

第十四届全国粒子物理学术会议

高能物理和量子物理交叉

曹庆宏

北京大学物理学院

北京大学高能物理研究中心

会议主题：人工智能、量子信息、量子计算
在粒子物理、核物理和宇宙学等学科前沿的应用

第28届 LHC Mini-Workshop 会议

通化师范学院 2024年7月10日



第28届LHC Mini-Workshop报告

2024年7月9日-10日

人工智能

罗民兴: Some observations on AI assisted theoretical studies

马滢青: 人工智能回顾及愚见

机器学习

马伯强: Application of machine learning method with cosmic photons

李 亮: 高能物理的下一场革命: 从深度学习走向通用人工智能

李. 靖: HEP ML Lab: an end-to-end framework for machine learning application in high energy physics

郭禹辰: Using machine learning to optimize the measurements of anomalous gauge couplings

朱永峰: Jet Origin Identification & Quantum-based Jet Clustering

量子计算

李英英: HEP Opportunities in the Quantum Computing Era

刘晓辉: Partonic Collinear Structure by Quantum Computing

杨冀翀: Using machine learning method suitable for quantum computing in the phenomenological study of new physics

量子信息

肖明磊: Emergent Symmetry from Entanglement Suppression

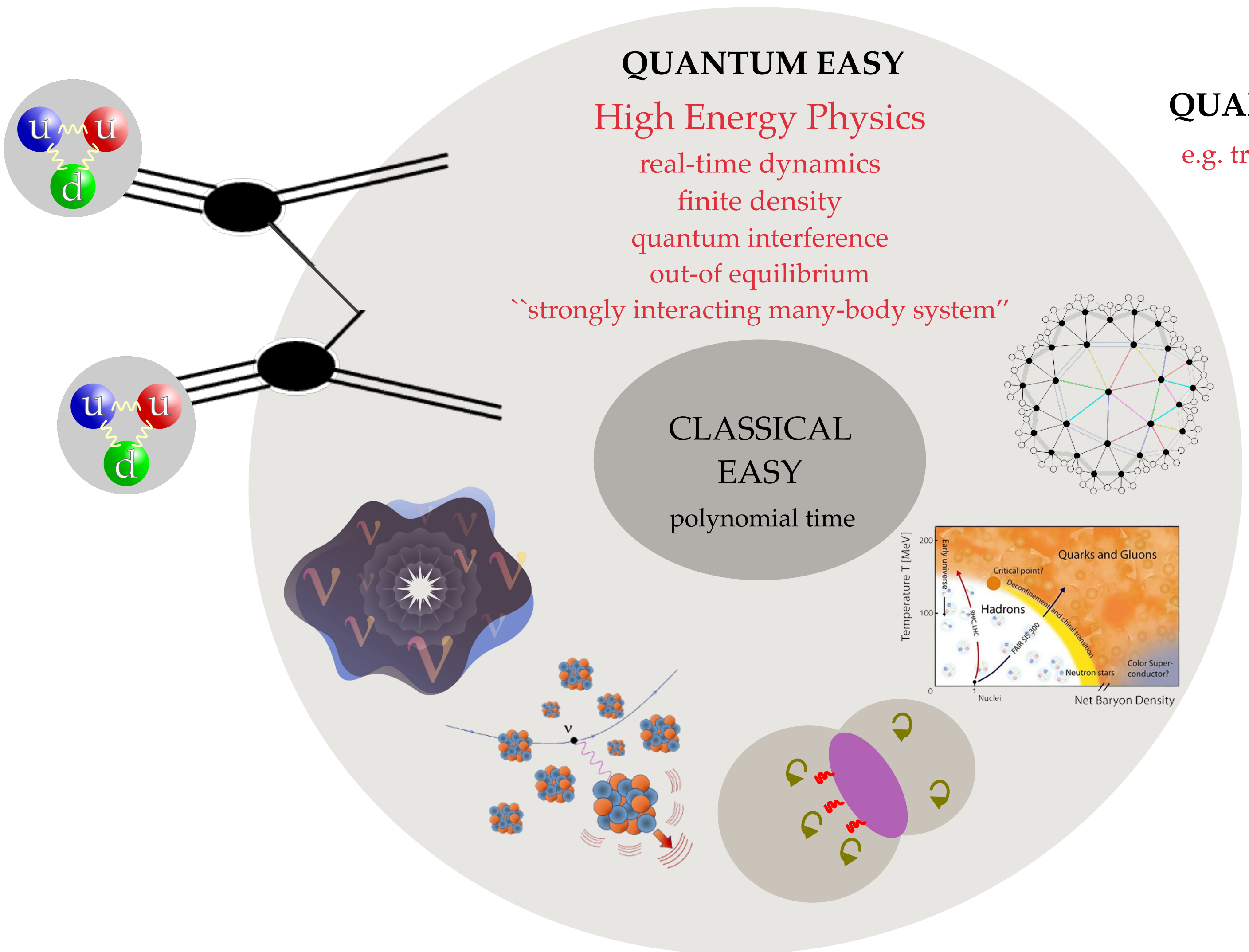
施 郁: Quantum Entanglement in High Energy Physics

程 焜: Quantum Entanglement at High-energy Colliders

量子计算

1、量子计算

李英英



Quantum Computing

Now - Noisy Intermediate Scale Quantum (NISQ) era
 more than 50 well controlled qubits, not error-corrected yet

<p>superconducting processor</p> <p>176 qubits 54 qubits 1121 qubits access to 133 qubits</p>	<p>multi-chip quantum processor</p> <p>80 qubits</p>
<p>photon qubits</p> <p>Jiuzhang - 255 qubits</p>	<p>trapped ion qubits</p> <p>22 qubits</p>

48 logical qubits

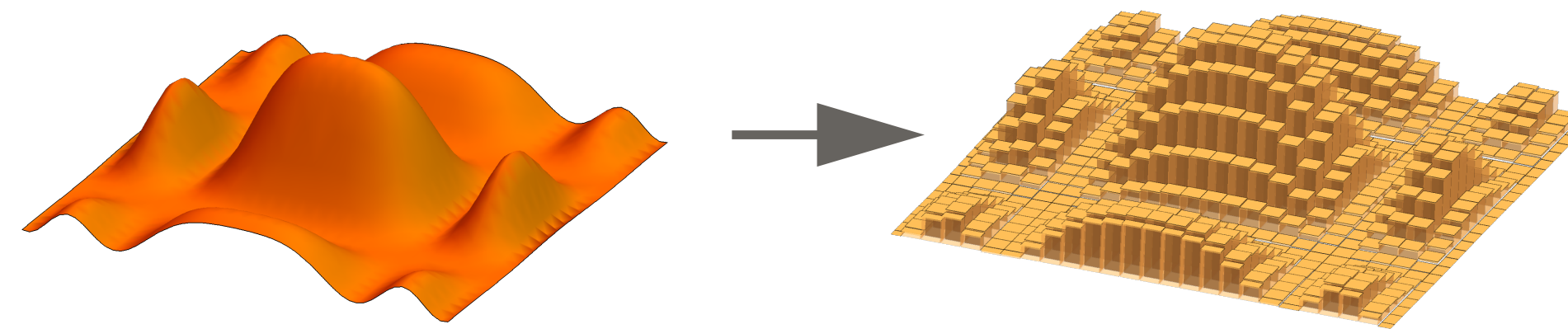
Problems in HEP that are beyond classical easy but are “QUANTUM EASY”

Quantum Computing for HEP

[Jordan, Lee, Preskill, 2011]

$$\int \mathcal{D}\phi e^{iS} = \langle x | e^{-iHt} | y \rangle$$

Discretization



infinities in space

Carena, Lamm, YYL, Liu,
Gustafson, Water,...

Digitization

$$|q\rangle^N \rightarrow |G\rangle$$

infinities in field variables

Bauer, Davoudi, Gustafson, Meurice,
Lamm, YYL, Savage,...

Initialization

$$\mathcal{U} |G\rangle^L \rightarrow |\psi_0\rangle$$

ground/thermal/bound state prep

Karsen, Davoudi, Lawrence, YYL, Xu, Liu, Xing...

Propagation

$$\mathcal{U} |\psi_0\rangle \rightarrow |\psi(t)\rangle$$

efficiency of time evolutions

Davoudi, Gustafson, YYL, Stryker, Wang, Zohar...

Evaluation

$$\langle \mathcal{O} \rangle$$

parton distribution function,

Lamm, Liu, Yamauchi, Xing...

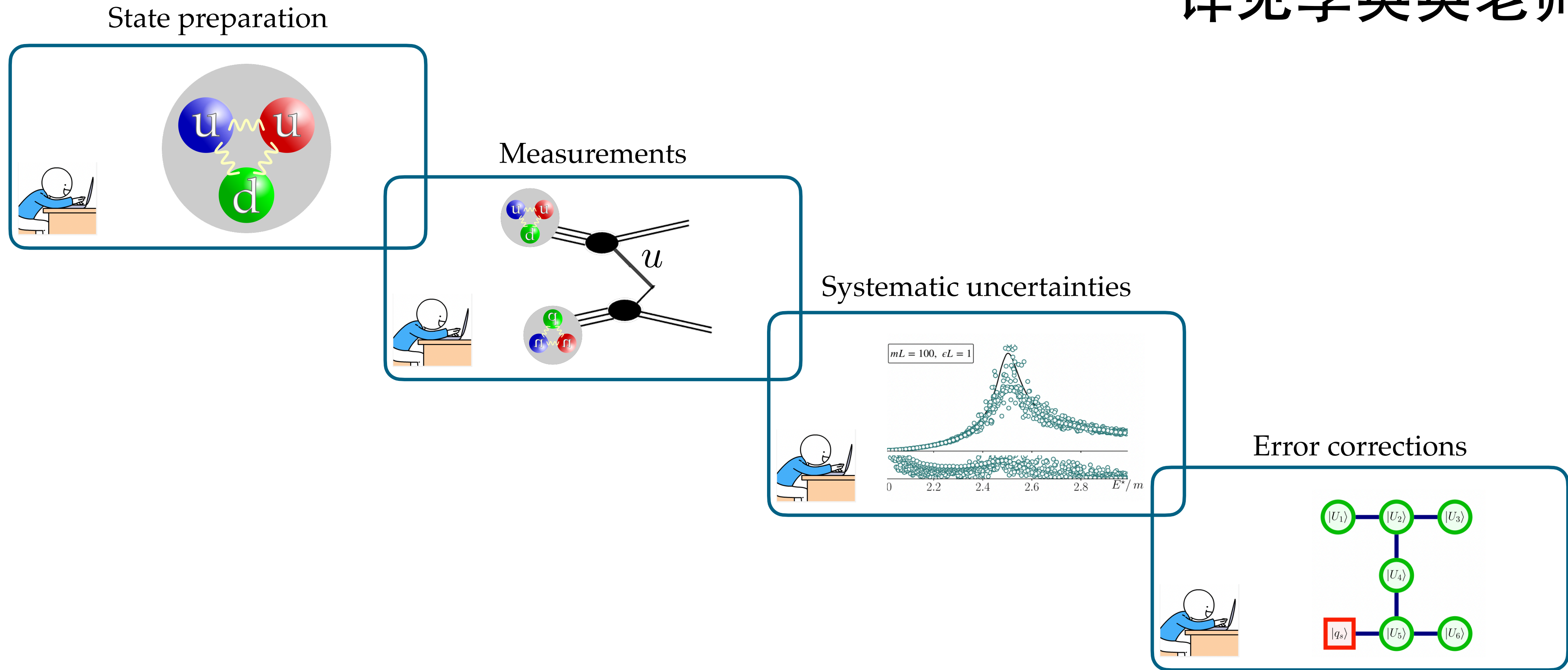
Error mitigation / corrections

gauge symmetry for error corrections

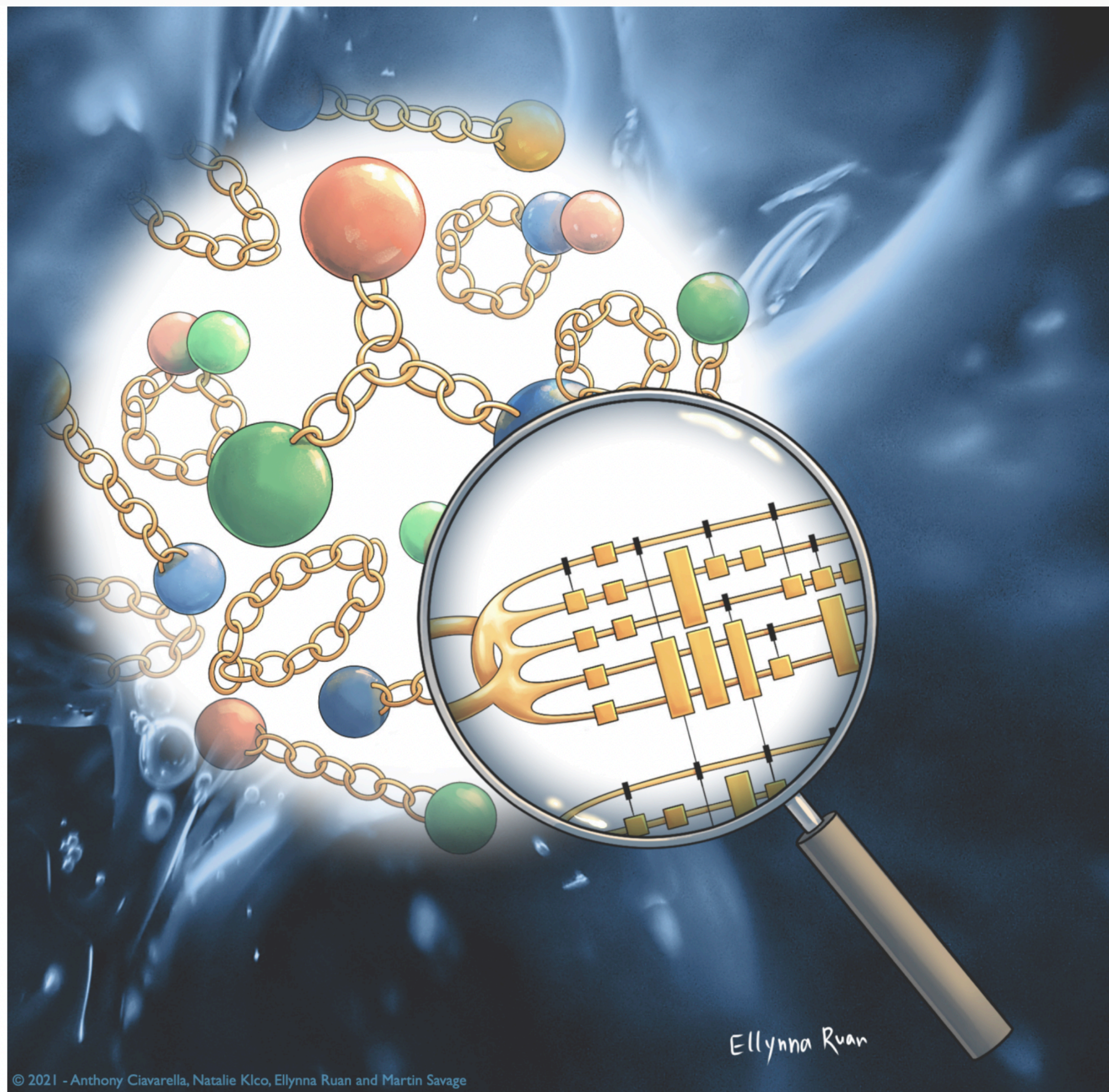
Bauer, Carena, Halimeh, Lamm, YYL,...

To reach the observables — How to do...

详见李英英老师报告



and reach the continuum limit



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Partonic Collinear Structure by Quantum Computing

第28届LHC Mini Workshop

刘晓辉

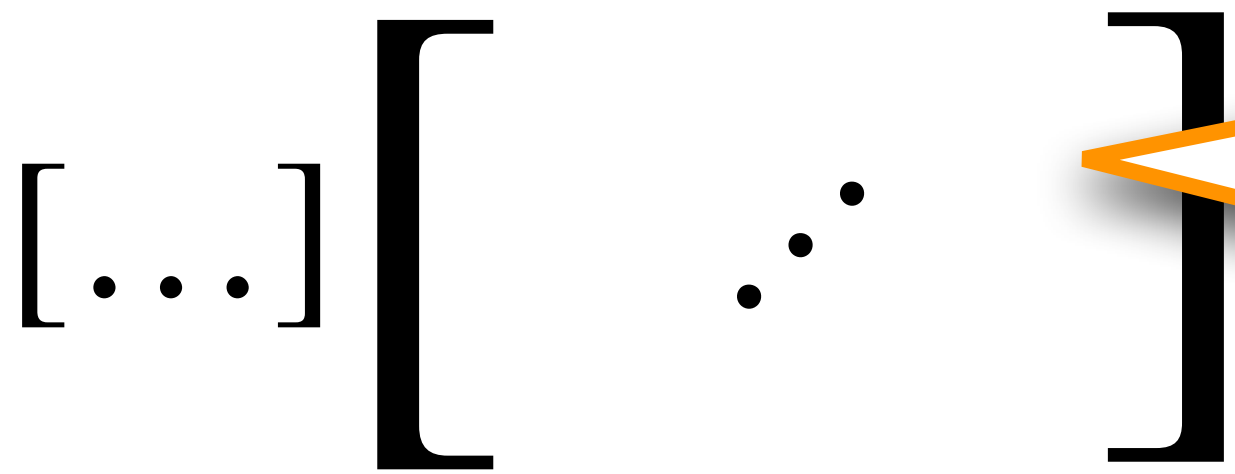


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Quantum Computing for HEP

In HEP, we are dealing with n-pt correlates, i.e. S-matrix

$$\langle \text{out} | e^{-iH[\psi]t} | \text{in} \rangle$$

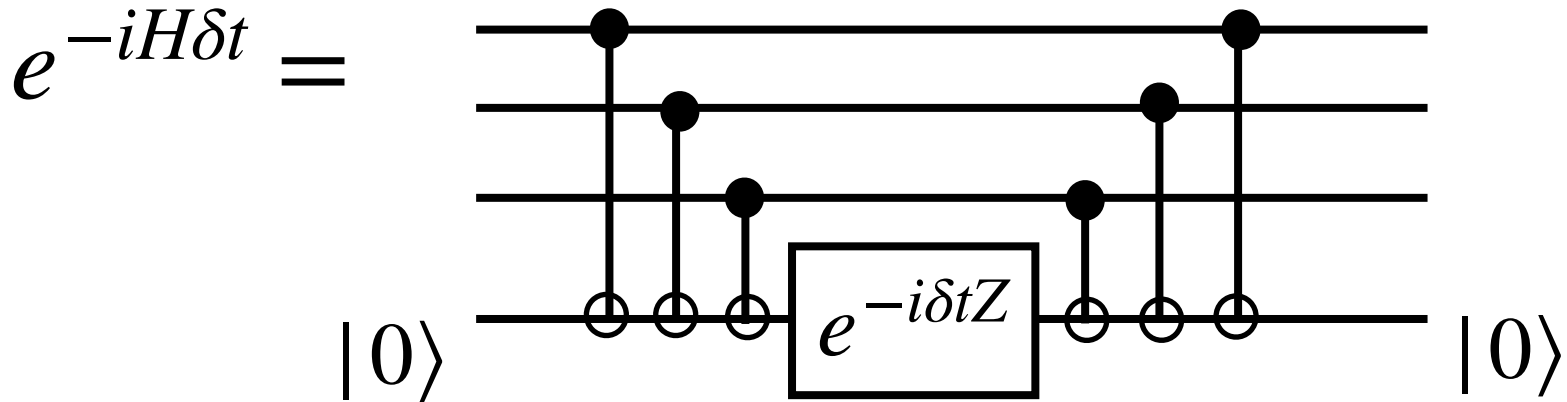


$$\text{dim} = (n_\psi)^{n^{D_{sp}}}$$

decompose to a set of gates, evolution much cheaper

e.g. $H = Z_1 \otimes Z_2 \otimes Z_3$ $e^{-iHt} \approx \lim_{\delta t \rightarrow 0, N \rightarrow \infty} [e^{-iH\delta t}]_N$ Trotter, 1959

Use the fact that $|\phi_1\rangle |\phi_2\rangle \dots |0\rangle \rightarrow |\phi_1\rangle |\phi_2\rangle \dots |\phi_1 \oplus \phi_2 \dots\rangle$



Others can be realized similarly

e.g. by using $e^{-i\delta t X_1 \otimes Z_2 \otimes \dots} = H_1 e^{-i\delta t Z_1 \otimes Z_2 \otimes \dots} H_1$

Quantum computing: reasonable size and operations (scales logarithmically)

Jordan, Lee, Preskill, Science 336, 1130-1133 (2012)

A toy model: Map QFT on to a qubits+gates system

$$\mathcal{L} = \bar{\psi}(i\partial - m)\psi + g(\bar{\psi}\psi)^2 \quad (\text{no gauge, } 1+1)$$

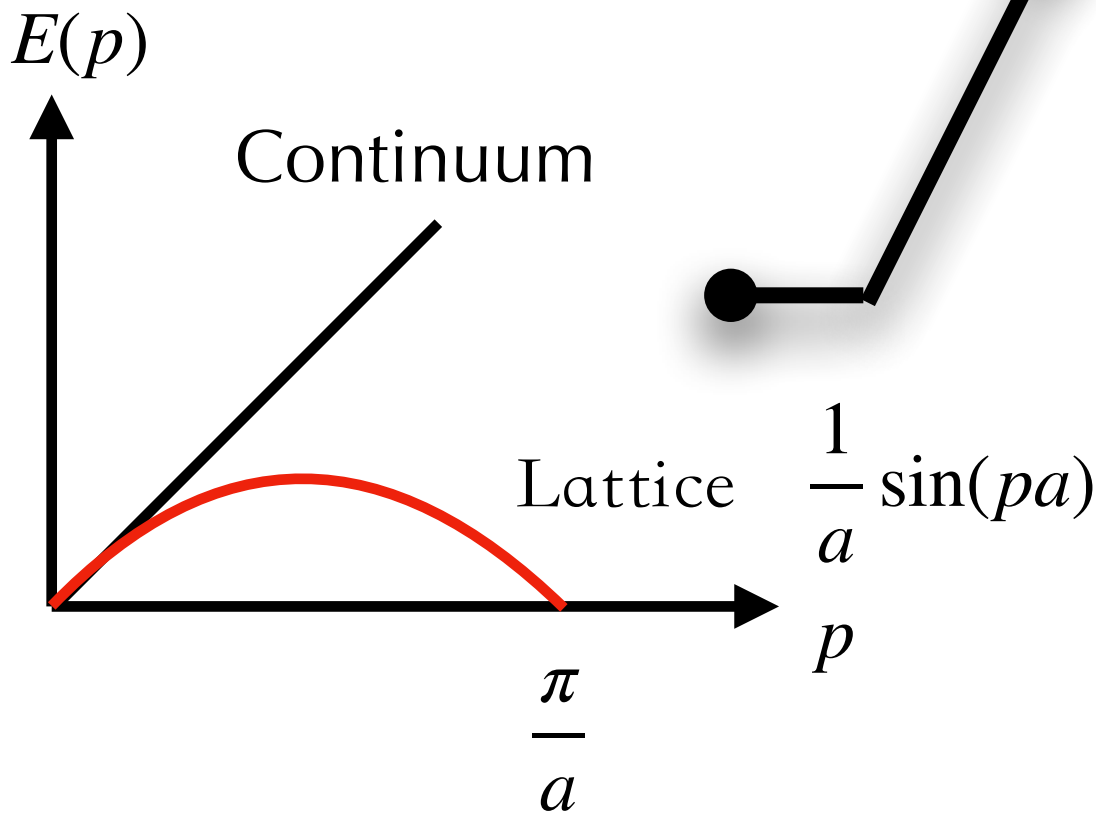
Gross, Neveu, 1974

$$f(x) = \int dz^- e^{-ixM_h z^-} \langle h | \bar{\psi}(z^-) \gamma^+ \psi(0) | h \rangle = \int dz^- e^{-ixM_h z^-} \langle h | e^{iHz} \bar{\psi}(0, -z) e^{-iHz} \gamma^+ \psi(0) | h \rangle$$

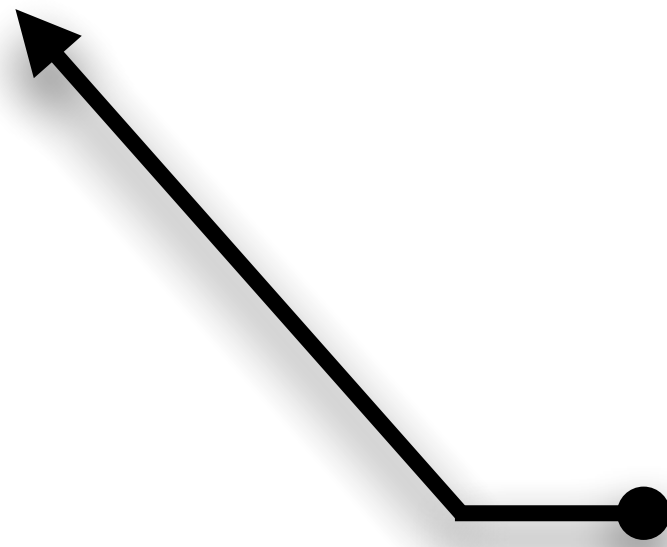
$$\psi = \begin{pmatrix} \psi_1 \\ \psi_2 \end{pmatrix} \rightarrow \begin{pmatrix} \phi_{2n} \\ \phi_{2n+1} \end{pmatrix}$$

Staggered fermion,
Put different fermion
components, flavors on
different sites

$$\phi_n = \prod_{i < n} Z_i (X + iY)_n \quad \text{Jordan-Wigner}$$



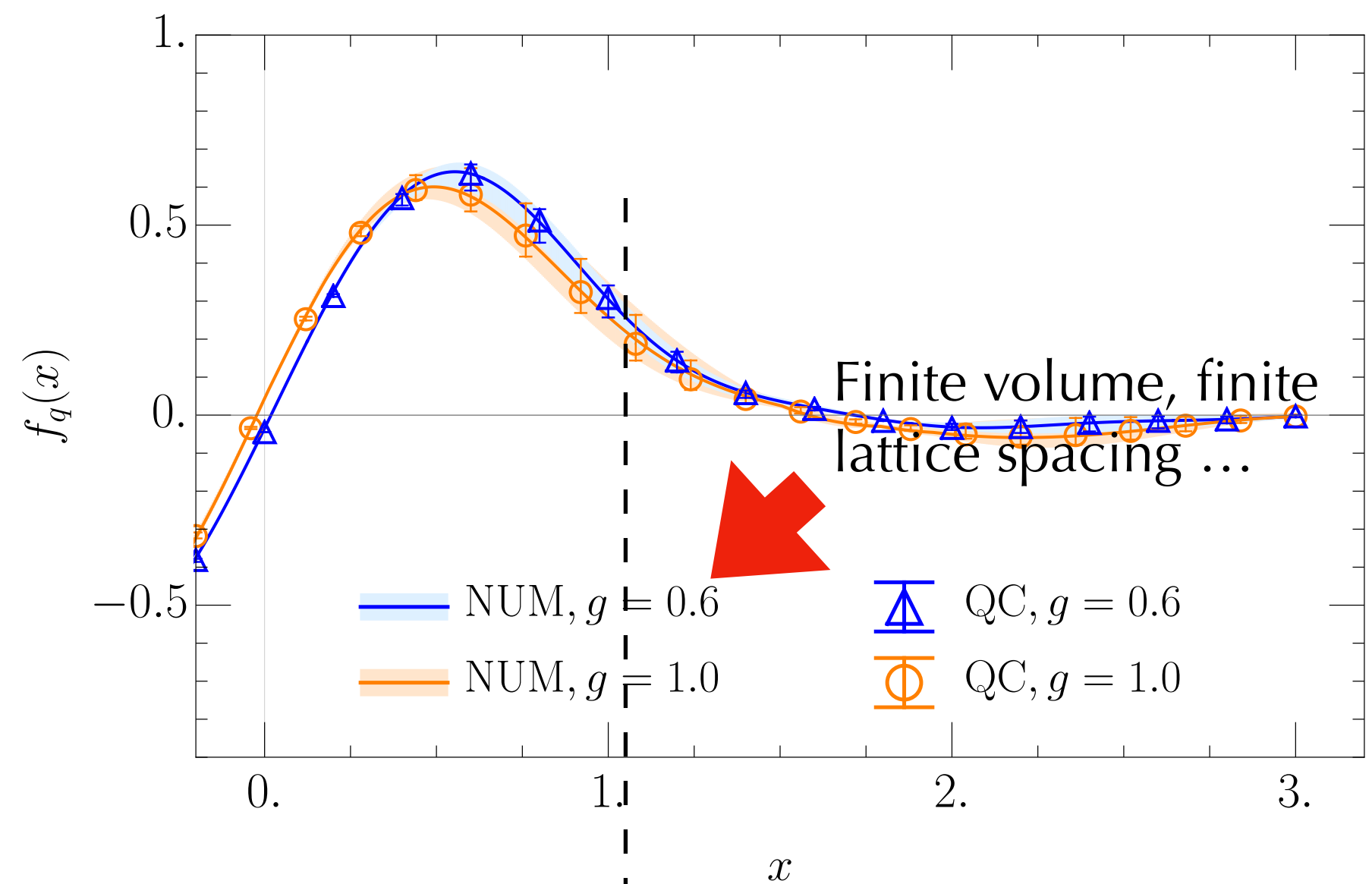
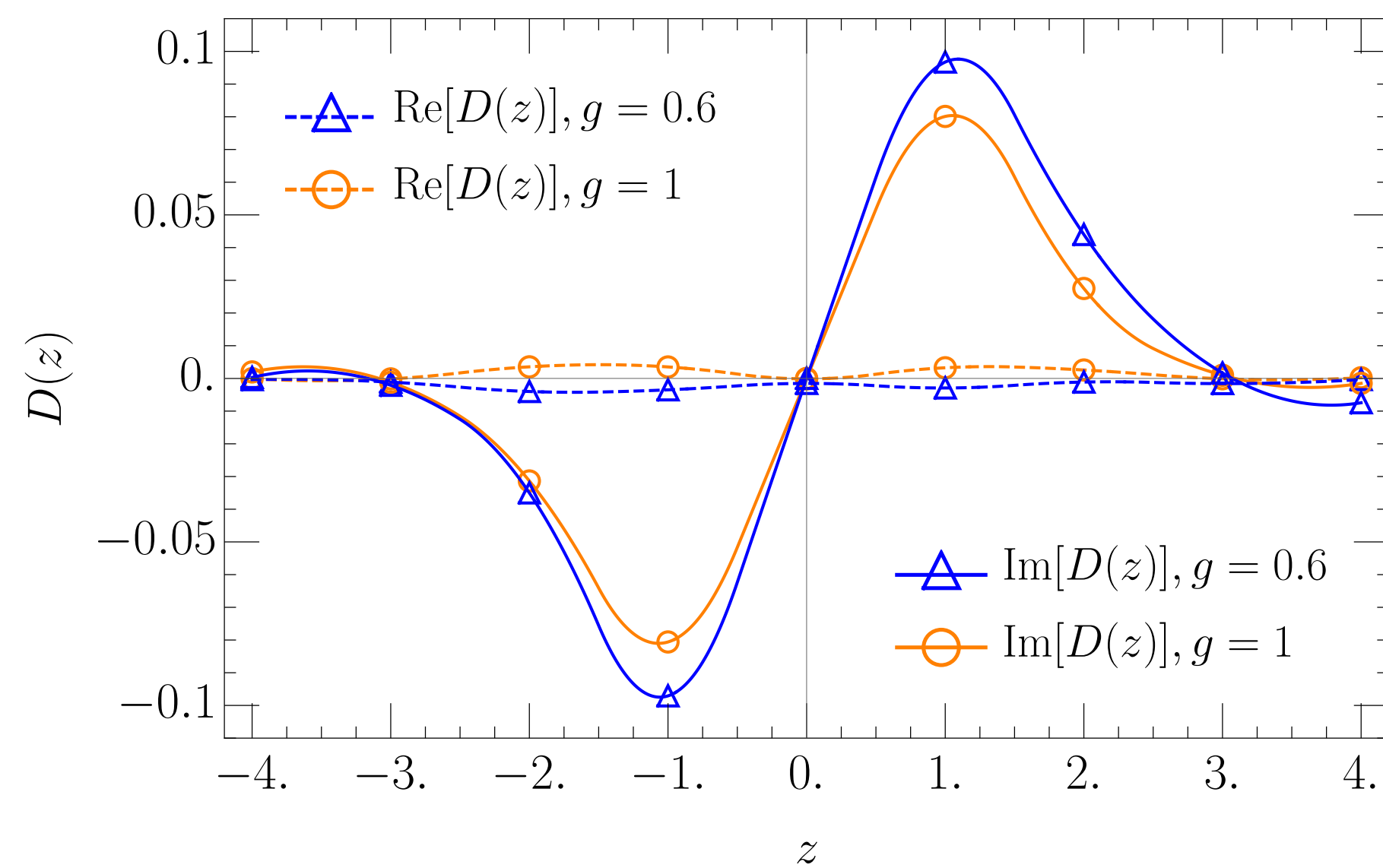
Fermion doubling!!



fields will be
represented by
a set of gates.

Li, et al, PRD letter 22

A toy model: Results



$${}_H \langle h | T[\psi_H(x) \psi_H(y)] | h \rangle_H =$$

The diagram shows two Feynman diagrams representing the two-point function $T[\psi_H(x) \psi_H(y)]$. The first diagram is a loop diagram with two external legs labeled x and y . The second diagram is a tree-level diagram with two external legs labeled x and y . The diagrams are connected by a plus sign and a multiplication sign.

$$f(x) \rightarrow \sum_{i,j} \sum_z \frac{1}{4\pi} e^{-ixM_h z} \langle h | e^{iHz} \phi_{-2z+i}^\dagger e^{-iHz} \phi_j | h \rangle$$

量子计算在模式识别中的作用



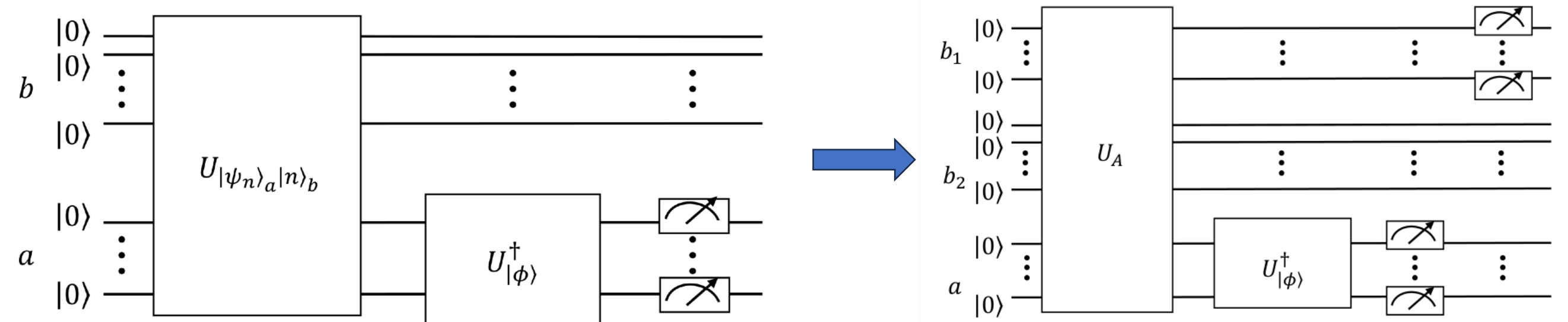
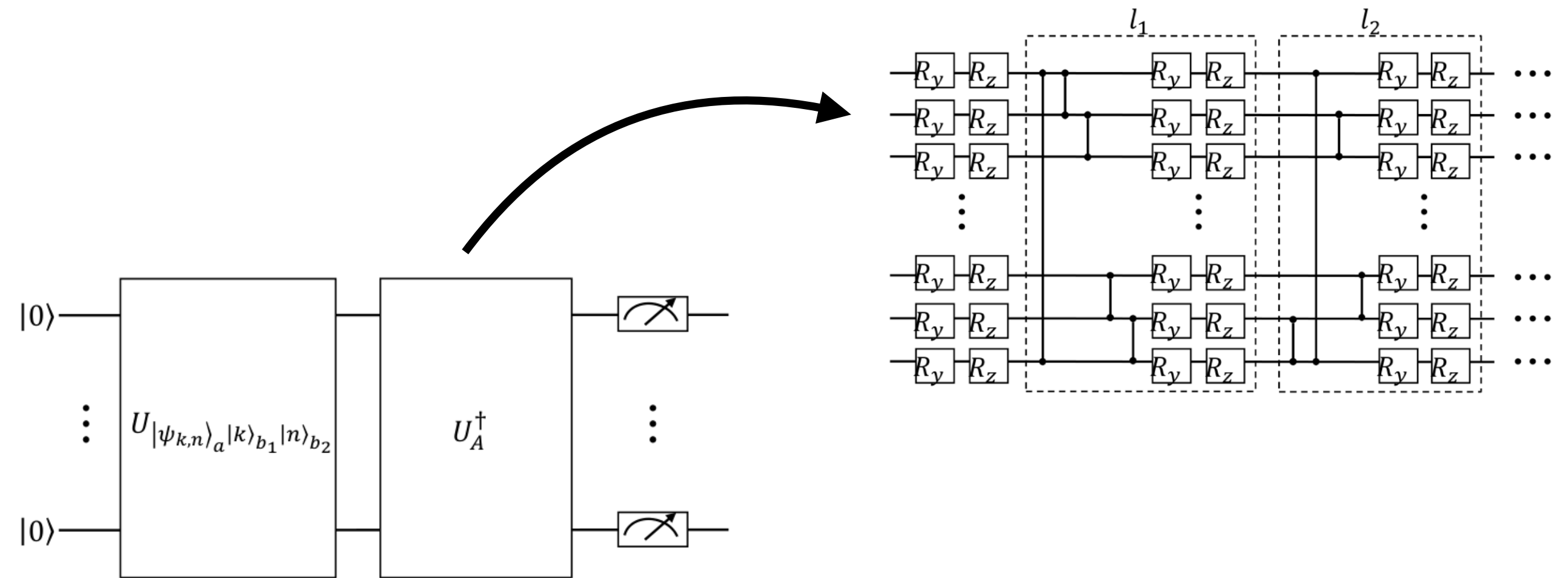
Using machine learning method suitable for quantum computing in the phenomenological study of new physics

Ji-Chong Yang

(yangjichong@lnnu.edu.cn)

In preparation, collaborator: Chong-Xing Yue

Variational quantum classifier



量子信息

粒子物理散射和衰变过程中的纠缠

Quantum Entanglement in High Energy Physics

Yu Shi (施郁)

2024.7.9.

28th LHC MiniWorkshop



Quantum Entanglement at High-energy Colliders

Kun Cheng 程焜
28th Mini-workshop on the frontier of LHC

粒子物理散射和衰变过程中的纠缠

Bell non-locality

Fabbrichesi et al,
PRL 127, 161801(2021)

∩

Quantum Steering

Afki et al,
PRL 130, 11801(2023)

∩

Quantum Entanglement

Afki et al,
EPJP 136, 907(2021)

∩

Quantum Discord

Afki et al,
PRL 130, 11801(2023)

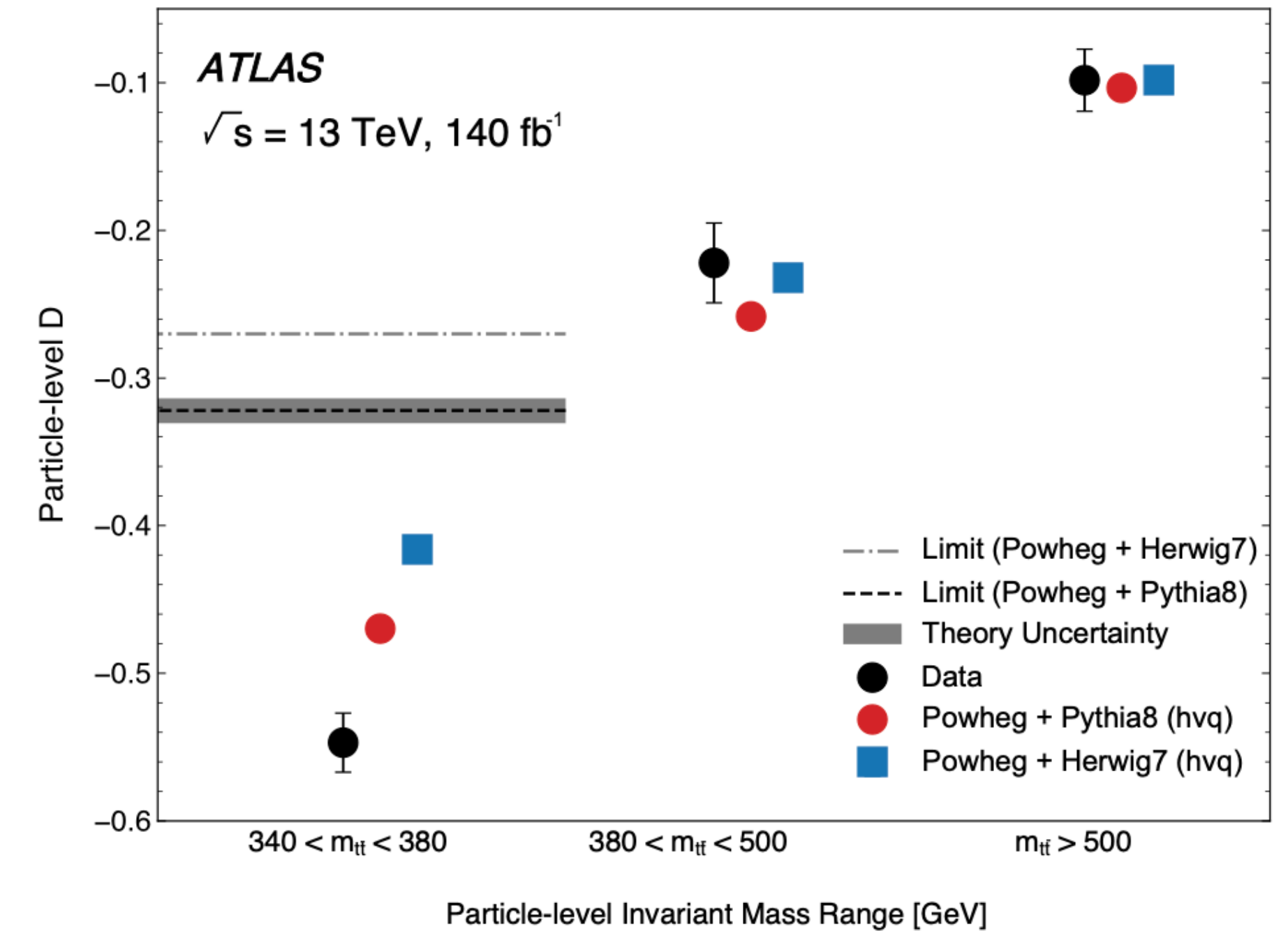


Fig: $D < -1/3$ indicates entanglement
[ATLAS, 2311.07288]

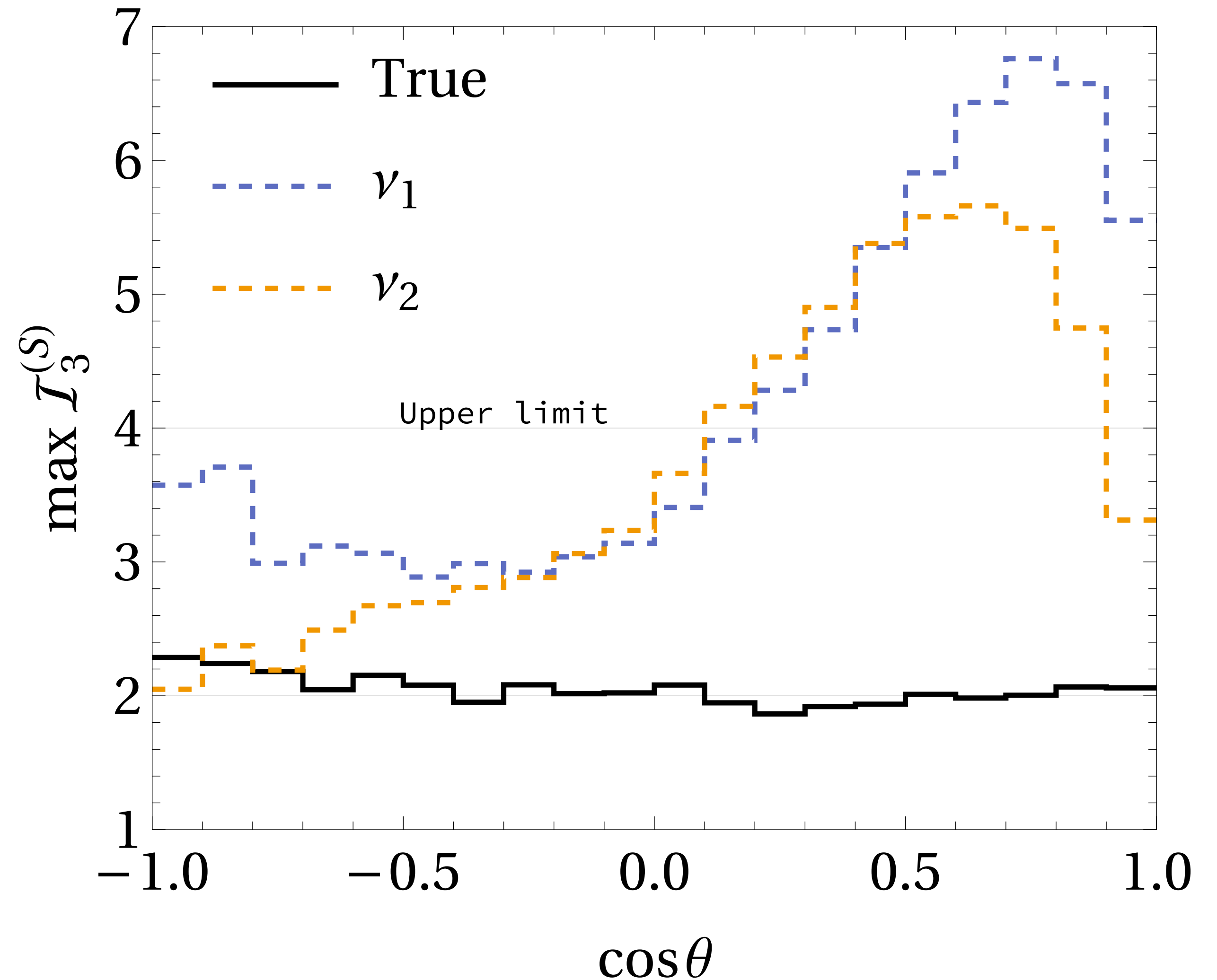
- Test the principle of QM at the highest energy we can achieve
- The entanglement between fundamental particles: sensitive to NP?
- The entanglement of unstable particles: some new properties?
- *More than spin correlation*

贝尔不等式破坏的检验

希格斯工厂上的W玻色子对产生过程

- W^\pm 的衰变：
 - 轻子衰变道：100%自旋关联可知自旋投影方向，但伴随中微子丢失能量；
 - 强子衰变道：没有丢失能量，但无法确定自旋投影方向（无法测量夸克电荷）
- 唯象学困难：若利用自旋关联确定 W^\pm 的自旋投影方向，需要 W^\pm 均轻子衰变，此时末态有两个中微子，无法重建 W^\pm 的四动量（重根）。

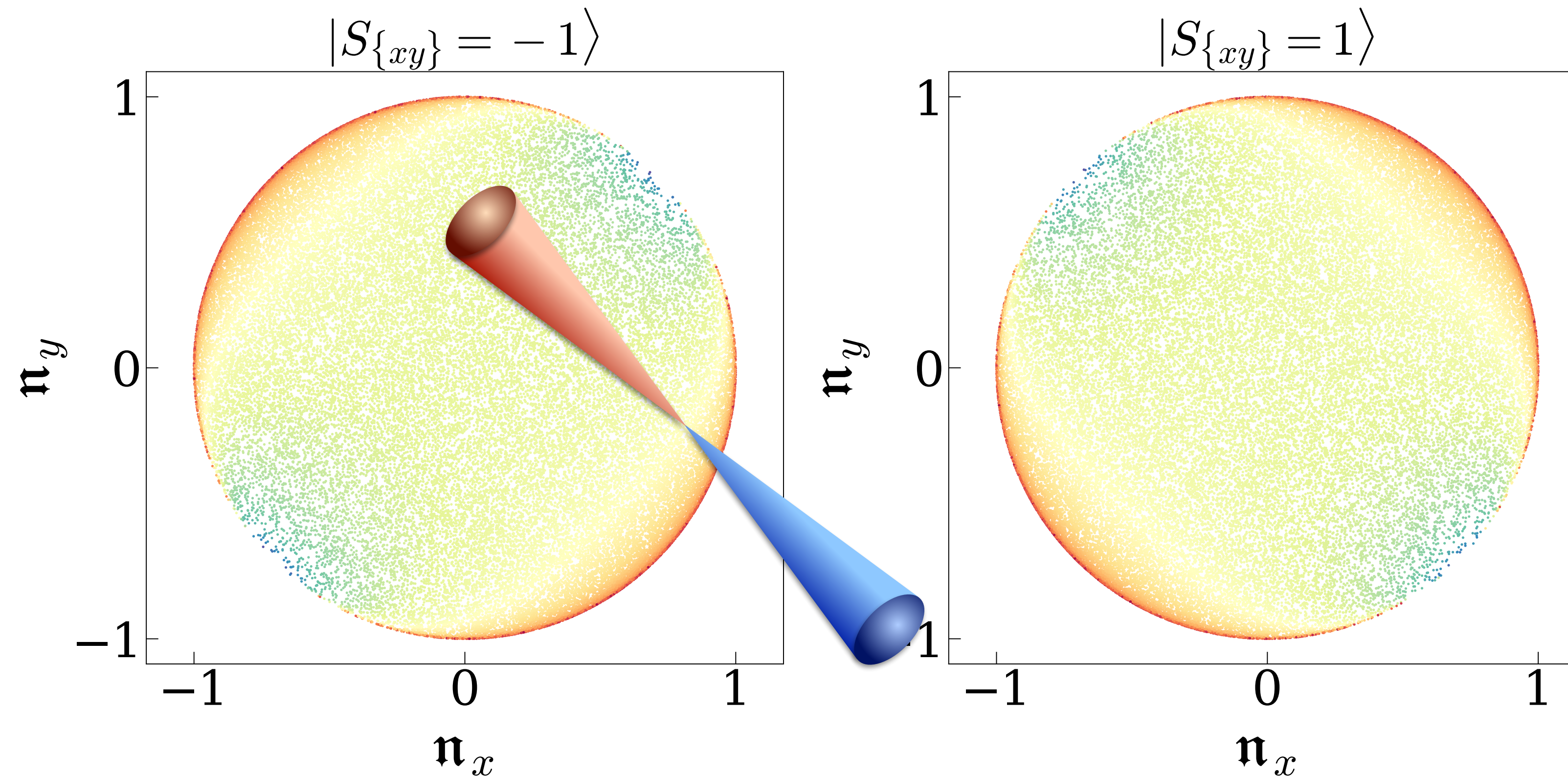
错误选取中微子动量将导致“超越”量子力学的结果



贝尔不等式破坏的检验

希格斯工厂上的W玻色子对产生过程

- 新贝尔观测量： W^+ 玻色子在其静止系中的自旋状态和 W^- 玻色子在其静止系中的线偏振状态。
- 好处： W^+ 玻色子轻子衰变， W^- 玻色子强子衰变，末态可重建。



$$\begin{aligned}\vec{\epsilon}_{|S_{\{xy\}}=-1\rangle} &= \frac{1}{\sqrt{2}}(1, 1, 0), \\ \vec{\epsilon}_{|S_{\{xy\}}=1\rangle} &= \frac{1}{\sqrt{2}}(1, -1, 0), \\ \vec{\epsilon}_{|S_{\{xy\}}=0\rangle} &= (0, 0, 1),\end{aligned}$$

详见张昊的报告

Emergent Symmetry from Entanglement Suppression

Ming-Lei Xiao

Sun Yat-Sen University



with M. Carena, I. Low, and C.E. Wagner [arXiv:2209.00198]

July 9, 2024 @ 28th Mini-workshop on the frontier of LHC, Tonghua

It from (qu)bit

J. A. Wheeler: Every **it**—every particle, every field of force, even the space-time continuum itself— derives its function, its meaning, its very existence entirely—even if in some contexts indirectly— **from** the apparatus-elicited answers to yes-or-no questions, binary choices, **bits**.



it	from	(qu)bit
Symmetry		Entanglement
Black Hole		Complexity
(OTOC) Correlation		Quantum Chaos
Holography		Error-Correcting Code
⋮		⋮

Information-theoretic properties may provide insights on the origin of physical principles.

Two-Higgs-Doublet Model

Two flavors of $SU(2)_L$ doublet $\Phi_a = (\underbrace{\Phi_{a=1,2}^+}_{p_{\uparrow,\downarrow}}, \underbrace{\Phi_{a=1,2}^0}_{n_{\uparrow,\downarrow}})$.

$$\begin{aligned} \mathcal{V}(\Phi_1, \Phi_2) = & m_1^2 \Phi_1^\dagger \Phi_1 + m_2^2 \Phi_2^\dagger \Phi_2 - \left[m_{12}^2 \Phi_1^\dagger \Phi_2 + \text{h.c.} \right] \\ & + \frac{\lambda_1}{2} (\Phi_1^\dagger \Phi_1)^2 + \frac{\lambda_2}{2} (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) \\ & + \left[\frac{\lambda_5}{2} (\Phi_1^\dagger \Phi_2)^2 + \lambda_6 (\Phi_1^\dagger \Phi_1) (\Phi_1^\dagger \Phi_2) + \lambda_7 (\Phi_2^\dagger \Phi_2) (\Phi_1^\dagger \Phi_2) + \text{h.c.} \right] . \end{aligned}$$

Consider tree-level scattering $\mathcal{S}(\Phi_a^+, \Phi_b^0 \rightarrow \Phi_c^+, \Phi_d^0) \equiv 1 + iM_{ab,cd} \delta^{(4)}(p)$

$$\text{Unbroken Phase: } M_{ab,cd} = \begin{pmatrix} \lambda_1 & \lambda_6^* & \lambda_6^* & \lambda_5^* \\ \lambda_6 & \lambda_3 & \lambda_4 & \lambda_7^* \\ \lambda_6 & \lambda_4 & \lambda_3 & \lambda_7^* \\ \lambda_5 & \lambda_7 & \lambda_7 & \lambda_2 \end{pmatrix} \begin{matrix} 11 \\ 12 \\ 21 \\ 22 \end{matrix}$$

Carena, Low, Wagner and Xiao [2307.08112]

11 12 21 22

Entanglement suppression in the broken phase impose constraints even on the spectrum, leading to the maximal symmetry in 2HDM.

Enhanced Symmetry in the Unbroken Phase

The Bose symmetry for the s-wave amplitude imposes $\vec{\alpha} = \vec{\beta} = \vec{r}$

Enhanced Symmetry: $SO(2)$ rotation along \vec{r}

Redefine $\Phi'_a = U_a{}^b \Phi_b$ such that $U \in SU(2)$ brings $\vec{r} \parallel \hat{z}$.

$$\mathcal{V}(\Phi'_1, \Phi'_2) = \dots + \frac{\lambda'_1}{2} (\Phi'_1{}^\dagger \Phi'_1)^2 + \frac{\lambda'_2}{2} (\Phi'_2{}^\dagger \Phi'_2)^2 + \lambda_3 (\Phi'_1{}^\dagger \Phi'_1) (\Phi'_2{}^\dagger \Phi'_2)$$

- $\Phi'_{1,2}$ may have independent phase symmetries $e^{i\phi_0}$ and $e^{i\phi_z \sigma^z}$.
- In the original basis, $U^{-1} e^{i\phi_z \sigma^z} U$ is the new $SO(2)$ rotation around \vec{r} .

Emergent Maximal Symmetry

A last chance: if the two charged scalars H_1^+ and H_2^+ are degenerate

$$P_1^s = P_2^s, \quad P_1^u = P_2^u \quad \Rightarrow \quad m_{H^+}^2 = 0, \quad Y_2 = -\frac{Z_3}{2} v^2,$$

we only need the combinations $M_1^s + M_2^s$ and $M_1^u + M_2^u$ to satisfy the entanglement suppression condition!

$$\begin{pmatrix} Z_1^2 + Z_6^2 & Z_1 Z_6 & (Z_1 + Z_3) Z_6 & Z_6^2 \\ Z_1 Z_6 & Z_6^2 & Z_6^2 & 0 \\ (Z_1 + Z_3) Z_6 & Z_6^2 & Z_6^2 + Z_3^2 & Z_3 Z_6 \\ Z_6^2 & 0 & Z_3 Z_6 & Z_6^2 \end{pmatrix} \Rightarrow Z_1 = Z_3 \text{ and } Z_6 = 0$$

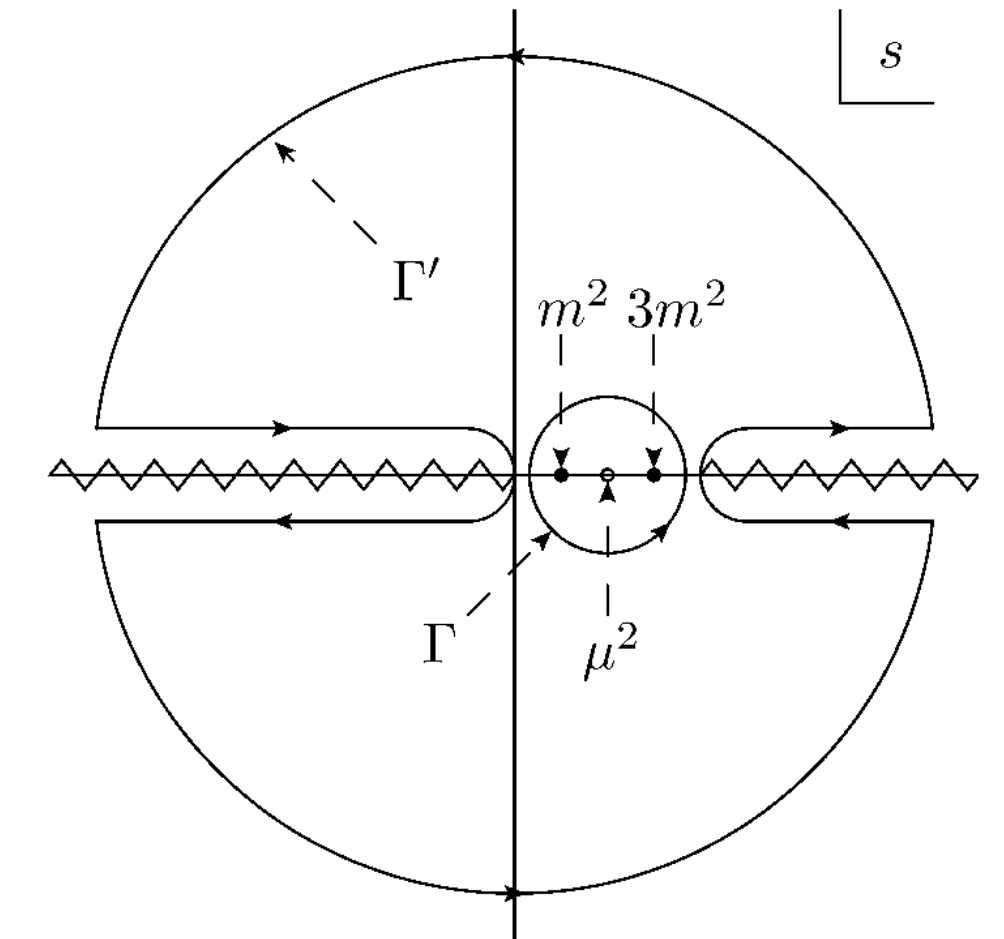
The resulting potential: $\mathcal{V} = \frac{Z_1}{2} \left[H_1^\dagger H_1 + H_2^\dagger H_2 - \frac{v^2}{2} \right]^2$ has $SO(8)$ symmetry!

Positivity from elastic scattering

$$A_{2 \rightarrow 2}(s, t = 0) = c_0 + c_2 s^2 + c_4 s^4 + \dots$$

◆ Analyticity: $f = \frac{1}{2\pi i} \oint_{\Gamma} ds \frac{A(s, 0)}{(s - \mu^2)^3}$

◆ Unitarity + Locality: $A(s, 0) < \mathcal{O}(s \ln^2 s)$



[Cheung, Remmen, 1601.04068]

$$f = \frac{1}{2\pi i} \oint_{\Gamma} ds \frac{A(s, 0)}{(s - \mu^2)^3} = \frac{1}{2\pi i} \left(\int_{-\infty}^0 + \int_{4m^2}^{\infty} \right) ds \frac{\text{Disc} A(s, 0)}{(s - \mu^2)^3}$$

IR ↑

Calculable
in EFT

↑ UV

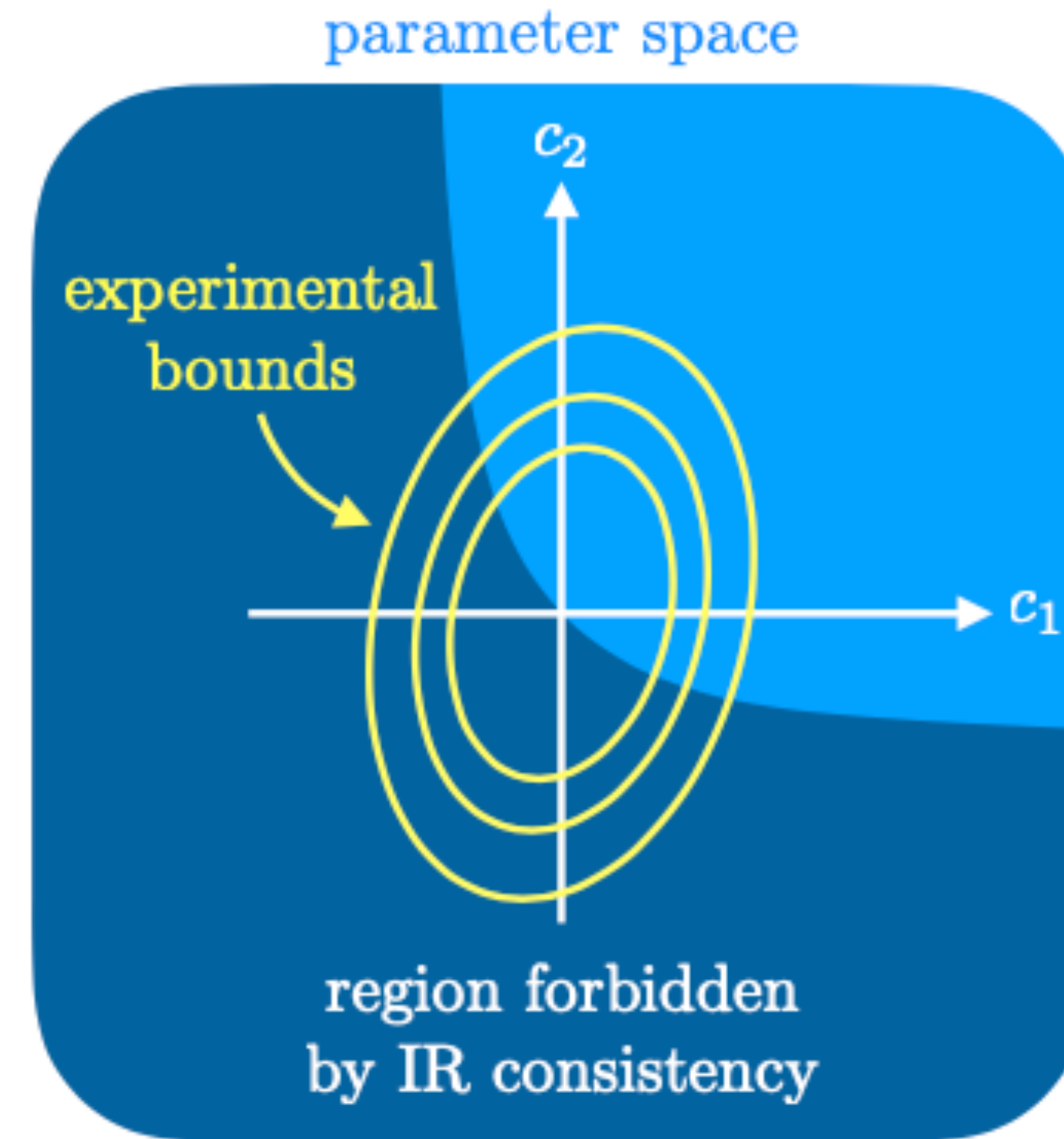
Disc > 0 by
optical theorem
+ (s-u) crossing

$$A_{2 \rightarrow 2}(s, t = 0) = c_0 + c_2 s^2 + c_4 s^4 + \dots$$

Positivity bounds in SM EFT

Dim-8 operators in $SU(N)$ gauge theory

$$\begin{aligned}
 \mathcal{O}_1^{F^4} & (F^a F^a)(F^b F^b) \\
 \mathcal{O}_2^{F^4} & (F^a \tilde{F}^a)(F^b \tilde{F}^b) \\
 \mathcal{O}_3^{F^4} & (F^a F^b)(F^a F^b) \\
 \mathcal{O}_4^{F^4} & (F^a \tilde{F}^b)(F^a \tilde{F}^b) \\
 \mathcal{O}_5^{F^4} & d^{abe} d^{cde} (F^a F^b)(F^c F^d) \\
 \mathcal{O}_6^{F^4} & d^{abe} d^{cde} (F^a \tilde{F}^b)(F^c \tilde{F}^d) \\
 \mathcal{O}_7^{F^4} & d^{ace} d^{bde} (F^a F^b)(F^c F^d) \\
 \mathcal{O}_8^{F^4} & d^{ace} d^{bde} (F^a \tilde{F}^b)(F^c \tilde{F}^d) \\
 \tilde{\mathcal{O}}_1^{F^4} & (F^a F^a)(F^b \tilde{F}^b) \\
 \tilde{\mathcal{O}}_2^{F^4} & (F^a F^b)(F^a \tilde{F}^b) \\
 \tilde{\mathcal{O}}_3^{F^4} & d^{abe} d^{cde} (F^a F^b)(F^c \tilde{F}^d) \\
 \tilde{\mathcal{O}}_4^{F^4} & d^{ace} d^{bde} (F^a F^b)(F^c \tilde{F}^d)
 \end{aligned}$$



$$3c_1^{G^4} + 3c_3^{G^4} + c_5^{G^4} > 0$$

$$3c_3^{G^4} + 2c_5^{G^4} > 0$$

$$3c_2^{G^4} + 3c_4^{G^4} + c_6^{G^4} > 0$$

$$3c_4^{G^4} + 2c_6^{G^4} > 0$$

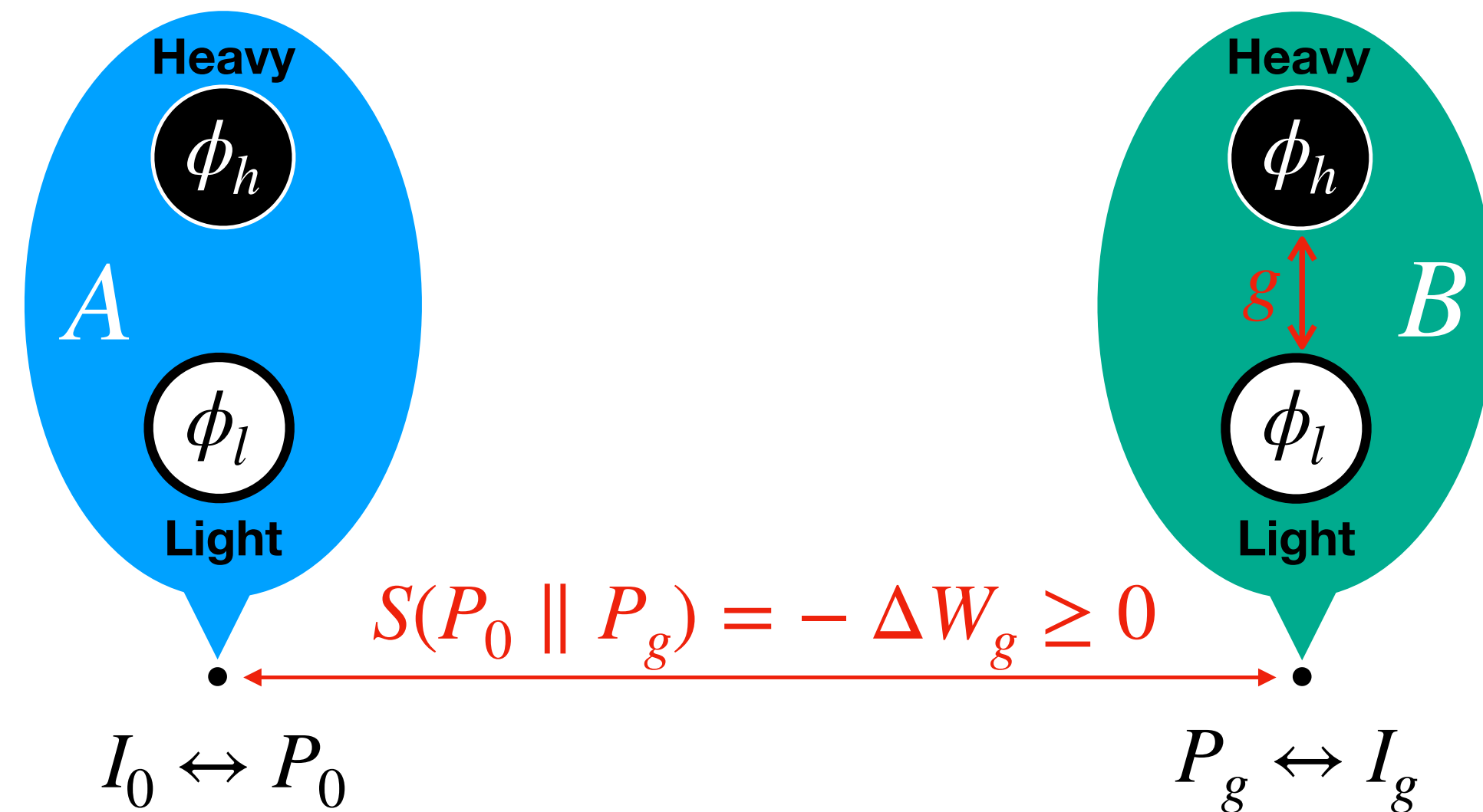
$$(3\tilde{c}_1^{G^4} + 3\tilde{c}_2^{G^4} + \tilde{c}_3^{G^4})^2 < 4(3c_1^{G^4} + 3c_3^{G^4} + c_5^{G^4})(3c_2^{G^4} + 3c_4^{G^4} + c_6^{G^4})$$

$$(3\tilde{c}_2^{G^4} + 2\tilde{c}_3^{G^4})^2 < 4(3c_3^{G^4} + 2c_5^{G^4})(3c_4^{G^4} + 2c_6^{G^4}).$$

For SU(3) gauge theory

Relative entropy

We defined a **distance** between two different theories in the approach of information theory



QHC, Ueda, PhysRevD.108.025011
 QHC, Kan, Ueda, JHEP07(2023)111)

I_0 without interaction b/w heavy and light

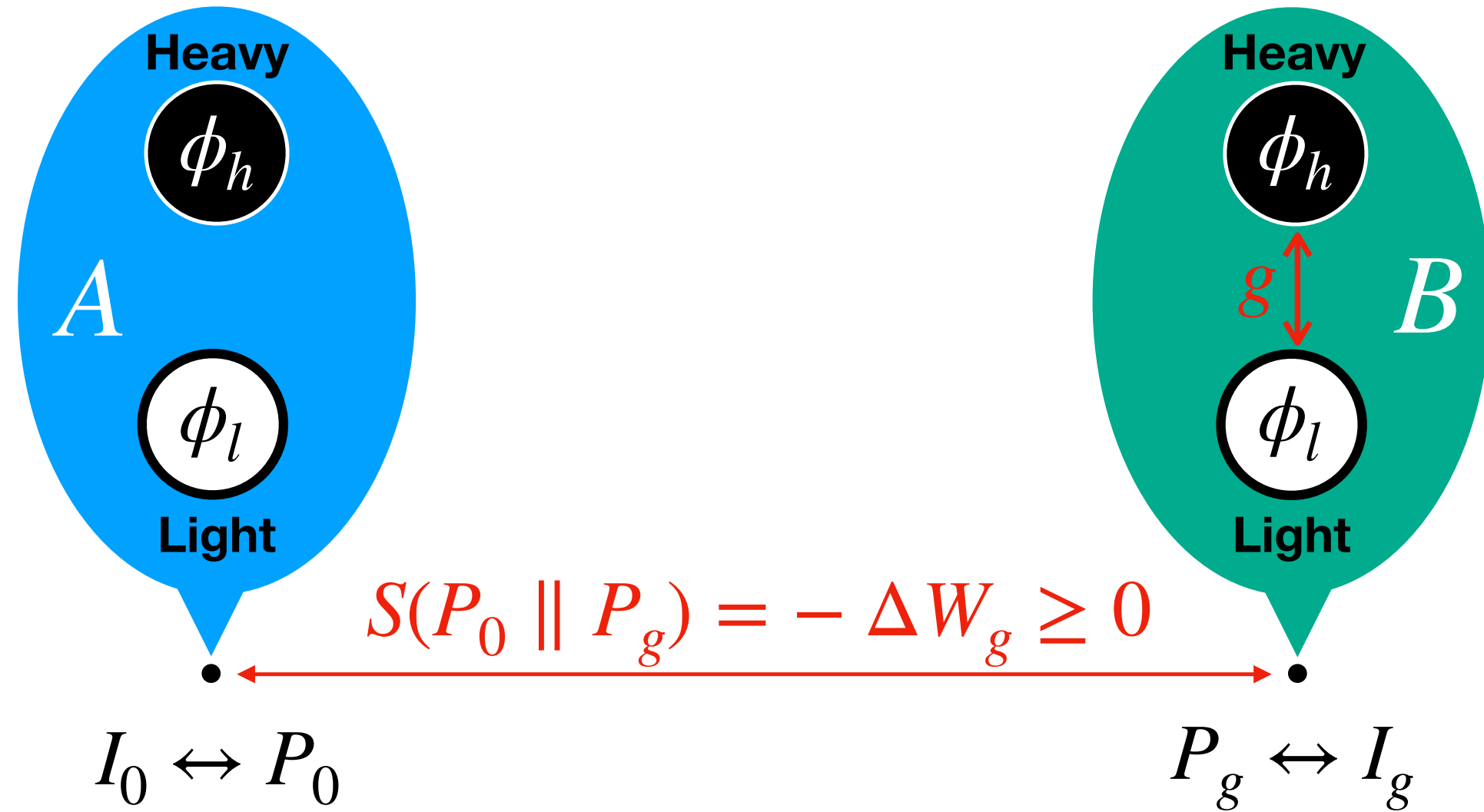
$I_g = I_0 + g \cdot I_I$ with the interaction

Distance defined by relative entropy

$$S(P_0 \parallel P_g) = \int d[\phi_h] \left(P_0 \ln P_0 - P_0 \ln P_g \right) \geq 0$$

$$P_0 = e^{-I_0/Z_0} \quad P_g = e^{-I_g/Z_g}$$

Non-negativity of relative entropy



U(1) and SU(2) bounds are the **same** as positivity bounds from unitarity and causality

- $U(1)_Y$:

$$c_1^{B^4} \geq 0, \quad c_2^{B^4} \geq 0, \quad 4c_1^{B^4} c_2^{B^4} \geq (\tilde{c}_1^{B^4})^2,$$

- $SU(2)_L$:

$$c_1^{W^4} + c_3^{W^4} \geq 0, \quad c_2^{W^4} + c_4^{W^4} \geq 0, \quad 4(c_1^{W^4} + c_3^{W^4})(c_2^{W^4} + c_4^{W^4}) \geq (\tilde{c}_1^{W^4} + \tilde{c}_2^{W^4})^2,$$

We show that the positive distance yields

- Positivity bounds on SMEFT dim-8 gauge bosonic operators
- WGC-like behavior in extremal relation holds in wide class of black holes
- Any UV theory violating second law of thermodynamics yields pathological EFTs

SU(3) bounds are **stronger** than positivity bounds from unitarity and causality

- $SU(3)_C$:

$$2c_1^{G^4} + c_3^{G^4} \geq 0, \quad 3c_2^{G^4} + 2c_5^{G^4} \geq 0,$$

$$4(3c_1^{G^4} + 3c_3^{G^4} + c_5^{G^4})(3c_2^{G^4} + 3c_4^{G^4} + c_6^{G^4}) \geq (3\tilde{c}_1^{G^4} + 3\tilde{c}_2^{G^4} + \tilde{c}_3^{G^4})^2$$

$$4(3c_3^{G^4} + 2c_5^{G^4})(3c_4^{G^4} + 2c_6^{G^4}) \geq (3\tilde{c}_2^{G^4} + 2\tilde{c}_3^{G^4})^2$$

结论

高能量

新重粒子
干涉效应
有效场论

新物理

高亮度

精确检验

新技术和新视角

人工智能和高能物理的结合
量子计算和场论的结合
量子信息和高能物理的结合

谢谢!

Theoretical Advanced Study Institute in Elementary Particle Physics (TASI)

TASI 2024

The Frontiers of Particle Theory

University of Colorado Boulder | June 3–28, 2024

ORGANIZERS: NATHANIEL CRAIG (UC SANTA BARBARA), TONGYAN LIN (UC SAN DIEGO), JESSE THALER (MIT)

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- T.-T. Yu: Dark Matter: Evidence and Opportunities