

Electromagnetic form factor of proton from the perturbative QCD



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Outline

- Introduction
 - LCDAs
 - Electromagnetic form factors
 - PQCD approach
- Preliminary Result
- Outlook
- Summary

- Baryons play an important role in the evolution of the Universe, big-bang nucleosynthesis, matter-antimatter asymmetry, ...
- However, our knowledge of the baryons are limited, even the simplest baryon proton, such as mass and spin of proton.
- The inner structure of baryon is highly related to perturbative and non-perturbative predictions, understanding of QCD dynamics.



LCDAs

• Light-cone distribution amplitudes (LCDAs) are the fundamental structure of hadrons, describe hadron inner structure in exclusive processes.

LCDAs: exclusive processes



3-parton distribution





V.Braun, R.J.Fries, N.Mahnke, E.Stein, 2001

PDF: inclusive processes



1-parton distribution



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LCDAs

- Light-cone distribution amplitudes (LCDAs) are the fundamental structure of hadrons, describe hadron inner structure in exclusive processes.
- LCDAs are important inputs in the theoretical approaches—QCDF, PQCD, Sum rules, SCET,...
- LCDAs are non-perturbative quantities, thus difficult for predictions.
- Baryon LCDAs are less known. They dominate the theoretical uncertainties.

$$\begin{split} 4 \left\langle 0 \right| \epsilon^{ijk} u_{\alpha}^{i}(a_{1}z) u_{\beta}^{j}(a_{2}z) d_{\gamma}^{k}(a_{3}z) \left| P \right\rangle = \\ &= S_{1} M C_{\alpha\beta} \left(\gamma_{5} N^{+} \right)_{\gamma} + S_{2} M C_{\alpha\beta} \left(\gamma_{5} N^{-} \right)_{\gamma} + P_{1} M \left(\gamma_{5} C \right)_{\alpha\beta} N_{\gamma}^{+} + P_{2} M \left(\gamma_{5} C \right)_{\alpha\beta} N_{\gamma}^{-} \\ &+ V_{1} \left(\mathscr{P} C \right)_{\alpha\beta} \left(\gamma_{5} N^{+} \right)_{\gamma} + V_{2} \left(\mathscr{P} C \right)_{\alpha\beta} \left(\gamma_{5} N^{-} \right)_{\gamma} + \frac{V_{3}}{2} M \left(\gamma_{\perp} C \right)_{\alpha\beta} \left(\gamma^{\perp} \gamma_{5} N^{-} \right)_{\gamma} \\ &+ \frac{V_{4}}{2} M \left(\gamma_{\perp} C \right)_{\alpha\beta} \left(\gamma^{\perp} \gamma_{5} N^{-} \right)_{\gamma} + V_{5} \frac{M^{2}}{2pz} \left(\mathscr{E} C \right)_{\alpha\beta} \left(\gamma_{5} N^{+} \right)_{\gamma} + \frac{M^{2}}{2pz} V_{6} \left(\mathscr{E} C \right)_{\alpha\beta} \left(\gamma_{5} N^{-} \right)_{\gamma} \\ &+ A_{1} \left(\mathscr{P} \gamma_{5} C \right)_{\alpha\beta} N_{\gamma}^{+} + A_{2} \left(\mathscr{P} \gamma_{5} C \right)_{\alpha\beta} N_{\gamma}^{-} + \frac{A_{3}}{2} M \left(\gamma_{\perp} \gamma_{5} C \right)_{\alpha\beta} \left(\gamma^{\perp} N^{+} \right)_{\gamma} \\ &+ \frac{A_{4}}{2} M \left(\gamma_{\perp} \gamma_{5} C \right)_{\alpha\beta} \left(\gamma^{\perp} N^{-} \right)_{\gamma} + A_{5} \frac{M^{2}}{2pz} \left(\mathscr{E} \gamma_{5} C \right)_{\alpha\beta} N_{\gamma}^{+} + \frac{M^{2}}{2pz} A_{6} \left(\mathscr{E} \gamma_{5} C \right)_{\alpha\beta} N_{\gamma}^{-} \\ &+ T_{1} \left(i \sigma_{\perp p} C \right)_{\alpha\beta} \left(\gamma^{\perp} N^{-} \right)_{\gamma} + T_{2} \left(i \sigma_{\perp p} C \right)_{\alpha\beta} \left(\gamma^{\perp} \gamma_{5} N^{-} \right)_{\gamma} + T_{3} \frac{M}{pz} \left(i \sigma_{pz} C \right)_{\alpha\beta} \left(\gamma^{\perp} \gamma_{5} N^{-} \right)_{\gamma} \\ &+ T_{4} \frac{M}{pz} \left(i \sigma_{zp} C \right)_{\alpha\beta} \left(\gamma^{\perp} \gamma_{5} N^{+} \right)_{\gamma} + M \frac{T_{8}}{2} \left(\sigma_{\perp \perp'} C \right)_{\alpha\beta} \left(\sigma^{\perp \perp'} \gamma_{5} N^{-} \right)_{\gamma} , \end{split}$$
Stein, 2001

V.Braun, R.J.Fries, N.Mahnke, E.Stein, 200

proton LCDAs

• Light-cone distribution amplitudes (LCDAs) are the fundamental structure of hadrons, describe hadron inner structure in exclusive processes.

	Twist-3	Twist-4	Twist-5	Twist-6
Vector	V_1	V_2, V_3	V_4, V_5	V_6
Pseudo-vector	A_1	A_2, A_3	A_4, A_5	A_6
Tensor	T_1	T_2, T_3, T_7	T_4, T_5, T_8	T_6
Scalar		S_1	S_2	
Pesudoscalar		P_1	P_2	

Table 1 Twist classification of the proton LCDAs in Eq. (16)

Model	Method	$\begin{array}{c} f_N \cdot 10^3 \\ \text{Gev}^2 \end{array}$	$\lambda_1 \cdot 10^3$ Gev ²	$\lambda_2 \cdot 10^3$ Gev ²	A_1^u	V_1^d	f_1^u	f_1^d	f_2^d	Ref.
	QCDSR	5.0(5)	-27(9)	54(19)						
ASY		-	-	-	0	1/3	1/10	3/10	4/15	
CZ	QCDSR	5.3(5)	-	-	0.47	0.22	-	-	-	[1]
KS	QCDSR	5.1(3)	-	-	0.34	0.24	. – .	-	-	[2]
COZ	QCDSR	5.0(3)	-	-	0.39	0.23	-	-	-	[3]
SB	QCDSR	-	-	-	0.38	0.24	-	-	-	[4]
BK	PQCD	6.64	-	-	0.08	0.31	-	-	-	[5]
BLW	QCDSR	-	-	-	0.38(15)	0.23(3)	0.07(5)	0.40(20)	0.22(5)	[6]
BLW	LCSR (LO)		-	-	0.13	0.30	0.09	0.33	0.25	[6]
ABO1	LCSR (NLO)	-	-	-	0.11	0.30	0.11	0.27	-	[7]
ABO2	LCSR (NLO)				0.11	0.30	0.11	0.29	-	[7]
LAT09	LATTICE	3.23 (63)	-35.57 (65)	70.02 (13)	0.19 (2)	0.20 (1)	-	-	-	[8]
LAT14	LATTICE	3.07 (36)	-38.77 (18)	77.64 (37)	0.07 (4)	0.31 (2)	-	-	-	[9]
LAT19	LATTICE	3.54 (6)	-44.9 (42)	93.4 (48)	0.30 (32)	0.192 (22)	-	-	-	[10]

Slide from Ke-Sheng Huang

Electromagnetic Form factors

- Understanding the internal structure of nucleons from quantum chromodynamics, the underlying theory of strong interaction, is the fundamental challenge faced by the contemporary hadron and nuclear physics.
- Nucleon electromagnetic form factors, which gauge the distributions of the electric charge and magnetization inside the nucleon, are the fundamental probes to the internal structure of nucleon. F. Gross, E. Klempt, S. J. Brodsky, A. J. Buras, et al. Eur. Phys. J. C 83, 1125 (2023)
- Nucleon electromagnetic form factors are among the most fundamental quantities for exploring diverse facets of the non-perturbative QCD dynamics and for advancing our understanding towards the perturbative factorization formalism.

arXiv:2407.18724v1 [hep-ph]

• The nucleon electromagnetic form factors can be defined as follows.

$$\mathcal{M} = \overline{u}(p') \left[F_1(Q^2) \gamma^{\mu} - i \frac{F_2(Q^2)}{2m_p} \sigma^{\mu\nu} q_{\nu} \right] u(p)$$
$$Q^2 \to \infty, F_1(Q^2) \sim O\left(\frac{1}{Q^4}\right), F_2(Q^2) \sim O\left(\frac{1}{Q^6}\right)$$

$$p \longrightarrow p$$

Electromagnetic Form factors



arXiv:2407.18724v1 [hep-ph]

Electromagnetic Form factors



arXiv:2406.19994v1 [hep-ph]

• Based on k_T factorization, the PQCD approach provides a framework applied to hard exclusive processes.

$$\mathcal{A} = \int_0^1 [dx][dx'] \int [d^2 \mathbf{k}_T] \int [d^2 \mathbf{k}_T] \psi_p'(x', \mathbf{k}_T', p', \mu) H(x, x', \mathbf{k}_T, \mathbf{k}_T', \mu) \psi_p(x, \mathbf{k}_T, p, \mu)$$

• PQCD approach is powerful and helpful.

直接CP破坏 (%)	GFA	QCDF	PQCD	Exp.
$B \to \pi^+\pi^-$	-5 ± 3	-6 ± 12	+30 ± 20	+32 ± 4
$B\to K^+\pi^-$	+10 ± 3	+5 ± 9	−17 ± 5	-8.3 ± 0.4

Y.Y.Keum, H.n.Li, A.I.Sanda., et.al., 2001 C.D.Lu, K.Ukai, M.Z.Yang, et.al., 2001

• PQCD approach can calculate high-twist contributions. The high-twist contributions of the proton's LCDAs are significant and cannot be neglected.

Table 4 Form factor $f_1(0)$ from various twist combinations of the Λ_b baryon and proton LCDAs. The first (second) theoretical errors of the total results come from the variations of the relevant parameters in the Λ_b baryon (proton) LCDAs

	Twist-3	Twist-4	Twist-5	Twist-6	Total
Exponential					
Twist-2	0.0007	-0.00007	-0.0005	-0.000003	0.0001
Twist-3 ⁺⁻	-0.0001	0.002	0.0004	-0.000004	0.002
Twist-3 ⁻⁺	-0.0002	0.0060	0.000004	0.00007	0.006
Twist-4	0.01	0.00009	0.25	0.0000007	0.26
Total	0.01	0.008	0.25	0.00007	$0.27 {\pm} 0.09 {\pm} 0.07$

J.J.Han, Y.Li, H.n.Li, Y.L.Shen, Z.J.Xiao, F.S.Yu, Eur.Phys.J.C82,686 (2022)

• The hard-scattering $H(x, x', k_T, k'_T, \mu)$ can be calculated perturbatively. To the lowest order of α_s with two hard exchanged gluons, 42 diagrams can be drawn.



 If only the leading twist is considered, the number of diagrams can be reduced to just 11 by symmetries. However, high twists break these symmetries.





Preliminary result

• This is the result after considering the high twists of the proton's LCDA. The uncertainties stem from the errors in the 8 non-perturbative parameters of the proton's LCDA.



- Within the margin of error, the theoretical and experimental results are in agreement, demonstrating the reliability of the PQCD method.
- The uncertainties in the non-perturbative parameters are relatively large, and the precision needs to be further improved.

Preliminary result

• The high-twist contributions of the proton's LCDA are significant and cannot be neglected.

$Q^4 F_1^p$	twist-3	twist-4	twist-5	twist-6
twist-3	0.0035	-0.0510	0.2151	0.0002
twist-4	0.0010	0.2192	0.0429	0.0047
twist-5	0.2154	-0.1558	0.0877	0.0118
twist-6	-0.0001	0.0049	-0.0119	0.0050

$$Q^2 = 10 GeV^2$$

• This conclusion is consistent with the findings in the $\Lambda_b \rightarrow p$ form factor, suggesting that it may be a general conclusion.

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Twist-3 ⁺⁻	-0.0001	0.002	0.0004	-0.000004	0.002
Twist-3 ⁻⁺	-0.0002	0.0060	0.000004	0.00007	0.006
Twist-4	0.01	0.00009	0.25	0.0000007	0.26
Total	0.01	0.008	0.25	0.00007	$0.27 {\pm} 0.09 {\pm} 0.07$

- We should gain new insights into the internal structure of the proton.
- At not high energy scales, the proton's LCDA is expected to be dominated by high-twist contributions.

Preliminary result



Outlook

- The $ep \rightarrow ep$ process or $e^+e^- \rightarrow p\bar{p}$ process are very clean channels for constraining the proton's LCDA.
- Four non-perturbative parameters have a significant impact on F1.



Outlook

• In future high-precision and high-energy measurements of electron-proton scattering or electron-positron annihilation, we expect to extract non-perturbative parameters with greater precision using the PQCD approach.



- Baryon LCDAs are crucial inputs for baryonic exclusive research in QCD methods, and play an very important role in understanding the inner structure of baryons
- We should gain new insights into the internal structure of the proton. The hightwist contributions of the proton's LCDA are significant and cannot be neglected.
- Determine LCDAs from experimental data in PQCD is a good complement for PDF, also can be a good motivation for STCF, EIC,...

