



Heavy-ion collision experiment (Lecture 1)

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- 1. Introduction to relativistic heavy-ion collisions
- 2. Particle yields and statistical model
- 3. Hard probes of QGP: jets
- 4. Hard probes of QGP: heavy flavors
- 5. Quarkonia
- 6. Highlights from small systems
- 7. Summary and outlook





1. Introduction to relativistic heavy-ion collisions



Strong-Interaction Physics



- Strong interaction
 - binds quarks into hadrons
 - binds protons and neutrons into nuclei
- QCD is a very successful theory...
 - e.g. for jet production at high p_{T} and heavy-flavour production

... with some open puzzles

Confinement	Hadron Masses	() 10 ⁵ () 10,	Higgs Vacuum Electroweak.symmetry.breaking	•
Impossible to find an isolated quark or gluon	Proton consists of 2u+1d quark m _{uud} ~10 MeV, m _p = 938 MeV	ggs quark mass 101 101		
Why?	Where is the extra mass generated?	∃, ∃10	d R U U	CD Vacuum



Fundamental Questions



- How do "free" quarks and gluons behave?
- How do quarks and gluons behave when chiral symmetry is restored?
- What generates the constituent masses?
- How does matter behave at very high temperature and large density?
- In the early universe a phase with free quarks and gluons and restored chiral symmetry has existed
 - Quark-gluon plasma (QGP)
 - Recreate in the laboratory with heavy-ion collisions



What is the purpose of heavy-ion collision?



- A way to study QCD
 - ... without confinement
 - ... with quarks at their bare masses
- A way to study matter
 - ... at temperatures 10⁵ times larger than in the sun core
 - \ldots at energy densities like 10 μs after the Big Bang



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Phases of QED matter



Vapor















Ultra-dense QCD matter



Increase the Temperature (T)





Ultra-dense QCD matter



Increase the Temperature (T)



Increase the Density (ρ)







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 $\rho \sim \rho_{\rm C}$





Fundamental properties of **QCD**

Confining property of QCD

Quarks and gluons exist in nature as confined in colorless hadrons: ordinary matter at room temperature

Asymptotic freedom

The strong coupling, i.e. the effective strength of the interaction, becomes weaker for processes involving larger quadrimomentum exchange

Mass of hadrons is a consequence of the strong interaction acting among their constituents

Chiral symmetry is explicitly broken by non-zero quark masses: (light) quarks acquire mass dynamically





Understanding confinement

ALICE

Properties of the QCD vacuum:

- Gluon-gluon self-interaction (non-abelian)
- QCD field lines compressed in flux tube (or "string")

The q-qbar potential is of the form (Cornell potential):



- The potential grows with distance.
- If pulled apart, the energy in the string increases
- A new q-qbar pair is created once the energy is above production threshold.
- No free quark can be obtained by breaking a flux tube \rightarrow confinement



[[]illustration from Fritzsch]

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Phenomenology of Confinement

- QCD vacuum can be seen as liquid of gluon-gluon pairs
- Why does this create confinement?
- MIT bag model : hadrons are confined in bubbles of QCD vacuum
 - Surrounded by QCD vacuum exerting pressure





The MIT Bag Model



Inside the bag. quarks have very small masses and the interaction is weak

Outside the bag. quarks are not allowed to propagate, no colored partons, but quark and gluon condensates

The equilibrium between the kinetic pressure of the quarks inside the hadron vs the pressure of the surrounding QCD vacuum ("bag pressure", *B*) defines the radius *R* of the hadron.

If the hadron can be modeled as N massless Dirac fermions in a spherical cavity,

$$E = \frac{2.04N}{R} + \frac{4\pi}{3}R^3B$$

By asking $\partial E/\partial R = 0$ and $R(p) \sim 0.8$ fm $\rightarrow B_{MIT} \sim (200 \text{ MeV})^4$

If kinetic pressure exceeds bag pressure? \rightarrow deconfinement



"Empty" vacuum



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Deconfinement



For a gas of massless, relativistic partons, the pressure can be calculated from the Stefan-Boltzmann law

$$P = \left(n_g + \frac{7}{8}n_f\right)\frac{\pi^2 T^4}{90}$$

where the degrees of freedom of the system are $n_g = 8$ gluons x 2 spin = 16 $n_f = 2$ quark flavors x 2 spin x 3 colors + anti-q = 24

The systems gets **deconfined** if the kinetic pressure exceeds the bag pressure

P > B → T⁴ > (200 MeV)⁴ * 90 / (16+7/3) / π^2 → T_c > 141 MeV (critical temperature)

Above T_c , the system undergoes a phase transition to a state of matter where quark and gluons are (quasi) free, the Quark-Gluon Plasma (QGP)

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A gas of relativistic partons







QCD phase diagram



the early universe The phases of QCD matter can be Phase transition Heavy ion collisions at particle accelerators summarized in a phase diagram as a 200 function of two parameters: 175 Quark-gluon plasma temperature T and baryochemical deconfined 150 potential μ_B [emperature [MeV] 125 chiral symmetry $\mu_B = \frac{\partial E}{\partial n_B}, \quad n_B = n(B) - n(\overline{B})$ 100 confined 75 $\mu_{\rm B} = 0 \rightarrow \text{antimatter} / \text{matter} = 1$ Hadron phase color superconductivity as at the LHC and in the Early Universe! (several phases) 50 25 chiral symmetry broken **Neutron stars** The quark-gluon plasma is the deconfined phase of strongly-interacting matter. 250 500 750 1000 1500 2000 1250 1750 Baryon chemical potential [MeV] Increase the Density (ρ) (CCN⊕ baryons - antibaryons 2024复旦大学暑期学校/周代翠 2024年8月9日-12日 14



Lattice QCD



- Lattice QCD (IQCD) is a well-established non-perturbative method to solve QCD starting ICE from first principles, i.e. the QCD Lagrangian.
- discretization on a space-time lattice
 - \rightarrow ultraviolet (i.e. large-momentum scale) divergencies can be avoided

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- around critical temperature (T_c): rapid change of
 - energy density ε
 - entropy density s
 - pressure density p
- due to activation of partonic degrees of freedom
- at zero baryon density \rightarrow smooth crossover
- \circ T_C = (156.5 ± 1.5) MeV [A Bazavov et al. Phys.Lett.B 795 (2019) 15]

U, (~')

 $\circ \quad \varepsilon \sim O(GeV/fm^3)$

Lattice QCD predicts a **continuous cross-over** between hadron gas and quark gluon plasma



相对论重离子碰撞的演化图像



 No direct observation of the QGP is possible → rely on emerging particles as "probes" 2024年8月9日-12日
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How to probe the QGP: hard probes





 $1 \text{ fm/c} = 3 \times 10^{-24} \text{ s}, 1 \text{ MeV} \sim 10^{10} \text{ K}$

Charm and beauty quarks (\rightarrow open HF, quarkonia), high-p_T partons (\rightarrow jets)

produced in the early stages in hard processes,

traverse the QGP interacting with its constituents

- \rightarrow rare, calibrated probes, perturbative QCD
- \rightarrow in-medium interaction (energy loss) and transport properties
- \rightarrow in-medium modification of the strong force and of fragmentation









 $1 \text{ fm/c} = 3 \times 10^{-24} \text{ s}, 1 \text{ MeV} \sim 10^{10} \text{ K}$

Low-p_T **particles, light flavour hadrons (u,d,s, + exotic states)** produced from hadronization of the strongly-interacting, thermalized QGP constitutes

- \rightarrow non-perturbative QCD regime
- \rightarrow thermodynamical and transport properties



Evolution of the heavy-ion collisions











With symmetric proton beams with energy E, the centre-of-mass energy is $\sqrt{s} = 2E$.

With heavy-nuclei, only protons can be accelerated, but neutrons are there too:

 $p_{A} = Z/Ap_{proton}$

<u>At the LHC</u>, the rigidity of accelerated particles is fixed by the magnet field configuration ($B_{max} = 8.3 \text{ T}$).

For the ²⁰⁸Pb⁸²⁺ ions used at the LHC: $p_{Pb} = 82 / 208 * p_{proton}$

$$p_{\text{proton}} = 6.5 \text{ TeV} (\text{Run 2}) \rightarrow p_{\text{Pb}} = 2.56 \text{ TeV}$$

 $\rightarrow \sqrt{s_{\text{NN}}} = 5.02 \text{ TeV} \rightarrow \sqrt{s} \sim 1.04 \text{ PeV}$



RHIC @ BNL (2000-) $\sqrt{s_{NN}} < 200 \text{ GeV}$ [beam energy scan $\sqrt{s_{NN}} = 3$, 7.7, 11.5, 19.6, 27, 39, and 62.4 GeV] LHC @ CERN (Run I, 2009-2013) $\sqrt{s_{NN}} = 2.76 \text{ TeV}$ LHC @ CERN (Run II, 2015-2018) $\sqrt{s_{NN}} = 5.02 \text{ TeV}$ HL-LHC @ CERN (Run III+IV, 2022-2030) $\sqrt{s_{NN}} = 5.5 \text{ TeV}$ NICA @ JINR (2021) 3 < $\sqrt{s_{NN}} < 11 \text{ GeV}$



世界上正在运行的高能重离子加速器





Super Proton Syncrotron and Large Hadron Collider **CERN (Switzerland/France)**



Brookhaven RHIC

- Operating since 2000
- Circumference 3.83 km, 2 rings
- Superconducting magnets
- $\sqrt{s_{NN}} = 3 200 \text{ GeV}$ in Au-Au
- Beam energy scan I: 2010-11
- Beam energy scan II: 2019-22
- Ongoing exp: STAR, PHENIX

F. Bellini | SSL 2022 | Heavy lons 2024年8月9日-12日



Operating since 1986 Circumference 6.9 Km max p = 450 A/Z GeV

Ongoing: NA61/Shine

 $\sqrt{s_{NN}} < 20 \, GeV$

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CERN LHC

- **Operating since 2009**
- Run III: started in 2022
- Circumference: 27 km
- B-field: 8 T, superconducting
- pp $\sqrt{s} = 0.9 13.6$ TeV
- Pb-Pb $\sqrt{s_{NN}} = 2.76-5.5 \text{ TeV}$
- Main ongoing: ALICE, ATLAS, CMS, LHCb •



世界上正在建造中的低能重离子加速器





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Heavy-ion physics at the LHC







Dedicated heavy-ion detector: ALICE

MUON-Arm

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V0.





Length: 26 meters Height: 16 meters Weight: 10,000 tons

Solenoid: magnetic field B = 0.5 T

Central Barrel (|η|<0.9)

- ITS, TPC: vertexing + 2π tracking and PID down to very low pT~0.1 GeV/c
- **EMCal/Dcal, PHOS:** high-*p*_T electron trigger, PID, photon, high pT pi0 and electrons

EMCAL

TPC

ITS

TDR

TOF

DCA

PHOS

F

TOF, TRD, HMPID, etc.: Particle identification detectors Dedicated to measure hadrons, electrons, muons and photons to cope with very high multiplicities

Muon-Arm (-4<η<-2.5): Muon trigger, tracking, PID; Heavy-floavor hadrons and W/Z⁰

Forward rapidity detector (**V0, T0, ZDC, SPD)**: Trigger, centrality selection, event plane rec.







LHC上运行的粒子束流



Comprehensive study of Pb-Pb, p-Pb, pp as well as the collision of lighter ions: Xe (done), O (planned for 2024)

pp



p-Pb







Pb-Pb



 $[\]sqrt[(*)]{s_{\rm NN}} = 2.76, 5.02 \, TeV$

 $(\ensuremath{^*})$ collisions energy in Run 1 and 2

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ALICE shopping list



- Soft physics
- Event multiplicity and particle production
- Correlations and fluctuations

Hard probes

- Heavy quarks, jets and high*p*_T photons
- Ultra-peripheral collisions

New physics: magnetic field effects, exotic particles, light nuclei, antimatter…

- Heavy-flavors and quarkouniums
- Collectivity
- Exotic particle properties and CME
- Cold nuclear effects

ICE



Geometry of heavy-ion collisions



We can control the geometry of the collision only by selecting in centrality.

Centrality = fraction of the total hadronic cross section of a nucleus-nucleus collision, typically expressed in percentile, and related to the impact parameter (b)



Other variables related to centrality:

- N_{coll} , number of binary nucleon-nucleon collisions
- N_{part}, number of participating nucleons

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Geometry of heavy-ion collisions

More **central**, ie. "head-on" collisions

- \rightarrow smaller impact parameter
- \rightarrow larger overlap region
- \rightarrow more participants
- \rightarrow more particles produced

More **peripheral** collision

- \rightarrow larger impact parameter
- \rightarrow smaller overlap region
- \rightarrow less participants
- \rightarrow fewer particles produced

Centrality is determined by **counting the number of particles (multiplicity)** or **measuring the energy deposition** in a region of phase space *independent* from the measurement, to avoid biases/autocorrelations in the results.



ALICE, PRL 106 (2011) 032301, PRC 91 (2015) 064905



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Bjorken's Estimate of the Energy Density





πR²



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2cτ





How many particles are created in a collision?



In a central Pb-Pb collision at the LHC, more than 20000 charged tracks must be reconstructed.

→ High granularity tracking systems, primary importance of tracking, vertexing calibration



Charged particle density in central AA collisions





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Particle production per participant in AA collisions follows a steeper power law than in pp, pA and increased by **2-3x** from RHIC to the LHC

AA collisions are more efficient in transferring energy from beam- to mid- rapidity than pp



- ALICE: Pb–Pb at 5.02 TeV highest energy so far
 - → For 0–5% most central collisions, confirms trend from lower energies

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Temperature of the QGP



• Low- p_T : 2.6 σ excess w. r. t. models in 0–20% central — thermal contribution

- $T_{\text{eff}} = 297 \pm 12(\text{stat.}) \pm 41 \text{ (syst.) MeV}$ in central collisions above $T_c \sim 170 \text{ MeV}$
- 30% higher than at RHIC (Au–Au at √SNN=200 GeV) ALICE Phys. Lett. B754 (2016) 235 2024年8月9日-12日 2024复旦大学暑期学校/周代翠 (CCNU) 34





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Dedicated heavy-ion detector: ALICE



Countries: 40 Institutes: 176 Members: 2002

Detector:

Length: 26 meters Height: 16 meters Weight: 10,000 tons

Solenoid: magnetic field B = 0.5 T



Central Barrel (|n|<0.9)

- ITS, TPC: vertexing + 2π tracking and PID down to very low pT~0.1 GeV/c
- EMCal/Dcal, PHOS: high- $p_T \pi 0$, and electron trigger, PID

Chinese participation: CCNU, CIAE, CGU, Fudan, HUST, HTU, USTC

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TOF, TRD, HMPID, etc.: Particle identification detectors

Hardware contribution from China:

- PHOS: FEE;
- DCAL/EMCAL: one supermodule
- ITS2: 520 modules, MAPS design
- MFT: five readout discs

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Dedicated to measure

hadrons, electrons, muons and photons to cope with very high multiplicities

Collisions systems (so far) : Pb-Pb, pp, p-Pb, Pb-p, Xe-Xe

Muon-Arm (-4<η<-2.5): Muon trigger, tracking, PID

Forward rapidity detectors (**V0, T0, ZDC, SPD)**: Trigger, centrality selection, event plane rec.

Physics on:

- Photons and jets
- Heavy-flavors and quarkouniums
- Collectivity
- Exotic particle properties
- Cold nuclear effects