

# Momentum space oscillation properties in vortex states collision

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## 1 Background and Motivation

## 2 Main contents

- Bessel vortex collision
- General monochromatic vortex collision
- Laguerre-Gaussian vortex collision

## 3 Conclusions and Prospects

## Vortex states (scalar)

Coordinate space:  $\psi(r_{\perp}, \varphi, z) = f(r_{\perp}, z) e^{i\ell\varphi}$

Momentum space:  $\phi(k_{\perp}, \varphi_k, k_z) = g(k_{\perp}, k_z) e^{i\ell\varphi_k}$

### The states:

- Solutions for equation of motion (*i.e.* Schrodinger, Klein-Gordon);
- Angular momentum eigenstates  $\rightarrow L_z = \hbar\ell$ .

### Distribution $f(r_{\perp}, z)$ and $g(k_{\perp}, k_z)$ :

- Rotational-symmetric along propagation direction;
- Multiple choices  $\rightarrow$  Real (*i.e.* Bessel) or complex (*i.e.* Laguerre-Gaussian).

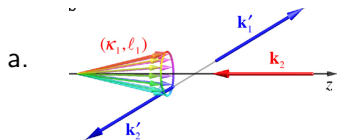
### ★ Topological charge $\ell$ :

- An integer quantum number  $\rightarrow$  infinite intrinsic freedoms;
- Proportional to angular momentum.

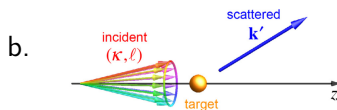
# Background: Collision/decay processes involving vortex particles

**Final particles:** Detected only by plane wave basis nowadays!

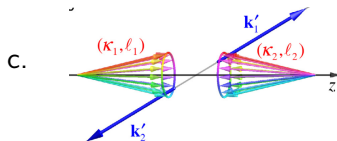
**Initial particles:** Several different schemes with different phenomena.



Collision of vortex particle and plane wave particle **or** decay of vortex particle: **no interference.**



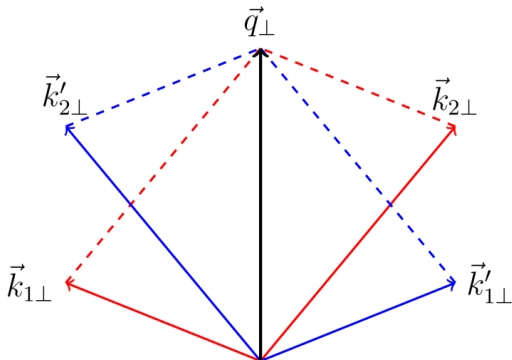
Collision of vortex particle with a target (size  $\neq 0$ ) **or** wave packet (*i.e.* Gaussian): **there is interference.**



Collision of two vortex particles: **there is interference.**

Pictures from [arXiv:1703.06879].

The **interference** (in momentum space) means:  
For the same final plane wave particle, there is **more than one combination of initial plane wave components** which satisfy momentum conservation law.



Example:  $\vec{k}_{1\perp}$  and  $\vec{k}'_{1\perp}$  come from plane wave components in initial state "1",  
 $\vec{k}_{2\perp}$  and  $\vec{k}'_{2\perp}$  come from plane wave components in initial state "2".

# Case a: Collision of vortex particle with plane wave particle or decay of vortex particle

Special kinetic phenomenons induced by special transverse momentum distribution (*i.e.* two-peak spectrum).

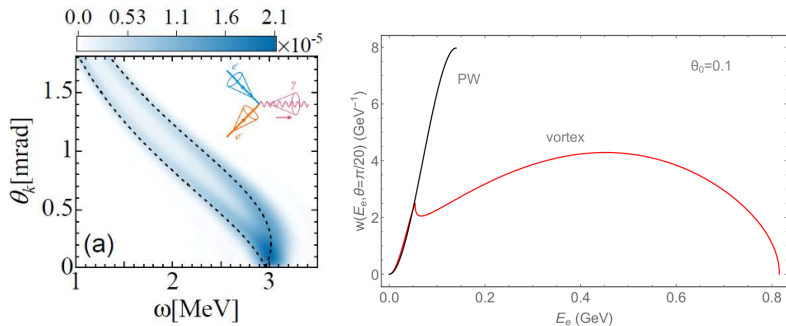


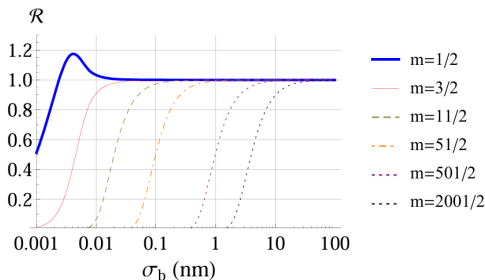
Figure 1: Two-peak spectrum of final particle: Radiation of charged vortex particle (Left: [arXiv:2302.05065]) & Decay of vortex particle (Right: [arXiv:2106.00345]).

No effect relating to topological charge.

## Case b: Collision of vortex particle with a target or Gaussian wave packet

### Effects which are dependent on topological charge

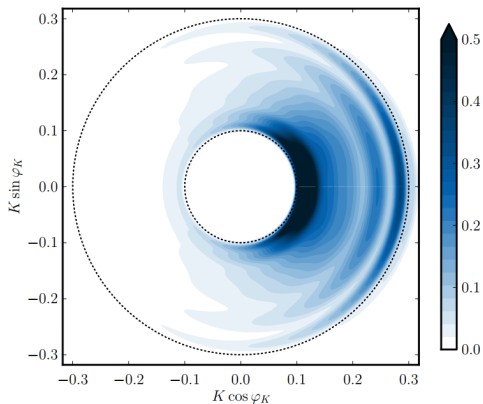
- Packet size effect [arXiv:2305.12419];
- Super-kick effect [arXiv:2109.01323];
- And so on.



Packet size effect.

## Case c: Collision of two vortex particles

**Oscillatory distribution in momentum space**, which is dependent on topological charge. [arXiv:1608.06551]



Oscillation of momentum distribution for final particle.



# Motivation and Research plan

## Motivation

Trying to find general observable quantity that:

- show common features which are dependent on topological charge;
- are not sensitive to transverse and longitudinal distributions.

## Research plan

- Choosing scheme of two vortex particle collision: Calculating  $S$  matrix for vortex particle collision (Scalar  $\phi^4$  model);
- Angular momentum eigenstates  $\rightarrow$  Rotational-symmetric total momentum distribution.
- Comparison: Bessel vortex states vs General monochromatic vortex states vs Laguerre-Gaussian vortex states.

# Bessel vortex collision

Bessel vortex state:

$$\psi_{E,\kappa,\ell}^B(\mathbf{x}, t) = \int \frac{d^3\mathbf{k}}{(2\pi)^3} a_{\kappa\ell}^B(\mathbf{k}) \exp(i\mathbf{k}\mathbf{x} - iEt),$$
$$a_{\kappa\ell}^B(\mathbf{k}) = N^B \delta(\vec{k}_\perp - \kappa) \delta(k_z - p_z) \exp(i\ell\varphi_k)$$

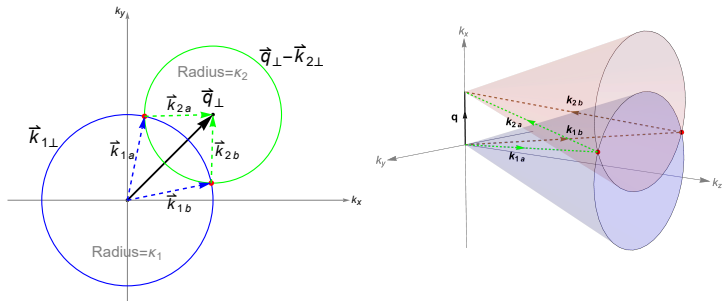
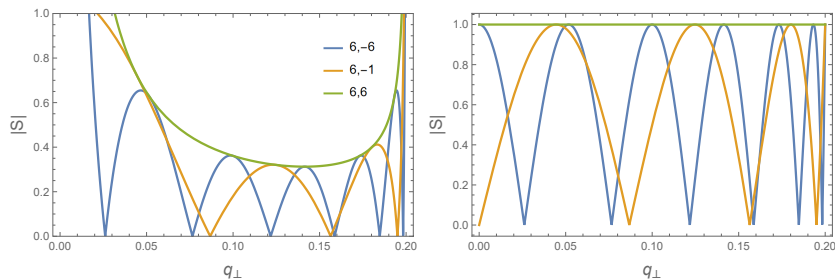


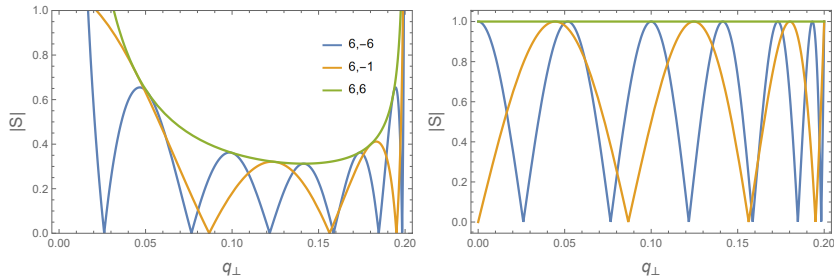
Figure 2: Interference space for Bessel vortex collision: two points



**Figure 3:** Oscillatory distribution for Bessel vortex collision with  $\Delta$ (left) and without  $\Delta$ (right).  $\ell = 6$  and  $\ell_2$  is chosen to be  $-6$ ,  $-1$  or  $6$ .

**Momentum value range:**  $|\kappa_1 - \kappa_2| < |\vec{q}_{\perp}| < |\kappa_1 + \kappa_2|$ .

$$\mathcal{S}(B_1 + B_2 \rightarrow P_3 + P_4) \propto \frac{\cos(\ell_1 \varphi_1 + \ell_2 \varphi_2)}{\Delta} \delta(k_{1z} + k_{2z} - q_z) \delta(E_1 + E_2 - E_q),$$



**Figure 3:** Oscillatory distribution for Bessel vortex collision with  $\Delta$ (left) and without  $\Delta$ (right).  $\ell = 6$  and  $\ell_2$  is chosen to be  $-6$ ,  $-1$  or  $6$ .

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$$S(B_1 + B_2 \rightarrow P_3 + P_4) \propto \frac{\cos(\ell_1 \varphi_1 + \ell_2 \varphi_2)}{\Delta} \delta(k_{1z} + k_{2z} - q_z) \delta(E_1 + E_2 - E_q),$$

## Oscillation rule

$$N^B = \text{Integer}[\ell_1 - \ell_2 / 2 + 1]$$

# General monochromatic vortex collision

General monochromatic vortex state:

$$\psi_{E,\ell}^M(\mathbf{x}, t) = \int \frac{d^3\mathbf{k}}{(2\pi)^3} a_\ell^M(\mathbf{k}) \exp(i\mathbf{k}\mathbf{x} - iEt),$$
$$a_\ell^M(\mathbf{k}) = N^M f(\theta) \delta(\sqrt{\mathbf{k}^2 + m^2} - E) \exp(i\ell\varphi_k)$$

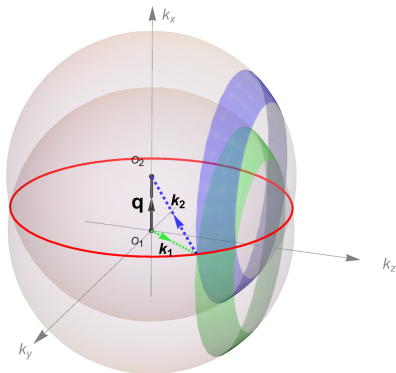


Figure 4: Interference space for general monochromatic vortex collision: a circle.

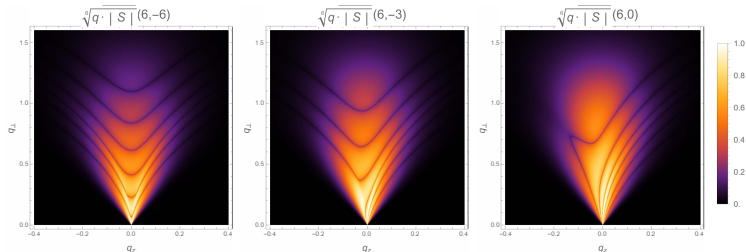


Figure 5: Oscillatory distribution for general monochromatic vortex collision with different  $(\ell_1, \ell_2)$ .

**Momentum value range:**  $||\mathbf{k}_1| - |\mathbf{k}_2|| < q_z < ||\mathbf{k}_1| + |\mathbf{k}_2||$ .

$$\mathcal{S}(M_1 + M_2 \rightarrow P_3 + P_4) \propto \int_0^\pi d\hat{\varphi}_1 \frac{k_1 k_2}{q} f_1(\theta_1) f_2(\theta_2) \cos(\ell_1 \varphi_1 + \ell_2 \varphi_2) \mathcal{M} \delta(E_1 + E_2 - E_q)$$

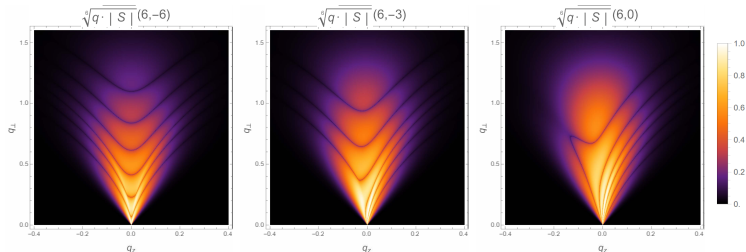


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**Momentum value range:**  $\|\mathbf{k}_1\| - \|\mathbf{k}_2\| < q_z < \|\mathbf{k}_1\| + \|\mathbf{k}_2\|$ .

$$S(M_1 + M_2 \rightarrow P_3 + P_4) \propto \int_0^\pi d\hat{\varphi}_1 \frac{k_1 k_2}{q} f_1(\theta_1) f_2(\theta_2) \cos(\ell_1 \varphi_1 + \ell_2 \varphi_2) \mathcal{M} \delta(E_1 + E_2 - E_q)$$

## Oscillation rule

$$N_{\perp}^M = (|\ell_1| + |\ell_2| - |\ell_1 + \ell_2|)/2 + 1, \quad N_z^M = |\ell_1 + \ell_2| + 1$$

## Limitation at $q_{\perp} = 0$

$$\mathcal{S}(M_1 + M_2 \rightarrow P_3 + P_4) \propto \int_0^{\pi} d\varphi_1 \cos[(\ell_1 + \ell_2)\varphi_1] \mathcal{M},$$

$$\lim_{q_{\perp} \rightarrow 0} |\mathcal{S}| = \begin{cases} \text{Nonzero} & \text{if } \ell_1 + \ell_2 \equiv 0 \text{ \& } \|\mathbf{k}_1\| - \|\mathbf{k}_2\| < q < \|\mathbf{k}_1\| + \|\mathbf{k}_2\|; \\ 0 & \text{if } \ell_1 + \ell_2 \neq 0. \end{cases}$$

## Limitation at $q = 0$

$$\lim_{q \rightarrow 0} |\mathcal{S}| = \begin{cases} \infty & \text{if } \ell_1 + \ell_2 \equiv 0; \\ 0 & \text{if } \ell_1 + \ell_2 \neq 0. \end{cases}$$



# Laguerre-Gaussian vortex collision

Laguerre-Gaussian(LG) vortex state:

$$\psi_{n,\ell}^L(\mathbf{x}, t) = \int \frac{d^3\mathbf{k}}{(2\pi)^3} a_{n\ell}^L(\mathbf{k}) \exp(i\mathbf{k}\mathbf{x} - iEt),$$

$$a_{n\ell}^L(\mathbf{k}) = N^L k_{\perp}^{|\ell|} L_n^{|\ell|} \left( \frac{k_{\perp}^2}{\sigma_{\perp}^2} \right) \exp\left(-\frac{k_{\perp}^2}{2\sigma_{\perp}^2} - \frac{(k_z - k_{z0})^2}{2\sigma_z^2} + i\ell\varphi_k\right)$$

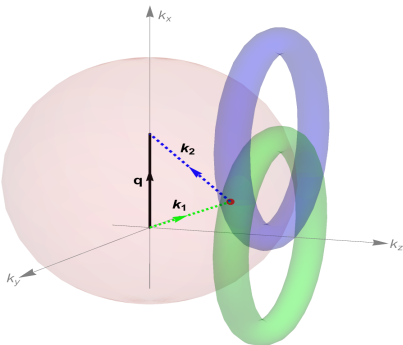


Figure 6: Interference space for LG vortex collision: a surface.

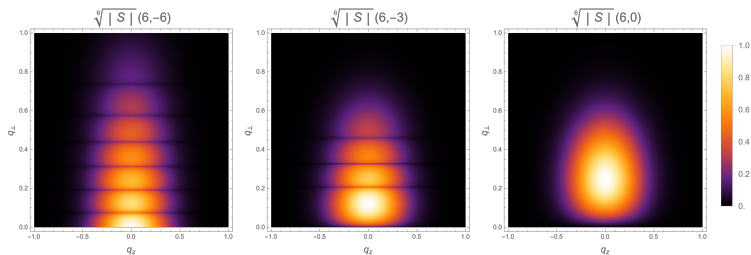


Figure 7: Oscillatory distribution for LG vortex collision

$$\begin{aligned}
 & S(L_1 + L_2 \rightarrow P_3 + P_4) \\
 & \propto \int_0^\pi d\hat{\varphi}_1 \int_0^\pi d\hat{\theta}_1 \frac{\sin \hat{\theta}_1 k_1^2 E_1 E_2 g_1(k_{1\perp}, k_{1z}) g_2(k_{2\perp}, k_{2z})}{|E_q k_1 - E_1 q \cos \hat{\theta}_1|} \cos(\ell_1 \varphi_1 + \ell_2 \varphi_2) \mathcal{M}
 \end{aligned}$$

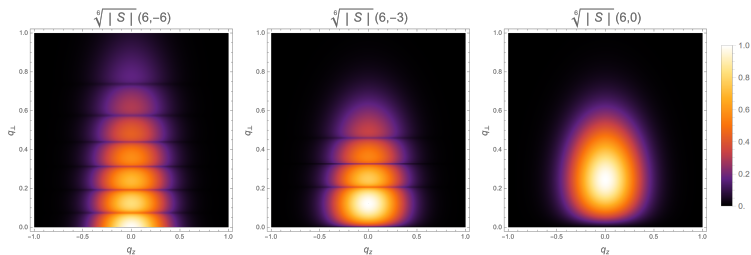


Figure 7: Oscillatory distribution for LG vortex collision

$$\begin{aligned}
 & S(L_1 + L_2 \rightarrow P_3 + P_4) \\
 & \propto \int_0^\pi d\hat{\varphi}_1 \int_0^\pi d\hat{\theta}_1 \frac{\sin \hat{\theta}_1 k_1^2 E_1 E_2 g_1(k_{1\perp}, k_{1z}) g_2(k_{2\perp}, k_{2z})}{|E_q k_1 - E_1 q \cos \hat{\theta}_1|} \cos(\ell_1 \varphi_1 + \ell_2 \varphi_2) \mathcal{M}
 \end{aligned}$$

oscillation rule

$$N_{\perp}^L = (|\ell_1| + |\ell_2| - |\ell_1 + \ell_2|)/2 + 1, \quad N_z^L = 1$$

# LG vortex collision: 1D oscillatory distributions

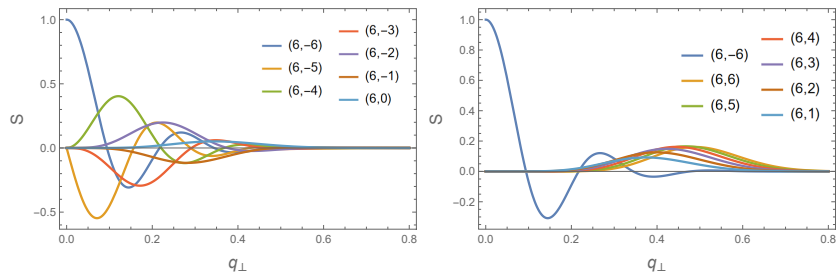
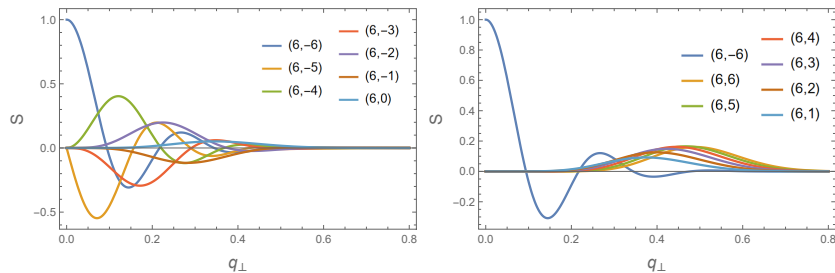


Figure 8: Distribution of total transverse momentum at fixed total longitudinal momentum ( $q_z = 0$ ).

# LG vortex collision: 1D oscillatory distributions



**Figure 8:** Distribution of total transverse momentum at fixed total longitudinal momentum ( $q_z = 0$ ).

## Uniqueness of zero total angular momentum case

- $l_1 + l_2 = 0 \rightarrow$  Maximum of  $|\mathcal{S}|$  appears on collision axis.
- $l_1 + l_2 \neq 0 \rightarrow \mathcal{S} = 0$  on  $q_{\perp} = 0$ .

# Conclusions and Prospects

## Conclusions

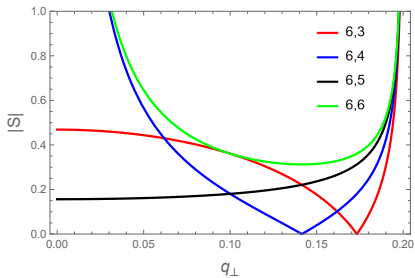
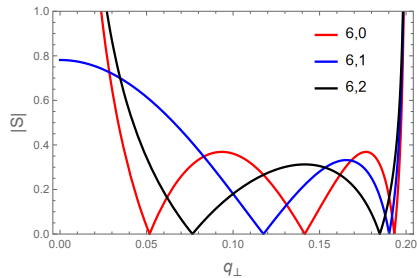
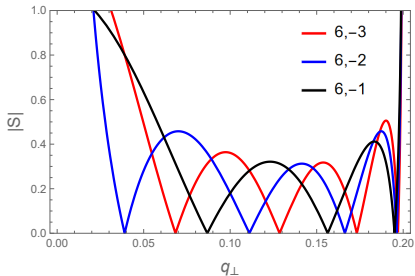
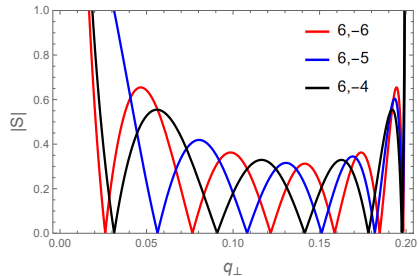
- Final total momentum distributions of vortex particle collision (in the scalar model) have **topological structures** that are determined by initial topological charges.
- The longitudinal **topological number** is different for different vortex states, while the transverse one show general dependence on topological charges ( $N_{\perp} = (|\ell_1| + |\ell_2| - |\ell_1 + \ell_2|)/2 + 1$ ).
- Distributions of two cases with  $\ell_1 + \ell_2 = 0$  or  $\ell_1 + \ell_2 \neq 0$  have significant difference when **total transverse momentum tends to zero**. This may also be a general feature that is dependent on topological charges.

## Prospects

Both features mentioned above can be used for detection of vortex states' topological charge.

Thanks!

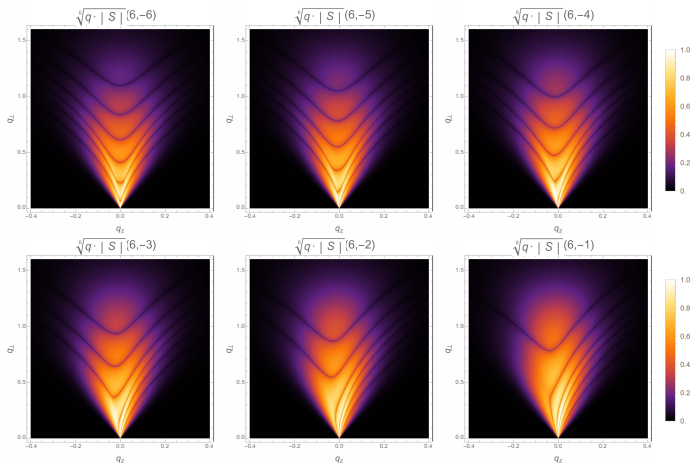
# More results of Bessel vortex collision



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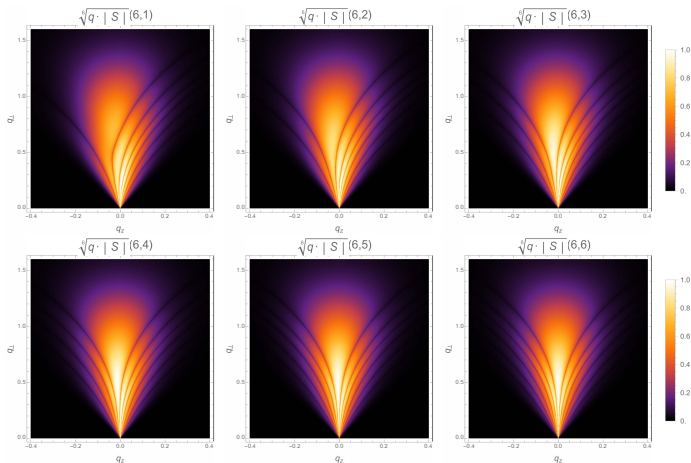


# More results of monochromatic vortex collision



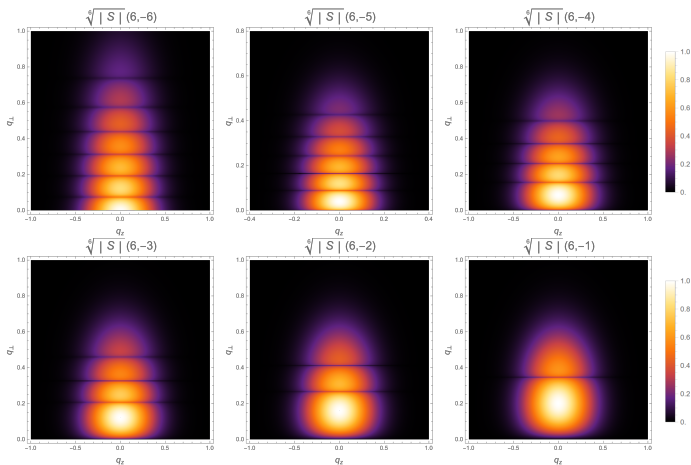
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# More results of monochromatic vortex collision



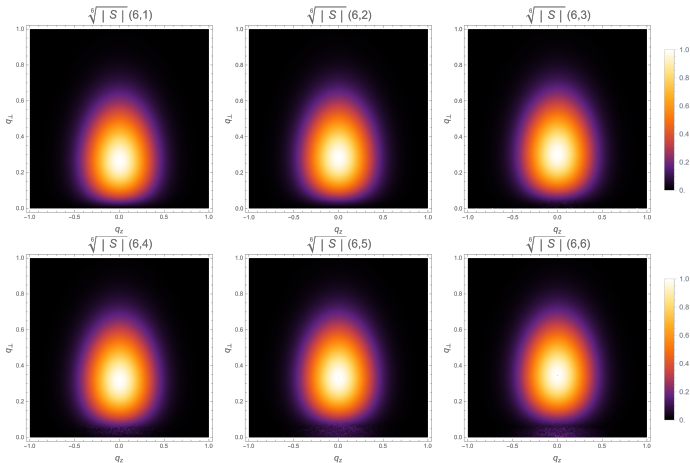
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# More results of LG vortex collision



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# More results of LG vortex collision



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