

Search for ultralight new particles with spin-based sensors

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Matter AntiMater Section

2024.07.14

贵阳



New Physics (Beyond the Standard Model)

Do spins couple to gravity?

How to include Gravity?

Dark Matter
and
Dark Energy

Why the universe
look like this

Strong CP problem

Why more matter
than antimatter

Why do we exist?

Fifth force?

2.2 M u $+2/3$ $1/2$ up	1.28 G c $+2/3$ $1/2$ charm	173.1 G t $+2/3$ $1/2$ top	Mass: eV/c^2 Charge Spin Name	125.09 G H 0 0 higgs
4.7 M d $-1/3$ $1/2$ down	96 M s $-1/3$ $1/2$ strange	4.18 G b $-1/3$ $1/2$ bottom	g 0 1 gluon	electromagnetic force
0.51 M e -1 $1/2$ electron	105.66 M μ -1 $1/2$ muon	1.78 G τ -1 $1/2$ tau	γ 0 1 photon	
< 2 ν_e 0 $1/2$ e neutrino	< 0.19 M ν_μ 0 $1/2$ μ neutrino	< 18.2 M ν_τ 0 $1/2$ τ neutrino	80.38 G W ± 1 1 W boson	weak nuclear force
			91.19 G Z 0 1 Z boson	
FERMIONS			BOSONS	

strong nuclear force

electromagnetic force

weak nuclear force



Axion: Peccei-Quinn-Weinberg-Wilczek (PQWW) Model

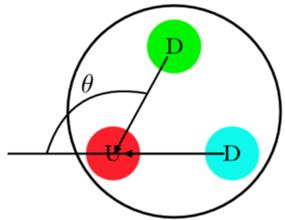
- Strong CP problem.

$$L_{CP_{vio}} = \frac{\alpha_s}{8\pi} \theta \text{tr}(G_{\mu\nu} \tilde{G}^{\mu\nu})$$

$$\bar{\theta} = \theta + \arg \det M_q$$

- Neutron EDM Prediction

$$d_n(\bar{\theta}) \approx \frac{e\bar{\theta}m_u m_d}{(m_u + m_d)m_n^2} \approx 6 \times 10^{-17} \bar{\theta} e \cdot \text{cm}$$



Experiment

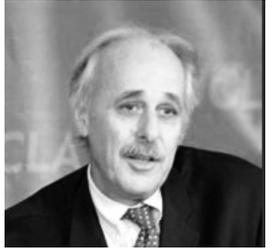
$$d_n = \left(0.0 \pm 1.1_{stat} \pm 0.2_{sys} \right) \times 10^{-26} e \cdot \text{cm}.$$

PRL 124, 081803 (2020)

$$|\bar{\theta}| < 10^{-9}$$

- Peccei-Quinn mechanism

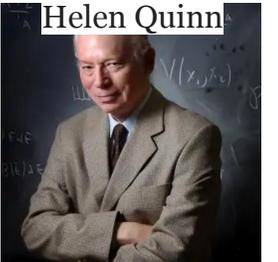
$a + \bar{\theta} f_a = \bar{a}$ The $\bar{\theta}$ term can be absorbed into the axion field and it can explain why $\bar{\theta}$ is so small.



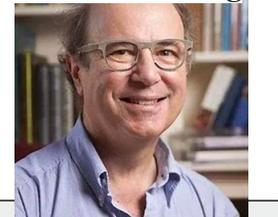
Roberto Peccei



Helen Quinn

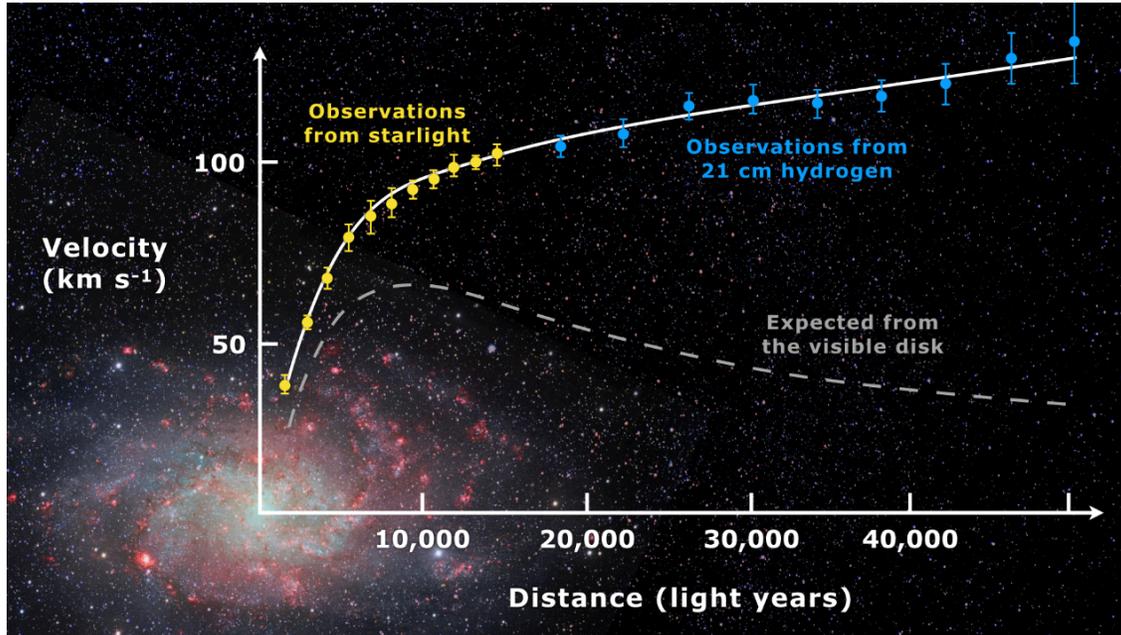


Steven Weinberg

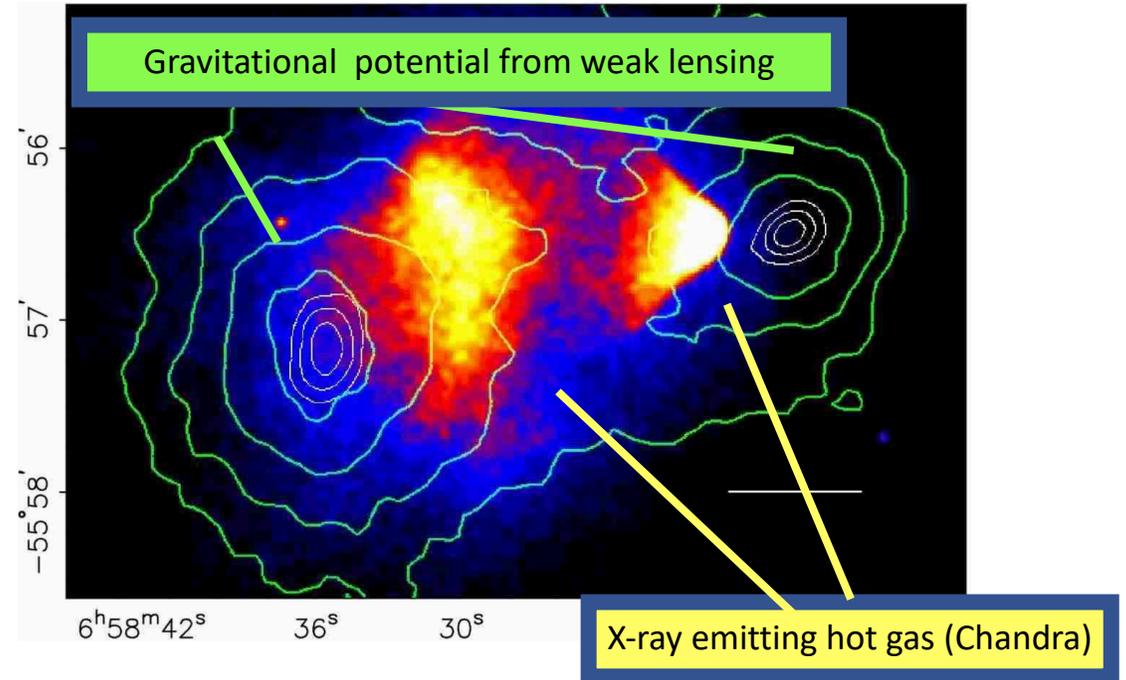


Frank Wilczek

Dark Matter



M33 Galaxy, Velocity v.s. Distance from the center



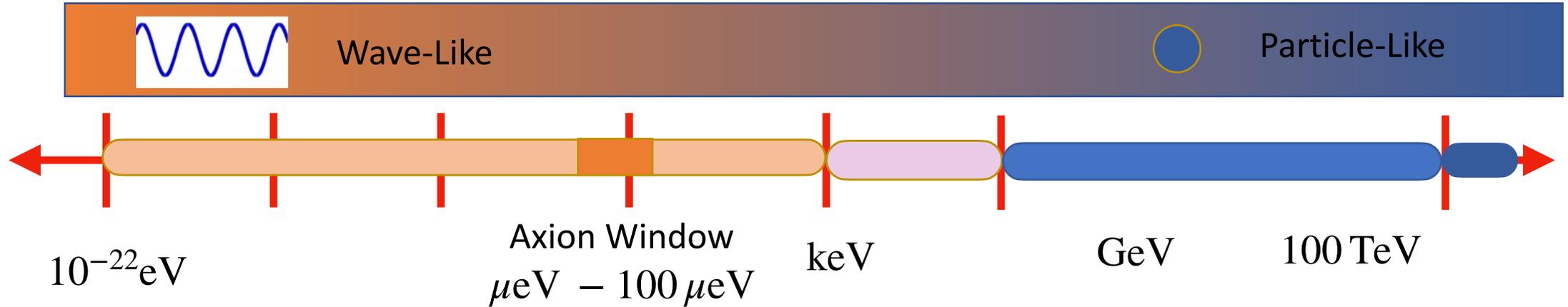
D. Clowe, et al. *Astrophys. J. Lett.* 648, L109 (2006)



Fritz Zwicky
Coma Cluster
1930s

5 times more than normal matter!

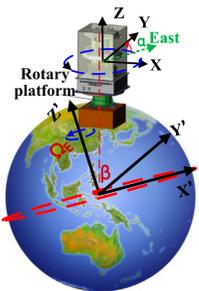
Dark Matter Mass Range



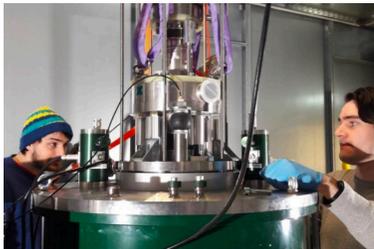
Ultralight Bosonic DM

Table-top Precision
Measurement experiments

ChangeE



CASPER



Cavity
Experiments

ADMX

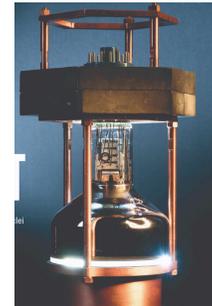


SRF



Light DM

Sterile ν

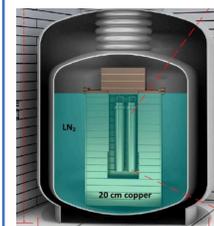


WIMPs

ALETHEIA



CDEX

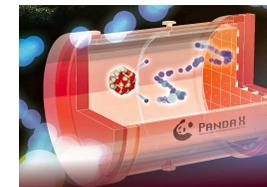


XENON, LUX, Accelerator

DOPMPE



PandaX



etc

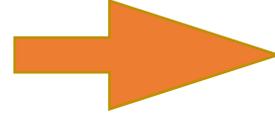
Composite
DMs &
Primordial
Black holes

James Webb
Telescope



Axion (and Axionlike Particles) Dark Matter Detection

- Ultralight
- Weakly interacting
- Huge number density

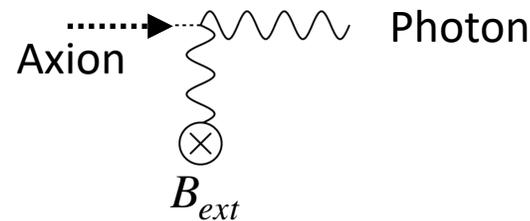


Perfect DM candidate

- Axion-Photon Coupling:

$$\mathcal{L}_{EM} \approx g_{a\gamma\gamma} a(\vec{r}, t) F_{\mu\nu} \tilde{F}^{\mu\nu}$$

Inverse Primakoff Effect

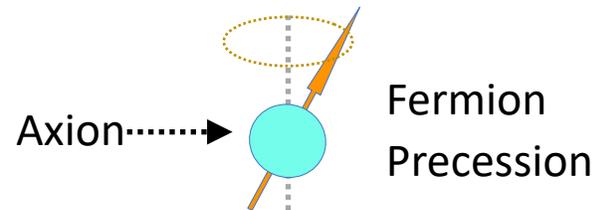


ADMX

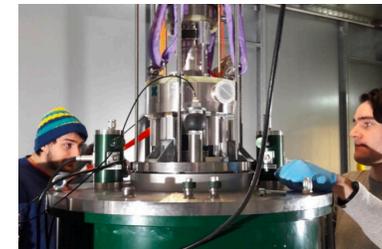


- Axion-Fermion Coupling:

$$\mathcal{L}_{spin} \approx g_{aNN} \left[\partial_{\mu} a(\vec{r}, t) \right] \bar{\Psi}_n \gamma^{\mu} \gamma_5 \Psi_n$$

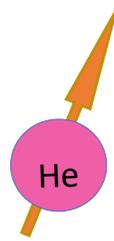
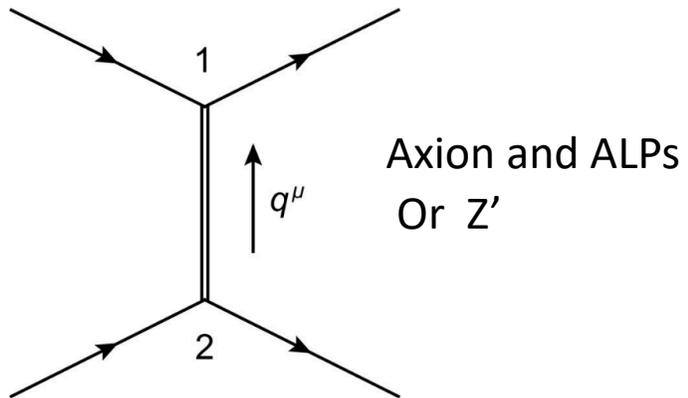


CASPEr

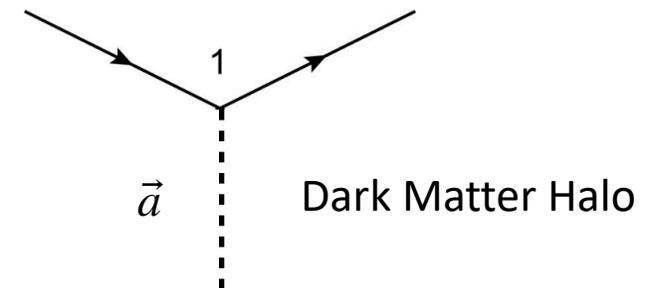


Exotic Spin-Dependent Force

New Interaction



Axion and ALPs



Spin-0

$$\mathcal{L}_\phi = \phi \sum_\psi \bar{\psi} \left(g_\psi^S + i\gamma_5 g_\psi^P \right) \psi,$$

Spin-1

$$\mathcal{L}_{Z'} = Z'_\mu \sum_\psi \bar{\psi} \gamma^\mu \left(g_\psi^V + \gamma_5 g_\psi^A \right) \psi,$$

Axion Gradient coupling

$$\mathcal{L}_a = (\partial_\mu a) \bar{\psi} \gamma^\mu \gamma_5 \psi,$$

Exotic Spin-Dependent Force

Scalar and Pseudoscalar Coupling (Axion)

$$V_{9+10} = \frac{g_s g_p \hbar^2}{8\pi m_1} (\hat{\sigma}_1 \cdot \hat{r}) \left(\frac{1}{r\lambda} + \frac{1}{r^2} \right) e^{-r/\lambda},$$

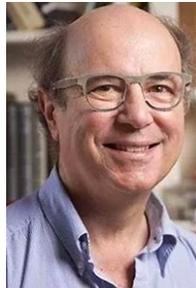
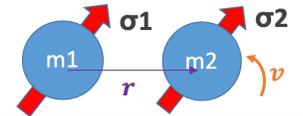
P Violation

CP violation

Vector and Axial Vector Coupling

$$V_{4+5} = \frac{(g_A^2 - 3g_V^2)\hbar^2}{16\pi m c} [(\hat{\sigma} \cdot (\mathbf{v} \times \hat{r}))] \left(\frac{1}{\lambda r} + \frac{1}{r^2} \right) e^{-r/\lambda},$$

$$V_8 = \frac{g_A g_A \hbar}{4\pi c} [(\hat{\sigma}_1 \cdot \mathbf{v})(\hat{\sigma}_2 \cdot \mathbf{v})] \left(\frac{1}{r} \right) e^{-\frac{r}{\lambda}}$$



New macroscopic forces?

Phys. Rev. D 30, 130

J. E. Moody* and Frank Wilczek

Institute for Theoretical Physics, University of California, Santa Barbara, California 93106

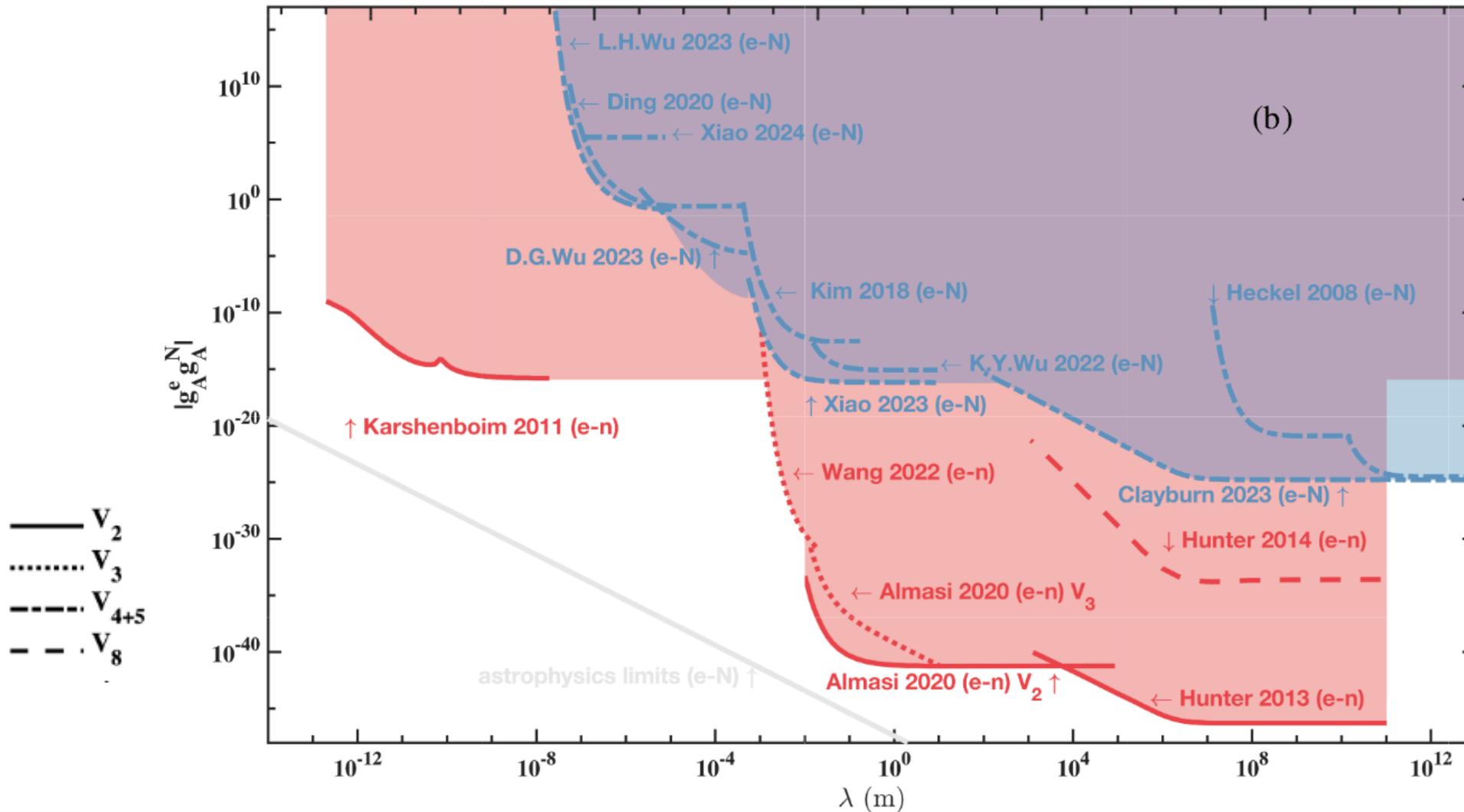
(Received 17 January 1984)



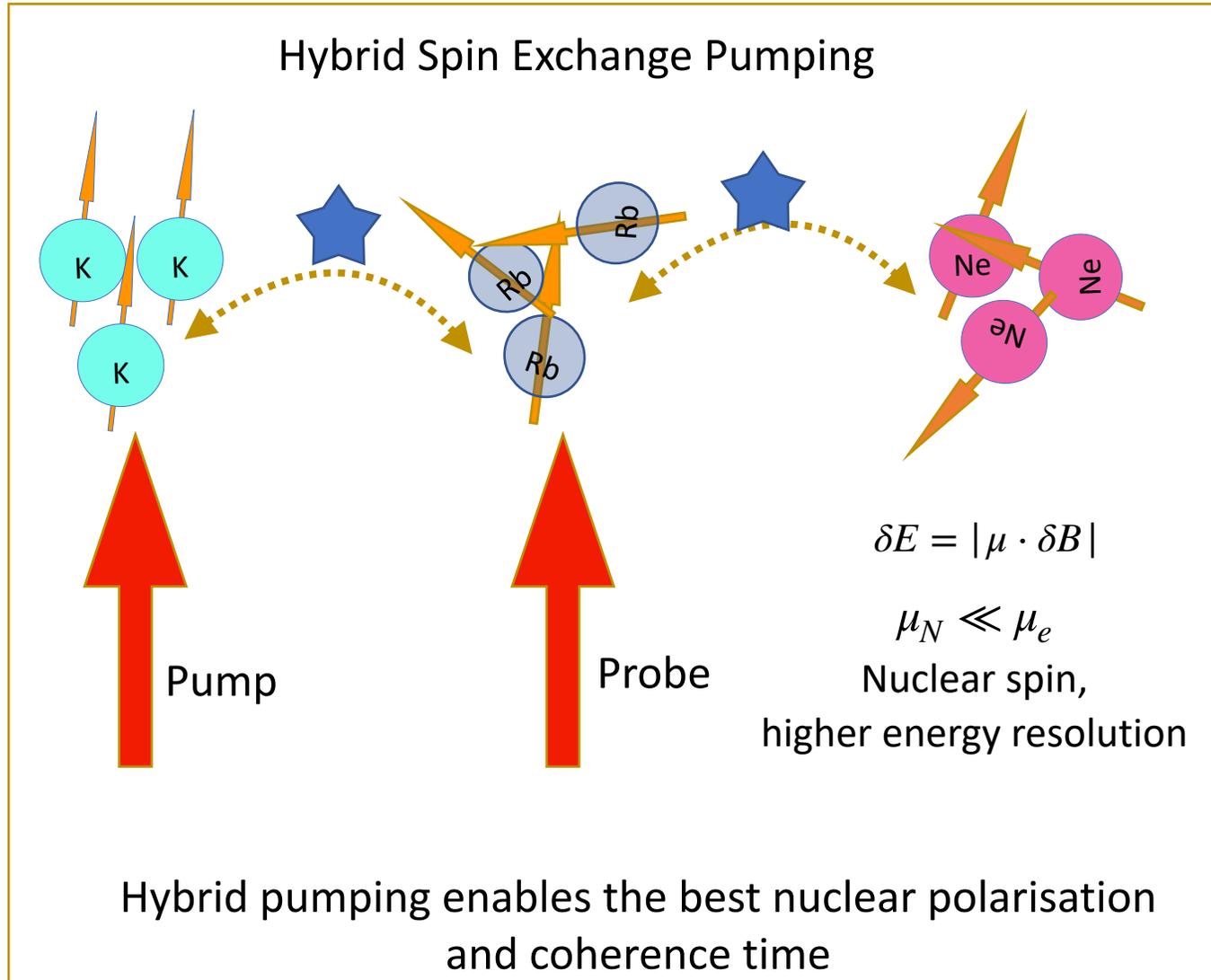
Spin-dependent macroscopic forces from new particle exchange

B. A. Dobrescu and I. Mocioiu, J. High Energy Phys. 11 (2006) 005.

On 4/17/23, 8:05 AM, "rmp@aps.org" <rmp@aps.org> wrote:



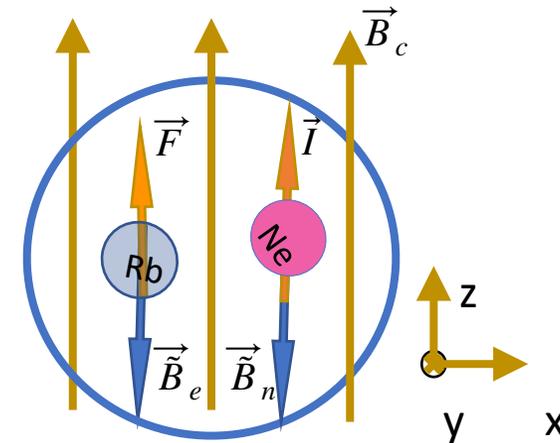
Alkali-Noble Comagnetometer: Sword for New Physics



Dynamics of the spins

$$\frac{\partial \mathbf{P}^e}{\partial t} = \frac{\gamma_e}{Q(P^e)} [\mathbf{b}^e + \mathbf{B} + \lambda M^n \mathbf{P}^n + \mathbf{L}] \times \mathbf{P}^e + \frac{P_{0z}^e - P^e}{T_e Q(P^e)}$$

$$\frac{\partial \mathbf{P}^n}{\partial t} = \gamma_n [\mathbf{b}^n + \mathbf{B} + \lambda M^e \mathbf{P}^e] \times \mathbf{P}^n + \frac{P_{0z}^n - P^n}{\{T_{2n}, T_{2n}, T_{1n}\}}$$

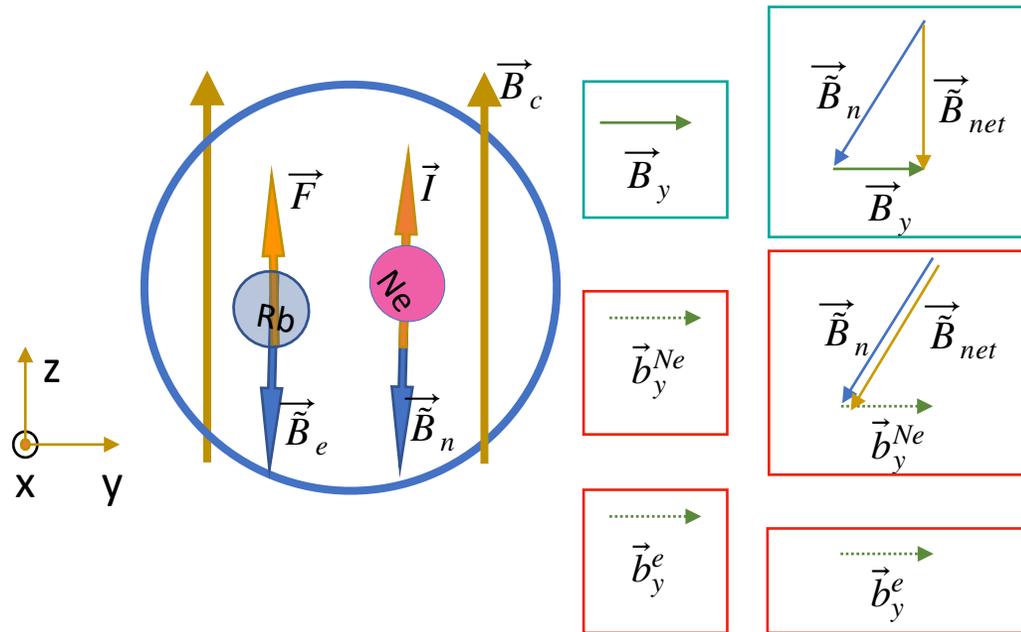


Self-Compensation (SC) Regime

Wei et al, Phys Rev Lett.130.063201 (2023)

$$S = K \frac{\gamma_e P_z^e}{R_{\text{tot}}^e} \left(\boxed{(B_y^{\text{Ne}} - B_y^e)} + \boxed{(b_y^{\text{Ne}} - b_y^e)} + \frac{\Omega_y}{\gamma_{\text{Ne}}} \right),$$

The SC can suppress normal magnetic field \mathbf{B} , rather than new physics field $\mathbf{b}^{e,n}$



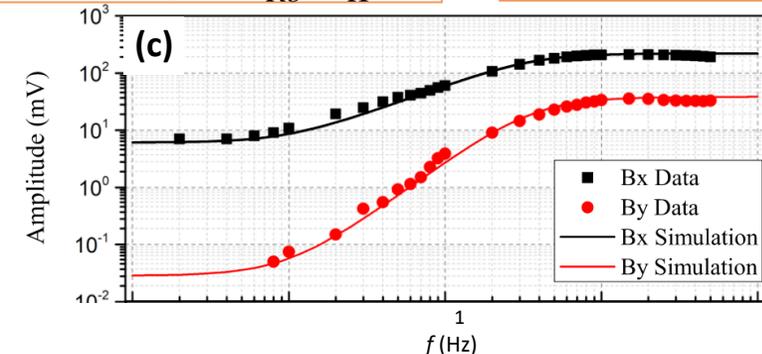
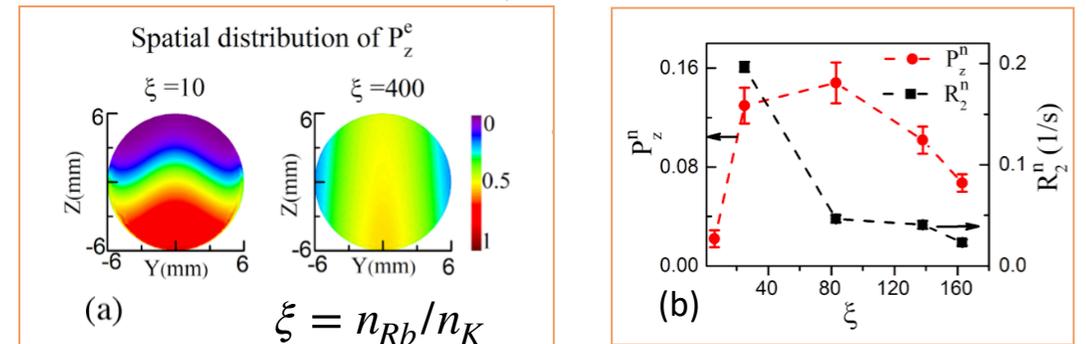
The y axis (sensitive axis) fields feel by Rb

$$SF_x = \frac{R_2^n + \omega/2 + \omega^2 \hat{\omega}_0^e / (R_2^e \hat{\omega}_0^n)}{\sqrt{(\hat{\omega}_0^n)^2 + (\hat{\omega}_0^e)^2 \omega^2 / R_2^e}},$$

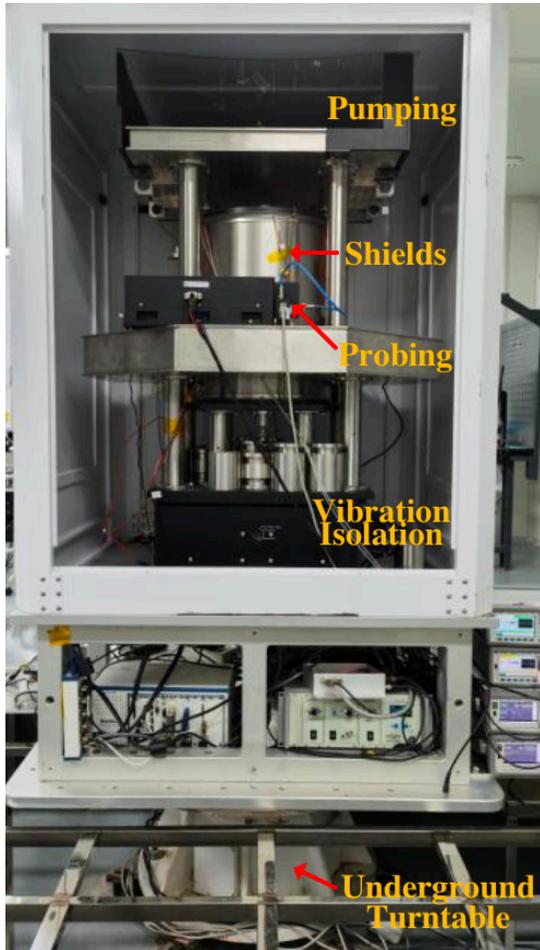
$$SF_y = \frac{(R_2^n)^2 + \omega^2}{\hat{\omega}_0^n \sqrt{(\hat{\omega}_0^n)^2 + (\hat{\omega}_0^e)^2 \omega^2 / R_2^e}},$$

Suppression Factor $SF_{x,y} \propto R_2 \propto \nabla B$ (smaller the better)

Reduce Alkali gradient ∇B_y^e to optimise suppression factor

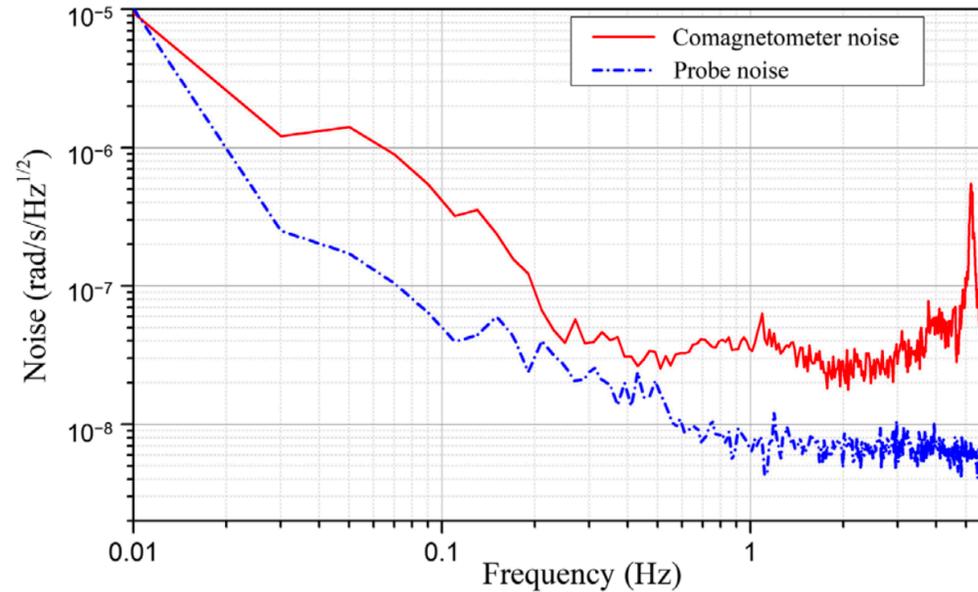


Ultrahigh sensitivity Self-Compensation Comagnetometry



K-Rb-²¹Ne SERF Comagnetometer

Wei, et al. Physical Review Letters, 130(6), p.063201



$$\delta\Omega = 3 \times 10^{-8} \text{ rad/s Hz}^{-1/2}$$

$$\delta b^n = 1.5 \text{ fT Hz}^{-1/2}$$

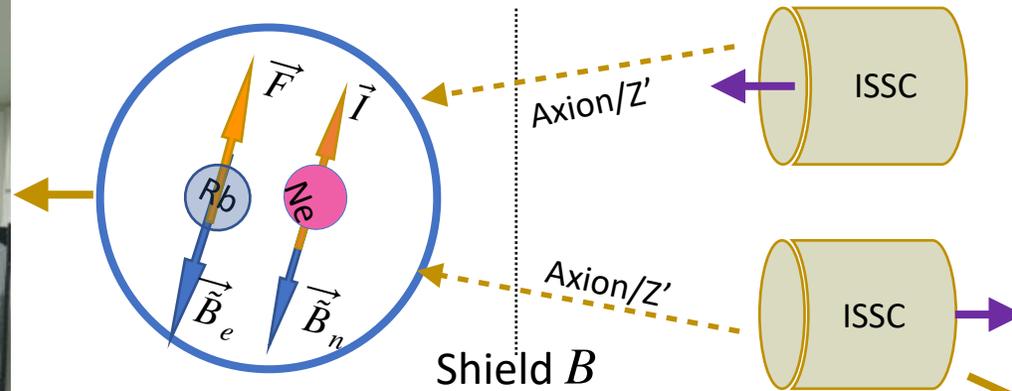
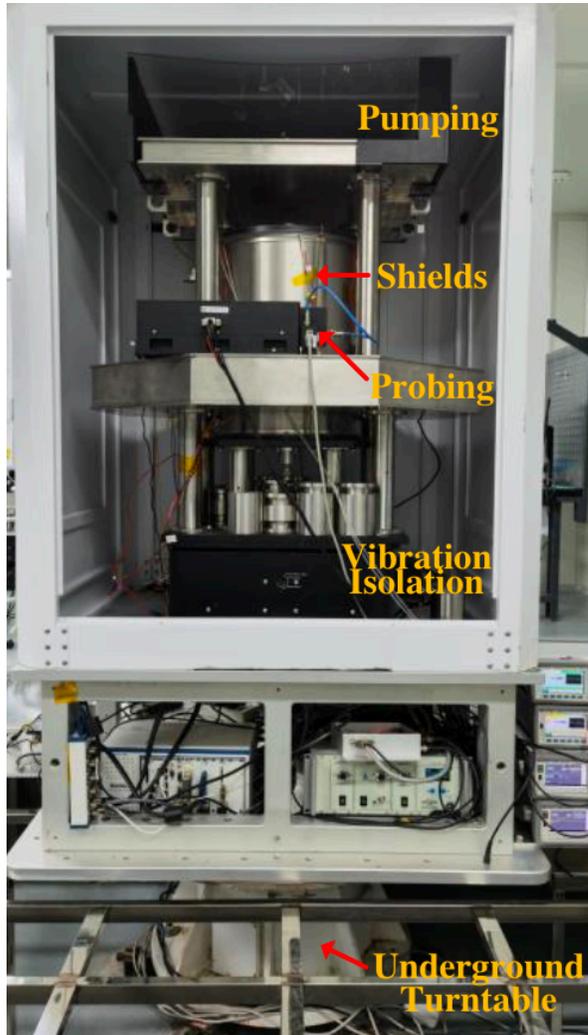
Highest Energy Resolution

$$\delta E_{\text{Ne}} = 3.1 \times 10^{-23} \text{ eV Hz}^{-1/2}$$

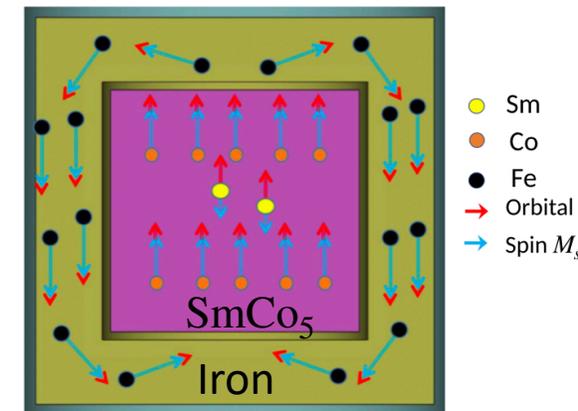
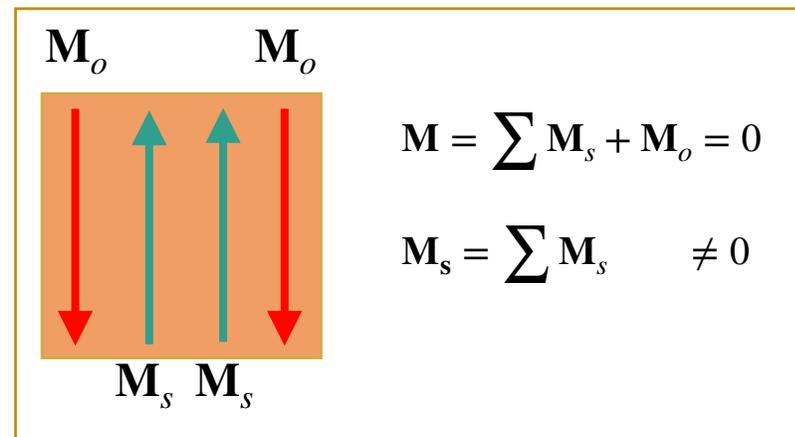
• Recipe

- Magnetic Noise: Ferrite Shielding
- Temperature gradient: Vacuum Isolation
- Vibration Isolation: Additional foundation
- Laser polarisation stability: Polarisation Modulation

Fifth Force Experiment

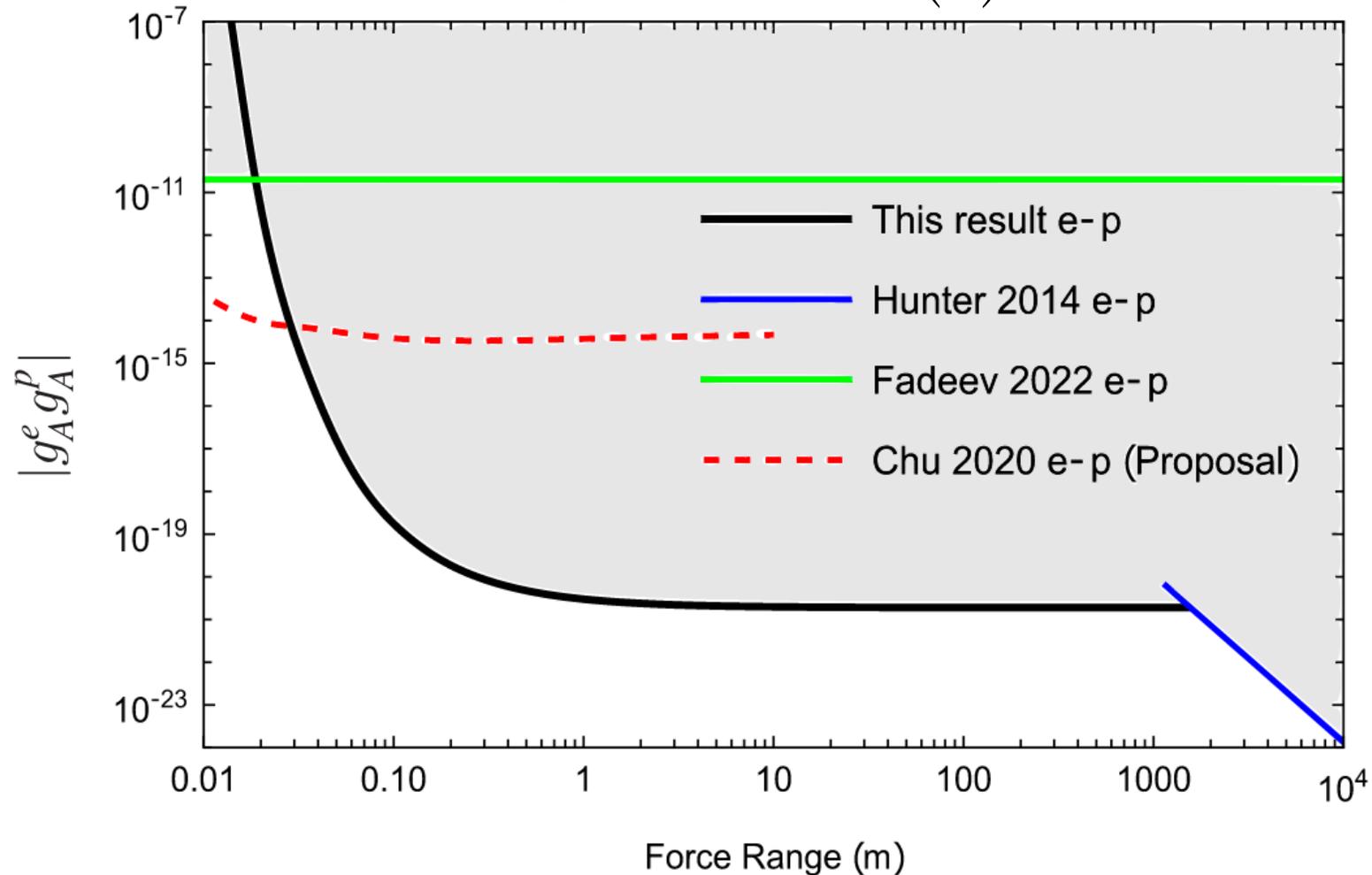


- Ji et al, **PRL** 130, 133202(2023).
- Wei et al, **PRL** 130, 063201 (2023).
- Ji et al, **PRL**, 121(26): 261803, 2018.
- Ji et al, **PRD**, 95(7): 075014, 2017.
- Wei et al, **NC** 13, 7387, (2022)

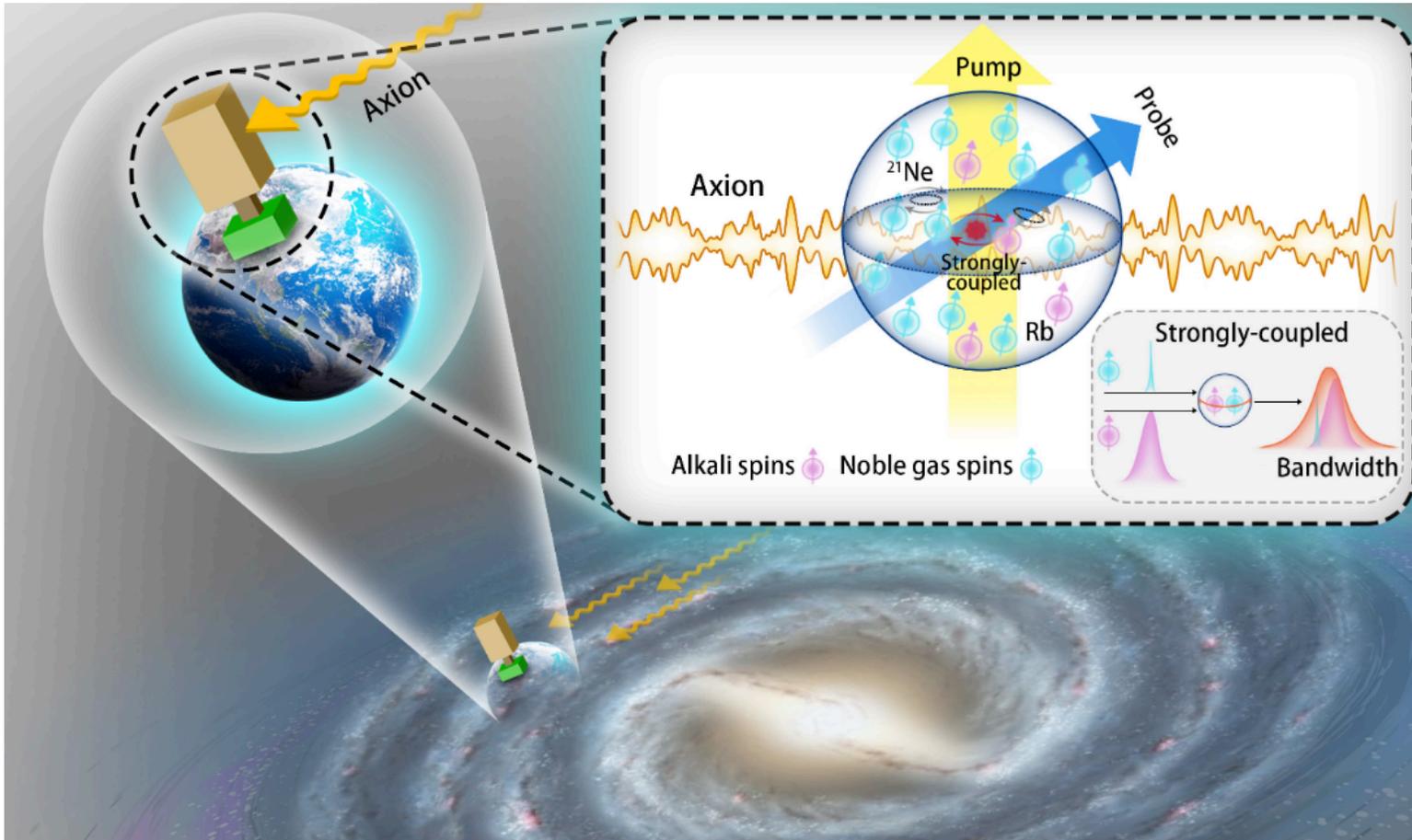


Fifth Force Experiment

$$V_8 = \frac{g_A g_A \hbar}{4\pi c} [(\hat{\sigma}_1 \cdot \mathbf{v})(\hat{\sigma}_2 \cdot \mathbf{v})] \left(\frac{1}{r}\right) e^{-\frac{r}{\lambda}}$$



Axion Dark Matter Wind



Axion wave

$$\nabla a(x) = \sum_{\mathbf{p}} \sqrt{\frac{2N_{\mathbf{p}}}{V\omega_{\mathbf{p}}}} \cos(\omega_{\mathbf{p}}t - \mathbf{p} \cdot \mathbf{x} + \phi_{\mathbf{p}}) \mathbf{p}$$

Axion-Nucleon spin coupling

$$\mathcal{H} = g_{aNN} \nabla a \cdot \boldsymbol{\sigma}_N,$$

Signal

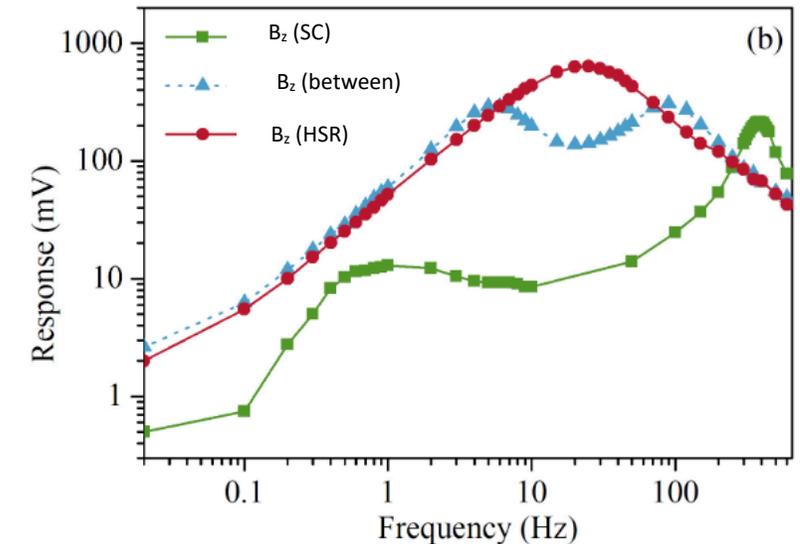
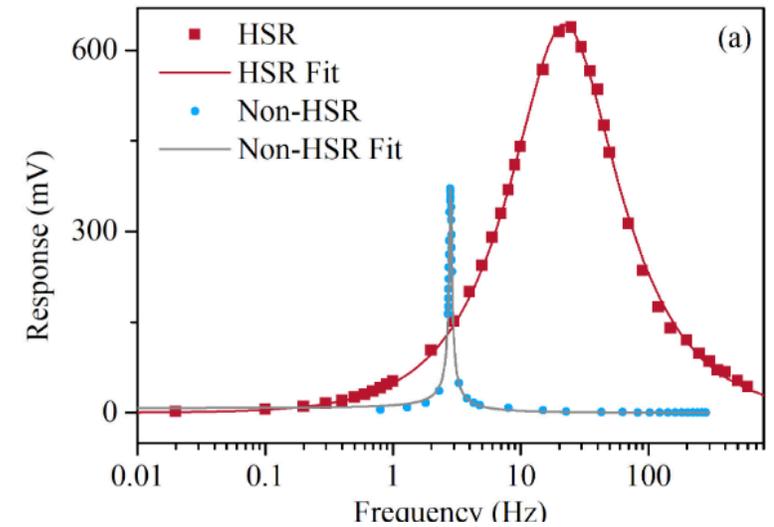
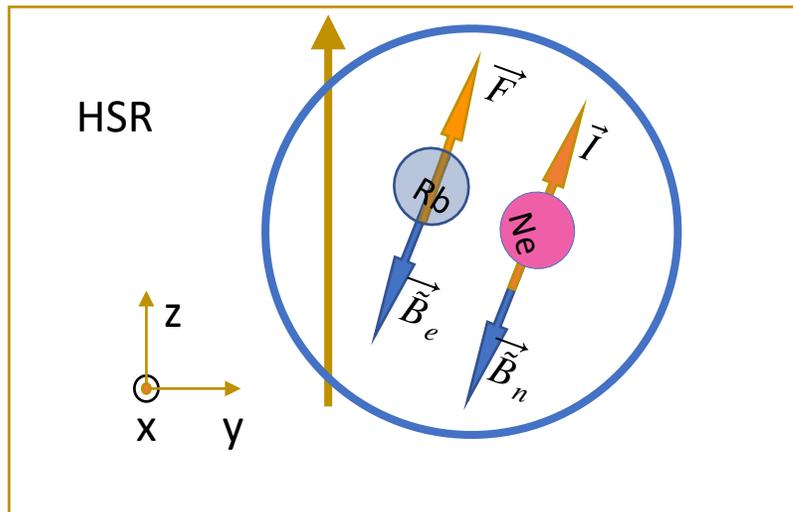
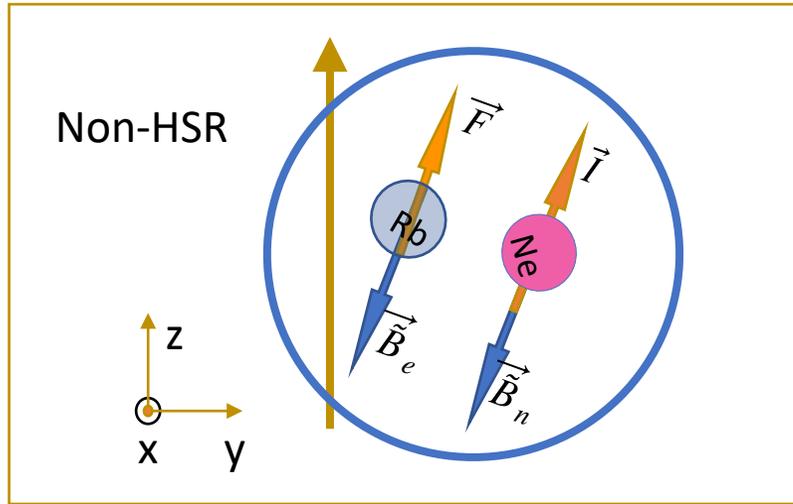
$$\beta_j = \frac{g_{aN}}{\gamma_N} \nabla a(j\Delta t) \cdot \hat{\mathbf{m}}(j\Delta t),$$

Daily Modulation

$$\hat{\mathbf{m}}_i(j\Delta t) \approx \mathbf{C}_i \cos(\omega_e j\Delta t + \theta_i) + \mathbf{D}_i,$$

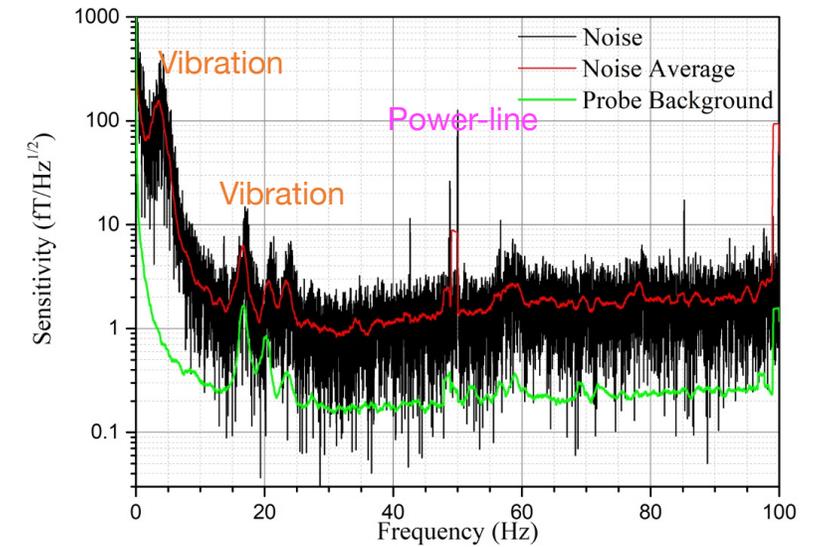
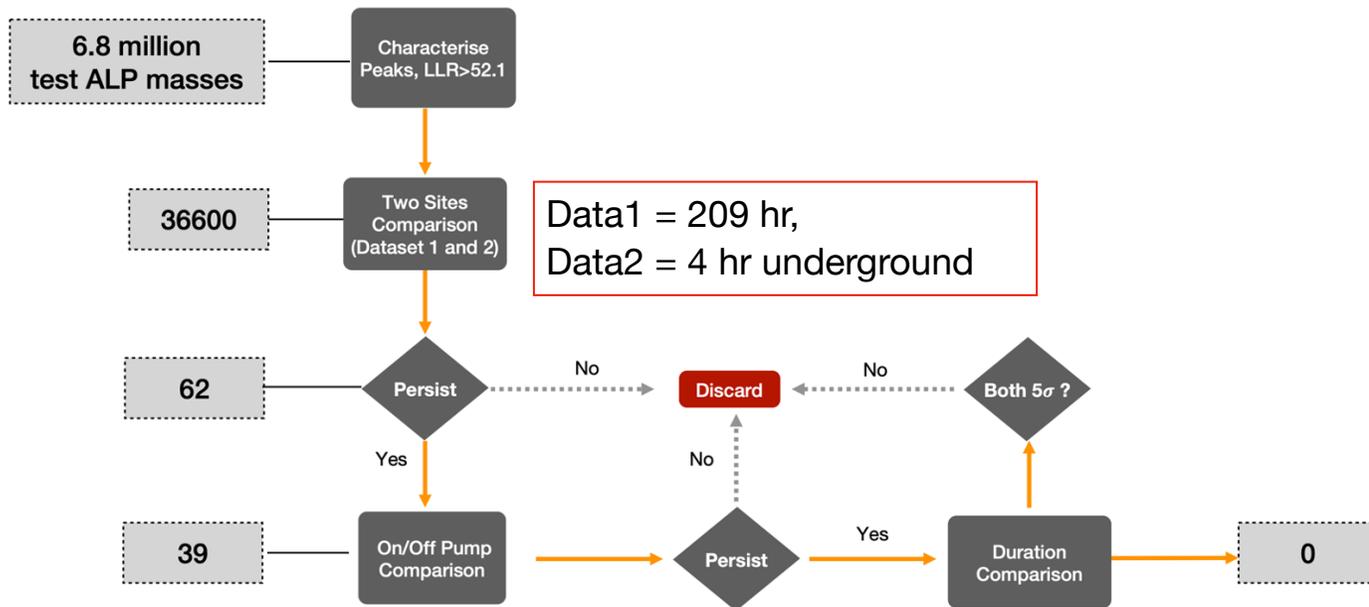
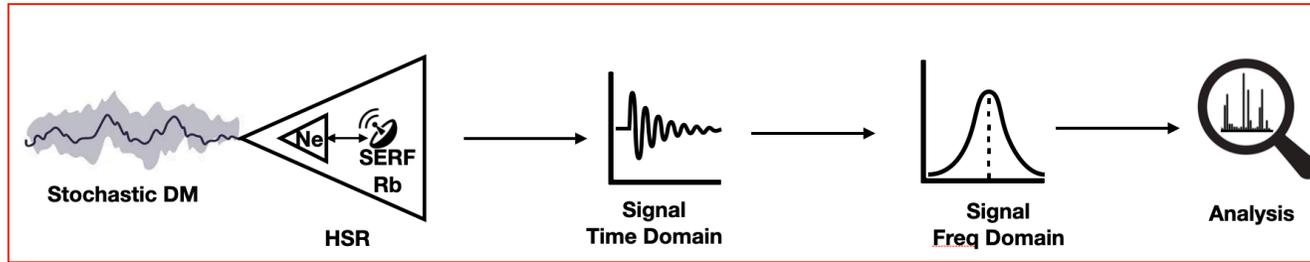
NMR resonance line width only 0.02 Hz, need broadband!

Hybrid Spin Resonance: High bandwidth



3 order of magnitude bandwidth improvement!

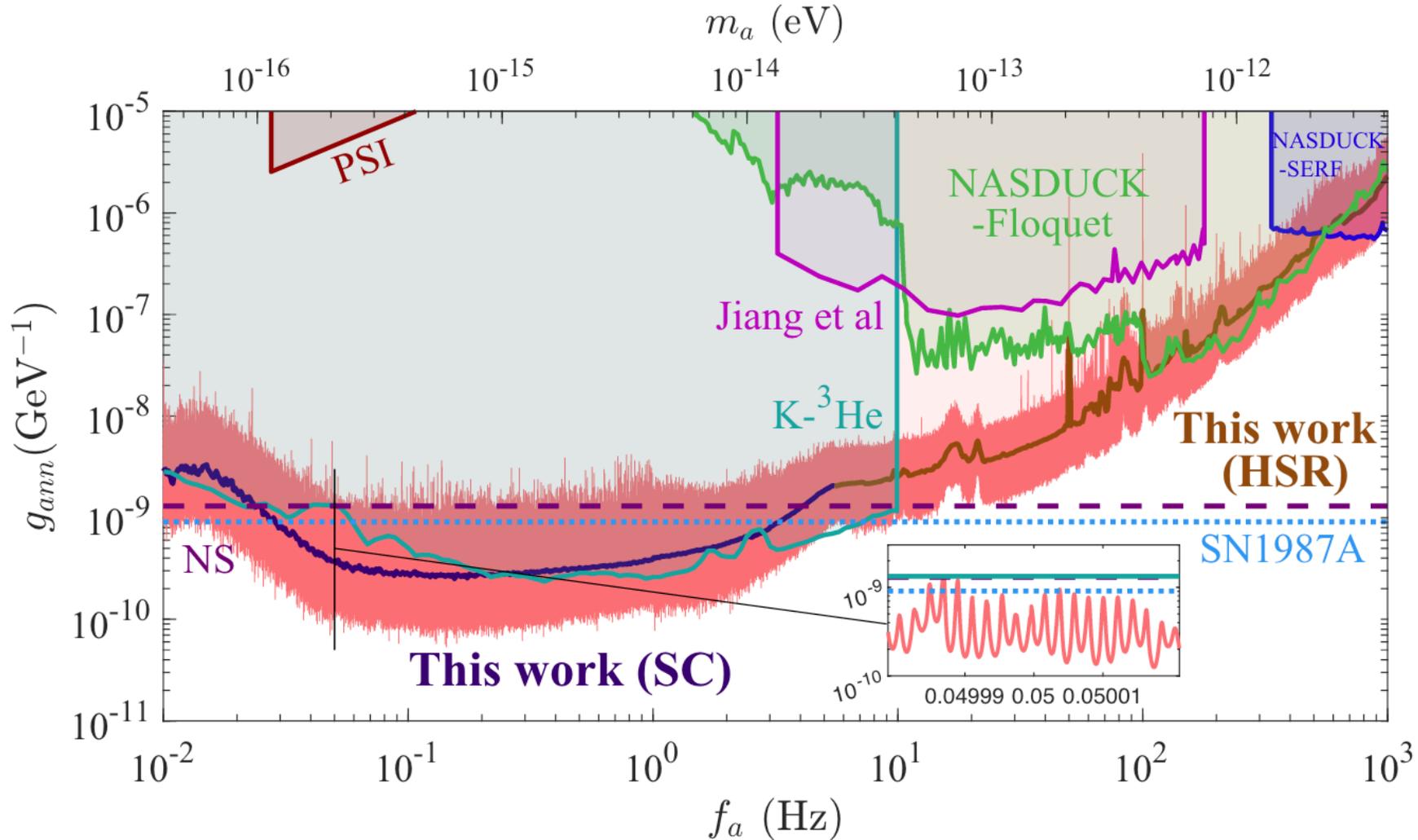
Broadband Axion Search



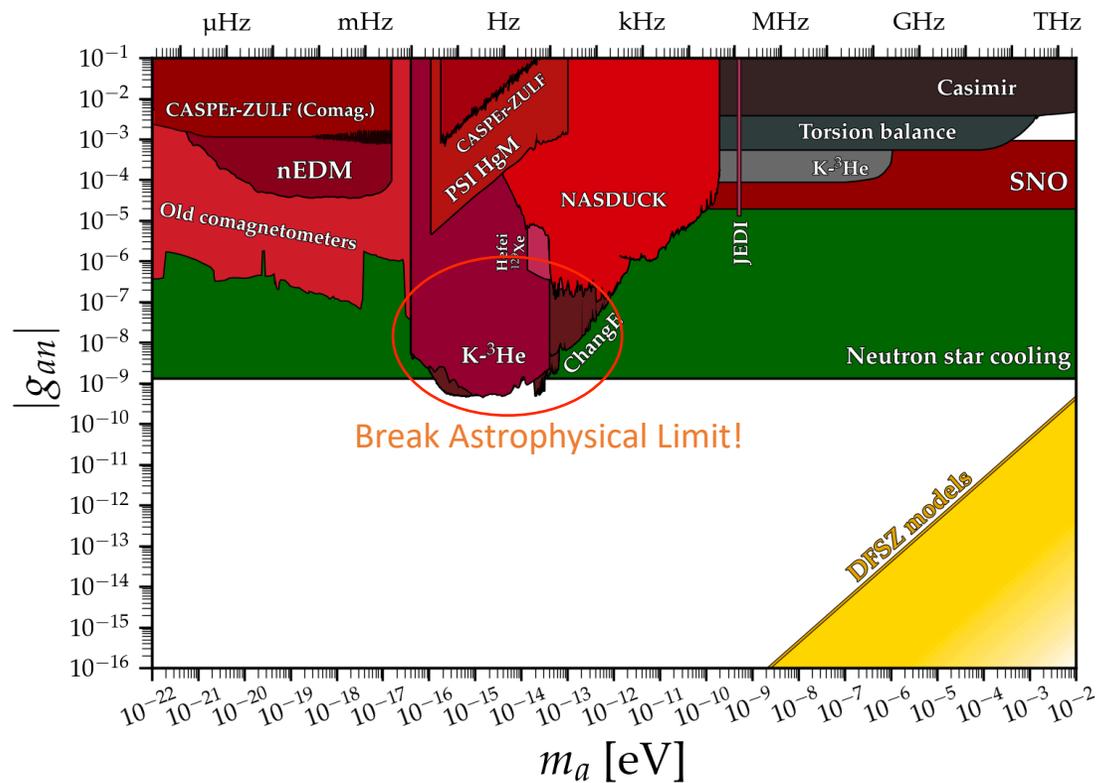
- Signal should be very narrow, $\frac{\delta f}{f} \approx 10^{-6}$
- Signal should persist
- Axion DM frequency unknown

Axion Limits

Wei et al., arXiv:2306.08039



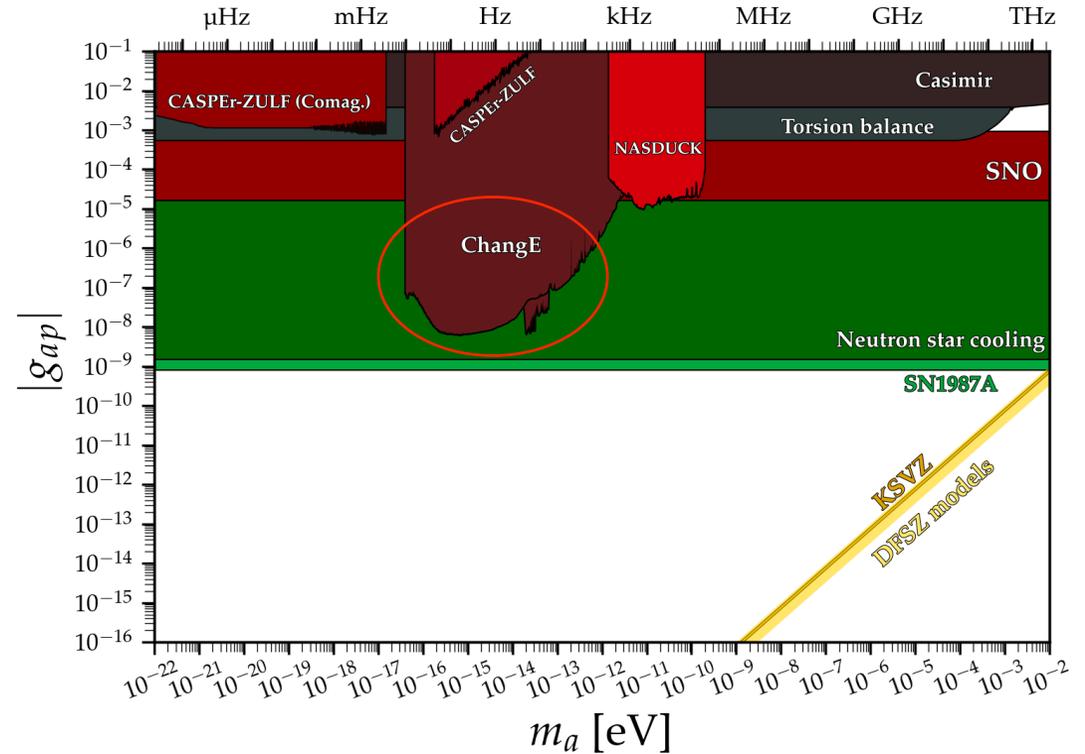
ChangE (Coupled Hot Atom eNsembles to search for liGht dark mattEr)



Xu et al, Commun. Phys. 7, 226, (2024)

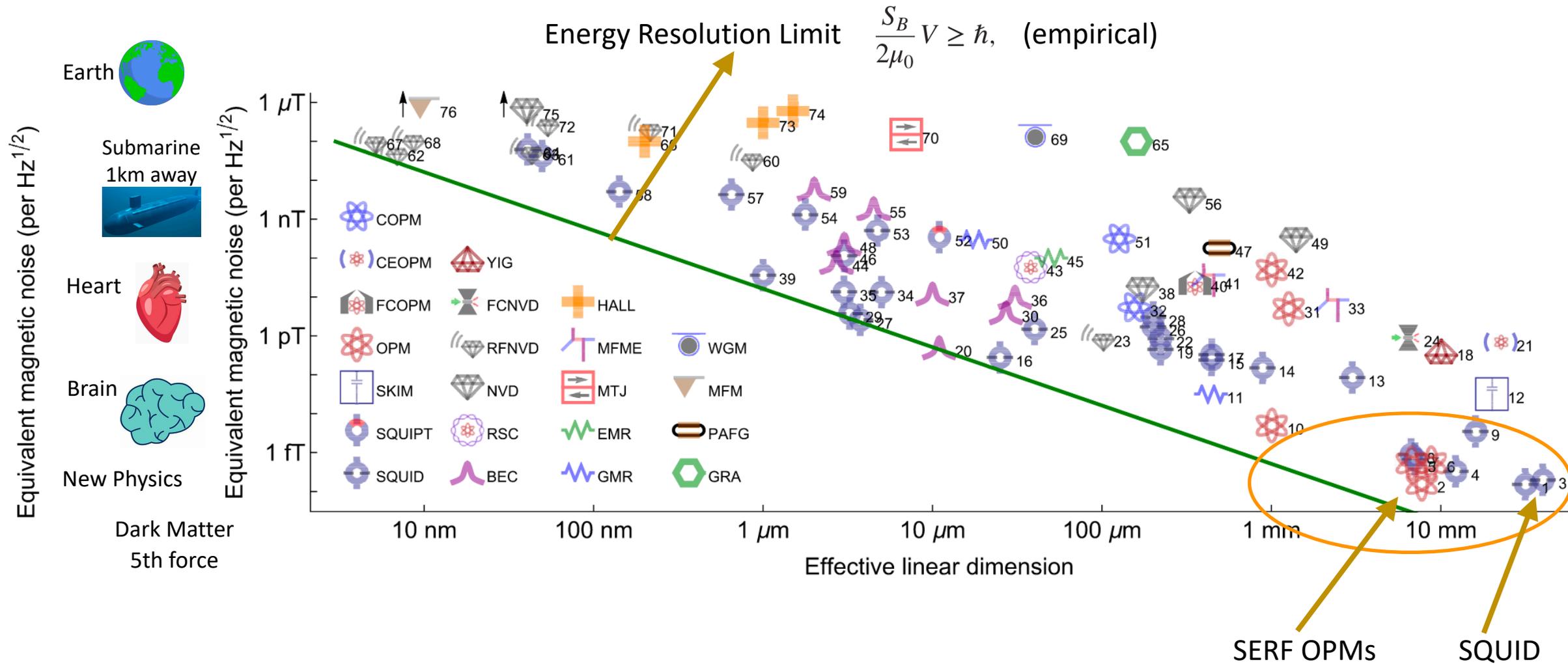
Wei et al., arXiv:2306.08039

$$b^{Ne} = b^n \zeta^n + b^p \zeta^p$$



<https://github.com/cajohare/AxionLimits/blob/master/docs/app.md>

LEvitated MAgnets for QUantum METrology (LEMAQUME)



REVIEWS OF MODERN PHYSICS, VOLUME 92, APRIL–JUNE 2020

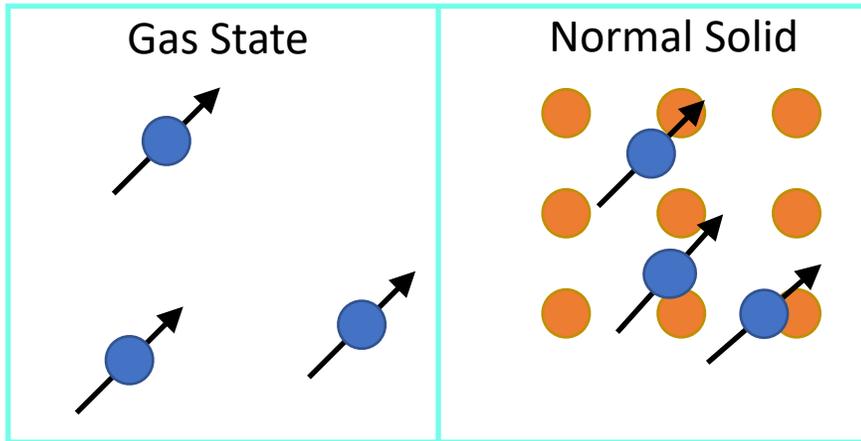
Ferromagnetism and Precession

PRL 116, 190801 (2016)

PHYSICAL REVIEW LETTERS

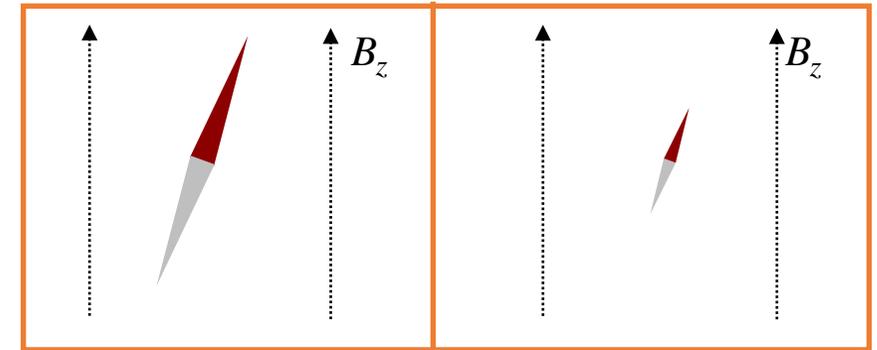


Precessing Ferromagnetic Needle Magnetometer

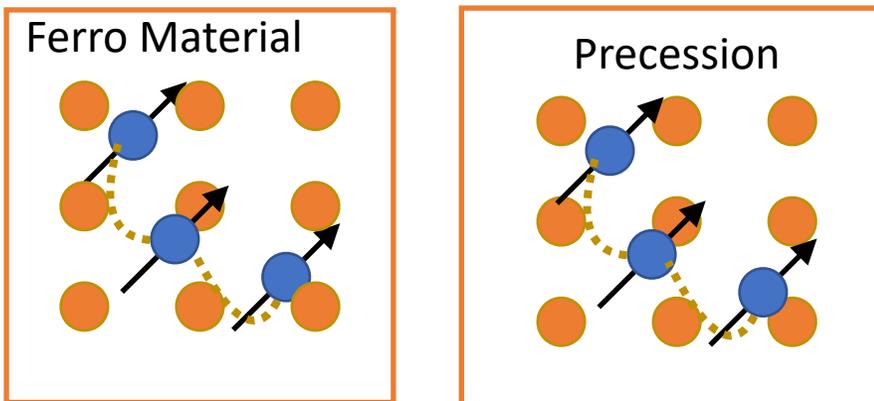
 Derek F. Jackson Kimball,¹ Alexander O. Sushkov,² and Dmitry Budker^{3,4,5}


$$\Delta\phi = \frac{\Delta S_{\perp}}{S_{\parallel}}$$

$$\approx \sqrt{\frac{R_{rel}}{Nt}}$$

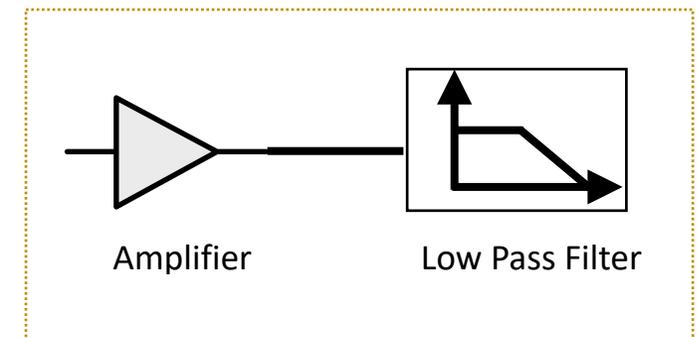
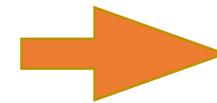


Libration v.s Precession



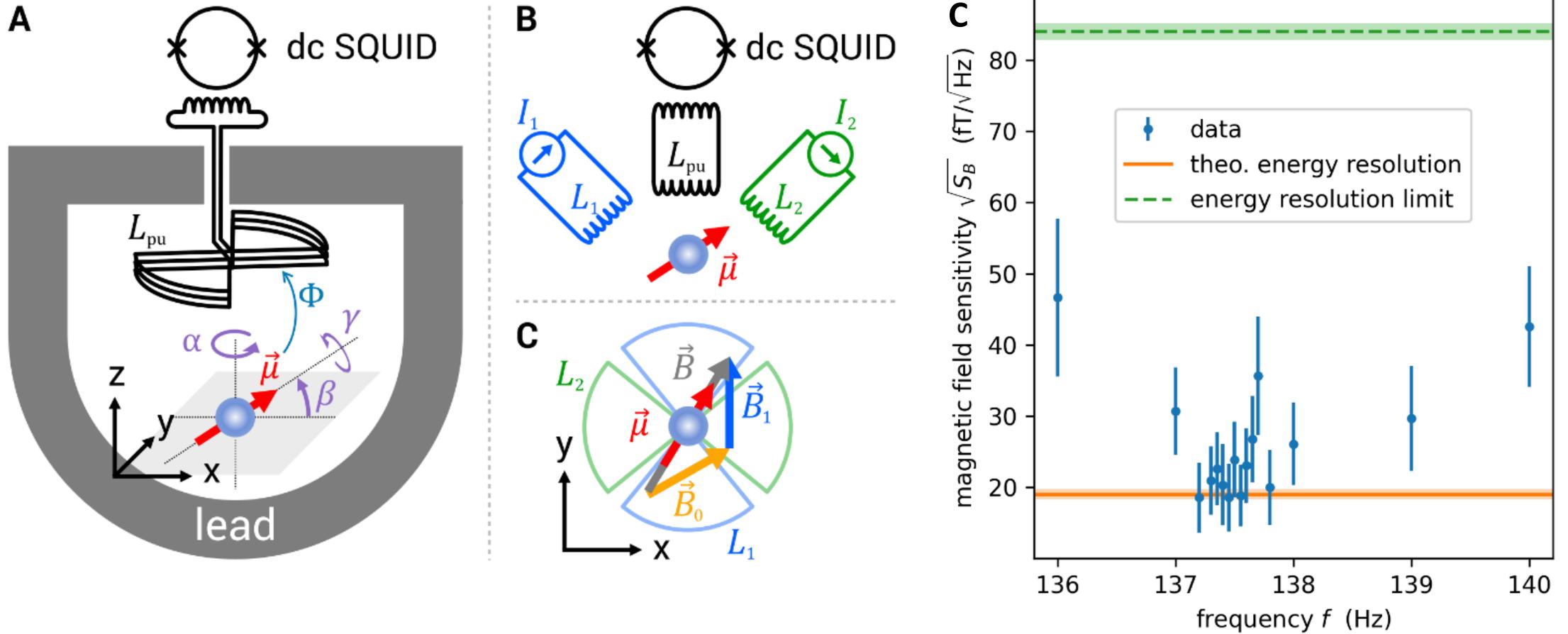
$$\Delta\phi = \frac{\Delta S_{\perp}}{S_{\parallel}} \frac{1}{\sqrt{t}}$$

$$\approx \sqrt{\frac{2\alpha k_B T}{N\hbar\omega_0^2 t}}$$

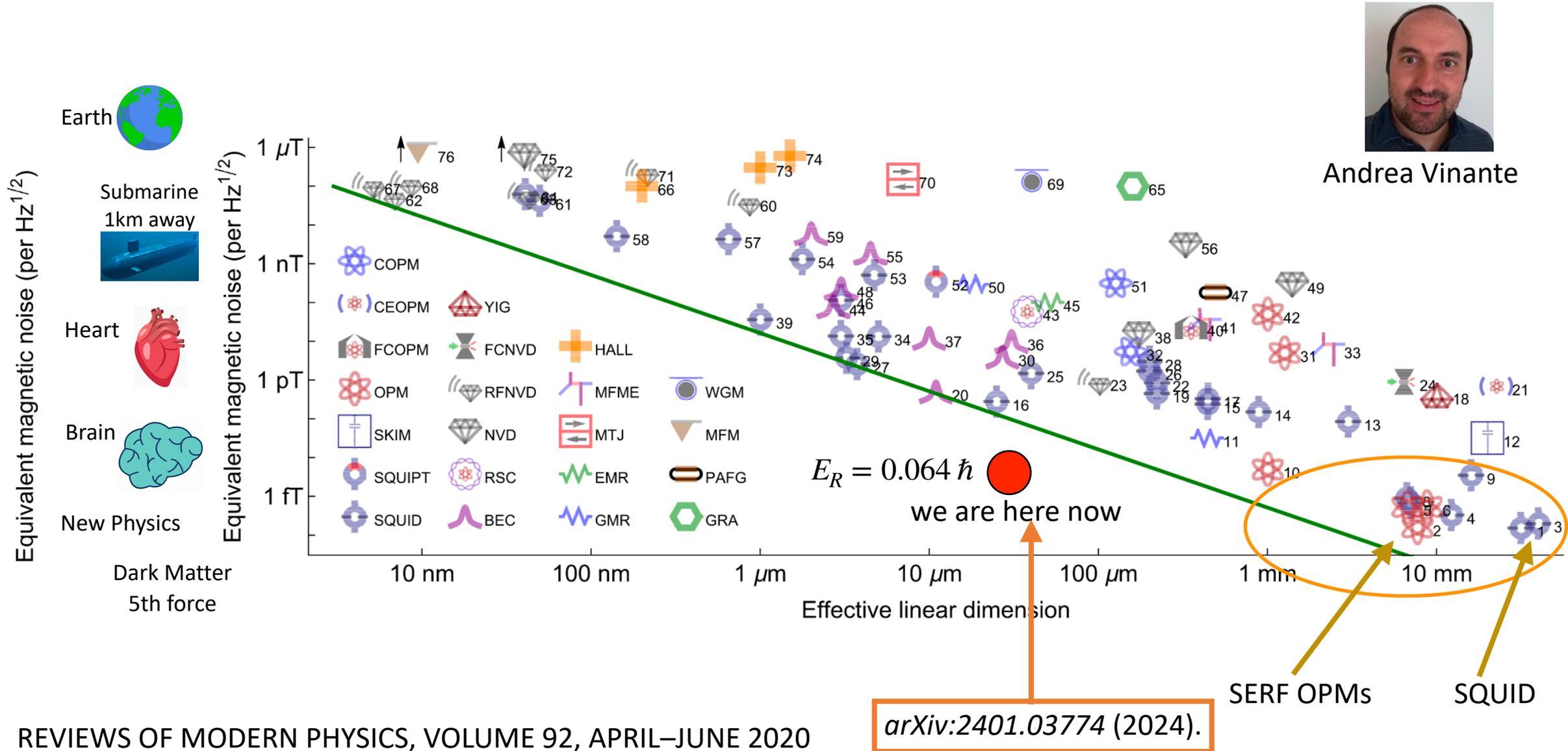


Ferromagnetic Resonance (FMR) v.s Precession

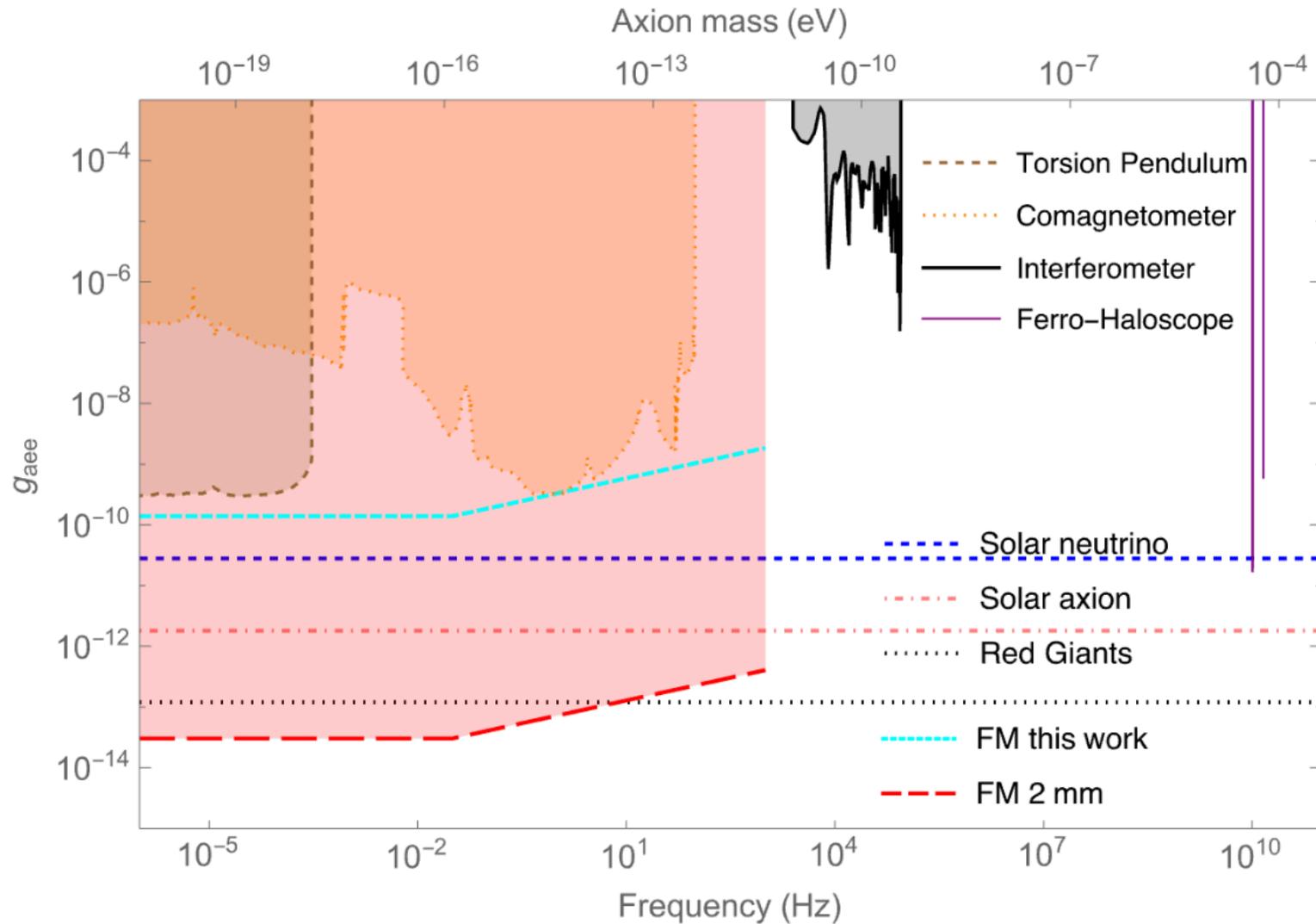
LEMAQUME: Magnetometry



LEMAQUME: Magnetometry



LEMAQUME: Search For Axions



Thanks!

New Interaction Collaboration



Dmitry Budker
Mainz
& UC Berkeley



高海燕
Duke
& BNL



房建成
北航



魏凯
北航



彭新华
中科大



符长波
复旦



刘永椿
清华



丛磊
Mainz

Dark Matter Collaboration



Dmitry Budker



刘佳



房建成



魏凯



王小平

And many others

LEMAQUME



Dmitry Budker



Andrea Vinante



Changhao Xu



Guofeng Qu



Pavel Fadeev

Backup Slides

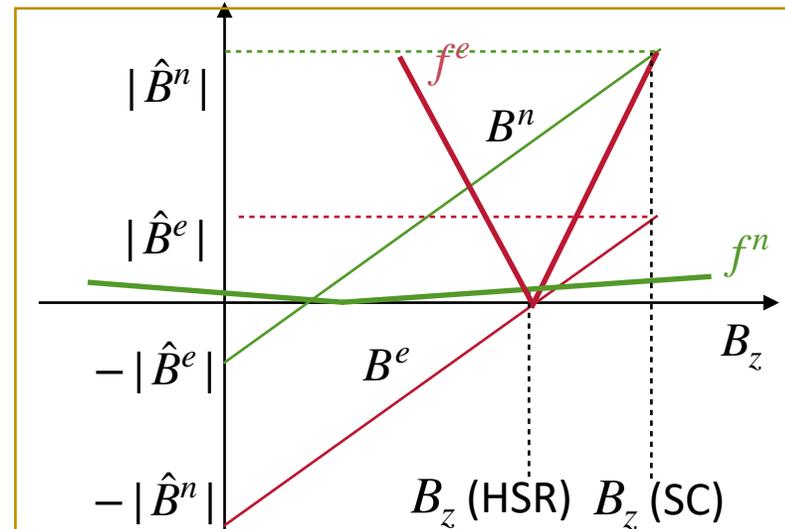
Screening Effect

$$B^e = -|\hat{B}^n| + B_z$$

$$B^n = -|\hat{B}^e| + B_z$$

$$B_z(\text{HSR}) \approx |\hat{B}^n|$$

$$B_z(\text{SC}) \approx |\hat{B}^n| + |\hat{B}^e|$$



Optically Pumped Magnetometer (OPM)

nature physics

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Published: April 2007

Optical magnetometry

 Dmitry Budker  & Michael Romalis

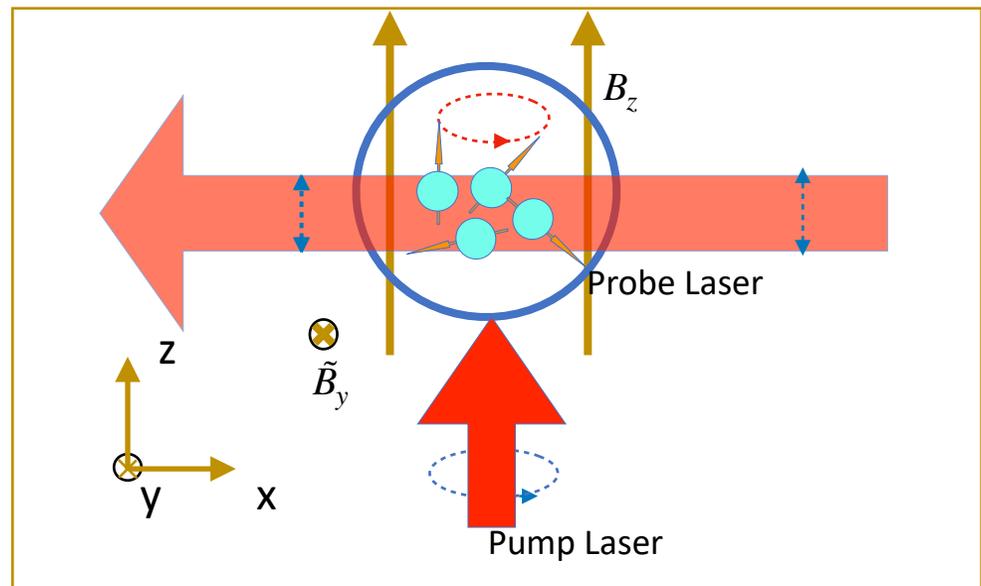
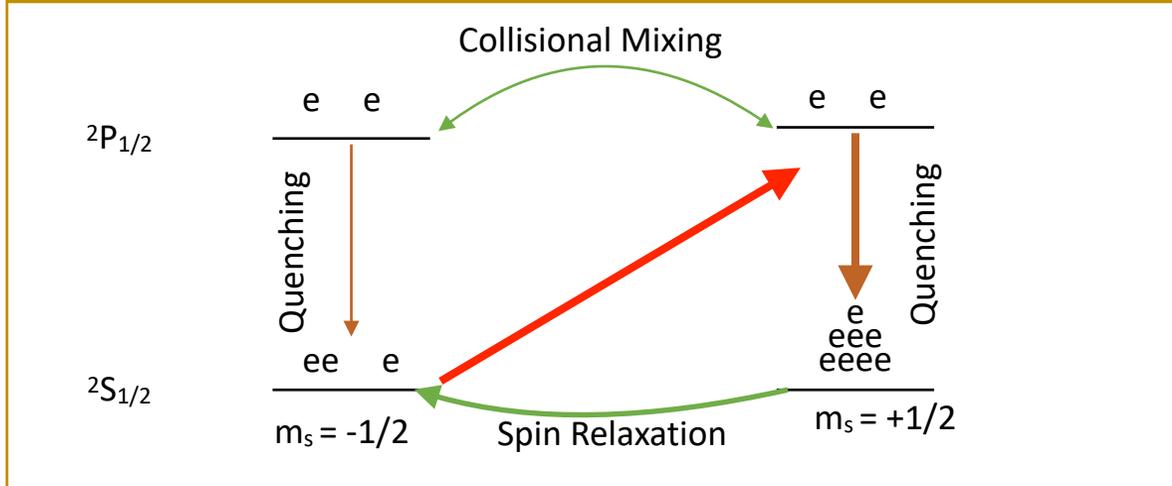
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Dmitry Budker **18k** Accesses | **1209** Citations | **11** Altmetric | [Metrics](#)

UC Berkeley and Mainz

Michael Romalis
Princeton

- Optical Pumping
- Spin Precession
- Probe (Faraday Rotation)



OPM: Alkali Magnetometer

Budker's famous Equation 1:

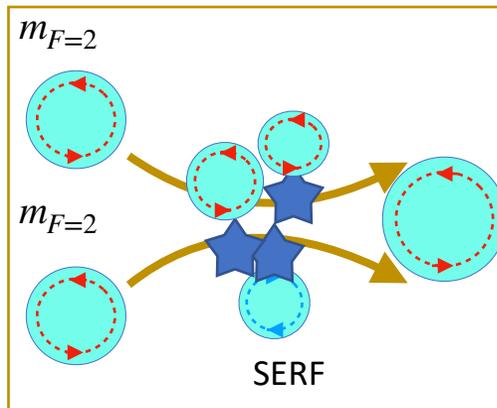
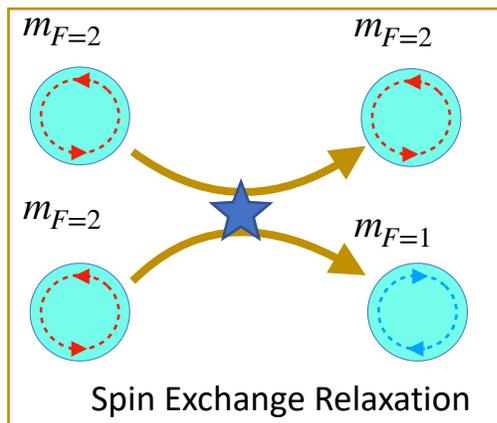
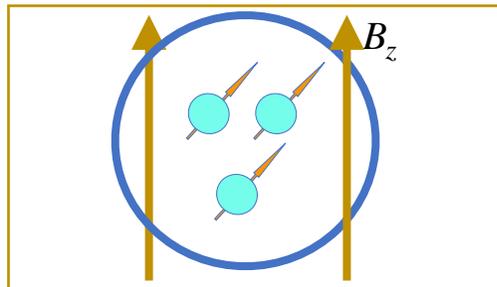
$$\delta B = \frac{1}{\gamma \sqrt{n T_2 V t}} \quad (\text{Spin Projection Noise})$$

Spin Exchange Relaxation

$$\frac{1}{T_2^{SE}} = \frac{\omega^2}{R_{SE}} \left[\frac{1}{2} - \frac{(2I+1)^2}{Q(P)^2} \right] Q(P)^2$$

SERF (Spin Exchange Relaxation Free):

$$R_{SE} \gg \omega^2 \Rightarrow T_2^{SE} \rightarrow \infty$$



High-Sensitivity Atomic Magnetometer Unaffected by Spin-Exchange Relaxation

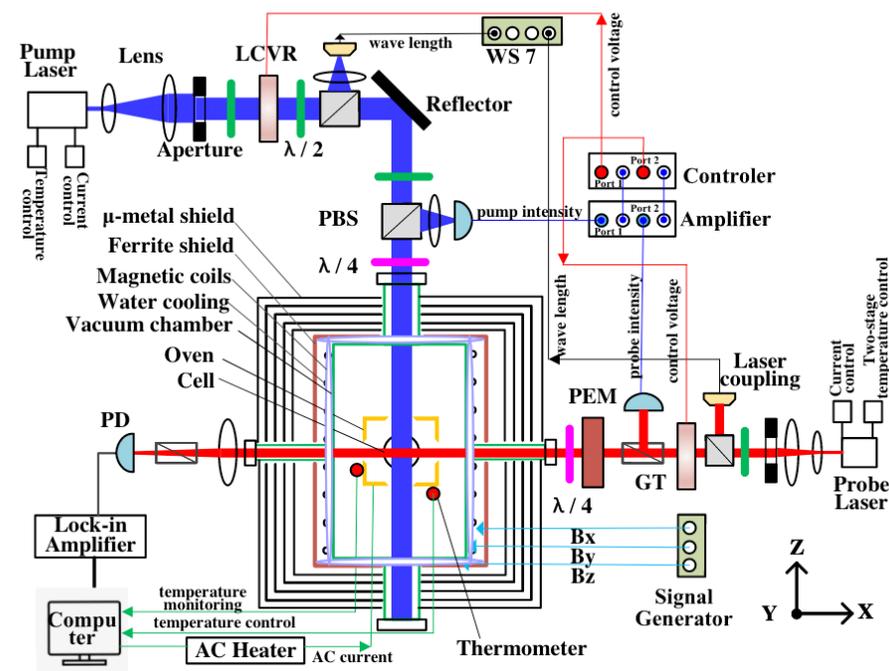
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$$V_2 = -g_A^X g_A^Y \frac{\hbar c}{4\pi} \boldsymbol{\sigma}_X \cdot \boldsymbol{\sigma}_Y \frac{1}{r} e^{-r/\lambda}. \quad (46)$$

$$V_3|_{AA} = -\hbar c g_A^X g_A^Y \lambda^2 \left[\boldsymbol{\sigma}_X \cdot \boldsymbol{\sigma}_Y \left[\frac{1}{r^3} + \frac{1}{\lambda r^2} + \frac{4\pi}{3} \delta(\mathbf{r}) \right] - (\boldsymbol{\sigma}_X \cdot \hat{\mathbf{r}}) (\boldsymbol{\sigma}_Y \cdot \hat{\mathbf{r}}) \left(\frac{3}{r^3} + \frac{3}{\lambda r^2} + \frac{1}{\lambda^2 r} \right) \right] \frac{e^{-r/\lambda}}{4\pi}, \quad (47)$$

$$V_{4+5}|_{AA} = -\frac{g_A^X g_A^Y \hbar^2}{4c} \left\{ \boldsymbol{\sigma}_X \cdot \left(\frac{\mathbf{p}_Y}{m_Y^2} \times \hat{\mathbf{r}} \right), \left(\frac{1}{r^2} + \frac{1}{\lambda r} \right) \frac{e^{-r/\lambda}}{8\pi} \right\} \Rightarrow g_A^X g_A^Y \frac{\hbar^2}{16\pi c} \frac{m_X}{m_Y(m_X + m_Y)} \boldsymbol{\sigma}_X \cdot (\mathbf{v} \times \hat{\mathbf{r}}) \left(\frac{1}{r^2} + \frac{1}{\lambda r} \right) e^{-r/\lambda}, \quad (48)$$

$$V_8 = g_A^X g_A^Y \frac{\hbar}{c} \left[\left\{ \frac{\boldsymbol{\sigma}_X \cdot \mathbf{p}_X}{m_X}, \left\{ \frac{\boldsymbol{\sigma}_Y \cdot \mathbf{p}_Y}{m_Y}, \frac{e^{-r/\lambda}}{16\pi r} \right\} \right\} - \frac{1}{2} \left\{ \frac{\boldsymbol{\sigma}_X \cdot \mathbf{p}_Y}{m_Y}, \left\{ \frac{\boldsymbol{\sigma}_Y \cdot \mathbf{p}_Y}{m_Y}, \frac{e^{-r/\lambda}}{16\pi r} \right\} \right\} - \frac{1}{2} \left\{ \frac{\boldsymbol{\sigma}_X \cdot \mathbf{p}_X}{m_X}, \left\{ \frac{\boldsymbol{\sigma}_Y \cdot \mathbf{p}_X}{m_X}, \frac{e^{-r/\lambda}}{16\pi r} \right\} \right\} \right] \\ \Rightarrow -g_A^X g_A^Y \frac{\hbar}{8\pi c} (\boldsymbol{\sigma}_X \cdot \mathbf{v}) (\boldsymbol{\sigma}_Y \cdot \mathbf{v}) \frac{1}{r} e^{-r/\lambda}. \quad (49)$$