

A triple Z' signal via light Higgs interaction in Z -factories

Takaaki Nomura (Sichuan University)

Collaborated with Kei Yagyu (Osaka University)

Appearing soon: arXiv: 2410.xxxxx

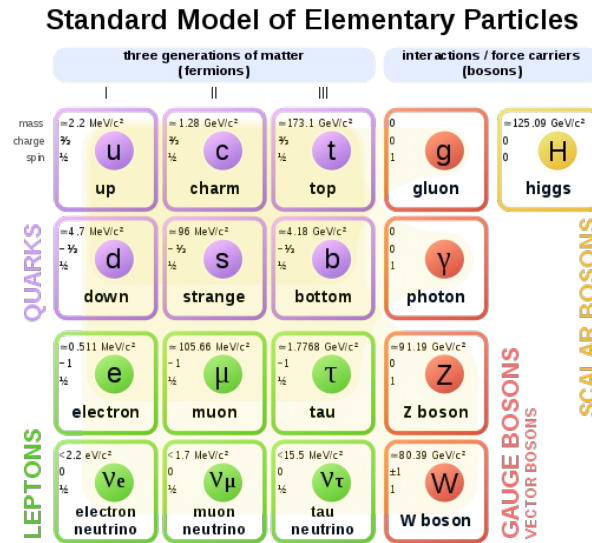


Outline of the talk

1. Introduction
2. Extra $U(1)'$ and Z' boson
3. Triple Z' signatures
4. Summary

1. Introduction

The standard model (SM) of particle physics is successful



The SM is based on gauge symmetry $SU(3)_C \times SU(2)_L \times U(1)_Y$

However, there should be beyond the SM (BSM) physics

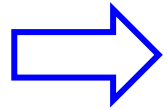
⇒ Dark matter, neutrino mass etc.

What kind of BSM could be tested in near future?

⇒ New particle(s) from extra $U(1)'$ sector is one candidate

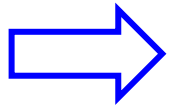
1. Introduction

One simple extension of the SM



A model with extra $U(1)'$ gauge symmetry

◆ The SM is based on gauge symmetry



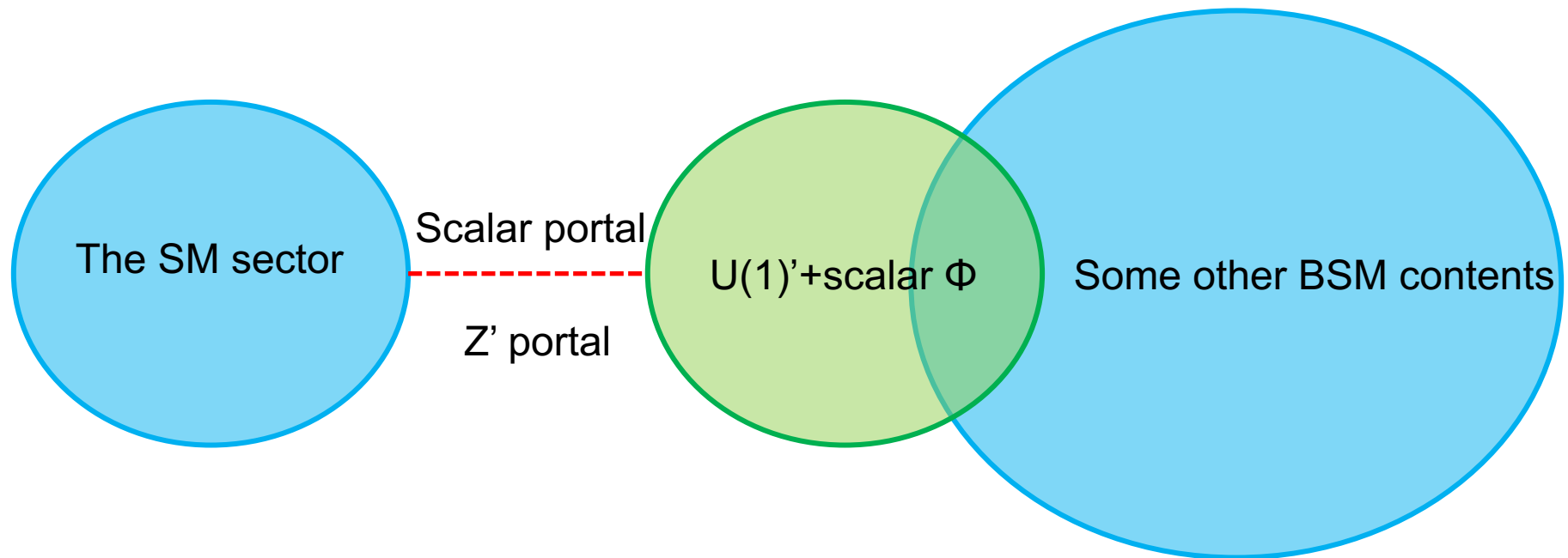
The BSM would be also described by gauge symmetry

◆ It is plausible new physics
providing rich **phenomenology**, e.g.

- Stabilizing dark matter and providing DM-SM mediator
- Application to flavor structure (flavor dependent $U(1)$)
- $U(1)$ breaking scalar VEV \rightarrow Higgs physics
- It would appear from high scale theory (like GUTs, string theory)
- Can be tested if new Z' (scalar) is light (e.g. dark photon(Higgs))

1. Introduction

If $U(1)'$ is spontaneously broken introduction of new scalar is natural



Scalar + Z' portal may connect BSM and the SM

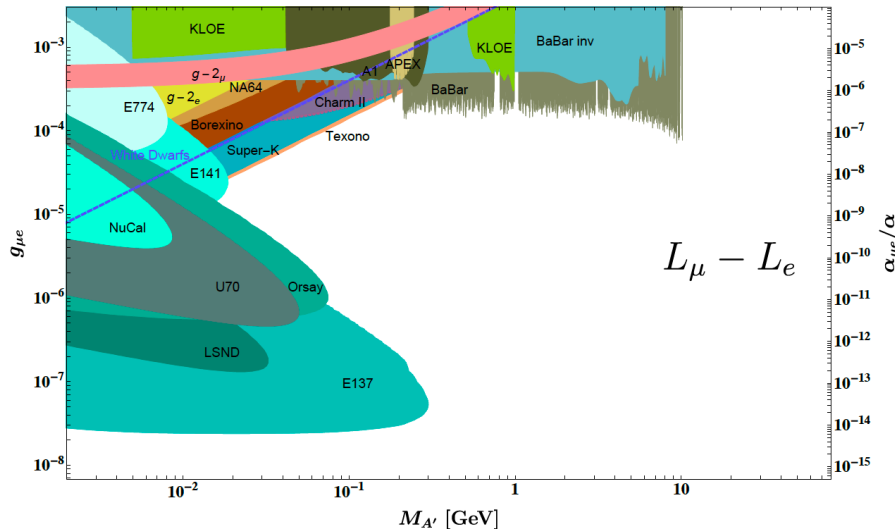
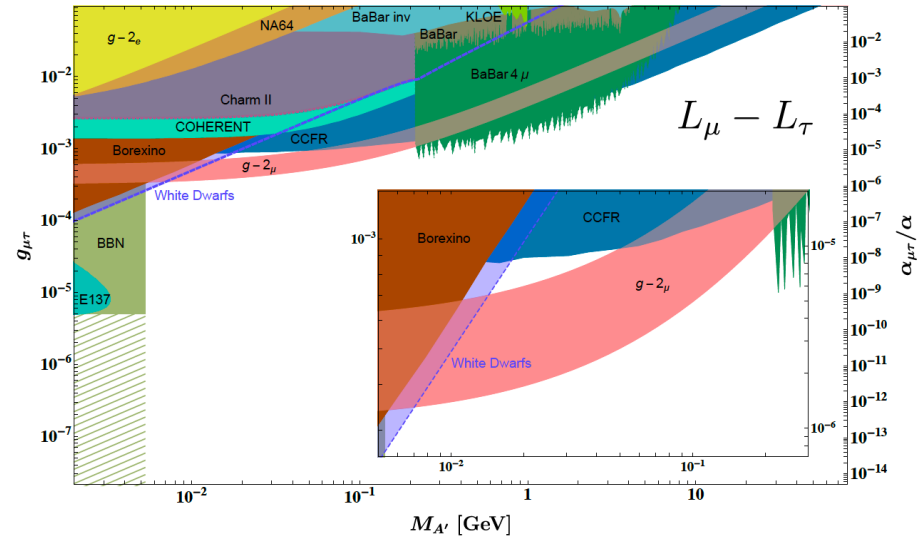
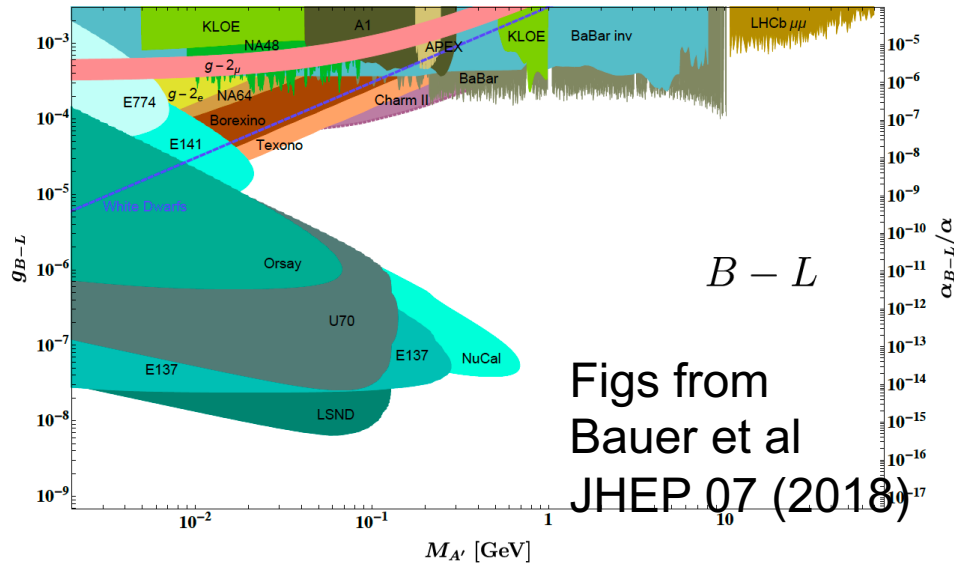
It is more natural to consider scalar and Z' at the same time



Signatures of spontaneously broken $U(1)'$ gauge symmetry

1. Introduction

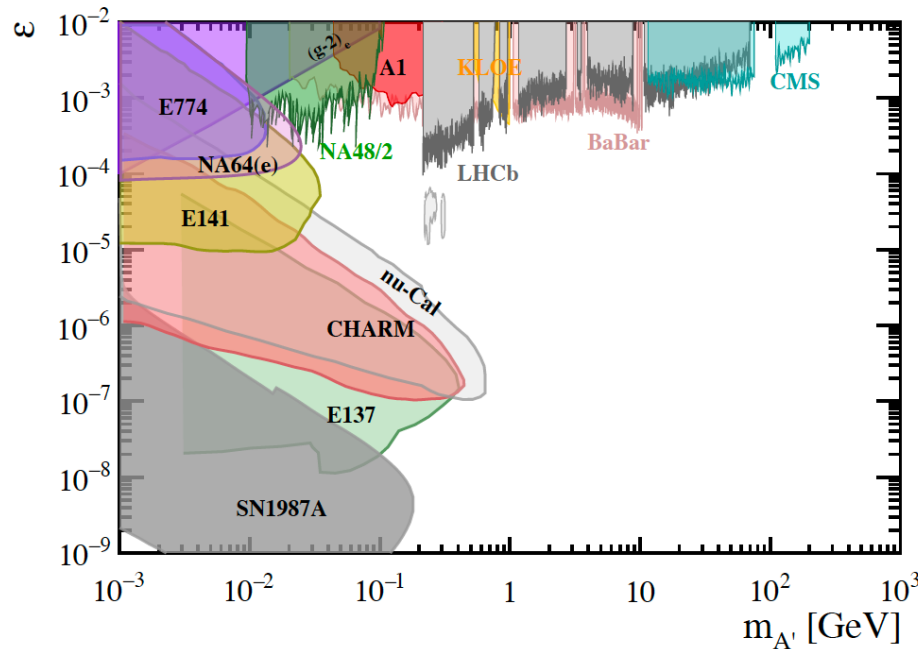
Many Z' boson searches (below electroweak scale)



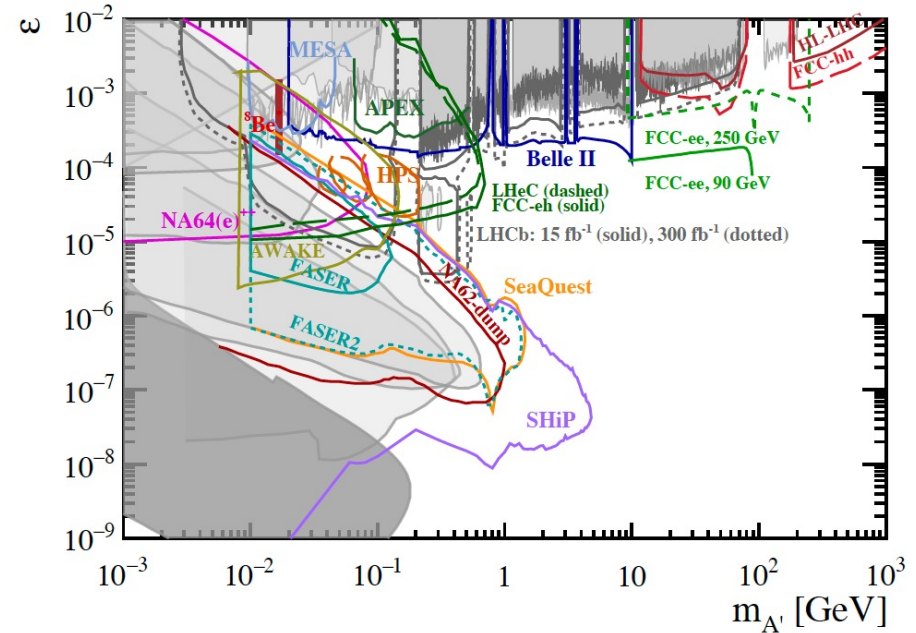
- Beam dump experiments,
- Neutrino scattering,
- Lepton collider experiments,
- Hadron collider experiments,
- Etc.

1. Introduction

Dark photon searches are active (below electroweak scale)



Figs from 2005.01515



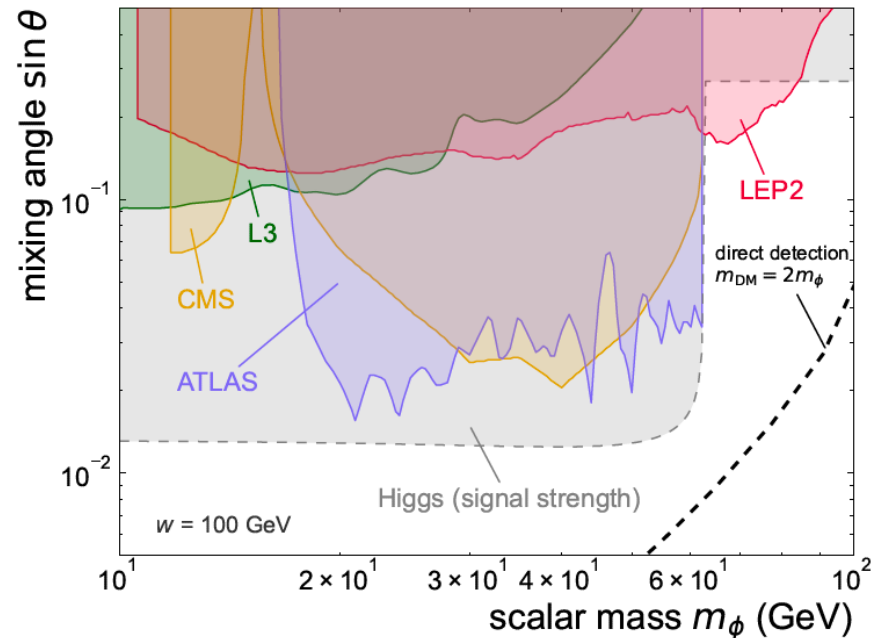
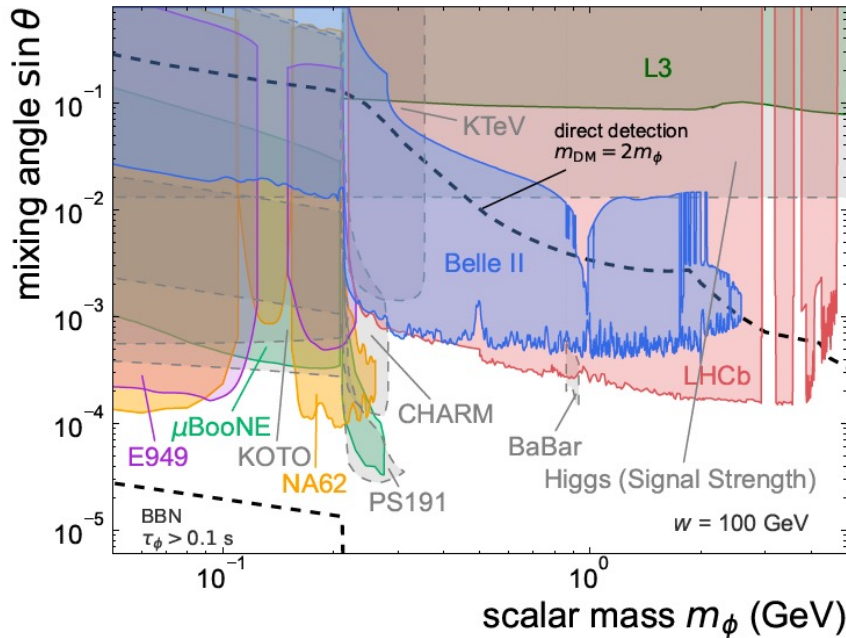
- ✓ Kinetic mixing parameter can be tested
- ✓ Many constraints and future prospects

$$- \frac{\sin \epsilon'}{2} B_{\mu\nu} X^{\mu\nu}$$

$$\mathcal{L}_{A' f \bar{f}} \simeq e \epsilon Q_f \bar{f} \gamma^\mu f A'_\mu$$

1. Introduction

Light scalar is also searched for



Figs from 2305.16169

Constraint on scalar mixing (new scalar and SM Higgs)

In particular, meson decays constrain the mixing strongly ($m_{\text{scalar}} < 5$ GeV)

1. Introduction

New signal would appear considering both scalar and Z'

We discuss new Z boson decay chain

$$\Rightarrow Z \rightarrow Z' \phi \rightarrow Z' Z' Z' \quad \text{Providing triple } Z'$$

It is a good target at Z -factories (e^+e^- collision with $\sqrt{s} \simeq m_Z$)

\Rightarrow Can be realized at CEPC, FCC-ee, etc.

In Z -factories we expect $O(10^{12})$ Z boson production (**Tera- Z**)

\Rightarrow High sensitivity to Z boson decay

Outline of the talk

1. Introduction
2. Extra $U(1)'$ and Z' boson
3. Triple Z' signatures
4. Summary

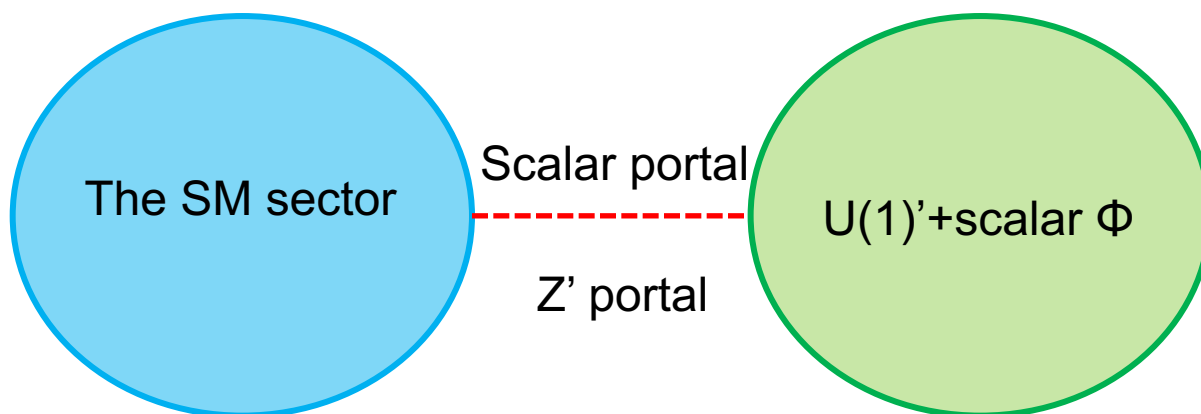
2. Extra $U(1)$ and Z' boson

We consider simple setting

⇒ **New $U(1)'$ gauge symmetry**
+
SM singlet scalar Φ with non-zero $U(1)'$ charge

The scalar field develops a vacuum expectation value (VEV) to break $U(1)'$

It is **the simplest** field contents for spontaneously broken local $U(1)'$ scenario



2. Extra U(1) and Z' boson

For SM + extra U(1)_X gauge symmetry


U(1) Gauge sector

$$\mathcal{L} \supset -\frac{1}{4}B_{\mu\nu}B^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} - \frac{\sin \epsilon'}{2}B_{\mu\nu}X^{\mu\nu}$$

U(1)_Y U(1)' U(1) kinetic mixing

Scalar potential

$$V = -\mu_H^2|H|^2 - \mu_\Phi^2|\Phi|^2 + \frac{\lambda_H}{2}|H|^4 + \frac{\lambda_\Phi}{2}|\Phi|^4 + \lambda_{H\Phi}|H|^2|\Phi|^2.$$
$$H = \begin{pmatrix} G^+ \\ \frac{1}{\sqrt{2}}(v + \tilde{h} + iG) \end{pmatrix}, \quad \Phi = \frac{1}{\sqrt{2}}(v_\Phi + \tilde{\phi} + iG_\Phi)$$

Scalar develops VEVs: $\langle H(\Phi) \rangle = v(v_\Phi)/\sqrt{2}$  Electroweak and U(1)' break

Scalar mass and mixing

$$\begin{pmatrix} \tilde{h} \\ \tilde{\phi} \end{pmatrix} = R(\alpha) \begin{pmatrix} h \\ \phi \end{pmatrix}, \quad R(\theta) = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \quad \tan 2\alpha = \frac{2\lambda_{H\Phi}vv_\Phi}{\lambda_H v^2 - \lambda_\Phi v_\Phi^2}$$

$$m_h^2 = \lambda_H v^2 \cos^2 \alpha + \lambda_\Phi v_\Phi^2 \sin^2 \alpha + 2\lambda_{H\Phi}vv_\Phi \sin \alpha \cos \alpha,$$

$$m_\phi^2 = \lambda_\Phi v_\Phi^2 \cos^2 \alpha + \lambda_H v^2 \sin^2 \alpha - 2\lambda_{H\Phi}vv_\Phi \sin \alpha \cos \alpha.$$

2. Extra U(1) and Z' boson

Kinetic term of U(1) gauge fields can be diagonalized by

$$\begin{pmatrix} X_\mu \\ B_\mu \end{pmatrix} = \begin{pmatrix} \operatorname{cosec} \epsilon' & 0 \\ -\tan \epsilon' & 1 \end{pmatrix} \begin{pmatrix} \tilde{Z}'_\mu \\ \tilde{B}_\mu \end{pmatrix}$$

General kinetic term after the transformation

$$D_\mu \Psi = \left[\partial_\mu - ig(T^+ W_\mu^+ + \text{c.c.}) - ieQ_\Psi A_\mu - ig_Z(T_\Psi^3 - s_W^2 Q_\Psi) \tilde{Z}_\mu - ig_X X_\Psi \tilde{Z}'_\mu \right] \Psi$$

$$X_\Psi = \tilde{X}_\Psi - Y_\Psi \frac{g'}{g_X} \tan \epsilon'$$

\tilde{X}_Ψ : U(1)' charge

Z-Z' mass term

$$M_{ZZ'} = \begin{pmatrix} m_Z^2 - \delta_Z^2 & -m_Z v \delta \\ -m_Z v \delta & m_{Z'}^2 - \delta_{Z'}^2 + v^2 \delta^2 \end{pmatrix} + \mathcal{O}(\delta^3)$$

$$\begin{cases} \delta_Z^2 = \frac{g_Z^2 v^4}{g_Z^2 v^2 - 4g_X^2 X_\Phi^2 v_\Phi^2} \delta^2, \\ \delta_{Z'}^2 = -\frac{4g_X^2 X_\Phi^2 v^2 v_\Phi^2}{g_Z^2 v^2 - 4g_X^2 X_\Phi^2 v_\Phi^2} \delta^2. \end{cases}$$

$$\delta \equiv g_X X_H$$

Mass eigenstates

$$\begin{pmatrix} \tilde{Z}_\mu \\ \tilde{Z}'_\mu \end{pmatrix} = R(\zeta) \begin{pmatrix} Z_\mu \\ Z'_\mu \end{pmatrix}, \quad \sin 2\zeta = \frac{g_Z v^2 \delta}{m_{Z'}^2 - m_Z^2}$$

($\delta \ll 1$ to suppress Z-Z' mixing)

2. Extra U(1) and Z' boson

Gauge-Gauge-Scalar interaction

$$\mathcal{L}_{\text{int}} = (Z_\mu, Z'_\mu) R^T(\zeta) M_{\text{int}} R(\zeta) \begin{pmatrix} Z^\mu \\ Z'^\mu \end{pmatrix} + \mathcal{O}(\delta^3),$$

$$M_{\text{int}} = \begin{pmatrix} (m_Z^2 - \delta_Z^2) \frac{\tilde{h}}{v} & -m_Z v \delta \frac{\tilde{h}}{v} \\ -m_Z v \delta \frac{\tilde{h}}{v} & (m_{Z'}^2 - \delta_{Z'}^2) \frac{\tilde{\phi}}{v_\Phi} + v \delta^2 \tilde{h} \end{pmatrix},$$



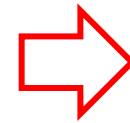
For tiny Z-Z' mixing ($\zeta \ll 1$)

$$\mathcal{L}_{\text{int}} = (Z_\mu, Z'_\mu) \begin{pmatrix} \Lambda_{11} & \Lambda_{12} \\ \Lambda_{12} & \Lambda_{22} \end{pmatrix} \begin{pmatrix} Z^\mu \\ Z'^\mu \end{pmatrix}$$

$$\Lambda_{11} = \frac{m_Z^2}{v} \tilde{h} + \frac{m_Z^2 m_{Z'}^2 v^2}{(m_Z^2 - m_{Z'}^2)^2} \left(\frac{\tilde{\phi}}{v_\Phi} - \frac{\tilde{h}}{v} \right) \delta^2 + \mathcal{O}(\delta^4),$$

$$\Lambda_{22} = \frac{m_{Z'}^2}{v_\Phi} \tilde{\phi} - \frac{m_{Z'}^4 v^2}{(m_Z^2 - m_{Z'}^2)^2} \left(\frac{\tilde{\phi}}{v_\Phi} - \frac{\tilde{h}}{v} \right) \delta^2 + \mathcal{O}(\delta^4),$$

$$\Lambda_{12} = \frac{m_Z m_{Z'}^2 v}{m_Z^2 - m_{Z'}^2} \left(\frac{\tilde{\phi}}{v_\Phi} - \frac{\tilde{h}}{v} \right) \delta + \mathcal{O}(\delta^3).$$



ZZ'Φ interaction

$$\begin{pmatrix} \tilde{h} \\ \tilde{\phi} \end{pmatrix} = R(\alpha) \begin{pmatrix} h \\ \phi \end{pmatrix}, \quad R(\theta) = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix}$$

Outline of the talk

1. Introduction
2. Extra $U(1)'$ and Z' boson
3. Triple Z' signatures
4. Summary

3. Triple Z' signatures

New signal of spontaneously broken U(1)' at Z factories

⇒ Decay of Z boson into Z'Φ mode

It is induced by interactions after spontaneous symmetry breaking

$$\mathcal{L}_{ZZ'\phi} \simeq \frac{m_Z m_{Z'}^2 v}{m_{Z'}^2 - m_Z^2} \frac{\cos \alpha}{v_\Phi} \delta Z_\mu Z'^\mu \phi$$

We consider the BR of the process in focusing on **dark photon** case

$$X_H = -\frac{1}{2} \frac{g'}{g_X} \tan \epsilon' \quad (\tilde{X}_H = 0) \quad \Rightarrow \quad \delta = g_X X_H = -\frac{1}{2} e \frac{\cos \theta_W \tan \epsilon'}{\epsilon}$$

$$\mathcal{L}_{Z'f\bar{f}} \simeq e\epsilon Q_f \bar{f} \gamma^\mu f Z'_\mu \quad \text{Dark photon interaction with SM fermions}$$

⇒ Explore sensitivity to kinetic mixing ϵ (and new gauge coupling g_X)

3. Triple Z' signatures

Decay BRs of $Z \rightarrow Z'\Phi$ process

Decay width

$$\Gamma(Z \rightarrow Z'\phi) = \frac{m_Z}{48\pi} \left(s_\alpha + \frac{v}{v_\Phi} c_\alpha \right)^2 \frac{x_{Z'} \delta^2}{(1 - x_{Z'})^2} \\ \times \left[(1 + x_{Z'} - x_\phi)^2 + 8x_{Z'} \right] \lambda^{1/2}(x_{Z'}, x_\phi)$$

$$x_{Z'} = m_{Z'}^2/m_Z^2 \text{ and } x_\phi = m_\phi^2/m_Z^2 \quad \lambda(x, y) = 1 + x^2 + y^2 - 2x - 2y - 2xy$$

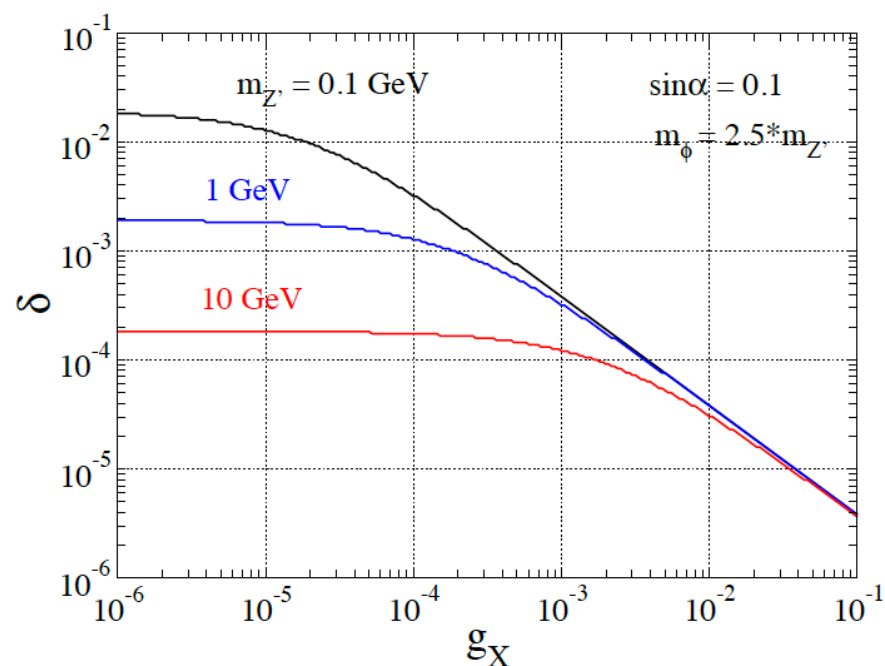
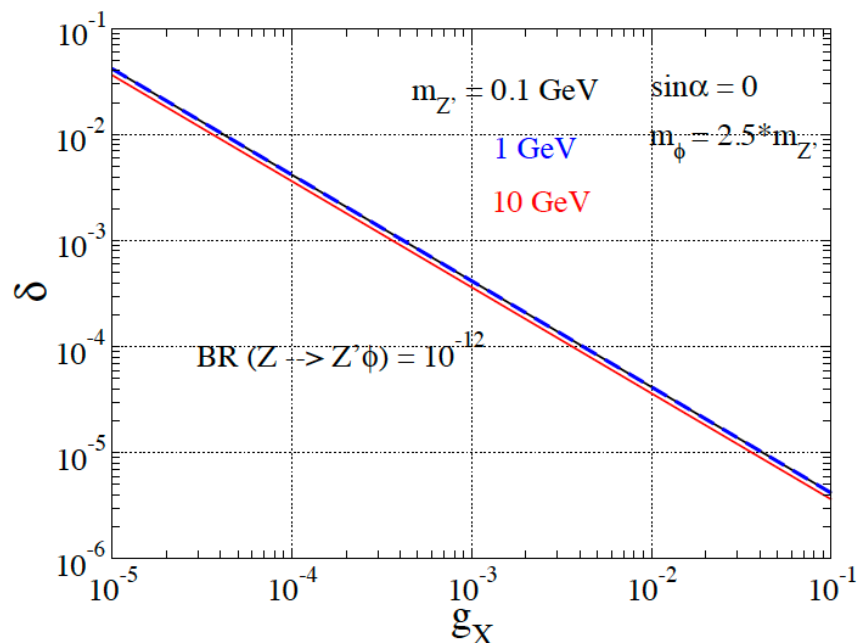
Decay BR

$$\text{BR}(Z \rightarrow Z'\phi) \sim \frac{m_Z}{48\pi} \left(s_\alpha + \frac{v}{v_\Phi} c_\alpha \right)^2 \frac{x_{Z'} \delta^2}{\Gamma_Z^{\text{obs}}} \\ \sim 0.24 \times \left(s_\alpha + \frac{v}{v_\Phi} c_\alpha \right)^2 x_{Z'} \delta^2$$

$$\Rightarrow \left\{ \begin{array}{ll} \text{BR}(Z \rightarrow Z'\phi) \sim 1.76 \times g_X^2 X_\Phi^2 \delta^2 & \text{for } s_\alpha \rightarrow 0 \\ \text{BR}(Z \rightarrow Z'\phi) \sim 0.24 \times s_\alpha^2 x_{Z'} \delta^2 & \text{for } |s_\alpha| \gg v/v_\Phi \end{array} \right.$$

3. Triple Z' signatures

Decay BRs of $Z \rightarrow Z'\Phi$ process



- ✓ We have sensitivity for small δ if g_X is sizable
- ✓ Contribution from scalar mixing for tiny g_X but not very large
 - ❖ We consider $\alpha \rightarrow 0$ limit for simplicity

Dark photon case is promising since we can have $g_X = O(1)$ - $O(0.1)$

Note: $m_\phi \geq m_{Z'}$, and $g_X = O(1)$ - $O(0.1)$ is plausible case (SM: $m_h \geq m_Z$, $0.1 < g, g' < 1$)

3. Triple Z' signatures

Decay BRs of scalar bosons

Decay widths of new scalar boson

$$\Gamma(\phi \rightarrow Z'Z') = \frac{m_{Z'}^4 \cos^2 \alpha}{8\pi v_\Phi^2 m_\phi} \beta(x_{Z'}) \left[2 + \frac{1}{4x_{Z'}^2} (1 - 2x_{Z'})^2 \right]$$

$$\Gamma(\phi \rightarrow hh) = \frac{\lambda_{\phi hh}^2}{8\pi m_\phi} \beta(x_h), \quad x_i = m_i^2/m_\phi^2 \text{ and } \beta(x) = \sqrt{1 - 4x}.$$

- ✓ Decay widths for the SM particles modes also exist via H- Φ mixing

Decay widths of SM Higgs boson for new modes

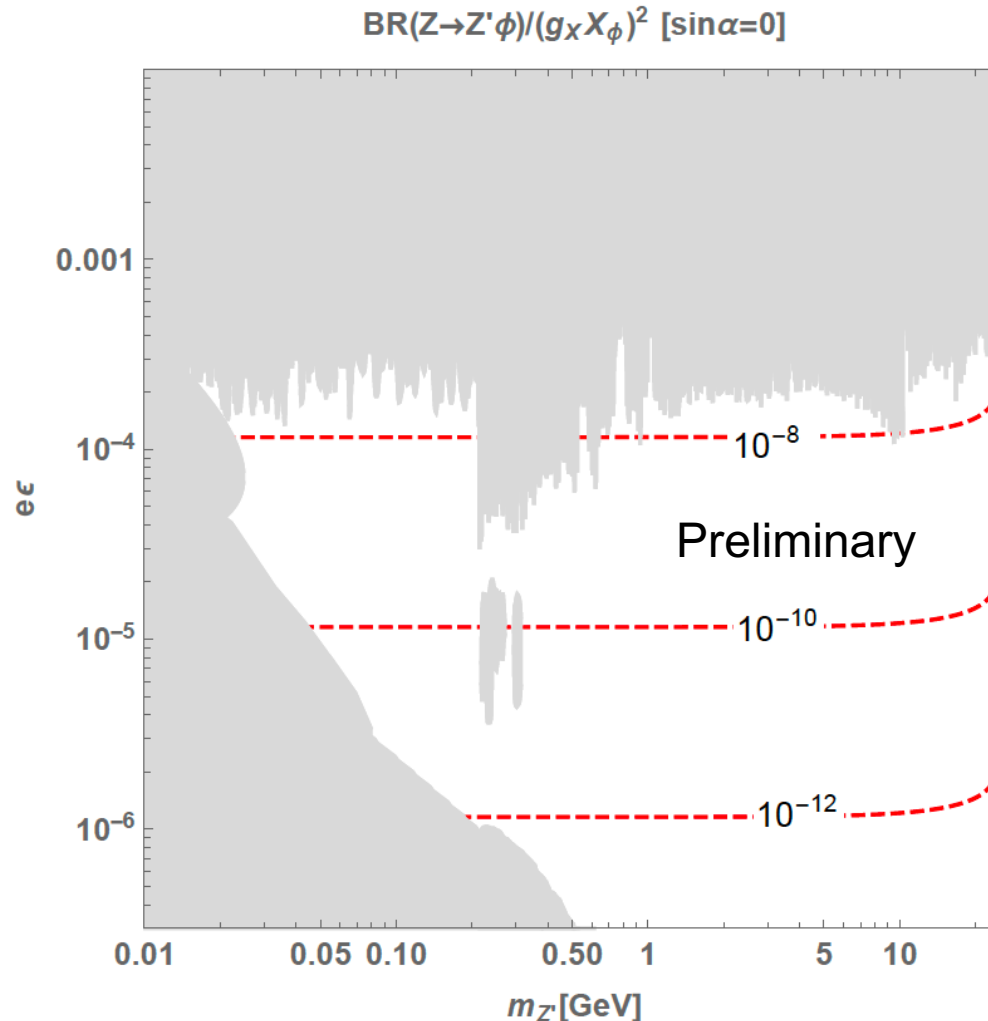
$$\Gamma(h \rightarrow Z'Z') = \frac{m_{Z'}^4 \sin^2 \alpha}{8\pi v_\Phi^2 m_h} \beta(z_{Z'}) \left[2 + \frac{1}{4z_{Z'}^2} (1 - 2z_{Z'})^2 \right]$$

$$\Gamma(h \rightarrow \phi\phi) = \frac{\lambda_{\phi\phi h}^2}{8\pi m_h} \beta(z_\phi),$$

- ✓ For $\alpha \rightarrow 0$, new scalar decays into $Z'Z'$ with 100% BR
- ✓ New Higgs decay modes also vanish in the limit

3. Triple Z' signatures

The branching ratio (dark photon case)



$$(m_\phi = 2.5 m_{Z'})$$

✓ Gray region is excluded by dark photon searches

We expect sizable number of event at Z-factories: a few $\times 10^{12}$ Z boson

3. Triple Z' signatures

Benchmark points

BP1 : $m_{Z'} = 0.1$ GeV, $e\epsilon = 10^{-5}$, $g_X = 0.5$,

BP2 : $m_{Z'} = 10$ GeV, $e\epsilon = 10^{-5}$, $g_X = 0.5$,

Z' (dark photon) decaying into SM fermions and BRs are

$$BR(Z' \rightarrow e^+e^-) \simeq 1.0$$

BP1

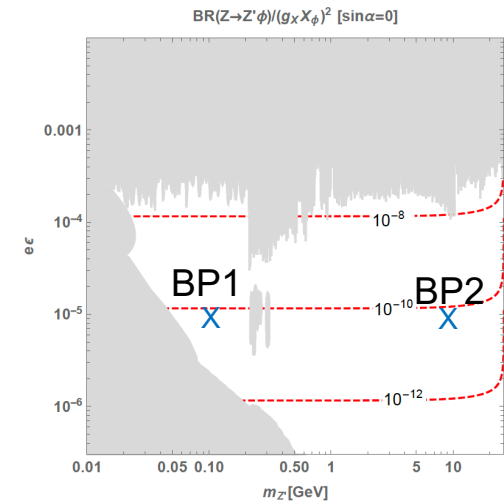
$$BR(Z' \rightarrow e^+e^-/\mu^+\mu^-/\tau^+\tau^-) \simeq 0.15, \quad BR(Z' \rightarrow \text{hadron}) \simeq 0.55$$

BP2

Expected number of events for BPs with 10^{12} Z

	6ℓ	$4\ell + \text{had.}$	$2\ell + \text{had.}$	had.
BP1	46.	0	0	0
BP2	1.1	6.1	11.	6.6

$\left(\begin{array}{l} \text{had.} = \text{hadrons} \\ \ell = e, \mu \end{array} \right)$



We expect clear signal at Z-factories

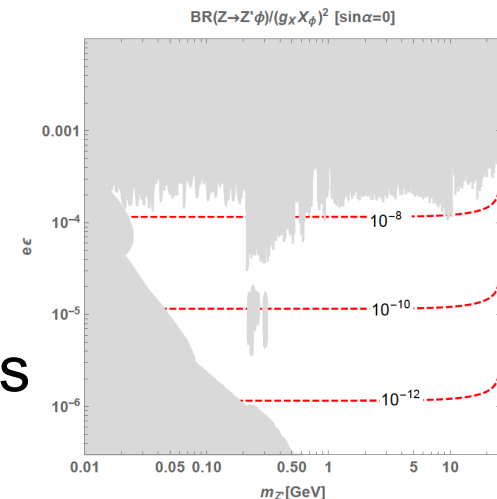
Summary and Discussions

Extension of the SM with extra $U(1)'$ gauge symmetry

- ✓ Z' boson from extra $U(1)'$
- ✓ We also would have new scalar boson via SSB
- ✓ Z' and scalar boson would be similar mass scale
- ✓ Z decay into scalar + Z' can happen

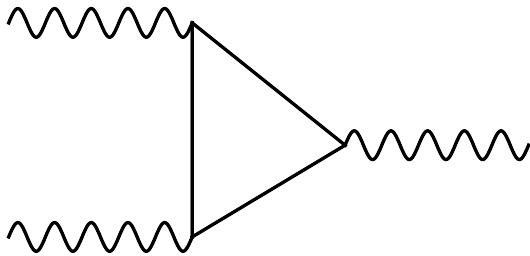
□ Triple Z' (dark photon) signal at Z-factories

- ✓ $Z \rightarrow Z' \Phi \rightarrow Z' Z' Z'$ decay chain
- ✓ Sensitivity to kinetic mixing $> 10^{-6}$ at Tera Z-factories
- ✓ Clear signals are expected in dark photon case
- ✓ Potential of direct BSM discovery at Tera Z-factories



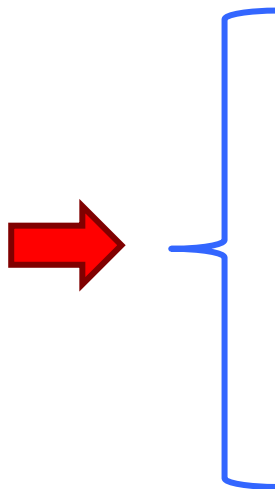
Gauge anomaly free conditions

In constructing an extra $U(1)'$ gauge symmetric model charge assignment of fermion contents should be anomaly free



$$\sum_f \left(\text{Tr}[T_i T_j T_k]_R - \text{Tr}[T_i T_j T_k]_L \right) = 0$$

T_i : generator of gauge group


$$\begin{aligned} & [\text{SU}(3)_c]^2 U(1)_X \quad [\text{SU}(2)_L]^2 U(1)_X \\ & [U(1)_Y]^2 U(1)_X \quad [U(1)_X]^2 U(1)_Y \\ & [U(1)_X]^3 \quad [\text{gravity}]^2 U(1)_X \end{aligned}$$

Conditions in addition to the SM gauge anomaly free conditions