

PROSPECT FOR MEASUREMENT OF CP-VIOLATION PHASES WITH B_s DECAYS AT FUTURE Z FACTORIES

[EPJC 84 (2024) 859]

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Date: 10/25/24

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CEPC Workshop 2024



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INTRODUCTION

- CKM parameter:
 - In neutral B meson decays to a final state the interference between the amplitude for the direct decay and the amplitude for decay after oscillation, leads to a time-dependent CP-violating asymmetry between the decay time distributions of B and anti-B mesons.
 - $\phi_s = -\arg(-V_{ts}V_{tb}^*/V_{cs}V_{cb}^*)$
 - $\gamma = \arg(V_{ud}V_{ub}^*/V_{cd}V_{cb}^*)$
 - Contributions from physics beyond the SM could lead to much larger values of ϕ_s , insensitive to γ .
- B_s decay parameters:
 - $\Delta\Gamma_s \equiv \Gamma_L - \Gamma_H, \Gamma_s \equiv (\Gamma_L + \Gamma_H)/2$.
 - Able to be calculated with heavy quark expansion (HEQ) theory.

ϕ_s MEASUREMENTS WITH $B_s \rightarrow$
 $J/\psi\phi$

Measurement of ϕ_s ($\Delta\Gamma$, Γ_s) in experiments

Extract the observables $\phi_s, \Gamma_s, \Delta\Gamma_s$ from the time dependent angular distribution.

$$\frac{d^4\Gamma(B_s \rightarrow J/\psi\phi)}{dt d\Omega} \propto \sum_{k=1}^{10} h_k(t) f_k(\Omega),$$

where

$$h_k(t|B_s) = N_k e^{-\Gamma_s t} \left[a_k \cosh\left(\frac{1}{2}\Delta\Gamma_s t\right) + b_k \sinh\left(\frac{1}{2}\Delta\Gamma_s t\right) + c_k \cos(\Delta m_s t) + d_k \sin(\Delta m_s t) \right]$$

$$h_k(t|\bar{B}_s) = N_k e^{-\Gamma_s t} \left[a_k \cosh\left(\frac{1}{2}\Delta\Gamma_s t\right) + b_k \sinh\left(\frac{1}{2}\Delta\Gamma_s t\right) - c_k \cos(\Delta m_s t) - d_k \sin(\Delta m_s t) \right]$$

$f_k(\Omega)$: amplitude function.

$$b_k \sim \pm |\lambda| \cos(\phi_s), \quad d_k \sim \pm |\lambda| \sin(\phi_s)$$

Projection to future Z-factories

$$\sigma(\phi_s) \propto 1/\sqrt{N_{\text{eff}}}$$

- $N_{\text{eff}} \propto N_{b\bar{b}}$
- $N_{\text{eff}} \propto \text{Efficiency}$
- $N_{\text{eff}} \propto \text{Tagging power}$
- $\sigma_{\phi_s} \propto 1/e^{-\frac{1}{2}\Delta m_s^2 \sigma_t^2}$

Define:

$$\xi = 1/ \left(\sqrt{N_{b\bar{b}} \times \varepsilon \times Br} \times \sqrt{p} \times \exp\left(-\frac{1}{2}\Delta m_s^2 \sigma_t^2\right) \right)$$

Then: $\sigma(\phi_s, FE) = \xi_{FE} \times \frac{\sigma(\phi_s, EE)}{\xi_{EE}}$

ξ for LHCb Run2 and LHCb on HL-LHC

Numbers are quoted from *Eur.Phys.J.C79(2019)706*

$$\xi = 1 / \left(\sqrt{N_{b\bar{b}} \times \varepsilon \times Br} \times \sqrt{p} \times \exp\left(-\frac{1}{2} \Delta m_s^2 \sigma_t^2\right) \right)$$

- $N_{b\bar{b}} \times \varepsilon \times Br = 11700$. Avoid considering the efficiency on LHCb.
 - $\mathcal{L}_{\text{int}} = 1.9 \text{ fb}^{-1}$, $b\bar{b}$ cross-section: $144 \mu\text{b}$, $Br = 20\% \times 0.001 \times 0.06 \times 0.5$.
 - $\varepsilon = 7\%$, where the $b\bar{b}$ is already in the acceptance, reasonable estimation.
- Tagging power $p = 4.73\%$.
- Decay time resolution: 45.5 fs.

ξ :

- $\xi_{\text{LHCb}} = 0.018, \sigma(\phi_s, \text{LHCb}) = 0.041 \text{ rad}$.
- $\xi_{\text{LHCb}} = 0.0014, \sigma(\phi_s, \text{HL-LHCb}) = \xi_{\text{HL-LHCb}} \times \sigma(\phi_s, \text{lhcb}) / \xi_{\text{LHCb}} = 3.3 \text{ mrad}$
(HL-LHC: $N_{\text{HL-LHCb}} = N_{\text{LHCb}} \times \frac{300 \text{ fb}^{-1}}{1.9 \text{ fb}^{-1}}$)

$$\xi = 1 / \left(\sqrt{N_{b\bar{b}} \times \varepsilon \times Br} \times \sqrt{p} \times \exp\left(-\frac{1}{2} \Delta m_s^2 \sigma_t^2\right) \right)$$

- Tera-Z: 0.152×10^{12} , 10-Tera-Z: 1.52×10^{12}
- $Br = 20\% \times 0.001 \times 0.06 \times 0.5 \times 2$. (J/ψ can also be reconstructed from e^+e^- on CEPC)

$$\xi = 1 / \left(\sqrt{N_{b\bar{b}} \times \epsilon \times Br} \times \sqrt{p} \times \exp\left(-\frac{1}{2} \Delta m_s^2 \sigma_t^2\right) \right)$$

Reconstruction:

- Assume that we can distinguish $b\bar{b}$ events from other events.
- Assume that we have perfect ability to distinguish leptons with hadrons.
- ϕ candidates: $1.017 - 1.023 \text{ GeV}/c^2$, two hadron tracks.
- J/ψ candidates: $3.07 - 3.14 \text{ GeV}/c^2$, two lepton tracks.
- B_s^0 candidates: $5.28 - 5.46 \text{ GeV}/c^2$, combination of all J/ψ ϕ candidates.

Extraction of ϕ_s require a clean background.

The number of background events are 1.7×10^5 times larger than the number of signal events.

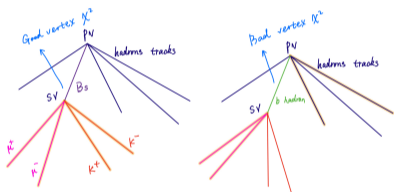
In pure background (from simulation):

- The probability to find a J/ψ candidate is 0.4%.
- The probability to find a ϕ candidate is 3.6%.
- The probability to get a B_s^0 candidate from $J/\psi \phi$ combination is 4.6%.
- Total: 6.7×10^{-6} .

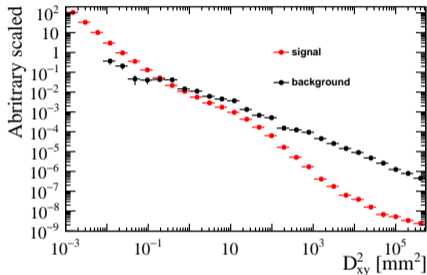
The background is of same magnitude with the signal.

ξ for CEPC (Efficiency)

Vertex χ^2 : reject background.



Signal χ^2 distribution:



- Background: very large spread χ^2 distribution.
- $\chi^2 < 0.1$ keeps 95% of the signal and reject 99.2% of the background.

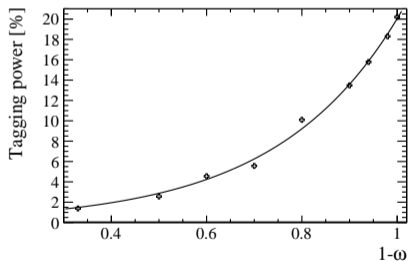
$\varepsilon = 75\%$ with 1% background level.

ξ for CEPC (Tagging power)

$$\xi = 1 / \left(\sqrt{N_{b\bar{b}} \times \varepsilon \times Br} \times \sqrt{p} \times \exp\left(-\frac{1}{2} \Delta m_s^2 \sigma_t^2\right) \right)$$

20% of the tagging power can be easily achieved with a naive algorithm and with assumption of perfect pid. (Same side + Opposite side algorithm)

Dependence on particle identification:



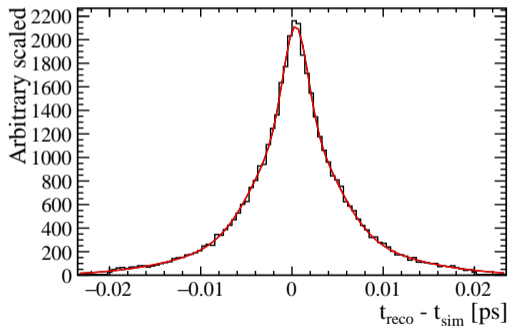
Correctly identification rate: $1 - \omega$,
misidentification probability: $\omega/2$

Considering the particle identification from the detector simulation, the tagging power is:

- Intrinsic tagging power (without considering the effects from the readout electronics): 19.1%.
- Realistic/conservative tagging power (if the particle identification resolution is degraded by 30% with respect to the intrinsic case): 17.4%.

ξ for CEPC (Time resolution)

$$\xi = 1 / \left(\sqrt{N_{b\bar{b}} \times \varepsilon \times Br} \times \sqrt{p} \times \exp\left(-\frac{1}{2} \Delta m_s^2 \sigma_t^2\right) \right)$$



Obtained from detector simulation.

Proper decay time: $t = \frac{m l_{xy}}{p_T}$

Fit with sum of three gaussian.

$$\sigma_{\text{eff}} = \sqrt{-\frac{2}{\Delta m_s^2} \ln\left(\sum_i f_i e^{-\frac{1}{2} \sigma_i^2 \Delta m_s^2}\right)} = 4.7 \text{ fs.}$$

(Reminder LHCb: 45 fs)

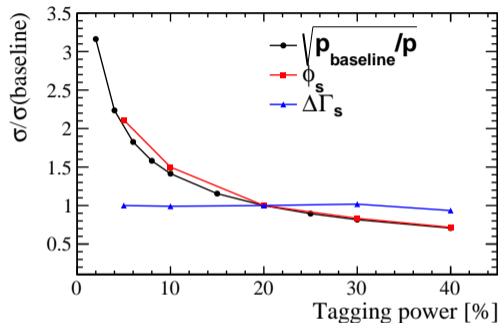
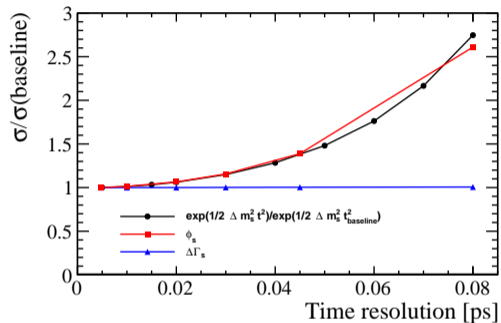
The excellent time resolution benefits from the precise vertex reconstruction and large energy of B_s .

ξ for CEPC (Summary)

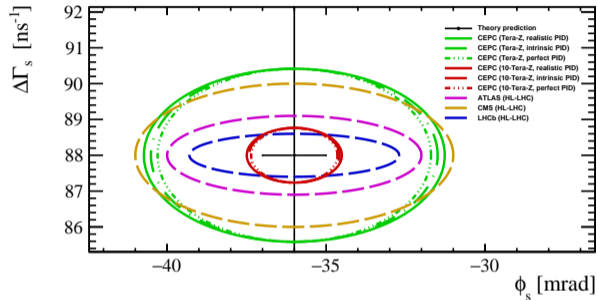
Putting all the components together: $\xi_{\text{CEPC}} = 0.0019$ (Tera-Z), $\sigma(\phi_s, \text{CEPC}) = 4.3\text{mrad}$.

	LHCb(HL-LHC)	CEPC(Tera-Z)	CEPC/LHCb
$b\bar{b}$ statics	43.2×10^{12}	0.152×10^{12}	1/284
Acceptance \times efficiency	7%	75%	10.7
Br	6×10^{-6}	12×10^{-6}	2
Flavour tagging (perfect pid)	4.7%	20%	4.3
Time resolution ($\exp(-\frac{1}{2}\Delta m_s^2 \sigma_t^2)$)	0.52	1	1.92
scaling factor ξ	0.0014	0.0019	0.8
$\sigma(\phi_s)$	3.3 mrad	4.3 mrad	
Flavour tagging (realistic/conservative pid)	4.7%	17.3%	3.7
$\sigma(\phi_s)$	3.3 mrad	4.6 mrad	

Impact from time resolution and flavour tagging



- Time resolution and tagging power dependence for observables.
- ϕ_s resolution has potential to be improved with better tagging power.
- $\Delta\Gamma_s$ (and also Γ_s) has weak dependence: lose the factor of 4.3×1.92 .



- Black point: SM global fit (CKMfitter group/UTfit collaboration) + HQE
([Proc.Int.Sch.Phys.Fermi 137\(1998\)329](#), [Adv.Ser.Direct.High Energy Phys.15\(1998\)239](#))
prediction.
- ATLAS and CMS results are from [ATL-PHYS-PUB-2018-041](#) and [CMS-PAS-FTR-18-041](#).

PENGUIN CONTROL FOR β_s

If the penguin diagram is considered in the B_s decay, the relation between ϕ_s and β_s should be corrected as

$$\phi_s = -2\beta_s + \Delta\phi_s(a, \theta). \quad (1)$$

The shift $\Delta\phi_s$ could be expressed as

$$\tan(\Delta\phi_s) = \frac{2\epsilon a \cos \theta \sin \gamma + \epsilon^2 a^2 \sin(2\gamma)}{1 + 2\epsilon a \cos \theta \cos \gamma + \epsilon^2 a^2 \cos(2\gamma)}, \quad (2)$$

where a and θ are penguin parameters, $\epsilon = \lambda^2/(1 - \lambda^2)$ is defined through a Wolfenstein parameter λ , and γ is the angle γ of the Unitarity Triangle.

Control channels: $B_s^0 \rightarrow J/\Psi K^*$: determine the penguin parameters a and θ .

The observables in $B_s^0 \rightarrow J/\Psi K^*$ measurements:

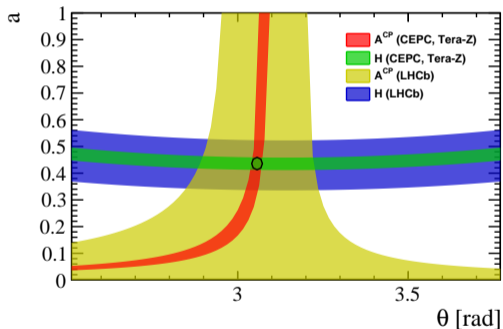
$$A^{\text{CP}} = -\frac{2a \sin \theta \sin \gamma}{1 - 2a \cos \theta \cos \gamma + a^2}, \quad (3)$$

and

$$H = \frac{1 - 2a \cos \theta \cos \gamma + a^2}{1 + 2\epsilon a \cos \theta \cos \gamma + \epsilon^2 a^2}, \quad (4)$$

where A^{CP} is the CP asymmetry and H is an observable constructed containing the branching fraction information, assuming the $SU(3)$ symmetry.

Follow a similar projection method as for ϕ_s and Γ_s :



The expected uncertainty of a and θ is obtained by a χ^2 fit, resulting in

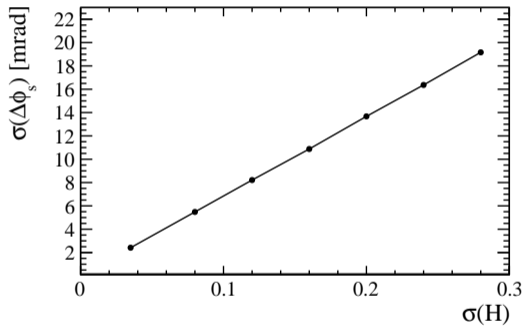
$$a = 0.436 \pm 0.023, \theta = 3.057 \pm 0.016^\circ.$$

. With an error propagation neglecting the correlation between a and θ , the precision of the penguin shift is estimated as

$$\sigma(\Delta\phi_s) = 2.4 \text{ mrad.}$$

(note: $\sigma(\Delta\phi_s) = 4.6 \text{ mrad}$)

SU(3) symmetry violation



- The SU(3) symmetry does not always hold.
- The rightmost point corresponds to $\sigma(H) = 0.28$ (current theory uncertainty).
- $\sigma(\Delta\phi_s)$ is roughly linearly dependent on $\sigma(H)$.
- Without improved theoretical input, the control of penguin contamination will be far from satisfactory.

γ MEASUREMENTS WITH B_s

Extraction of γ

γ is extracted by fitting the time distribution:

$$P_{B_s^0 \rightarrow D_s^+ K^-}(t) \propto e^{-\Gamma_s t} \left(\cosh \left(\frac{\Delta\Gamma_s}{2} t \right) - C \cos(\Delta m_s t) + D_{\bar{f}} \sinh \left(\frac{\Delta\Gamma_s}{2} t \right) - S_{\bar{f}} \sin(\Delta m_s t) \right)$$

$$P_{B_s^0 \rightarrow D_s^- K^+}(t) \propto e^{-\Gamma_s t} \left(\cosh \left(\frac{\Delta\Gamma_s}{2} t \right) + C \cos(\Delta m_s t) + D_f \sinh \left(\frac{\Delta\Gamma_s}{2} t \right) - S_f \sin(\Delta m_s t) \right)$$

$$P_{\bar{B}_s^0 \rightarrow D_s^+ K^-}(t) \propto e^{-\Gamma_s t} \left(\cosh \left(\frac{\Delta\Gamma_s}{2} t \right) + C \cos(\Delta m_s t) + D_{\bar{f}} \sinh \left(\frac{\Delta\Gamma_s}{2} t \right) + S_{\bar{f}} \sin(\Delta m_s t) \right)$$

$$P_{\bar{B}_s^0 \rightarrow D_s^- K^+}(t) \propto e^{-\Gamma_s t} \left(\cosh \left(\frac{\Delta\Gamma_s}{2} t \right) - C \cos(\Delta m_s t) + D_f \sinh \left(\frac{\Delta\Gamma_s}{2} t \right) + S_f \sin(\Delta m_s t) \right)$$

The γ parameters are in the D and S parameters, eg. $D_f = \frac{-2 r_{D_s K} \cos(\delta - (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}$.

Projection to future Z-factories

$$\sigma(\gamma) \propto 1/\sqrt{N_{\text{eff}}}$$

- $N_{\text{eff}} \propto N_{b\bar{b}}$
- $N_{\text{eff}} \propto \text{Efficiency}$
- $N_{\text{eff}} \propto \text{Tagging power}$
- $\sigma_\gamma \propto 1/e^{-\frac{1}{2}\Delta m_s^2 \sigma_t^2}$

Use the similar equation as for ϕ_s to estimate the resolution of γ .

$$\xi = 1/\left(\sqrt{N_{b\bar{b}} \times \varepsilon \times Br} \times \sqrt{p} \times \exp\left(-\frac{1}{2}\Delta m_s^2 \sigma_t^2\right)\right)$$

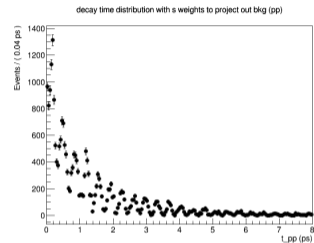
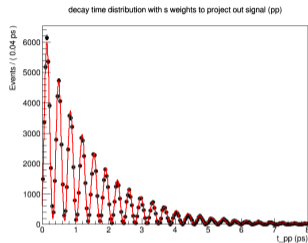
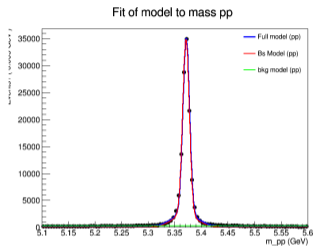
- Stat:

$$\begin{aligned}
 N_{exp}(Z \rightarrow b\bar{b} \rightarrow B_s^0(\rightarrow D_s^-(\rightarrow K^-K^+\pi^-)K^+)X) \\
 &= 10^{12} \times \mathcal{B}(Z \rightarrow b\bar{b}) \times \mathcal{B}(\bar{b} \rightarrow B_s^0) \times \mathcal{B}(B_s^0 \rightarrow D_s^-K^+) \times \mathcal{B}(D_s^- \rightarrow K^-K^+\pi^-) \\
 &= 149804
 \end{aligned} \tag{5}$$

- For the specific $D_s^- \rightarrow K^-K^+\pi^-$ subdecay in the signal samples, the events of $B2DK$ in total should be:

$$\begin{aligned}
 N_{exp}(B_s^0 \rightarrow D_s^\mp(KK\pi)K^\pm) \\
 &= N_{exp}(B_s^0 \rightarrow D_s^+(KK\pi)K^-) + N_{exp}(B_s^0 \rightarrow D_s^-(KK\pi)K^+) \\
 &+ N_{exp}(\bar{B}_s^0 \rightarrow D_s^+(KK\pi)K^-) + N_{exp}(\bar{B}_s^0 \rightarrow D_s^-(KK\pi)K^+) \\
 &= 4 \times N_{exp}(Z \rightarrow b\bar{b} \rightarrow B_s^0(\rightarrow D_s^-(\rightarrow K^-K^+\pi^-)K^+)X) \\
 &= 599216
 \end{aligned} \tag{6}$$

Parameters extraction at MC Truth level



- Perfect flavour tagging and time resolution.
- Resolution: $\sigma(\gamma) = 0.35^\circ$.

$$\xi = 1 / \left(\sqrt{N_{b\bar{b}} \times \varepsilon \times Br} \times \sqrt{p} \times \exp\left(-\frac{1}{2} \Delta m_s^2 \sigma_t^2\right) \right)$$

- Temporarily ignore the time resolution effects, considering the time resolution of B_s is excellent (from $B_s \rightarrow J/\psi\phi$ study).
- Tagging power: 40%
- Resulting $\sigma(\gamma) = 0.55^\circ$
- Expectation from HL-LHC LHCb: $\sigma(\gamma) = 0.35^\circ$.

SUMMARY

Summary

- Competitive ϕ_s resolution for CEPC(Tera-Z) and LHCb(HL-LHC).
 - Expected ϕ_s resolution: CEPC(Tera-Z) is a little worse than LHCb(HL-LHC).
 - CEPC has potential to improve the flavour tagging to get better ϕ_s resolution with better algorithm.
- Only in the 10-Tera-Z configuration, can Z factories be competitive to the LHCb(HL-LHC) for $\Delta\Gamma_s$ and Γ_s measurements.
- Expect good resolution for γ , but more to investigate.
- Particle identification is critical.
 - Hadron pid is not used in reconstruction. With the information, a better efficiency is expected.
 - Tagging power drop fast with particle misidentification.
- Vertex reconstruction is critical for background suppression.

Thank you for your attention!