

Collision Energy Dependence of Hypertriton Production in Au+Au Collisions at RHIC

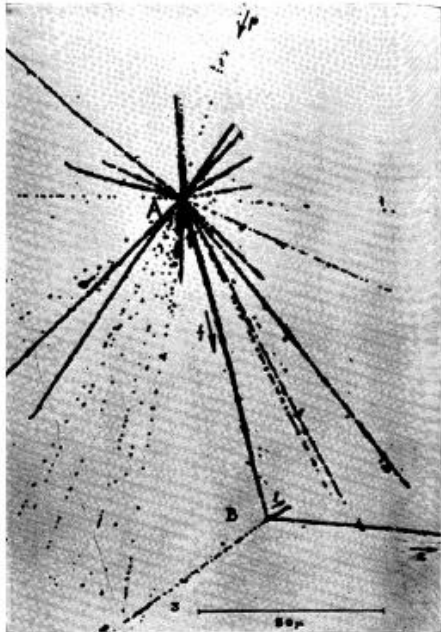
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Aug 16, 2024

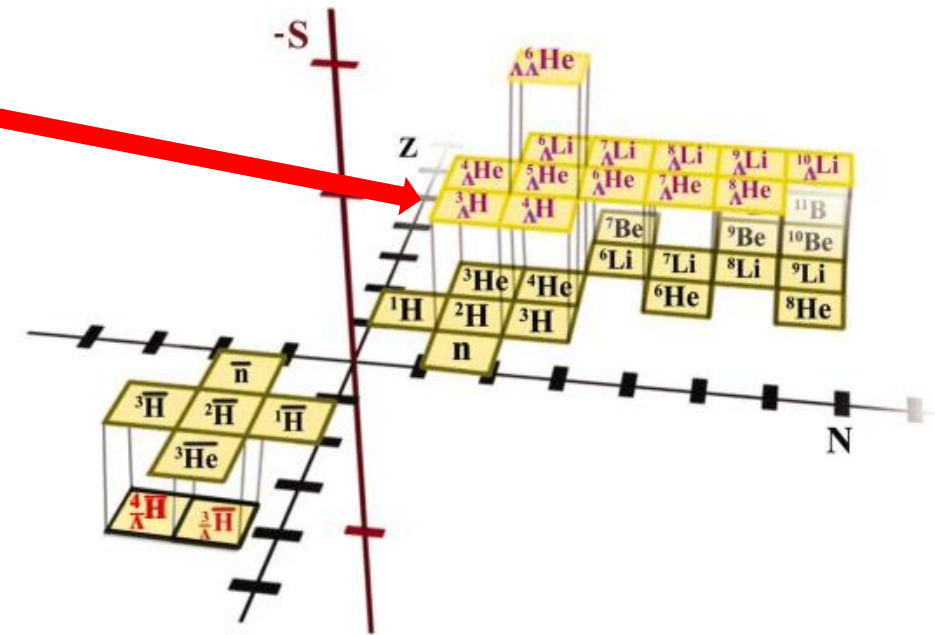
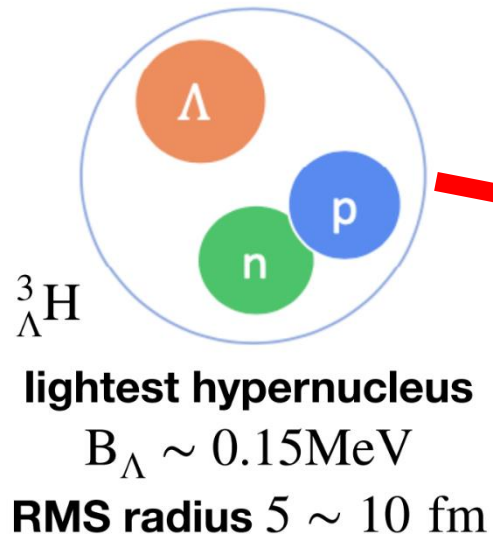
Introduction: hypernuclei

- Hypernuclei : bound nuclear systems of non-strange and strange baryons
 - Natural hyperon-baryon correlation system



The first discovery of hypernucleus by Marian Danysz (right) and Jerzy Pniewski (left) in 1952

M. Danysz, J. Pniewski, Philos. Mag. 44 (1953) 348.



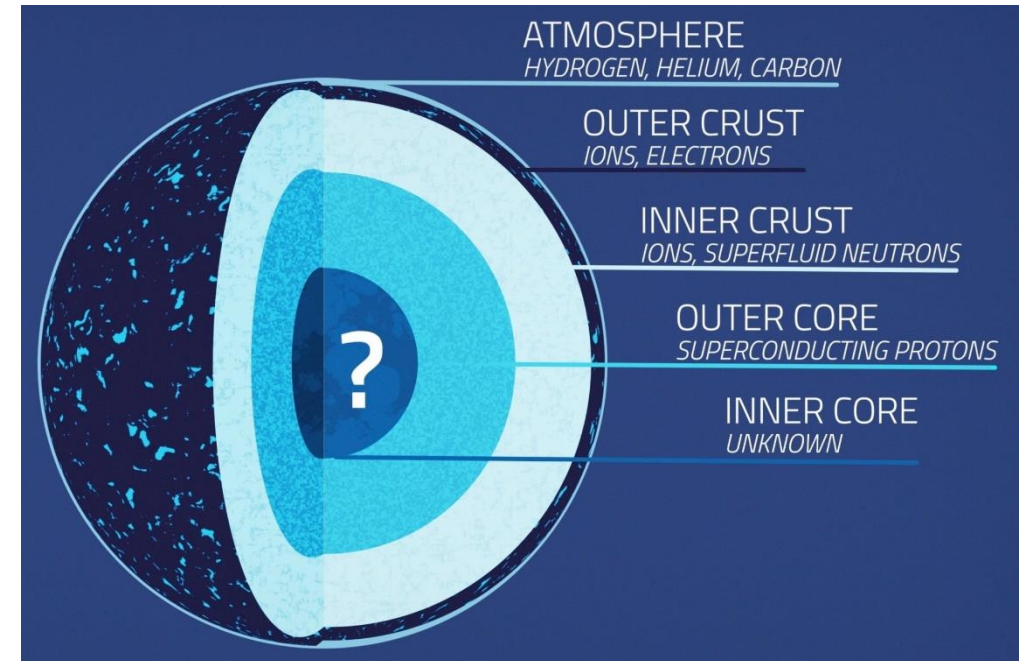
Introduction: YN interaction in dense matter

- Hypernuclei serve as a laboratory to study **the hyperon–nucleon (YN) interaction**

- YN interaction is essential in probing neutron star inner core

- **Hyperon puzzle:** do hyperons exist in the dense inner core of neutron stars?

- No direct probe method
- Rely on theoretical models
- Lack of direct experimental data of YN, YNN, YY interactions to constrain theoretical models of the dense matter equation of state (EoS)



Introduction: RHIC BES-II

- RHIC beam energy scan Phase II (BES-II)

- Specific focus on low $\sqrt{s_{NN}}$

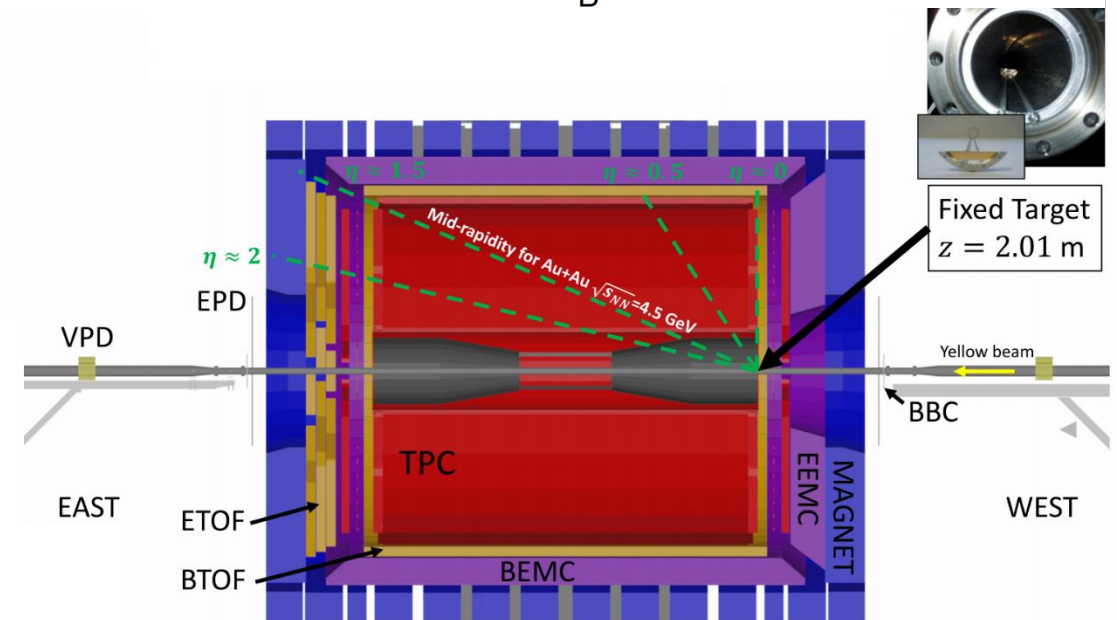
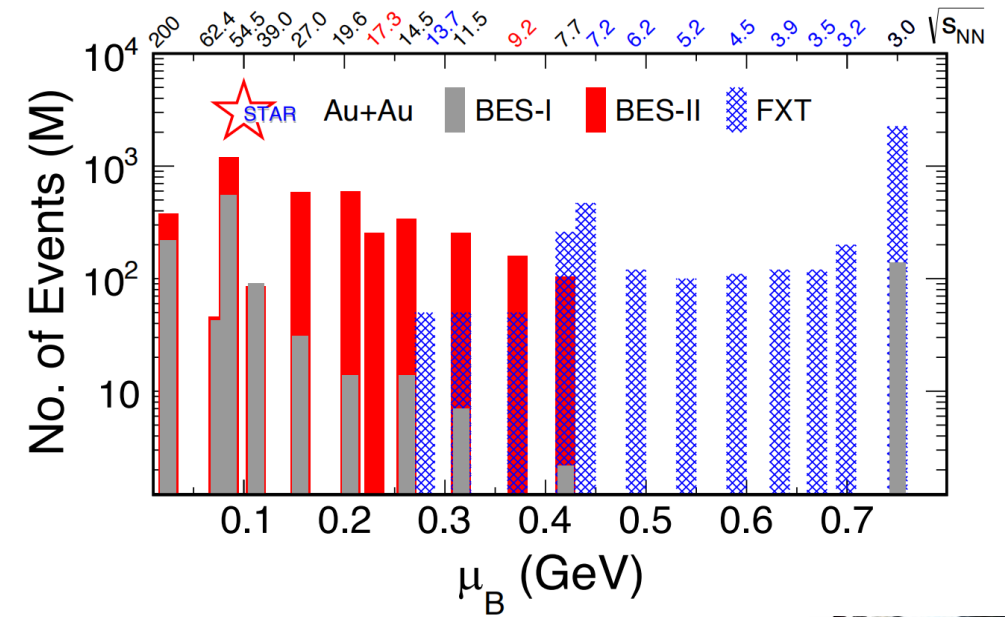
Include fixed target (FXT) mode extends down to 3.0 GeV, increase μ_B range from ~ 400 MeV to ~ 700 MeV

- High statistics data

- Improve systematics

- Detector upgrade: iTPC, EPD, eTOF

Enhances the capability of various measurements with excellent precision

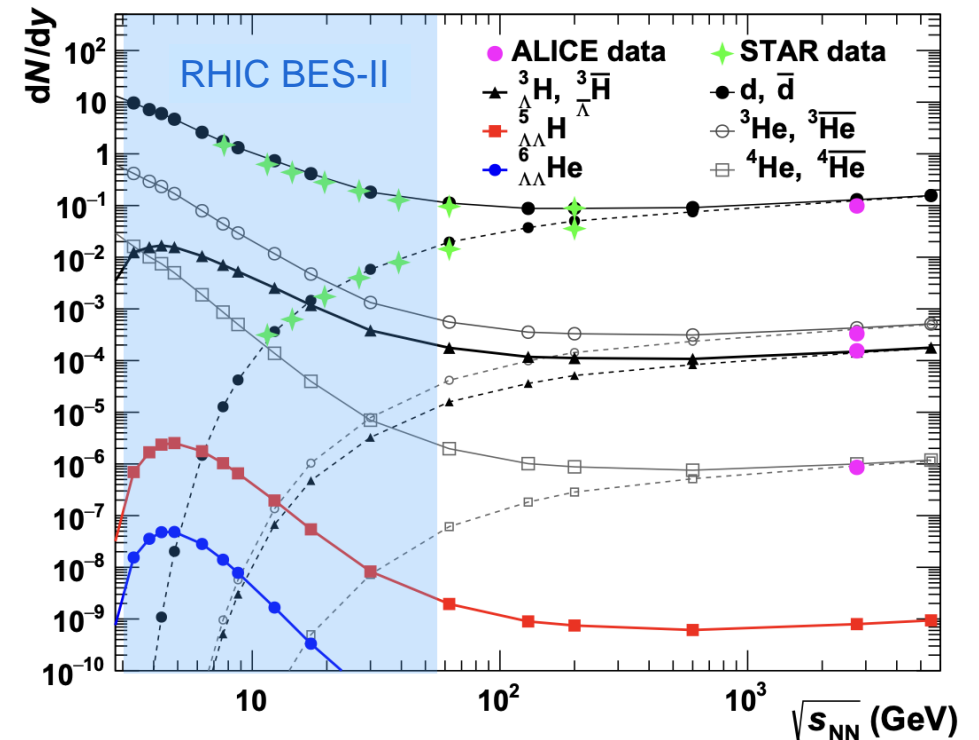


Production mechanism of hypernuclei is still not well understood.

Hypernuclei formation process in relativistic heavy-ion (HI) collisions can be studied through measurements related to **spectra and collective flow**.

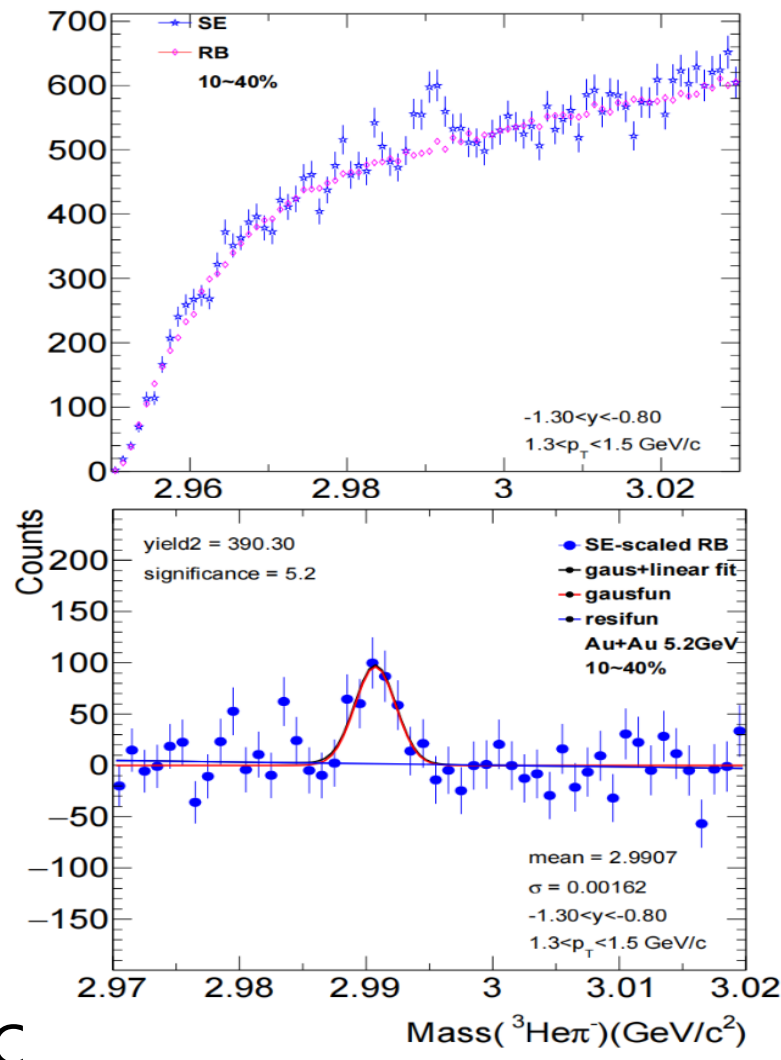
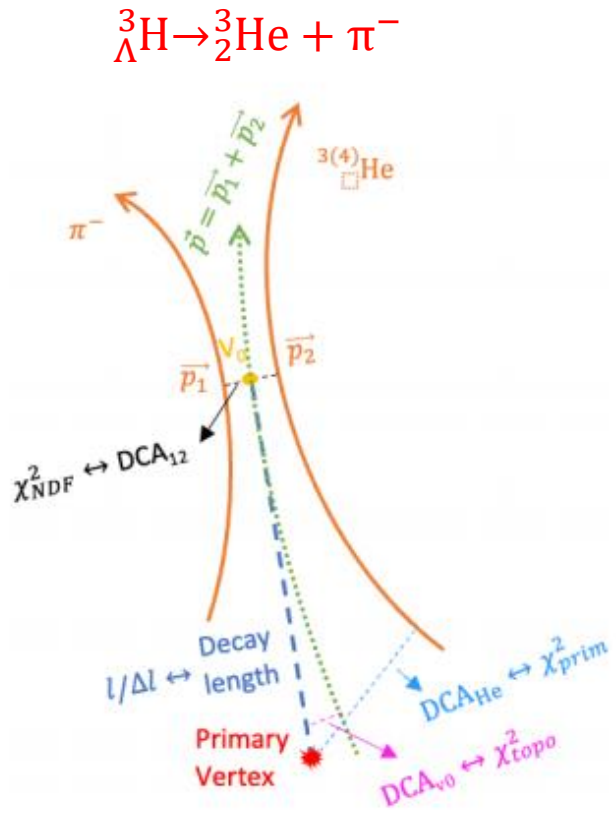
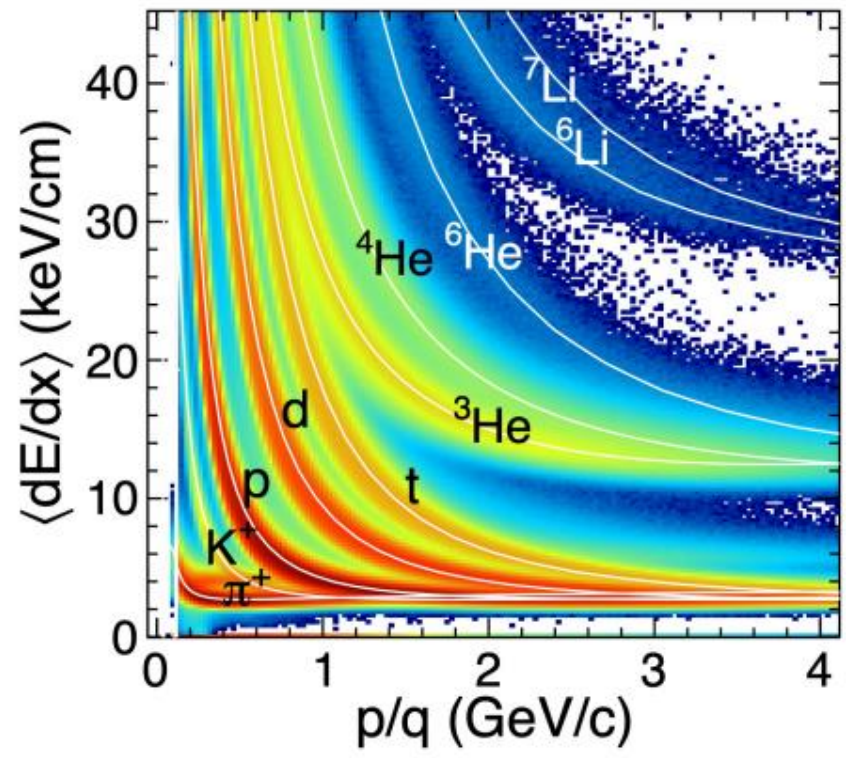
- Hypernuclei measurements are scarce in HI collision experiments
- At **low beam energies**, hypernuclei production is expected to be **enhanced** due to high baryon density

RHIC BES-II offers great opportunity for hypernuclei measurements.



B. Dönigus, Eur. Phys. J. A (2020) 56:280
A. Andronic et al. PLB (2011) 697:203–207

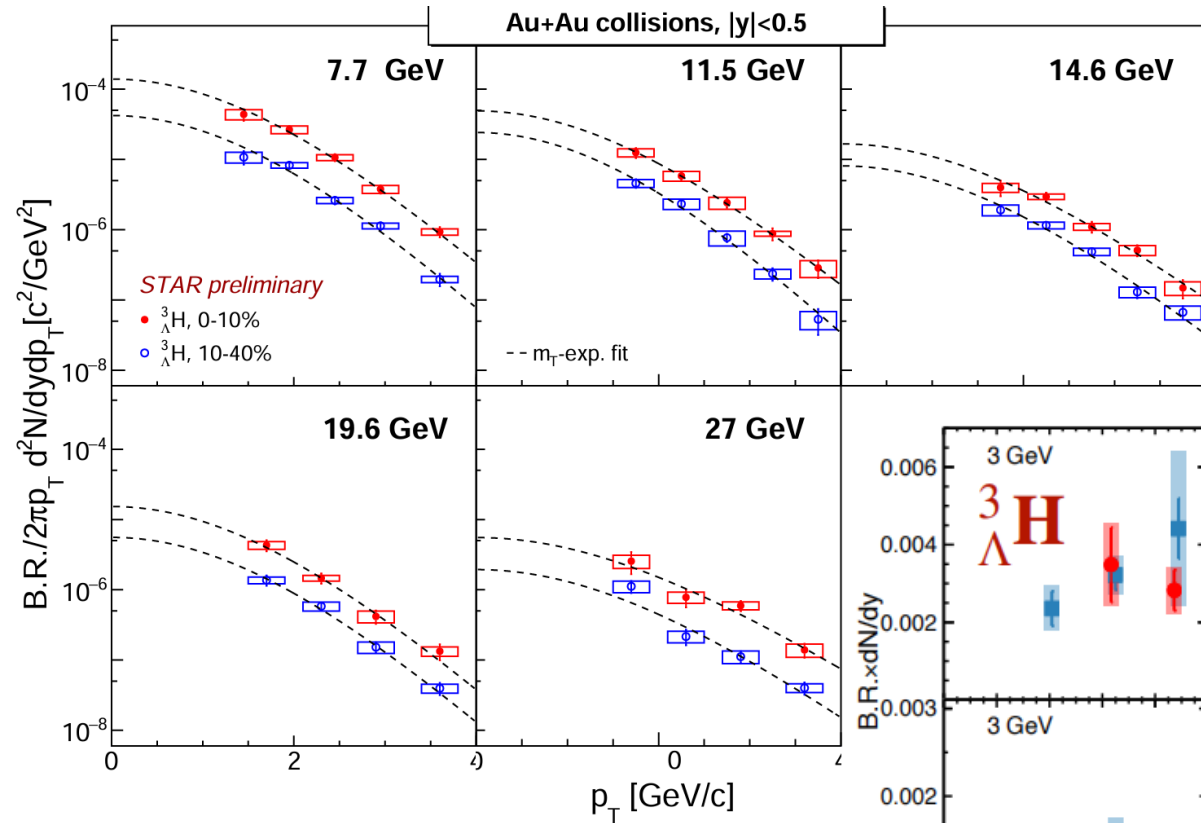
${}^3_{\Lambda}\text{H}$ reconstruction



- Reconstruction channel: ${}^3_{\Lambda}\text{H} \rightarrow {}^3_2\text{He} + \pi^-$
- Particle identification from energy loss measurement using TPC
- KF particle package * is used for signal reconstruction

* XY. Ju et al. Nucl.Sci.Tech. 34 (2023) 10, 158

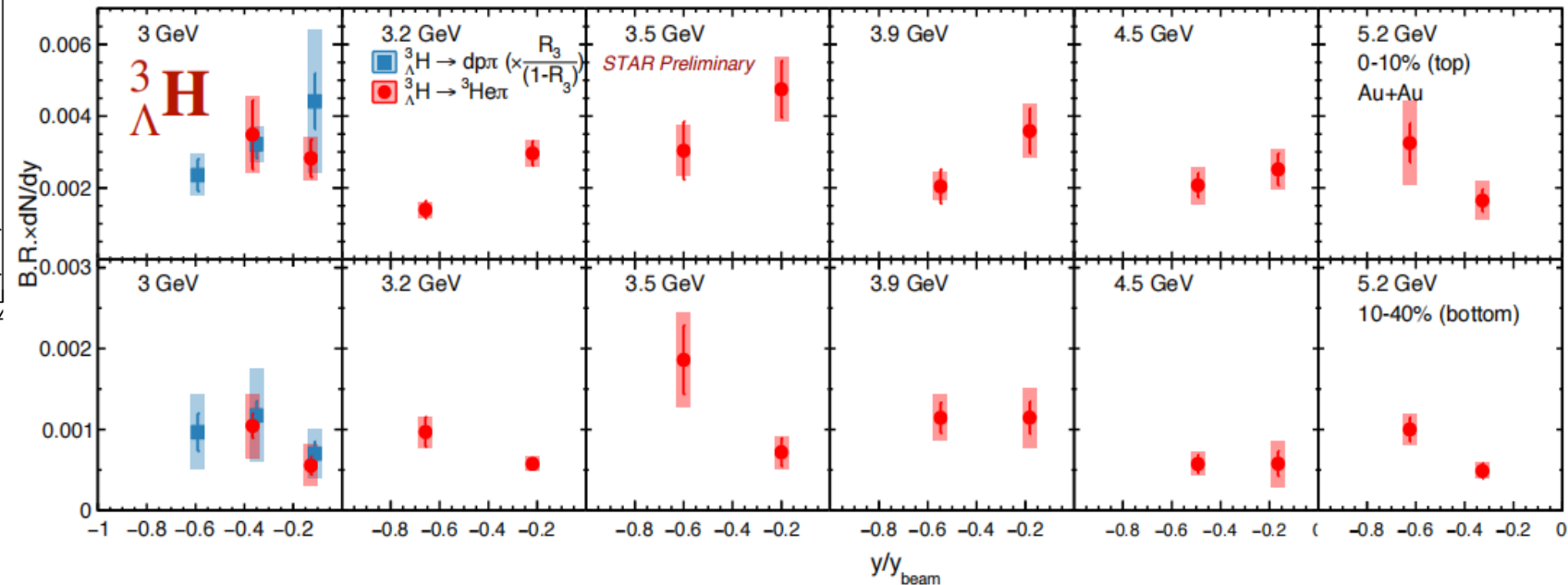
${}^3\Lambda$ H rapidity and p_T spectra



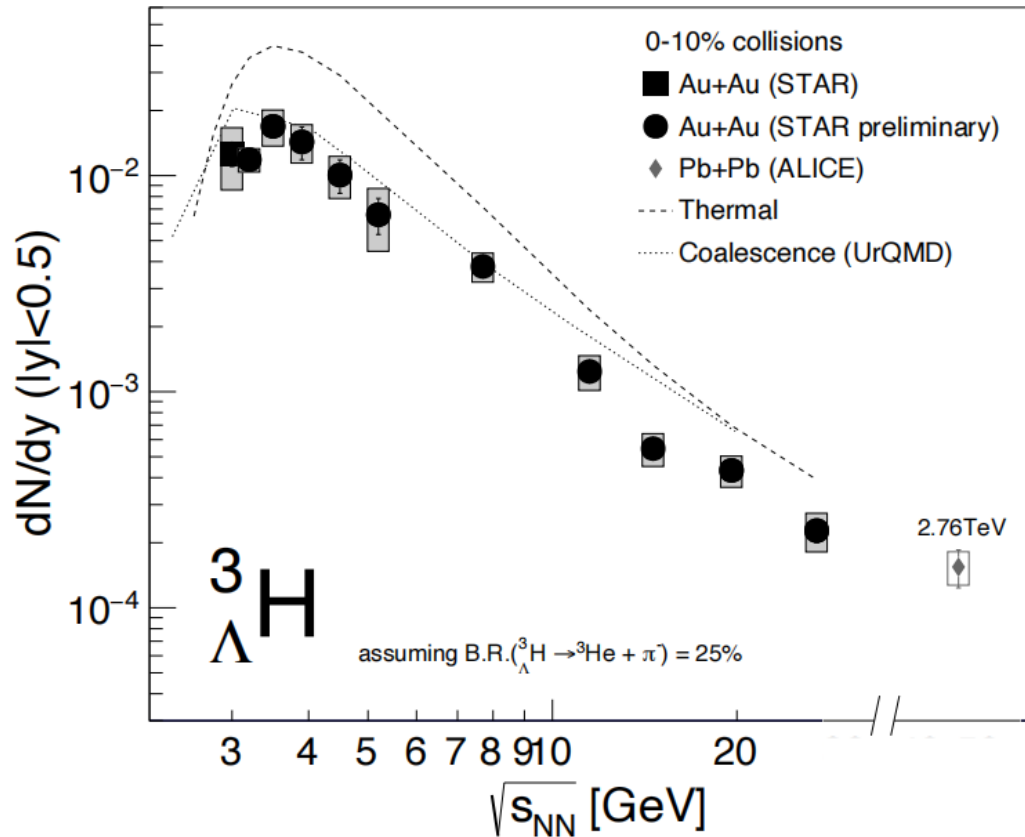
- Measurements cover 11 different energies

Collider: 7.7, 11.5, 14.6, 19.6, 27 GeV

Fixed Target: 3.0, 3.2, 3.5, 3.9, 4.5, 5.2 GeV



Energy dependence of ${}^3_{\Lambda}\text{H}$ production

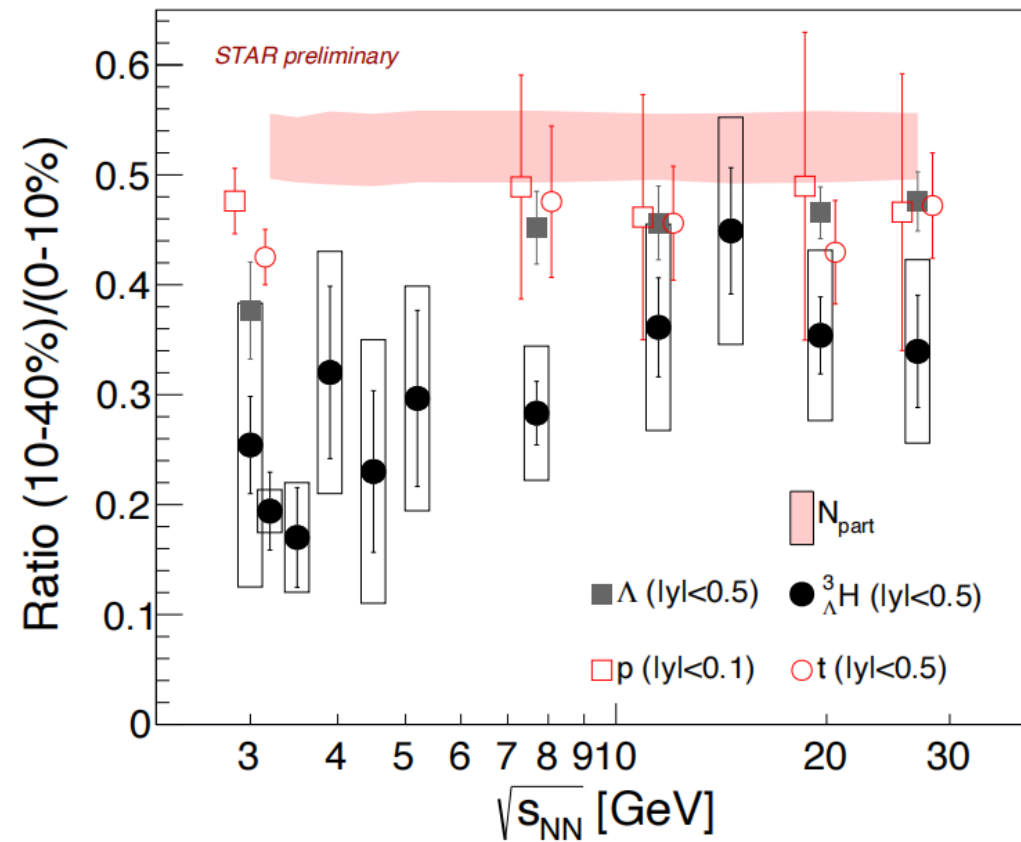
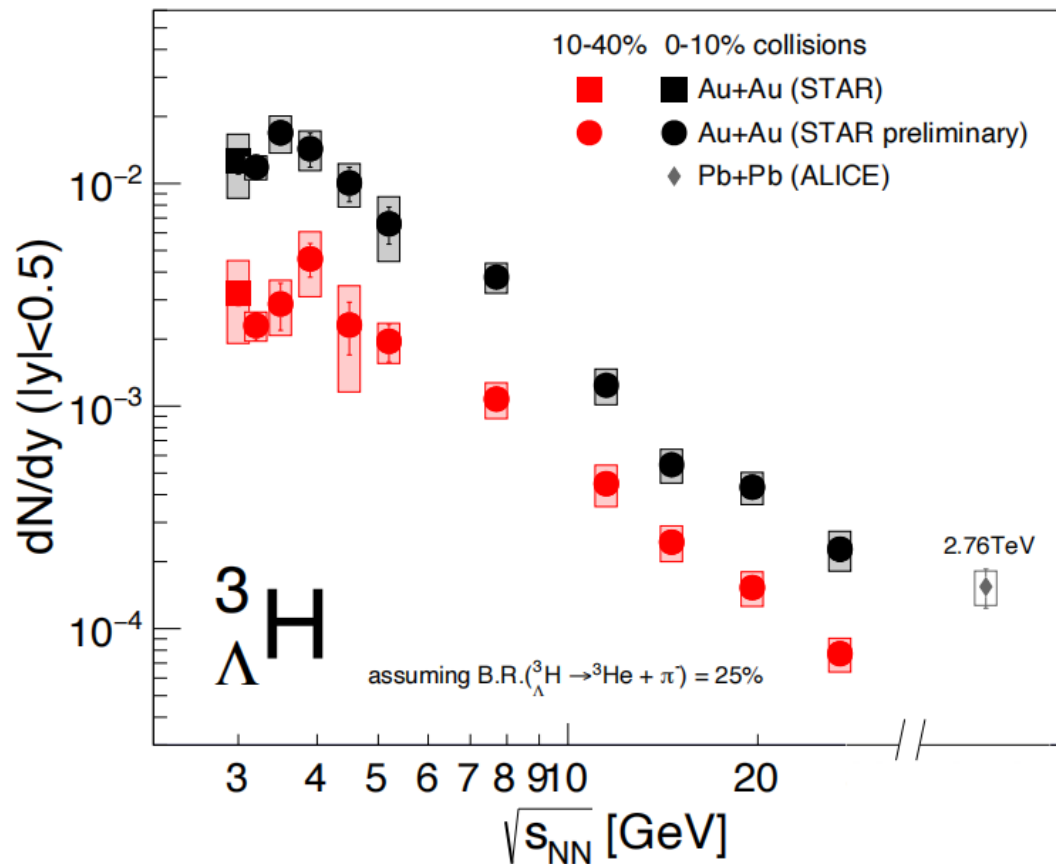


- Yields increase strongly from $\sqrt{s_{\text{NN}}} = 27$ GeV to ~ 4 GeV
- Peak at 3-4 GeV
- Hadronic transport + coalescence models qualitatively describe the data
- Thermal model overestimates the data

First energy dependence of ${}^3_{\Lambda}\text{H}$ production yields in the high-baryon-density region

STAR, PRL 128 (2022) 202301
ALICE, PLB 754 (2016) 360
T. Reichert, et al, PRC 107 (2023) 014912

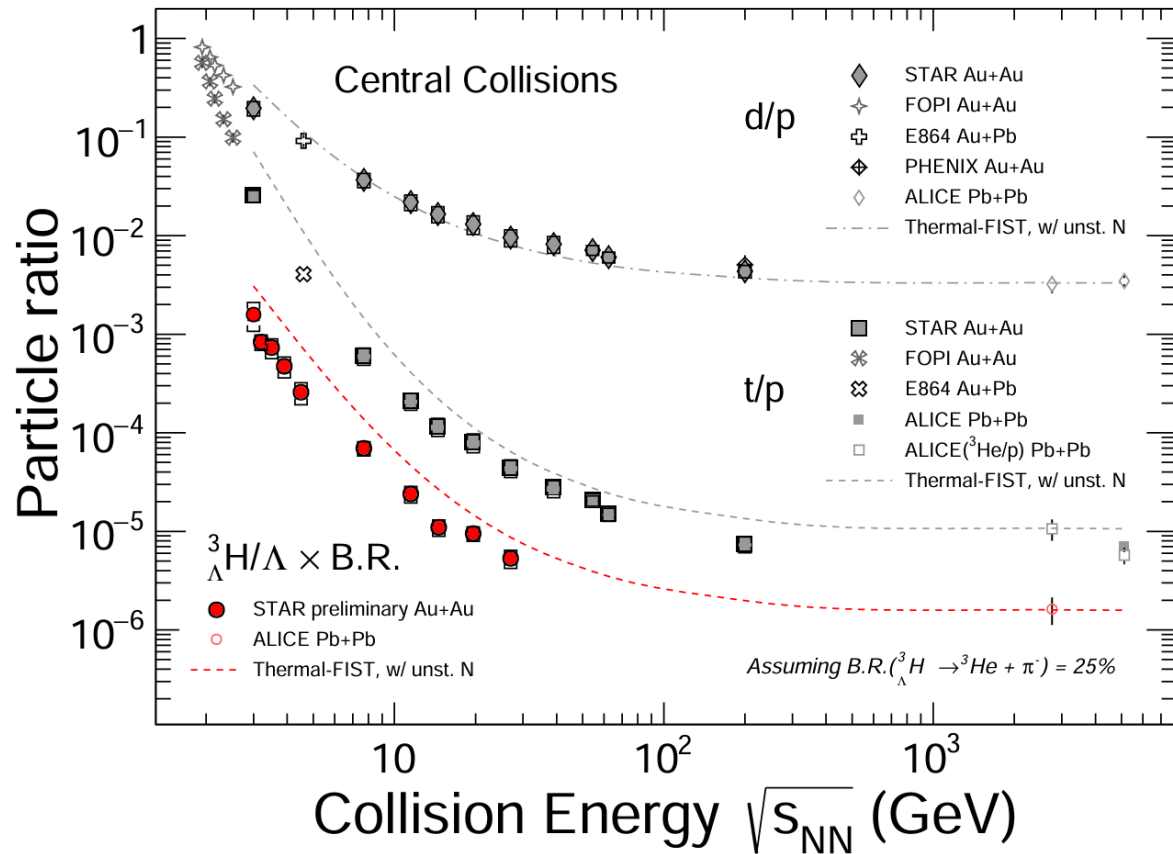
Centrality dependence of ${}^3\Lambda\text{H}$ production



- Similar trend in central (0-10%) and mid-central (10-40%) collisions
- Suppression of mid-central/central ${}^3\Lambda\text{H}$ yield ratio seems more apparent below $\sqrt{s_{NN}} = 7.7$ GeV
- ${}^3\Lambda\text{H}$ yield ratio tends to increase more steeply than proton, Λ , triton at low energies

Suppression of ${}^3\Lambda\text{H}$ production in mid-central collisions at low energies compared to central collisions

Nuclei-to-Hadron ratios



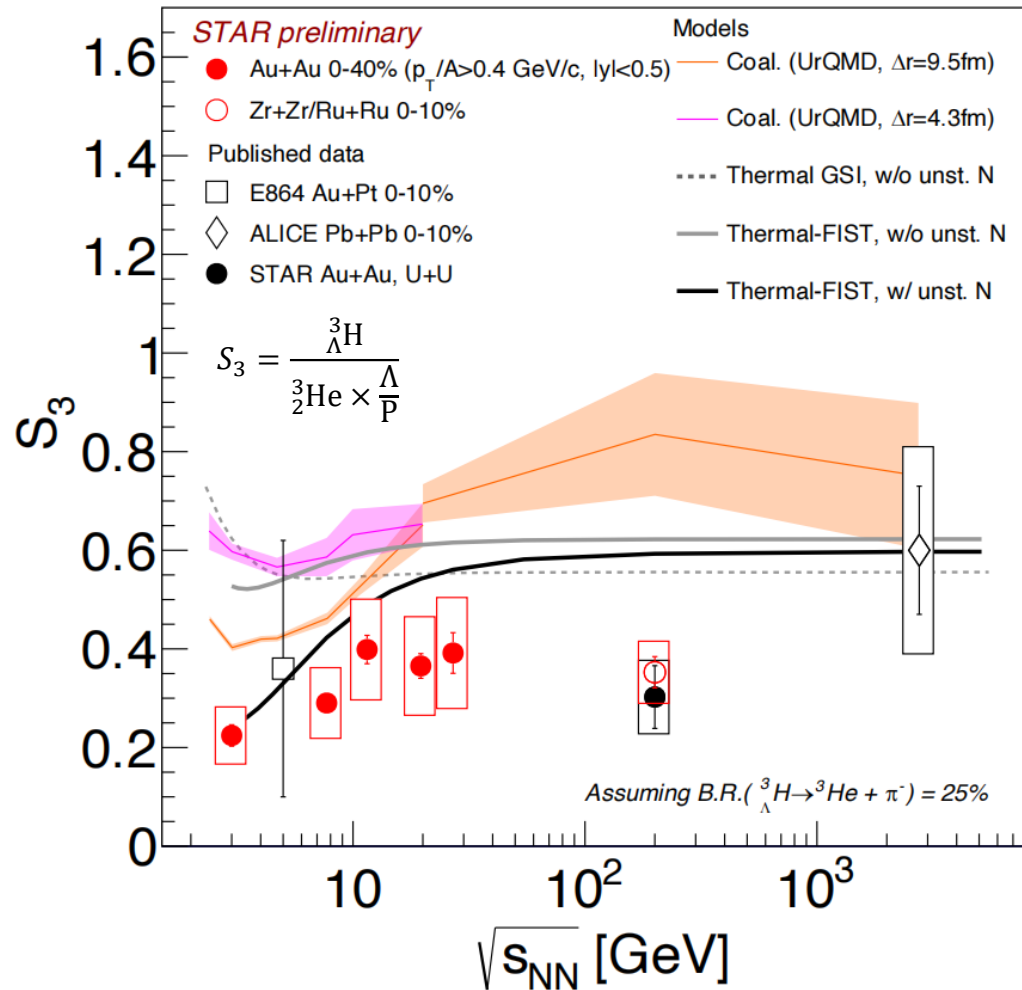
- Thermal model, assuming that chemical freeze-out of light/ hypernuclei happens at same time with hadrons, **overestimates** ${}^3\text{H}/\Lambda$ by a factor of ~ 2 , as well as t/p

- In thermal model, particle yield ratio is independent of volume. ${}^3\text{H}/\Lambda$ yield ratio is dependent of strangeness correlation length

Might suggest ${}^3\text{H}$ and t yields are not in equilibrium and fixed at chemical freeze-out simultaneously with other hadrons

STAR, PRL 130 (2023) 202301
STAR, arXiv: 2311.11020
T. Reichert, et al, PRC 107 (2023) 014912

Energy dependence of S_3

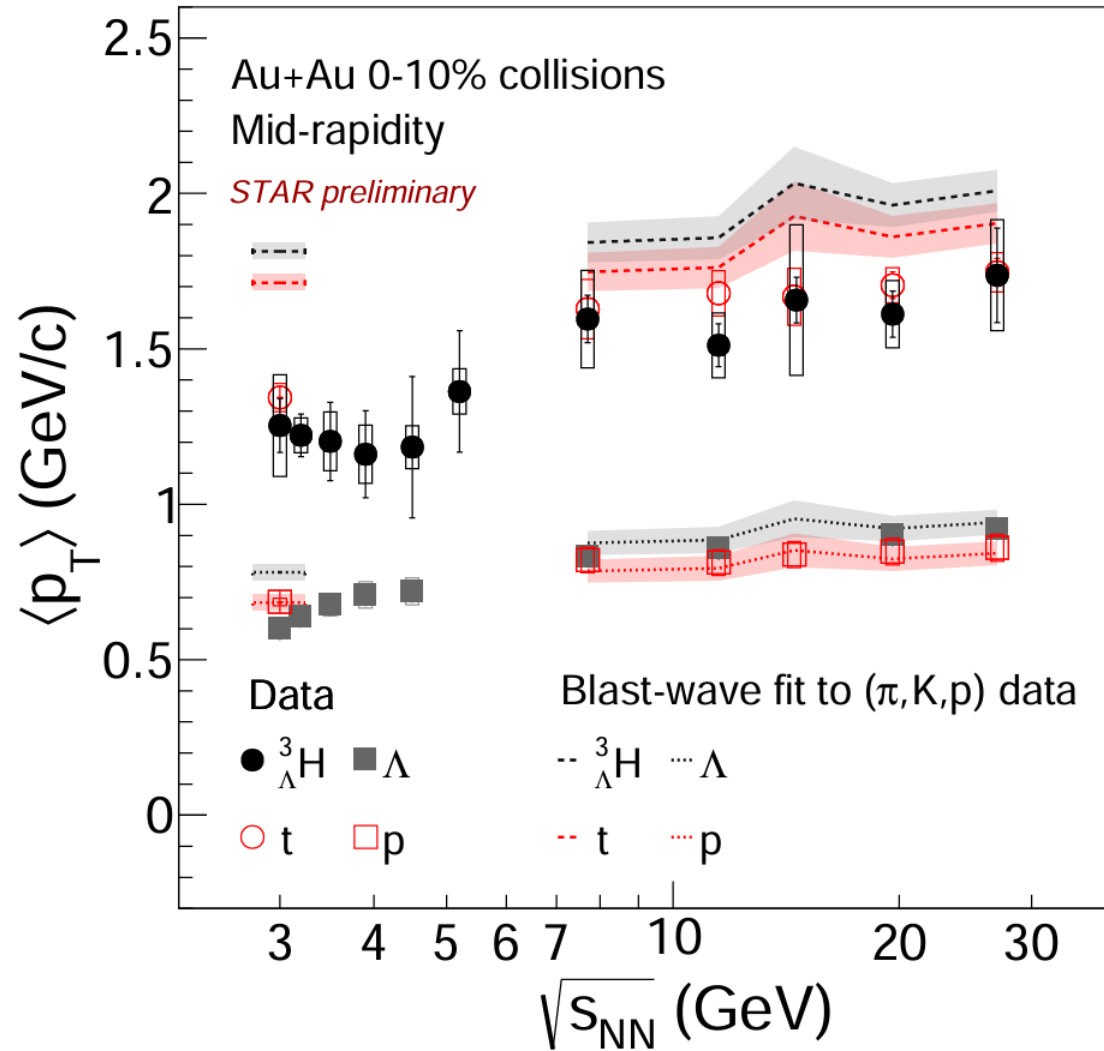


- A prominent enhancement of S_3 was proposed as a probe for deconfinement
- Data shows a mild increasing trend from $\sqrt{s_{NN}} = 3.0$ GeV to 2.76 TeV
- For coalescence (UrQMD) models, the energy dependence is sensitive to the **source radius** (Δr)
 - Data favor larger radius
- Thermal-FIST, which includes **feed-down** from unstable nuclei to stable $p, {}^3He$, describes the S_3 data better
 - Possible feed-down should be accounted

STAR, Science 328 (2010) 58
 STAR, arXiv: 2310.12674
 ALICE, PLB 754 (2016) 360
 E864, PRC 70 (2004) 024902

A. Andronic et al, PLB 697 (2011) 203 (Thermal (GSI))
 S. Zhang, PLB 684 (2010) 224 (Coal.+AMPT)
 T. Reichert, et al, PRC 107 (2023) 014912 (UrQMD, Thermal-FIST)

Energy dependence of $\langle p_T \rangle$ for ${}^3_\Lambda\text{H}$



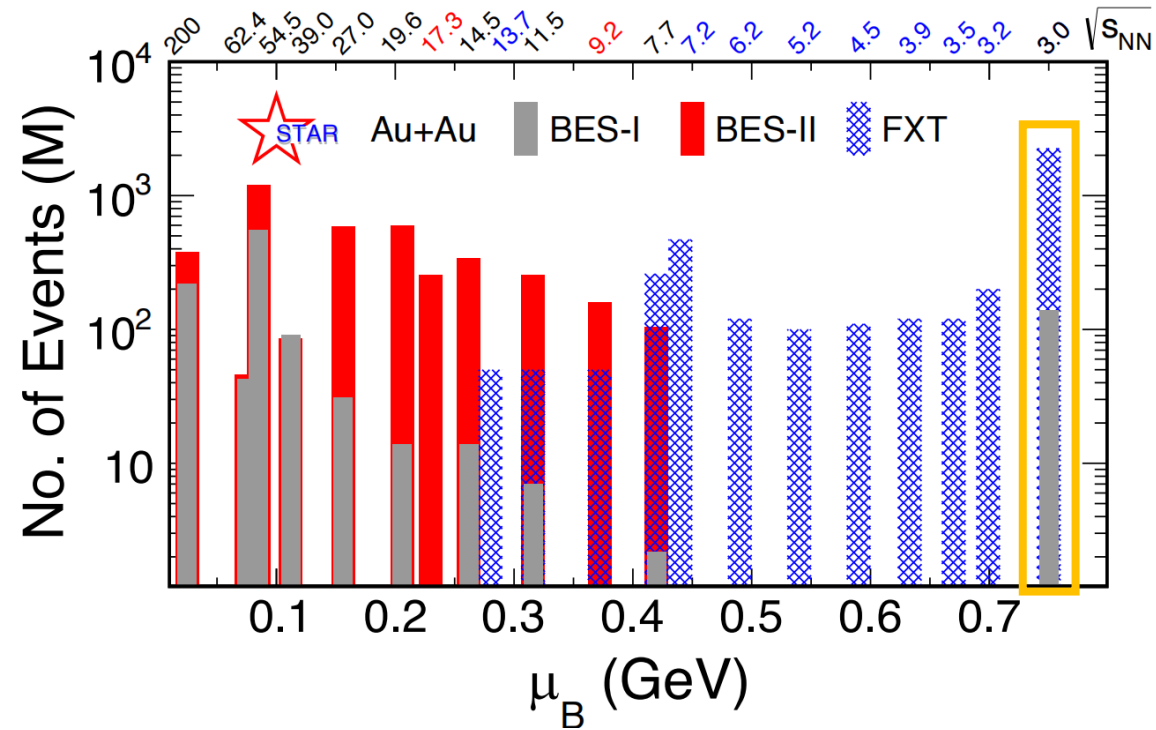
- Similar $\langle p_T \rangle$ for ${}^3_\Lambda\text{H}$ and t
- Hint of ${}^3_\Lambda\text{H}$ and t $\langle p_T \rangle < \langle p_T \rangle^{BW} > 7.7\text{GeV}$
Blast-wave expectation calculated using measured kinetic freeze-out parameters from light hadrons (π, K, p) spectra

${}^3_\Lambda\text{H}$ and t might do not follow same collective expansion as light hadrons. Can be interpreted as ${}^3_\Lambda\text{H}$ and t decoupling at different times compared to light hadrons

- Different trend around $\sqrt{s_{NN}} = 4.5\text{ GeV}$
→ Suggest different expansion dynamics?

- First measurements of ${}^3_{\Lambda}\text{H}$ yields vs. energy at high baryon density region presented, consistent with coalescence model
- Suppression of ${}^3_{\Lambda}\text{H}$ in 10-40% collisions at low collisions energies observed
- ${}^3_{\Lambda}\text{H}/\Lambda$ ratio vs. energy seems in contradiction to that ${}^3_{\Lambda}\text{H}$ freeze out at the same time as light hadrons
- S_3 vs. collision energy favor large ${}^3_{\Lambda}\text{H}$ radius
- ${}^3_{\Lambda}\text{H} \langle p_T \rangle$ overestimated by Blast-wave fit parameterization from light hadrons
 - ${}^3_{\Lambda}\text{H}$ are likely formed at or decouples from the system at a different time compared to the light hadrons

Outlook



- Huge datasets enable precision hypernuclei measurements
 - **Run 21, Au+Au 3 GeV, ~2 billion events**
 - Run 18, Isobar 200 GeV, ~6 billion events
- Opportunities for heavier hypernuclei: ${}^4_{\Lambda}\text{H}$, ${}^4_{\Lambda}\text{He}$, ${}^5_{\Lambda}\text{He}$, ${}^6_{\Lambda}\text{H}$, ${}^A_{\Lambda\Lambda}\text{H}$, ${}^A_{\Lambda\Lambda}\text{He}$

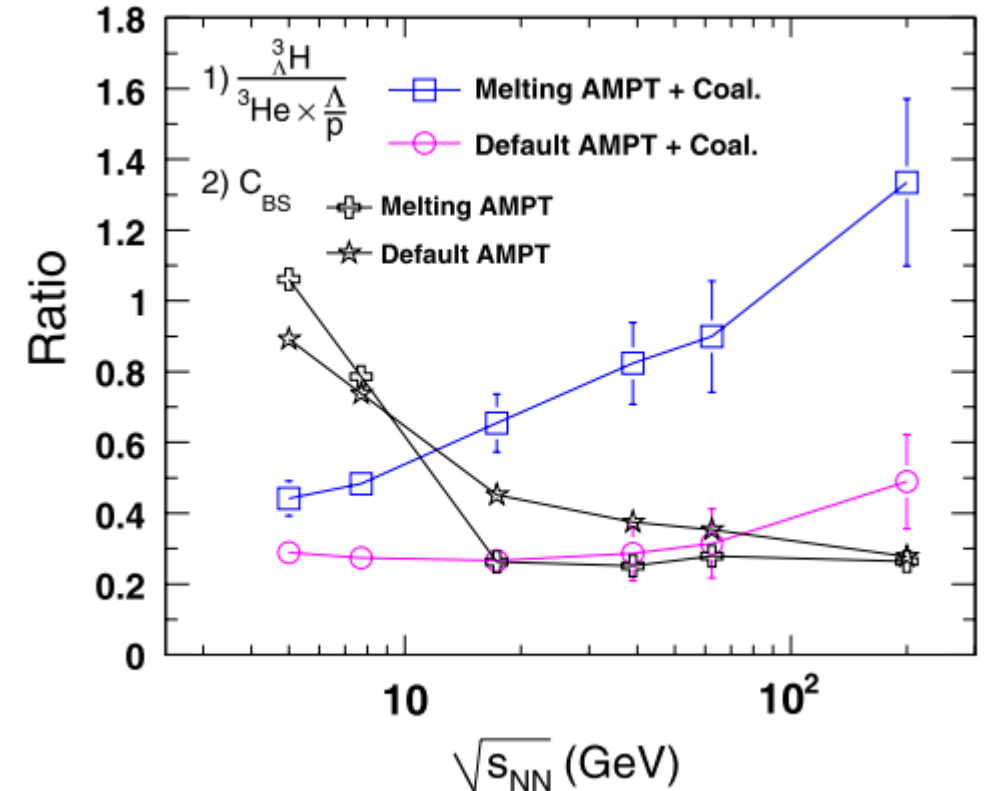
Back up

Introduction: the strangeness population factor S_3

- S_3 may be sensitive to the **onset of deconfinement**

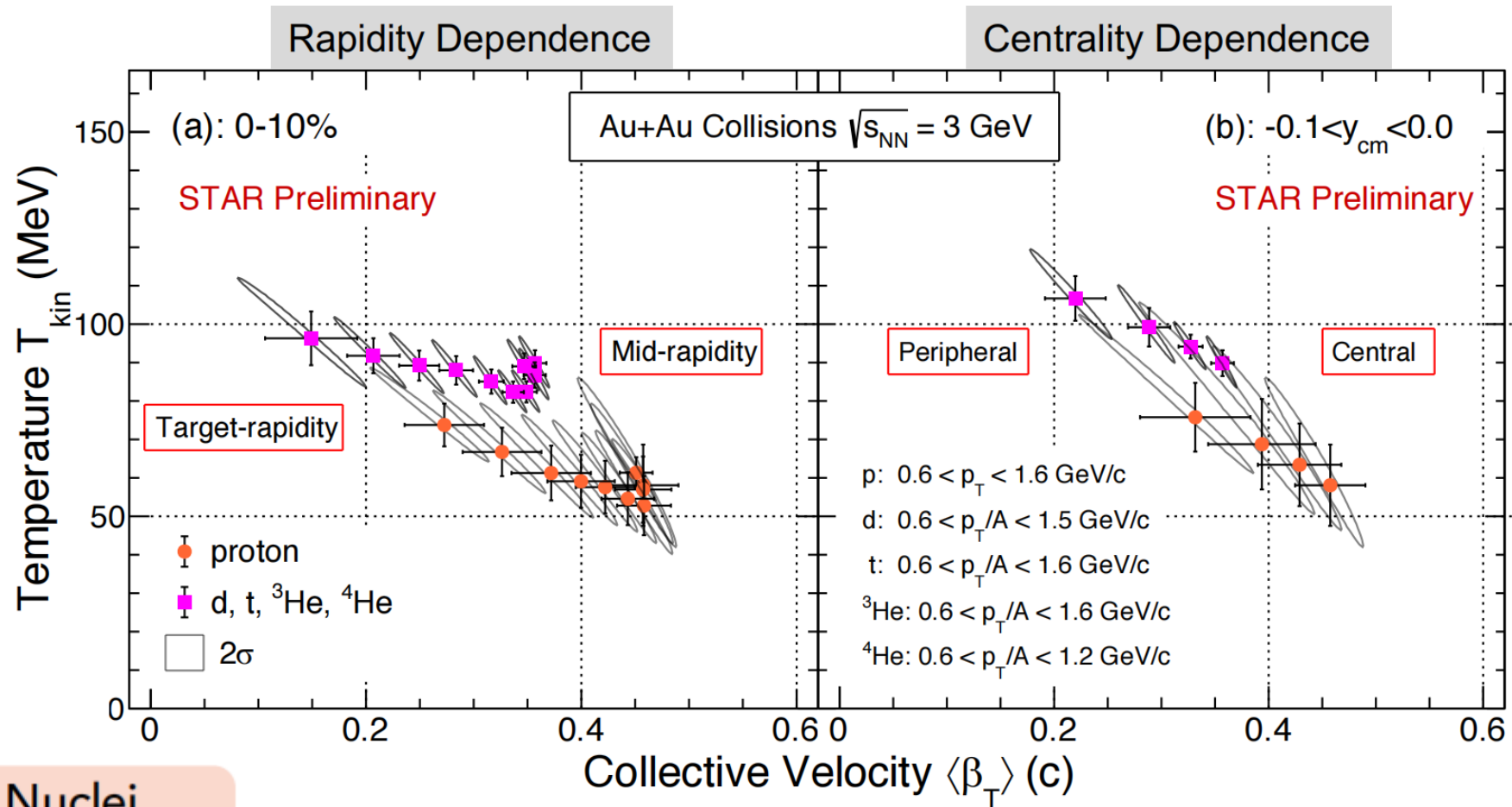
$$S_3 = \frac{{}^3\Lambda\text{H}}{{}^3_2\text{He} \times \frac{\Lambda}{\bar{p}}}$$

- S_3 maybe enhanced in a system involving partonic interactions
- Models suggest S_3 is more sensitive to the local baryon-strangeness correlation than the global baryon-strangeness correlation coefficient (C_{BS})



S. Zhang et al. PLB 684 (2010) 224–227

Kinetic Freeze-out Dynamics at 3 GeV



Light Nuclei freeze out earlier?

J. Adams et al. [STAR Collaboration] Nucl.Phys.A 757 (2005) 102-183

- The kinetic freeze-out parameters (T_{kin} and $\langle \beta_T \rangle$) show rapidity and centrality dependence
- The freeze-out parameter (T_{kin}) of the light nuclei is systematically higher than that of the protons

Hui Liu, QM2022 @ Krakow, Poland