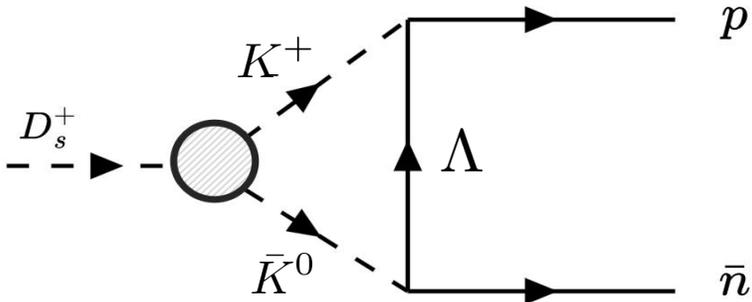




Hidden strangeness in meson weak decays to baryon pair

arXiv: 2409.11374



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Oct. 26, 2024

Research Motivation

1. Long distance contributions at charm scale

A. Calculations of CPV

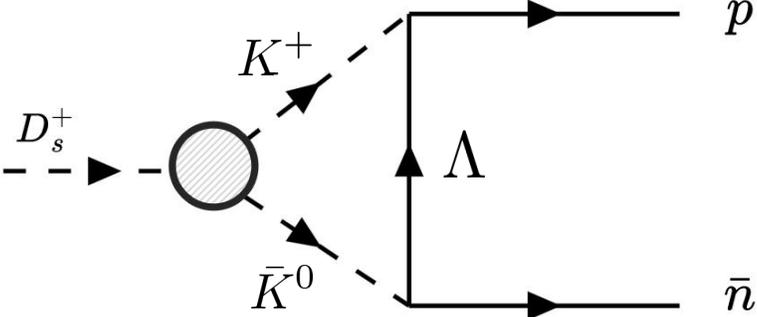


$$a_{CP}^{\pi\pi} = (23.2 \pm 6.1) \times 10^{-4} \quad a_{CP}^{KK} = (7.7 \pm 5.7) \times 10^{-4}$$

PRL 131, 091802 (2023) PRL 122, 211803 (2019)

H. Y. Cheng and C. W. Chiang, Phys.Rev.D109, no.7, 073008(2024)

H. n. Li, C. D. Lu and F. S. Yu, Phys.Rev.D86, 036012(2012)



B. Charm meson mixing

$$D^0 - \bar{D}^0 \text{ mixing} \quad \longrightarrow$$

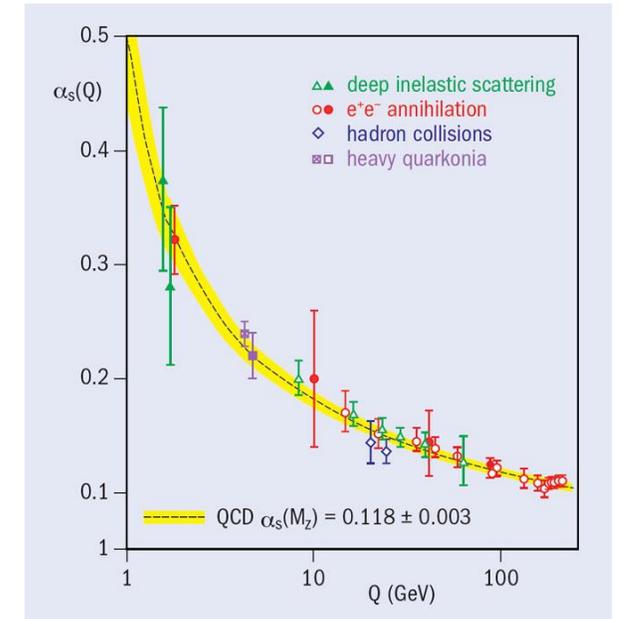
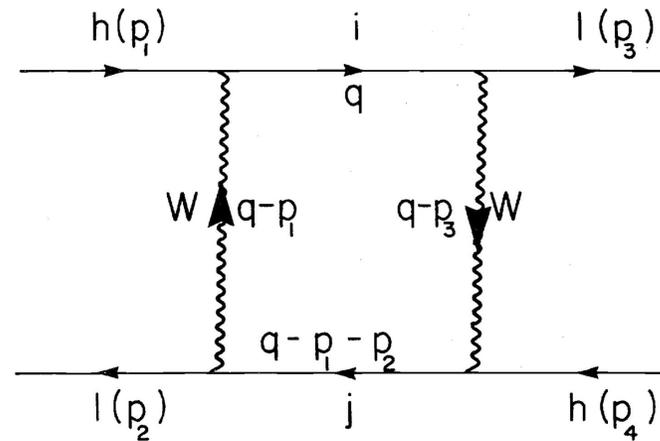
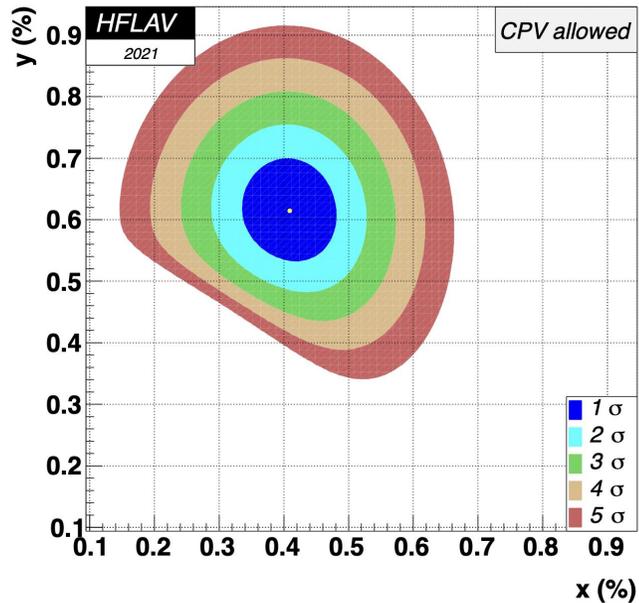
Research Motivation

Global fit : $\frac{\Delta m}{\Gamma} = (0.409^{+0.048}_{-0.049})\%$

$\frac{\Delta\Gamma}{2\Gamma} = (0.615^{+0.056}_{-0.055})\%$

Large cancellation from GIM

Theory from short distance $\sim 10^{-6}$



H. Y. Cheng, Phys. Rev. D 26, 143 (1982)

Peter Zerwas 2004 *Physik Journal* 3 12 31.

[HFLAV], Phys. Rev. D 107, 052008 (2023)

Research Motivation

2. Experimental progress

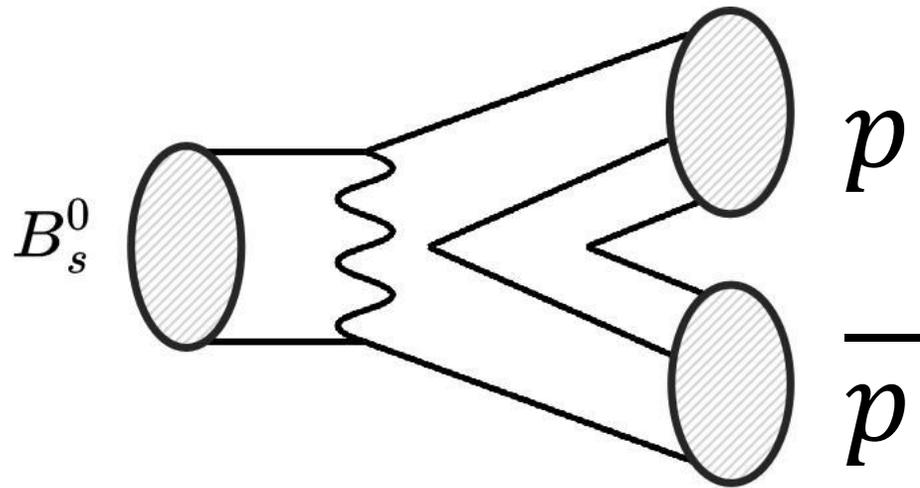


LHCb: $\mathcal{B}(B_s^0 \rightarrow p\bar{p}) < 10^{-10}$

X.N.Jin, C.W.Liu and C.Q.Geng, Phys.Rev.D 105 (2022) 5, 053005

$\mathcal{B}(B_s^0 \rightarrow p\bar{p}) < 5.1 \times 10^{-9}$

R. Aaij et al. [LHCb], Phys. Rev. D 108, no.1, 012007(2023)



Annihilation-diagram \rightarrow suppression from $m_{u,d,s}$

$$\mathcal{A} = \frac{G_F}{\sqrt{2}} V_{ub} V_{us}^* f_{B_s^0} (c_2 + c_1/3) 2m_u \langle p\bar{p} | \bar{u} \gamma_5 u | 0 \rangle$$

$$\approx 0$$

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Both channels are expected to be proportional to $m_{u,d} \dots$

Chiral suppression does not hold in D_s^+ ?

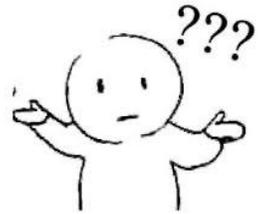
✓ $\mathcal{B}(B_s^0 \rightarrow p \bar{p})$

Data

$$< 5.1 \times 10^{-9}$$

Theory

$$\sim 10^{-10}$$



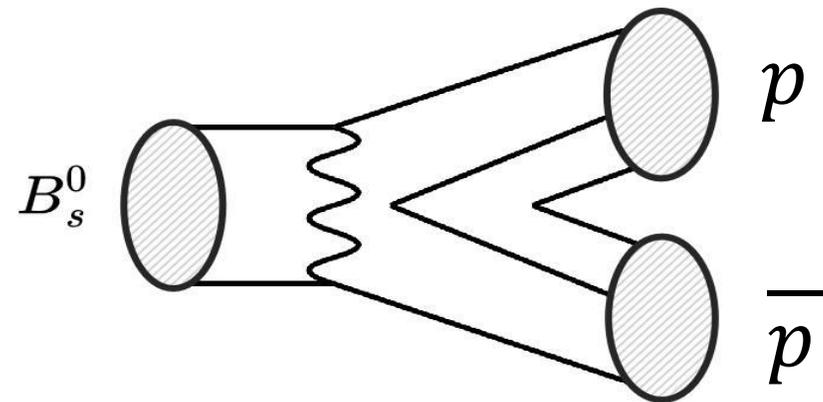
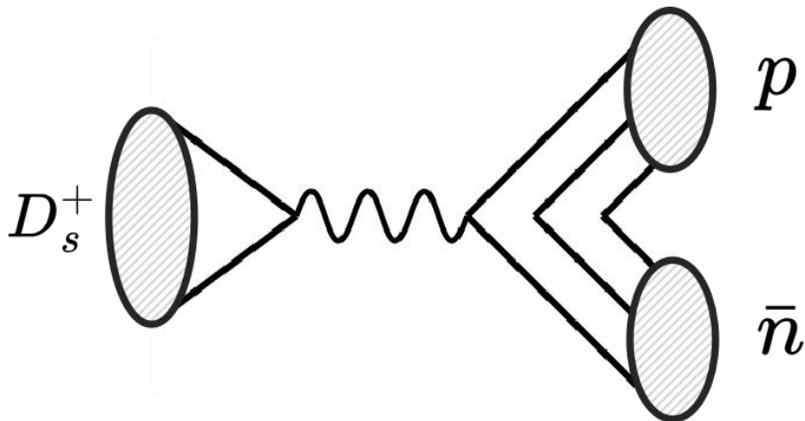
$\mathcal{B}(D_s^+ \rightarrow p \bar{n})$

$$(1.22 \pm 0.11) \times 10^{-3}$$

pdg

$$(0.4_{-0.3}^{+1.1}) \times 10^{-6}$$

Phys. Lett. B **663**, 326 (2008)



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$$D_s^+ \rightarrow p \bar{n}$$

C.H.Chen, H.Y.Cheng and Y.K.Hsiao, Phys. Lett. B **663**, 326-329 (2008)

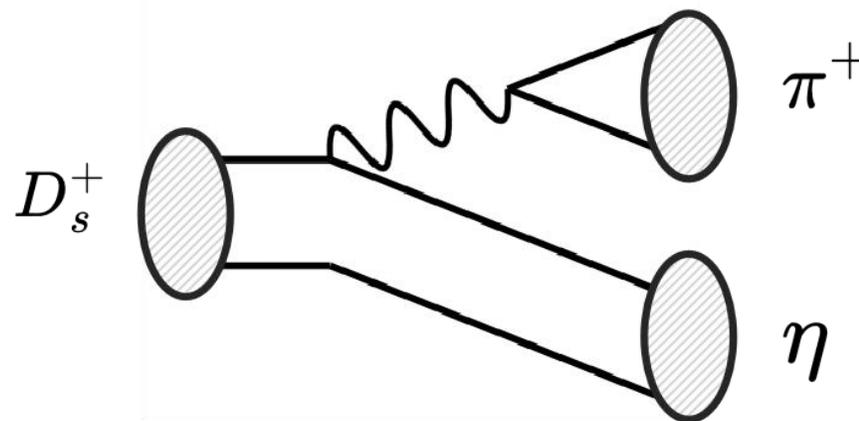
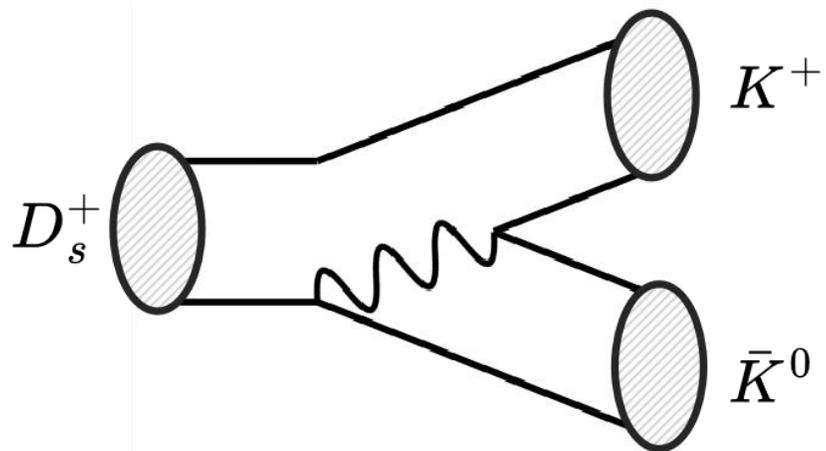
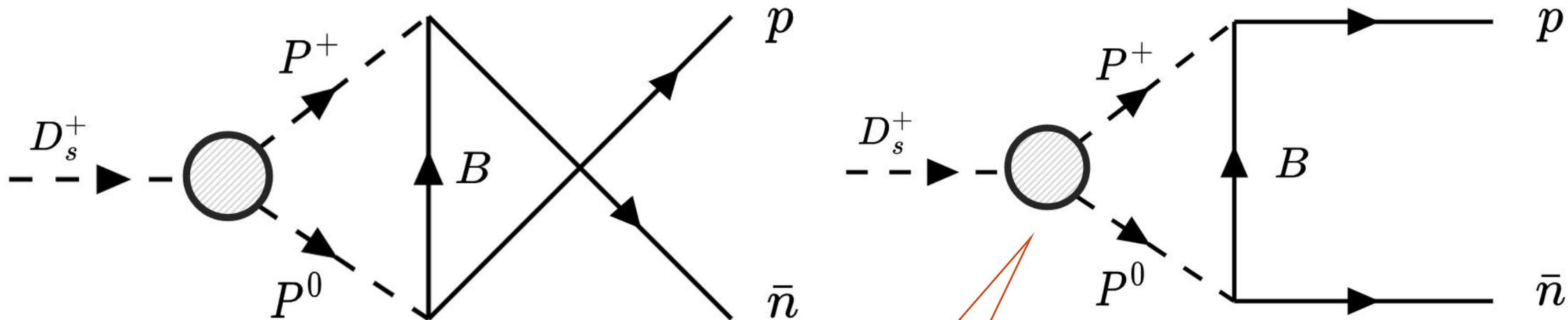
The **only** two-body baryonic decays in D mesons !

Tiny released energy; $m_{D_s} - m_p - m_n \approx m_\pi/2$

Short-distance QCD contributions are **suppressed** by light quark mass

At this energy scale, **the degree of freedom is hadrons** instead of gluons and quarks.

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Hidden strangeness in meson weak decays to baryon pair

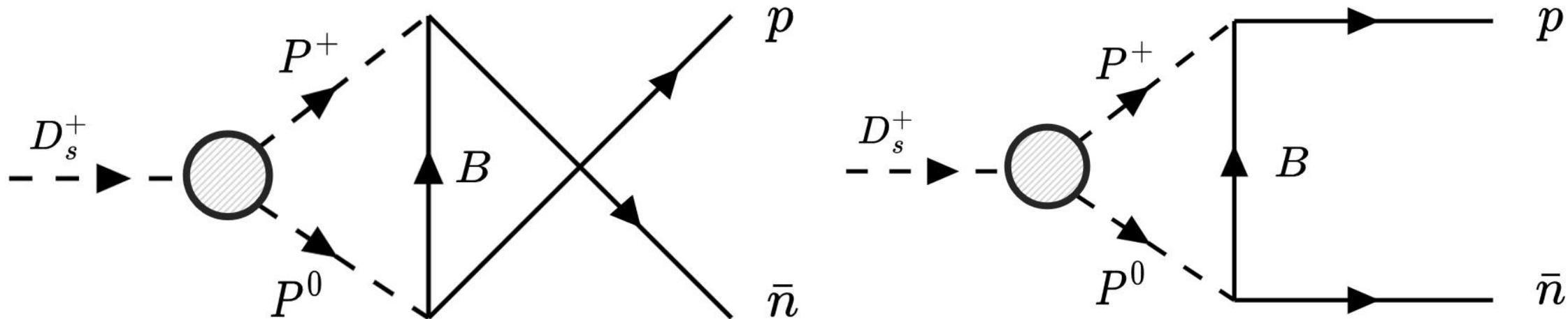


FIG. 2. The LD triangular diagrams for $D_s^+ \rightarrow p \bar{n}$. The blob stands for the insertion of \mathcal{H}_{eff} with the underlying SD transition at quark level. The considered intermediate hadrons are $(P^0, P^+) = (\bar{K}^0, K^+)$ or (η, π^+) , and $B \in (\Lambda, \Sigma^{0,+}, p, n)$.

$$\sum_{P^0, P^+, B} T_{D_s^+ P^0 P^+}^{\text{weak}} g_{P^0 B n} g_{P^+ p B} \int \frac{d^4 q}{(2\pi)^4} \frac{-(q - p_1)^\mu \gamma_\mu + m_2}{(q^2 - m_1^2)[(q - p_1)^2 - m_2^2][(q - p_1 - p_2)^2 - m_3^2]}$$

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$$\sum_{P^0, P^+, B} T_{D_s^+ P^0 P^+}^{\text{weak}} \mathcal{G}_{P^0 B n} \mathcal{G}_{P^+ p B} \int \frac{d^4 q}{(2\pi)^4} \frac{-(q - p_1)^\mu \gamma_\mu + m_2}{(q^2 - m_1^2)[(q - p_1)^2 - m_2^2][(q - p_1 - p_2)^2 - m_3^2]}$$

Weak SD transitions;
PDG branching fractions

Strong couplings;
Goldberger-Treiman relations

$$(T_{D_s^+ \bar{K}^0 K^+}^{\text{weak}}, T_{D_s^+ \eta \pi^+}^{\text{weak}}) = (0.180 \pm 0.02, -0.132 \pm 0.003) G_F \text{GeV}^3$$

$$g_{\bar{K}^0 \Sigma^0 n} = g_{K^+ p \Sigma^0} = \frac{1}{\sqrt{2}} g_{K^+ \Sigma^+ n} = -\frac{1}{\sqrt{2}} g_{\bar{K}^0 \Sigma^+ p} = 3.215 \pm 0.163,$$

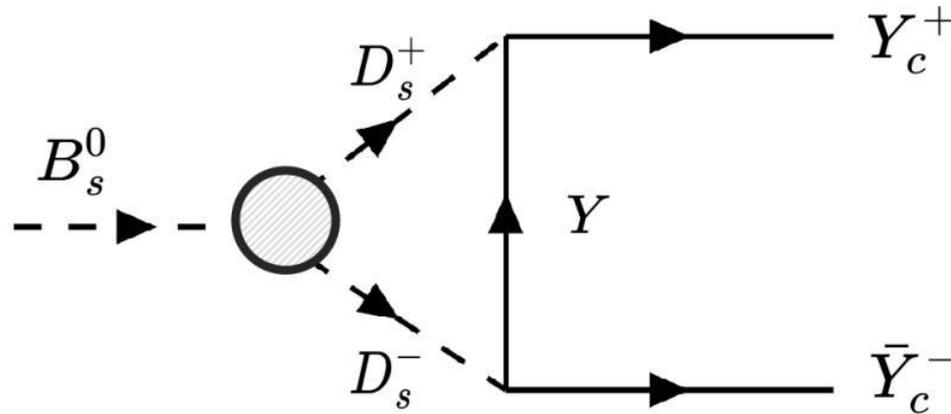
$$g_{\bar{K}^0 \Lambda n} = -g_{K^+ p \Lambda} = 9.228 \pm 0.209, \quad \frac{1}{\sqrt{2}} g_{\pi^+ p n} = 12.897 \pm 0.047.$$

$\mathcal{B}(D_s^+ \rightarrow p \bar{n}) = (1.43 \pm 0.10) \times 10^{-3},$ This work	$\mathcal{B}(D_s^+ \rightarrow p \bar{n}) = (1.30 \pm 0.36_{-0.16}^{+0.12}) \times 10^{-3}$	CLEO
		Phys. Rev. Lett. 100 , 181802 (2008).
	$\mathcal{B}(D_s^+ \rightarrow p \bar{n}) = (1.21 \pm 0.10 \pm 0.05) \times 10^{-3}$	BesIII
		Phys. Rev. D 99 , 031101 (2019).

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CP symmetry leads to a selection rule in partial waves:

$$\mathcal{B}(B_{sH}^0 \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-) = 2\mathcal{B}_P(B_s^0 \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-), \mathcal{B}(B_{sL}^0 \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-) = 2\mathcal{B}_S(B_s^0 \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-).$$



The LD triangular diagram for $B_s^0 \rightarrow Y_c^+ \bar{Y}_c^-$ with $Y = \Lambda, \Sigma$.

pQCD: Rui, Zhou and Zou, Zhi-Tian and Li, Ying, arXiv:2409.16113 [hep-ph]

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$$A_{LD} = T_{B_s^0 D_s^+ D_s^-}^{\text{weak}} (g_{D_s^+ Y_c^+ Y})^2 \int \frac{d^4 q}{(2\pi)^4} \frac{-(q - p_1)^\mu \gamma_\mu + m_2}{(q^2 - m_1^2)[(q - p_1)^2 - m_2^2][(q - p_1 - p_2)^2 - m_3^2]}$$

$\Gamma \propto (g_{D^+ \Lambda_c^+ n})^4$ The coupling is poorly known

$g_{D^+ \Lambda_c^+ n} = 10.7_{-4.3}^{+5.3}$ A. Khodjamirian, C. Klein, T. Mannel and Y. M. Wang, JHEP **09**, 106 (2011).

$\mathcal{B}(B_s^0 \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-) > 4.7 \times 10^{-5}$ this work

$\mathcal{B}(B_s^0 \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-) < 8 \times 10^{-5}$ R. Aaij *et al.* [LHCb], Phys. Rev. Lett. **112**, 202001 (2014)

$$\frac{\Gamma(B^0 \rightarrow \Xi_c^+ \bar{\Xi}_c^-)}{\Gamma(B_s^0 \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-)} = 1.4\% \quad 5.3\% \text{ in the exact SU(3)!}$$

$$m_{B^0} - 2m_{\Xi_c} = 0.34\text{GeV}$$

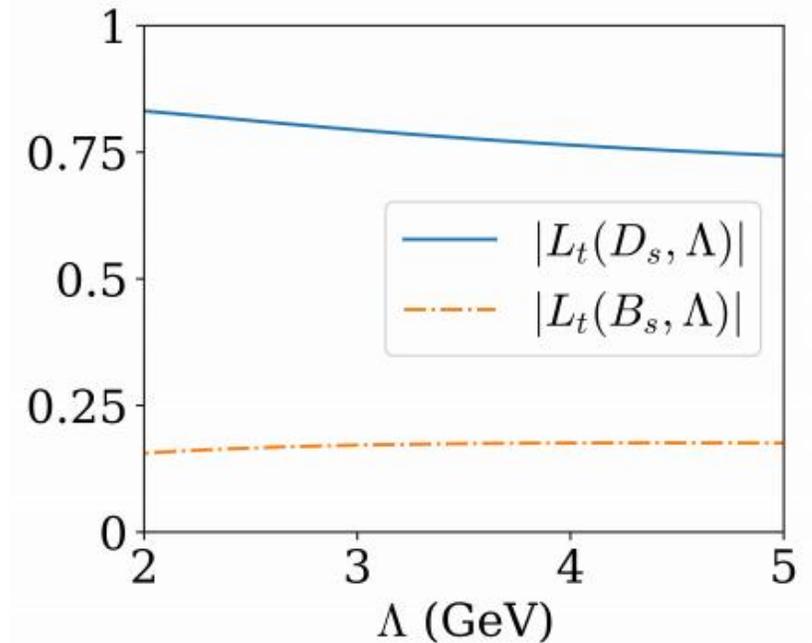
$$m_{B_s^0} - 2m_{\Lambda_c} = 0.76\text{GeV}$$

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- At high energy, the picture of hadron scattering may not be valid.

$$L_t(\Lambda) = \int \frac{d^4q}{(2\pi)^4} \frac{-(q-p_1)^\mu \gamma_\mu + m_2}{(q^2 - m_1^2)[(q-p_1)^2 - m_2^2][(q-p_1-p_2)^2 - m_3^2]} \frac{\Lambda^2}{\Lambda^2 - (q-p_2)^2}$$

- A **cut-off** has been included to study the high energy dependence.
- The **cut-off** dependence of the integrals are tiny! It shows that our numerical results are **stable** and provide an **univocal** prediction.



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TABLE I. Collections of the branching fractions compared to the data [22]. The prediction of $\mathcal{B}_P(D_s^+ \rightarrow p\bar{n})$ is taken from Ref. [7]. Here, $\mathcal{B}_{S,P}$ stand for the branching S and P wave parts of the branching fractions, respectively.

Channel	\mathcal{B}_S	\mathcal{B}_P	\mathcal{B}_{tot}	$\mathcal{B}_{\text{tot}}^{\text{exp}}$
$D_s^+ \rightarrow p\bar{n}$	$(1.43 \pm 0.10) \times 10^{-3}$	$(0.4_{-0.3}^{+1.1}) \times 10^{-6}$	$(1.43 \pm 0.10) \times 10^{-3}$	$(1.22 \pm 0.11) \times 10^{-3}$
$B_s^0 \rightarrow \Lambda_c^+ \bar{\Lambda}_c^-$	$> 4.2 \times 10^{-5}$	$(3.9 \pm 1.5) \times 10^{-6}$	$> 4.7 \times 10^{-5}$	$< 8 \times 10^{-5}$
$B_s^0 \rightarrow \Sigma_c^+ \bar{\Sigma}_c^-$	$< 10^{-6}$	$(3.4 \pm 1.6) \times 10^{-6}$	$(3.9 \pm 2.1) \times 10^{-6}$	—
$B^0 \rightarrow \Xi_c^+ \bar{\Xi}_c^-$	$> 7 \times 10^{-5}$	$(1.3 \pm 0.5) \times 10^{-7}$	$> 7 \times 10^{-5}$	—
$B^0 \rightarrow \Xi_c'^+ \bar{\Xi}_c'^-$	$< 3 \times 10^{-8}$	$(8 \pm 3) \times 10^{-8}$	$(1.0 \pm 0.4) \times 10^{-7}$	—

We are working on CPV in these decays. To be appeared on arXiv soon.

Thank you!

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