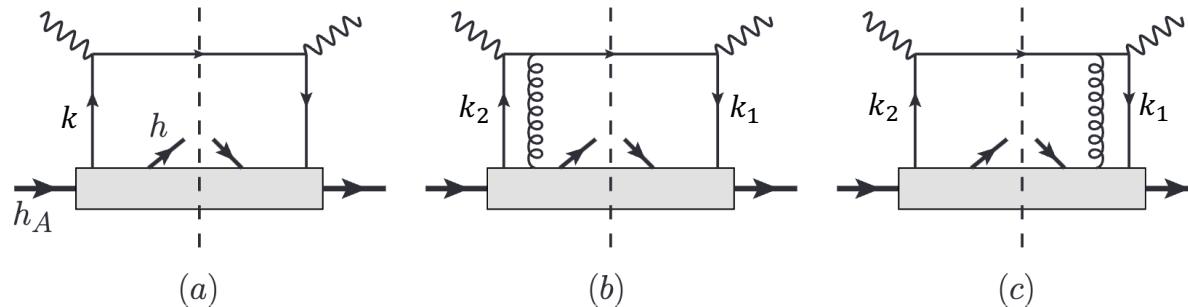
$$\xi_h = \frac{P_h \cdot q}{P \cdot q} \approx \frac{P_h^+}{P^+},$$

$$z_h = \frac{P \cdot P_h}{P \cdot q} \ll 1.$$

Twist-3 contributions for TFR SIDIS

■ Hadronic tensor

Hadronic tensor at the tree level



.....

Multiple gluon scattering in the final state.

$$W^{\mu\nu} \Big|_a = \int \frac{d^3 k}{(2\pi)^3} \left[(\gamma^\mu (\not{k} + \not{q}) \gamma^\nu)_{ij} 2\pi \delta((k+q)^2) \right] \sum_X \int \frac{d^3 \eta}{(2\pi)^4} e^{-ik \cdot \eta} \langle h_A | \bar{\psi}_i(\eta) | hX \rangle \langle Xh | \psi_j(0) | h_A \rangle,$$

$$W^{\mu\nu} \Big|_b = \int \frac{d^3 k_1 d^3 k_2}{(2\pi)^6} \left[\left(\gamma^\mu (\not{k}_1 + \not{q}) \gamma_\alpha \frac{i(\not{k}_2 + \not{q})}{(k_2 + q)^2 + i\epsilon} \gamma^\nu \right)_{ij} 2\pi \delta((k_1 + q)^2) \right] \\ \times (-ig_s) \sum_X \int \frac{d^3 \eta d^3 \eta_1}{(2\pi)^4} e^{-ik_1 \cdot \eta} e^{i(k_1 - k_2) \cdot \eta_1} \langle h_A | \bar{\psi}_i(\eta) | hX \rangle \langle Xh | G^\alpha(\eta_1) \psi_j(0) | h_A \rangle,$$

$$W^{\mu\nu} \Big|_c = \int \frac{d^3 k_1 d^3 k_2}{(2\pi)^6} \left[\left(\gamma^\mu \frac{i(\not{k}_1 + \not{q})}{(k_1 + q)^2 - i\epsilon} \gamma_\alpha (\not{k}_2 + \not{q}) \gamma^\nu \right)_{ij} 2\pi \delta((k_2 + q)^2) \right] \\ \times (-ig_s) \sum_X \int \frac{d^3 \eta d^3 \eta_1}{(2\pi)^4} e^{-ik_1 \cdot \eta} e^{i(k_1 - k_2) \cdot \eta_1} \langle h_A | \bar{\psi}_i(\eta) G^\alpha(\eta_1) | hX \rangle \langle Xh | \psi_j(0) | h_A \rangle,$$

$$k_i \sim Q(1, \lambda^2, \lambda), \text{ with } \lambda \sim \frac{\Lambda_{QCD}}{Q}$$

Carrying out the collinear expansion procedures



Twist-3 contributions for TFR SIDIS

■ Hadronic tensor

$$W^{\mu\nu} = \xi_h (\gamma^\mu \gamma^+ \gamma^\nu)_{ij} \mathcal{M}_{ji}(x_B) + \left[\frac{-i\xi_h}{2q^-} (\gamma^\mu \gamma^+ \gamma_{\perp\alpha} \gamma^- \gamma^\nu)_{ij} \mathcal{M}_{\partial,ji}^\alpha(x_B) + (\mu \leftrightarrow \nu)^* \right] \\ + \left\{ \frac{-i\xi_h}{2q^-} (\gamma^\mu \gamma^+ \gamma_{\perp\alpha} \gamma^- \gamma^\nu)_{ij} \int dx_2 \left[P \frac{1}{x_2 - x_B} - i\pi\delta(x_2 - x_B) \right] \mathcal{M}_{F,ji}^\alpha(x_B, x_2) + (\mu \leftrightarrow \nu)^* \right\}$$

twist-2 *twist-3*

$$\mathcal{M}_{ij}(x) = \int \frac{d\eta^-}{2\xi_h(2\pi)^4} e^{-ixP^+\eta^-} \sum_X \langle h_A | \bar{\psi}_j(\eta^-) \mathcal{L}_n^\dagger(\eta^-) | hX \rangle \langle Xh | \mathcal{L}_n(0) \psi_i(0) | h_A \rangle,$$

$$\mathcal{M}_{\partial,ij}^\alpha(x) = \int \frac{d\eta^-}{2\xi_h(2\pi)^4} e^{-ixP^+\eta^-} \sum_X \langle h_A | \bar{\psi}_j(\eta^-) \mathcal{L}_n^\dagger(\eta^-) | hX \rangle \langle Xh | \partial_\perp^\alpha (\mathcal{L}_n \psi_i)(0) | h_A \rangle,$$

$$\mathcal{M}_{F,ij}^\alpha(x_1, x_2) = \int \frac{d\eta^- d\eta_1^-}{4\pi\xi_h(2\pi)^4} e^{-ix_1 P^+ \eta^- - i(x_2 - x_1) P^+ \eta_1^-} \sum_X \langle h_A | \bar{\psi}_j(\eta^-) | hX \rangle \langle Xh | g_s F^{+\alpha}(\eta_1^-) \psi_i(0) | h_A \rangle.$$

Gauge invariant matrix elements!

$$\mathcal{L}_n(x) = \mathcal{P} \exp \left\{ -ig_s \int_0^\infty d\lambda G^+(\lambda n + x) \right\}$$

$$g_s F^{+\alpha} = g_s [\partial^+ G_\perp^\alpha - \partial_\perp^\alpha G^+] + \mathcal{O}(g_s^2)$$



Twist-3 contributions for TFR SIDIS

■ Fracture functions defined via the correlation matrices

$$\begin{aligned}\mathcal{M}_{ij}(x) = & \frac{(\gamma_\rho)_{ij}}{2N_c} \left[\bar{n}^\rho \left(\textcolor{red}{u_1} - \frac{P_{h\perp} \cdot \tilde{S}_\perp}{M} \textcolor{red}{u_{1T}^h} \right) + \frac{1}{P^+} \left(P_{h\perp}^\rho \textcolor{green}{u^h} - M \tilde{S}_\perp^\rho \textcolor{green}{u_T^h} - S_L \tilde{P}_{h\perp}^\rho \textcolor{green}{u_L^h} - \frac{P_{h\perp}^{\langle\rho} P_{h\perp}^{\beta\rangle}}{M} \tilde{S}_{\perp\beta} \textcolor{green}{u_T^h} \right) \right] \\ & - \frac{(\gamma_\rho \gamma_5)_{ij}}{2N_c} \left[\bar{n}^\rho \left(S_L \textcolor{red}{l_{1L}} - \frac{P_{h\perp} \cdot S_\perp}{M} \textcolor{red}{l_{1T}^h} \right) + \frac{1}{P^+} \left(\tilde{P}_{h\perp}^\rho \textcolor{green}{l^h} + M S_\perp^\rho \textcolor{green}{l_T^h} + S_L P_{h\perp}^\rho \textcolor{green}{l_L^h} - \frac{P_{h\perp}^{\langle\rho} P_{h\perp}^{\beta\rangle}}{M} S_{\perp\beta} \textcolor{green}{l_T^h} \right) \right] + \dots,\end{aligned}$$

$$\begin{aligned}\mathcal{M}_{\partial,ij}^\alpha(x) = & \frac{(\gamma^-)_{ij}}{2N_c} i \left(-P_{h\perp}^\alpha \textcolor{green}{u_\partial^h} + M \tilde{S}_\perp^\alpha \textcolor{green}{u_{\partial T}^h} + S_L \tilde{P}_{h\perp}^\alpha \textcolor{green}{u_{\partial L}^h} + \frac{P_{h\perp}^{\langle\alpha} P_{h\perp}^{\beta\rangle}}{M} \tilde{S}_{\perp\beta} \textcolor{green}{u_{\partial T}^h} \right) \\ & + \frac{(\gamma^- \gamma_5)_{ij}}{2N_c} i \left(\tilde{P}_{h\perp}^\alpha \textcolor{green}{l_\partial^h} + M S_\perp^\alpha \textcolor{green}{l_{\partial T}^h} + S_L P_{h\perp}^\alpha \textcolor{green}{l_{\partial L}^h} - \frac{P_{h\perp}^{\langle\alpha} P_{h\perp}^{\beta\rangle}}{M} S_{\perp\beta} \textcolor{green}{l_{\partial T}^h} \right) + \dots,\end{aligned}$$

$$\begin{aligned}\mathcal{M}_{F,ij}^\alpha(x_1, x_2) = & \frac{(\gamma^-)_{ij}}{2N_c} \left(P_{h\perp}^\alpha \textcolor{green}{w^h} - M \tilde{S}_\perp^\alpha \textcolor{green}{w_T^h} - S_L \tilde{P}_{h\perp}^\alpha \textcolor{green}{w_L^h} - \frac{P_{h\perp}^{\langle\alpha} P_{h\perp}^{\beta\rangle}}{M} \tilde{S}_{\perp\beta} \textcolor{green}{w_T^h} \right) \\ & - \frac{(\gamma^- \gamma_5)_{ij}}{2N_c} i \left(\tilde{P}_{h\perp}^\alpha \textcolor{green}{v^h} + M S_\perp^\alpha \textcolor{green}{v_T^h} + S_L P_{h\perp}^\alpha \textcolor{green}{v_L^h} - \frac{P_{h\perp}^{\langle\alpha} P_{h\perp}^{\beta\rangle}}{M} S_{\perp\beta} \textcolor{green}{v_T^h} \right) + \dots,\end{aligned}$$

Fracture functions
up to twist-3

$$\tilde{a}_\perp^\mu \equiv \epsilon_\perp^{\mu\nu} a_{\perp\nu}$$

$$P_{h\perp}^{\langle\alpha} P_{h\perp}^{\beta\rangle} \equiv P_{h\perp}^\alpha P_{h\perp}^\beta + g_\perp^{\alpha\beta} \vec{P}_{h\perp}^2 / 2$$

- Red: twist-2.
 - Green: twist-3
 - Functions of $(x, \xi_h, P_{h\perp}^2)$
 - Only chiral-even terms contribute
 - Similar form ($k_\perp \rightarrow P_{h\perp}$) and naming rules as TMDs
- See e.g., S.Y. Wei, Y.K. Song, KBC, Z.T. Liang, PRD 95, 074017*



Twist-3 contributions for TFR SIDIS

■ Fracture functions defined via the correlation matrices

The twist-3 fracture functions are **not independent** from each other!

Applying QCD equation of motion $i\gamma \cdot D\psi = 0$, gives:

$$\begin{aligned} & x[u_S^K(x) + i l_S^K(x)] \\ &= u_{\partial S}^K(x) + i l_{\partial S}^K(x) + i \int dy \left[P \frac{1}{y-x} - i\pi\delta(y-x) \right] [w_S^K(x,y) - v_S^K(x,y)] \end{aligned}$$

Four sets of equations

The relations take a **unified form**!

The hadronic tensor will be expressed by those fracture functions defined **only from the quark-quark correlator \mathcal{M}_{ij}** .

S	K
null	h
L	h
T	null
T	h



Twist-3 contributions for TFR SIDIS

■ Results for structure functions and azimuthal asymmetries

Structure functions

$$F_{UU,T} = x_B u_1, \quad F_{UT,T}^{\sin(\phi_h - \phi_S)} = \frac{|\vec{P}_{h\perp}|}{M} x_B u_{1T}^h,$$

$$F_{LL} = x_B l_{1L}, \quad F_{LT}^{\cos(\phi_h - \phi_S)} = \frac{|\vec{P}_{h\perp}|}{M} x_B l_{1T}^h.$$

Twist-2

Azimuthal/spin asymmetries

$$\langle \sin(\phi_h - \phi_S) \rangle_{UT} = \frac{|\vec{P}_{h\perp}|}{2M} \frac{u_{1T}^h}{u_1},$$

$$\langle \cos(\phi_h - \phi_S) \rangle_{LT} = \frac{|\vec{P}_{h\perp}| C(y)}{2MA(y)} \frac{l_{1T}^h}{u_1}.$$

$$F_{UU}^{\cos \phi_h} = -\frac{2|\vec{P}_{h\perp}|}{Q} x_B^2 u^h, \quad F_{LU}^{\sin \phi_h} = \frac{2|\vec{P}_{h\perp}|}{Q} x_B^2 l^h,$$

$$F_{UL}^{\sin \phi_h} = -\frac{2|\vec{P}_{h\perp}|}{Q} x_B^2 u_L^h, \quad F_{LL}^{\cos \phi_h} = -\frac{2|\vec{P}_{h\perp}|}{Q} x_B^2 l_L^h,$$

$$F_{UT}^{\sin \phi_S} = -\frac{2M}{Q} x_B^2 u_T, \quad F_{LT}^{\cos \phi_S} = -\frac{2M}{Q} x_B^2 l_T,$$

$$F_{UT}^{\sin(2\phi_h - \phi_S)} = -\frac{\vec{P}_{h\perp}^2}{QM} x_B^2 u_T^h, \quad F_{LT}^{\cos(2\phi_h - \phi_S)} = -\frac{\vec{P}_{h\perp}^2}{QM} x_B^2 l_T^h.$$

Twist-3

$$\langle \cos \phi_h \rangle_{UU} = -\frac{|\vec{P}_{h\perp}|}{Q} \frac{B(y)}{A(y)} \frac{x_B u^h}{u_1}, \quad \langle \sin \phi_h \rangle_{LU} = \frac{|\vec{P}_{h\perp}|}{Q} \frac{D(y)}{A(y)} \frac{x_B l^h}{u_1},$$

$$\langle \sin \phi_h \rangle_{UL} = -\frac{|\vec{P}_{h\perp}|}{Q} \frac{B(y)}{A(y)} \frac{x_B u_L^h}{u_1}, \quad \langle \cos \phi_h \rangle_{LL} = -\frac{|\vec{P}_{h\perp}|}{Q} \frac{D(y)}{A(y)} \frac{x_B l_L^h}{u_1},$$

$$\langle \sin \phi_S \rangle_{UT} = -\frac{M}{Q} \frac{B(y)}{A(y)} \frac{x_B u_T}{u_1}, \quad \langle \cos \phi_S \rangle_{LT} = -\frac{M}{Q} \frac{D(y)}{A(y)} \frac{x_B l_T}{u_1},$$

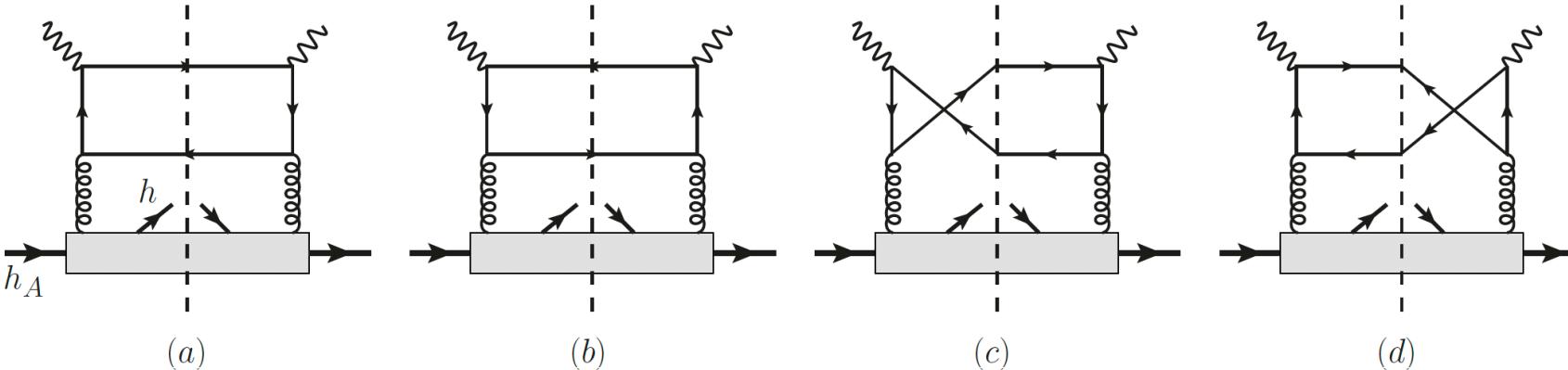
$$\langle \sin(2\phi_h - \phi_S) \rangle_{UT} = -\frac{\vec{P}_{h\perp}^2}{2MQ} \frac{B(y)}{A(y)} \frac{x_B u_T^h}{u_1}, \quad \langle \cos(2\phi_h - \phi_S) \rangle_{LT} = -\frac{\vec{P}_{h\perp}^2}{2MQ} \frac{D(y)}{A(y)} \frac{x_B l_T^h}{u_1}.$$



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■ Gluonic contribution



$$W^{\mu\nu}(q, P, S, P_h) = \alpha_s T_F \sum_f e_f^2 \int \frac{dx}{x} \int d\Phi_{k_1 k_2} H^{\mu\nu\alpha\beta}(k_g, k_1, k_2) \mathcal{M}_{G,\alpha\beta}(x, \xi_h, P_{h\perp}),$$

$$\begin{aligned} \mathcal{M}_G^{\alpha\beta}(x, \xi_h, P_{h\perp}) &= \frac{1}{2\xi_h(2\pi)^3} \frac{1}{xP^+} \int \frac{d\lambda}{2\pi} e^{-i\lambda xP^+} \sum_X \langle h_A(P) | (G^{+\alpha}(\lambda n) \mathcal{L}_n^\dagger(\lambda n))^a | X h(P_h) \rangle \\ &\times \langle h(P_h) X | (\mathcal{L}_n(0) G^{+\beta}(0))^a | h_A(P) \rangle, \end{aligned}$$



One-loop contributions for TFR SIDIS

■ Gluonic contribution

Gluon fracture functions

$$\begin{aligned}\mathcal{M}_G^{\alpha\beta} = & -\frac{1}{2-2\epsilon}g_{\perp}^{\alpha\beta}u_{1g} + \frac{1}{2M^2}\left(P_{h\perp}^{\alpha}P_{h\perp}^{\beta} + \frac{1}{2-2\epsilon}g_{\perp}^{\alpha\beta}P_{h\perp}^2\right)t_{1g}^h + S_L\left[i\frac{\varepsilon_{\perp}^{\alpha\beta}}{2}l_{1gL} + \frac{\tilde{P}_{h\perp}^{\{\alpha}P_{h\perp}^{\beta\}}}{4M^2}t_{1gL}^h\right] \\ & + \frac{g_{\perp}^{\alpha\beta}}{2-2\epsilon}\frac{P_{h\perp} \cdot \tilde{S}_{\perp}}{M}u_{1gT}^h + \frac{P_{h\perp} \cdot S_{\perp}}{M}\left[i\frac{\varepsilon_{\perp}^{\alpha\beta}}{2}l_{1gT}^h - \frac{\tilde{P}_{h\perp}^{\{\alpha}P_{h\perp}^{\beta\}}}{4M^2}t_{1gT}^{hh}\right] + \frac{\tilde{P}_{h\perp}^{\{\alpha}S_{\perp}^{\beta\}} + \tilde{S}_{\perp}^{\{\alpha}P_{h\perp}^{\beta\}}}{8M}t_{1gT}^h + \dots.\end{aligned}$$

The hard parts

$$H_{(a)+(b)}^{\mu\nu\alpha\beta}(k_g, k_1, k_2) = \text{Tr} \left[\not{k}_1 \gamma^{\nu} \frac{(\not{k}_g - \not{k}_2)}{(k_g - k_2)^2} \gamma^{\beta} \not{k}_2 \gamma^{\alpha} \frac{(\not{k}_g - \not{k}_2)}{(k_g - k_2)^2} \gamma^{\mu} \right] + (k_1 \leftrightarrow k_2),$$

$$H_{(c)+(d)}^{\mu\nu\alpha\beta}(k_g, k_1, k_2) = \text{Tr} \left[\not{k}_1 \gamma^{\nu} \frac{(\not{k}_g - \not{k}_2)}{(k_g - k_2)^2} \gamma^{\beta} \not{k}_2 \gamma^{\mu} \frac{(\not{k}_1 - \not{k}_g)}{(k_1 - k_g)^2} \gamma^{\alpha} \right] + (k_1 \leftrightarrow k_2).$$



One-loop contributions for TFR SIDIS

■ Gluonic contribution

$$F_{UU}^{\cos 2\phi_h} = -\frac{\alpha_s T_F}{2\pi} \frac{P_{h\perp}^2}{2M^2} x_B \sum_{q,\bar{q}} e_q^2 \int_{x_B/\xi_h}^1 \frac{dz}{z} z^2 t_{1g}^h(x_B/z, \xi_h, P_{h\perp}),$$

$$F_{UL}^{\sin 2\phi_h} = \frac{\alpha_s T_F}{2\pi} \frac{P_{h\perp}^2}{2M^2} x_B \sum_{q,\bar{q}} e_q^2 \int_{x_B/\xi_h}^1 \frac{dz}{z} z^2 t_{1gL}^h(x_B/z, \xi_h, P_{h\perp}),$$

$$F_{UT}^{\sin(3\phi_h - \phi_s)} = \frac{\alpha_s T_F}{2\pi} \frac{P_{h\perp}^3}{4M^3} x_B \sum_{q,\bar{q}} e_q^2 \int_{x_B/\xi_h}^1 \frac{dz}{z} z^2 t_{1gT}^{hh}(x_B/z, \xi_h, P_{h\perp}),$$

$$F_{UT}^{\sin(\phi_h + \phi_s)} = \frac{\alpha_s T_F}{2\pi} \frac{P_{h\perp}}{2M} x_B \sum_{q,\bar{q}} e_q^2 \int_{x_B/\xi_h}^1 \frac{dz}{z} z^2 \left[t_{1gT}^h(x_B/z, \xi_h, P_{h\perp}) + \frac{P_{h\perp}^2}{2M^2} t_{1gT}^{hh}(x_B/z, \xi_h, P_{h\perp}) \right].$$

Four structure functions generated uniquely by the gluon fracture functions!



One-loop contributions for TFR SIDIS

■ Quark and gluon contribution

$$F_{UU,L} = \frac{\alpha_s}{2\pi} x_B \sum_{q,\bar{q}} e_q^2 \int_{x_B/\xi_h}^1 \frac{dz}{z} \left[4T_F z \bar{z} u_{1g}(x_B/z, \xi_h, P_{h\perp}) + 2C_F z u_1(x_B/z, \xi_h, P_{h\perp}) \right],$$

$$F_{UT,L}^{\sin(\phi_h - \phi_S)} = \frac{\alpha_s}{2\pi} \frac{P_{h\perp}}{M} x_B \sum_{q,\bar{q}} e_q^2 \int_{x_B/\xi_h}^1 \frac{dz}{z} \left[4T_F z \bar{z} u_{1gT}^h(x_B/z, \xi_h, P_{h\perp}) + 2C_F z u_{1T}^h(x_B/z, \xi_h, P_{h\perp}) \right].$$

Twist-4 if at
the tree level

$$F_{UU,T} = x_B \sum_{q,\bar{q}} e_q^2 \int_{x_B/\xi_h}^1 \frac{dz}{z} \left[\mathcal{H}_g(z) u_{1g}(x_B/z, \xi_h, P_{h\perp}) + \mathcal{H}_q(z) u_1(x_B/z, \xi_h, P_{h\perp}) \right],$$

$$F_{UT,T}^{\sin(\phi_h - \phi_S)} = \frac{P_{h\perp}}{M} x_B \sum_{q,\bar{q}} e_q^2 \int_{x_B/\xi_h}^1 \frac{dz}{z} \left[\mathcal{H}_g(z) u_{1gT}^h(x_B/z, \xi_h, P_{h\perp}) + \mathcal{H}_q(z) u_{1T}^h(x_B/z, \xi_h, P_{h\perp}) \right],$$

$$F_{LL} = x_B \sum_{q,\bar{q}} e_q^2 \int_{x_B/\xi_h}^1 \frac{dz}{z} \left[\Delta \mathcal{H}_g(z) l_{1gL}(x_B/z, \xi_h, P_{h\perp}) + \Delta \mathcal{H}_q(z) l_{1L}(x_B/z, \xi_h, P_{h\perp}) \right],$$

$$F_{LT}^{\cos(\phi_h - \phi_S)} = \frac{P_{h\perp}}{M} x_B \sum_{q,\bar{q}} e_q^2 \int_{x_B/\xi_h}^1 \frac{dz}{z} \left[\Delta \mathcal{H}_g(z) l_{1gT}^h(x_B/z, \xi_h, P_{h\perp}) + \Delta \mathcal{H}_q(z) l_{1T}^h(x_B/z, \xi_h, P_{h\perp}) \right],$$

LO + NLO
corrections

$$\mathcal{H}_q(z) = \delta(\bar{z}) + \frac{\alpha_s}{2\pi} \left\{ P_{qq}(z) \ln \frac{Q^2}{\mu^2} + C_F \left[2 \left(\frac{\ln \bar{z}}{\bar{z}} \right)_+ - \frac{3}{2} \left(\frac{1}{\bar{z}} \right)_+ - (1+z) \ln \bar{z} - \frac{1+z^2}{\bar{z}} \ln z + 3 - \left(\frac{\pi^2}{3} + \frac{9}{2} \right) \delta(\bar{z}) \right] \right\}, \quad \mathcal{H}_g(z) = \frac{\alpha_s}{2\pi} \left[P_{qg}(z) \ln \frac{Q^2 \bar{z}}{\mu^2 z} - T_F (1-2z)^2 \right],$$

$$\Delta \mathcal{H}_q(z) = \delta(\bar{z}) + \frac{\alpha_s}{2\pi} \left\{ \Delta P_{qq}(z) \ln \frac{Q^2}{\mu^2} + C_F \left[(1+z^2) \left(\frac{\ln \bar{z}}{\bar{z}} \right)_+ - \frac{3}{2} \left(\frac{1}{\bar{z}} \right)_+ - \frac{1+z^2}{\bar{z}} \ln z + 2 + z - \left(\frac{\pi^2}{3} + \frac{9}{2} \right) \delta(\bar{z}) \right] \right\}, \quad \Delta \mathcal{H}_g(z) = \frac{\alpha_s}{2\pi} \left[\Delta P_{qg}(z) \left(\ln \frac{Q^2 \bar{z}}{\mu^2 z} - 1 \right) + 2T_F \bar{z} \right].$$



One-loop contributions for TFR SIDIS

Structure functions	Twist	Order (α_S)
$F_{UU,T}$	2	0
$F_{UU,L}$		1
$F_{UU}^{\cos \phi_h}$	3	
$F_{UU}^{\cos 2\phi_h}$		1
$F_{LU}^{\sin \phi_h}$	3	
$F_{UL}^{\sin \phi_h}$	3	
$F_{UL}^{\sin 2\phi_h}$		1
F_{LL}	2	0
$F_{LL}^{\cos \phi_h}$	3	

Structure functions	Twist	Order (α_S)
$F_{UT,T}^{\sin(\phi_h - \phi_S)}$	2	0
$F_{UT,L}^{\sin(\phi_h - \phi_S)}$		1
$F_{UT}^{\sin(\phi_h + \phi_S)}$		1
$F_{UT}^{\sin \phi_S}$	3	
$F_{UT}^{\sin(2\phi_h - \phi_S)}$	3	
$F_{UT}^{\sin(3\phi_h - \phi_S)}$		1
$F_{LT}^{\cos \phi_S}$	3	
$F_{LT}^{\cos(\phi_h - \phi_S)}$	2	0
$F_{LT}^{\cos(2\phi_h - \phi_S)}$	3	

All 18 structure functions are non-zero up to twist-3 or one-loop level.



Contents

- Introduction
- Twist-3 contributions for TFR SIDIS
- One-loop contributions for TFR SIDIS
- Summary



Summary

- TFR SIDIS is factorized with fracture functions. TFR is complementary to CFR for describing the SIDIS process and studying the nucleon structure.
- We calculate the TFR SIDIS up to twist-3 at the tree level of pQCD. Structure functions and azimuthal asymmetries are given using the gauge-invariant fracture functions.
- By adding one-loop contributions at twist-2, all 18 structure functions for TFR SIDIS are non-zero.

谢谢！