Time Dependent Measurements of D⁰ Decays at LHCb

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Content

Motivation

Charm@LHCb

Measurements using 2010 data:

- y_{CP} and A_{Γ}
- wrong sign $D^0 \to K\pi$ decays

Prospects for 2011



Measurements of mixing and (possible) CP parameters in the charm sector provide a complementary way to search for New Physics, as

- they are sensitive to the up sector.
- due to the dominance of the first 2 generations in the SM CP is conserved to first order.

Contributions to CP violation may come from:

- **Mixing:** different rates for $D^0 \to \overline{D^0}$ and $\overline{D^0} \to D^0$
- **Decay:** different amplitudes for a process and its charge conjugate
- Interference between mixing and decay



Time dependent measurements give valuable information to identify the sources for CPV in the charm sector

The $D^0 - \overline{D}^0$ system



Some definitons: $x = \frac{\Delta m_D}{\Gamma_D}, \quad y = \frac{\Delta \Gamma_D}{2\Gamma_D}$ $\lambda_f = \frac{q}{p} \frac{\overline{A_f}}{A_f}, \quad R_m = |\frac{q}{p}|$ $r_f = \frac{\overline{A_f}}{A_f}$

Parametrization of CPV sources:

Direct CPV : $a_f^d = 2r_f \sin q_f^{CP} \frac{x}{2}$ CPV from interference: $a^i = \eta_f^{CP} \frac{x}{2}$ CP violation in mixing: $a^m = -\eta_f^{CP}$ (assuming $x, y, r_f \ll 1$)

$$a_f^a = 2r_f \sin \Phi_f \sin \delta_f$$
$$a^i = \eta_f^{CP} \frac{x}{2} (R_m + R_m^{-1}) \sin \Phi$$

$$a^m = -\eta_f^{CP} \frac{y}{2} (R_m - R_m^{-1}) \cos \Phi$$

Measure time evolution of D^0 decay rates for different final states :

1. DCS $D^0 \to K^+\pi^-$ with respect to CF $D^0 \to K^-\pi^+$ allows to extract $\delta_{K\pi}$: Strong phase

$$x' = x \cos \delta_{K\pi} + y \sin \delta_{K\pi}$$
$$y' = -x \sin \delta_{K\pi} + y \cos \delta_{K\pi}$$

 $\delta_{K\pi}: \text{Strong phase} \\ for D^0 \to K\pi.$

To extract x and y we need a precise measurement of $\delta_{K\pi}$. This can be extracted from quantum correlations in the $D^0 - \overline{D}^0$ system. Valuable input from BES III possible.

2. Time dependent Dalitz plot analysis of $D^0 \to K_s h^+ h^$ gives direct access to x and y.

Experimental access to CPV parameters

1. Compare decay time of CP eigenstates f_{CP} relative to CP non-eigenstates f_{non-CP} :

$$y_{CP} = \frac{\Gamma(D^0 \to f_{CP})}{\Gamma(D^0 \to f_{non-CP})} - 1 = y \cos \Phi - \frac{1}{2} y R_M \sin \Phi$$

A measurement of $y_{CP} \neq y$ signals CP violation! Possible for $D^0 \to KK/\pi\pi$ wrt. $D^0 \to K\pi$ and in time dependent Dalitz analysis of $D^0 \to K_s KK$

2. Compare decay rate evolution of D^0 and \overline{D}^0 , e.g.

$$-A_{\Gamma} = \frac{\Gamma_{D^0 \to f_{CP}} - \Gamma_{\overline{D}^0 \to f_{CP}}}{\Gamma_{D^0 \to f_{CP}} + \Gamma_{\overline{D}^0 \to f_{CP}}} = a_m + a_i$$

- DCS $D^0 \to K^+\pi^-$ wrt. to the cc mode $\overline{D}^0 \to K^-\pi^+$

Experimental status







Within the sensitivity of LHCb there is still a lot of room for both, direct and indirect CPV...





The LHCb collaboration





- few times $10^8 D^0 \rightarrow K\pi$ decays on tape
- use slow pion in $D^* \! \to D^0 \pi_{\sc s}$ for tagging

LHCb key features for charm measurements:

- Excellent vertex resolution and proper time resolution
 - IP resolution: 38µm@1GeV
 - decay time resolution: ~50fs (12% of D⁰ decay time)
- Very good particle identification via two RICH detectors
 - clean separation of π & K over broad momentum range
- Dedicated exclusive trigger lines for charm
 - 1kHz of charm events written to tape







- ${\sim}10^7~D^0 \to K\pi$ decays on tape
- use slow pion in $D^* \to D^0 \pi_{\sc s}$ for tagging

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 D^0 -vertex

h



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We use Charm produced at the proton-proton primary vertex for our analysis.

But: ~50% of B-hadrons decay into charm (so called secondary charm)

- misidentification of charm from B-decays as prompt charm results in a significant bias of the lifetime measurement.
- physical BG: Can not be distinguished from invariant mass distribution

Use simultanoues fit of lifetime and Impact Parameter (IP_{D^0}) to discriminate prompt charm and charm from B-decays:





Knowledge of trigger & selection acceptance is crucial for all lifetime-dependent measurements:

- can be determined on an event-per-event basis by the so called swimming method (initially developed by CDF*)
- LHCb offline software allows to rerun the trigger for different configurations

Elegant implementation of 'swimming' with improved performance possible

Measure acceptance as a function of lifetime by iterative procedure:



- 1. Moving PV along D0 momentum
- 2. Evaluation of trigger & selection decision

* J. Rademacker, Nucl. Instrum. Meth. A **570** (2007) 525, T. Aaltonen **et al** [CDF Collaboration], Phys.Rev.**D 83**(2011)032008





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Lifetime acceptance cross-check: Determination of lifetime for $D^0 \to K\pi$

Determine D⁰ lifetime as proof of principle:

- Acceptance well modelled by swimming method.
- 220k $D^0 \rightarrow K\pi$ events.
- Mean value consistent with PDG!
- Error is statistical only.

2010 data $\tau = 410.3 \pm 0.9 \text{ fs}$ $(\tau_{PDG} = 410.1 \pm 1.5 \text{ fs})$ Entries 10⁴ LHCb Preliminary L= 28 pb⁻¹ 10 data 10² prompt secondary 10 2 Proper Time [ps]

Time dependent observables for $D^0 \to hh$

1. Ratio of DCS (wrong sign, WS) over CF (right sign, RS) decays in $D^0\to K\pi$

$$r(t) = \frac{\Gamma(D^0 \to K^+ \pi^-)}{\Gamma(D^0 \to K^- \pi^+)} = R_D + \sqrt{R_D} y' \Gamma t + \frac{x'^2 + y'^2}{4} t^2$$

$$k_D : \text{DCS decay rate relative to CF}$$

$$x' = x \cos \delta_{K\pi} + y \sin \delta_{K\pi}$$

$$y' = -x \sin \delta_{K\pi} + y \cos \delta_{K\pi}$$

2.
$$y_{CP} = \frac{\Gamma(D^0 \to KK)}{\Gamma(D^0 \to K\pi)} - 1 = y \cos \Phi - \frac{1}{2} y R_M \sin \Phi$$

3.
$$A_{\Gamma} = \frac{\Gamma(D^0 \to KK) - \Gamma(\overline{D^0} \to KK)}{\Gamma(D^0 \to KK) + \Gamma(\overline{D^0} \to KK)} = \frac{1}{2} R_M y \cos \Phi - x \sin \Phi$$

Nota bene:

All equations assume |x| << 1, |y| << 1 and absence of direct CP-violation

Signal yield for $D^0 \to K\pi$

- Signal yield is extracted from an unbinned maximum likelihood fit. in the $(m_{D^0}, \Delta m)$ -plane
- WS decays need due to higher BG stronger selection criteria.
- Identical selection for WS and RS decays.



The time integrated WS/RS ratio from 2010 data

Due to limited statistics in 2010 data we measure only time integrated WS/RS ratio:

$$R = \int_0^\infty r(t)dt = R_D + \sqrt{R_D}y' + \frac{x'^2 + y'^2}{4}$$

relimination	Y	
rest		WS/RS of $D \to K\pi$ decays (%)
$R_{\rm meas}$		$0.442 \pm 0.033 \text{ (stat.)} \pm 0.042 \text{ (sys.)}$
$R_{\rm corr}^{\star}$		$0.409 \pm 0.031 \text{ (stat.) } \pm 0.039 \text{ (sys.)}$
R(PI	DG)	0.380 ± 0.018

*A correction for the non-flat lifetime acceptance has been applied.

Systematic error sources:

- 1. fit model: 0,035%
- 2. size of signal box: 0,023%
- 3. double mis-id of K and π : negligible





Analysis for y_{CP} and A_{Γ} very similar, based on direct lifetime measurements:

- Selection of events in signal region of two-dimensional $(m_{D0}, \Delta m)$ plane.
- Combinatorial BG is estimated from side bands to $\mathcal{O}(1\%)$ for $D^0 \to K\pi$ and $\mathcal{O}(3\%)$ for $D^0 \to KK$.
- We use a maximum likelihood fit for extraction of lifetime.



Result for y_{CP} with 2010 data

 y_{CP} measures the ratio of lifetimes between CP eigenstates $D^0\to K\pi$ and CP-mixed states $D^0\to\pi\pi$.



Budget for systematic error:

- Contribution from Combinatorical BG: ±0.08%
- Fit model: ± 0.08%
- Impact of secondary charm: ±0.39%

Measurement of A_{Γ} with 2010 data

 A_{Γ} measures the lifetime difference between D^0 and $\overline{D^0}$ into CP-eigenstates K^+K^-

- Mistag rate determined from sidebands in $\Delta m(D^* D^0(hh))$
- Systematic error mainly from combinatorical and secondary charm BG



Error budget for A_{Γ}

Effect	$A_{\Gamma} (10^{-3})$
VELO length scale	negligible
Turning point bias	negligible
Turning point scaling	± 0.1
Combinatorial background	± 1.3
Proper time resolution	± 0.1
Minimum proper-time cut	± 0.1
Maximum proper-time cut	± 0.2
Secondary charm background	± 1.6
Total	± 2.1

Summary of 2010 results



October 3rd 2011



Prospects for 2011

In 2011 we increased already now the amount of data by a factor 30! (i.e. already now the world highest statistics for many channels)

Signal yield per pb⁻¹ approximately the same or even improved (e.g. $3x10^5$ untagged D0 \rightarrow K π decays per pb⁻¹)

Trigger settings have been improved And more channels have been added To the trigger..

Modelling of combinatorical BG and BG from B-decays will be improved.

Larger data sample + improved trigger allow to reduce systematic error significantly.



But there is much more to come

Higher statistics also allows to tackle new channels and new analysis techniques with high sensitivity to seek for new physics.

Studies are going on to

• exploit more channels to measure mixing parameters

 $D^* \to D^0(K\pi\pi\pi)\pi_s$ \neg needs 2011 data for sensitive measurement $D^* \to D^0(K^+\mu^-\overline{\nu}_\mu)\pi_s$ \neg needs 2011&2012 data for sensitive measurement

• search for CPV using T-odd correlations in $D^0 \to K^+ K^- \pi^+ \pi^$ construct T-odd observables like $C_T^{\pm} = \vec{p}_{K^{\pm}} \cdot (\vec{p}_{\pi^+} \times \vec{p}_{\pi^-})$

these observables allow to construct T-violating asymmetries between C_T^+ and C_T^- . I.I.Bigi hep/ph-0107102 (2001)

• perform a time dependent Dalitz Plot analysis of $D^0 \to K_S hh$ see next slides

Time dependent Dalitz analysis of $D^0 \to K_s hh$

- Presence of many resonant modes result in a rich Dalitz structure for $D^0 \to K_s hh$.
- Decay amplitudes for $|\mathcal{D}^{0} > \text{and } |\overline{\mathcal{D}}^{0} > \text{are given by:}$ $\mathcal{M}(m_{-}^{2}, m_{+}^{2}, t) = \mathcal{A}(m_{-}^{2}, m_{+}^{2}) \frac{e_{H}(t) + e_{L}(t)}{2} + \frac{q}{p} \overline{\mathcal{A}}(m_{-}^{2}, m_{+}^{2}) \frac{e_{H}(t) - e_{L}(t)}{2}$ $\overline{\mathcal{M}}(m_{-}^{2}, m_{+}^{2}, t) = \overline{\mathcal{A}}(m_{-}^{2}, m_{+}^{2}) \frac{e_{H}(t) + e_{L}(t)}{2} + \frac{p}{q} \mathcal{A}(m_{-}^{2}, m_{+}^{2}) \frac{e_{H}(t) - e_{L}(t)}{2}$

sum over resonant and n.r. decay amplitudes

- Squaring of \mathcal{M} and $\overline{\mathcal{M}}$ give decay rates $\exp(-\Gamma t) \cos(x\Gamma t)$, $\exp(-\Gamma t) \sin(x\Gamma t)$ and $\exp(-(1 \pm y)\Gamma t)$
 - time dependent Dalitz analysis gives direct access to x and yand CPV parameters! ... contrary to WS decays in $D^0 \to K\pi$ (see there)

 $e_{H,L}(t) = \exp[-i(m_{H,L} - i\Gamma_{H,L}/2)t]$

• Sensitivity stems mostly from regions in Dalitz plots where CF and DCS amplitudes from CP eigenstates interfere.

Prospects for $D^0 \to K_s hh$

- LHCb was running in 2010 without a dedicated trigger for D⁰->K_shh
 - despite that a nice signal in 35pb⁻¹ is seen!
- In 2011 LHCb is running with a dedicated trigger for that channel.

This offers LHCb a method to measure x, y and CPV parameters complementary to the decay channel $D^0 \rightarrow hh!$





- With the analysis of data from 2010 LHCb has proven to be able to perform 38pb⁻¹ competitive measurements in the field of mixing CPV in the charm sector.
- ◆ Already now, LHCb has collected ~30 x more data than in 2010.
- Trigger settings have been improved, new channels have been added in the high level trigger.
- High statistics available and the very good performance of LHCb allows us to improve the sensitivity for our existing measurements and to exploit new channels.



We expect many charming results from the 2011 data with unprecedented sensitivity to New Physics in the charm sector.