



Exploring the heavy meson distribution amplitudes from first principles

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In collaboration with:

Xue-Ying Han, Jun Hua, Xiangdong Ji, Cai-Dian Lü, Wei Wang, Ji Xu, Shuai Zhao.

Qi-An Zhang, Beihang University

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Outline

- **Motivation**
- **Heavy meson LCDAs:**
 - **Challenges and recent progresses**
 - **Two-step factorization for heavy meson LCDAs**
 - **LCDAs in QCD and HQET**
- **Summary and outlook**

Motivation

➤ Weak decays of B meson are critical for:

- **Precise tests of SM**

- **Searching for NP**

- **Understanding the origins of CPV**

-

- $B \rightarrow \pi\pi$: *Beneke, Buchalla, Neubert, Sachrajda, 1999; 1422 citations*

- Lu, Ukai, Yang, 2001; 568 citations*

- $B \rightarrow \pi K$: *Beneke, Buchalla, Neubert, Sachrajda, 2001; 1177 citations*

- $B \rightarrow \pi\ell\nu$: *Becher, Hill, 2005; 215 citations*

- Khodjamirian, Mannel, Offen, Wang, 2011; 192 citations*

- $B \rightarrow K^{(*)}\ell\ell$: *Khodjamirian, Mannel, Pivavov, Wang, 2010; 486 citations*

- $B \rightarrow D\ell\nu$: *HPQCD Collaboration, 2015; 387 citations*

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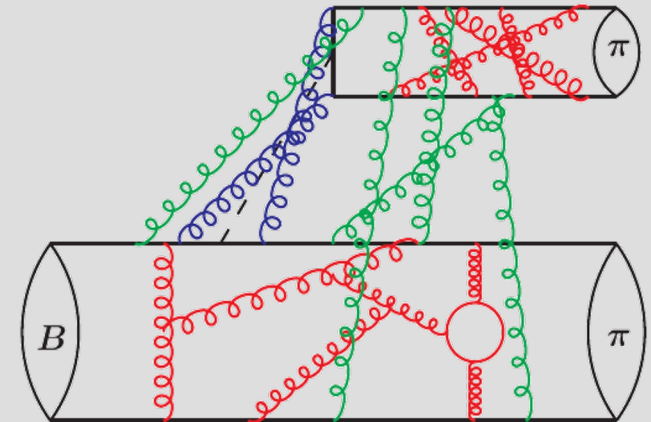
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➤ Factorization: categories by different characteristic scales

$$\langle \pi(p') \pi(q) | Q_i | \bar{B}(p) \rangle = f^{B \rightarrow \pi}(q^2) \int_0^1 dx T_i^I(x) \phi_\pi(x) + \int_0^1 d\xi dx dy T_i^{II}(\xi, x, y) \phi_B(\xi) \phi_\pi(x) \phi_\pi(y)$$

Form factor =
Hard kernel + LCDAs

Hard kernel (Perturbative) **Meson LCDAs (Nonperturbative)**



Light Meson LCDAs: Research Progresses and Challenges

➤ Light meson LCDAs have been extensively pursued: (1970s - now)

- **Asymptotic LCDAs**

Chernyak, Zhitnitsky, 1977; Lepage, Brodsky, 1979; Efremov, Radyushkin, 1980

- **Dyson-Schwinger Equation**

Chang, Cloet, Cobos-Martinez, Roberts, Schmidt, 2013; Gao, Chang, Liu, Roberts, Schmidt, 2014; Roberts, Richards, Chang, 2021

- **Sum rules**

Chernyak, Zhitnitsky, 1982; Braun, Filyanov, 1989; Ball, Braun, Koike, Tanaka, 1998; Ball, Braun, 1998; Khodjamirian, Mannel, Melcher, 2004; Ball, Lenz, 2007

- **Inverse Problem**

Li, 2022

- **Models**

Arriola, Broniowski, 2002, 2006; Zhong, Zhu, Fu, Wu, Huang, 2021;

- **Global Fits**

Stefanis, 2020; Cheng, Khodjamirian, Rusov, 2020; Hua, Li, Lu, Wang, Xing, 2021

- **Lattice with current-current correlation**

Bali, Braun, Gläfle, Göckeler, Gruber, 2017, 2018;

- **Lattice with OPE**

Martinelli, Sachrajda, 1987; Braun, Bruns, et al., 2016; RQCD collaboration, 2019, 2020

- **Lattice with LaMET**

Zhang, Chen, Ji, Jin, Lin, 2017; LP3 Collaboration, 2019; Zhang, Honkala, Lin, Chen, 2020; Lin, Chen, Fan, Zhang², 2021; LPC Collaboration, 2021, 2022

- **Quantum Computing**

QuNu Collaboration, 2023, 2024

Heavy Meson LCDAs: Research Progresses and Challenges

- The HQET matrix element of heavy meson [Grozin, Neubert, 1997; Beneke, Feldmann, 2000]

$$\langle 0 | \bar{q}_\beta(\xi)[\xi, 0] h_{v\alpha}(0) | \bar{B}(v) \rangle = -\frac{i\tilde{f}_B m_B}{8} \left\{ \left[\underbrace{\tilde{\phi}_B^+(t, \mu) v_+ \gamma_-}_{\text{Leading twist}} + \underbrace{\tilde{\phi}_B^-(t, \mu) v_- \gamma_+}_{\text{Sub-leading twist}} \right] \gamma_5 \right\}_{\alpha\beta}$$

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Leading twist Sub-leading twist

- **Evolution of φ_B^+ and φ_B^- :** *[Lange, Neubert, 2003; Bell, Feldmann, 2008]*

$$\frac{d}{d \ln \mu} \varphi_B^+(\omega, \mu) = -\frac{\alpha_s C_F}{4\pi} \int_0^\infty d\omega' \gamma_+^{(1)}(\omega, \omega', \mu) \varphi_B^+(\omega', \mu) + \mathcal{O}(\alpha_s^2)$$

$$\frac{d}{d \ln \mu} \varphi_B^-(\omega, \mu) = -\frac{\alpha_s C_F}{4\pi} \int_0^\infty d\omega' \gamma_-^{(1)}(\omega, \omega', \mu) \varphi_B^-(\omega', \mu) + \mathcal{O}(\alpha_s^2)$$

with

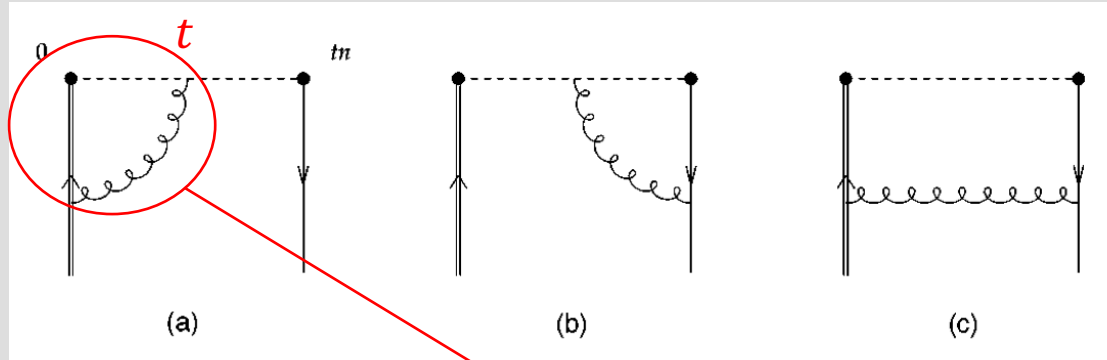
$$\gamma_+^{(1)}(\omega, \omega', \mu) = \left(\Gamma_{\text{cusp}}^{(1)} \ln \frac{\mu}{\omega} - 2 \right) \delta(\omega - \omega') - \Gamma_{\text{cusp}}^{(1)} \omega \left[\frac{\theta(\omega' - \omega)}{\omega'(\omega' - \omega)} + \frac{\theta(\omega - \omega')}{\omega(\omega - \omega')} \right]_+$$

$$\gamma_-^{(1)}(\omega, \omega'; \mu) = \gamma_+^{(1)}(\omega, \omega'; \mu) - \Gamma_{\text{cusp}}^{(1)} \frac{\theta(\omega' - \omega)}{\omega'}$$

- Solution of evolution equations: *[Bell, Feldmann, Wang and Yip, 2013; Braun, Manashov, 2014]*
- RG equations of $\varphi_B^+(\omega, \mu)$ at two-loops: *[Braun, Ji, Manashov, 2019; Liu, Neubert, 2020]*
- RG equations of the higher-twist B-meson distribution amplitudes: *[Braun, Ji, Manashov, 2017]*
- NNLO QCD correction to relevant hadronic B-meson decays: *[Bell, Beneke, Huber, Li, 2020]*

Heavy Meson LCDAs: Research Progresses and Challenges

➤ Difficulties in calculating the moments of heavy meson LCDA:



Diverge at $t \rightarrow 0$!

$$O_+^{\text{ren}}(t, \mu) = O_+^{\text{bare}}(t) + \frac{\alpha_s C_F}{4\pi} \left\{ \left(\frac{4}{\hat{\epsilon}^2} + \frac{4}{\hat{\epsilon}} \ln(it\mu) \right) O_+^{\text{bare}}(t) - \frac{4}{\hat{\epsilon}} \int_0^1 du \frac{u}{1-u} [O_+^{\text{bare}}(ut) - O_+^{\text{bare}}(t)] \right\}$$

[Braun, Ivanov, Korchemsky, 2004]

- Diverge at $t \rightarrow 0 \Leftrightarrow$ **No local limit**
- Non-negative moments $\int dk k^n \varphi_+(k)$ for $n=0, 1, 2, \dots$ are **not related to OPE**, and actually they **diverge**
- Cannot obtain φ_B from lattice QCD through their moments.

Heavy Meson LCDAs: Research Progresses and Challenges

- A theoretical attempt to extract the heavy meson LCDAs in the framework of LaMET:

$$\tilde{\varphi}_B^+(\xi, \mu) \propto \int_{-\infty}^{+\infty} \frac{d\tau}{2\pi} e^{in_z \cdot v \xi \tau} \langle 0 | (\bar{q} W_c) (\tau n_z) \eta_z \gamma_5 (W_c^\dagger h_v) (0) | \bar{B}(v) \rangle \Rightarrow \varphi_B^+(\omega, \mu)$$

Connecting the between equal-time HQET correlator and light-cone one.

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Connecting the between equal-time HQET correlator and light-cone one.

- Leading twist matching @ 1-loop: [Wang², Xu, Zhao, 2020; Xu, Zhang, 2022]
- Sub-leading twist matching: [Hu, Wang, Xu, Zhao, 2024]
- Inverse moment and log moments: [Xu, Zhang, Zhao, 2022; Hu, Xu, Zhao, 2024]
- Ioffe-time distributions: [Zhao, Radyushkin, 2021]

🤔 Difficult to realize the boosted HQET field (time-like Wilson link) on lattice QCD.

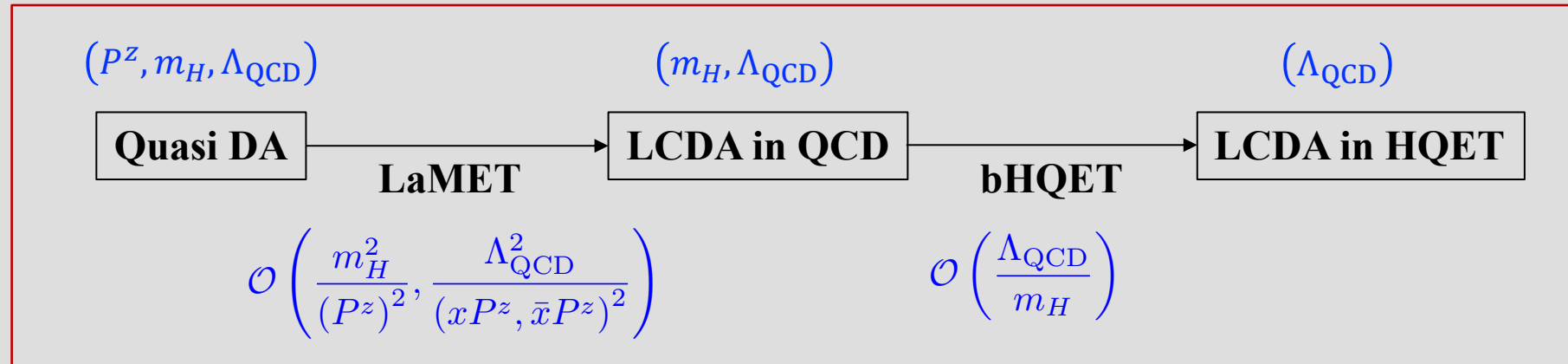
Two-step factorization to access heavy meson LCDA

- Focus on the Quasi DA with QCD fields, which is easy to implement in LQCD:
 - Light meson (2 characteristic scales: $P^Z, \Lambda_{\text{QCD}}$) → Heavy meson (3 scales: $P^Z, m_H, \Lambda_{\text{QCD}}$)

Two-step factorization to access heavy meson LCDA

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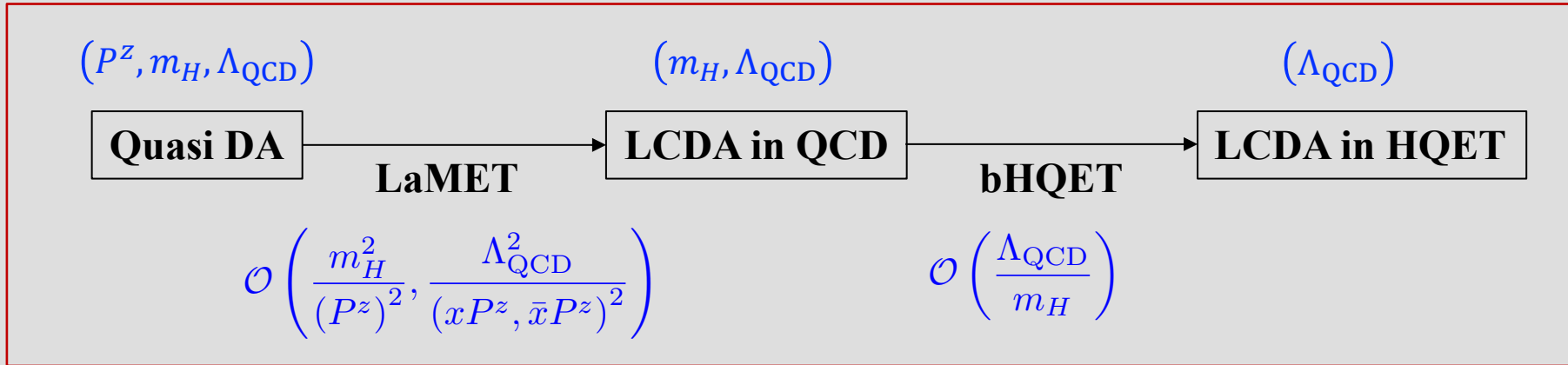


• A multi-scale processes:

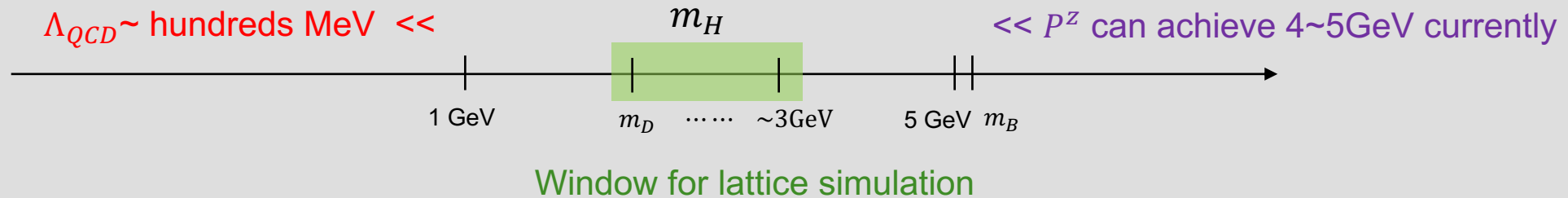
1. LaMET requires $\Lambda_{\text{QCD}}, m_H \ll P^z$ and finally integrate out P^z ;
2. bHQET requires $\Lambda_{\text{QCD}} \ll m_H$ and integrate out m_H ;

⇒ **Hierarchy** $\Lambda_{\text{QCD}} \ll m_H \ll P^z$.

Two-step factorization to access heavy meson LCDA



⇒ **Hierarchy $\Lambda_{\text{QCD}} \ll m_H \ll P^z$** : A big challenge for lattice simulation



At this stage, the heavy meson could be D , but hard to directly implement B !

Matching I: from quasi DAs to LCDAs in QCD

- D meson quasi DA $\tilde{\phi}(x, P^z)$, include the scales $\Lambda_{\text{QCD}} \ll m_D \ll P^z$

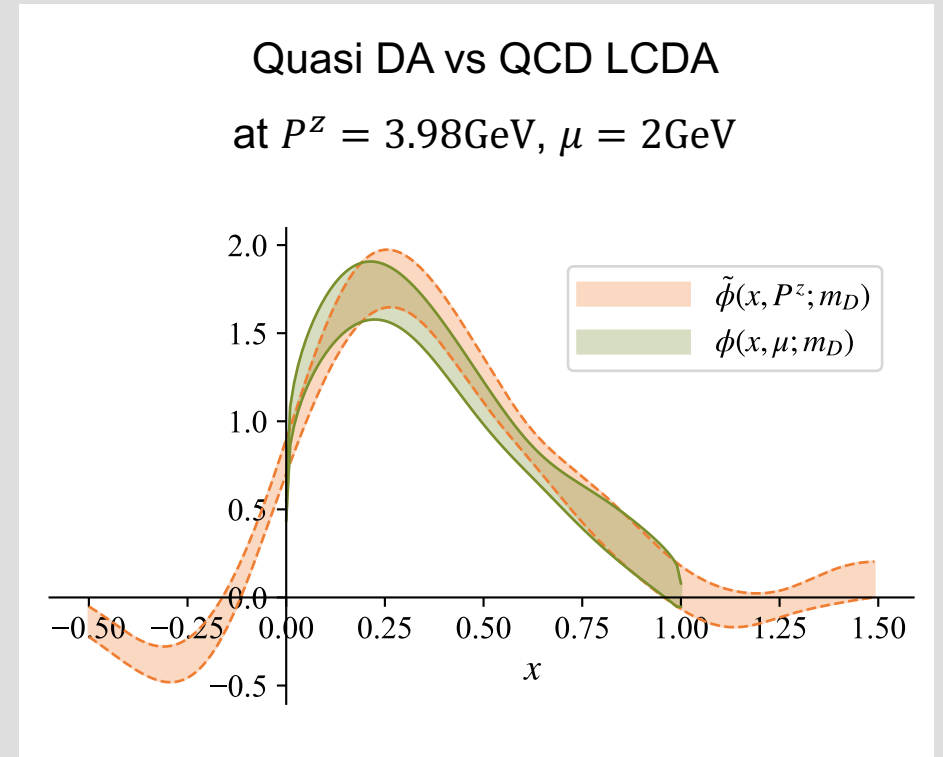
$$\tilde{\phi}(x, P^z) = \int \frac{dz}{2\pi} e^{-ixP^z z} \tilde{M}(z, P^z)$$

- Matching formula in LaMET:

$$\tilde{\phi}(x, P^z) = \int_0^1 C\left(x, y, \frac{\mu}{P^z}\right) \phi(y, \mu) + \mathcal{O}\left(\frac{m_H^2}{(P^z)^2}, \frac{\Lambda_{\text{QCD}}^2}{(xP^z, \bar{x}P^z)^2}\right)$$

Liu, Wang, Xu, QAZ, Zhao, 2019;
Han, Hua, Ji, Lu, Wang, Xu, QAZ, Zhao, 2024

This matching integrate out P^z , obtain the LCDAs in QCD.

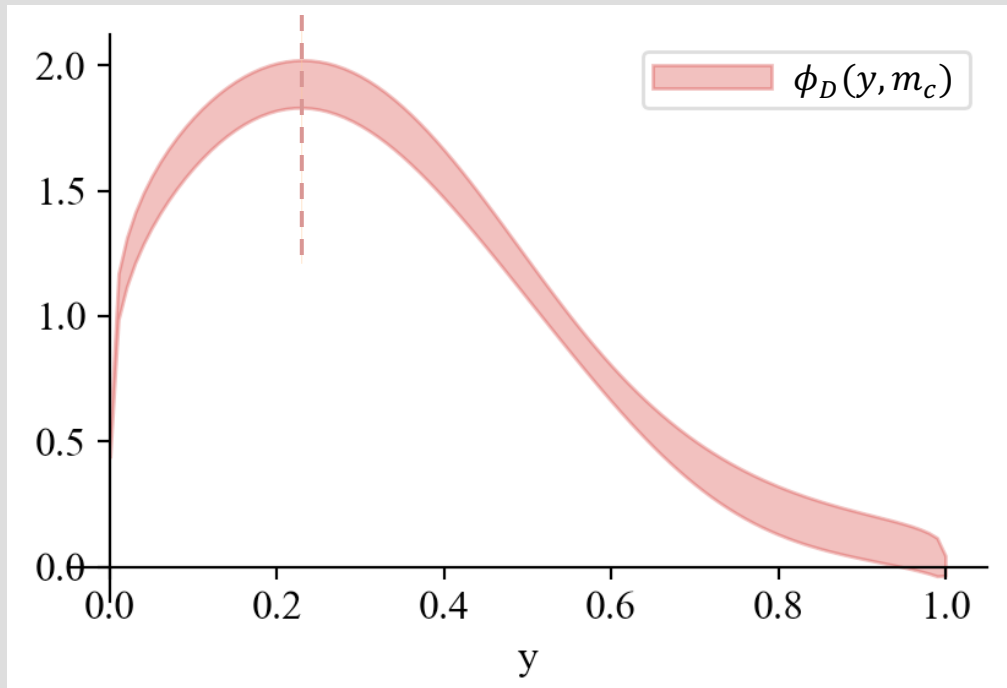


LCDAs in QCD

➤ Heavy meson LCDAs in QCD

$$\begin{aligned} \phi(y, \mu) = & \frac{1}{if_H} \int_{-\infty}^{+\infty} \frac{d\tau}{2\pi} e^{iyP_H\tau n_+} \\ & \times \langle 0 | \bar{q}(\tau n_+) \not{n}_+ \gamma_5 W_c(\tau n_+, 0) Q(0) | H(P_H) \rangle \end{aligned}$$

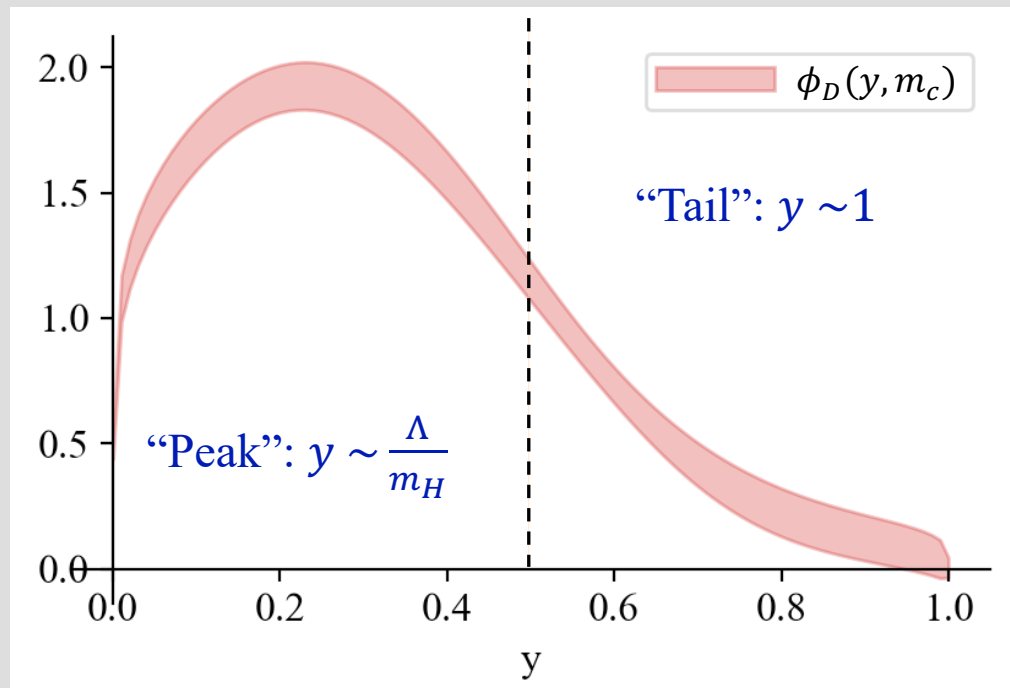
- The peak position dominated by m_H and μ ;
- At very large scale $\mu \gg m_H$, asymptotic form;



LCDAs in QCD

➤ Heavy meson LCDAs in QCD

$$\phi(y, \mu) = \frac{1}{if_H} \int_{-\infty}^{+\infty} \frac{d\tau}{2\pi} e^{iyP_H\tau n_+} \times \langle 0 | \bar{q}(\tau n_+) \not{n}_+ \gamma_5 W_c(\tau n_+, 0) Q(0) | H(P_H) \rangle$$



- The peak position dominated by m_H and μ ;
- At very large scale $\mu \gg m_H$, asymptotic form;
- For the scale $\mu \lesssim m_Q$,
 - ⇒ Light quark carries small momentum fraction $y \sim \Lambda/m_H$ peak region, related to the HQET LCDA;
[Ishaq, Jia, Xiong, Yang, 2020; Beneke, Finauri, Vos, Wei, 2023]
 - ⇒ $y \sim O(1)$ region be suppressed in LCDA:
Only hard-collinear physics (perturbative), and starts at the one-loop level.

Matching II: connecting LCDAs in QCD and HQET

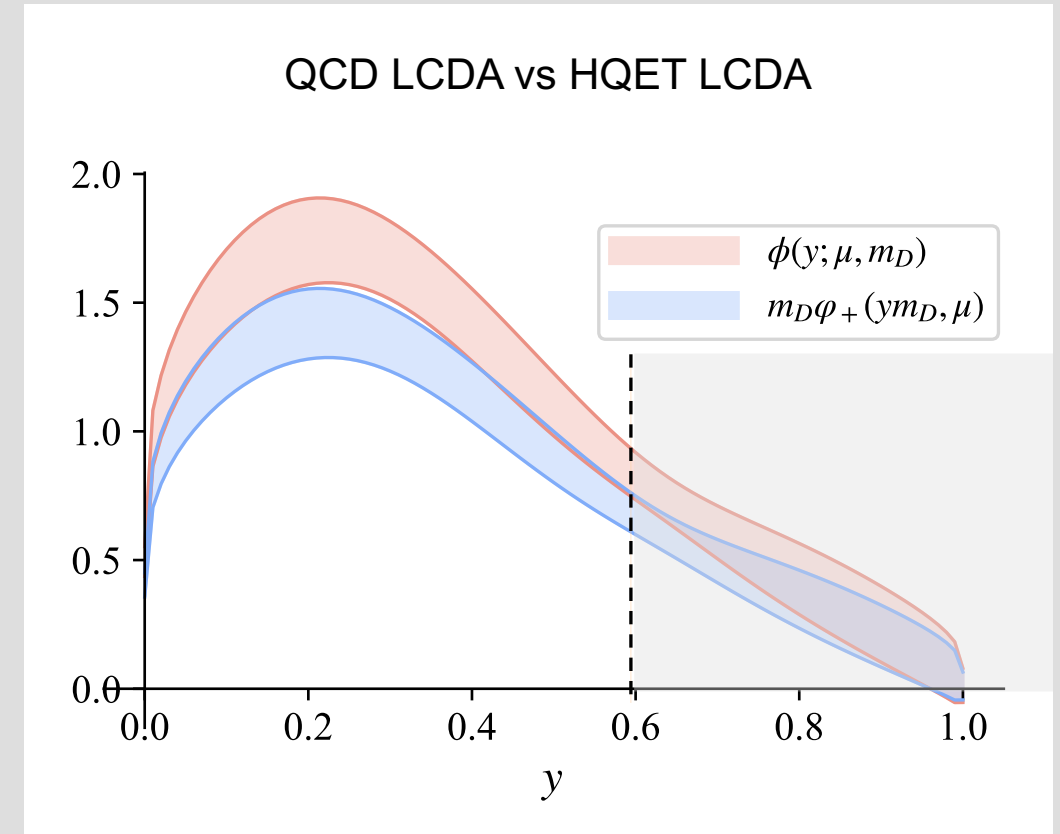
- Leading twist heavy meson LCDA in HQET

$$\varphi^+(\omega, \mu) = \frac{1}{i\tilde{f}_H(\mu)m_H} \int_{-\infty}^{+\infty} \frac{d\eta}{2\pi} e^{i\omega n_+ \cdot v\eta} \times \langle 0 | \bar{q}(\eta n_+) / n_+ \gamma_5 W_c(\eta n_+, 0) h_v(0) | H(v) \rangle$$

connected with the QCD LCDA through a multiplicative factorization in the peak region:

[Beneke, Finauri, Vos, Wei, 2023]

$$\phi(y, \mu; m_H) = \frac{\tilde{f}_H}{f_H} J_{\text{peak}} m_H \varphi^+(\omega, \mu) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{m_H}\right)$$



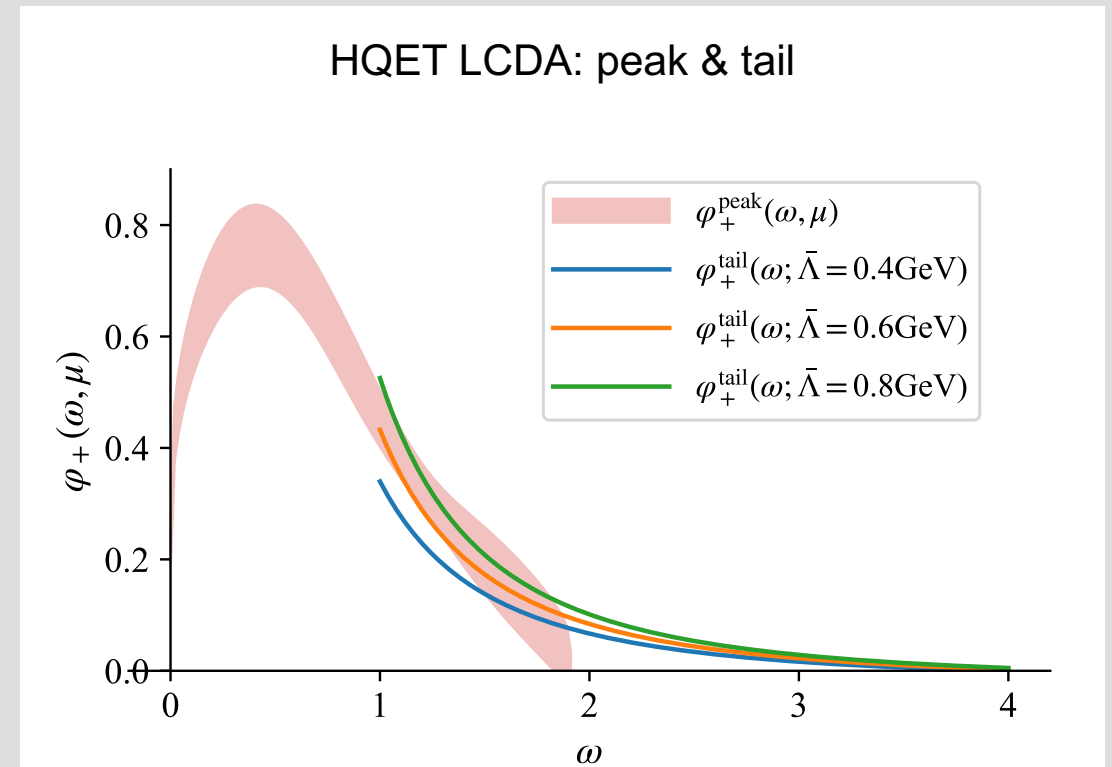
Tails of HQET LCDA

- The tail region of HQET LCDA is perturbative: *[Lee, Neubert, 2005]*

$$\varphi_{\text{tail}}^+(\omega, \mu) = \frac{\alpha_s C_F}{\pi\omega} \left[\left(\frac{1}{2} - \ln \frac{\omega}{\mu} \right) + \frac{4\bar{\Lambda}}{3\omega} \left(2 - \ln \frac{\omega}{\mu} \right) \right]$$

where $\bar{\Lambda} \equiv m_H - m_Q^{\text{pole}}$ reflect the power correction, and usually be chosen as 400~600MeV.

The final results of HQET LCDA will merge the **peak (from LQCD)** and **tail region (from 1-loop calculation)**.



Comparison with phenomenological models

➤ Several commonly used models:

[Wang, Shen, 2015; Beneke, Braun, Ji, Wei, 2018; Gao, Huber, Ji, Wang², 2022]

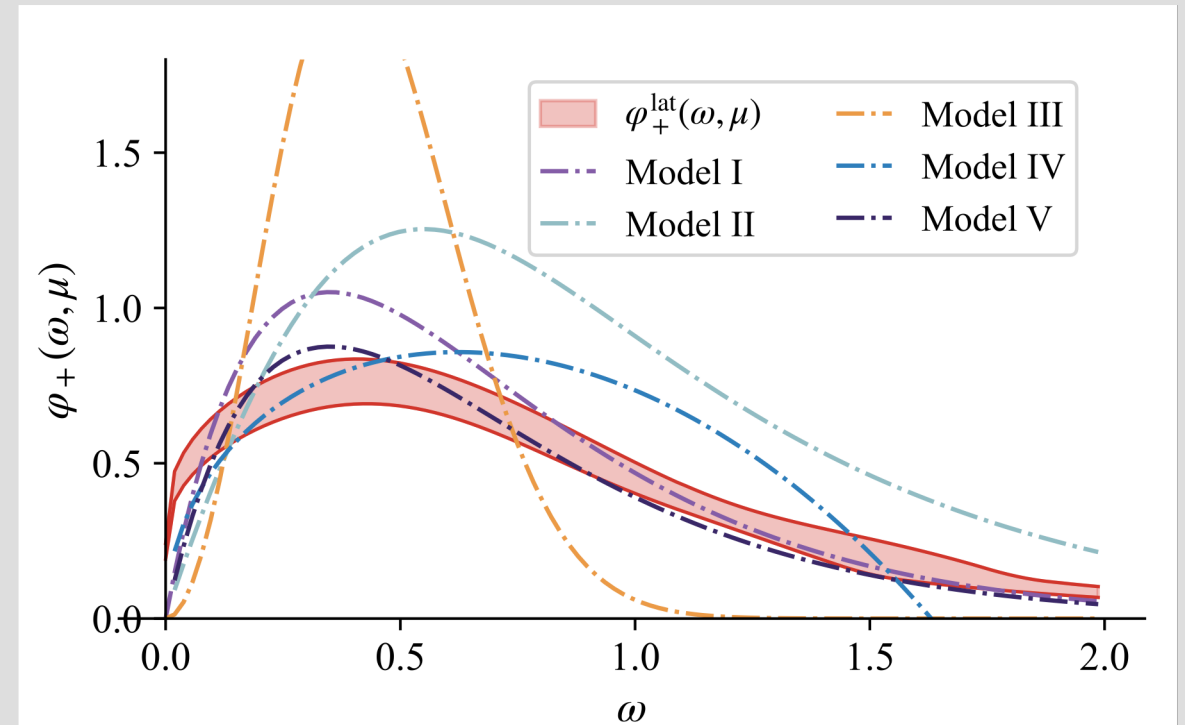
$$\varphi_{\text{I}}^+(\omega, \mu_0) = \frac{\omega}{\omega_0^2} e^{-\omega/\omega_0},$$

$$\varphi_{\text{II}}^+(\omega, \mu_0) = \frac{4}{\pi\omega_0} \frac{k}{k^2 + 1} \left[\frac{1}{k^2 + 1} - \frac{2(\sigma_B^{(1)} - 1)}{\pi^2} \ln k \right],$$

$$\varphi_{\text{III}}^+(\omega, \mu_0) = \frac{2\omega^2}{\omega_0\omega_1^2} e^{-(\omega/\omega_1)^2},$$

$$\varphi_{\text{IV}}^+(\omega, \mu_0) = \frac{\omega}{\omega_0\omega_2} \frac{\omega_2 - \omega}{\sqrt{\omega(2\omega_2 - \omega)}} \theta(\omega_2 - \omega),$$

$$\varphi_{\text{V}}^+(\omega, \mu_0) = \frac{\Gamma(\beta)}{\Gamma(\alpha)} \frac{\omega}{\omega_0^2} e^{-\omega/\omega_0} U(\beta - \alpha, 3 - \alpha, \omega/\omega_0),$$



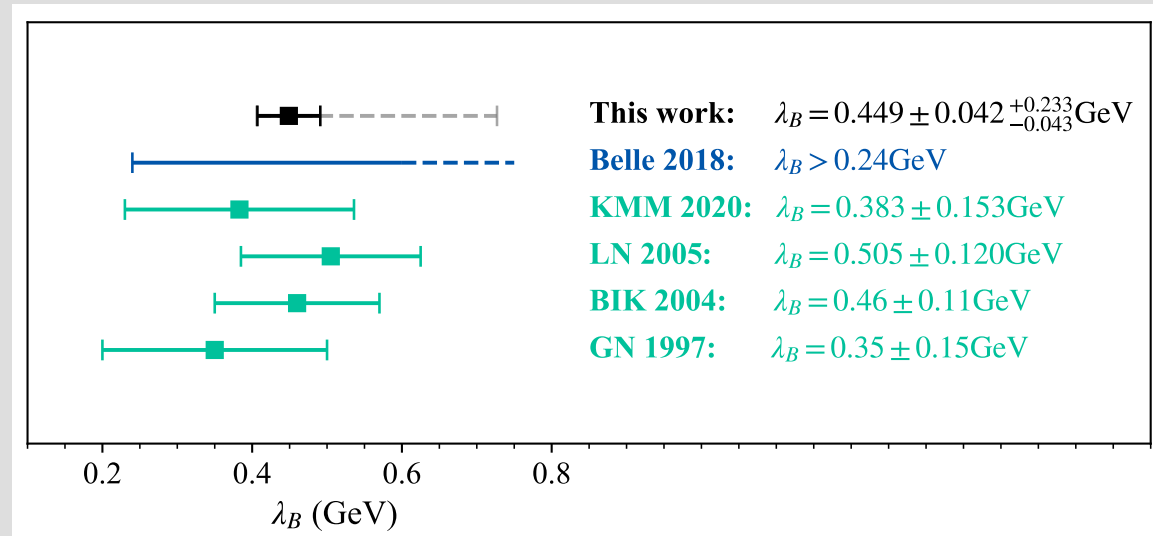
First inverse moment

➤ The first inverse moment

$$\lambda_B^{-1}(\mu) = \int_{-\infty}^{\infty} d\omega \frac{\varphi^+(\omega, \mu)}{\omega}$$

Models	I	II	III	IV	V
Parameters	$\omega_0 = 0.433(23)\text{GeV}$	$\omega_0 = 0.682(45)\text{GeV}$	—	$\omega_0 = 0.427(21)\text{GeV}$	$\omega_0 = 0.449(42)\text{GeV}$
		$\sigma_B^{(1)} = 2.78(48)$			
fit range	$\omega \in [0.2, 1.4]\text{GeV}$	$\omega \in [0.2, 1.4]\text{GeV}$		$\omega \in [0.4, 0.8]\text{GeV}$	$\omega \in [0.2, 1.4]\text{GeV}$
$\chi^2/\text{d.o.f}$	1.4	1.2		2.1	1.0

- The current numerical results are unable to accomplish the integration over full- ω range;
- We determine the λ_B^{-1} by fitting the parameterization forms of different model.



Belle collaboration, 2018
Khodjamirian, Mandal, Mannel, 2020 *Lee, Neubert, 2005*
Braun, Ivanov, Korchemsky, 2004
Grozin, Neubert, 1997

Summary and outlook

- ✓ **Heavy meson LCDAs**: the first **lattice-implementable scheme**, which can continually be **improved**.

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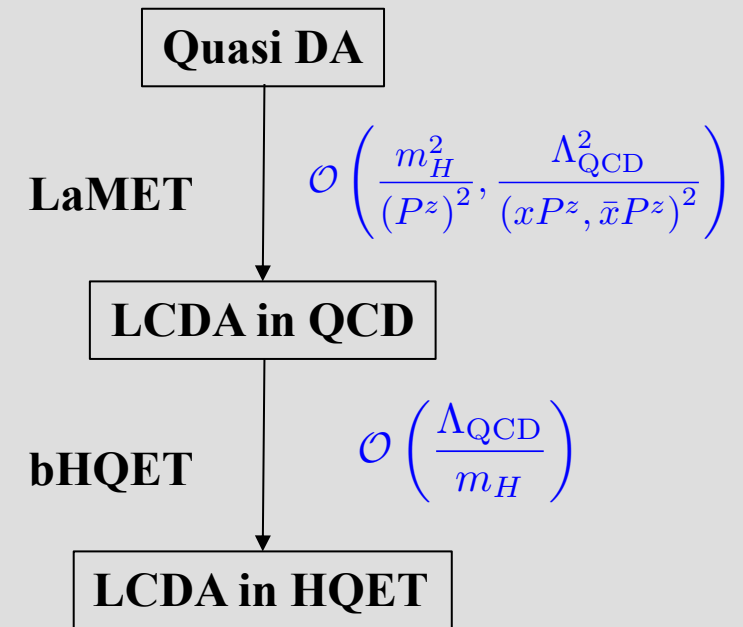
□ The most urgent is to properly control the power corrections within two step factorization:

- Larger P^z to increase the window for lattice calculations;
- $m_D \rightarrow m_B$ is hard, LQCD can simulate some unphysical m_H which satisfy the hierarchy better;

□ More systematic lattice QCD calculations:

- Nonperturbative renormalization, continuum and physical mass extrapolation, operator mixing effects,
- More reliable method to merge the peak and tail regions

□ Realize the HQET quasi DA on lattice QCD directly?



Thanks