

Based on arXiv: 2403.07499

Wen-Fei Wang, Li-Fei Yang, Ai-Jun Ma, Angels Ramos

The low-mass enhancement of kaon pair in the decays

$$B^+ \rightarrow \bar{D}^{(*)0} K^+ \bar{K}^0 \text{ and } B^0 \rightarrow D^{(*)-} K^+ \bar{K}^0$$

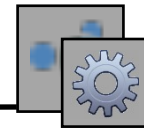
王文飞 (山西大学)

第六届重味物理与量子色动力学研讨会

中国海洋大学 2024-04-21

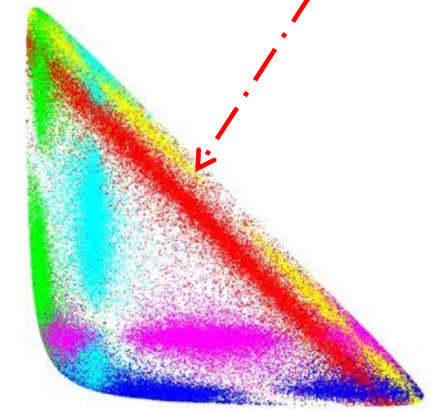
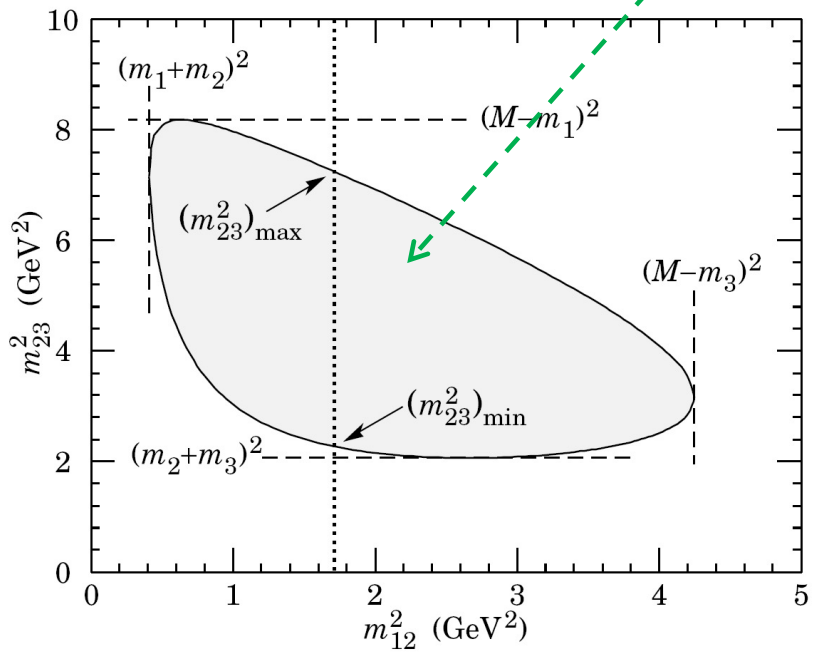
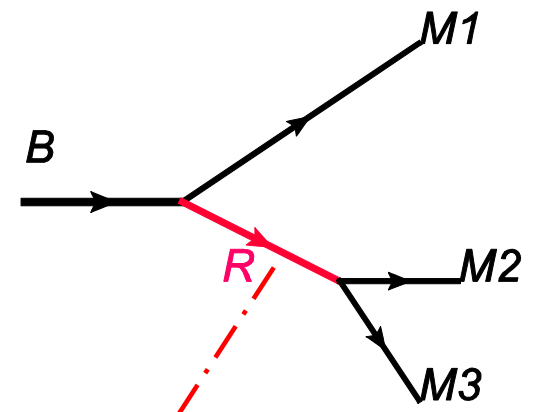
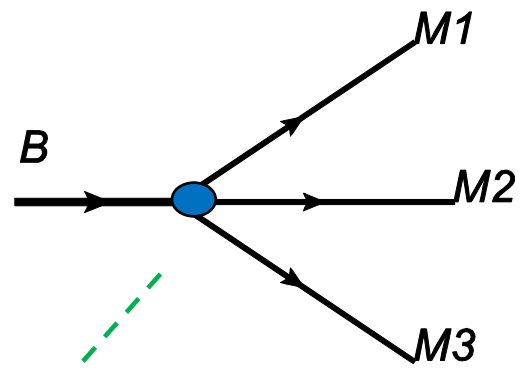


Outline



- **3-body B meson decays and the virtual contribution**
- **Kaon pair in $B^+ \rightarrow \bar{D}^{(*)0} K^+ \bar{K}^0$ and $B^0 \rightarrow D^{(*)-} K^+ \bar{K}^0$ decays**
- **Results and discussions**
- **Conclusions**

$B \rightarrow M1 M2 M3$

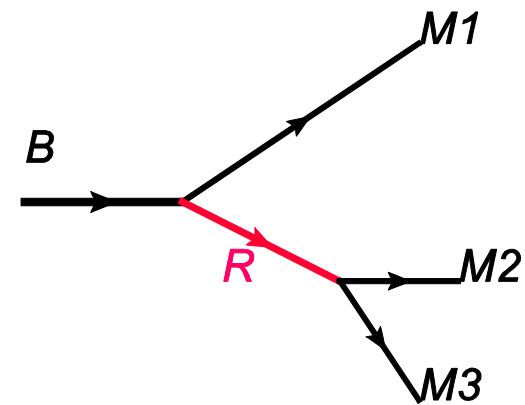
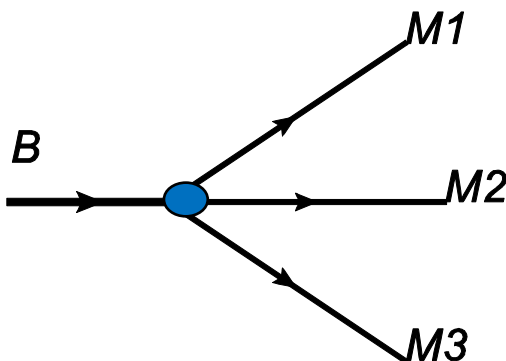


Dalitz Plot



3-body B meson decays and the virtual contribution

$$B \rightarrow M1 M2 M3$$



The weak effective Hamiltonian:

1.

$$\mathcal{H}_{\text{eff}} = \frac{G_F}{\sqrt{2}} \left\{ V_{ub} V_{uq}^* [C_1(\mu) Q_1^u(\mu) + C_2(\mu) Q_2^u(\mu)] - V_{tb} V_{tq}^* \left[\sum_{i=3}^{10} C_i(\mu) Q_i(\mu) \right] \right\} + \text{H.c.}, \quad (2)$$

where $q = d, s$. The functions Q_i ($i = 1, \dots, 10$) are the local four-quark operators:

The total amplitude within isobar approach:

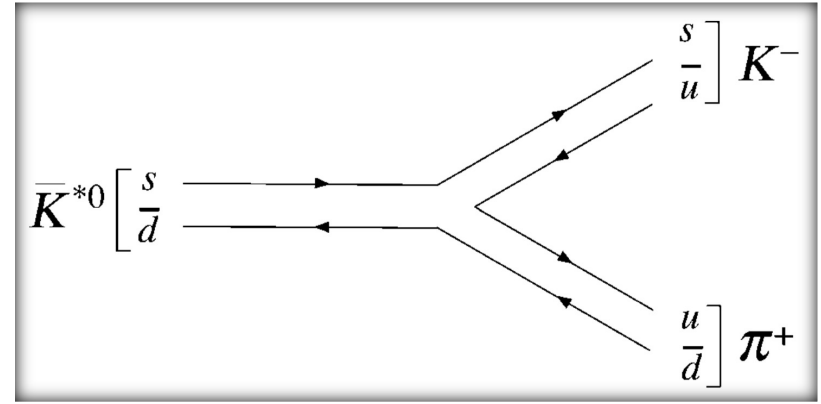
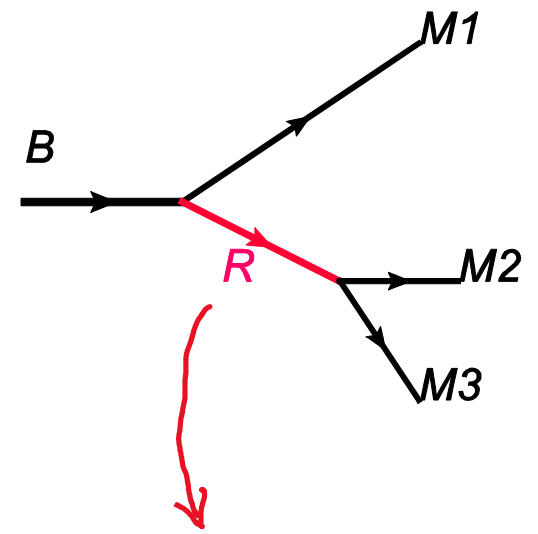
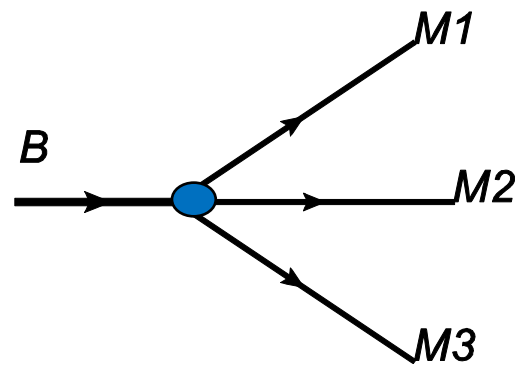
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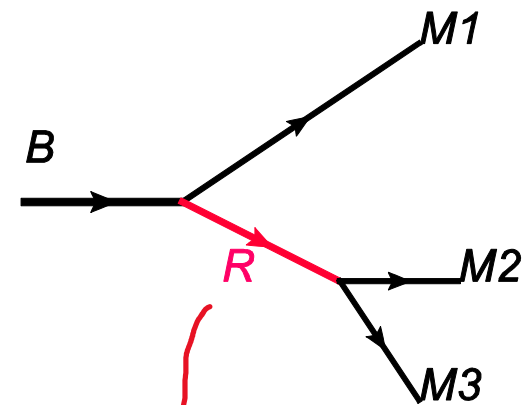
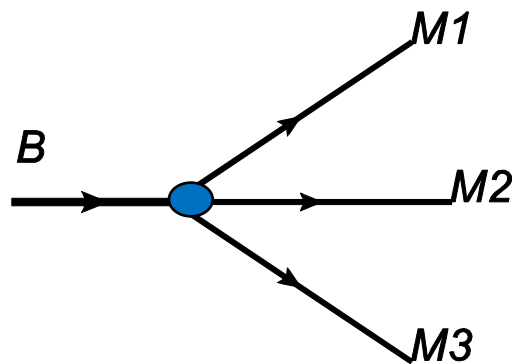
$$\mathcal{A}_{\text{total}} = \sum_i \mathcal{A}_{NR}^i + \sum_j \mathcal{A}_R^j$$



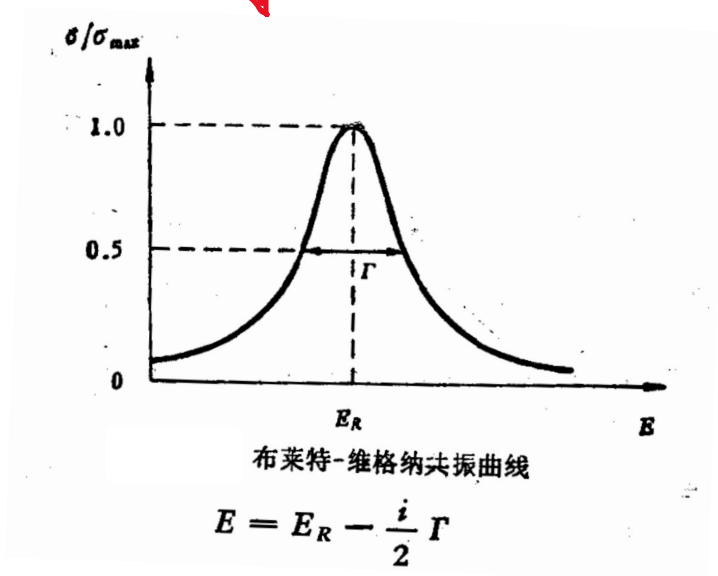
3-body B meson decays and the virtual contribution

$$B^+ \rightarrow K^+ K^- \pi^+$$



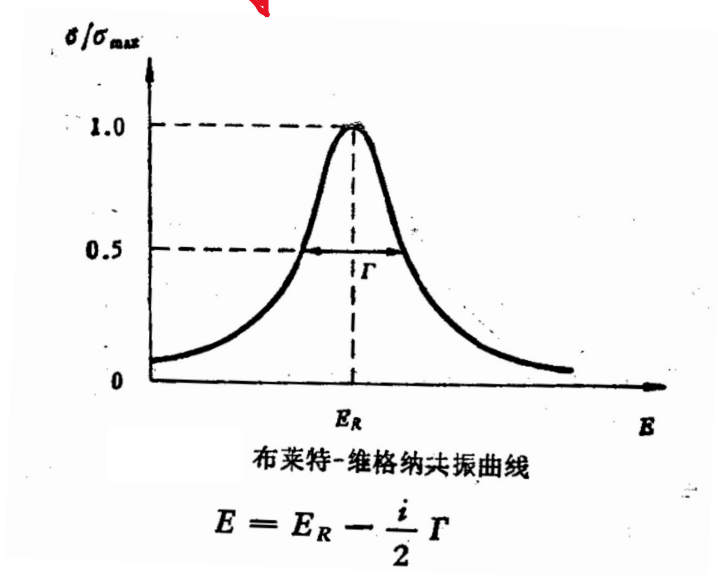
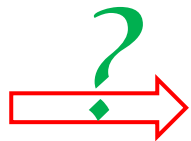
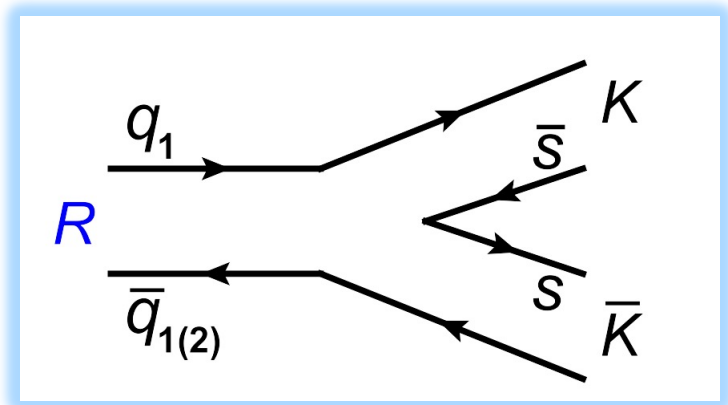
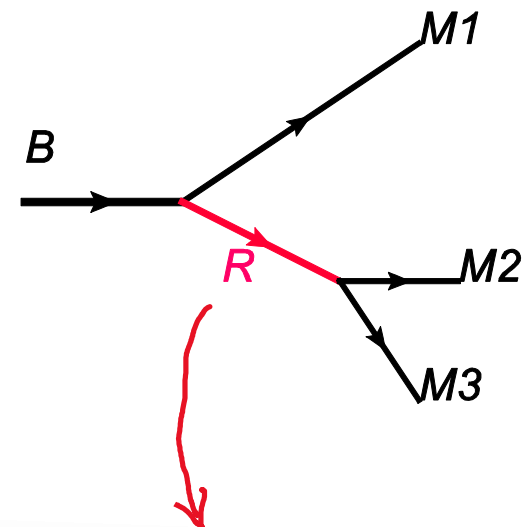
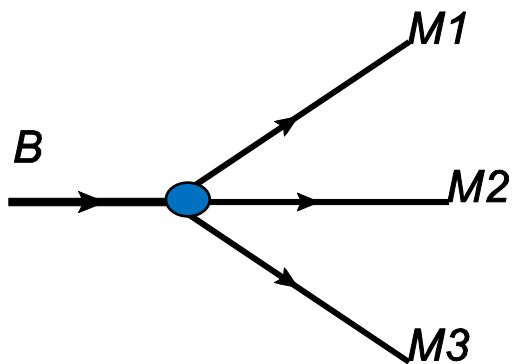


$$BW_R = \frac{m_R^2}{m_R^2 - s - im_R \Gamma_R(s)}$$



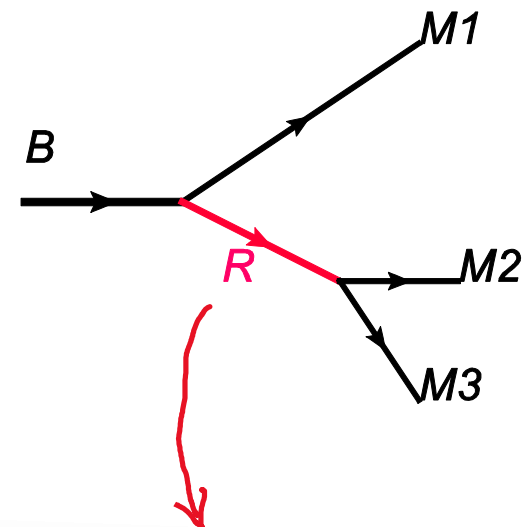
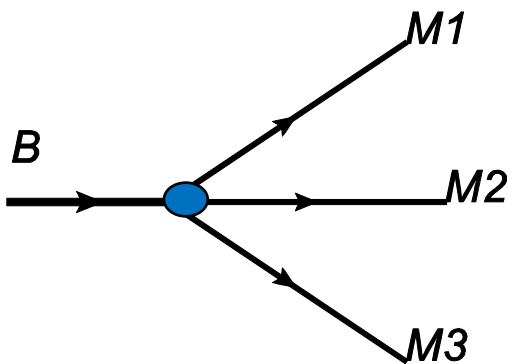
$$B^+ \rightarrow K^+ K^- \pi^+$$

$$\rho(770)^0 \rightarrow K^+ K^-$$



$$B^+ \rightarrow K^+ K^- \pi^+$$

$$\rho(770)^0 \rightarrow K^+ K^-$$

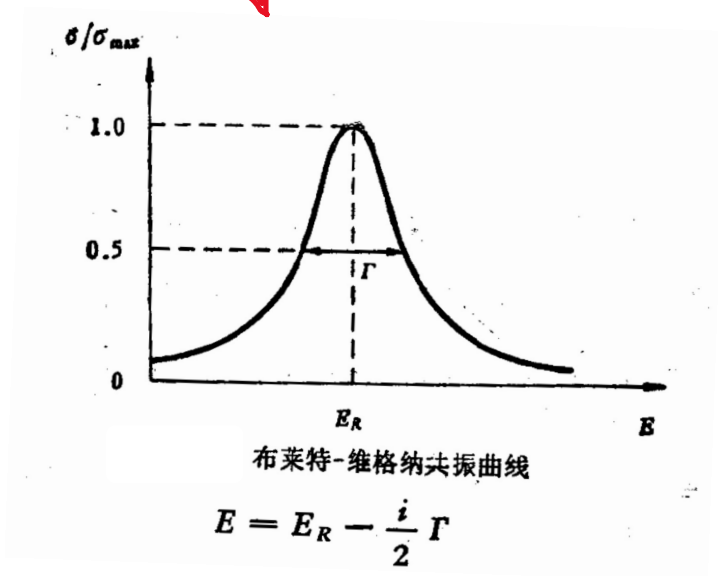


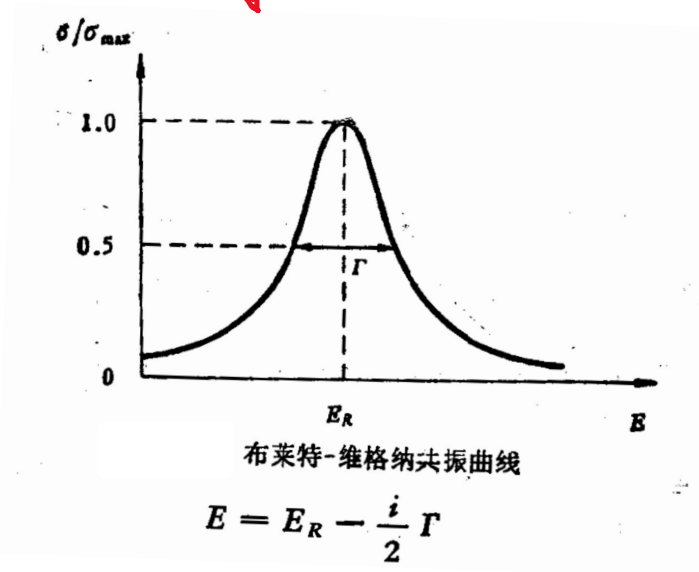
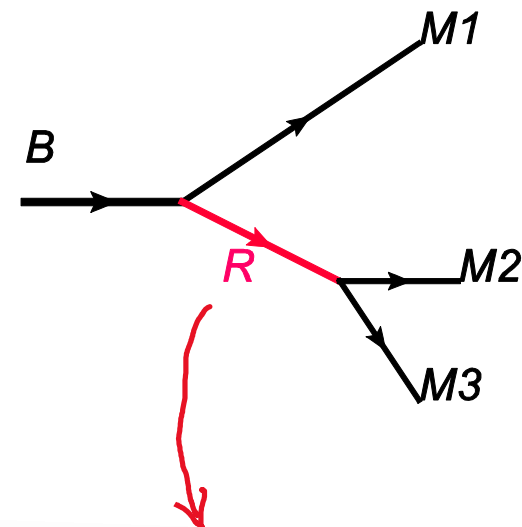
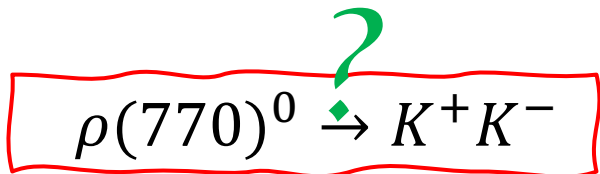
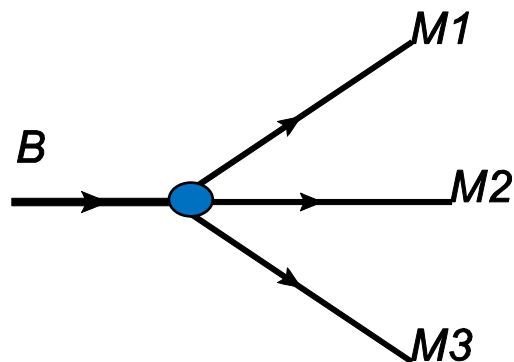
PHYSICAL REVIEW LETTERS **123**, 231802 (2019)

Amplitude Analysis of $B^\pm \rightarrow \pi^\pm K^+ K^-$ Decays

R. Aaij *et al.*
(LHCb Collaboration)

(Received 12 June 2019; revised manuscript received 15 October 2019; published 6 December 2019)





PHYSICAL REVIEW LETTERS **123**, 231802 (2019)

PHYSICAL REVIEW D **101**, 111901(R) (2020)
 Rapid Communications

Will the subprocesses $\rho(770, 1450)^0 \rightarrow K^+ K^-$ contribute large branching fractions for $B^\pm \rightarrow \pi^\pm K^+ K^-$ decays?

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(Received 21 April 2020; accepted 28 May 2020; published 9 June 2020)



PHYSICAL REVIEW D **94**, 072001 (2016)

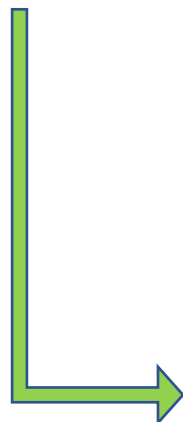
Amplitude analysis of $B^- \rightarrow D^+ \pi^- \pi^-$ decays

R. Aaij *et al.**

(LHCb Collaboration)

(Received 4 August 2016; published 5 October 2016)

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Contribution	Fit fraction (%)
$D_2^*(2460)^0$	35.7 ± 0.6
$D_1^*(2680)^0$	8.3 ± 0.6
$D_3^*(2760)^0$	1.0 ± 0.1
$D_2^*(3000)^0$	0.23 ± 0.07
$D_v^*(2007)^0$	10.8 ± 0.7
B_v^{*0}	2.7 ± 1.0
Total S wave	57.0 ± 0.8
Total fit fraction	115.7



PHYSICAL REVIEW D

VOLUME 15, NUMBER 11

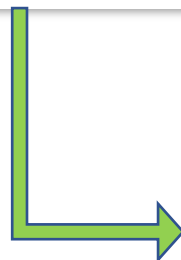
1 JUNE 1977

High-statistics study of the reactions $\pi^- p \rightarrow K^- K^+ n$ and $\pi^+ n \rightarrow K^- K^+ p$ at 6 GeV/c*

A. J. Pawlicki, D. S. Ayres, D. Cohen, R. Diebold, S. L. Kramer, and A. B. Wicklund

Argonne National Laboratory, Argonne, Illinois 60439

(Received 23 December 1976; revised manuscript received 15 March 1977)



ambiguities. As discussed above, the P wave is consistent with the tail of the ρ^0 decaying into $K^- K^+$, with a $\rho K K$ coupling that agrees with SU(3), including the sign. Only one of the ambiguous sol-

PHYSICAL REVIEW D

VOLUME 15, NUMBER 11

1 JUNE 1977

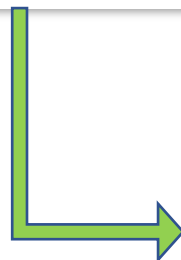
High-statistics study of the reactions $\pi^- p \rightarrow K^- K^+ n$ and $\pi^+ n \rightarrow K^- K^+ p$ at 6 GeV/c*

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BW-Tail



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Physics Letters B

Volume 25, Issue 4, 4 September 1967, Pages 294-297

ELSEVIER



Further study of the $I = 1 K \bar{K}$ structure near threshold

A. Astier, J. Cohen-Ganouna, M. Della Negra, B. Maréchal, L. Montanet, M. Tomas[†], M. Baubillier, J. Duboc



3) The $I = 1 K \bar{K}$ channel is dominated, at threshold, by a virtual bound state resonance

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Kaon pair in $B^+ \rightarrow \bar{D}^{(*)0} K^+ \bar{K}^0$ and $B^0 \rightarrow D^{(*)-} K^+ \bar{K}^0$ decays

MOTIVATION:

Belle-II arXiv:2305.01321

$$\mathcal{B}(B^- \rightarrow D^0 K^- K_S^0) = (1.89 \pm 0.16 \pm 0.10) \times 10^{-4},$$

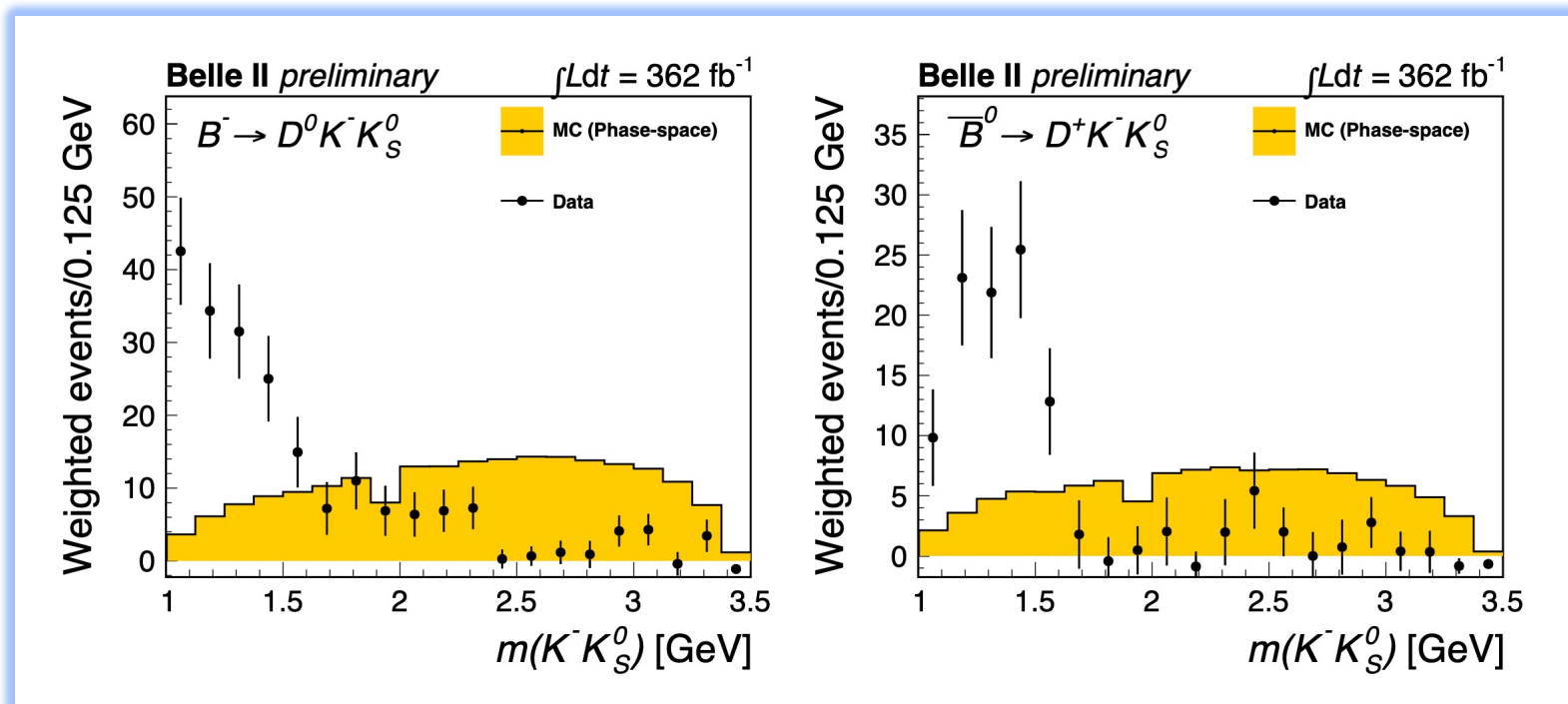
$$\mathcal{B}(\bar{B}^0 \rightarrow D^+ K^- K_S^0) = (0.85 \pm 0.11 \pm 0.05) \times 10^{-4},$$

$$\mathcal{B}(B^- \rightarrow D^{*0} K^- K_S^0) = (1.57 \pm 0.27 \pm 0.12) \times 10^{-4},$$

$$\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} K^- K_S^0) = (0.96 \pm 0.18 \pm 0.06) \times 10^{-4},$$

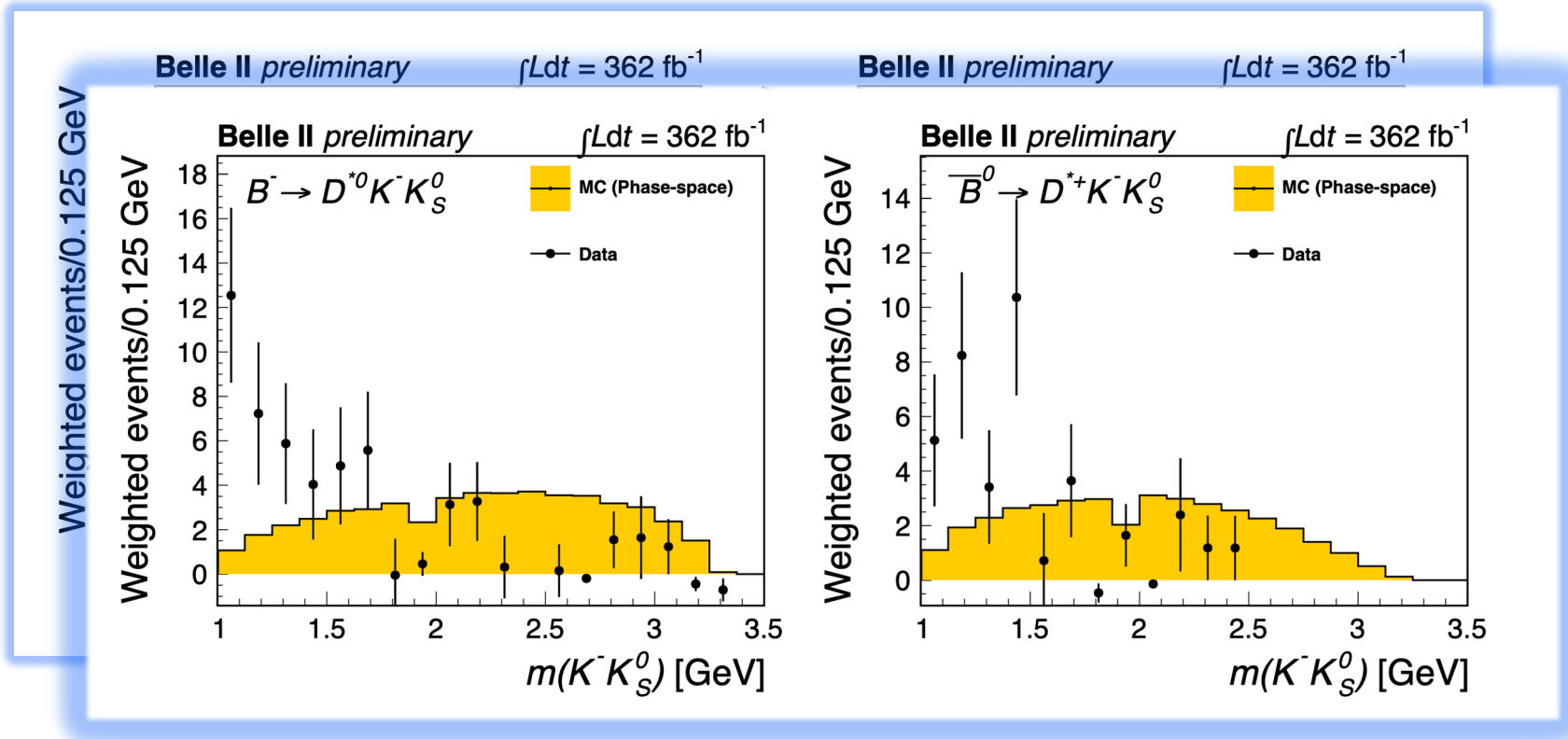
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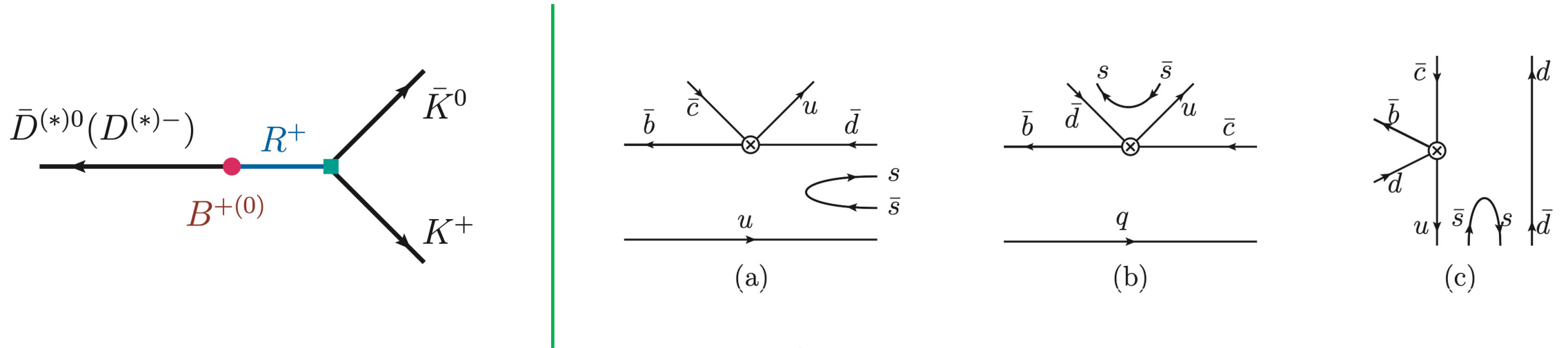
Belle-II arXiv:2305.01321





Kaon pair in $B^+ \rightarrow \bar{D}^{(*)0} K^+ \bar{K}^0$ and $B^0 \rightarrow D^{(*)-} K^+ \bar{K}^0$ decays

APPROACH:



$$\mathcal{A}_V = M_{2B} \cdot \frac{\langle hh' | \rho^+ \rangle}{\mathcal{D}_{\rho^+}(s)}$$

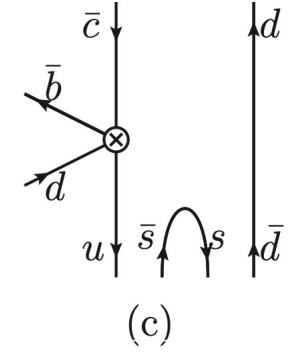
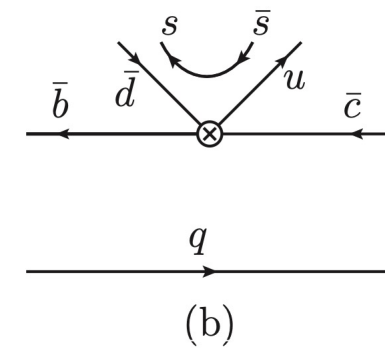
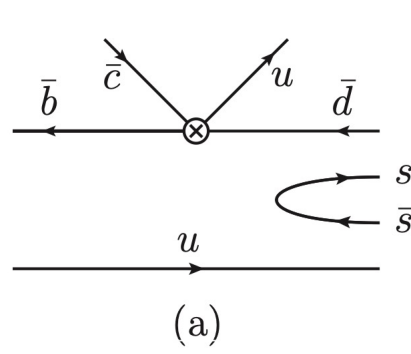
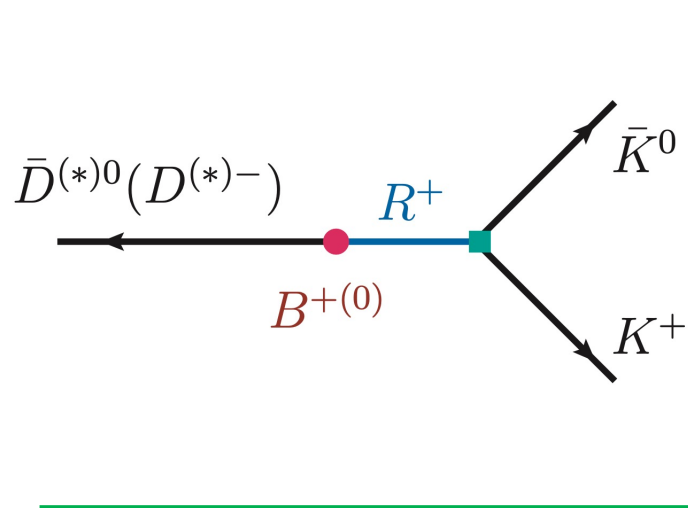
$$F_{\pi, K}^R(s) = c_R^{\pi, K} \text{BW}_R(s) \equiv c_R^{\pi, K} \frac{m_R^2}{\mathcal{D}_R(s)}$$

$$\mathcal{H}_{\text{eff}} = \frac{G_F}{\sqrt{2}} V_{cb}^* V_{ud} [C_1(\mu) O_1^c(\mu) + C_2(\mu) O_2^c(\mu)],$$



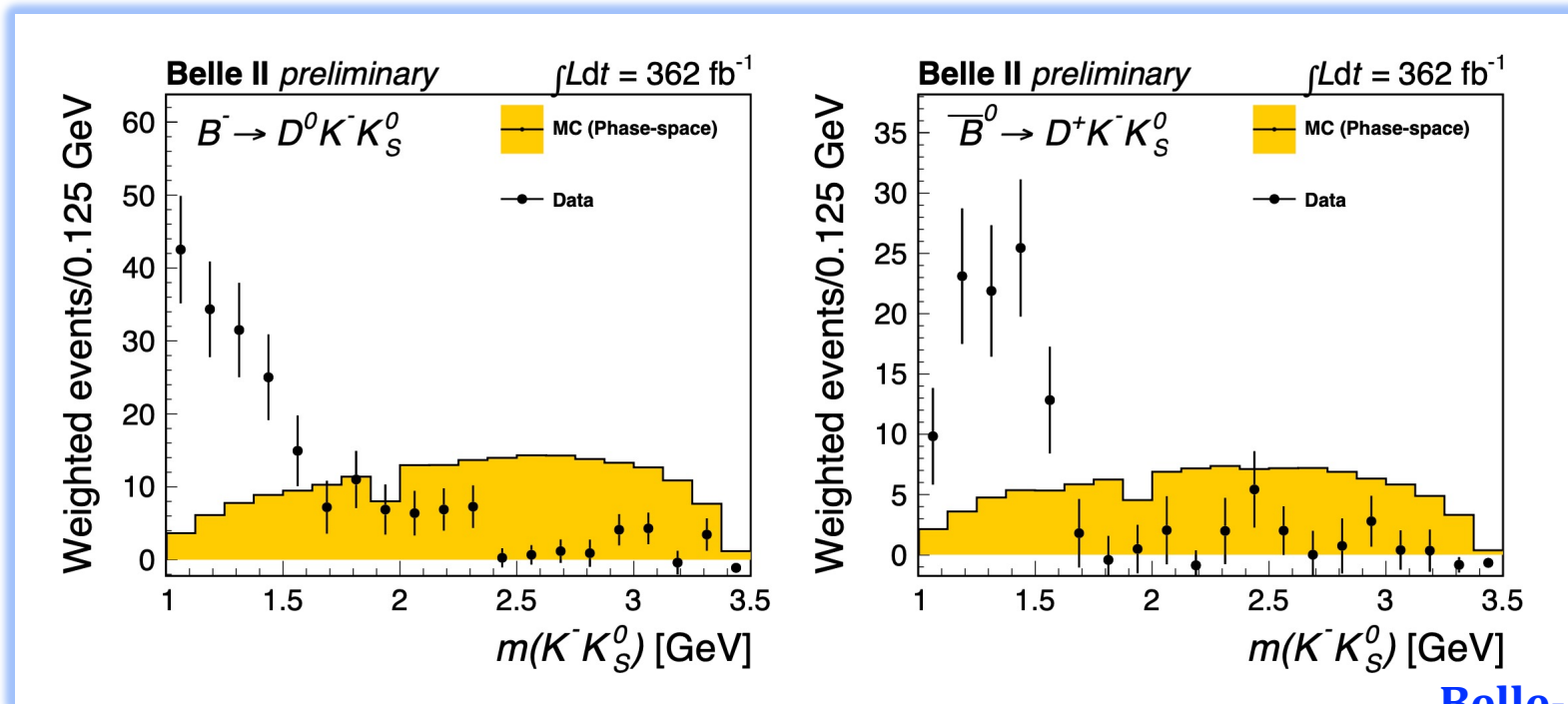
Kaon pair in $B^+ \rightarrow \bar{D}^{(*)0} K^+ \bar{K}^0$ and $B^0 \rightarrow D^{(*)-} K^+ \bar{K}^0$ decays

APPROACH:



Isospin $I = 1$ $K^+ \bar{K}^0$ { $\rho(770)^+ & \rho(1450)^+ & \rho(1700)^+$
 $a_0(980)^+ & a_0(1450)^+$
 $a_2(1320)^+ & a_2(1700)^+$

APPROACH:



[Belle-II arXiv:2305.01321](https://arxiv.org/abs/2305.01321)

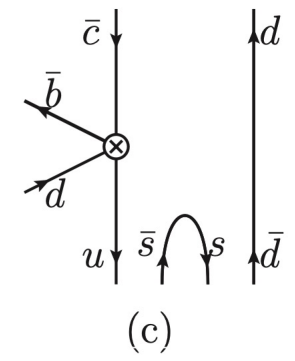
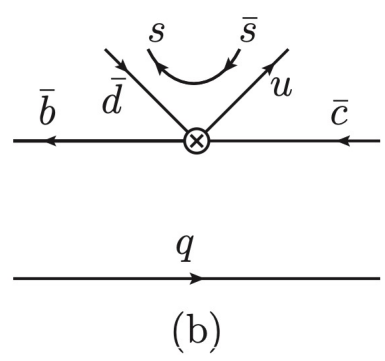
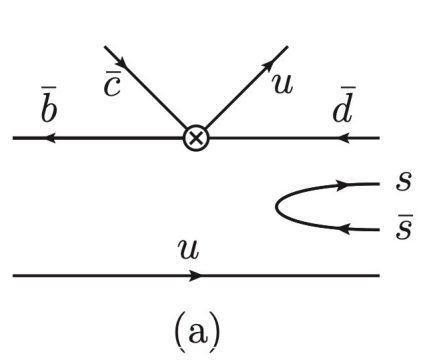
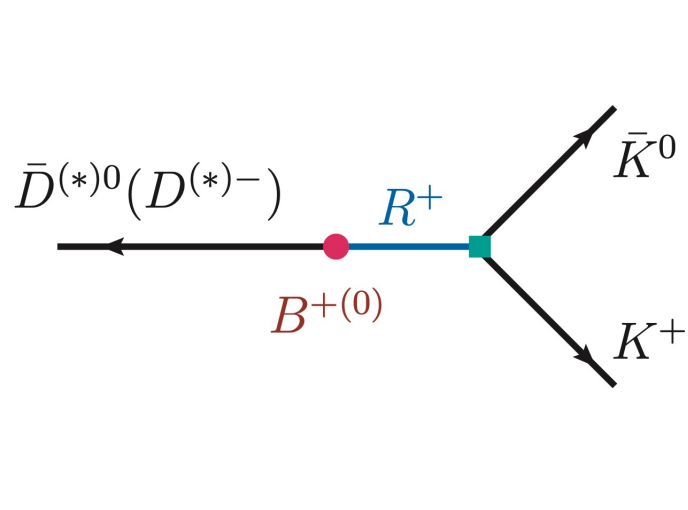
Isospin $I = 1$ $K^+ \bar{K}^0$

- $\rho(770)^+ & \rho(1450)^+ & \rho(1700)^+$
- $a_0(980)^+ & a_0(1450)^+$
- $a_2(1320)^+ & a_2(1700)^+$



Kaon pair in $B^+ \rightarrow \bar{D}^{(*)0} K^+ \bar{K}^0$ and $B^0 \rightarrow D^{(*)-} K^+ \bar{K}^0$ decays

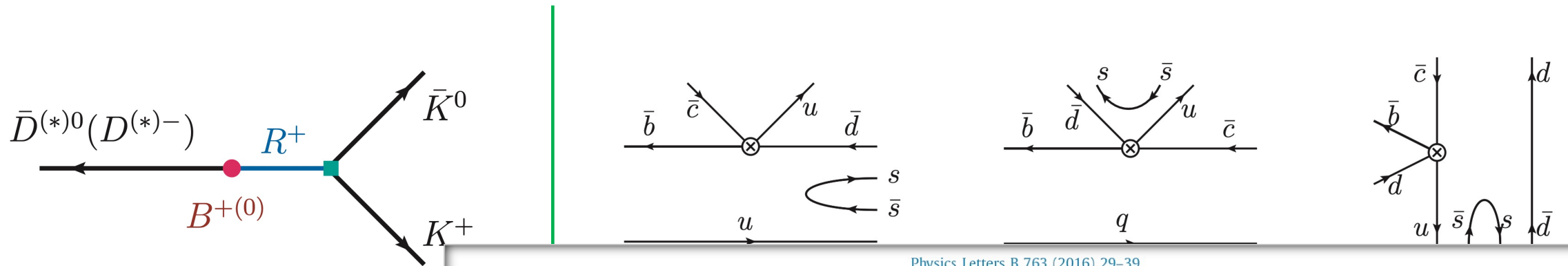
APPROACH:



With PQCD approach

$$\begin{aligned} \mathcal{A} &= \langle (\omega\pi)_{P\text{-wave}} D^{(*)} | \mathcal{H}_{\text{eff}} | B \rangle \\ &= \phi_B \otimes \mathcal{H} \otimes \phi_{\omega\pi}^{P\text{-wave}} \otimes \phi_{D^{(*)}} \end{aligned}$$

APPROACH:



Physics Letters B 763 (2016) 29–39

With PQCD approach



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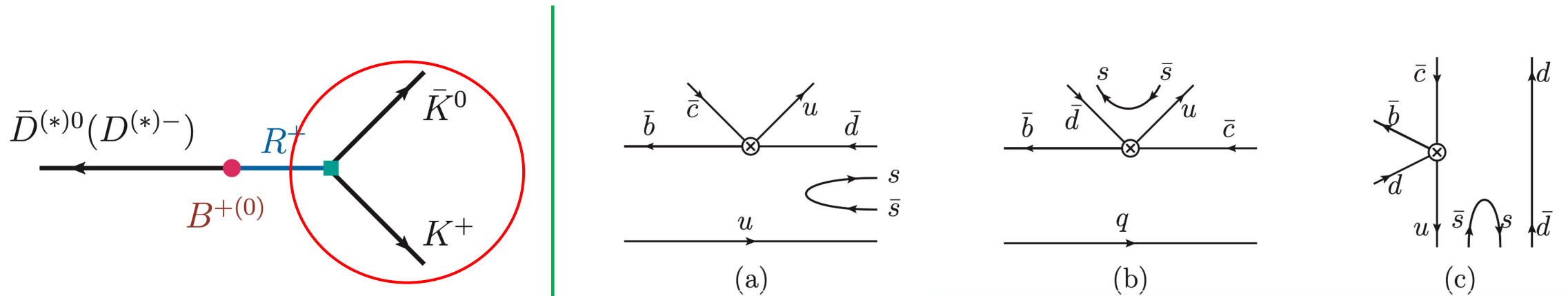
Quasi-two-body decays $B \rightarrow K \rho \rightarrow K \pi \pi$ in perturbative QCD approach

Wen-Fei Wang^{a,b}, Hsiang-nan Li^{a,*}



Kaon pair in $B^+ \rightarrow \bar{D}^{(*)0} K^+ \bar{K}^0$ and $B^0 \rightarrow D^{(*)-} K^+ \bar{K}^0$ decays

APPROACH:



With PQCD approach

$$F_{\pi,K}^R(s) = c_R^{\pi,K} \text{BW}_R(s) \equiv c_R^{\pi,K} \frac{m_R^2}{\mathcal{D}_R(s)}$$

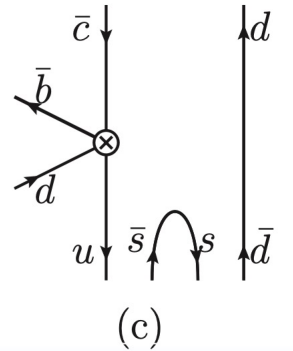
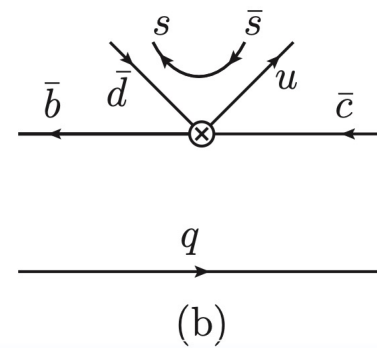
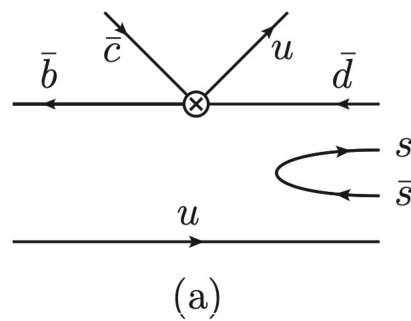
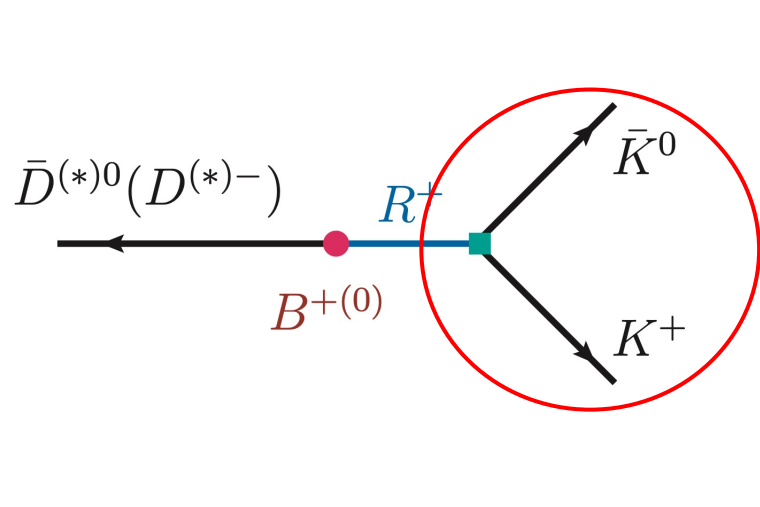
$$\phi_{K\bar{K},L}^{P\text{-wave}}(x, s) = \frac{-1}{\sqrt{2N_c}} [\sqrt{s} \not{\epsilon}_L \phi^0(x, s) + \not{\epsilon}_L \not{p} \phi^t(x, s) + \sqrt{s} \phi^s(x, s)]$$

$$\phi_{K\bar{K},T}^{P\text{-wave}}(x, s) = \frac{-1}{\sqrt{2N_c}} [\sqrt{s} \not{\epsilon}_T \phi^v(x, s) + \not{\epsilon}_T \not{p} \phi^T(x, s) + \sqrt{s} i \epsilon_{\mu\nu\rho\sigma} \gamma_5 \gamma^\mu \epsilon_T^{*\nu} n^\rho v^\sigma \phi^a(x)]$$



Kaon pair in $B^+ \rightarrow \bar{D}^{(*)0} K^+ \bar{K}^0$ and $B^0 \rightarrow D^{(*)-} K^+ \bar{K}^0$ decays

APPROACH:



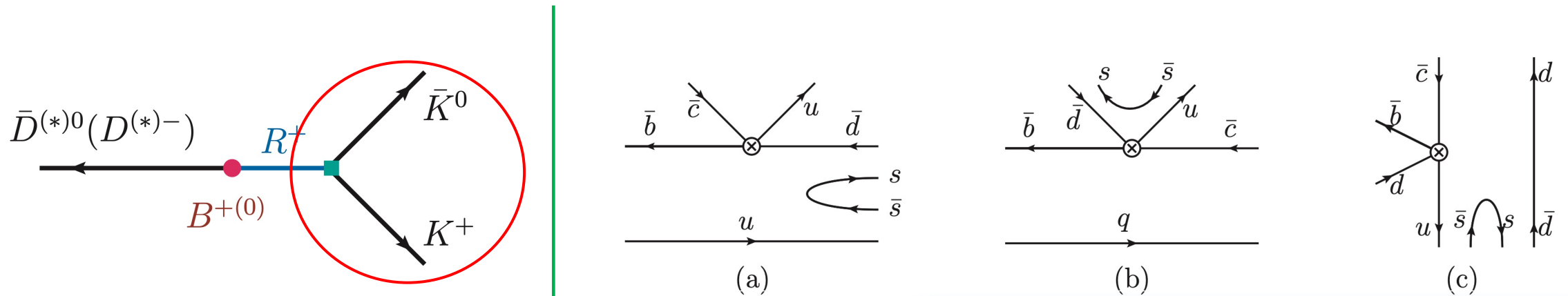
With PQCD approach

$$F_{\pi,K}^R(s) = c_R^{\pi,K} \text{BW}_R(s) \equiv c_R^{\pi,K} \frac{m_R^2}{\mathcal{D}_R(s)}$$

$$f_R g_{R\pi\pi} / (\sqrt{2} m_R)$$

$$g_{\rho(770)^0 K^+ K^-} = g_{\rho(770)^0 \pi^+ \pi^-} / 2$$

APPROACH:



With PQCD approach

$$F_{\pi,K}^R(s) = c_R^{\pi,K} \text{BW}_R(s) \equiv c_R^{\pi,K} \frac{m_R^2}{\mathcal{D}_R(s)}$$

PHYSICAL REVIEW D **103**, 056021 (2021)

Contributions for the kaon pair from $\rho(770)$, $\omega(782)$ and their excited states in the $B \rightarrow K\bar{K}h$ decays

Wen-Fei Wang^{1,2,*}



$$c_{\rho(770)}^K = 1.247 \pm 0.019$$

$$c_{\rho(1450)}^K = -0.156 \pm 0.015$$

TABLE II: PQCD results for the branching fractions of the quasi-two-body decays $B^+ \rightarrow \bar{D}^{(*)0}[\rho(770)^+ \rightarrow]\pi^+\pi^0$ and $B^0 \rightarrow D^{(*)-}[\rho(770)^+ \rightarrow]\pi^+\pi^0$.

Decay modes	Units	PQCD
$B^+ \rightarrow \bar{D}^0[\rho(770)^+ \rightarrow]\pi^+\pi^0$	%	$1.21^{+0.16+0.10+0.05}_{-0.16-0.12-0.06}$
$B^0 \rightarrow D^-[\rho(770)^+ \rightarrow]\pi^+\pi^0$	10^{-3}	$7.63^{+0.58+0.97+0.34}_{-0.58-0.73-0.21}$
$B^+ \rightarrow \bar{D}^{*0}[\rho(770)^+ \rightarrow]\pi^+\pi^0$	10^{-3}	$9.03^{+1.55+0.73+0.51}_{-1.55-0.64-0.46}$
$B^0 \rightarrow D^{*-}[\rho(770)^+ \rightarrow]\pi^+\pi^0$	10^{-3}	$8.15^{+1.31+0.64+0.03}_{-1.31-0.62-0.07}$

$$\begin{aligned} \mathcal{B}(B^+ \rightarrow \bar{D}^0 \rho(770)^+) &= (1.34 \pm 0.18)\%, \\ \mathcal{B}(B^0 \rightarrow D^- \rho(770)^+) &= (7.6 \pm 1.2) \times 10^{-3}, \\ \mathcal{B}(B^+ \rightarrow \bar{D}^{*0} \rho(770)^+) &= (9.8 \pm 1.7) \times 10^{-3}, \\ \mathcal{B}(B^0 \rightarrow D^{*-} \rho(770)^+) &= (6.8 \pm 0.9) \times 10^{-3}, \end{aligned}$$

PDG-2022

TABLE III: PQCD predictions for the branching fractions of the concerned quasi-two-body decays with the subprocess $\rho^+ \rightarrow K^+ \bar{K}^0$, here $\rho^+ = \rho(770)^+ + \rho(1450)^+$.

Mode	Unit	\mathcal{B}
$B^+ \rightarrow \bar{D}^0[\rho^+ \rightarrow]K^+ \bar{K}^0$	10^{-4}	$1.68^{+0.20+0.17+0.12}_{-0.20-0.15-0.12}$
$B^0 \rightarrow D^-[\rho^+ \rightarrow]K^+ \bar{K}^0$	10^{-4}	$0.98^{+0.06+0.13+0.06}_{-0.06-0.12-0.06}$
$B^+ \rightarrow \bar{D}^{*0}[\rho^+ \rightarrow]K^+ \bar{K}^0$	10^{-4}	$1.33^{+0.21+0.11+0.05}_{-0.21-0.11-0.07}$
$B^0 \rightarrow D^{*-}[\rho^+ \rightarrow]K^+ \bar{K}^0$	10^{-4}	$1.16^{+0.19+0.08+0.02}_{-0.19-0.09-0.02}$

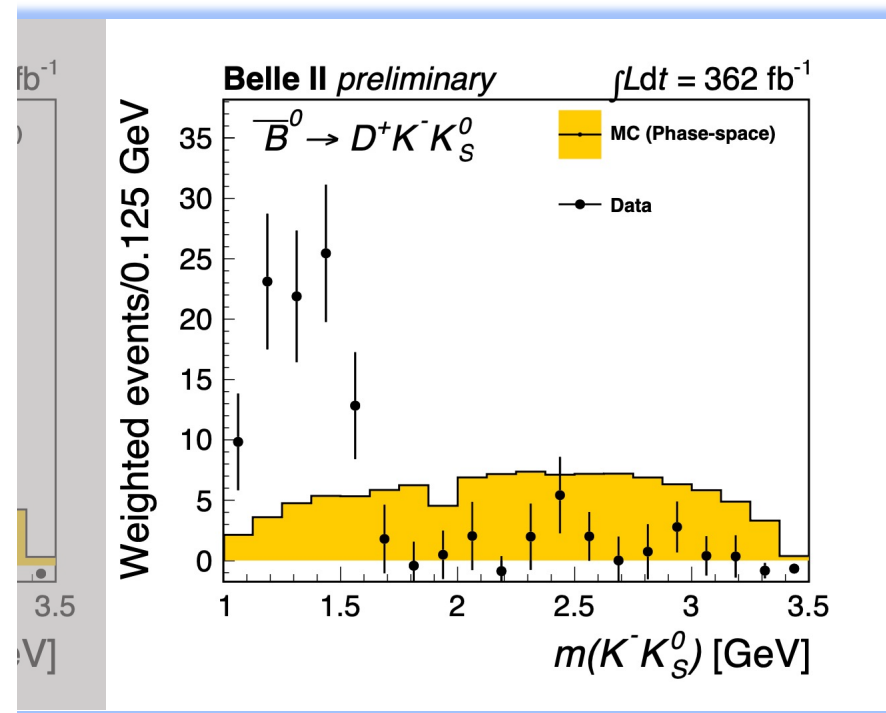
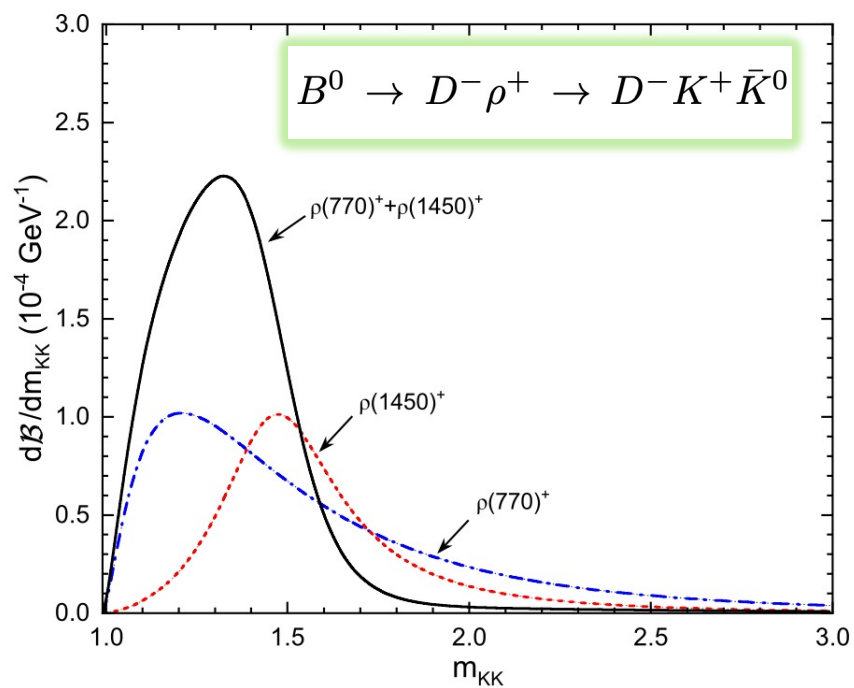
Belle-II arXiv:2305.01321

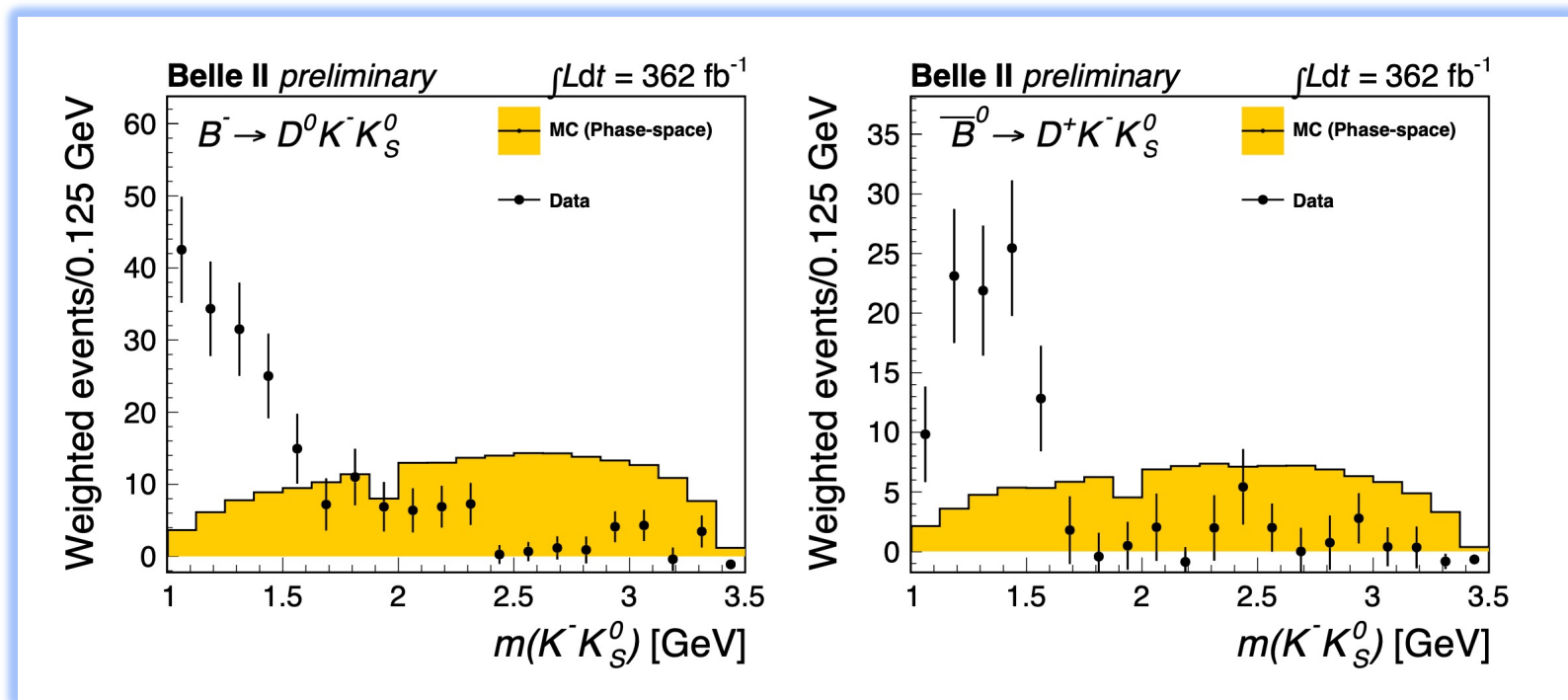
$$\mathcal{B}(B^- \rightarrow D^0 K^- K_S^0) = (1.89 \pm 0.16 \pm 0.10) \times 10^{-4},$$

$$\mathcal{B}(\bar{B}^0 \rightarrow D^+ K^- K_S^0) = (0.85 \pm 0.11 \pm 0.05) \times 10^{-4},$$

$$\mathcal{B}(B^- \rightarrow D^{*0} K^- K_S^0) = (1.57 \pm 0.27 \pm 0.12) \times 10^{-4},$$

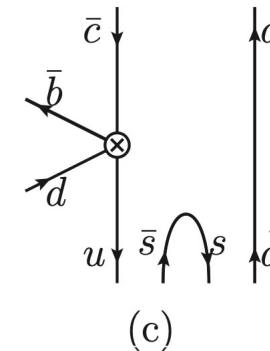
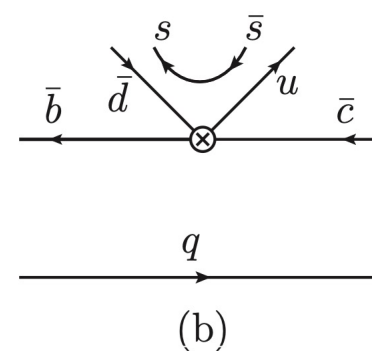
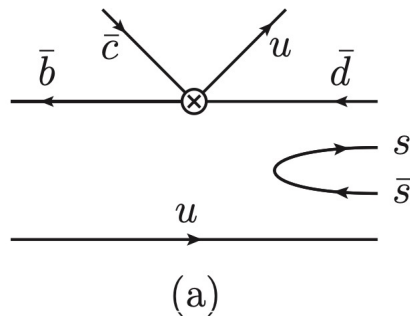
$$\mathcal{B}(\bar{B}^0 \rightarrow D^{*+} K^- K_S^0) = (0.96 \pm 0.18 \pm 0.06) \times 10^{-4},$$





CONTRIBUTIONS:

$a_0(980)^+ & a_0(1450)^+$



$$\langle K^+ \bar{K}^0 | (\bar{u}d)_{V-A} | 0 \rangle \implies f_{a_0(980)} \approx 1.1 \text{ MeV}$$

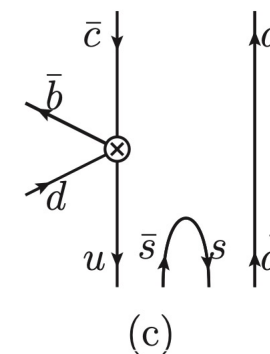
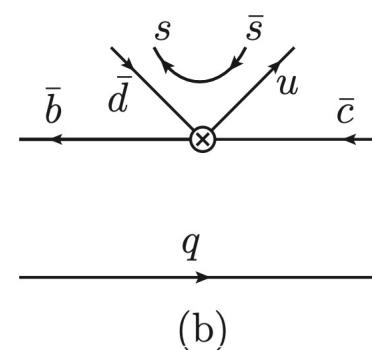
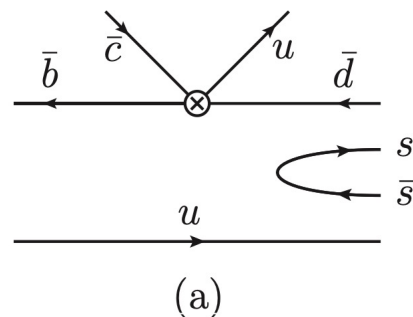
$$\mu_S f_S = \bar{f}_S, \quad \text{with} \quad \mu_S = \frac{m_S}{m_2(\mu) - m_1(\mu)}$$

$$\bar{f}_{a_0}(1 \text{ GeV}) \approx 385 \text{ MeV}$$

PRD73-014017

CONTRIBUTIONS:

$a_0(980)^+ & a_0(1450)^+$



$$\mathcal{B}(B^0 \rightarrow D_s^+ a_0(980)^-) = 1.93 \times 10^{-5}$$

upper limit 1.9×10^{-5} at 90% C.L. presented by the *BABAR* $\Rightarrow 2.24 \times 10^{-5}$

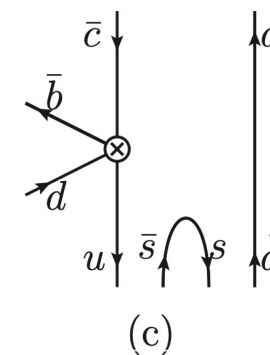
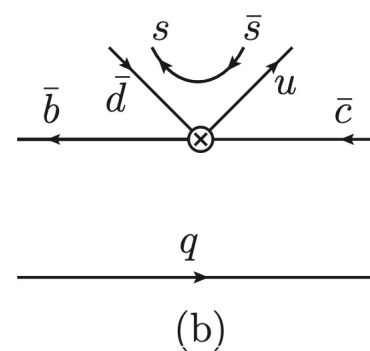
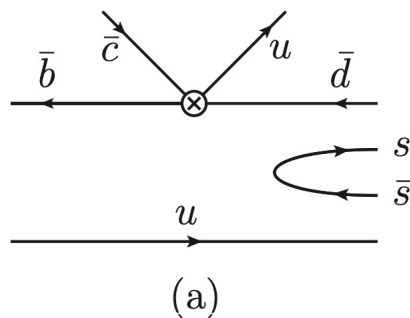
PRD73-071103

$$\mathcal{B} = 0.72 \times 10^{-5} \text{ for } B^+ \rightarrow \bar{D}^0 a_0(1450)^+ \rightarrow \bar{D}^0 K^+ \bar{K}^0$$

$$\mathcal{B} = 1.56 \times 10^{-5} \quad B^+ \rightarrow \bar{D}^0 a_0(980)^+ \rightarrow \bar{D}^0 K^+ \bar{K}^0$$

CONTRIBUTIONS:

$a_2(1320)^+$ & $a_2(1700)^+$



$a_2(1700)$

$$I^G(J^{PC}) = 1^-(2^{++})$$

Mass $m = 1698 \pm 40$ MeV

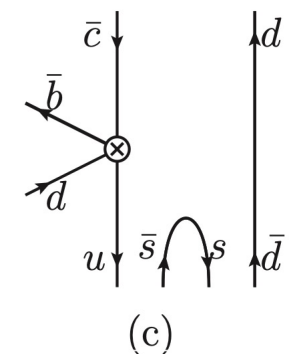
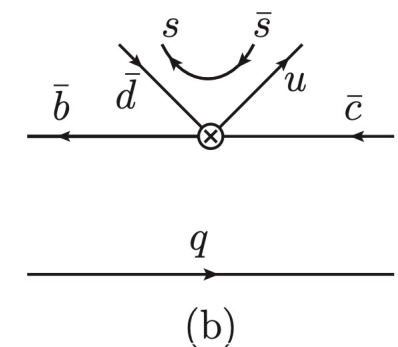
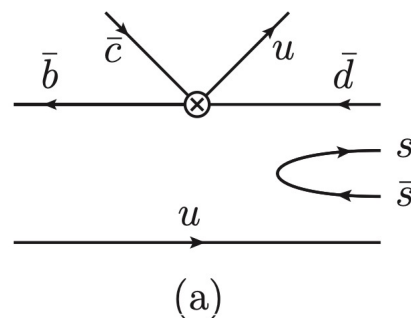
Full width $\Gamma = 265 \pm 60$ MeV

$a_2(1700)$ DECAY MODES

	Fraction (Γ_i/Γ)
$\eta\pi$	$(3.6 \pm 1.1) \%$
$\gamma\gamma$	$(1.13 \pm 0.30) \times 10^{-6}$
$K\bar{K}$	$(1.9 \pm 1.2) \%$

CONTRIBUTIONS:

$a_2(1320)^+$ & ~~$a_2(1700)^+$~~



$a_2(1320)$

$$I^G(J^{PC}) = 1^-(2^{++})$$

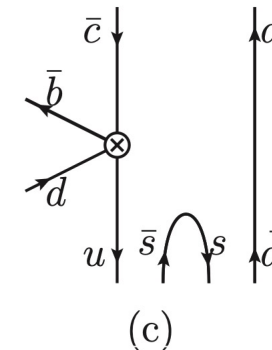
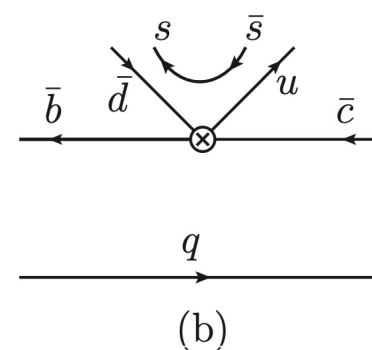
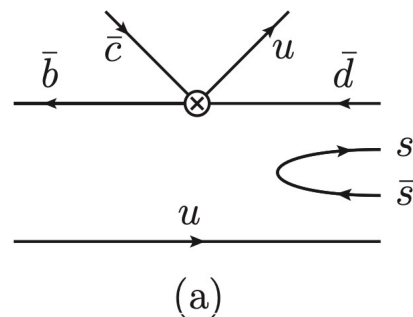
Mass $m = 1318.2 \pm 0.6$ MeV (S = 1.2)

Full width $\Gamma = 107 \pm 5$ MeV [i]

$a_2(1320)$ DECAY MODES	Fraction (Γ_i/Γ)	Scale factor/ Confidence level	p (MeV/c)
3π	(70.1 \pm 2.7) %	S=1.2	624
$\eta\pi$	(14.5 \pm 1.2) %		535
$\omega\pi\pi$	(10.6 \pm 3.2) %	S=1.3	366
$K\bar{K}$	(4.9 \pm 0.8) %		437

CONTRIBUTIONS:

$a_2(1320)^+$ & ~~$a_2(1700)^+$~~

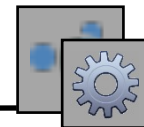


$$B_s^0 \rightarrow \bar{D}^0 [\bar{K}_2^*(1430)^0 \rightarrow] K^- \pi^+ \quad \mathcal{B} = (3.7 \pm 1.4) \times 10^{-5}$$

PRD90-072003 LHCb

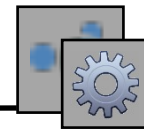
$s \rightarrow u$

$$B^+ \rightarrow \bar{D}^0 a_2(1320)^+ \quad \mathcal{B} = (0.99 \pm 0.37) \times 10^{-4}$$



For the Kaon pair in $B^+ \rightarrow \bar{D}^{(*)0} K^+ \bar{K}^0$ and $B^0 \rightarrow D^{(*)-} K^+ \bar{K}^0$ decays

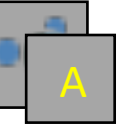
- (i). The intermediate states $\rho(770)^+$ & $\rho(1450)^+$ dominate the branching fractions for the four corresponding decay channels;
- (ii). The role of $a_2(1320)^+$ in these four decay channels is negligible;
- (ii). The state $a_0(980)^+$ turned out to be less important than expected for the kaon pair near the threshold;



For the Kaon pair in $B^+ \rightarrow \bar{D}^{(*)0} K^+ \bar{K}^0$ and $B^0 \rightarrow D^{(*)-} K^+ \bar{K}^0$ decays

- (i). The intermediate states $\rho(770)^+$ & $\rho(1450)^+$ dominate the branching fractions for the four corresponding decay channels;
- (ii). The role of $a_2(1320)^+$ in these four decay channels is negligible;
- (ii). The state $a_0(980)^+$ turned out to be less important than expected for the kaon pair near the threshold;

Thank You !



Appendix

Belle-II arXiv:2305.01321

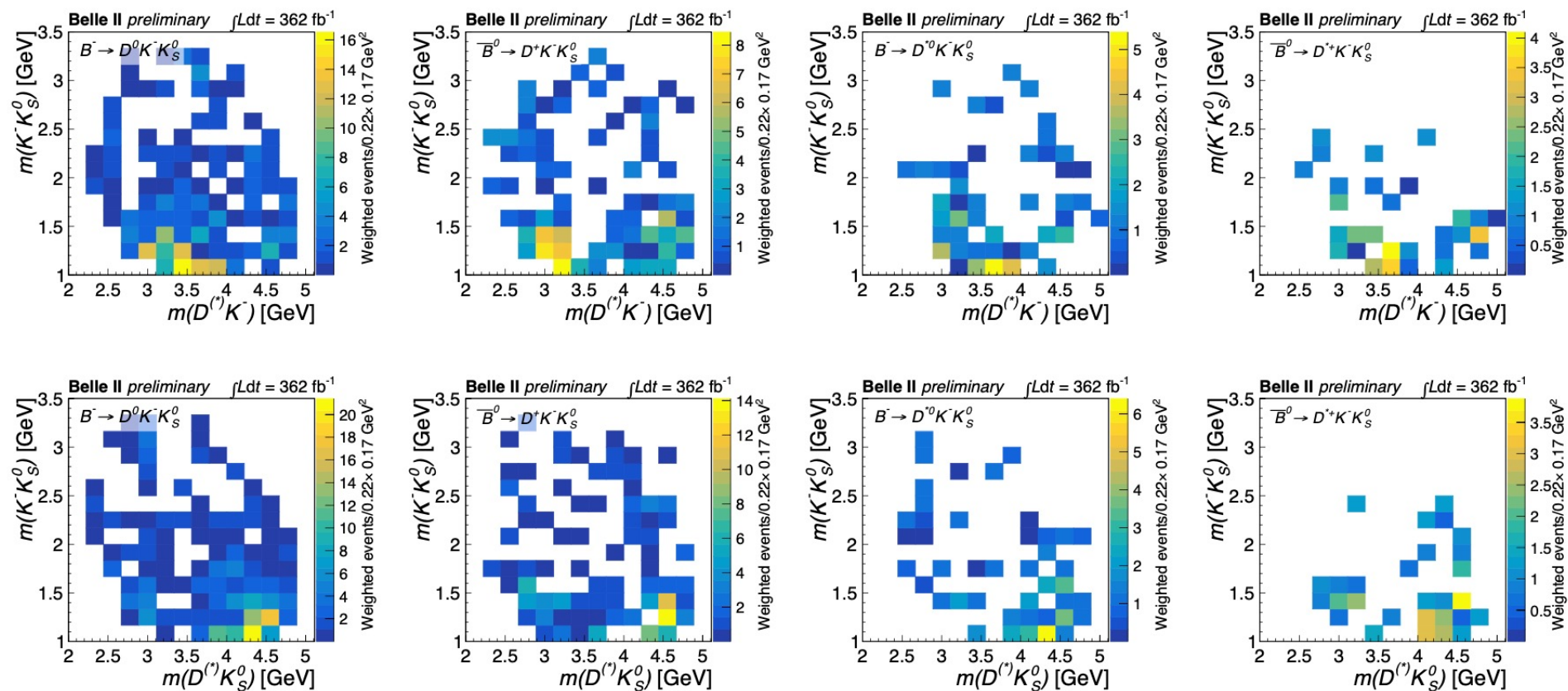


Figure 4. Dalitz distributions of $(m(D^*K^-), m(K^-K_S^0))$ (upper panels) and $(m(D^*K_S^0), m(K^-K_S^0))$ (lower panels) for (left to right) $B^- \rightarrow D^0K^-K_S^0$, $\bar{B}^0 \rightarrow D^+K^-K_S^0$, $B^- \rightarrow D^{*0}K^-K_S^0$, and $\bar{B}^0 \rightarrow D^{*+}K^-K_S^0$ channels. The background is subtracted by applying the signal *sWeight*.



Belle-II preliminary

Channel	Yield (K_S^0 / K^{*0})	Average ϵ (K_S^0 / K^{*0})	\mathcal{B} [10^{-4}]
$B^- \rightarrow D^0 K^- K_S^0$	209 ± 17	0.098	$1.82 \pm 0.16 \pm 0.08$
$\bar{B}^0 \rightarrow D^+ K^- K_S^0$	105 ± 14	0.048	$0.82 \pm 0.12 \pm 0.05$
$B^- \rightarrow D^{*0} K^- K_S^0$	51 ± 9	0.044	$1.47 \pm 0.27 \pm 0.10$
$\bar{B}^0 \rightarrow D^{*+} K^- K_S^0$	36 ± 7	0.046	$0.91 \pm 0.19 \pm 0.05$

first observation

13



Belle-II preliminary

