

Lam-Tung relation breaking in Z boson production as a probe of SMEFT effects

Bin Yan Institute of High Energy Physics

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Based on: Xu Li, Bin Yan, C.-P. Yuan, 2405.04069

The Drell-Yan process

A "Standard candle" at the LHC

- Large cross section
- Clean leptonic signature

Precise Z boson measurements:

test pQCD, constrain PDFs, weak mixing angle, strong coupling



Y. Fu, R. Brock, D. Hayden, C.-P. Yuan, PRD 2024, S. Yang, Y. Fu, M. Liu, L. Han, T. Hou, C.-P. Yuan, PRD 2022, S. Yang et al, EPJC 2022 ATLAS-CONF-2023-015



The most precision measurements

The Drell-Yan process

A "Standard candle" at the LHC

- Large cross section
- Clean leptonic signature



Precise Z boson measurements:

Detector calibration, luminosity monitor,...

W-mass @ CDF, Science 376, 170-176 (2022)

Searches for BSM physics

S. Alioli, W. Dekens, M. Girard, E. Mereghetti, JHEP 2018 S. Alioli, R. Boughzal, E. Mereghetti, F. Petriello, PLB 2020 R. Boughezal, Y. Huang, F. Petriello, PRD 2022, 2023 X. Li, K. Mimasu, K. Yamashita, C. Yang, C. Zhang, S.-Y. Zhou, JHEP 2022 S. Grossi and R. Torre 2404.10569

Retain full differential information on the leptons (spin information):

Can be encoded in eight angular coefficients

Novel angular dependence from NP: S. Alioli, R. Boughzal, E. Mereghetti, F. Petriello, PLB 2020

Directly probe production dynamics

R. Gauld et al, JHEP 2017, N3LO

Angular coefficients

Decomposition in terms of spherical harmonics in Collins-soper frame:



What's the physical meaning of these angular coefficients?

Angular coefficients



	$\left(-rac{1-\delta_L}{3} + rac{J_3}{2} ight)$	$rac{J_1+2Q_{xz}-i(J_2+2Q_{yz})}{2\sqrt{2}}$	$\lambda_T - i Q_{xy}$)
$ ho_{\lambda_Z\lambda'_Z} =$	$rac{J_1+2Q_{xz}+i(J_2+2Q_{yz})}{2\sqrt{2}}$	$rac{1+2\delta_L}{3}$	$rac{J_1 - 2Q_{xz} - i(J_2 - 2Q_{yz})}{2\sqrt{2}}$
	$ig\langle \lambda_T + i Q_{xy}$	$rac{J_1 - 2Q_{xz} + i(J_2 - 2Q_{yz})}{2\sqrt{2}}$	$rac{1-\delta_L}{3}-rac{J_3}{2} ight)$

$$\frac{\mathrm{d}\sigma}{\mathrm{d}^4 q \,\mathrm{d}\cos\theta \,\mathrm{d}\phi} = \frac{3}{16\pi} \,\frac{\mathrm{d}\sigma^{\mathrm{unpol.}}}{\mathrm{d}^4 q} \left\{ (1+\cos^2\theta) + \frac{1}{2} \,A_0 \,\left(1-3\cos^2\theta\right) \right. \\ \left. + A_1 \,\sin(2\theta)\cos\phi + \frac{1}{2} \,A_2 \,\sin^2\theta \,\cos(2\phi) \right. \\ \left. + A_3 \,\sin\theta \,\cos\phi + A_4 \,\cos\theta + A_5 \,\sin^2\theta \,\sin(2\phi) \right. \\ \left. + A_6 \,\sin(2\theta) \,\sin\phi + A_7 \,\sin\theta \,\sin\phi \right\},$$

$$\begin{split} \frac{\Gamma}{\Omega_f^*} \propto & \frac{|B_+|^2 + |B_-|^2}{2} \left[\frac{2}{3} + \frac{\delta_L}{3} \left(1 - 3\cos^2\theta_f^* \right) + \lambda_T \sin^2\theta_f^* \cos 2\phi_f^* \right. \\ & \left. + Q_{yz} \sin 2\theta_f^* \sin \phi_f^* + Q_{xz} \sin 2\theta_f^* \cos \phi_f^* + Q_{xy} \sin^2\theta_f^* \sin 2\phi_f^* \right] \\ & \left. + \frac{|B_+|^2 - |B_-|^2}{2} \left(J_1 \sin \theta_f^* \cos \phi_f^* + J_2 \sin \theta_f^* \sin \phi_f^* + J_3 \cos \theta_f^* \right) \right] \end{split}$$

$$A_{0} \rightarrow \delta_{L}$$

$$A_{1} \rightarrow Q_{xz}$$

$$A_{2} \rightarrow \lambda_{T}$$

$$A_{3} \rightarrow J_{1}$$

$$A_{4} \rightarrow J_{3}$$

$$A_{5} \rightarrow Q_{xy}$$

$$A_{6} \rightarrow Q_{yz}$$

$$A_{7} \rightarrow J_{2}$$

Angular coefficients and polarization



$$A_{0} \rightarrow \delta_{L}$$

$$A_{1} \rightarrow Q_{xz}$$

$$A_{3} \rightarrow J_{1}$$

$$A_{6} \rightarrow Q_{yz}$$

$$A_{7} \rightarrow J_{2}$$

$$A_{2} \rightarrow \lambda_{T}$$

$$A_{5} \rightarrow Q_{xy}$$

$$A_{4} \rightarrow J_{3}$$

Longitudinal polarization of Z boson

Interference between transverse and longitudinal Z boson

Linear polarization of Z boson

Parity violation effects



$$\frac{\mathrm{d}\sigma}{\mathrm{d}^4 q \,\mathrm{d}\cos\theta \,\mathrm{d}\phi} = \frac{3}{16\pi} \frac{\mathrm{d}\sigma^{\mathrm{unpol.}}}{\mathrm{d}^4 q} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0 \left(1 - 3\cos^2\theta\right) \right. \\ \left. + A_1 \,\sin(2\theta)\cos\phi + \frac{1}{2} A_2 \,\sin^2\theta \,\cos(2\phi) \right. \\ \left. + A_3 \,\sin\theta \,\cos\phi + A_4 \,\cos\theta + A_5 \,\sin^2\theta \,\sin(2\phi) \right. \\ \left. + A_6 \,\sin(2\theta) \,\sin\phi + A_7 \,\sin\theta \,\sin\phi \right\},$$

$$ho_{\lambda_Z\lambda'_Z} = egin{pmatrix} rac{1-\delta_L}{3}+rac{J_3}{2} & rac{J_1+2Q_{xz}-i(J_2+2Q_{yz})}{2\sqrt{2}} & \lambda_T-iQ_{xy} \ rac{J_1+2Q_{xz}+i(J_2+2Q_{yz})}{2\sqrt{2}} & rac{1+2\delta_L}{3} & rac{J_1-2Q_{xz}-i(J_2-2Q_{yz})}{2\sqrt{2}} \ \lambda_T+iQ_{xy} & rac{J_1-2Q_{xz}+i(J_2-2Q_{yz})}{2\sqrt{2}} & rac{1-\delta_L}{3}-rac{J_3}{2} \end{pmatrix}$$

$$\begin{split} \frac{\Gamma}{\Omega_f^*} \propto & \frac{|B_+|^2 + |B_-|^2}{2} \left[\frac{2}{3} + \frac{\delta_L}{3} \left(1 - 3\cos^2\theta_f^* \right) + \lambda_T \sin^2\theta_f^* \cos 2\phi_f^* \right. \\ & \left. + Q_{yz} \sin 2\theta_f^* \sin \phi_f^* + Q_{xz} \sin 2\theta_f^* \cos \phi_f^* + Q_{xy} \sin^2\theta_f^* \sin 2\phi_f^* \right] \\ & \left. + \frac{|B_+|^2 - |B_-|^2}{2} \left(J_1 \sin \theta_f^* \cos \phi_f^* + J_2 \sin \theta_f^* \sin \phi_f^* + J_3 \cos \theta_f^* \right). \end{split}$$

Lam-Tung relation: $A_0 = A_2$

Linear and Longitudinal polarization of Z boson

- ➢ Spin ½ nature of quarks @ LO
- Vector coupling of spin-1 gluon @ NLO
- Violation @ NNLO and beyond



$$\frac{\mathrm{d}\sigma}{\mathrm{d}^4 q \,\mathrm{d}\cos\theta \,\mathrm{d}\phi} = \frac{3}{16\pi} \frac{\mathrm{d}\sigma^{\mathrm{unpol.}}}{\mathrm{d}^4 q} \left\{ (1 + \cos^2\theta) + \frac{1}{2} A_0 \left(1 - 3\cos^2\theta\right) + A_1 \sin(2\theta)\cos\phi + \frac{1}{2} A_2 \sin^2\theta \cos(2\phi) + A_3 \sin\theta \cos\phi + A_4 \cos\theta + A_5 \sin^2\theta \sin(2\phi) + A_6 \sin(2\theta) \sin\phi + A_7 \sin\theta \sin\phi \right\},$$

$$p_{\lambda_Z\lambda'_Z} = egin{pmatrix} rac{1-\delta_L}{3}+rac{J_3}{2} & rac{J_1+2Q_{xz}-i(J_2+2Q_{yz})}{2\sqrt{2}} & \lambda_T-iQ_{xy} \ rac{J_1+2Q_{xz}+i(J_2+2Q_{yz})}{2\sqrt{2}} & rac{1+2\delta_L}{3} & rac{J_1-2Q_{xz}-i(J_2-2Q_{yz})}{2\sqrt{2}} \ \lambda_T+iQ_{xy} & rac{J_1-2Q_{xz}+i(J_2-2Q_{yz})}{2\sqrt{2}} & rac{1-\delta_L}{3}-rac{J_3}{2} \end{pmatrix}$$

$$\frac{\Gamma}{\Omega_f^*} \propto \frac{|B_+|^2 + |B_-|^2}{2} \left[\frac{2}{3} + \frac{\delta_L}{3} \left(1 - 3\cos^2\theta_f^* \right) + \lambda_T \sin^2\theta_f^* \cos 2\phi_f^* \right. \\ \left. + Q_{yz} \sin 2\theta_f^* \sin \phi_f^* + Q_{xz} \sin 2\theta_f^* \cos \phi_f^* + Q_{xy} \sin^2\theta_f^* \sin 2\phi_f^* \right] \\ \left. + \frac{|B_+|^2 - |B_-|^2}{2} \left(J_1 \sin \theta_f^* \cos \phi_f^* + J_2 \sin \theta_f^* \sin \phi_f^* + J_3 \cos \theta_f^* \right) \right]$$

Lam-Tung relation: $A_0 = A_2$

Linear and Longitudinal polarization of Z boson

 $A_0 \neq A_2$ @ NNLO in QCD non-coplanarity between the hadron and parton planes

J.C. Peng et al, PLB 758,384 (2016)



These results are confirmed by CMS (PLB750, 154 (2015)) and LHCb (PRL 129 (2022) 091801) collaborations

The discrepancy with the SM prediction NP effects or non-perturbative effects ?

Boer-Mulders function

Quark Spin

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The $cos2\phi$ dependence can be induced by the Boer-Mulders function

Nucleon Spin

Leading Quark TMDPDFs

			_		
		Quark Polarization			
		Un-Polarized (U)	Longitudinally Polarized (L)	Transversely Polarized (T)	
Nucleon Polarization	U	$f_1 = \underbrace{\bullet}_{\text{Unpolarized}}$		$h_1^\perp = \bigcirc - \bigcirc$ Boer-Mulders	
	L		$g_1 = \underbrace{\bullet \bullet}_{\text{Helicity}} \bullet \bullet \bullet$	$h_{1L}^{\perp} = \underbrace{ \checkmark}_{\text{Worm-gear}} - \underbrace{ \checkmark}_{\text{Worm-gear}}$	
	т	$f_{1T}^{\perp} = \underbrace{\bullet}^{\uparrow} - \underbrace{\bullet}_{Sivers}$	$g_{1T}^{\perp} = \underbrace{\stackrel{\uparrow}{\bullet \bullet}}_{\text{Worm-gear}} - \underbrace{\stackrel{\uparrow}{\bullet \bullet}}_{\text{Worm-gear}}$	$h_1 = \underbrace{\stackrel{\uparrow}{\blacktriangleright} - \underbrace{\stackrel{\uparrow}{\uparrow}}_{\text{Transversity}} \\ h_{1T}^{\perp} = \underbrace{\stackrel{\uparrow}{\checkmark} - \underbrace{\stackrel{\uparrow}{\checkmark}}_{\text{Pretzelosity}} \\ \end{pmatrix}$	





Transversely polarized quark

Lam-Tung relation and NP

Center-of-mass frame:

$$\frac{d\sigma}{d\Omega} = a\cos\hat{\theta} + b\cos^2\hat{\theta} + c\cos^3\hat{\theta} + d$$

$$\cos\hat{\theta} = \cos\theta\cos\theta_1 + \sin\theta\sin\theta_1\cos(\phi - \phi_1)$$

$$A_0 = \left\langle \frac{2(d-b) + 4b\sin^2\theta_1}{b+3d} \right\rangle,$$

$$A_2 = \left\langle \frac{4b\sin^2\theta_1\cos 2\phi_1}{b+3d} \right\rangle.$$

 $\langle P_l \left(\cos \theta, \phi \right) \rangle = \frac{\int P_l \left(\cos \theta, \phi \right) d\sigma d \cos \theta d\phi}{\int d\sigma d \cos \theta d\phi}$

J.C. Peng et al, PLB 758,384 (2016)



$$A_0 \neq A_2$$

- > Coplanarity case: $b \neq d$
- ➢ Non-coplanarity case: $\phi_1 ≠ 0$, NNLO and beyond or by the nonperturbative effects

Xu Li, Bin Yan, C.-P. Yuan, arxiv: 2405.04069



The discrepancy in Lam-Tung relation could be explained by electroweak dipole interactions (transversely polarized quark or lepton)

It could be more significant in high-invariant mass region



The accuracy from A0-A2 would be comparable to the results from cross section, but the violation effects will dominantly depend on the dipole interactions.

Xu Li, Bin Yan, C.-P. Yuan, arxiv: 2405.04069

Summary

- The violation of Lam-Tung relation in the Drell-Yan process can be emerged due to the non-coplanarity between the hadron plane and parton plane (NNLO and beyond)
- The discrepancy in ATLAS data could be explained by electroweak dipole interactions in coplanarity case

