

The contributions of $\rho \rightarrow \omega\pi$ in $B \rightarrow \bar{D}^{()}\omega\pi$ decays*

王文飞 (山西大学)

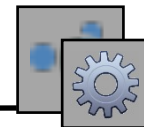
第二十一届全国重味物理和CP破坏研讨会
南华大学

2024-10-27

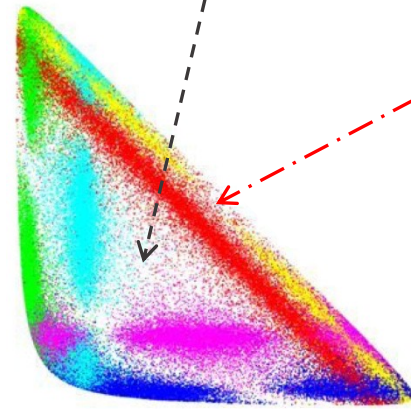
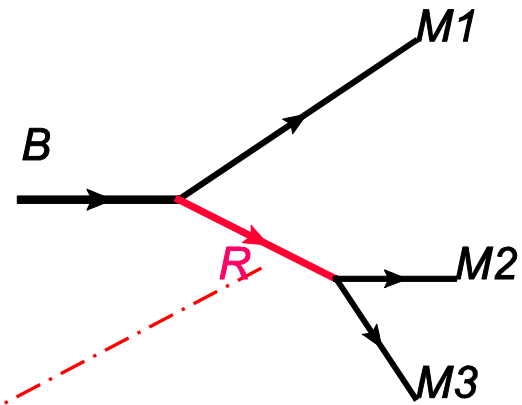
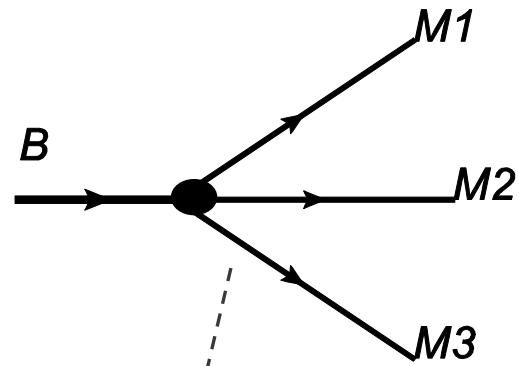
Based on JHEP 01 (2024) 047 with Yu-Shan Ren & Ai-Jun Ma



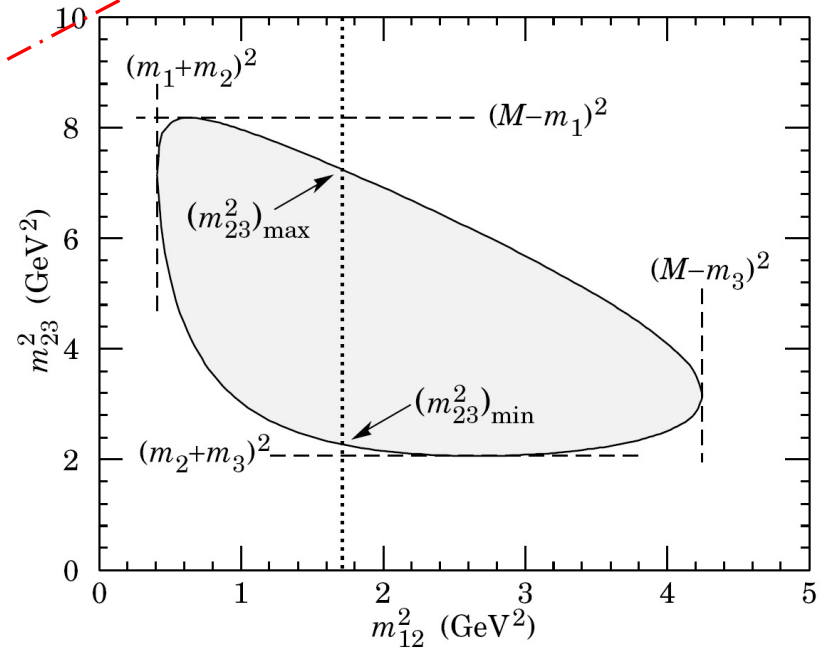
Outline



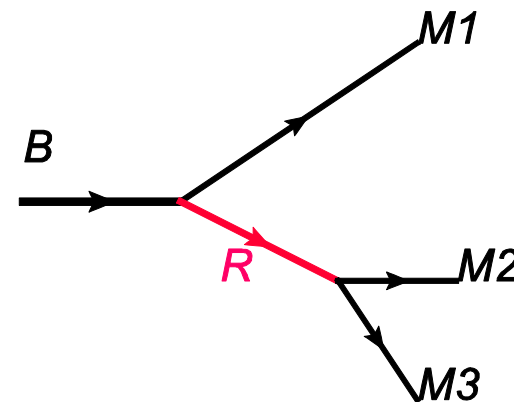
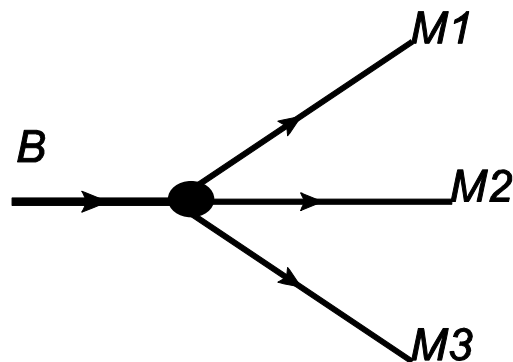
- **A brief introduction for 3-body B meson decays**
- **Contributions of the $\rho(770, 1450) \rightarrow \omega\pi$ in $B \rightarrow \bar{D}^{(*)}\omega\pi$**
- **Results and discussions**
- **Summary**



Dalitz Plot



$$B \rightarrow M1 M2 M3$$



The weak effective Hamiltonian:

1.

$$\mathcal{H}_{\text{eff}} = \frac{G_F}{\sqrt{2}} \left\{ V_{ub} V_{uq}^* [C_1(\mu) Q_1^u(\mu) + C_2(\mu) Q_2^u(\mu)] - V_{tb} V_{tq}^* \left[\sum_{i=3}^{10} C_i(\mu) Q_i(\mu) \right] \right\} + \text{H.c.}, \quad (2)$$

where $q = d, s$. The functions Q_i ($i = 1, \dots, 10$) are the local four-quark operators:

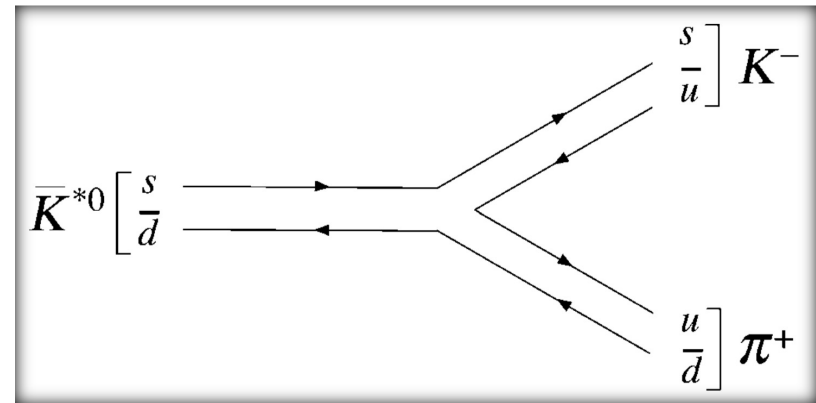
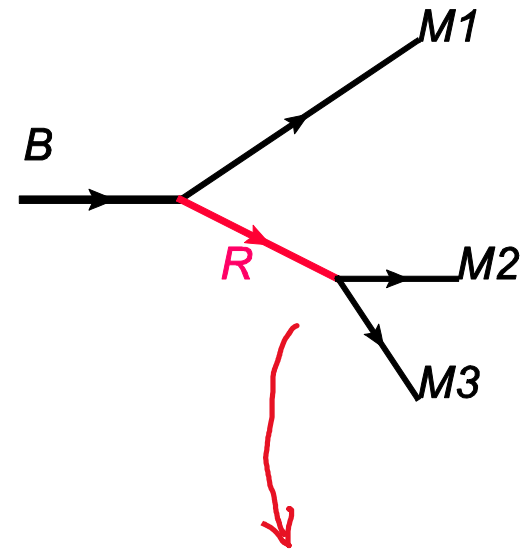
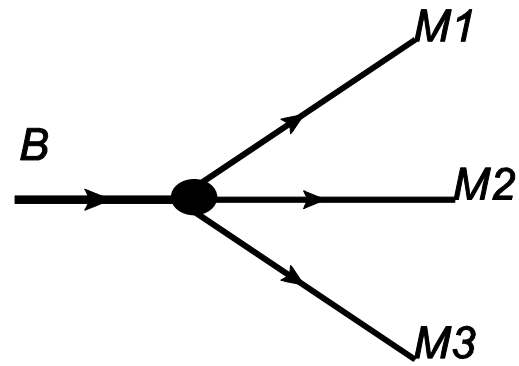
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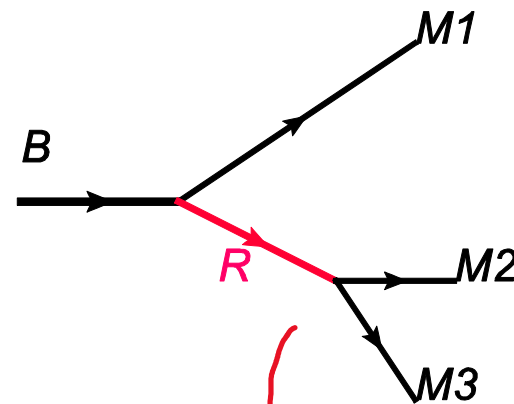
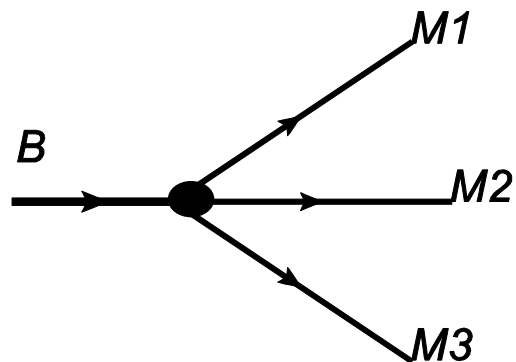
$$\mathcal{A}_{\text{total}} = \sum_i \mathcal{A}_{NR}^i + \sum_j \mathcal{A}_R^j$$



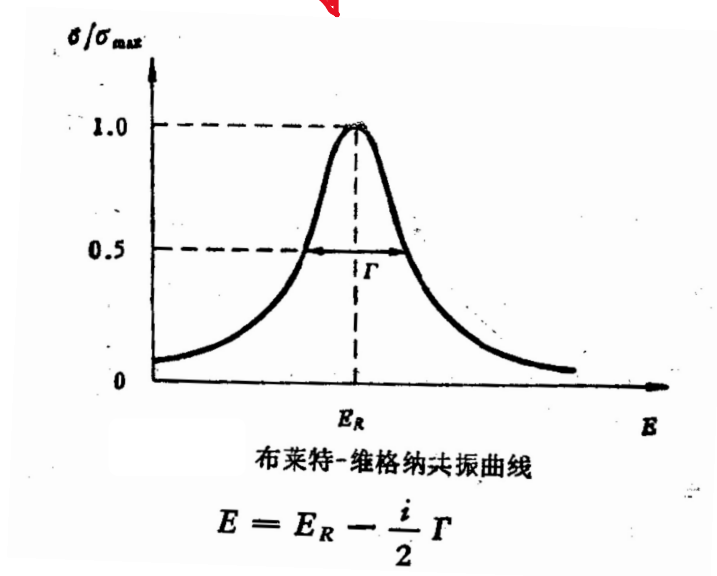
A brief introduction for 3-body B meson decays

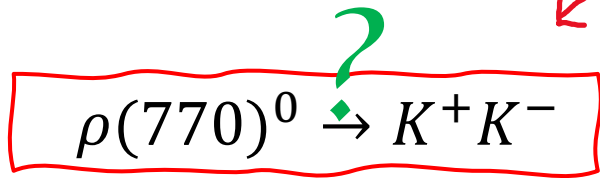
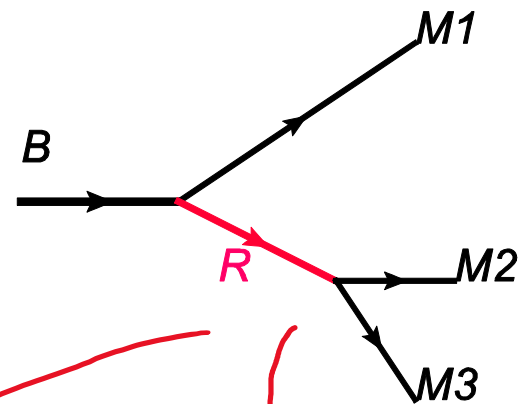
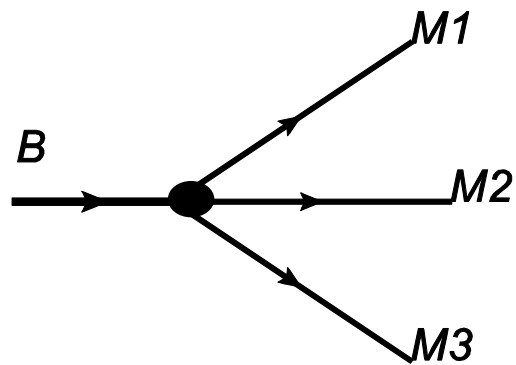
$$B^+ \rightarrow K^+ K^- \pi^+$$



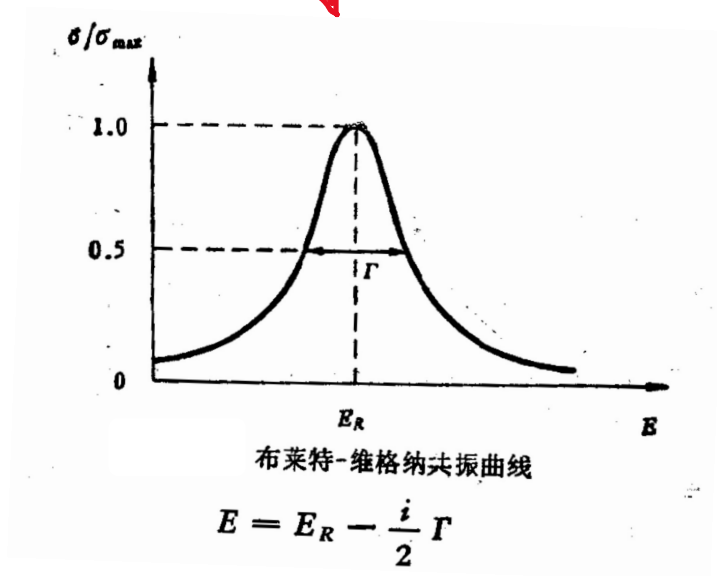


$$BW_R = \frac{m_R^2}{m_R^2 - s - im_R \Gamma_R(s)}$$





$$BW_R = \frac{m_R^2}{m_R^2 - s - im_R \Gamma_R(s)}$$





PHYSICAL REVIEW D

VOLUME 15, NUMBER 11

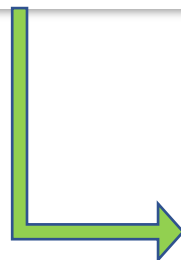
1 JUNE 1977

High-statistics study of the reactions $\pi^- p \rightarrow K^- K^+ n$ and $\pi^+ n \rightarrow K^- K^+ p$ at 6 GeV/c*

A. J. Pawlicki, D. S. Ayres, D. Cohen, R. Diebold, S. L. Kramer, and A. B. Wicklund

Argonne National Laboratory, Argonne, Illinois 60439

(Received 23 December 1976; revised manuscript received 15 March 1977)



ambiguities. As discussed above, the P wave is consistent with the tail of the ρ^0 decaying into $K^- K^+$, with a $\rho K K$ coupling that agrees with SU(3), including the sign. Only one of the ambiguous sol-



A brief introduction for 3-body B meson decays



PHYSICAL REVIEW D

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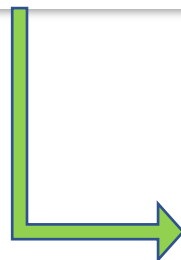
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BW-Tail



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ELSEVIER



Further study of the $I = 1 K \bar{K}$ structure near threshold

A. Astier, J. Cohen-Ganouna, M. Della Negra, B. Maréchal, L. Montanet, M. Tomas[†], M. Baubillier, J. Duboc



3) The $I = 1 K \bar{K}$ channel is dominated, at threshold, by a virtual bound state resonance

Page-4



PHYSICAL REVIEW D **94**, 072001 (2016)

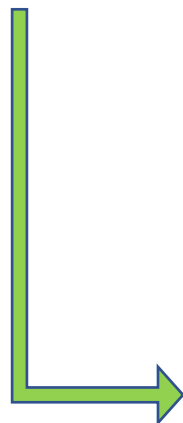
Amplitude analysis of $B^- \rightarrow D^+ \pi^- \pi^-$ decays

R. Aaij *et al.**

(LHCb Collaboration)

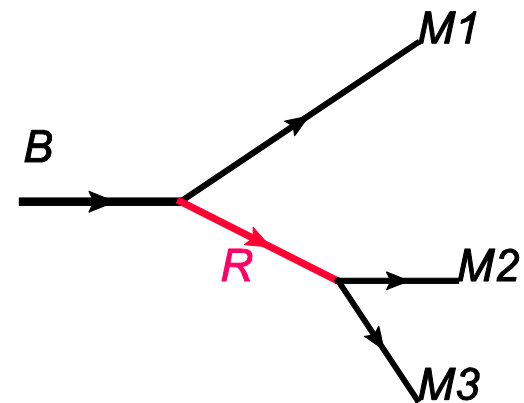
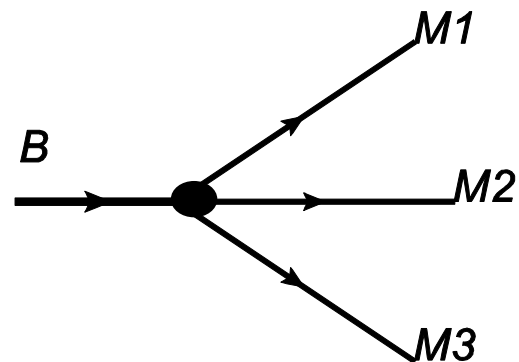
(Received 4 August 2016; published 5 October 2016)

虚 (共振态粒子的) 贡献



Contribution	Fit fraction (%)
$D_2^*(2460)^0$	35.7 ± 0.6
$D_1^*(2680)^0$	8.3 ± 0.6
$D_3^*(2760)^0$	1.0 ± 0.1
$D_2^*(3000)^0$	0.23 ± 0.07
$D_v^*(2007)^0$	10.8 ± 0.7
B_v^{*0}	2.7 ± 1.0
Total S wave	57.0 ± 0.8
Total fit fraction	115.7

$$B \rightarrow M1 M2 M3$$



The weak effective Hamiltonian:

1.

$$\mathcal{H}_{\text{eff}} = \frac{G_F}{\sqrt{2}} \left\{ V_{ub} V_{uq}^* [C_1(\mu) Q_1^u(\mu) + C_2(\mu) Q_2^u(\mu)] - V_{tb} V_{tq}^* \left[\sum_{i=3}^{10} C_i(\mu) Q_i(\mu) \right] \right\} + \text{H.c.}, \quad (2)$$

where $q = d, s$. The functions Q_i ($i = 1, \dots, 10$) are the local four-quark operators:

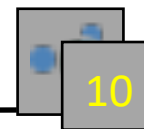
The total amplitude within isobar approach:

2.

$$\mathcal{A}_{\text{total}} = \sum_i \mathcal{A}_{NR}^i + \sum_j \mathcal{A}_R^j$$



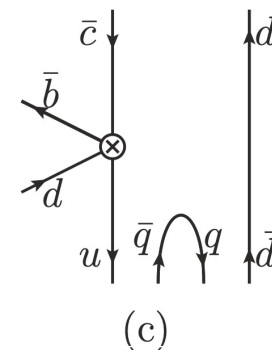
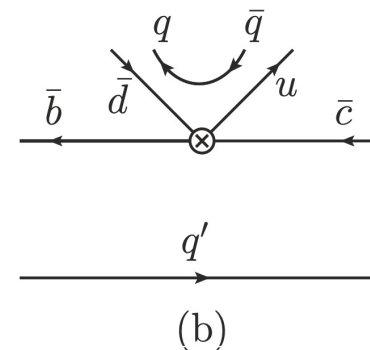
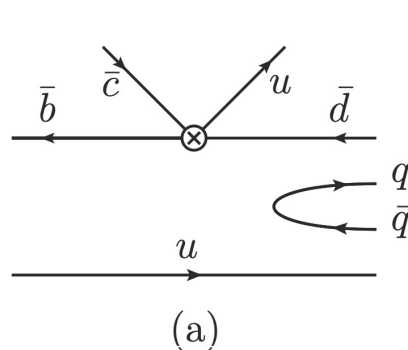
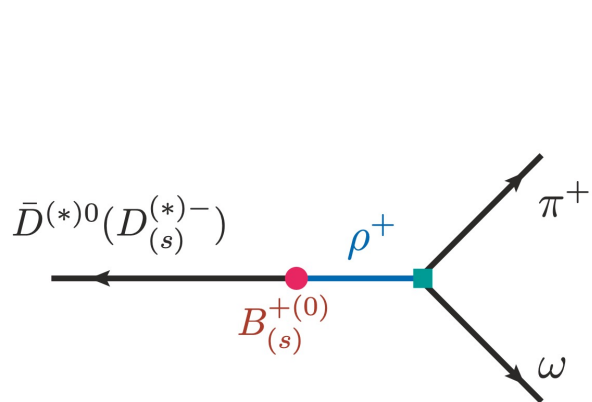
Contributions of the $\rho(770, 1450) \rightarrow \omega\pi$ in $B \rightarrow \bar{D}^{(*)}\omega\pi$



MOTIVATIONS:

- ❖ $B^+ \rightarrow \bar{D}^{(*)0}\omega\pi^+$ and $B^0 \rightarrow D^{(*)-}\omega\pi^+$ have been measured but without any theoretical predictions [PRD64-092001, PRD74-012001, PRD92-012013]
- ❖ The $\omega\pi^+$ is related to the resonances $\rho(1450)$ and $\rho(770)$
- ❖ Test the factorization hypothesis for B decays with $B^0 \rightarrow D^{(*)-}\omega\pi^+$ and $B_s^0 \rightarrow D_s^{*-}\omega\pi^+$ with the longitudinal polarization Γ_L/Γ [PLB89-105]

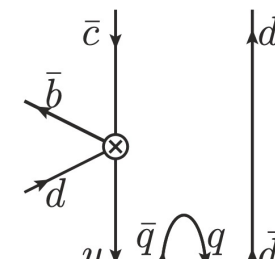
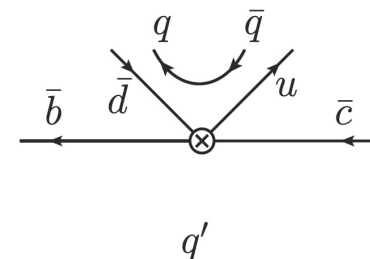
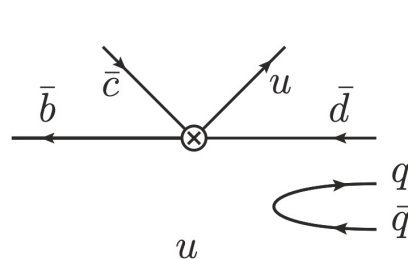
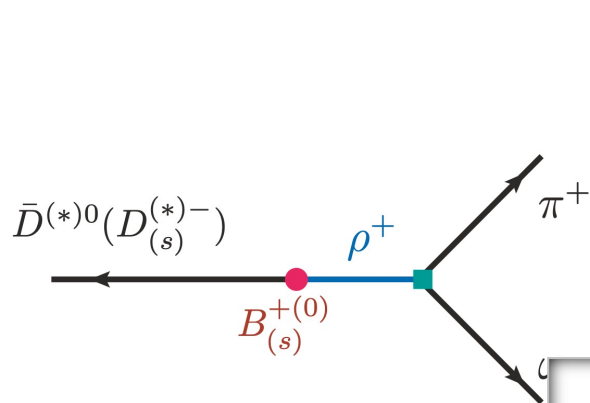
APPROACH:



With PQCD approach

$$\begin{aligned} \mathcal{A} &= \langle (\omega\pi)_{P\text{-wave}} D^{(*)} | \mathcal{H}_{\text{eff}} | B \rangle \\ &= \phi_B \otimes \mathcal{H} \otimes \phi_{\omega\pi}^{P\text{-wave}} \otimes \phi_{D^{(*)}} \end{aligned}$$

APPROACH:



Physics Letters B 763 (2016) 29–39

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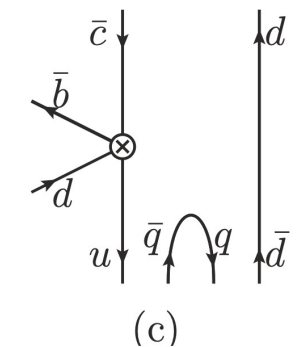
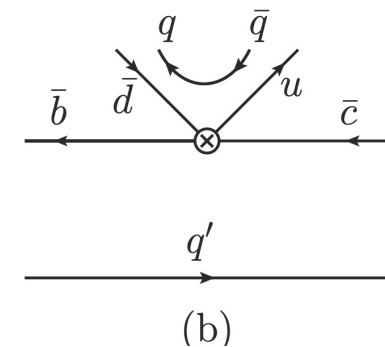
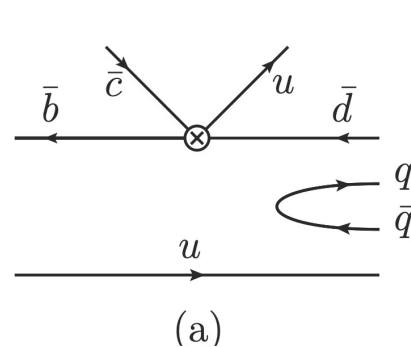
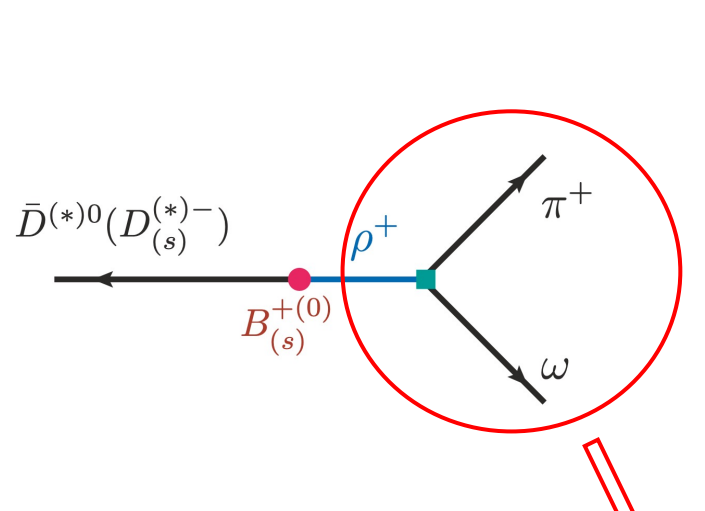
With PQCD approach

Quasi-two-body decays $B \rightarrow K\rho \rightarrow K\pi\pi$ in perturbative QCD approach

Wen-Fei Wang^{a,b}, Hsiang-nan Li^{a,*}



APPROACH:

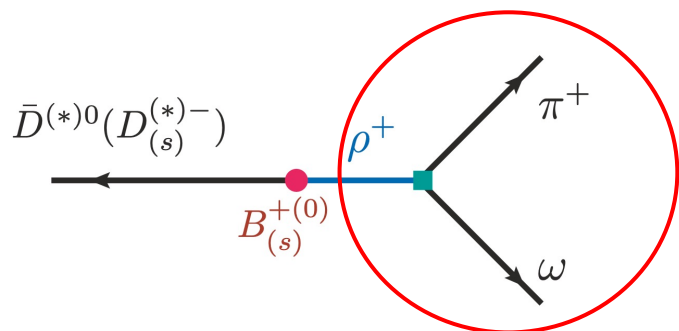


With PQCD approach

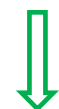
$$\phi_{\omega\pi,L}^{P\text{-wave}}(x, s) = \frac{-1}{\sqrt{2N_c}} [\sqrt{s} \not{\epsilon}_L \phi^0(x, s) + \not{\epsilon}_L \not{p} \phi^t(x, s) + \sqrt{s} \phi^s(x, s)],$$

$$\phi_{\omega\pi,T}^{P\text{-wave}}(x, s) = \frac{-1}{\sqrt{2N_c}} [\sqrt{s} \not{\epsilon}_T \phi^v(x, s) + \not{\epsilon}_T \not{p} \phi^T(x, s) + \sqrt{s} i \epsilon_{\mu\nu\rho\sigma} \gamma_5 \gamma^\mu \epsilon_T^{*\nu} n^\rho v^\sigma \phi^a(x)],$$

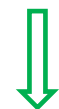
APPROACH:



$$\mathcal{L}_{\rho\omega\pi} = g_{\rho\omega\pi} \epsilon_{\mu\nu\alpha\beta} \partial^\mu \rho^\nu \partial^\alpha \omega^\beta \pi.$$



$$\langle \omega(p_a, \lambda) \pi(p_b) | j_\mu(0) | 0 \rangle = i \epsilon_{\mu\nu\alpha\beta} \epsilon^\nu(p_a, \lambda) p_b^\alpha p^\beta F_{\omega\pi}(s)$$

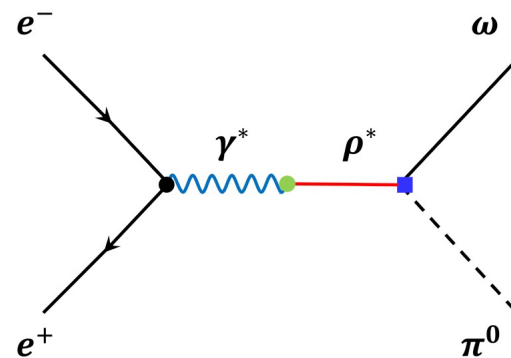
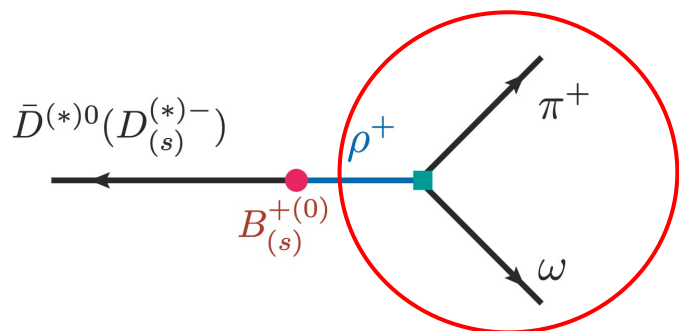


$$F_{\omega\pi}(s) = \frac{g_{\rho\omega\pi}}{f_\rho} \sum_{\rho_i} \frac{A_i e^{i\phi_i} m_{\rho_i}^2}{D_{\rho_i}(s)}$$

$$D_{\rho_i}(s) = m_{\rho_i}^2 - s - i\sqrt{s}\Gamma_{\rho_i}(s)$$

one has $A = 1$ and $\phi = 0$ for $\rho(770)$

APPROACH:



$$\langle \omega(p_a, \lambda)\pi(p_b) | j_\mu(0) | 0 \rangle = i\epsilon_{\mu\nu\alpha\beta} \epsilon^\nu(p_a, \lambda) p_b^\alpha p^\beta F_{\omega\pi}(s)$$



$$F_{\omega\pi}(s) = \frac{g_{\rho\omega\pi}}{f_\rho} \sum_{\rho_i} \frac{A_i e^{i\phi_i} m_{\rho_i}^2}{D_{\rho_i}(s)}$$

$$\frac{d\mathcal{B}}{ds} = \tau_B \frac{s |\mathbf{p}_\pi|^3 |\mathbf{p}_D|^3}{24\pi^3 m_B^7} |\mathcal{A}|^2$$

$$|\mathbf{p}_\pi| = \frac{\sqrt{[s - (m_\pi + m_\omega)^2] [s - (m_\pi - m_\omega)^2]}}{2\sqrt{s}},$$

$$|\mathbf{p}_D| = \frac{\sqrt{[m_B^2 - (s + m_D)^2] [m_B^2 - (s - m_D)^2]}}{2\sqrt{s}},$$

$$\sum_{\lambda=0,\pm} \varepsilon^\mu(p, \lambda) \varepsilon^\nu(p, \lambda) = -g^{\mu\nu} + \frac{p^\mu p^\nu}{p^2},$$

$$\sum_{\lambda=0,\pm} |\varepsilon_{\mu\nu\alpha\beta} p_3^\mu \varepsilon^\nu(p_\omega, \lambda) p_\pi^\alpha p^\beta|^2 = s |\mathbf{p}_\pi|^2 |\mathbf{p}_D|^2 (1 - \cos^2 \theta)$$

Six decay modes with large branching ratios (10^{-3})

Decay modes	Units	PQCD	Data [16]
$B^+ \rightarrow \bar{D}^0[\rho(770)^+ \rightarrow]\pi\pi^+$	%	$1.21^{+0.20}_{-0.21}$	1.34 ± 0.18
$B^0 \rightarrow D^-[\rho(770)^+ \rightarrow]\pi\pi^+$	10^{-3}	$7.63^{+1.18}_{-0.96}$	7.6 ± 1.2
$B_s^0 \rightarrow D_s^-[\rho(770)^+ \rightarrow]\pi\pi^+$	10^{-3}	$7.36^{+0.78}_{-0.82}$	6.8 ± 1.4
$B^+ \rightarrow \bar{D}^{*0}[\rho(770)^+ \rightarrow]\pi\pi^+$	10^{-3}	$9.03^{+1.79}_{-1.74}$	9.8 ± 1.7
$B^0 \rightarrow D^{*-}[\rho(770)^+ \rightarrow]\pi\pi^+$	10^{-3}	$8.15^{+1.46}_{-1.45}$	6.8 ± 0.9
$B_s^0 \rightarrow D_s^{*-}[\rho(770)^+ \rightarrow]\pi\pi^+$	10^{-3}	$7.12^{+1.09}_{-1.09}$	9.5 ± 2.0

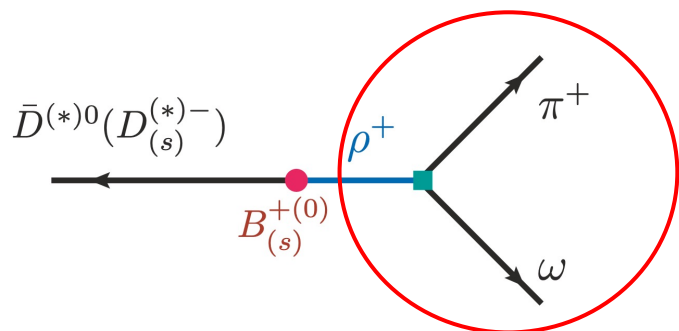
Decay modes	\mathcal{B} (in 10^{-3})
$B^+ \rightarrow \bar{D}^0 [\rho(770)^+ \rightarrow] \omega \pi^+$	$1.42^{+0.16+0.15+0.11+0.10}_{-0.16-0.13-0.09-0.10}$
$B^+ \rightarrow \bar{D}^0 [\rho(1450)^+ \rightarrow] \omega \pi^+$	$0.96^{+0.11+0.09+0.08+0.40}_{-0.11-0.09-0.08-0.40}$
$B^0 \rightarrow D^- [\rho(770)^+ \rightarrow] \omega \pi^+$	$0.80^{+0.06+0.12+0.06+0.07}_{-0.06-0.09-0.02-0.07}$
$B^0 \rightarrow D^- [\rho(1450)^+ \rightarrow] \omega \pi^+$	$0.52^{+0.03+0.06+0.03+0.22}_{-0.03-0.06-0.03-0.22}$
$B_s^0 \rightarrow D_s^- [\rho(770)^+ \rightarrow] \omega \pi^+$	$0.88^{+0.05+0.07+0.00+0.06}_{-0.05-0.07-0.01-0.06}$
$B_s^0 \rightarrow D_s^- [\rho(1450)^+ \rightarrow] \omega \pi^+$	$0.59^{+0.03+0.05+0.00+0.25}_{-0.03-0.04-0.00-0.25}$

Decay modes	\mathcal{B} (in 10^{-3})	Γ_L/Γ
$B^+ \rightarrow \bar{D}^{*0}[\rho(770)^+ \rightarrow]\omega\pi^+$	$1.21^{+0.17+0.09+0.05+0.07}_{-0.17-0.09-0.03-0.07}$	$0.74^{+0.02}_{-0.02}$
$B^+ \rightarrow \bar{D}^{*0}[\rho(1450)^+ \rightarrow]\omega\pi^+$	$0.87^{+0.12+0.07+0.03+0.37}_{-0.12-0.07-0.02-0.37}$	$0.67^{+0.02}_{-0.02}$
$B^0 \rightarrow D^{*-}[\rho(770)^+ \rightarrow]\omega\pi^+$	$1.20^{+0.18+0.09+0.02+0.07}_{-0.18-0.08-0.01-0.07}$	$0.68^{+0.02}_{-0.02}$
$B^0 \rightarrow D^{*-}[\rho(1450)^+ \rightarrow]\omega\pi^+$	$0.89^{+0.13+0.06+0.02+0.38}_{-0.13-0.06-0.02-0.38}$	$0.63^{+0.01}_{-0.01}$
$B_s^0 \rightarrow D_s^{*-}[\rho(770)^+ \rightarrow]\pi\pi^+$	$1.03^{+0.11+0.08+0.00+0.05}_{-0.11-0.08-0.00-0.05}$	$0.65^{+0.01}_{-0.01}$
$B_s^0 \rightarrow D_s^{*-}[\rho(1450)^+ \rightarrow]\pi\pi^+$	$0.77^{+0.08+0.06+0.00+0.32}_{-0.08-0.06-0.00-0.32}$	$0.59^{+0.01}_{-0.01}$

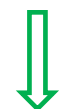
BELLE: PRD92-012013

$$\left\{ \begin{array}{l} \mathcal{B} = (1.48 \pm 0.27^{+0.15+0.21}_{-0.09-0.56}) \times 10^{-3} \\ \mathcal{B} = (1.07^{+0.15+0.06+0.40}_{-0.31-0.13-0.02}) \times 10^{-3} \end{array} \right.$$

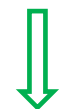
APPROACH:



$$\mathcal{L}_{\rho\omega\pi} = g_{\rho\omega\pi} \epsilon_{\mu\nu\alpha\beta} \partial^\mu \rho^\nu \partial^\alpha \omega^\beta \pi.$$



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$$F_{\omega\pi}(s) = \frac{g_{\rho\omega\pi}}{f_\rho} \sum_{\rho_i} \frac{A_i e^{i\phi_i} m_{\rho_i}^2}{D_{\rho_i}(s)}$$

$$D_{\rho_i}(s) = m_{\rho_i}^2 - s - i\sqrt{s}\Gamma_{\rho_i}(s)$$

one has $A = 1$ and $\phi = 0$ for $\rho(770)$



$$g_{\rho\omega\pi} = 16.0 \pm 2.0 \text{ GeV}^{-1}$$

$$g_{\rho\pi\pi} \approx 6.0$$

$$g_{\rho\omega\pi} = 16.0 \pm 2.0 \text{ GeV}^{-1}$$

$$F_{\omega\pi}(s) = \frac{g_{\rho\omega\pi}}{f_\rho} \sum_{\rho_i} \frac{A_i e^{i\phi_i} m_{\rho_i}^2}{D_{\rho_i}(s)}$$

The weight A_1 in Eq. (2.23) for the subprocess $\rho(1450) \rightarrow \omega\pi$ moves a lot in the literature, it has been measured to be 0.584 ± 0.003 and 0.164 ± 0.003 in [31], 0.175 ± 0.016 , 0.137 ± 0.006 and 0.251 ± 0.006 in [28], 0.26 ± 0.01 and 0.11 ± 0.01 in [27]

[31] : PRD108-092012

[28] : PRD94-112001

[27] : PRD88-054013

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$$A_1 = \frac{g_{\rho(1450)\omega\pi} f_{\rho(1450)} m_{\rho(770)}}{g_{\rho(770)\omega\pi} f_{\rho(770)} m_{\rho(1450)}}$$

$$A_1 = 0.171 \pm 0.036$$

$$f_{\rho(1450)} g_{\rho(1450)\omega\pi} = \sqrt{12\pi f_{\rho(1450)}^2 \mathcal{B}(\rho(1450) \rightarrow \omega\pi) \Gamma_{\rho(1450)} / p_c^3}$$

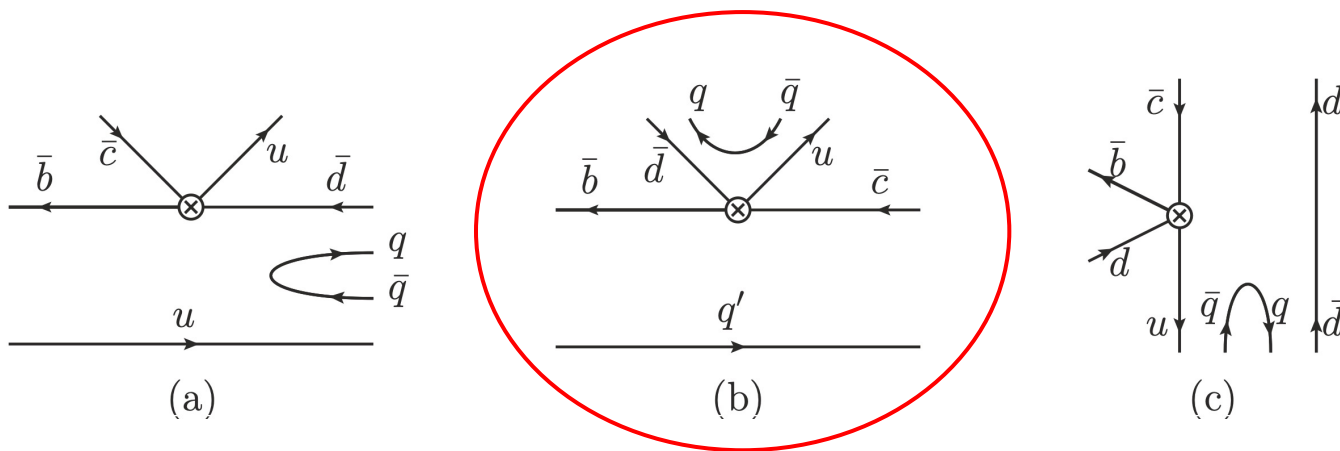
$$f_{\rho(1450)}^2 \mathcal{B}(\rho(1450) \rightarrow \omega\pi) = 0.011 \pm 0.003 \quad \text{PRD64-092001 CLEO}$$

- ❖ Test the factorization hypothesis for B decays with $B^0 \rightarrow D^{(*)-} \omega \pi^+$ and $B_s^0 \rightarrow D_s^{*-} \omega \pi^+$ with the longitudinal polarization Γ_L/Γ [PLB89-105]

pure annihilation decay

$$B^0 \rightarrow D_s^{*-} K^{*+}$$

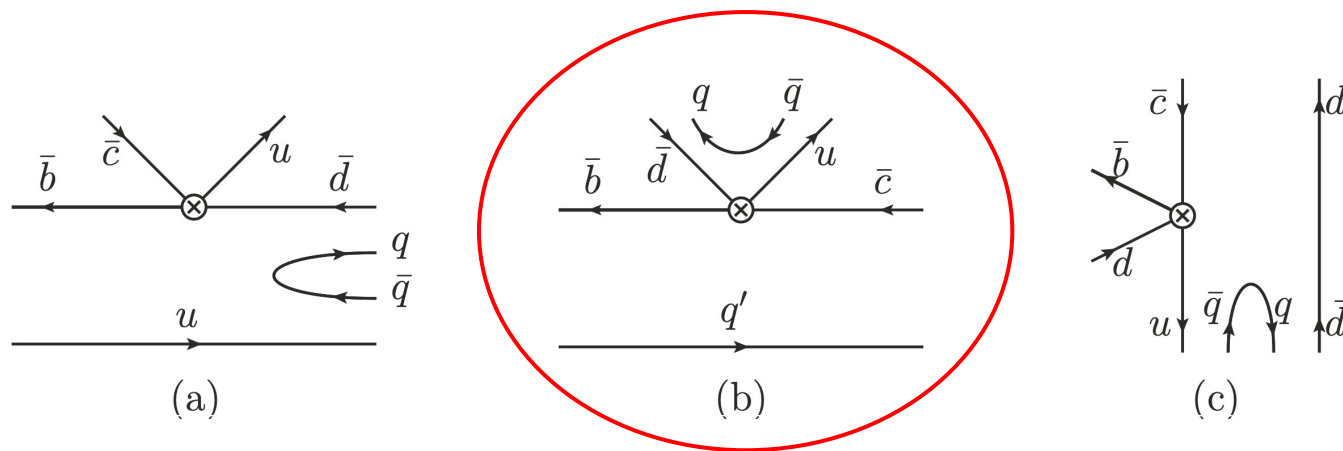
$$\mathcal{B} = (3.2_{-1.3}^{+1.5}) \times 10^{-5}$$

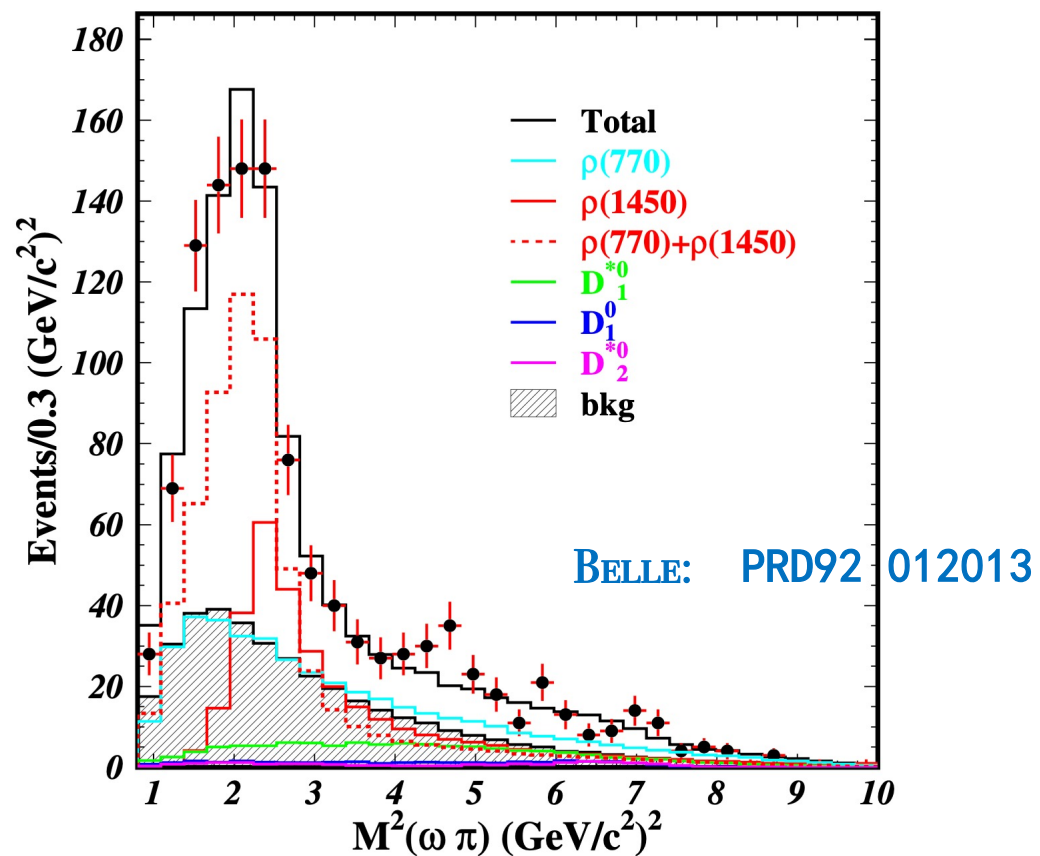
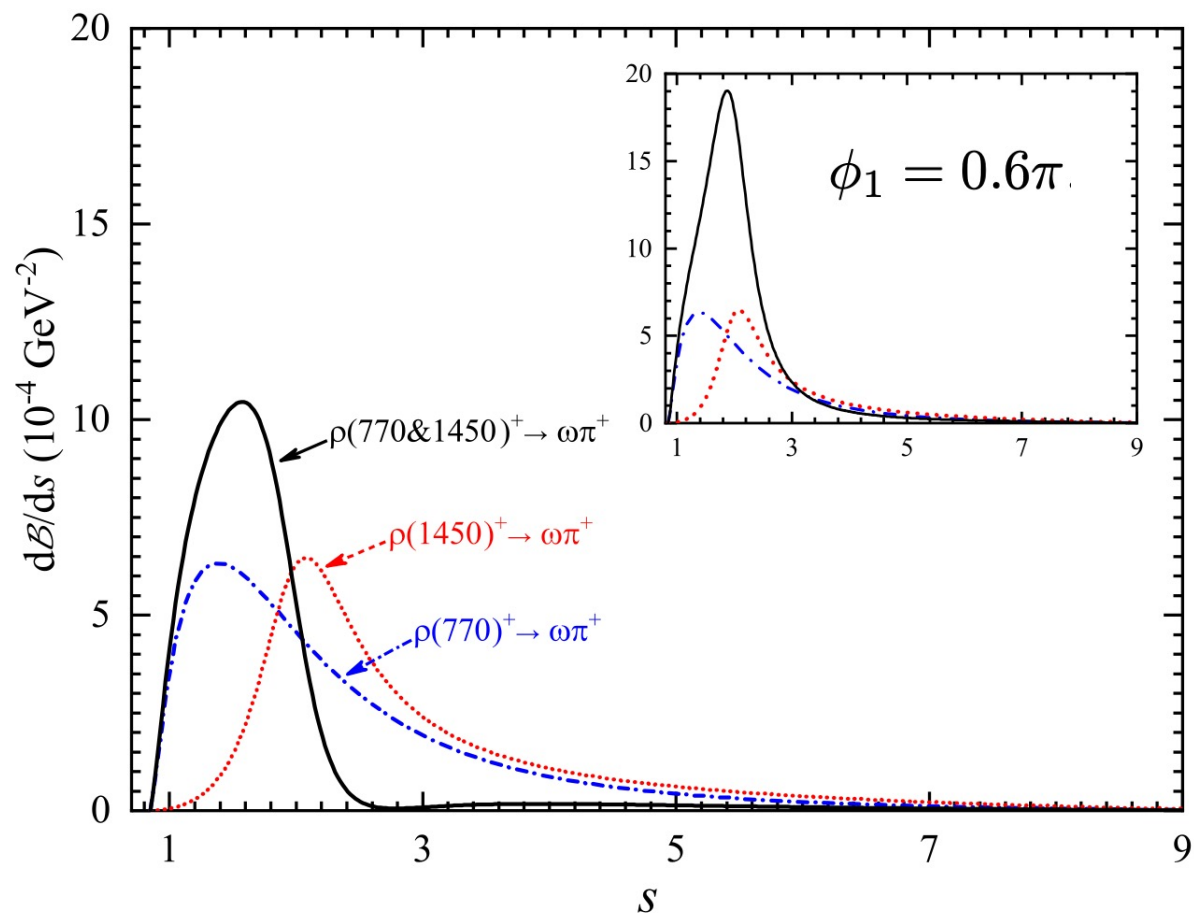


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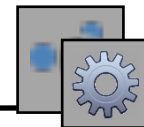
$$B_s^0 \rightarrow D_s^{*-} [\rho(770)^+ \rightarrow] \pi \pi^+$$

$$B_s^0 \rightarrow D_s^{*-} [\rho(1450)^+ \rightarrow] \pi \pi^+$$

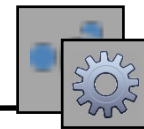




$$B^0 \rightarrow D^{*-} \rho^+ \rightarrow D^{*-} \omega \pi^+$$



- ✓ Introduction of the 3-body B decays / Virtual contributions
- ✓ Contributions of the $\rho(770, 1450) \rightarrow \omega\pi$ in $B \rightarrow \bar{D}^{(*)}\omega\pi$
- ✓ Results and discussions



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Thank You !