



# Doubly charmed $\Lambda_c\Lambda_c$ scattering from Lattice QCD

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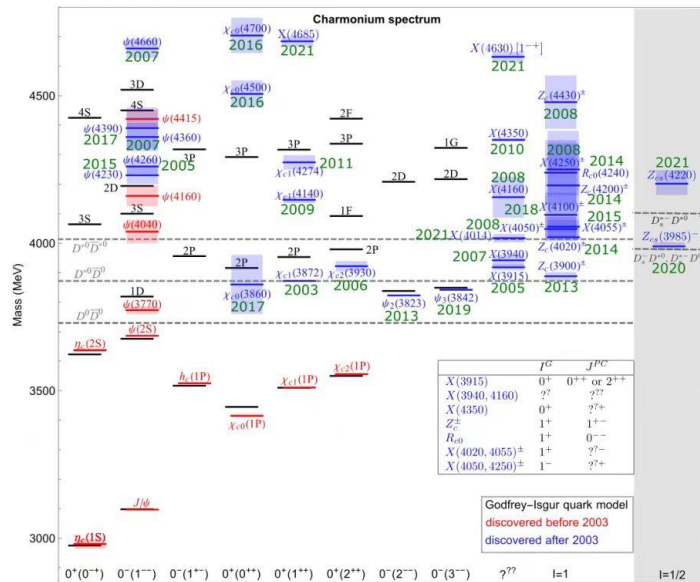
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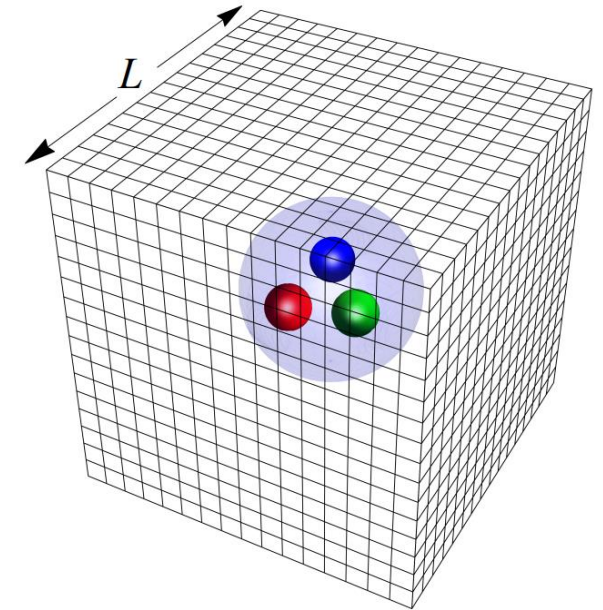
# Outline

- Motivation
- Lattice Setup : CLQCD ensembles
- Formalism : Finite-volume scattering
- Analyses and Result
- Summary

- Since X(3872) was found in 2003, many exotic states beyond conventional quark model were found in experiments one after another.
- Recently, many new exotic states were found, such as  $T_{cc}^+$ (3875),  $P_c$ (4440&4457), X(6900) and so on.

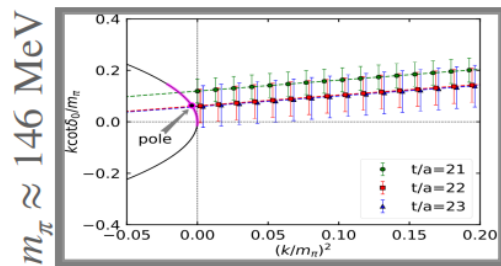


- Unquenched Lattice QCD pushes ahead many lattice study on hadronic spectrum.

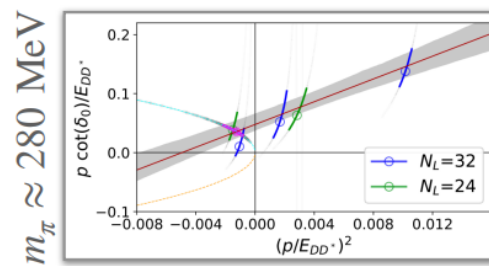


[F.-K. Guo et al, PoS LATTICE2022(2023)232]

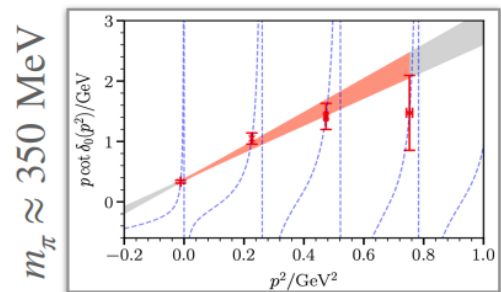
- Many lattice works on  $T_{cc}^+$  (3875) are trying to explain its structure.



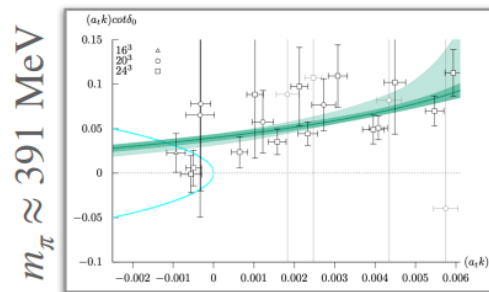
Lyu et al., PRL 131, 161901 (2023)



Padmanath, Prelovsek, PRL 129, 032002 (2022)  
Collins et al., PRD 109 (2024) 9, 094509



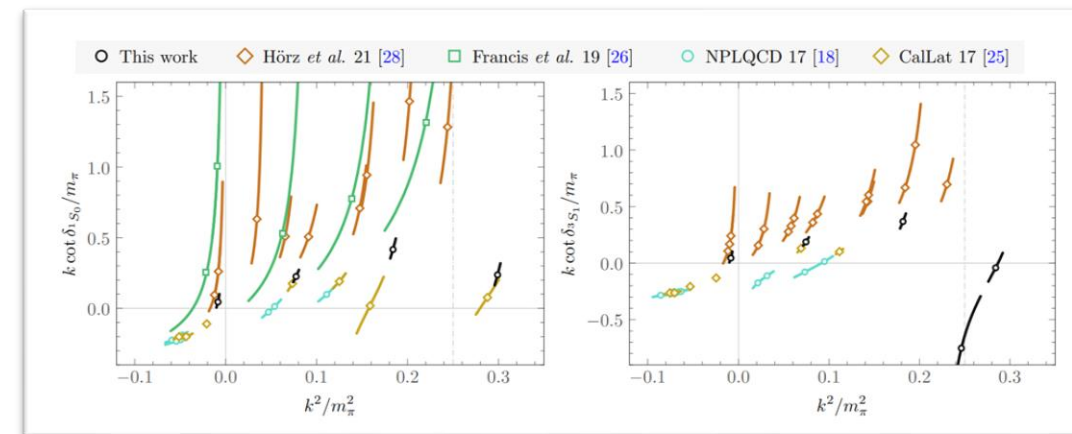
Chen et al. PLB 833, 137391 (2022)



Whyte, Wilson, Thomas, arXiv:2405.15741

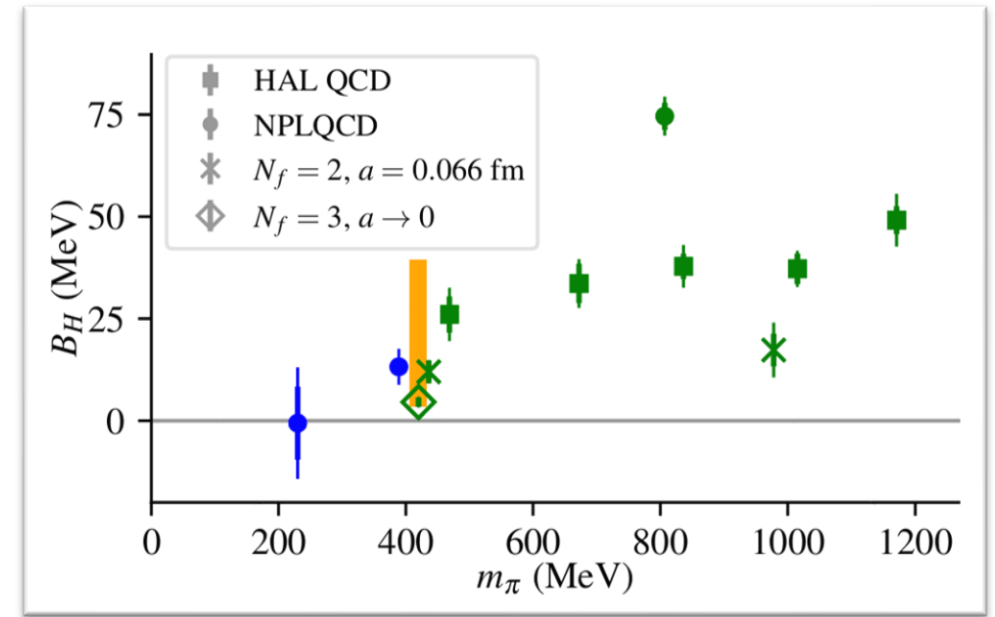
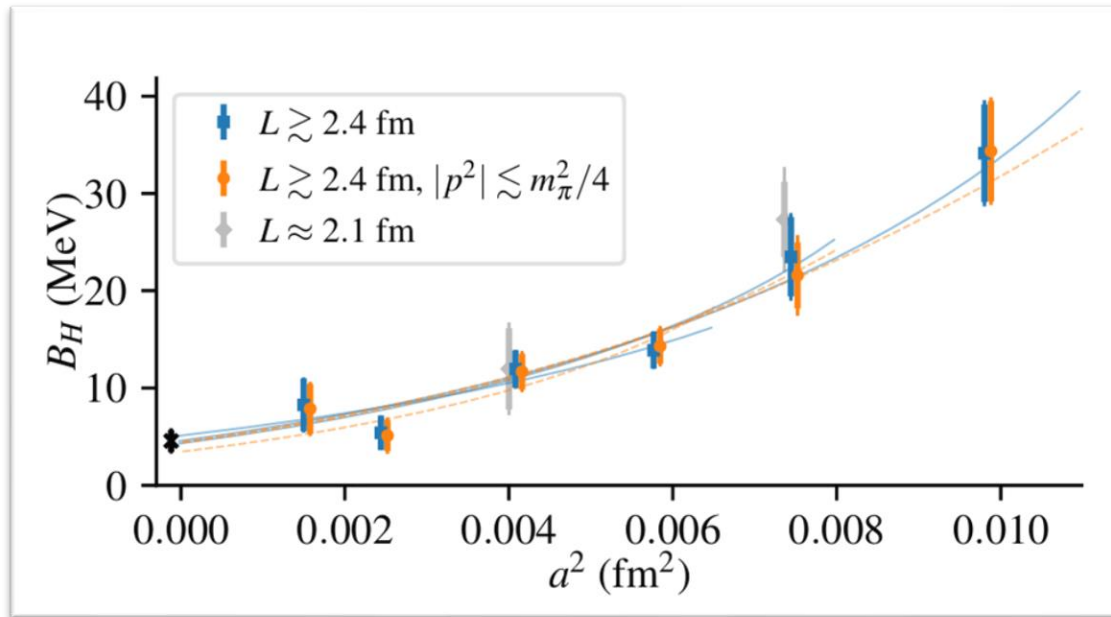
- Study of Baryon-Baryon interaction have many challenges on Lattice.

- Deuteron bound state hasn't been confirmed in the lattice calculation.



[Saman Amarasinghe et al. PRD 107 (2023) 9, 094508]

- A weekly bound state  $\Lambda\Lambda$ . However, it hasn't been found in experiment.

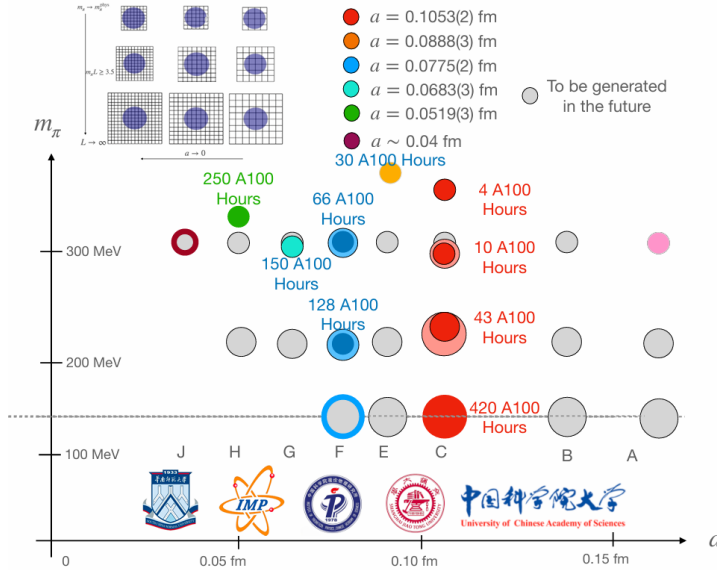


[Jeremy R. Green et al. Phys.Rev.Lett. 127 (2021) 24, 242003]

What about  $\Lambda_c\Lambda_c$ ?



## CLQCD ensembles



In this work,

- Two Wilson-Clover lattice ensembles are used.
- Space-time symmetrical and 2+1 flavor.
- Same  $m_\pi$  and lattice spacing, but different volume.

ensemble	$(L/a)^3 \times T/a$	$\beta$	a(fm)	$m_\pi$ (MeV)	$m_K$ (MeV)	$m_\pi \times L$	$N_{conf}$
F32P30	$32^3 \times 96$	6.41	0.07746(18)	303.2(1.3)	524.6(1.8)	3.81	567
F48P30	$48^3 \times 96$	6.41	0.07746(18)	303.4(0.9)	523.6(1.4)	5.72	201

## □ Operator construction

[S. Prelovsek et al, JHEP 01 (2017) 129]

$$\vec{k} = \frac{2\pi}{L} \vec{n}$$

- Single baryon

$$B(k, x^0) = \sum_{\vec{x}} P_+ \varepsilon_{abc} r_{ax} [s_{bx}^T (C\gamma_5)t_{cx}] e^{-i\vec{k}\cdot\vec{x}}$$

- Two baryons

$$\mathcal{O}_{B_1 B_2}^\Lambda(|\vec{k}|) = \sum_j c_j^\Lambda B_1^T(|\vec{k}_j|) C\gamma_5 B_2(-|\vec{k}_j|)$$

## □ Correlation function

$$C(t) = \langle \mathcal{O}(t) \mathcal{O}^\dagger(0) \rangle \approx A e^{-E t}$$

## □ Generalized eigenvalue problem(GEVP)

$$C(t)v_\alpha(t, t_0) = \lambda_\alpha(t, t_0)C(t_0)v_\alpha(t, t_0)$$

- eliminate the overlap between operators

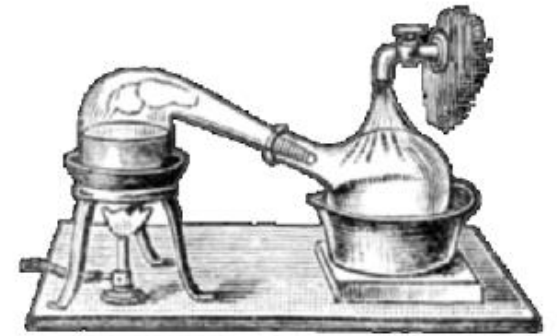
## □ Distillation quark smearing method

➤ Operate  $\square(t)=V(t)V^\dagger(t)$  on quarks  $\rightarrow P = VGV^\dagger$

- $G \sim$  propagator
- $P \sim$  perambulator
- $V \sim$  eigenvectors of  $\nabla^2$
- Dimension of  $V$  is  $[N_t, N_c \times N_x^3, N_{ev}]$

➤ Advantages

- Improve precision
- Efficient for large numbers of operations



[Hadron Spectrum Collaboration, Phys.Rev.D 80 (2009) 054506]

[Colin Morningstar, Phys.Rev.D 83 (2011) 114505]

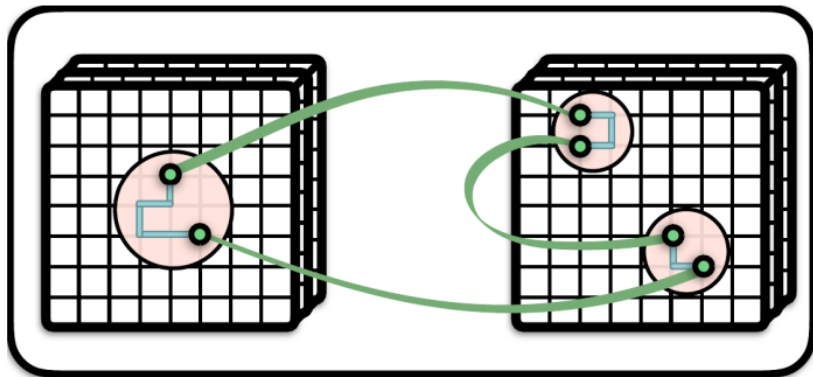


□ Lüscher's finite volume method [M. Lüscher, Nucl. Phys. B. 354 (1991) 531-578]

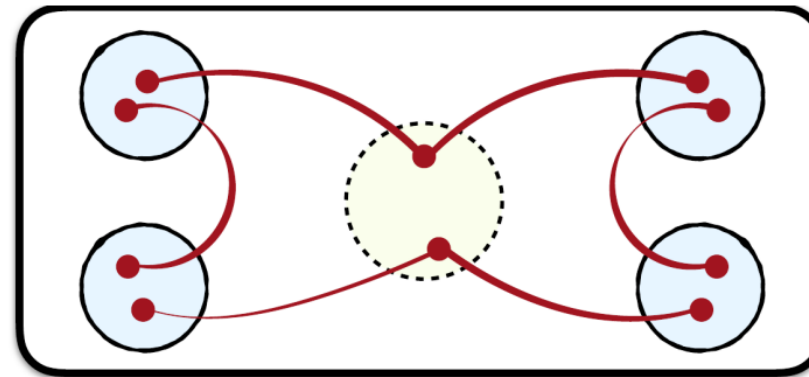
$$\det[1 + ipt(1 + i\mathcal{M})] = 0$$

- $\rho \sim$  phase space
- $t \sim$  scattering amplitude
- $\mathcal{M} \sim$  Lüscher matrix
- For  $S$  wave rest-frame case:

$$\mathcal{M}_{00}(\tilde{q}^2) = \frac{2}{\sqrt{\pi L}} Z_{00}(1; \tilde{q}^2)$$



finite volume spectrum



scattering amplitude

➤ A new modification on Lüscher equation is needed

- Caused by discretized effect with  $Z_{1(2)} \neq 1$  :

$$E(q) = \sqrt{m_1^2 + Z_1 q^2} + \sqrt{m_2^2 + Z_2 q^2}$$

$q$  is the interacting-momentum.  
 $k^* = \frac{2\pi}{L}n$  is center-mass one.

- Consider a translation of the limit

$$\begin{aligned} \lim_{k^* \rightarrow q} \frac{1}{4\omega_1(k^*)\omega_2(k^*)} \frac{q - k^*}{E^*(q) - \omega_1(k^*) - \omega_2(k^*)} &= \lim_{k^* \rightarrow q} \frac{1}{4q[Z_1\omega_2(q) + Z_2\omega_1(q)]} \\ \rightarrow \frac{1}{4\omega_1(k^*)\omega_2(k^*)} \frac{q - k^*}{E^*(q) - \omega_1(k^*) - \omega_2(k^*)} &= \frac{1}{2[Z_1\omega_2(q) + Z_2\omega_1(q)]} \frac{1}{q^2 - k^{*2}} \end{aligned}$$

$$\rightarrow \mathcal{M}_{lm,l'm'}(q, P) = \frac{16\pi^2}{q} \left( \frac{1}{L^3} \sum_k -P \int \frac{d^3k^r}{(2\pi)^3} \right) \mathcal{J}^r \frac{Y_{lm}(\hat{k}^*) Y_{l'm'}(\hat{k}^*) \left(\frac{|k^*|}{q}\right)^{l+l'}}{q^2 - k^{*2}} \frac{E^*(q)}{Z_1\omega_2(q) + Z_2\omega_1(q)}$$

- For  $S$  wave rest-frame case:  $\mathcal{M}_{00}(\tilde{q}^2) = \frac{2}{\sqrt{\pi}L} Z_{00}(1; \tilde{q}^2) \times \frac{E^*(q)}{Z_1\omega_2(q) + Z_2\omega_1(q)}$

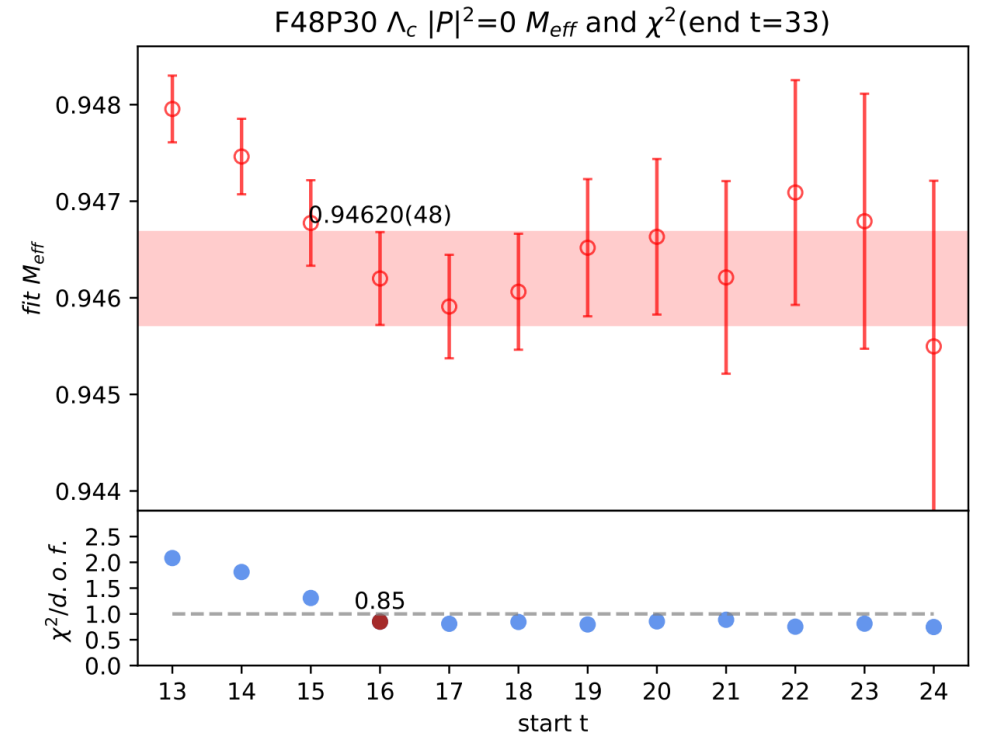
## □ Energy Fitting Method

### ➤ Fitting range selection

Single  $\Lambda_c$  correlation function is fitted with one exponential term:

$$C(t) = A \exp[-M_{eff} t]$$

- Fitting window is  $[t_{start}, t_{end}]$  while  $t_{end}$  is fixed.
- $t_{start}$  should give a stable fit result and a reasonable  $\chi^2$ .
- Fitting window should be as wide as possible.

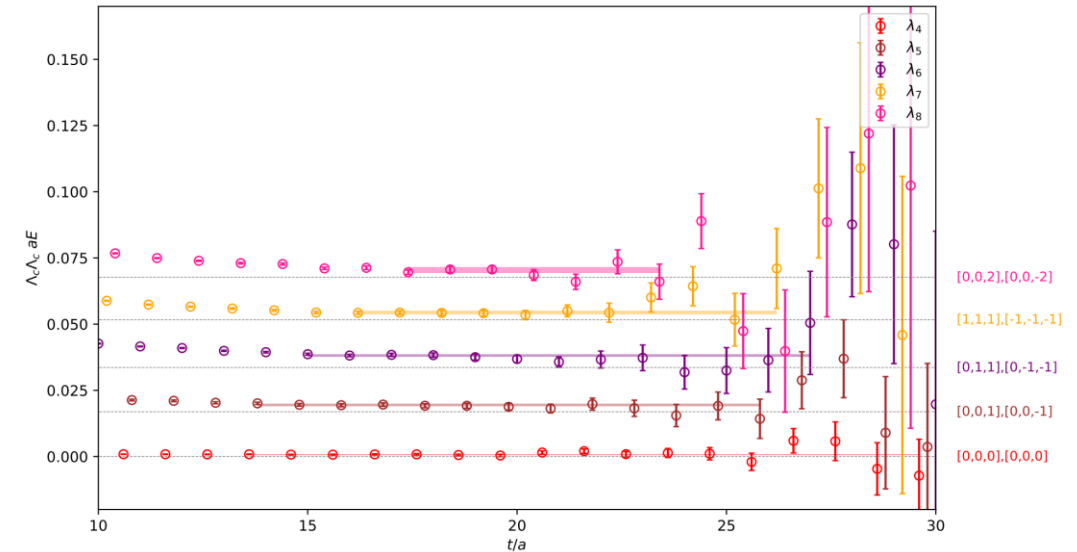
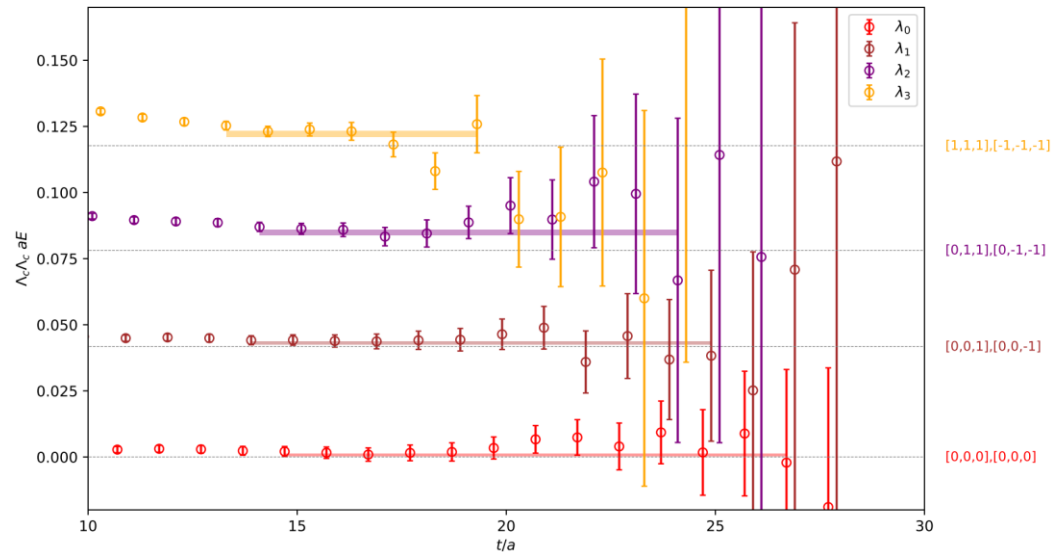


## □ Energy Fitting Method

### ➤ Two baryons correlation function

Since these energy levels are quite close to their corresponding free energies, a ratio form fitting could perform better

$$C_R(t) = \frac{C_{BB}(t)}{C_{thre.}(t)} \simeq Ae^{-(E_{BB}-2m_{\Lambda_c})t} = Ae^{-\Delta Et}$$



[Jozef J. Dudek et al, Phys.Rev.D 82 (2010) 034508]

$V=32^3 \times 96$  &  $48^3 \times 96$   
 $m_\pi \approx 303$  Mev  
 $a=0.07746$  fm

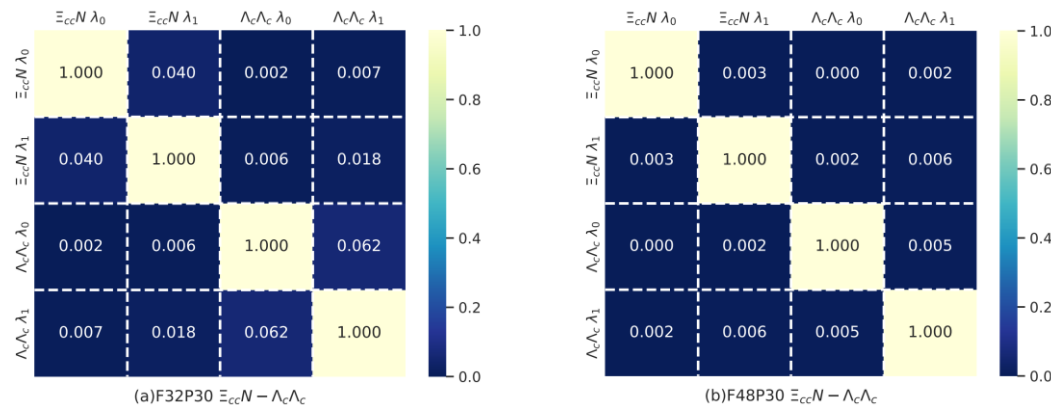
## ➤ Operators

$$\Lambda_c \Lambda_c^{I=0} = [\Lambda_c \Lambda_c]$$

$$\Xi_{cc} N^{I=0} = \frac{1}{2} ([p \Xi_{cc}^+] + [\Xi_{cc}^+ p] - [n \Xi_{cc}^{++}] - [\Xi_{cc}^{++} n])$$

## ➤ Cross-correlation matrix

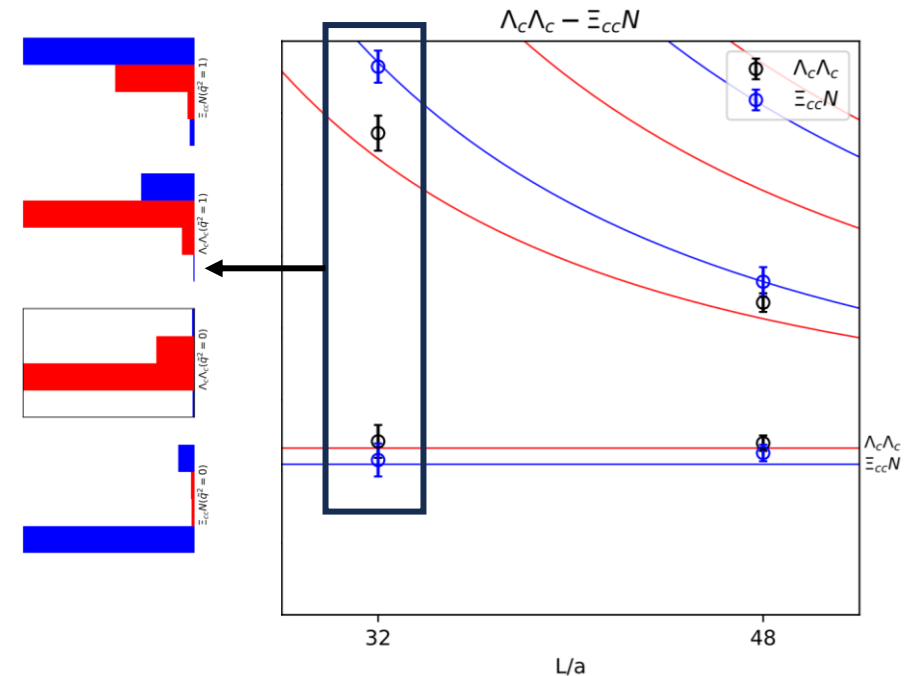
$$\tilde{c}_{ij}(\vec{P}, t) = c_{ij}(\vec{P}, t) / \sqrt{|c_{ii}(\vec{P}, t) c_{jj}(\vec{P}, t)|}$$



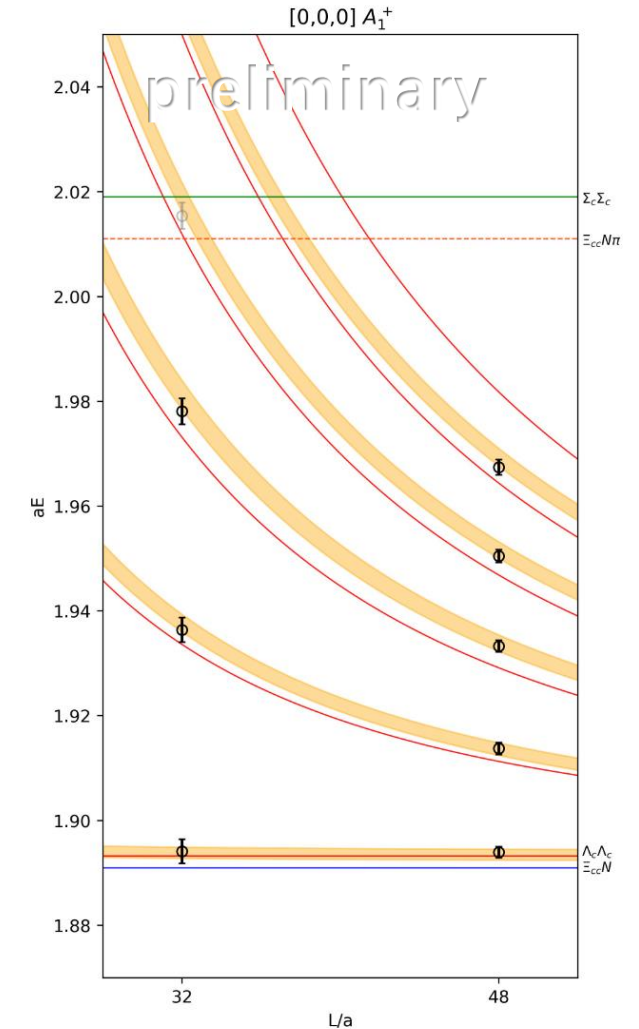
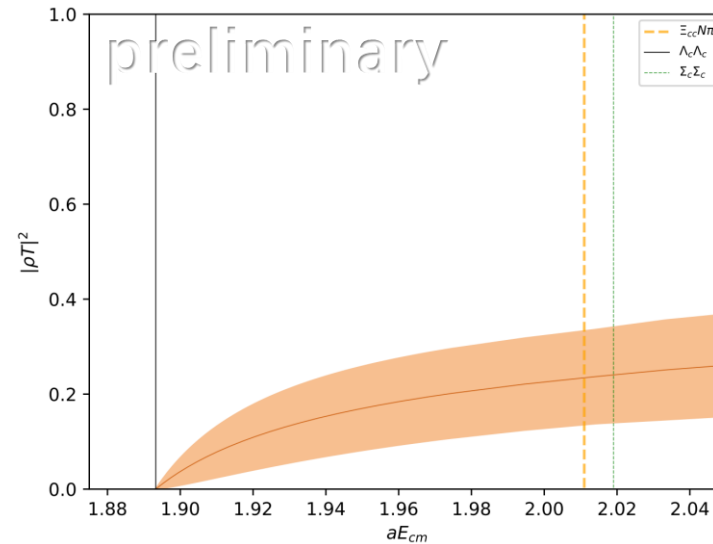
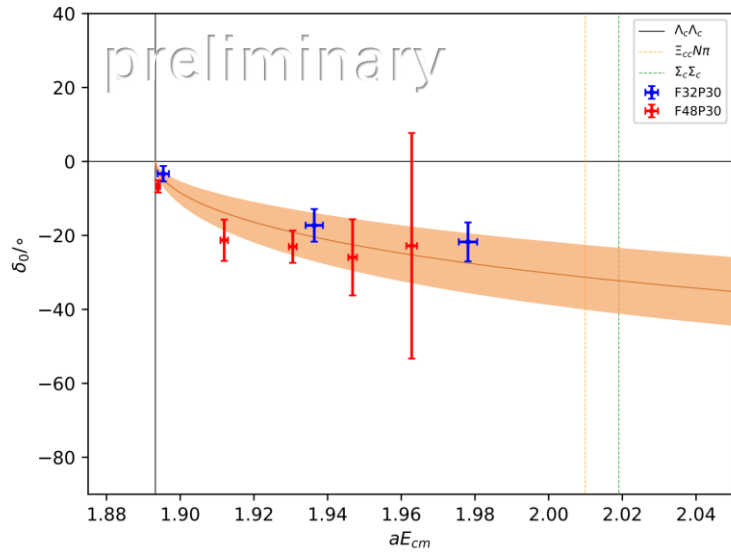
[C. Liu, L. Liu et al, Phys.Rev.D 101 (2020) 5, 054502]

## ➤ Overlap factor

$$Z_i^\alpha = \sqrt{2E_\alpha} e^{E_\alpha t_0/2} v_j^{\alpha*} c_{ji}(t_0)$$



The coupling between  $\Xi_{cc} N$  and  $\Lambda_c \Lambda_c$  is quite small. Single channel  $\Lambda_c \Lambda_c$  is contained latter.



- Repulsive interaction

Parameterization:  $k \cot \delta_0 = \frac{1}{a_0}$ .

Fitting result:  $a_0 = -0.143(49)$  fm, with  $\frac{\chi^2}{dof} = 0.86$ .

Conclusion: **No bound state** in S wave  $\Lambda_c \Lambda_c$  system.

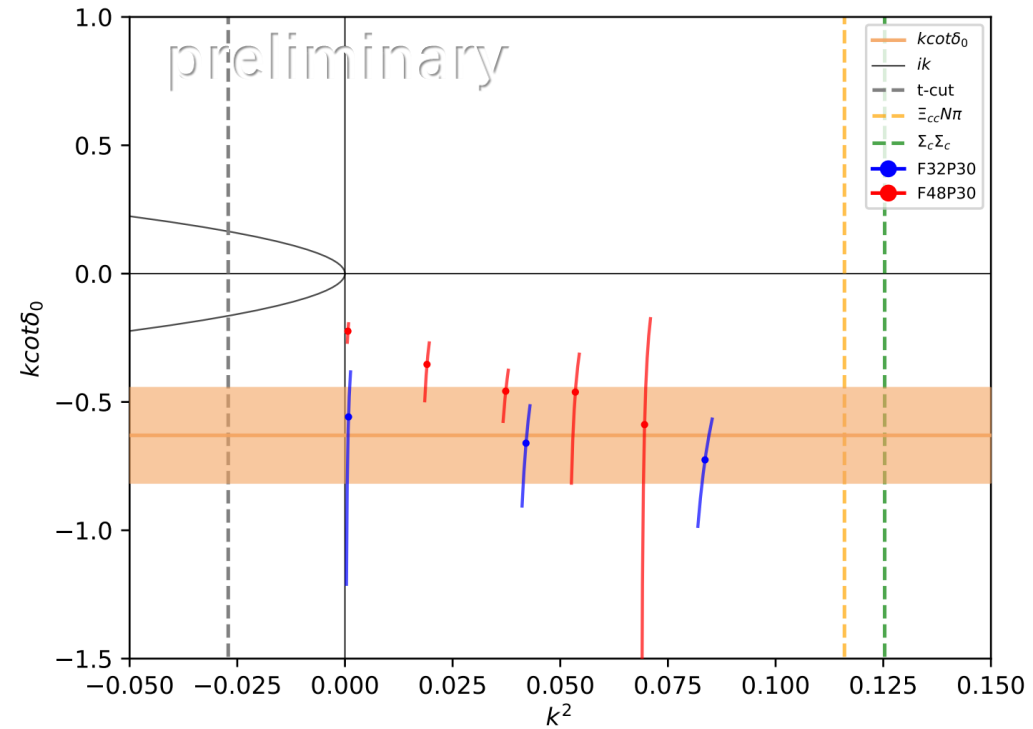
- Due to the **discretization** on lattice, a **modification on Lüscher equation** is proposed in this work.
- First calculation of  $\Lambda_c\Lambda_c$  system. Coupling between  $\Xi_{cc}N$  and  $\Lambda_c\Lambda_c$  is small. Therefore, **single channel** is contained.
- Showing a **repulsive interaction**, there's **no bound state** in this system.
- Scattering length  **$a_0 = -0.143(49)$  fm.**

Thanks!

Backup



● Left-hand cut



The left-hand cut gives the limit of pole extraction with  $k^2 = -\frac{m_\omega^2}{4}$  from the exchanging of omega boson.

[X.-K. Dong et al, Phys.Rev.D 105 (2022) 7, 074506]

[Also see talk in Aug 1<sup>st</sup> 11:30 and 11:50 in LT1]

- Finite volume effect

- The scattering length are consistent in the two ensembles.
- However, there's a little difference.

