



BESIII上粲介子半轻衰变到矢量介子和标量介子末态的研究

张书磊
湖南大学

中国格点量子色动力学研讨会
2024年10月13日@湖南师范大学

Email: zhangshulei@hnu.edu.cn

Content

01

Physics motivation ✓

02

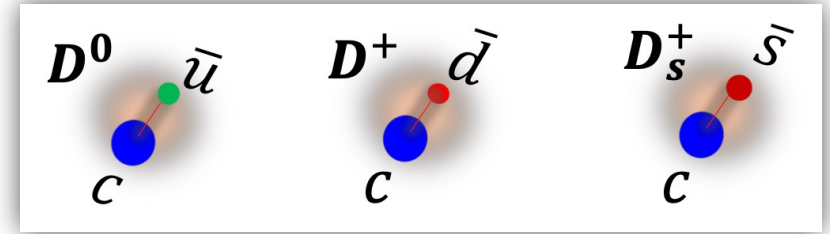
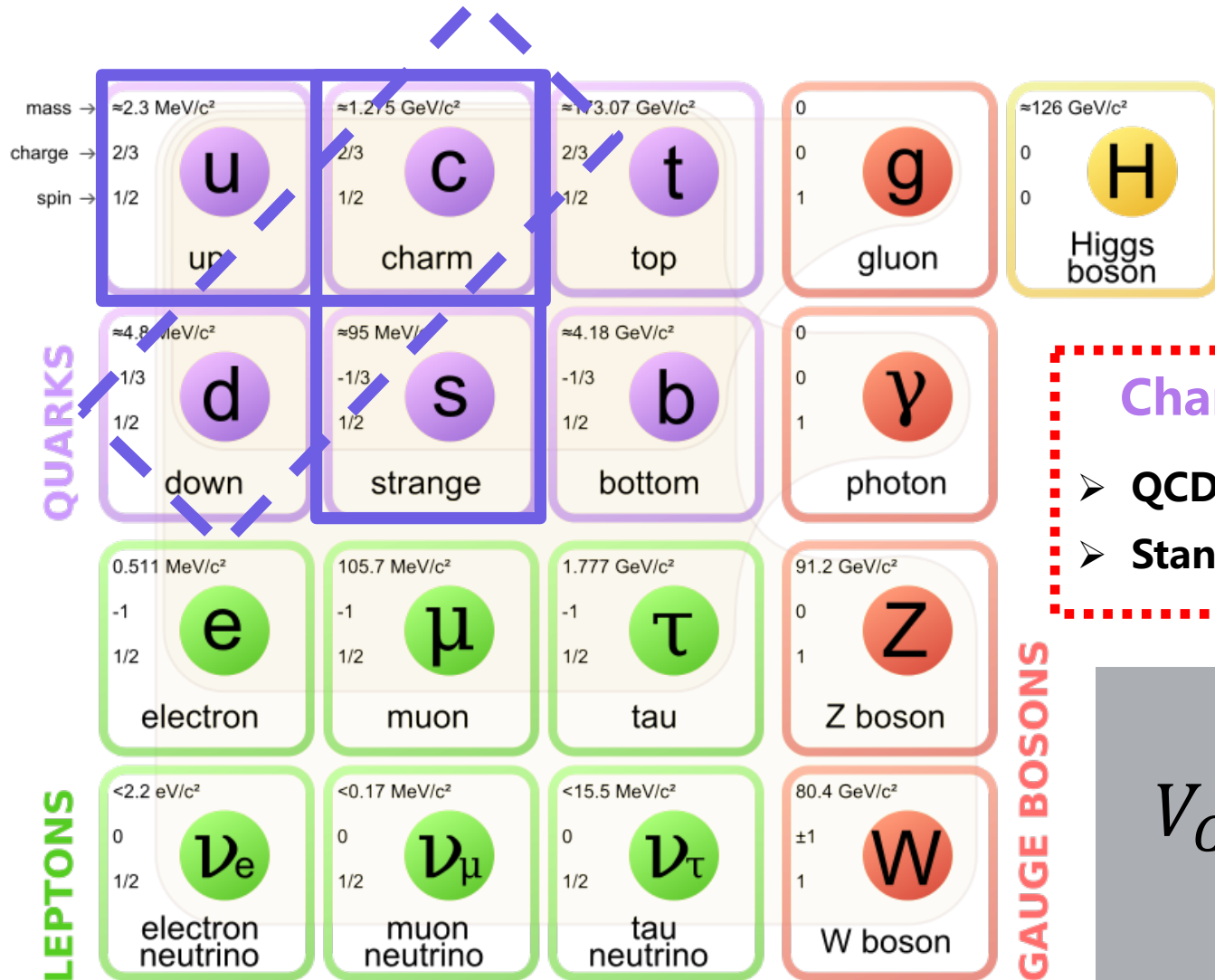
Data and analysis method

03

Results in review

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Summary and prospect



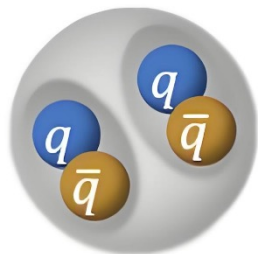
Charm physics

- **QCD**: Non-perturbative energy region
- **Standard Model test** : High-precision \rightarrow New Physics

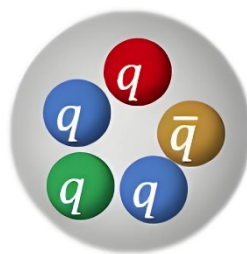
$$V_{CKM} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$



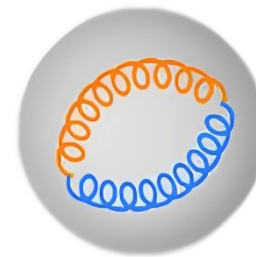
Tetraquark



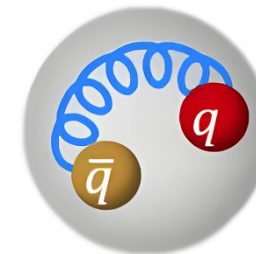
Hadronic molecules



Pentaquark



Glueball



hybrid

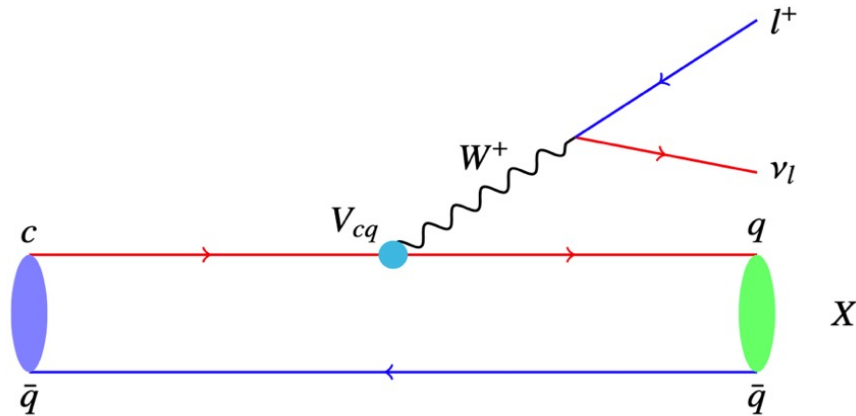
➤ Quark model allows for them.

➤ How about scalar mesons:

$f_0(500)$, $K^*(700)$, $f_0(980)$ and $a_0(980)$, etc

→ $q\bar{q}$ mixture, tetraquark, hadronic molecule or hybrid?

??? \Rightarrow Why is the semi-leptonic decay of charmed meson?



$$\Gamma(D_{(s)} \rightarrow S \ell^+ \nu_\ell) \propto |V_{cd(s)}|^2 |f_+(q^2)|^2 dq^2$$

$$\Gamma(D_{(s)} \rightarrow V \ell^+ \nu_\ell) \propto |V_{cd(s)}|^2 \mathfrak{T}(A_1(q^2), A_2(q^2), V(q^2)) dq^2$$

- **Clean environment:** hadrons X can be separated from leptons pair.
- **High statistics** of charmed meson at experiments.
 - CKM matrix elements $|V_{cs}|/|V_{cd}|$ measurements \rightarrow **Test the unitarity of the CKM matrix (New Physics).**
 - $\mathcal{R}_{\mu/e} = \mathcal{B}(D_{(s)} \rightarrow X \mu^+ \nu_\mu) / \mathcal{B}(D_{(s)} \rightarrow X e^+ \nu_e)$ measurement \rightarrow **Test lepton flavor universality (LFU)**
 - Hadronic Form factor (FF) measurements \rightarrow **Test different QCD models (LQCD)**

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➤ Symmetric e^+e^- collider @2 – 5GeV

➤ Pair-production near threshold

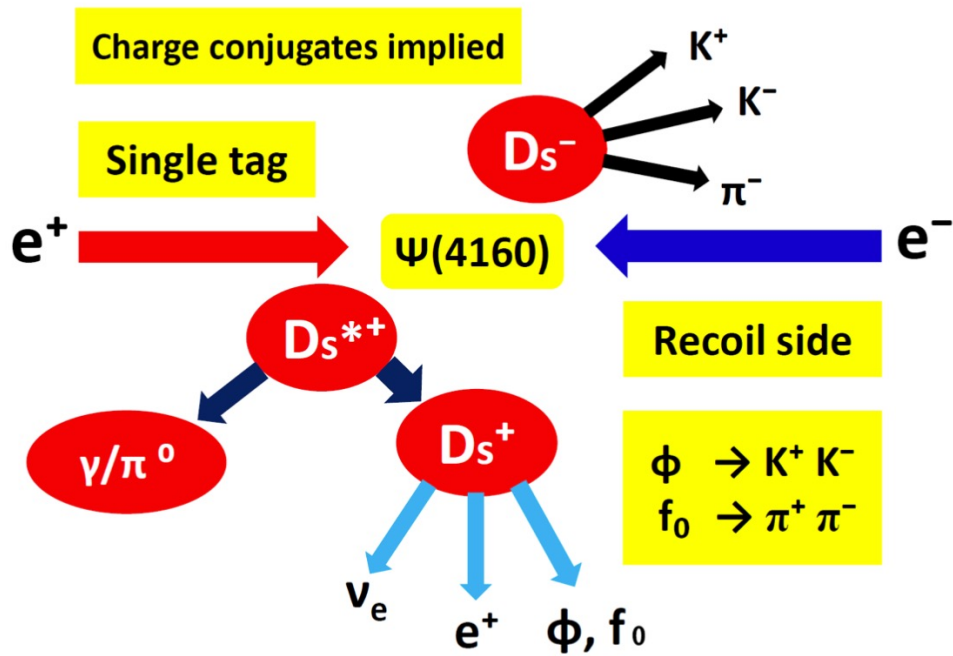
➤ $D\bar{D}$ @3.773GeV: **$\sim 20.3 \text{ fb}^{-1}$**

2.93 fb^{-1} @2010-2011; **4.99 fb^{-1} @2021-2022**; **8.16 fb^{-1} @2021-2022**; **4.19 fb^{-1} @2022-2024**

➤ $D_S D_S^*$ @4.13-4.23GeV: 7.33 fb^{-1}

$E_{\text{cm}}(\text{GeV})$	Data taking year	$\mathcal{L}(\text{fb}^{-1})$	ST D^0	ST D^+	ST D_S^+
3.773	2010-11 → 2022-24	2.93 → 20.3	2.7 M ($\sim 7\times$)	1.7 M ($\sim 7\times$)	0.8 M
4.13-4.23	2012, 2016-17, 2019	7.33			

Take Ds decay as an example (complicated case)



$B_\gamma(D_S^* \rightarrow \gamma D_S)$

$N_{tag} = 2N_{D_S^+ D_S^-} B_{tag} \epsilon_{tag}$

$N_{sig} = 2N_{D_S^+ D_S^-} B_{tag} B_{sig} B_\gamma \epsilon_{sig}$

$B_{sig} = \frac{N_{sig}}{B_\gamma N_{tag} \epsilon_{sig} / \epsilon_{tag}}$

$B_{sig} = \frac{N_{sig}}{B_\gamma \sum_\alpha N_{tag}^\alpha \epsilon_{sig}^\alpha / \epsilon_{tag}^\alpha}$

- Mature method
- Absolute BF measurement
- Low background
- Systematic cancellation (tag)

$U_{miss} = E_{miss} - |\vec{p}_{miss}|$

$M_{miss}^2 = E_{miss}^2 - |\vec{p}_{miss}|^2$

- Minimize the negative log-likelihood function.:

$$NLL = - \sum_{i=1}^N \ln \frac{\omega(\xi_i, \eta)}{\sigma_s}$$

$\omega(\xi_i, \eta)$ is the decay intensity, σ_s is the normalization factor, using signal Monte Carlo samples:

$$\sigma_s = \int d\xi \omega(\xi, \eta) \epsilon(\xi) \propto \frac{1}{N_{selected}} \sum_{k=1}^{N_{selected}} \frac{\omega(\xi_k, \eta)}{\omega(\xi_k, \eta_0)}$$

- When the background is low, it can be directly subtracted in the NLL. :

$$NLL = (-\ln L_{data}) - (-\ln L_{bkg})$$

- background is high :
$$-\sum_{i=1}^N \ln \left((1-f_b) \frac{\omega(\xi_i, \eta)}{\int d\xi_i \omega(\xi_i, \eta) \epsilon(\xi_i)} + f_b \frac{B_\epsilon(\xi_i)}{\int d\xi_i B_\epsilon(\xi_i) \epsilon(\xi_i)} \right)$$

➤ Background:

- High background level, especially for muon channels!
- How to describe background.

➤ Formula:

- Can the lepton mass be ignored? (e/μ)
- Amplitude formula with considering the lepton mass terms (include S/P/D wave).

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The differential decay rate of $D_{(s)} \rightarrow V \ell \nu_\ell$

$$\Gamma(D_{(s)} \rightarrow V \ell^+ \nu_\ell) \propto |V_{cd(s)}|^2 \mathfrak{I}(A_1(q^2), A_2(q^2), V(q^2), \dots) dm^2 dq^2 d\cos(\theta_h) d\cos(\theta_\ell) d\chi$$

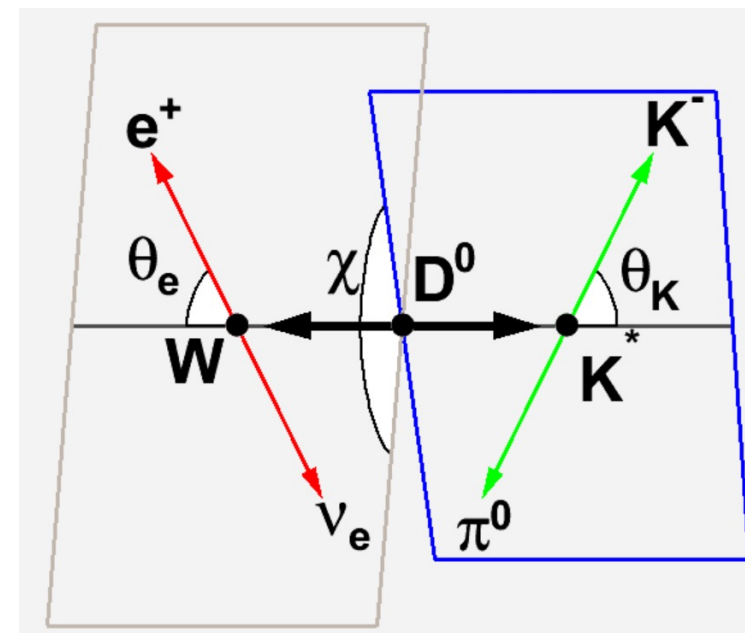
$$V: \rho, \omega, K^*, \phi$$

Theoretical : Phys. Rev. **137**, B438(1965), Phys.Rev.D 46,5040(1992)

- The decay intensity \mathfrak{I} can include components of S/P/D wave processes.
- Unbinned maximum likelihood method.
(implemented based on the RooFit framework).
- Form factor (single pole parameterization, double pole ...)

$$A_i(q^2) = \frac{A_i(0)}{1 - q^2/M_A^2} \quad V(q^2) = \frac{V(0)}{1 - q^2/M_V^2}$$

$$r_V = \frac{V(0)}{A_1(0)} \quad r_2 = \frac{A_2(0)}{A_1(0)}$$



Phys. Rev. D 94, 032001(2016)

- 2.93fb^{-1} data @ 3.773GeV
- $N_{sig} = 18262$ (background Level: 0.8%)
- $\mathcal{B}(D^+ \rightarrow K^- \pi^+ e^+ \nu_e) = (3.77 \pm 0.03 \pm 0.08)\%$

$$f_{S\text{-wave}} = (6.05 \pm 0.22 \pm 0.18)\%$$

- Form factor :

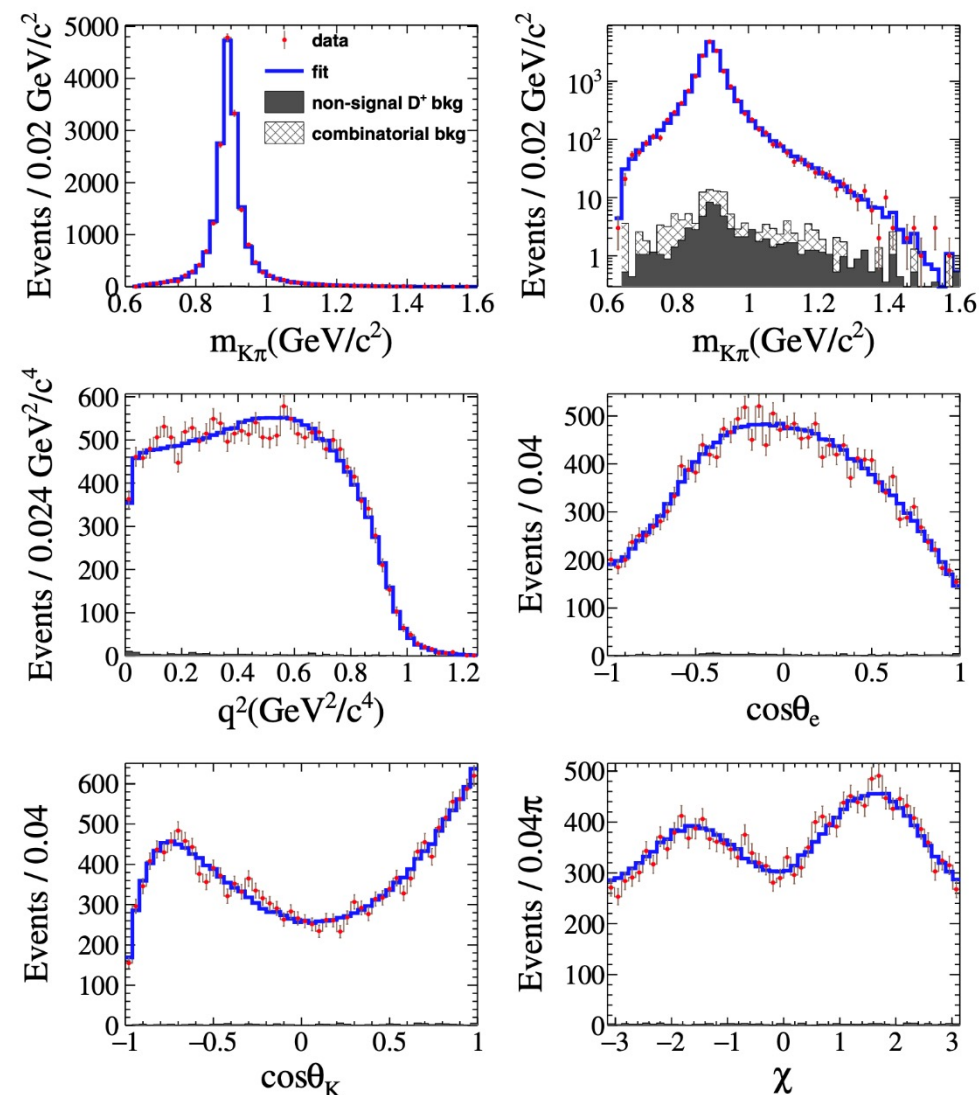
$$r_V = 1.411 \pm 0.058 \pm 0.007,$$

$$r_2 = 0.788 \pm 0.042 \pm 0.008$$

input $G_f, \tau_{D^+}, |V_{cs}| \rightarrow$

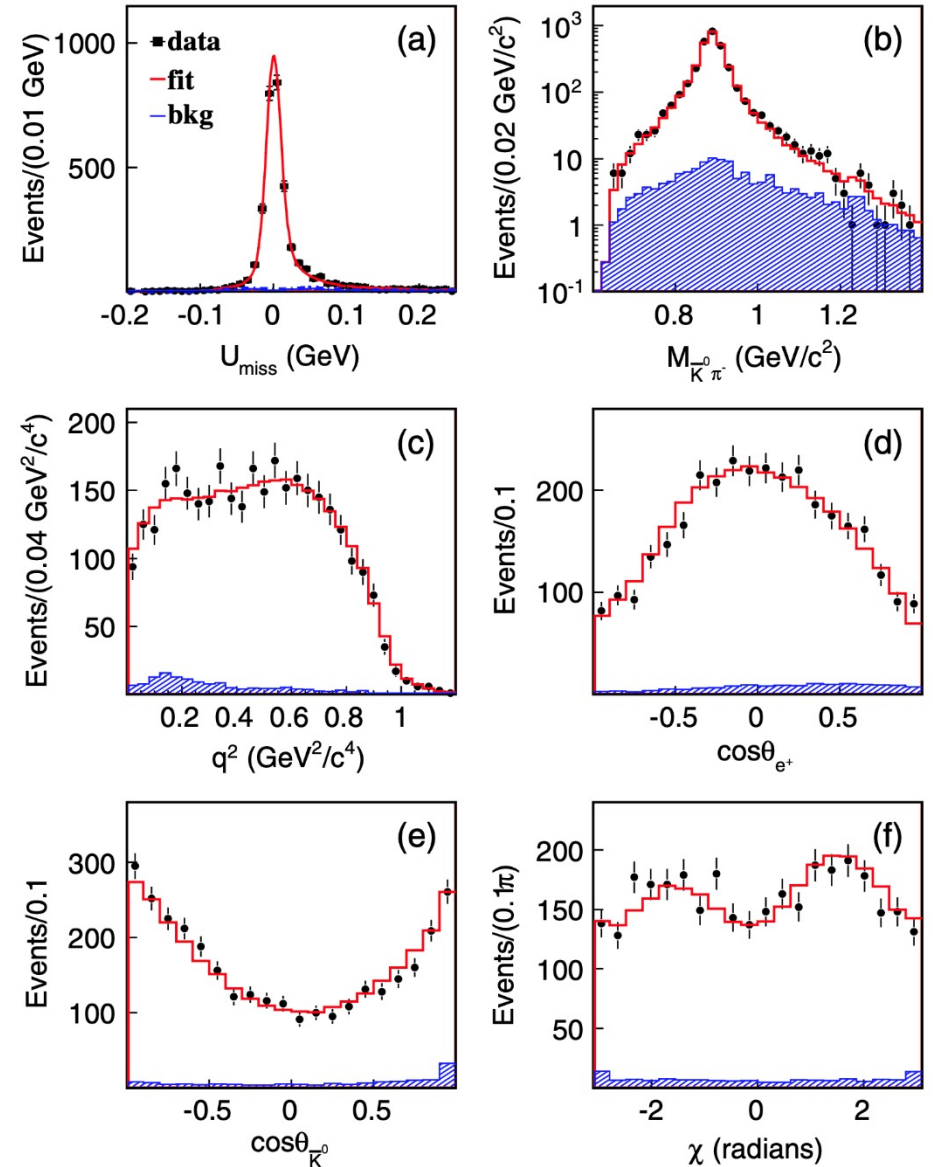
$$A_1(0) = 0.589 \pm 0.010 \pm 0.012 \text{ (Zero-width)}$$

$$A_1(0) = 0.619 \pm 0.011 \pm 0.013 \text{ (Considering width)}$$



Phys. Rev. D 99, 0111003(R)(2019)

- 2.93fb^{-1} data @3.773GeV
- $N_{sig} = 3112 \pm 64$ (background level: 0.6%)
- $\mathcal{B}(D^0 \rightarrow \bar{K}^0 \pi^+ e^+ \nu_e) = (1.434 \pm 0.029 \pm 0.032)\%$
 $f_{S\text{-wave}} = (5.51 \pm 0.97 \pm 0.62)\%$
- First measurement of the form factor :
 $r_V = 1.46 \pm 0.07 \pm 0.02,$
 $r_2 = 0.67 \pm 0.06 \pm 0.01$



JHEP12(2023)072

➤ 7.33 fb⁻¹ data @ 4.13-4.23 GeV

➤ $N_{\text{sig}} = 1725 \pm 68$ for BF measurement

$$\mathcal{B}(D_s^+ \rightarrow \phi \mu^+ \nu_\mu) = (2.25 \pm 0.09 \pm 0.07) \times 10^{-2}$$

$$\mathcal{B}(D_s^+ \rightarrow \phi \mu^+ \nu_\mu) / \mathcal{B}(D_s^+ \rightarrow \phi e^+ \nu_e) = 0.94 \pm 0.08 \rightarrow \text{No LFU violation}$$

$$\mathcal{B}(D_s^+ \rightarrow f_0(980) \mu^+ \nu_\mu) \cdot \mathcal{B}(f_0(980) \rightarrow K^+ K^-) < 5.45 \times 10^{-4} @ 90\% \text{ C.L. } (\sim 2.2\sigma)$$

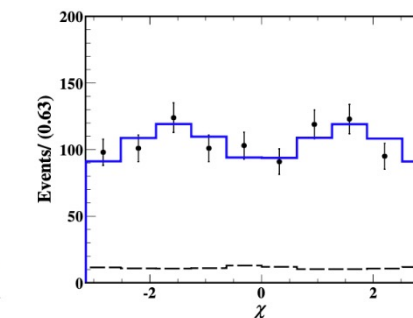
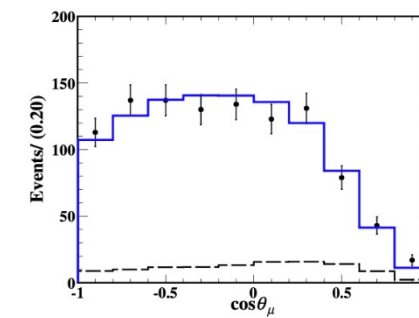
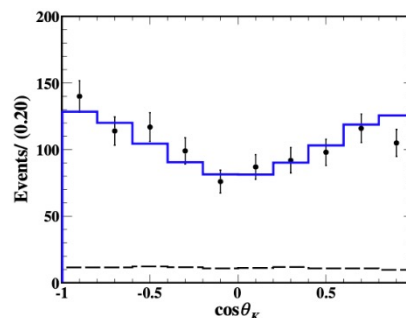
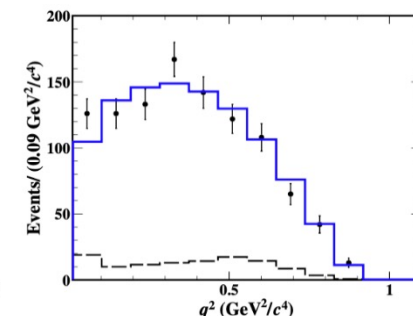
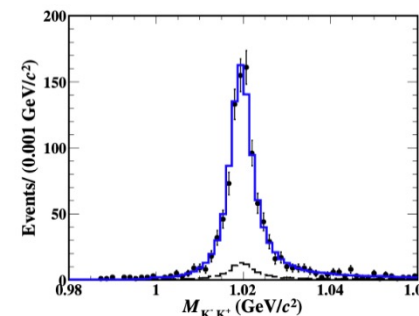
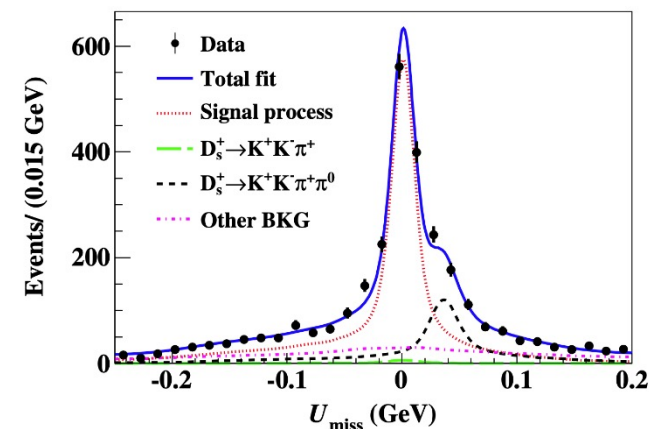
➤ First FF measurement based on single pole parameterization:

➤ Partial wave analysis is performed → ϕ dominate

➤ μ mass is considered in the formula

Table 5. Measured FF ratios and comparison with previous measurements.

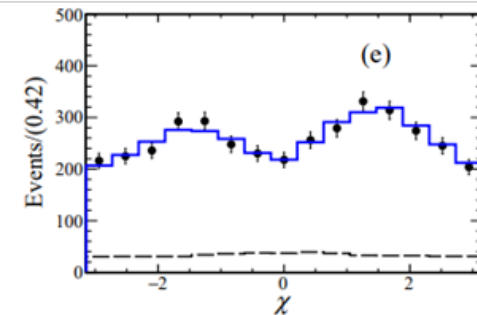
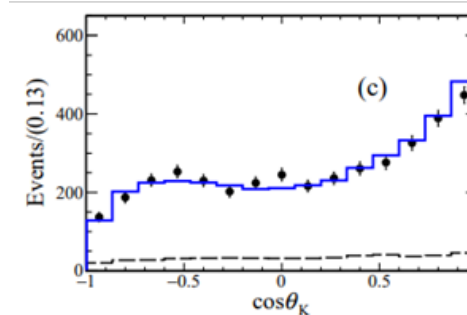
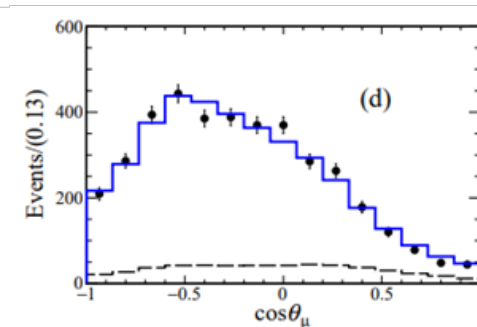
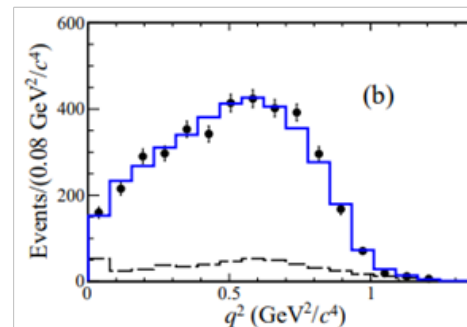
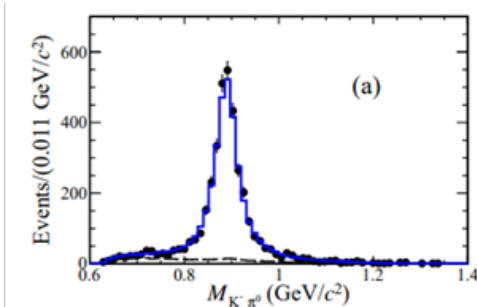
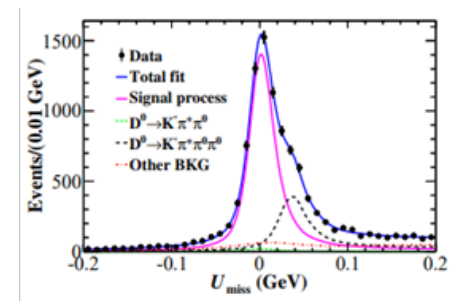
Experiments	r_V	r_2
PDG [42]	1.80±0.08	0.84±0.11
This analysis	1.58±0.17±0.02	0.71±0.14±0.02
BABAR [25]	1.807±0.046±0.065	0.816±0.036±0.030
FOCUS [58]	1.549±0.250±0.148	0.713±0.202±0.284
Theory	r_V	r_2
CCQM [5]	1.34±0.27	0.99±0.20
CQM [6]	1.72	0.73
LFQM [7]	1.42	0.86
LQCD [3]	1.72±0.21	0.74±0.12
HM χ T [8]	1.80	0.52



$|U_{\text{miss}}| < 0.02 \text{ GeV}$
 $N_{\text{Event}}^{D_s^+} = 939, \text{Bkg: } (9.8 \pm 0.7)\%$

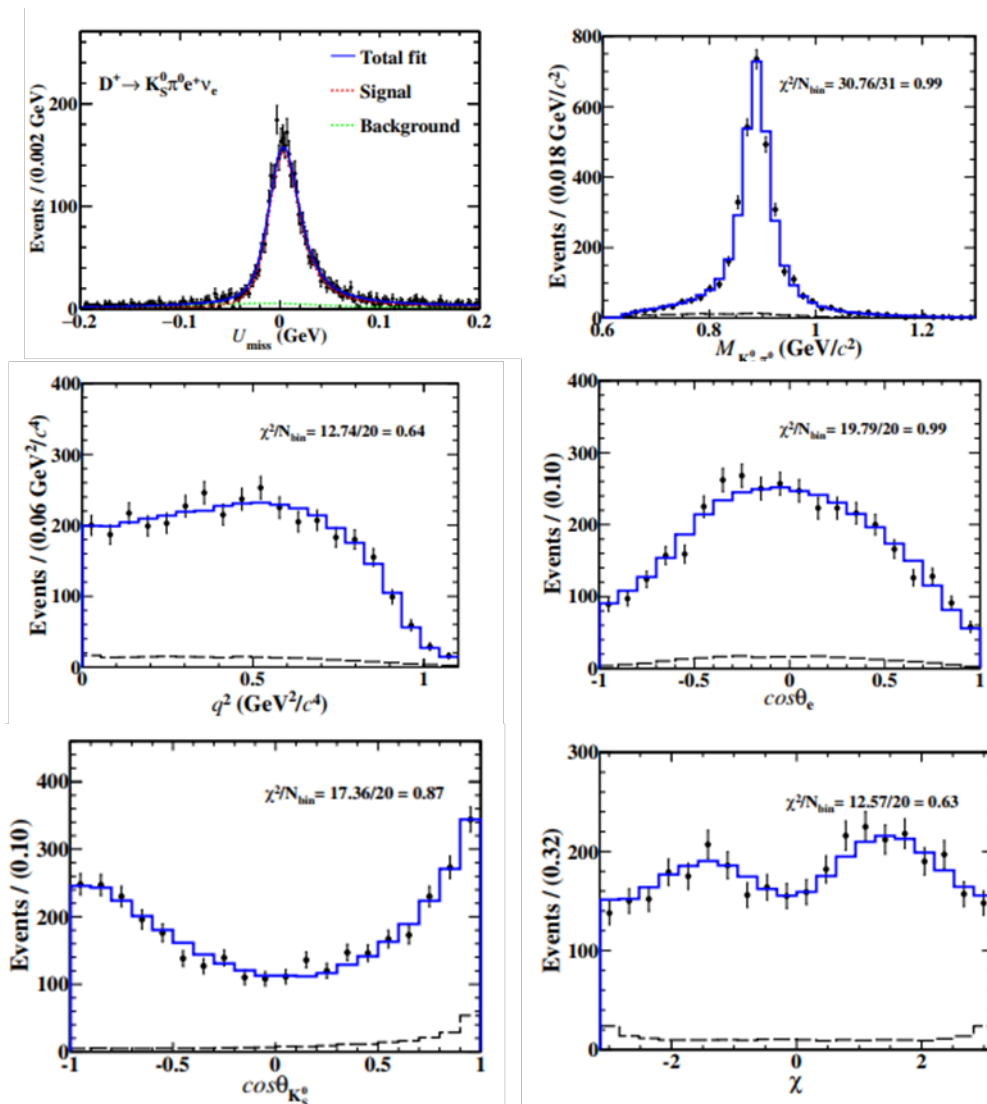
arXiv:2403.10877 (Submitted to PRL)

- 7.91 fb⁻¹ data @3.773GeV
- $N_{sig} = 6436 \pm 119$ (background level: 12.6%)
- $\mathcal{B}(D^0 \rightarrow K^- \pi^0 \mu^+ \nu_\mu) = (0.729 \pm 0.014 \pm 0.011)\%$
 $f_{S\text{-wave}} = (5.76 \pm 0.35 \pm 0.29)\%$
- $\mathcal{B}(D^0 \rightarrow K^{*-} \mu^+ \nu_\mu) / \mathcal{B}(D^0 \rightarrow K^{*-} e^+ \nu_e) = 0.96 \pm 0.08$
- First measurement of the form factor :
 $r_V = 1.37 \pm 0.09 \pm 0.03,$
 $r_2 = 0.76 \pm 0.06 \pm 0.02$



arXiv:2408.04422 (Accepted in JHEP)

- 7.91 fb⁻¹ data @3.773GeV
- $N_{sig} = 3852 \pm 75$ (background level: 6.54%)
- $\mathcal{B}(D^+ \rightarrow K_S^0 \pi^0 e^+ \nu_e) = (0.881 \pm 0.017 \pm 0.016)\%$
 $f_{S\text{-wave}} = (5.41 \pm 0.35 \pm 0.37)\%$
- Measurement of the form factor :
 $r_V = 1.43 \pm 0.07 \pm 0.03,$
 $r_2 = 0.72 \pm 0.06 \pm 0.02$

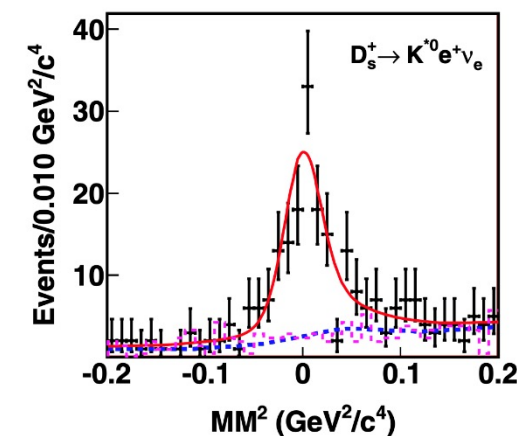
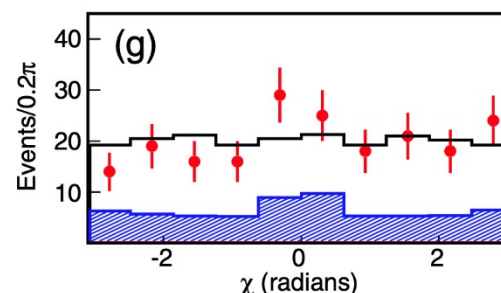
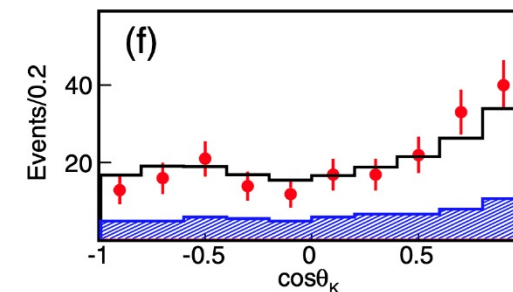
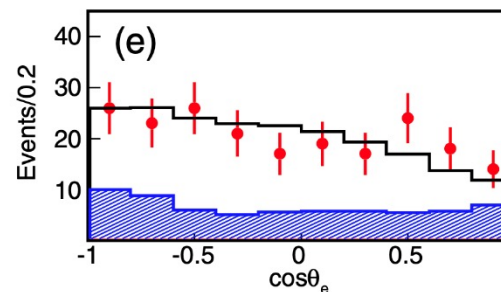
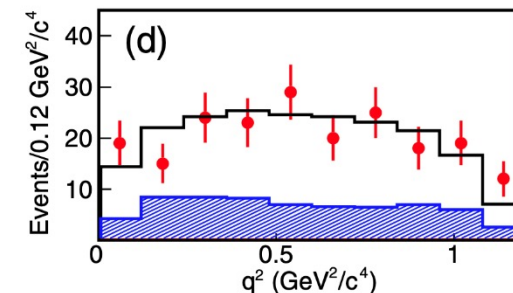
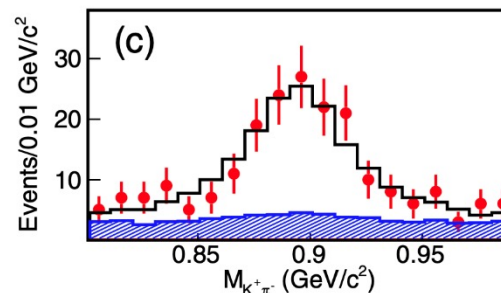


Phys. Rev. Lett. 122, 061801 (2019)

- 3.19 fb⁻¹ data @ 4.18 GeV
- $N_{sig}(D_s^+ \rightarrow K^{*0} e^+ \nu_e) = 155.0 \pm 17.2$
- $\mathcal{B}(D_s^+ \rightarrow K^{*0} e^+ \nu_e) = (2.37 \pm 0.26 \pm 0.20) \times 10^{-3}$
- Consistent with multiple theoretical calculations :
 $r_V = 1.67 \pm 0.34 \pm 0.16, r_2 = 0.77 \pm 0.28 \pm 0.07$
- Consistent with the expectations of Lattice QCD and U-spin
 (d ↔ s) symmetry: Form factors are insensitive to spectator quarks

Use **BESIII** and **CLEO** measurement

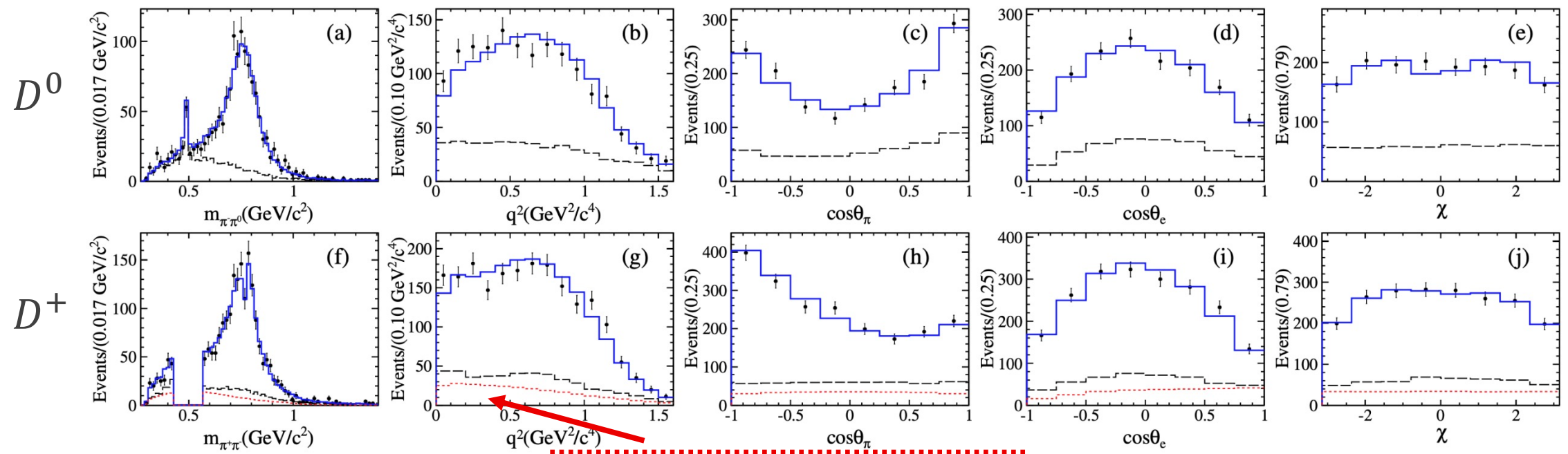
	Values
$f_+^{D_s^+ \rightarrow K^0}(0) / f_+^{D^+ \rightarrow \pi^0}(0)$	$1.16 \pm 0.14 \pm 0.02$
$r_V^{D_s^+ \rightarrow K^{*0}} / r_V^{D^+ \rightarrow \rho^0}$	$1.13 \pm 0.26 \pm 0.11$
$r_2^{D_s^+ \rightarrow K^{*0}} / r_2^{D^+ \rightarrow \rho^0}$	$0.93 \pm 0.36 \pm 0.10$



c -> d semi-leptonic decay : $D \rightarrow \rho e^+ \nu_e$

Phys. Rev. Lett. 122, 062001 (2019)

$|U_{\text{miss}}| < 0.06 \text{ GeV} \rightarrow$
 $N_{\text{Event}}^{D^0} = 1498, \text{Bkg: } (33.28 \pm 0.87)\%$
 $N_{\text{Event}}^{D^+} = 2017, \text{Bkg: } (23.82 \pm 0.69)\%$



➤ 2.93 fb⁻¹ data @ 3.773 GeV

➤ $R = \frac{B(D^+ \rightarrow f_0(500)e^+\nu_e) + B(D^+ \rightarrow f_0(980)e^+\nu_e)}{B(D^+ \rightarrow a_0(980)e^+\nu_e)} > 2.7 @ 90\% CL$

(W. Wang, PRD82, 034016(2010), Model independently way)

➔ BF results favor tetra-quark for f_0 and a_0

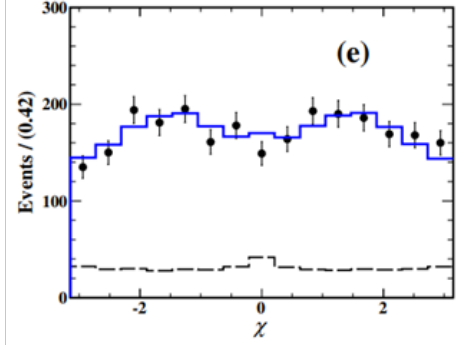
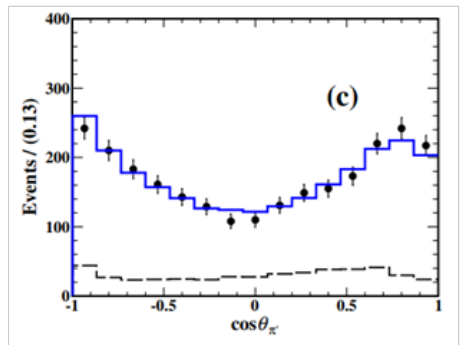
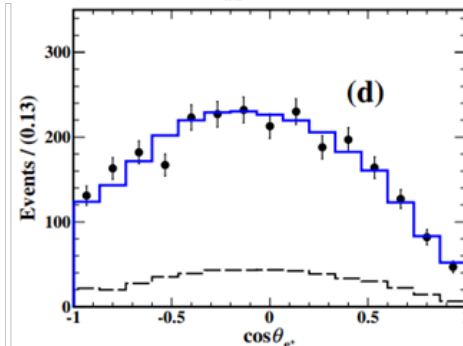
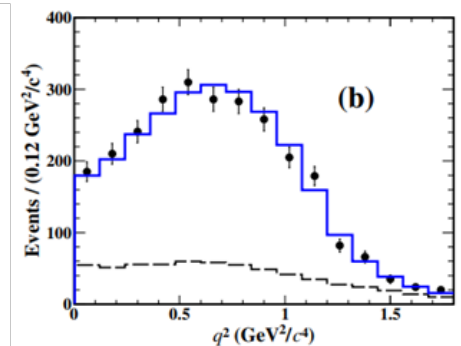
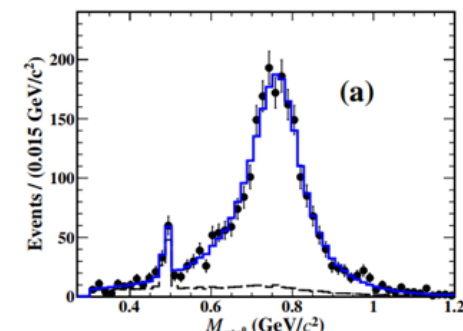
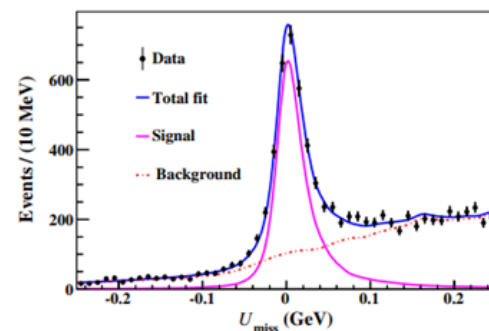
➤ $D \rightarrow \rho$: $r_V = 1.695 \pm 0.083 \pm 0.051, r_2 = 0.845 \pm 0.056 \pm 0.039$

$f_{f_0(500)} = (25.7 \pm 1.6 \pm 1.1)\%$

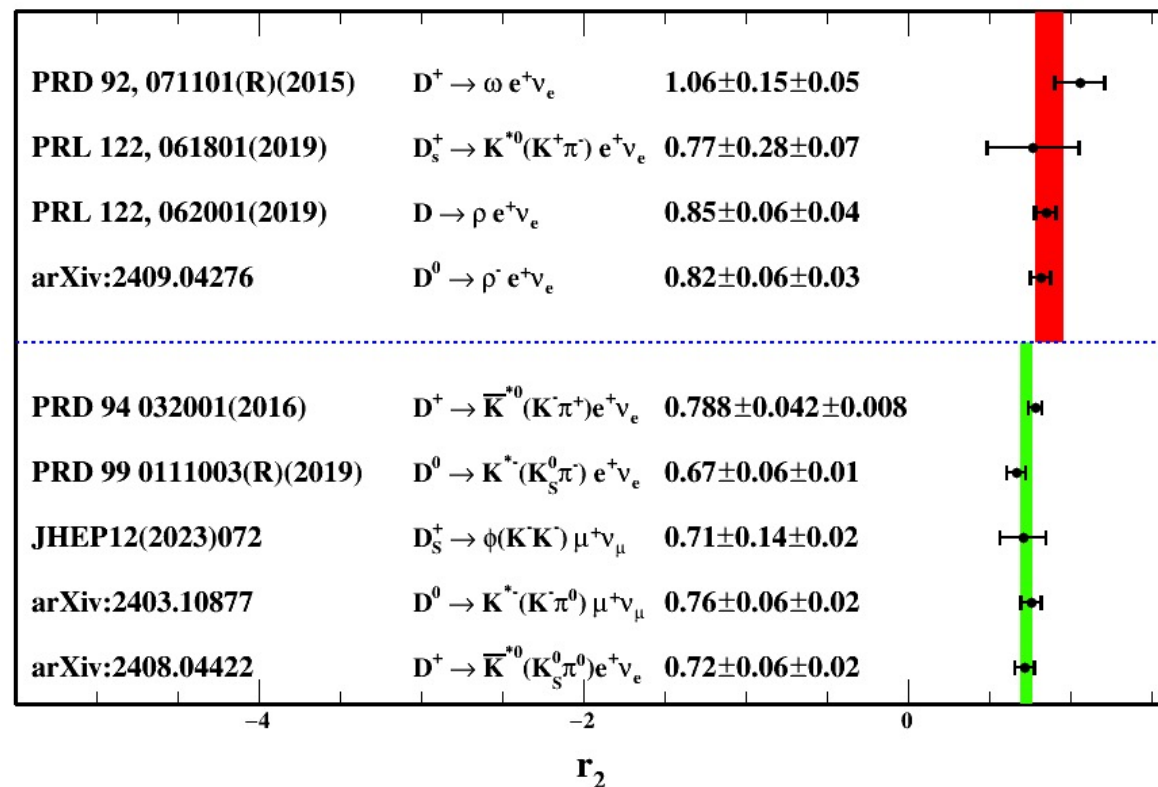
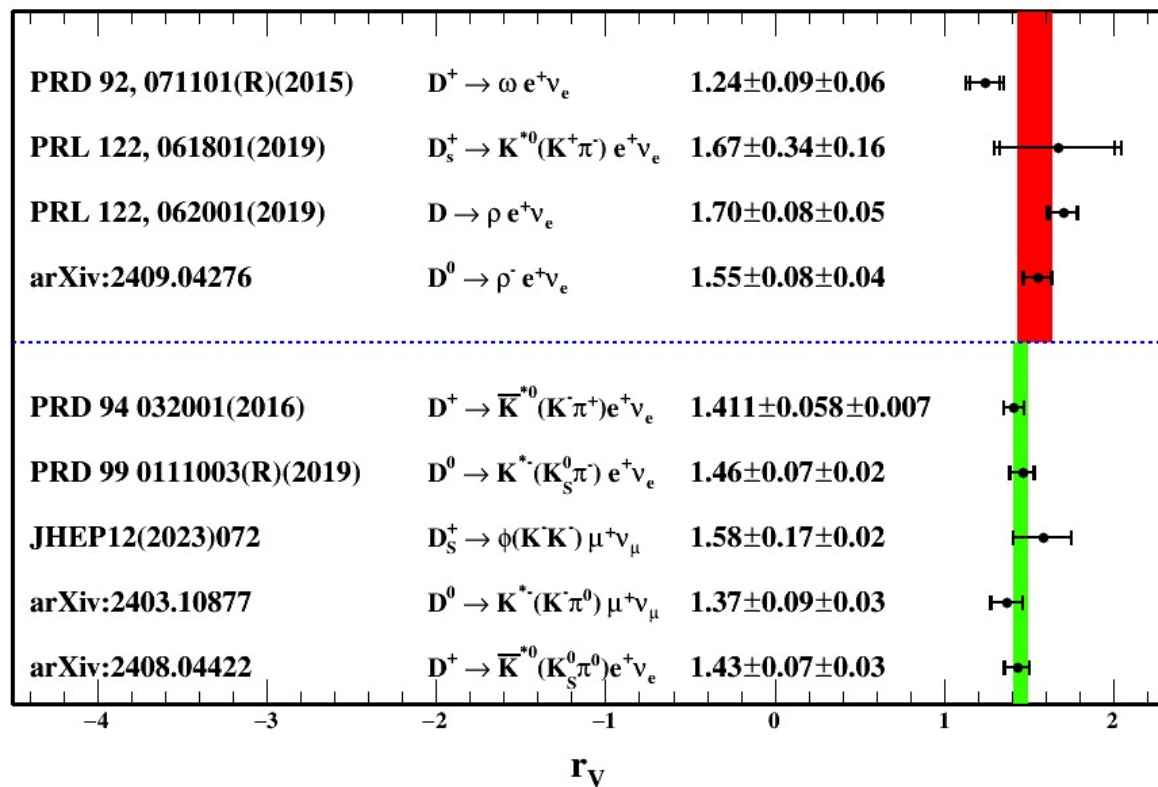
Signal mode	This analysis ($\times 10^{-3}$)
$D^0 \rightarrow \pi^- \pi^0 e^+ \nu_e$	$1.445 \pm 0.058 \pm 0.039$
$D^0 \rightarrow \rho^- e^+ \nu_e$	$1.445 \pm 0.058 \pm 0.039$
$D^+ \rightarrow \pi^- \pi^+ e^+ \nu_e$	$2.449 \pm 0.074 \pm 0.073$
$D^+ \rightarrow \rho^0 e^+ \nu_e$	$1.860 \pm 0.070 \pm 0.061$
$D^+ \rightarrow \omega e^+ \nu_e$	$2.05 \pm 0.66 \pm 0.30$
$D^+ \rightarrow f_0(500)e^+\nu_e, f_0(500) \rightarrow \pi^+\pi^-$	$0.630 \pm 0.043 \pm 0.032$
$D^+ \rightarrow f_0(980)e^+\nu_e, f_0(980) \rightarrow \pi^+\pi^-$	< 0.028

arXiv:2409.04276 (Submitted to PRD)

- 7.91 fb⁻¹ data @3.773GeV
- $N_{sig} = 3337 \pm 77$ (background level: 18.7%)
- $\mathcal{B}(D^0 \rightarrow \rho^- e^+ \nu_e) = (1.439 \pm 0.033 \pm 0.027) \times 10^{-3}$
- Form factor measurement:
 - $r_V = 1.548 \pm 0.079 \pm 0.041,$
 - $r_2 = 0.823 \pm 0.056 \pm 0.026$



Comparisons of r_V and r_2



$$\Gamma(D_{(s)} \rightarrow S \ell^+ \nu_\ell) / dq^2 \propto |V_{cd(s)}|^2 |f_+(q^2)|^2$$

$$S: a_0(980), f_0(500), f_0(980)$$

- Use least χ^2 method to fit the measured partial decay width in different q^2 bin.
- Taking the correlations among q^2 bins into account.
- FF in different form (The width needs to be considered ?)

– **Single pole form**

$$f_+(q^2) = \frac{f_+(0)}{1 - q^2/M_{pole}^2}$$

– **Modified pole model**

$$f_+(q^2) = \frac{f_+(0)}{\left(1 - \frac{q^2}{M_{pole}^2}\right) \left(1 - \alpha \frac{q^2}{M_{pole}^2}\right)}$$

– **ISGW2 model**

$$f_+(q^2) = f_+(q_{max}^2) \left(1 + \frac{r^2}{12} (q_{max}^2 - q^2)\right)^{-2}$$

– **Series expansion model**

$$f_+(t) = \frac{1}{P(t)\Phi(t, t_0)} a_0(t_0) \left(1 + \sum_{k=1}^{\infty} r_k(t_0) [z(t, t_0)]^k\right)$$

The differential decay rate of $D_{(s)} \rightarrow S \ell \nu_\ell$

➤ Point-like differential decay rate:

$$\frac{d\Gamma(D_{(s)} \rightarrow S \ell^+ \nu_\ell)}{dq^2} = \frac{G_F^2 |V_{cd(s)}|^2}{24\pi^3} p_{f_0}^3 |f_+(q^2)|^2$$

➤ Double differential decay rate:

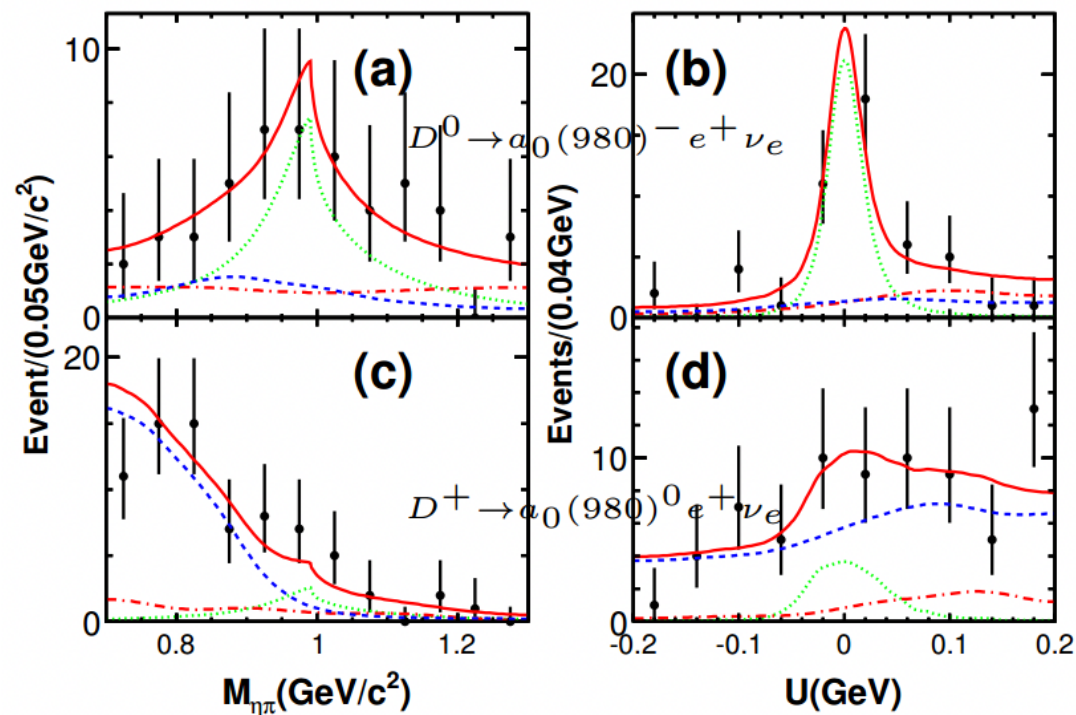
(N.N.Achasov *et al.*, PRD102,016022(2020); W. Wang, PLB759,501(2016))

$$\frac{d^2\Gamma(D_{(s)} \rightarrow S \ell^+ \nu_\ell)}{dsdq^2} = \frac{G_F^2 |V_{cd(s)}|^2}{192\pi^4 m_{D_{(s)}}^3} \lambda^{\frac{3}{2}} \left(m_{D_{(s)}}^2, s, q^2 \right) |f_+(q^2)|^2 P(s)$$

$$P(s) = \begin{cases} \frac{g_1 \rho_{\pi\pi/\pi\eta}}{|m_0^2 - s - i(g_1 \rho_{\pi\pi/\pi\eta} + g_1 \rho_{KK})|^2}, & \text{Flatte: } f_0(980)/a_0(980) \\ \frac{m_{f_0} \Gamma(s)}{(s - m_{f_0}^2)^2 + m_{f_0}^2 \Gamma^2(s)}, & \text{RBW: } f_0(500) \\ \frac{m_r \Gamma_{tot}(s)}{(m_r^2 - s - g_1^2 \frac{s - s_A}{m_r^2 - s_A} z(s))^2 + m_r^2 \Gamma_{tot}^2(s)}, & \text{Bugg: } f_0(500) \end{cases}$$

Phys. Rev. Lett. 121, 081802 (2018)

- 2.93 fb⁻¹ data @ 3.773 GeV
- $N_{\text{sig}}^{D^0} = 25.7^{+6.4}_{-5.7}$
- $N_{\text{sig}}^{D^+} = 10.2^{+5.0}_{-4.1}$
- Branching fraction (BF) help to understand the nature of the $a_0(980)$



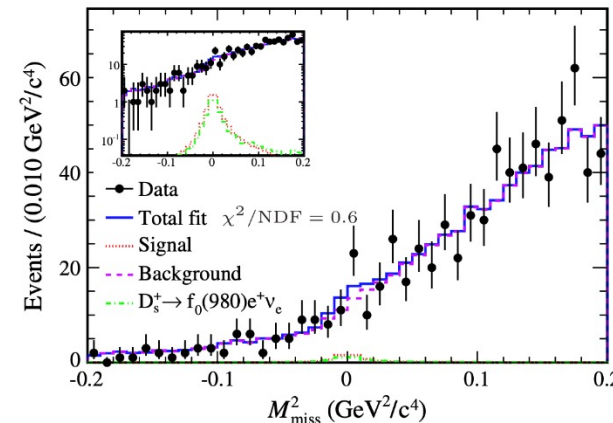
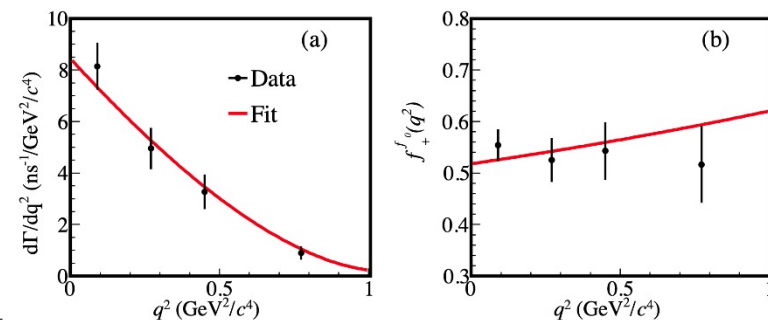
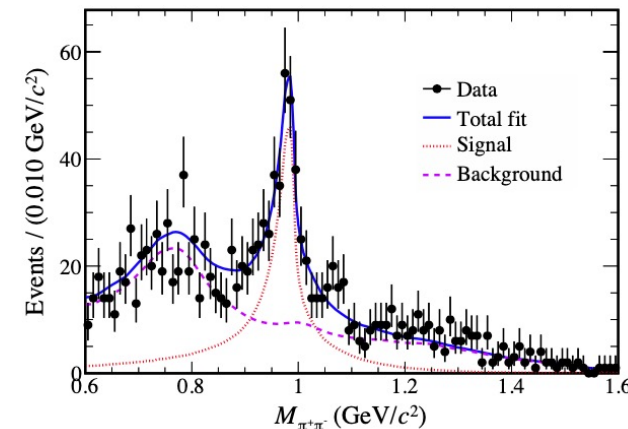
Decay	BF ($\times 10^{-4}$)	Significance
$D^0 \rightarrow a_0(980)^- e^+ \nu_e, a_0(980)^- \rightarrow \eta\pi^-$	$1.33^{+0.33}_{-0.29} \pm 0.09$	6.4σ
$D^+ \rightarrow a_0(980)^0 e^+ \nu_e, a_0(980)^0 \rightarrow \eta\pi^0$	$1.66^{+0.81}_{-0.66} \pm 0.11$ < 3.0 (90% C.L.)	2.9σ

Phys. Rev. Lett. 132, 141901 (2024)

- 7.33 fb^{-1} data @ 4.128-4.226 GeV $\rightarrow N_{\text{sig}} = 439 \pm 33$
- $\mathcal{B}(D_s^+ \rightarrow f_0(980)e^+ \nu_e, f_0(980) \rightarrow \pi^+ \pi^-) = (1.72 \pm 0.13 \pm 0.10) \times 10^{-3}$
 \rightarrow **$s\bar{s}$ is dominant** based on $|f_0(980)\rangle = \sin \phi |\frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d})\rangle + \cos \phi |s\bar{s}\rangle$
 $\phi = (19.7 \pm 12.8)^\circ$
- **First form factor measurement with simple pole form:**
 $\rightarrow f_+^{f_0}(0) |V_{cs}| = 0.504 \pm 0.017 \pm 0.035$
 $\rightarrow f_+^{f_0}(0) = 0.518 \pm 0.018 \pm 0.036$ ($|V_{cs}| = 0.97349 \pm 0.00016$ PDG2022)

	This work	CLFD [6]	DR [6]	QCDSR [7]	QCDSR [8]	LCSR [9]	LFQM [11]	CCQM [12]
$f_+^{f_0}(0)$	$0.518 \pm 0.018_{\text{stat}} \pm 0.036_{\text{syst}}$	0.45	0.46	0.50 ± 0.13	0.48 ± 0.23	0.30 ± 0.03	0.24 ± 0.05	0.36 ± 0.02
Difference (σ)	—	1.7	1.4	0.1	0.2	4.3	4.3	2.8
ϕ	$\phi = (19.7 \pm 12.8)^\circ$	$(32 \pm 4.8)^\circ$	$(41.3 \pm 5.5)^\circ$	35°	$(8_{-8}^{+21})^\circ$	—	$(56 \pm 7)^\circ$	31°

- **First search of $D_s^+ \rightarrow f_0(500)e^+ \nu_e, f_0(500) \rightarrow \pi^+ \pi^-$ ($M_{\pi^+ \pi^-} < 0.45 \text{ GeV}/c^2$)**
- $\mathcal{B}(D_s^+ \rightarrow f_0(500)e^+ \nu_e, f_0(500) \rightarrow \pi^+ \pi^-) < 3.3 \times 10^{-4}$

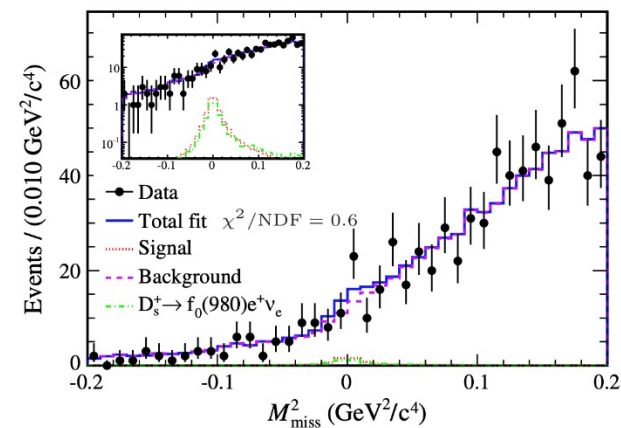
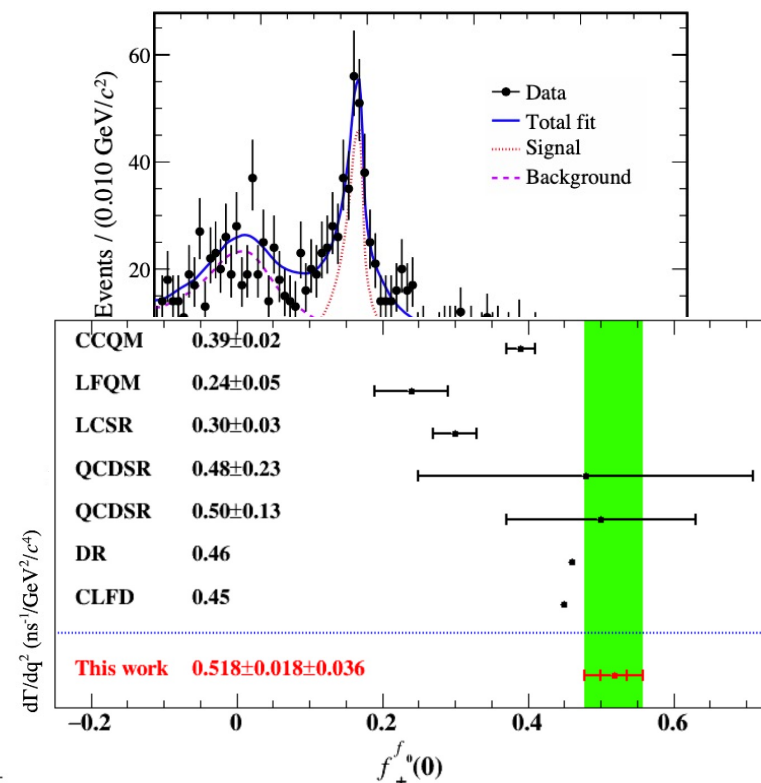


Phys. Rev. Lett. 132, 141901 (2024)

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Study of the decay $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$

arXiv: 2401.13225 (Accepted in PRD)

- 2.93 fb⁻¹ data @ 3.773 GeV
- First observation of $D^+ \rightarrow f_0(500)(\pi^+\pi^-)\mu^+\nu_\mu$.

Signal mode	N_{Obs}	S (σ)	ϵ_{sig} (%)	$\mathcal{B}_{\text{sig}} (\times 10^{-3})$
$f_0(500)\mu^+\nu_\mu$	209 ± 38	5.9	18.93 ± 0.13	0.72 ± 0.13
$\rho^0\mu^+\nu_\mu$	496 ± 38	> 10	19.86 ± 0.13	1.64 ± 0.13
$f_0(500)e^+\nu_e$	412 ± 43	> 10	44.76 ± 0.25	0.60 ± 0.06
$\rho^0e^+\nu_e$	1237 ± 47	> 10	44.12 ± 0.25	1.84 ± 0.07

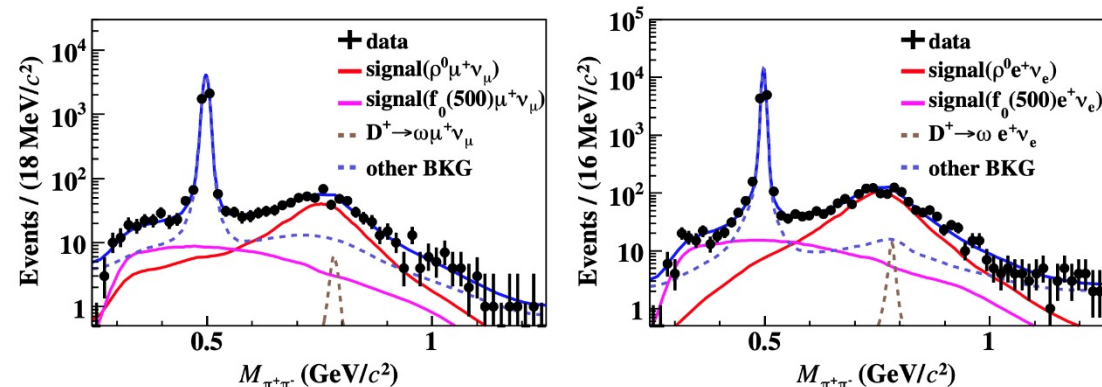
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Based Z series expansion for FF and Bugg form for $f_0(500)$

$$\rightarrow f_+^{f_0}(0)|V_{cd}| = 0.0787 \pm 0.0060 \pm 0.0033$$

$$\rightarrow f_+^{f_0}(0) = 0.350 \pm 0.027 \pm 0.015$$

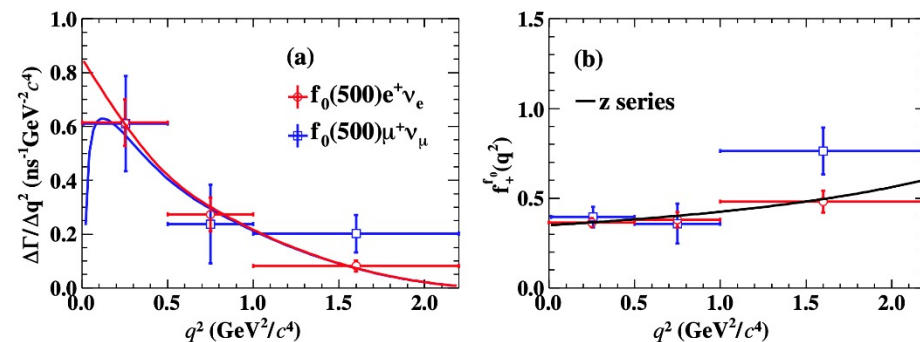
ps: $|V_{cd}| = 0.22438 \pm 0.00044$ from SM global fit (PDG2022)



➔ The measured BF of $D^+ \rightarrow f_0(500)\ell^+\nu_\ell$ are closer to **tetraquark assumption**.

R.M. Wang et al, PRD107,056022 (2023)

Y.K. Hsiao et al, arXiv:2306.06091



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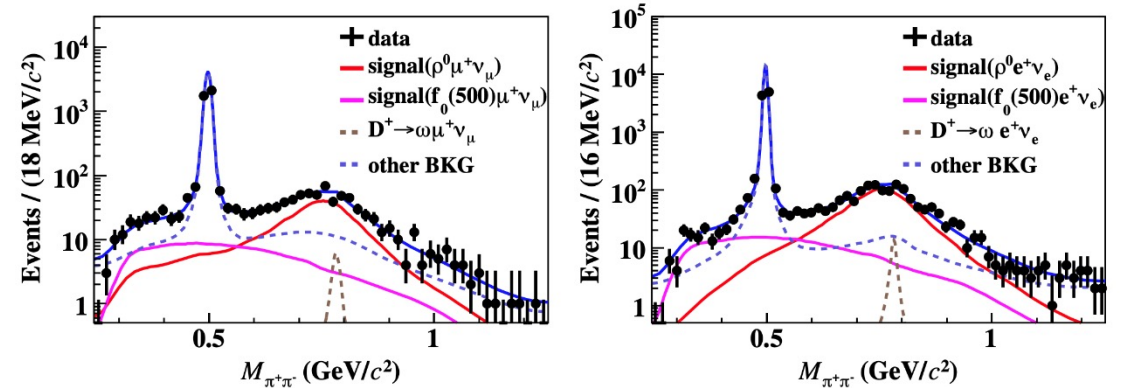
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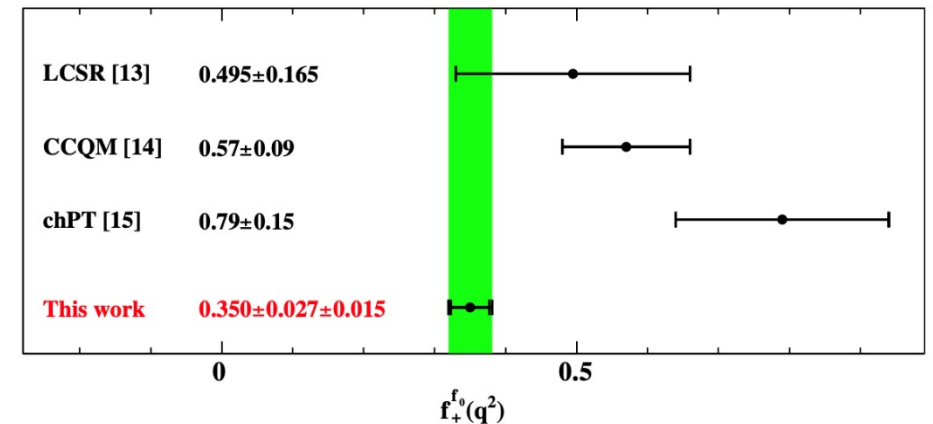
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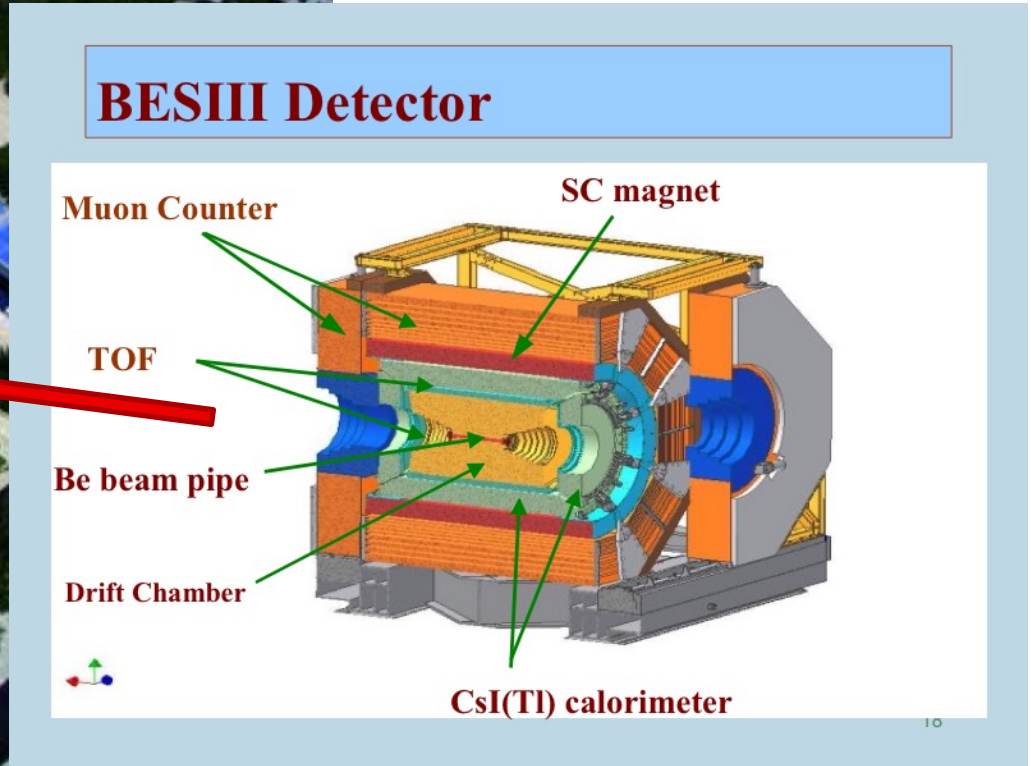
➤ Summary

BESIII have performed high-precision measurements on charm meson semi-leptonic decays to vector mesons (ρ, ω, K^*, ϕ) and scalar mesons ($f_0(500), f_0(980), a_0(980)$) using the unique advantage of charm meson pair produced at threshold .

- Absolute branching fraction measurements using double-tagging method
- Hadronic form factors measurements through amplitude analysis.
- Can help to measure CKM matrix elements $|V_{cs}|/|V_{cd}|$
- Help to understand the nature of light scalar mesons

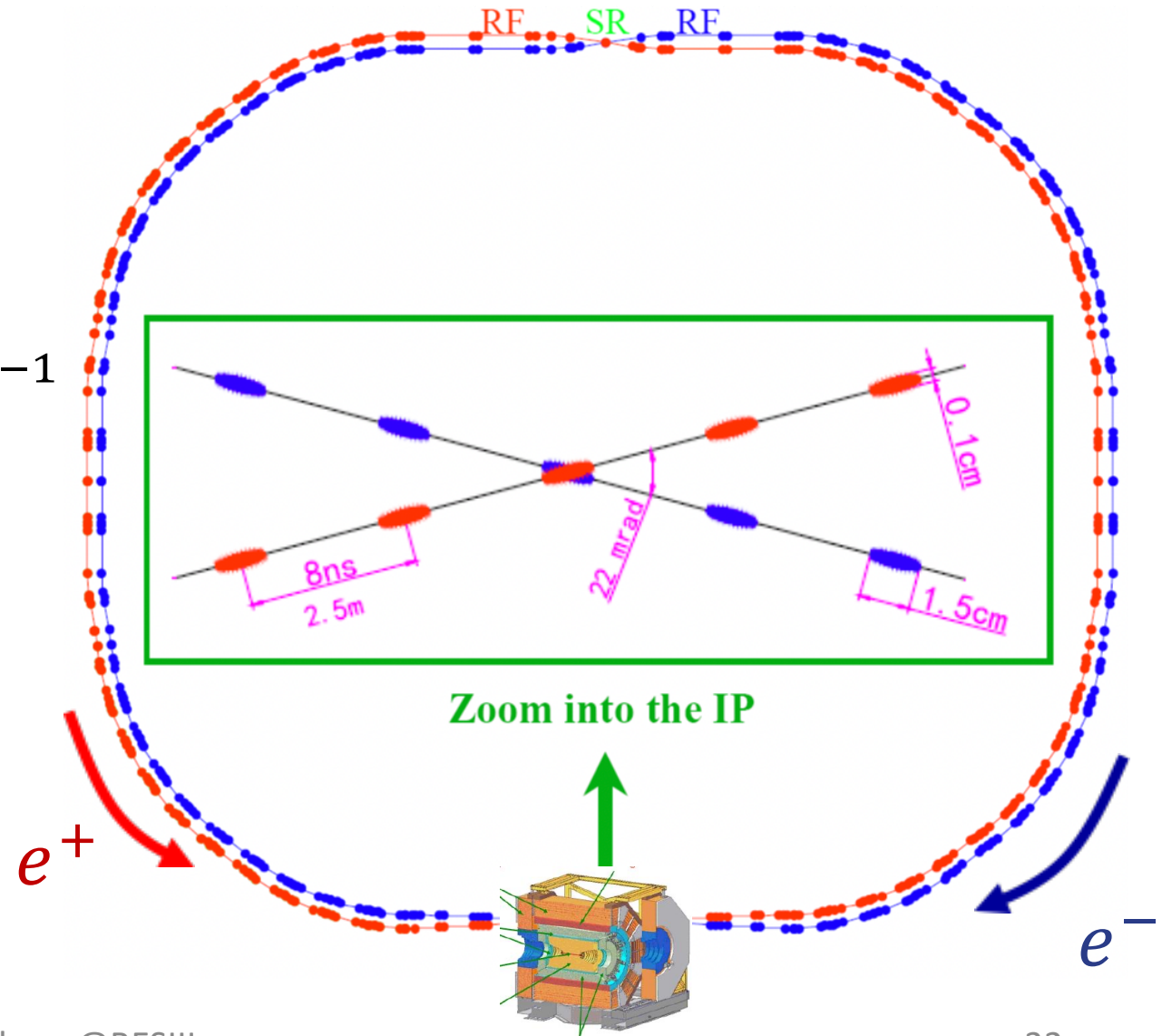
➤ prospect

- BESIII has $\sim 20 \text{ fb}^{-1}$ @3.773 GeV in total now.
- More studies are on the way, especially the muon channels.
- More jointed measurements of multiple channels

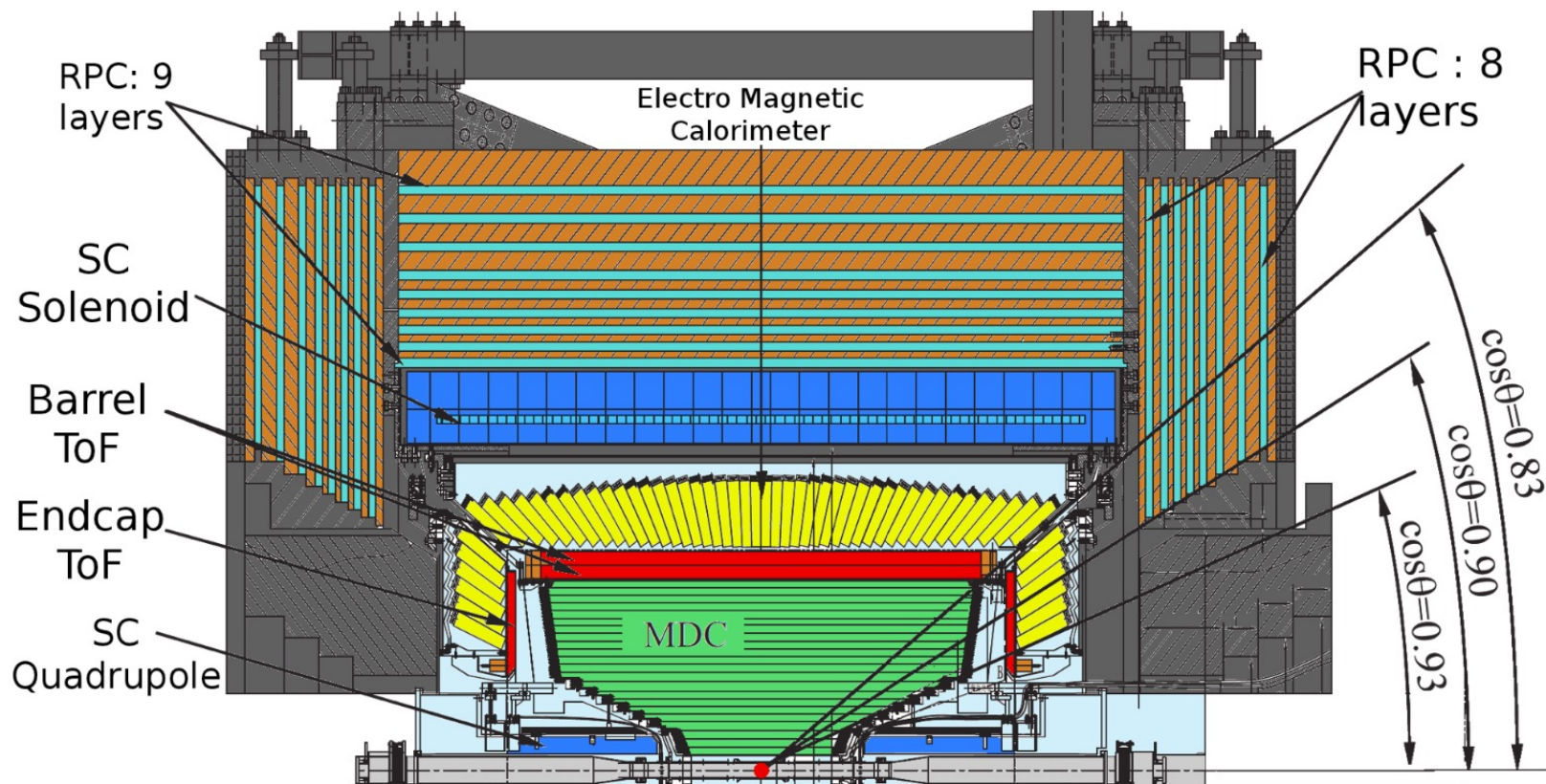


BEPCII collider

- Two ring symmetric e^+e^- collider
- Circumference: 240 m
- Design luminosity: $1 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$
- Achieved time: 5 April, 2016
- E_{cm} : 2 – 5 GeV
- Beam crossing angle: 22 mrad



BESIII detector



MDC

$$\frac{\delta p}{p} < 0.5\% \text{ @1 GeV}$$

$$\frac{\delta(dE/dx)}{dE/dx} < 6\%$$

TOF

$$\delta t \text{ 80 ps Barrel}$$

$$\delta t \text{ 110 ps Endcap}$$

EMC

$$\frac{\delta E}{E} < 2.5\% \text{ @1 GeV}$$

$$\delta z = 0.6/\sqrt{E}$$

MUC

$$\delta(xy) < 2 \text{ cm}$$