

高圈多外线费曼积分的解析计算

2024 全国粒子物理学术会议
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Based on

Henn, Peraro, Xu, YZ, *JHEP* 03 (2022) 056

Henn, Matijasic, Miczajka, Peraro, Xu, YZ, *JHEP* 08(2024) 027

Liu, Matijasic, Miczajka, Peraro, Xu, Xu, YZ, *to appear*

also the package ...

“**NeatIBP 1.0**, a package generating small-size integration-by-parts relations for Feynman integrals”

Wu, Boehm, Ma, Xu, YZ, *Comput. Phys. Commun.* 295 (2024), 108999

Outline

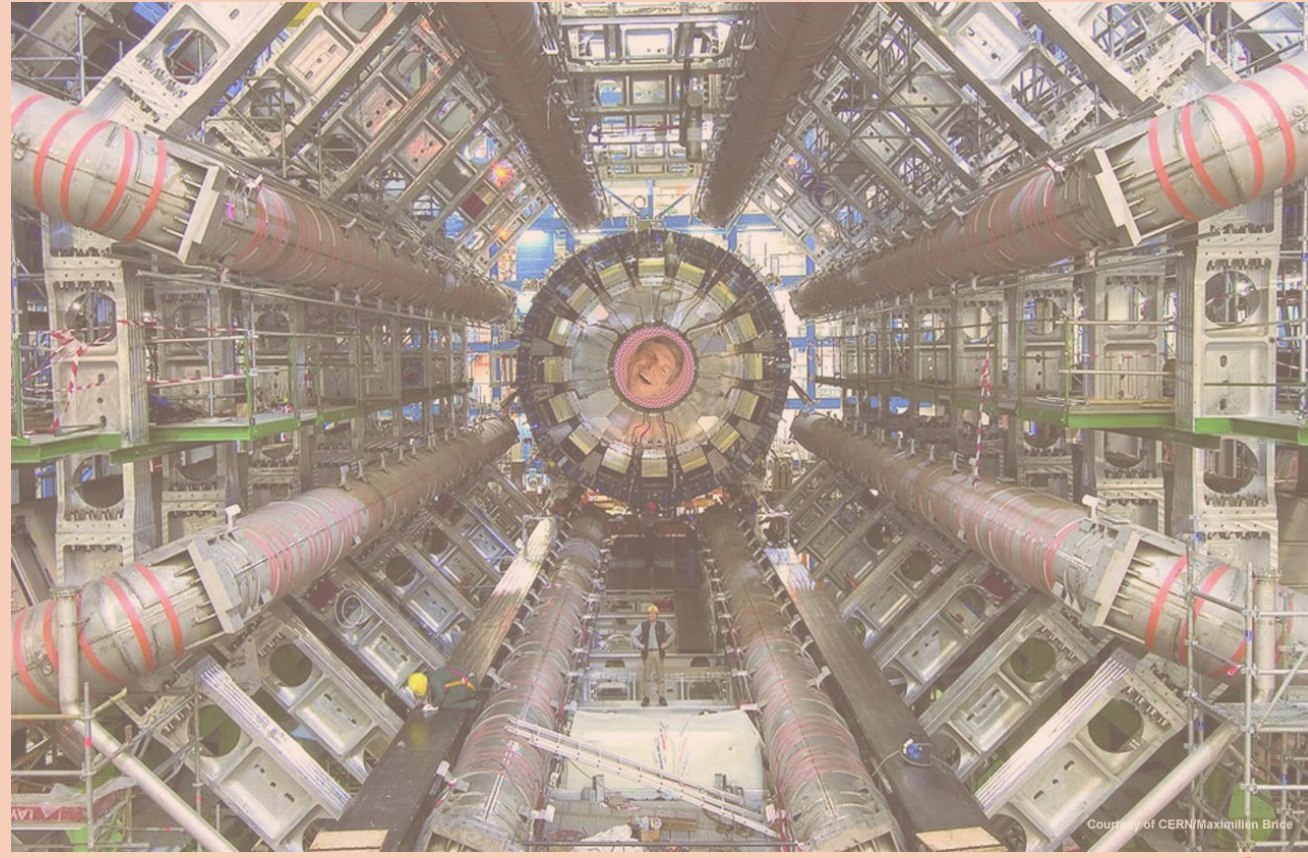
Why Feynman integrals? Why **analytic**?

Case 1: **2loop 6point Feynman integrals**

Case 2: **3loop 5point Feynman integrals**

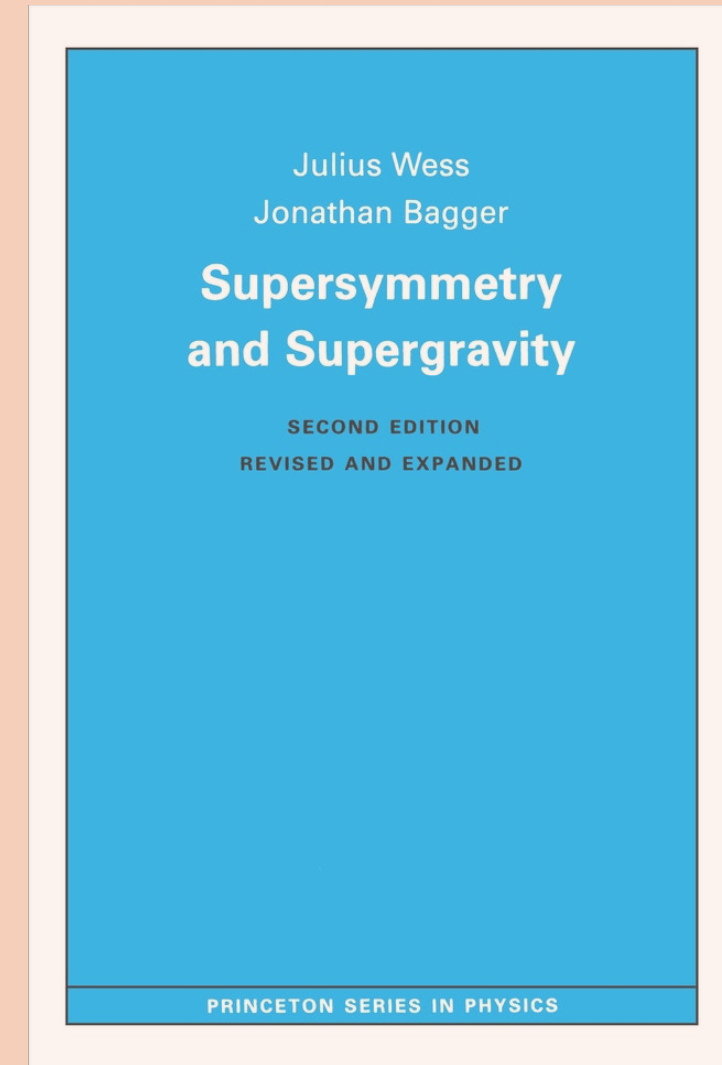
Summary and Outlook

Why Feynman integrals?



Precision physics

$$\sigma = \sigma^{LO} + \sigma^{NLO} + \sigma^{NNLO}$$



Formal theory

Feynman integrals



Gravitational wave template computations

Why analytic Feynman integrals?

- Auxiliary Mass Flow and Numeric Monte Carlo methods slow or not available yet for some multi-loop multi-leg Feynman integrals

for examples: **2loop 6point** and **3loop 5point** Feynman integrals

- Theoretical aspects of quantum field theory

for examples: 2loop Yang-Mills theory **collinear factorization violation**

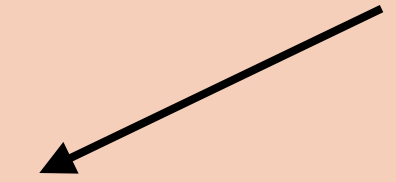
Henn, Ma, Xu, Yan, *YZ*, Zhu, arXiv 2406.14604

- Quantum field theory computation of gravitational wave

for the advance of AMFlow, refer to Yanqing Ma's plenary talk

Our Strategy

Guan, Liu, Ma, Wu, 2024



Integration-by-parts (IBP) reduction

Finite field techniques, Blade, NeatIBP

Uniformly transcendental (UT) basis determination
Canonical differential equation

$$\frac{\partial}{\partial x_i} I(x, \epsilon) = \epsilon A_i(x) I(x, \epsilon)$$

Alphabet searching

$$A_i = \frac{\partial}{\partial x_i} \tilde{A}, \quad \tilde{A} = \sum_k \tilde{a}_k \log(W_k)$$

Solving differential equation

$$\int d \log(W_{i_1}) \circ \dots \circ d \log(W_{i_k})$$

→ polylogarithm functions or one-fold integration

2loop 6point Feynman integrals

The status of art for analytic computations

Scale frontier

2loop 5point massless

Gehrmann, Henn, Lo Presti 2015

Chicherin, Gehrmann, Henn, Wasser, YZ, Zoia 2019

5 scales

2loop 5point one-mass

Papadopoulos, Tommasini, Wever 2019

Abreu, Ita, Moriello, Page, Tschernow, Zeng 2020

Abreu, Chicherin, Ita, Page, Sotnikov, Tschernow, Zoia 2023

6 scales

2loop 5point two-mass

Cordero, Figueiredo, Kraus, Page and Reina 2023

for leading-Color $pp \rightarrow ttH$ amplitudes with a light-quark loop

7 scales

2loop 6point massless

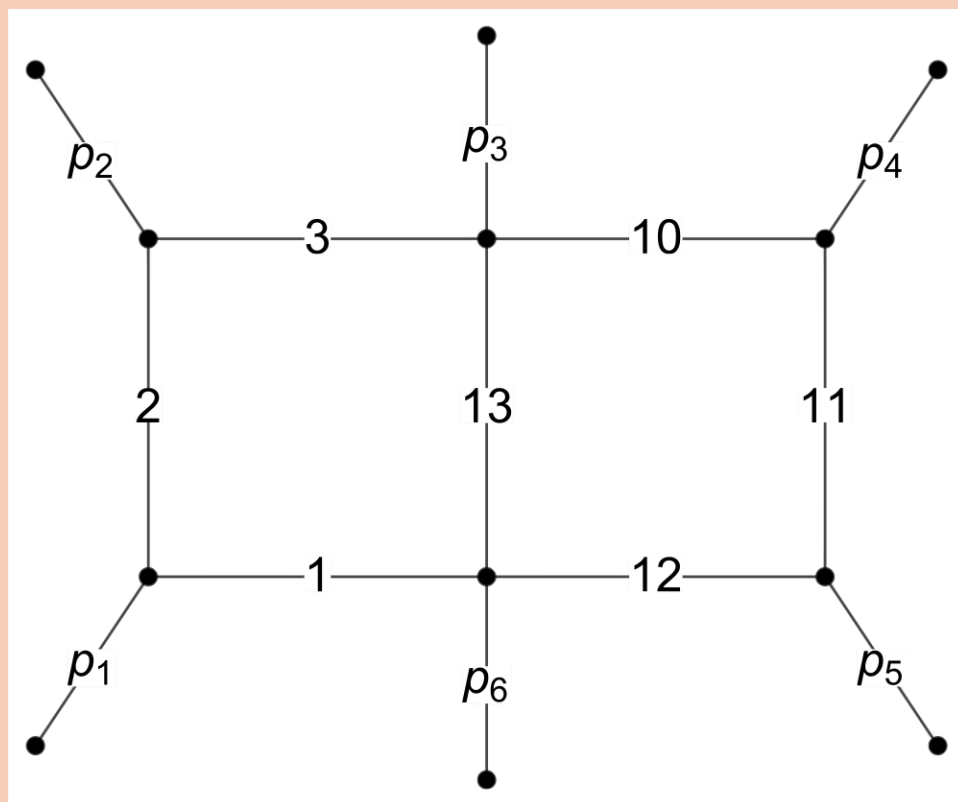
Henn, Matijasic, Miczajka, Peraro, Xu, YZ, 2024

for NNLO 4 jets production, 2 jets+ 2 photons

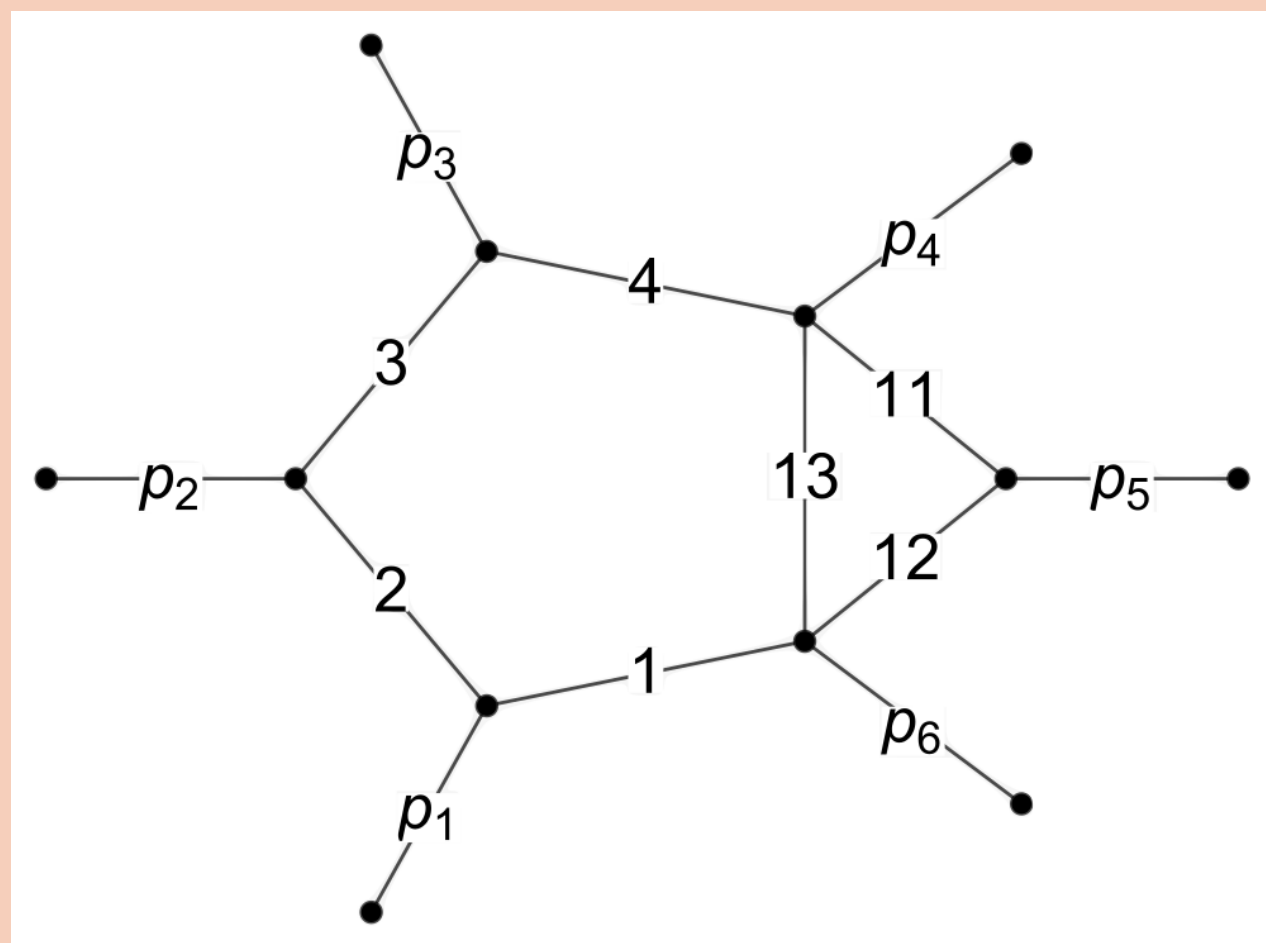
8 scales!

$S_{12}, S_{23}, S_{34}, S_{45}, S_{56}, S_{16}, S_{123}, S_{345}$

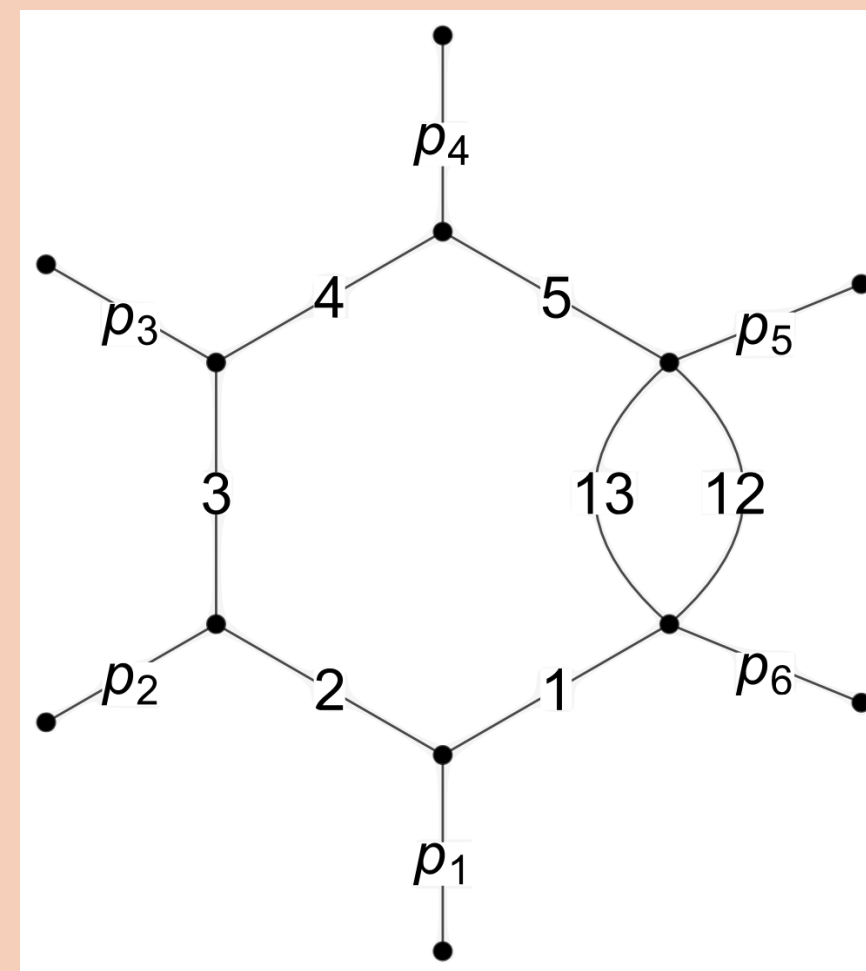
What we achieved



double box



pentagon triangle

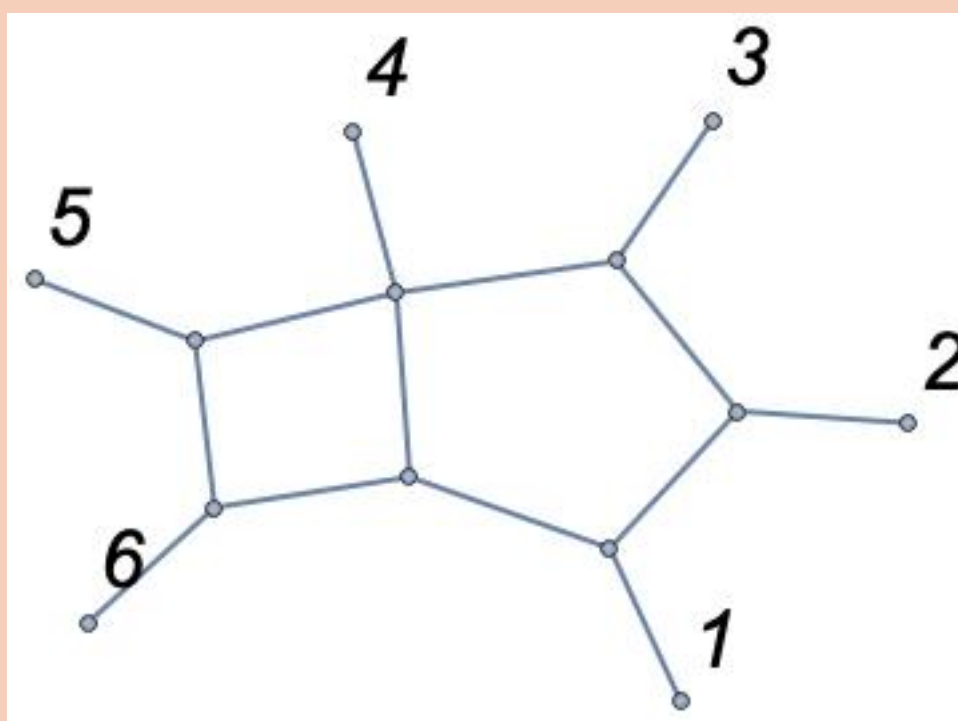


hexagon bubble

✓ UT basis found!

✓ Canonical differential equation solved!

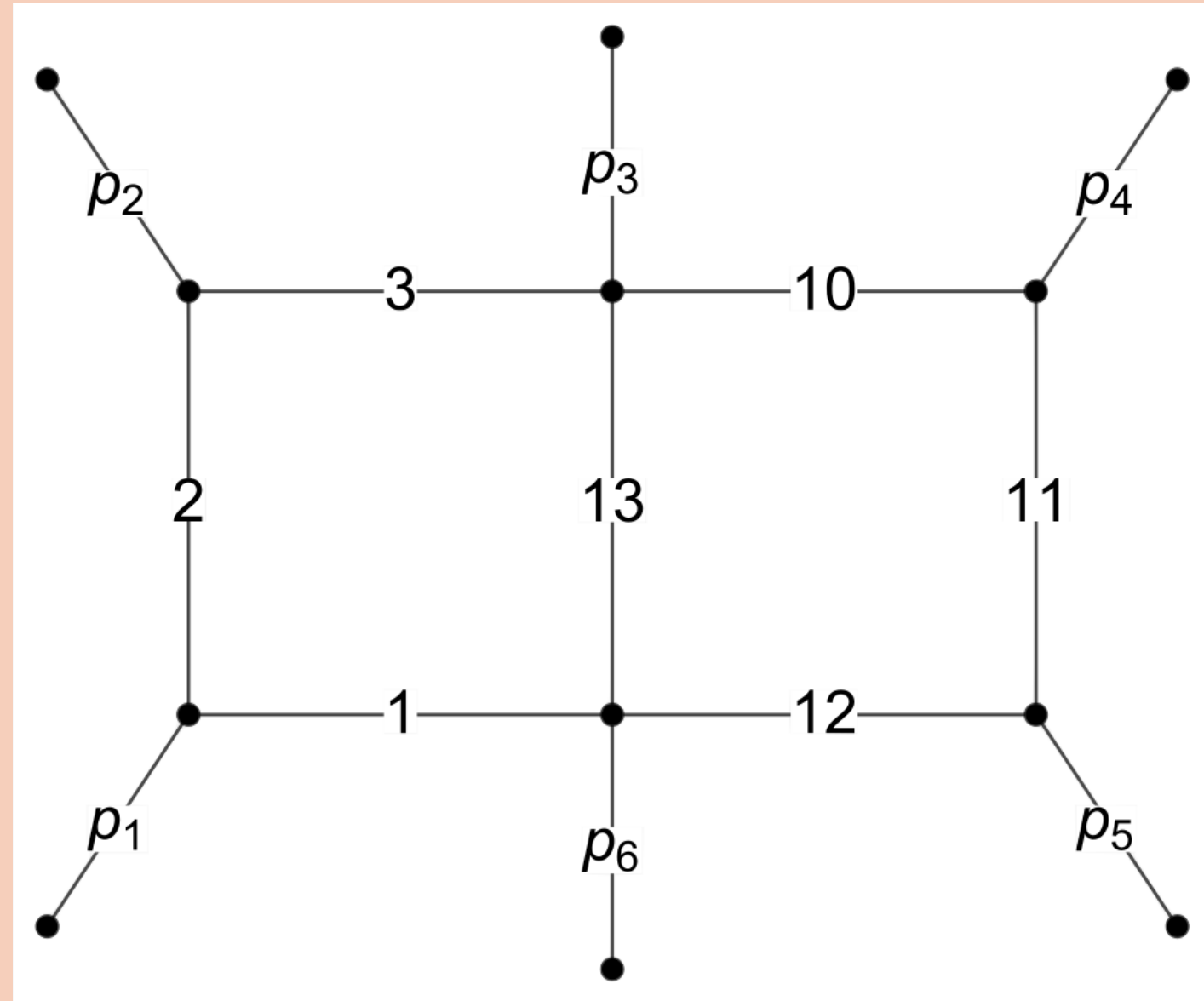
J. Henn, A. Matijasic, J. Miczajka, T. Peraro, Y. Xu, YZ, *JHEP08(2024)027*



canonical differential equation also solved

to appear as the next paper

Uniformal transcendental (UT) basis determination



Chiral numerator
 (Arkani-Hamed, Bourjaily, Cachazo, Trnka 2011)
 / Gram determinant
 correspondence

$$I_{\text{db},i} = \int \frac{d^{4-2\epsilon}l_1}{i\pi^{2-\epsilon}} \frac{d^{4-2\epsilon}l_2}{i\pi^{2-\epsilon}} \frac{N_i}{D_1 D_2 D_3 D_{10} D_{11} D_{12} D_{13}}, \quad i = 1, \dots, 7$$

$$N_1 = -s_{12}s_{45}s_{156},$$

$$N_2 = -s_{12}s_{45}(l_1 + p_5 + p_6)^2,$$

$$N_3 = \frac{s_{45}}{\epsilon_{5126}} G \begin{pmatrix} l_1 & p_1 & p_2 & p_5 + p_6 \\ p_1 & p_2 & p_5 & p_6 \end{pmatrix},$$

$$N_4 = \frac{s_{12}}{\epsilon_{1543}} G \begin{pmatrix} l_2 - p_6 & p_5 & p_4 & p_1 + p_6 \\ p_1 & p_5 & p_4 & p_3 \end{pmatrix},$$

$$N_5 = -\frac{1}{4} \frac{\epsilon_{1245}}{G(1, 2, 5, 6)} G \begin{pmatrix} l_1 & p_1 & p_2 & p_5 & p_6 \\ l_2 & p_1 & p_2 & p_5 & p_6 \end{pmatrix},$$

$$N_6 = \frac{1}{8} G \begin{pmatrix} l_1 & p_1 & p_2 \\ l_2 - p_6 & p_4 & p_5 \end{pmatrix} + \frac{D_2 D_{11} (s_{123} + s_{126})}{\delta},$$

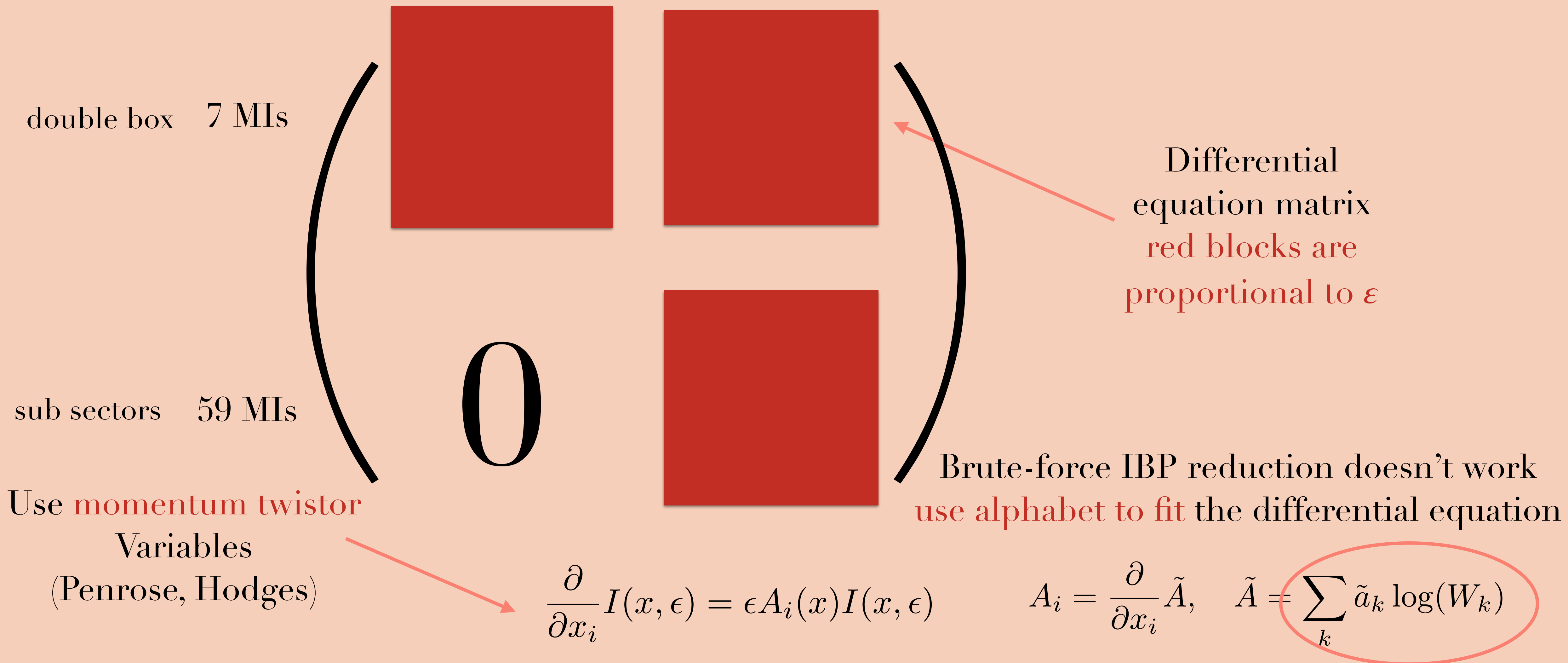
$$N_7 = -\frac{1}{2\epsilon} \frac{\Delta_6}{G(1, 2, 4, 5) D_{13}} G \begin{pmatrix} l_1 & p_1 & p_2 & p_4 & p_5 \\ l_2 & p_1 & p_2 & p_4 & p_5 \end{pmatrix}.$$

$$\Delta_6 = \langle 12 \rangle [23] \langle 34 \rangle [45] \langle 56 \rangle [61] - \langle 23 \rangle [34] \langle 45 \rangle [56] \langle 61 \rangle [12].$$

key step

J. Henn, A. Matijasic, J. Miczajka, T. Peraro, Y. Xu, YZ, JHEP08(2024)027

Complete canonical differential equation for 2l6p double box



A new algorithm to search for odd letters

Odd letter

$$\frac{P(s) - \sqrt{Q(s)}}{P(s) + \sqrt{Q(s)}}$$

$$P^2 - Q = c \prod_i W_i^{e_i}, \quad c \in \mathbb{Q}, \quad e_i \in \mathbb{N}$$

Even letter

An observation (and conjecture) from Heller, von Manteuffel, Schabinger 2020

Algorithm to solve for c, e_i Matijasic, J. Miczajka, to appear soon
package “Effortless”

Even letter, Odd letter and the more complicated ...

114 Even letter $F(s)$ a polynomial in Mandelstam variables and masses

Conjecture: a Feynman integrals' even letters are all from Landau singularity?

94 Odd letter $\frac{P(s) - \sqrt{Q(s)}}{P(s) + \sqrt{Q(s)}}$ $\log(W) \mapsto -\log(W)$ under the sign change of the square root

square roots are $\epsilon_{ijkl}, \Delta_6, \sqrt{\lambda(s_{12}, s_{34}, s_{56})}$

pseudo
scalar

leading
singularity
hexagon

Källin function
from massive triangle
diagrams

15 More
complicated
letter

$$\frac{P(s) - \sqrt{Q_1(s)}\sqrt{Q_2(s)}}{P(s) + \sqrt{Q_1(s)}\sqrt{Q_2(s)}}$$

Boundary Values

Numeric boundary values

It is fine to use the package `AMFlow` to get ~ 100 digits as the boundary value

Liu, Wang, Ma, 2018

Liu, Ma 2022

Analytic boundary values

It is still possible to *fully analytic* boundary values due to the kinematic symmetry

$$X_0 : \{s_{12}, s_{23}, s_{34}, s_{45}, s_{56}, s_{16}, s_{123}, s_{234}, s_{345}\} \rightarrow \{-1, -1, -1, -1, -1, -1, -1, -1, -1\}$$

UT integrals are not divergent at this point (spurious poles).

Solve the canonical DE on a curve starting with X_0 and require the finite solution

Some known integrals' boundary values

} analytic
boundary
value

Boundary Values

Analytic boundary values

Boundary values at the initial point, are combination of poly-logarithm of roots of unity

$$\epsilon^4 I_{\text{db},1}(X_0) = 1 + \frac{\pi^2}{6} \epsilon^2 + \frac{38}{3} \zeta_3 \epsilon^3 + \left(\frac{49\pi^4}{216} + \frac{32}{3} \text{Im} [\text{Li}_2(\rho)]^2 \right) \epsilon^4,$$

$$\epsilon^4 I_{\text{db},2}(X_0) = 1 + \frac{\pi^2}{6} \epsilon^2 + \frac{34}{3} \zeta_3 \epsilon^3 + \left(\frac{71\pi^4}{360} + 20 \text{Im} [\text{Li}_2(\rho)]^2 \right) \epsilon^4,$$

$$I_{\text{db},3}(X_0) = I_{\text{db},4}(X_0) = I_{\text{db},5}(X_0) = 0,$$

$$\epsilon^4 I_{\text{db},6}(X_0) = - \left(\frac{\pi^4}{540} + \frac{4}{3} \text{Im} [\text{Li}_2(\rho)]^2 \right) \epsilon^4,$$

$$\epsilon^4 I_{\text{db},7}(X_0) = 0.$$

from the differential equation
spurious pole asymptotic analysis

$$\rho = \frac{1}{2} + \frac{\sqrt{3}}{2}i$$

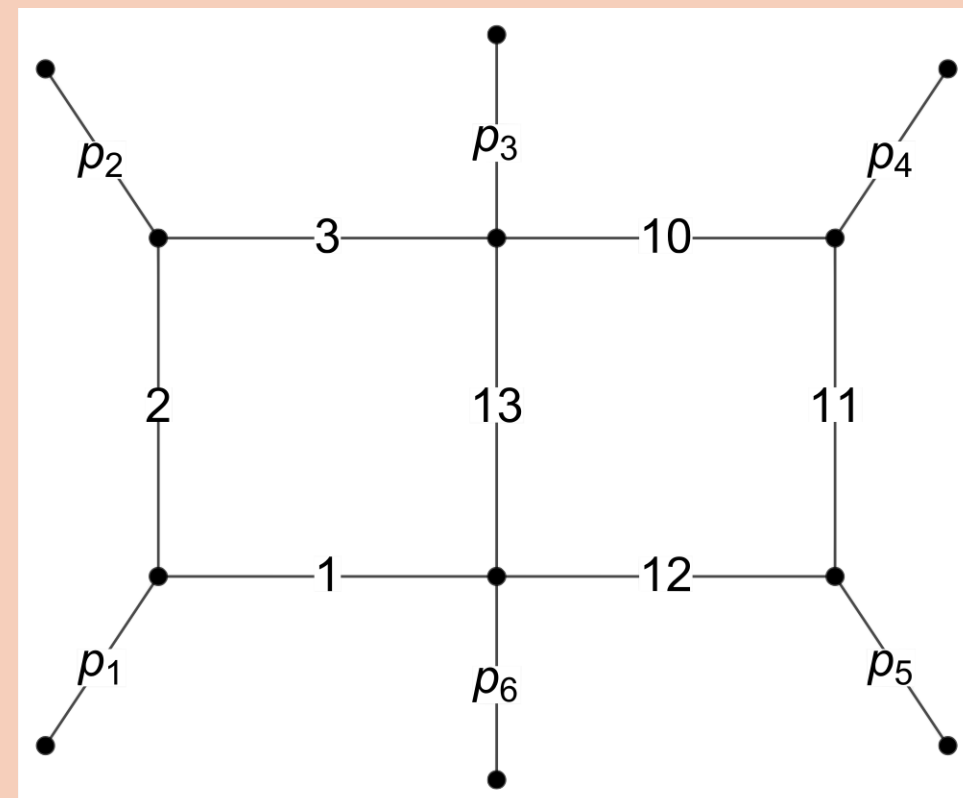
Solution of canonical DE

$$dI = \epsilon(d\tilde{A})I$$

$$I = \frac{1}{\epsilon^4} \left(I^{(0)} + \epsilon I^{(1)} + \epsilon^2 I^{(2)} + \epsilon^3 I^{(3)} + \epsilon^4 I^{(4)} + \dots \right) \quad I^{(n)} = I_{X_0}^{(n)} + \int_{\gamma} (d\tilde{A}) I^{(n-1)}$$

weight-1, weight-2

All in logarithm and classical poly-logarithm



$$I_{\text{db},1}^{(2)} =$$

$$\begin{aligned} & -\log(-v_1)\log(-v_2) - \log(-v_1)\log(-v_3) + \log(-v_1)\log(-v_4) - \log(-v_1)\log(-v_5) - \\ & \log(-v_1)\log(-v_6) + 4\log(-v_1)\log(-v_8) + \frac{1}{2}\log^2(-v_1) + \log(-v_2)\log(-v_3) - \\ & \log(-v_2)\log(-v_4) - \text{Li}_2\left(1 - \frac{v_2v_5}{v_7v_8}\right) + \log(-v_2)\log(-v_6) + \log(-v_2)\log(-v_7) - \\ & 2\text{Li}_2\left(1 - \frac{v_2}{v_8}\right) - \log(-v_2)\log(-v_8) - \log^2(-v_2) - \log(-v_3)\log(-v_4) + \log(-v_3)\log(-v_5) - \\ & \text{Li}_2\left(1 - \frac{v_3v_6}{v_8v_9}\right) - 2\text{Li}_2\left(1 - \frac{v_3}{v_8}\right) - \log(-v_3)\log(-v_8) + \log(-v_3)\log(-v_9) - \\ & \log^2(-v_3) - \log(-v_4)\log(-v_5) - \log(-v_4)\log(-v_6) + 4\log(-v_4)\log(-v_8) + \\ & \frac{1}{2}\log^2(-v_4) + \log(-v_5)\log(-v_6) + \log(-v_5)\log(-v_7) - 2\text{Li}_2\left(1 - \frac{v_5}{v_8}\right) - \log(-v_5)\log(-v_8) - \\ & \log^2(-v_5) - 2\text{Li}_2\left(1 - \frac{v_6}{v_8}\right) - \log(-v_6)\log(-v_8) + \log(-v_6)\log(-v_9) - \log^2(-v_6) - \\ & \log(-v_7)\log(-v_8) - \frac{1}{2}\log^2(-v_7) - \log(-v_8)\log(-v_9) + 3\log^2(-v_8) - \frac{1}{2}\log^2(-v_9) + \\ & \frac{\pi^2}{6} \end{aligned}$$

Solution of canonical DE

$$dI = \epsilon(d\tilde{A})I$$

$$I = \frac{1}{\epsilon^4} \left(I^{(0)} + \epsilon I^{(1)} + \epsilon^2 I^{(2)} + \epsilon^3 I^{(3)} + \epsilon^4 I^{(4)} + \dots \right) \quad I^{(n)} = I_{X_0}^{(n)} + \int_{\gamma} (d\tilde{A}) I^{(n-1)}$$

weight-3, weight-4

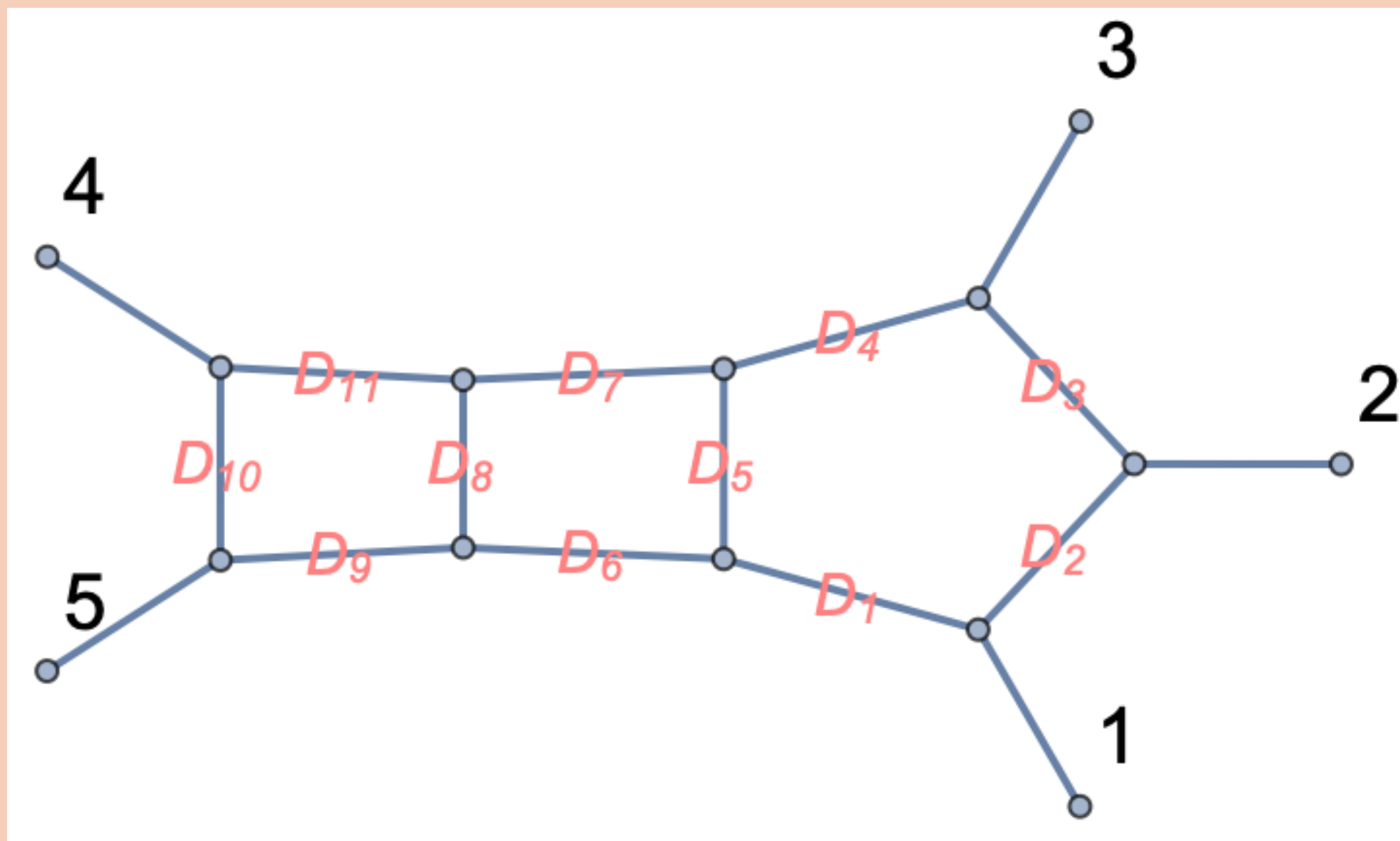
$$\begin{aligned} \vec{I}^{(4)} &= \vec{I}^{(4)}(\vec{x}_0) + \int_0^1 dt \frac{d\tilde{A}}{dt} \vec{I}^{(3)}(\vec{x}_0) + \int_0^1 dt_1 \int_0^{t_1} dt_2 \frac{d\tilde{A}}{dt_1} \frac{d\tilde{A}}{dt_2} \vec{f}^{(2)}(t_2) \\ &= \vec{I}^{(4)}(\vec{x}_0) + \int_0^1 dt \left(\frac{d\tilde{A}}{dt} \vec{I}^{(3)}(\vec{x}_0) + \left(\tilde{A}(1) - \tilde{A}(t) \right) \frac{d\tilde{A}}{dt} \vec{f}^{(2)}(t) \right). \end{aligned} \quad \text{one-fold integration}$$

It takes **minutes on a laptop** to get 20 digits from our analytic solution

3loop 5point Feynman integrals

from the request of John Ellis ...

What we achieved

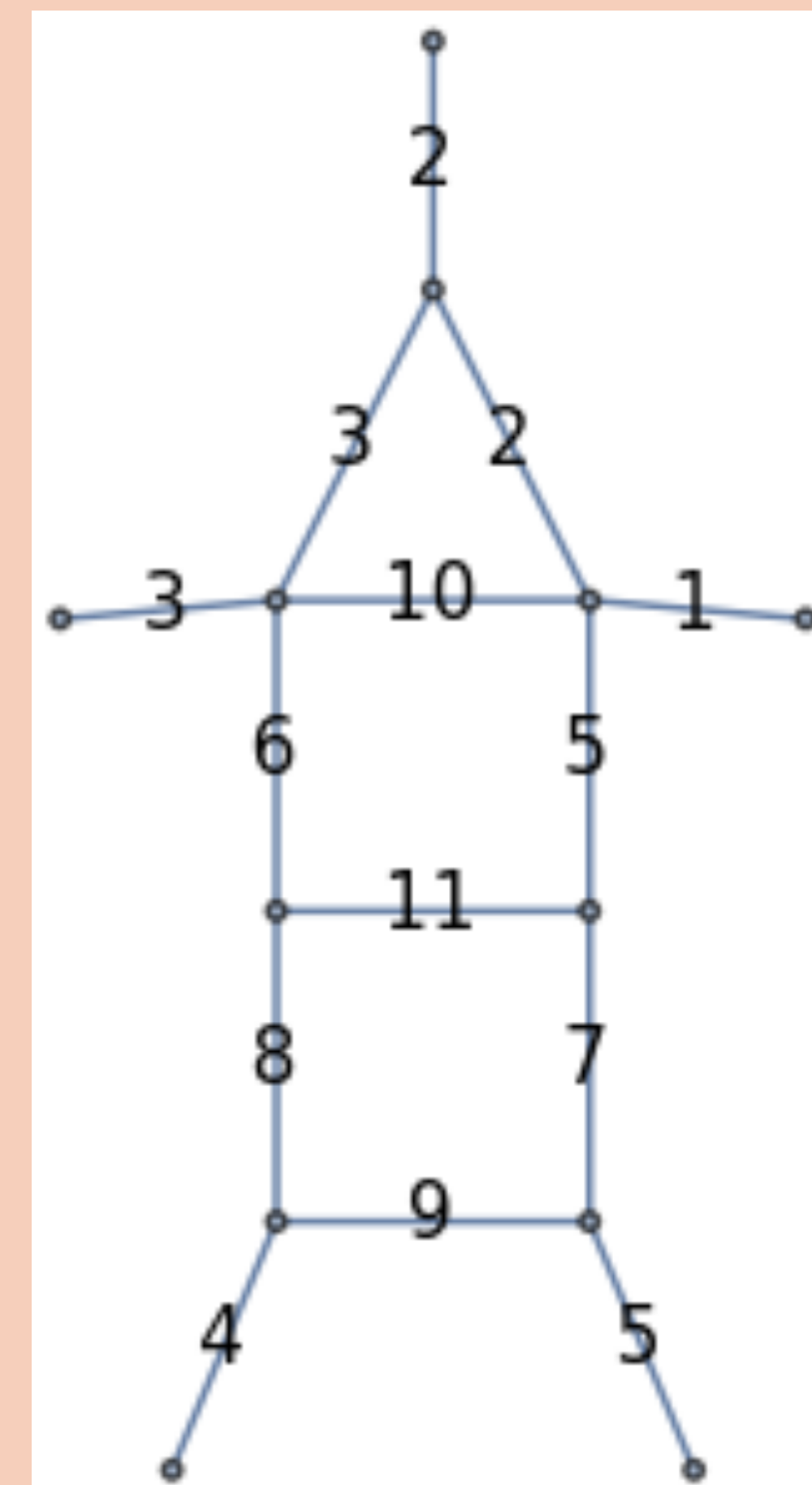


- ✓ UT basis found!
- ✓ Canonical differential derived from NeatIBP

5 scales $s_{12}, s_{23}, s_{34}, s_{45}, s_{15}$

316 Master Integrals

For this “rocket” subfamily, the analytic boundary values are obtained



Liu, Matijasic, Miczajka, Peraro, Xu, Xu, YZ, *to appear*

Summary

With the latest progress on IBP reduction, UT basis determination, alphabet searching, we analytically computed **2loop 6point** massless Feynman integrals.

This is the analytic computation of 2loop **8-scale** Feynman integrals in DR.

and the work on **3loop 5point** will appear soon.

The analytic computation of many more multi-loop multi-leg multi-scale Feynman integrals would be possible.

Thank you