

On-Shell Construction of Chiral Effective Theory

— supported by package ABC4EFT¹

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I. Low, J. Shu, **MLX** and Y.-H. Zheng, *JHEP* 01 (2023) 031

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¹<https://abc4eft.hepforge.org/>

- 1 Effective Field Theory of Goldstone Bosons
- 2 Building Soft Blocks
- 3 Results and Summary

Outline

- 1 **Effective Field Theory of Goldstone Bosons**
- 2 Building Soft Blocks
- 3 Results and Summary

The Non-Linear Symmetry

Symmetry breaking $\mathcal{G} \rightarrow \mathcal{H}$ with Goldstone fields $\xi = e^{i\pi^a X^a / f} \in \mathcal{G}/\mathcal{H}$.

- Although spontaneously broken, \mathcal{G} is realized **non-linearly** at low energy

$$\pi \rightarrow \pi + \epsilon$$

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$$\pi \rightarrow \pi + \epsilon - \frac{1}{3}\mathcal{T}(\pi)\epsilon + \dots$$

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- The bottom-up approach: **Adler Zero** (Low 14')

$$d_\mu = \frac{1}{f} \frac{\sin\sqrt{\mathcal{T}}}{\sqrt{\mathcal{T}}} \partial_\mu \pi, \quad e_\mu^i = \frac{2i}{f^2} \partial_\mu \pi \frac{\sin^2(\sqrt{\mathcal{T}}/2)}{\mathcal{T}} T^i \pi.$$

where $\mathcal{T}_{ab} = f^{-2} \sum_i (T^i \pi)_a (T^i \pi)_b$, without any knowledge of the broken group \mathcal{G} .

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- \mathcal{G} -invariant (\mathcal{H} -invariant + Adler Zero) Chiral Lagrangian

$$\mathcal{L}^{(2)} = \frac{f^2}{2} \text{Tr}(d_\mu d^\mu) \supset \pi^2 \partial^2, \pi^4 \partial^2, \pi^6 \partial^2, \dots$$

Chiral Perturbation Theory (ChPT)

- Consider $\mathcal{G} = SU(N_f)_L \times SU(N_f)_R \rightarrow SU(N_f)_V = \mathcal{H}$ by $\langle \bar{q}q \rangle \neq 0$
- External sources in QCD: $\mathcal{L} = \mathcal{L}_{\text{QCD}}^0 + \bar{q}\gamma^\mu(v_\mu + a_\mu\gamma^5)q - \bar{q}(s - ip\gamma^5)q$.

Hidden Local Symmetry: $\nabla_\mu X = \partial_\mu X - iA_\mu(X)$, $A \in \{v, a, e\}$

- Two ways to parameterize the chiral Lagrangian:
 - Left-right basis: $U = \xi^2$, $\chi \sim s + ip$, $F_L^{\mu\nu}$, $F_R^{\mu\nu}$.
 - u basis: $u_\mu \sim d_\mu$, χ_\pm , $f_\pm^{\mu\nu}$.
- Redundancy relations to build Lagrangian:
 - IBP: $(\nabla X_1)X_2 + X_1(\nabla X_2) \simeq 0$.
 - EOM (field redef.): $\nabla_\mu u^\mu = \frac{i}{2} \left(\chi_- - \frac{1}{N_f} \langle \chi_- \rangle \right)$.
 - CDC: $[\nabla_\mu, \nabla_\nu] = \frac{1}{4}[u_\mu, u_\nu] - \frac{i}{2}f_{+\mu\nu}$, $\nabla_{[\mu}u_{\nu]} + f_{-\mu\nu} = 0$.
 - Schouten Identity (specific for fixed $D = 4$): $\mathcal{O}_{[\mu_1\mu_2\mu_3\mu_4\mu_5]} = 0$.
- $O(p^4)$ (Gasser&Leutwyler 84'), $O(p^6)$ (Bijnens *et al.* 99'), $O(p^8)$ (Bijnens *et al.* 19').

On-Shell Construction of ChPT

- It is known for more than half a century that the non-linear constraint can be reproduced in the amplitude by Adler Zero: $\lim_{p_i=0} \mathcal{M} = 0$ (Susskind&Frye 69'):

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Algebraic Aspects of Pionic Duality Diagrams

LEONARD SUSSKIND* AND GRAHAM FRYE

Belfer Graduate School of Science, Yeshiva University, New York, New York 10033

(Received 9 May 1969)

Certain algebraic aspects are abstracted from the duality principle and are incorporated in a simple model of pion n -point functions. An algorithm for constructing the n -point function in the tree-graph approximation is based on the duality assumption and the Adler condition which states that the amplitudes vanishes if any pion four-momentum vanishes, all others remaining on shell. The resulting amplitudes satisfy the constraints of current algebra and partial conservation of axial-vector current for $n=4, 6,$ and $8,$ and (we conjecture) for all n . In addition, duality specifies a definite form for chiral symmetry breaking.

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$$\mathcal{L}^{(2)} = \frac{f^2}{2} \text{Tr}(d_\mu d^\mu) \sim \frac{1}{2} (\partial\pi)^2 + \frac{\pi^2}{2f^2} (\partial\pi)^2 + \dots$$

$$\begin{aligned} \mathcal{M}(p_1, \dots, p_6) &= \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} \\ &= \frac{1}{f^4} \left(\frac{s_{13}s_{46}}{s_{123}} + \frac{s_{24}s_{51}}{s_{234}} + \frac{s_{35}s_{62}}{s_{345}} \right) \end{aligned}$$

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- Extended to arbitrary n : **Soft Recursion Relation** (Cheung *et al.* 15’)

$$p_i \rightarrow \hat{p}_i = (1 - a_i z) p_i, \quad \mathcal{M}(\hat{p}_1, \dots, \hat{p}_n) \equiv \hat{\mathcal{M}}_n(z),$$

Poles at $\hat{s}_I(z_I^\pm) = 0$, thus the Cauchy's theorem gives

$$\mathcal{M}(p_1, \dots, p_n) = \hat{\mathcal{M}}_n(0) = \sum_{I, \pm} \frac{1}{s_I} \frac{\hat{\mathcal{M}}_L^{(I)}(z_I^\pm) \hat{\mathcal{M}}_R^{(I)}(z_I^\pm)}{F_n(z_I^\pm) (1 - z_I^\pm / z_I^\mp)}, \quad F_n(z) \equiv \prod_{i=1}^n (1 - a_i z).$$

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- How about higher derivatives?** Need new inputs (LEC)! (Low&Yin 19’)

$$\text{Single Trace: } \mathcal{S}_1^{(4)}(1, 2, 3, 4) = \frac{c_1}{\Lambda^2 f^2} s_{13}^2, \quad \mathcal{S}_2^{(4)}(1, 2, 3, 4) = \frac{c_2}{\Lambda^2 f^2} s_{13} s_{23},$$

$$\text{Double Trace: } \mathcal{S}_1^{(4)}(1, 2|3, 4) = \frac{d_1}{\Lambda^2 f^2} s_{13}^2, \quad \mathcal{S}_2^{(4)}(1, 2|3, 4) = \frac{d_2}{\Lambda^2 f^2} s_{13} s_{23}.$$

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(1) All soft blocks \mathcal{S} satisfy Adler Zero.

(2) 1-to-1 correspond to Lagrangian terms $\mathcal{L}^{(4)} = \sum_{i=1}^4 \frac{L_{4,i}}{\Lambda^2 f^2} \mathcal{O}_i$.

Outline

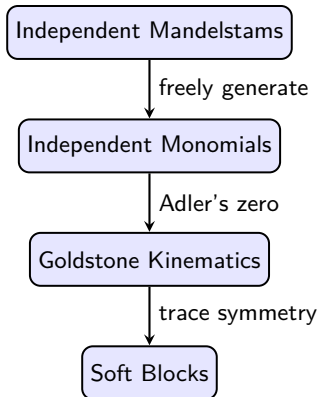
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Soft Block Construction: First Attempt

Systematic construction of soft blocks (Dai *et al.* 20'):

- Only P even Building block s_{ij} (no $\epsilon^{\mu\nu\rho\sigma}$).
- Mandelstams constrained by Gram determinant
 - $D = 4$ at $O(p^{10})$: $\det s_{ij} = 0$
- The resulting soft blocks are polynomials.
 - Do not correspond to monomial operator
- Assuming independent traces (large N_f).
 - Cayley-Hamilton when $N_f < N$.

$$\text{e.g. } \langle A^4 \rangle = \frac{1}{2} \langle A^2 \rangle^2, \quad \forall A \in \mathfrak{su}(3)$$



On-Shell Operator Construction in a Broader Sense

Amplitude/Operator basis correspondence (Shu, Ma, MLX 19')

- Operators \mathcal{O} are **isomorphic** to **local gauge invariant on-shell** amplitudes $\mathcal{B} = \mathcal{TM}$
 - IBP \simeq momentum conservation $\sum_i p_i = \sum_i |i\rangle[i] = 0$.
 - EOM \simeq on-shell condition $\langle ii\rangle = [ii] = 0$.
 - $D = 4$ Schouten Identity / Gram Determinants \simeq spinor Schouten Identities.

$$\langle ij\rangle\langle kl\rangle + \langle ik\rangle\langle lj\rangle + \langle il\rangle\langle jk\rangle = 0, \quad [ij][kl] + [ik][lj] + [il][jk] = 0.$$

An independent operator basis \simeq an independent amplitude basis.

- Young Tensor Method (Li, Ren, Shu, MLX, Yu, Zheng 20').

$$\mathcal{M} \xrightarrow{\text{reduce}} \sum_i c_i \mathcal{M}_i^{(y)}, \quad \mathcal{M}^{(y)} \sim \langle ij\rangle \dots [kl] \dots$$

- Mathematica Package ABC4EFT (Li, Ren, MLX, Yu, Zheng 22')

Soft Block Construction: Second Attempt

Young Tensor approach (Low, Shu, MLX, Zheng 22'):

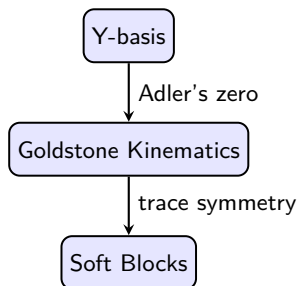
- ✓ Both parities are included (but mixed).

$$\begin{aligned} & \langle ij \rangle [jk] \langle kl \rangle [li] \pm [ij] \langle jk \rangle [kl] \langle li \rangle \\ &= \begin{cases} 4(s_{ij}s_{kl} - s_{ik}s_{jl} + s_{il}s_{jk}) \\ 4i\epsilon_{\mu\nu\rho\sigma} p_i^\mu p_j^\nu p_k^\rho p_l^\sigma \end{cases} \end{aligned}$$

- ✓ Free from Gram Det. and Schouten ($D = 4$)

	$\phi^6 p^8$	$\phi^6 p^{10}$	$\phi^6 p^{12}$	$\phi^7 p^8$	$\phi^7 p^{10}$	$\phi^7 p^{12}$	$\phi^8 p^{10}$	$\phi^8 p^{12}$
$f^{\text{even}}(s_{ij})$	–	1287	3003	–	8568	27132	42504	177100
y-basis	–	1286	2994	–	8547	26873	42308	173915
$f^{\text{odd}}(s_{ij}, \epsilon)$	225	825	2475	1575	8400	35700	53900	309925
y-basis	180	600	1650	1106	5019	18305	28196	132335

- The resulting soft blocks are still polynomials.
- Assuming independent traces (large N_f).

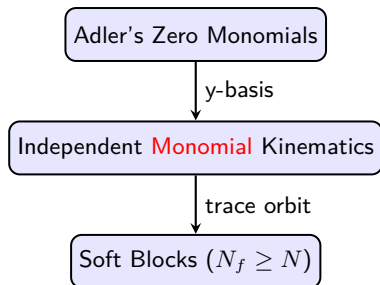


Soft Block Construction: Final Algorithm

Young Tensor approach (Low, Shu, MLX, Zheng 22'):

- ✓ Both parities are included separately.

$$s_{12834856} = \sum_{i=1}^{205} c_i^{(\text{even})} \mathcal{M}_i^y, \quad \epsilon_{1234856} = \sum_{i=1}^{205} c_i^{(\text{odd})} \mathcal{M}_i^y$$



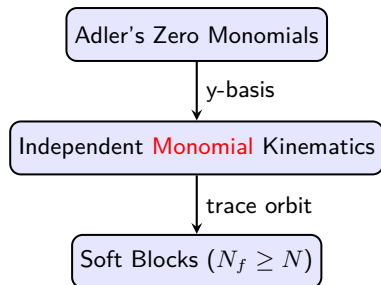
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- ✓ Free from Gram Det. and Schouten ($D = 4$)



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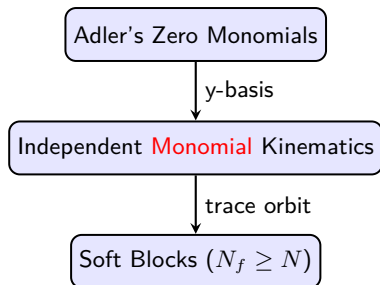
$$s_{12}s_{34}s_{56} = \sum_{i=1}^{205} c_i^{(\text{even})} \mathcal{M}_i^y, \quad \epsilon_{1234}s_{56} = \sum_{i=1}^{205} c_i^{(\text{odd})} \mathcal{M}_i^y$$

- ✓ Free from Gram Det. and Schouten ($D = 4$)

- Orbits of residual group $H^{(\mathcal{T})}$ (e.g. $H^{(6)} = Z_6$):

$$\mathcal{Y}(\mathcal{T}\mathcal{M}) = \mathcal{Y}(\mathcal{T}\mathcal{M}'), \quad \text{if } \mathcal{M}' \in H^{(\mathcal{T})}(\mathcal{M})$$

$$\mathcal{Y}(\mathcal{B}) = \mathcal{Y}(\mathcal{B}') \Leftrightarrow \mathcal{O} \simeq \mathcal{O}'$$



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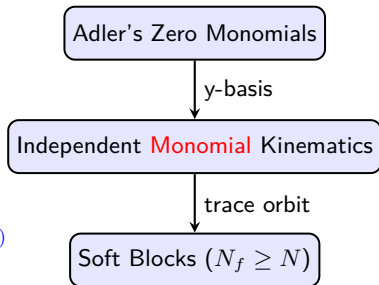
- Orbits of residual group $H^{(\mathcal{T})}$ (e.g. $H^{(6)} = Z_6$):

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$$\mathcal{Y}(\mathcal{B}) = \mathcal{Y}(\mathcal{B}') \Leftrightarrow \mathcal{O} \simeq \mathcal{O}'$$

$$\text{e.g. } \mathcal{Y}(\text{tr}[123456]s_{12}s_{34}s_{56}) = \mathcal{Y}(\text{tr}[123456]s_{23}s_{45}s_{61})$$

$$\Rightarrow \langle u_\mu u^\mu u_\nu u^\nu u_\rho u^\rho \rangle = \langle u_\rho u^\mu u_\mu u^\nu u_\nu u^\rho \rangle$$



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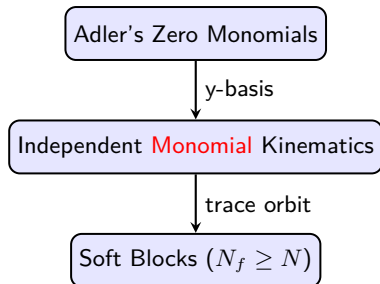
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- ✓ Free from Gram Det. and Schouten ($D = 4$)

- ✓ Linearly independent orbits $\{H^{(\mathcal{T})}(\mathcal{M}_a)\} \Leftrightarrow$
Monomial soft blocks $\mathcal{O}_a^{(\mathcal{T})} \simeq \mathcal{Y}(\mathcal{T}\mathcal{M}_a)$.

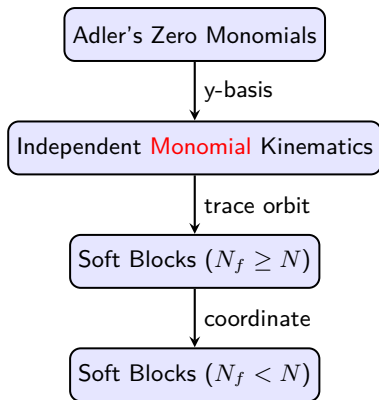
- They are independent at $N_f \geq N$.



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- ✓ Free from Gram Det. and Schouten ($D = 4$)
- ✓ Linearly independent orbits $\{H^{(\mathcal{T})}(\mathcal{M}_a)\} \Leftrightarrow$
Monomial soft blocks $\mathcal{O}_a^{(\mathcal{T})} \simeq \mathcal{Y}(\mathcal{T}\mathcal{M}_a)$.
- ✓ Reduce: $\mathcal{Y}(\mathcal{T}\mathcal{M}_a) = \sum_{i,j} c_{a,ij} \mathcal{T}_i \mathcal{M}_j^y$
 Find independent coordinates $c_{a,ij}$!



Independent Trace Structures

Group $SU(N_f)$	$SU(2)$	$SU(3)$	$SU(4)$	$SU(5)$	$SU(6)$	$SU(7)$	Trace
$\mathcal{T}^{a_1 a_2 a_3}$	1	2	2	2	2	2	2
$\mathcal{T}^{a_1 a_2 a_3 a_4}$	3	8	9	9	9	9	9
$\mathcal{T}^{a_1 a_2 a_3 a_4 a_5}$	6	32	43	44	44	44	44
$\mathcal{T}^{a_1 a_2 a_3 a_4 a_5 a_6}$	15	145	245	264	265	265	265

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$\mathcal{T}^{a_1 a_2 a_3 a_4}$	3	8	9	9	9	9	9
$\mathcal{T}^{a_1 a_2 a_3 a_4 a_5}$	6	32	43	44	44	44	44
$\mathcal{T}^{a_1 a_2 a_3 a_4 a_5 a_6}$	15	145	245	264	265	265	265

Inner Product Treatment

Define $g_{ij} \equiv \langle \mathcal{T}_i, \mathcal{T}_j \rangle = (\mathcal{T}_i^\dagger)_{a_1 \dots a_N} (\mathcal{T}_j)^{a_1 \dots a_N}$, we have

$$\det g(\mathbf{N}_f) \neq 0 \quad \Leftrightarrow \quad \{\mathcal{T}_i\} \text{ is independent}$$

$$\text{where} \quad T_{ij}^a T_{kl}^a = \delta_{il} \delta_{kj} - \delta_{ij} \delta_{kl} / N_f$$

Tensor reduction for $\mathcal{T} = \sum_i c^i \mathcal{T}_i$: $c^i = \langle (g^{-1})^{ij} \mathcal{T}_j, \mathcal{T} \rangle$.

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The Operator Basis

- ① Obtain independent orbits $\mathcal{Y}(\mathcal{TM}_a)$ (independent soft blocks for large N_f);
- ② Evaluate \mathcal{Y} and find the **unique coordinate on the $\mathcal{T}_i\mathcal{M}_j$** ;
- ③ Select independent monomials via linear algebra;
- ④ Translate to operators

$$p_{i\mu_1} p_{i\mu_2} \cdots p_{i\mu_n} \quad \Leftrightarrow \quad \nabla_{(\mu_1} \nabla_{\mu_2} \cdots \nabla_{\mu_{n-1}} u_{\mu_n})$$

$O(p^6)$ and $O(p^8)$ bases for 6-point soft blocks (agree with [Graf et al. 20'](#)):

6-pt $O(p^6)$	$SU(2)$	$SU(3)$	$SU(4)$	$SU(5)$	$SU(6)$
P-even	3	8	13	14	15
P-odd	0	3	4	4	4

6-pt $O(p^8)$	$SU(2)$	$SU(3)$	$SU(4)$	$SU(5)$	$SU(6)$
P-even	9	40	68	74	76
P-odd	2	20	33	35	35

The Operator Basis

$O(p^6)$ operators without external sources:

$SU(N_f)$	Operator Basis	Amplitude Basis
$SU(2)$	$\mathcal{O}_1 = \langle u^\mu u^\nu u^\rho u_\mu u_\nu u_\rho \rangle$	$\mathcal{B}_1 = \mathcal{Y} \circ \text{tr}[123456] s_{14} s_{25} s_{36}$
	$\mathcal{O}_2 = \langle u^\mu u^\nu u^\rho u_\mu u_\rho u_\nu \rangle$	$\mathcal{B}_2 = \mathcal{Y} \circ \text{tr}[123456] s_{14} s_{26} s_{35}$
	$\mathcal{O}_3 = \langle u^\mu u^\nu u^\rho u_\rho u_\mu u_\nu \rangle$	$\mathcal{B}_3 = \mathcal{Y} \circ \text{tr}[123456] s_{15} s_{26} s_{34}$
$SU(3)$	$\mathcal{O}_4 = \langle u^\mu u^\nu u^\rho u_\rho u_\nu u_\mu \rangle$	$\mathcal{B}_4 = \mathcal{Y} \circ \text{tr}[123456] s_{16} s_{25} s_{34}$
	$\mathcal{O}_5 = \langle u^\mu u^\nu u_\nu u^\rho u_\rho u_\mu \rangle$	$\mathcal{B}_5 = \mathcal{Y} \circ \text{tr}[123456] s_{16} s_{23} s_{45}$
	$\mathcal{O}_6 = \langle u^\mu u^\nu u^\rho u_\mu \rangle \langle u_\nu u_\rho \rangle$	$\mathcal{B}_6 = \mathcal{Y} \circ \text{tr}[1234 56] s_{14} s_{25} s_{36}$
	$\mathcal{O}_7 = \langle u^\mu u^\nu u^\rho u_\nu \rangle \langle u_\mu u_\rho \rangle$	$\mathcal{B}_7 = \mathcal{Y} \circ \text{tr}[1234 56] s_{15} s_{24} s_{36}$
	$\mathcal{O}_8 = \langle u^\mu u^\nu u_\mu u^\nu \rangle \langle u_\rho u_\rho \rangle$	$\mathcal{B}_8 = \mathcal{Y} \circ \text{tr}[1234 56] s_{13} s_{24} s_{56}$
$SU(4)$	$\mathcal{O}_9 = \langle u^\mu u^\nu u_\nu u_\mu \rangle \langle u^\rho u_\rho \rangle$	$\mathcal{B}_9 = \mathcal{Y} \circ \text{tr}[1234 56] s_{14} s_{23} s_{56}$
	$\mathcal{O}_{10} = \langle u^\mu u^\nu u^\rho \rangle \langle u_\mu u_\nu u_\rho \rangle$	$\mathcal{B}_{10} = \mathcal{Y} \circ \text{tr}[123 456] s_{14} s_{25} s_{36}$
	$\mathcal{O}_{11} = \langle u^\mu u^\nu u^\rho \rangle \langle u_\mu u_\rho u_\nu \rangle$	$\mathcal{B}_{11} = \mathcal{Y} \circ \text{tr}[123 456] s_{14} s_{26} s_{35}$
	$\mathcal{O}_{12} = \langle u^\mu u^\nu u_\mu \rangle \langle u^\rho u_\nu u_\rho \rangle$	$\mathcal{B}_{12} = \mathcal{Y} \circ \text{tr}[123 456] s_{13} s_{25} s_{46}$
	$\mathcal{O}_{13} = \langle u^\mu u^\rho \rangle \langle u^\nu u_\mu \rangle \langle u_\rho u_\nu \rangle$	$\mathcal{B}_{13} = \mathcal{Y} \circ \text{tr}[12 34 56] s_{14} s_{25} s_{36}$
$SU(5)$	$\mathcal{O}_{14} = \langle u^\mu u^\nu \rangle \langle u^\rho u_\rho \rangle \langle u_\mu u_\nu \rangle$	$\mathcal{B}_{14} = \mathcal{Y} \circ \text{tr}[12 34 56] s_{15} s_{26} s_{34}$
$SU(N_f \geq 6)$	$\mathcal{O}_{15} = \langle u^\mu u_\mu \rangle \langle u^\nu u_\nu \rangle \langle u^\rho u_\rho \rangle$	$\mathcal{B}_{15} = \mathcal{Y} \circ \text{tr}[12 34 56] s_{12} s_{34} s_{56}$

Induced Linear Relations

For example $N_f = 2$, choose $\{\mathcal{B}_1, \mathcal{B}_2, \mathcal{B}_3\}$ as the basis:

$$\begin{pmatrix} \mathcal{B}_4 \\ \mathcal{B}_5 \\ \mathcal{B}_6 \\ \mathcal{B}_7 \\ \mathcal{B}_8 \\ \mathcal{B}_9 \end{pmatrix} = \begin{pmatrix} 1 & 1 & -1 \\ 1 & 1 & -1 \\ 1 & 1 & 0 \\ 0 & 1 & 1 \\ 0 & 0 & 2 \\ 2 & 2 & -2 \end{pmatrix} \times \begin{pmatrix} \mathcal{B}_1 \\ \mathcal{B}_2 \\ \mathcal{B}_3 \end{pmatrix}, \quad \begin{pmatrix} \mathcal{B}_{10} \\ \mathcal{B}_{11} \\ \mathcal{B}_{12} \\ \mathcal{B}_{13} \\ \mathcal{B}_{14} \\ \mathcal{B}_{15} \end{pmatrix} = \begin{pmatrix} 0 & -1 & 1 \\ 0 & 1 & -1 \\ 0 & 0 & 0 \\ 1 & 2 & 1 \\ 2 & 2 & 0 \\ 4 & 4 & -4 \end{pmatrix} \times \begin{pmatrix} \mathcal{B}_1 \\ \mathcal{B}_2 \\ \mathcal{B}_3 \end{pmatrix}$$

Conclusion

- Features of our algorithm:
 - ① Monomial operators basis (consistent with Hilbert Series counting);
 - ② Parity-odd amplitude/operator basis;
 - ③ Gram determinant considered;
 - ④ Cayley-Hamilton theorem considered (finite N_f);
 - ⑤ Capable of deriving linear relations among the operators.
- Easy to add external sources: EWChEFT (NLO, NNLO) (Sun, MLX, Yu, 22' & 23').

Classes	$\mathcal{N}_{\text{type}}$	$\mathcal{N}_{\text{term}}$	$\mathcal{N}_{\text{operator}}$
UhD^4	3 + 6 + 0 + 0	15	15
X^2Uh	6 + 4 + 0 + 0	10	10
$XUhD^2$	2 + 6 + 0 + 0	8	8
X^3	4 + 2 + 0 + 0	6	6
ψ^2UhD	4 + 8 + 0 + 0	13(16)	$13n_f^2$ ($16n_f^2$)
ψ^2UhD^2	6 + 10 + 0 + 0	60(80)	$60n_f^2$ ($80n_f^2$)
ψ^2UhX	7 + 7 + 0 + 0	22(28)	$22n_f^2$ ($28n_f^2$)
ψ^4	12 + 24 + 4 + 8	117(160)	$\frac{1}{4}n_f^2(31 - 6n_f + 335n_f^2)$ ($n_f^2(9 - 2n_f + 125n_f^2)$)
Total	123	261(313)	$\frac{335n_f^4}{4} - \frac{3n_f^3}{2} + \frac{411n_f^2}{4} + 39$ ($39 + 133n_f^2 - 2n_f^2 - 2n_f^3 + 125n_f^4$) $\mathcal{N}_{\text{operators}}(n_f = 1) = 224(295)$, $\mathcal{N}_{\text{operators}}(n_f = 3) = 7704(11307)$

Thank you for your attention!

Gram Det Makes a Difference

ONLY at $O(p^{10})$ or beyond (# independent trace orbits):

Trace Class	(6)	(4 2)	(3 3)	(2 2 2)
General D	112	91	43	25
$D = 4$	111	90	42	24

	(8)	(6 2)	(5 3)	(4 4)	(4 2 2)	(3 3 2)	(2 2 2 2)
General D	435	320	226	129	149	117	26
$D = 4$	427	314	222	126	146	115	25

	(10)	(8 2)	(7 3)	(6 4)	(5 5)	(6 2 2)
General D	105	74	45	50	29	37
$D = 4$	99	71	43	47	27	35

	(5 3 2)	(4 4 2)	(4 3 3)	(4 2 2 2)	(3 3 2 2)	(2 2 2 2 2)
General D	35	30	21	18	18	7
$D = 4$	35	28	20	17	17	6