

Topological properties in electromagnetic fields

e-Print: [2312.14660](#) [hep-lat]

Mamiya

This is proceeding paper in lattice simulation.

Electromagnetic effects on topological observables in QCD

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This paper shows preliminary results about...

- Axion-photon coupling (without external electromagnetic fields)
- Magnetic field effect on topological susceptibility at low temperatures.

What are the new findings in this paper?

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- **Axion-photon coupling** (without external electromagnetic fields) (3 color)
- **Magnetic field effect on topological susceptibility** at low temperatures.

This study is the first attempt to evaluate the electromagnetic effect on topological property using lattice simulation.

Evaluation of Axion-photon coupling based on lattice is first attempt.

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- **Magnetic field effect on topological susceptibility** at low temperatures.

This study is the first attempt to evaluate the electromagnetic effect on topological property using lattice simulation.

In other situations regarding “topological susceptibility in lattice simulation”,

Temperature effect
(3 color)

- C. Bonati et al, JHEP 11, 170 (2018), 1807.07954.
- S. Borsanyi et al., Nature 539, no. 7627, 69 (2016).
- P. Petreczky et al, Phys. Lett. B 762, 498-505 (2016).

Density effect
(2 color)

- Itou et al. JHEP01(2020)181
- N. Astrakhantsev et al. PRD 102 (2020) 7, 074507

Magnetic effect

No previous research

QCD Lagrangian with θ term

$$\mathcal{L}_{\text{QCD}} = \bar{q}_L i \gamma^\mu D_\mu q_L + \bar{q}_R i \gamma^\mu D_\mu q_R - \bar{q}_L \mathbf{m}_f q_R - \bar{q}_R \mathbf{m}_f q_L + \frac{1}{2} \sum_{a=1}^8 \text{tr} [(G_{\mu\nu}^a T_c^a)^2] + \theta \frac{g^2}{64\pi^2} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu}^a G_{\rho\sigma}^a$$

U(1) axial anomaly

$$\partial_\mu j_A^{(f)\mu} = 2i \bar{q}^f m_f \gamma_5 q^f + \frac{g^2}{32\pi^2} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu}^a G_{\rho\sigma}^a + N_c \frac{e^2 [Q_{\text{em}}^f]^2}{32\pi^2} \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu} F_{\rho\sigma}$$

θ is promoted to a dynamical field (axion).

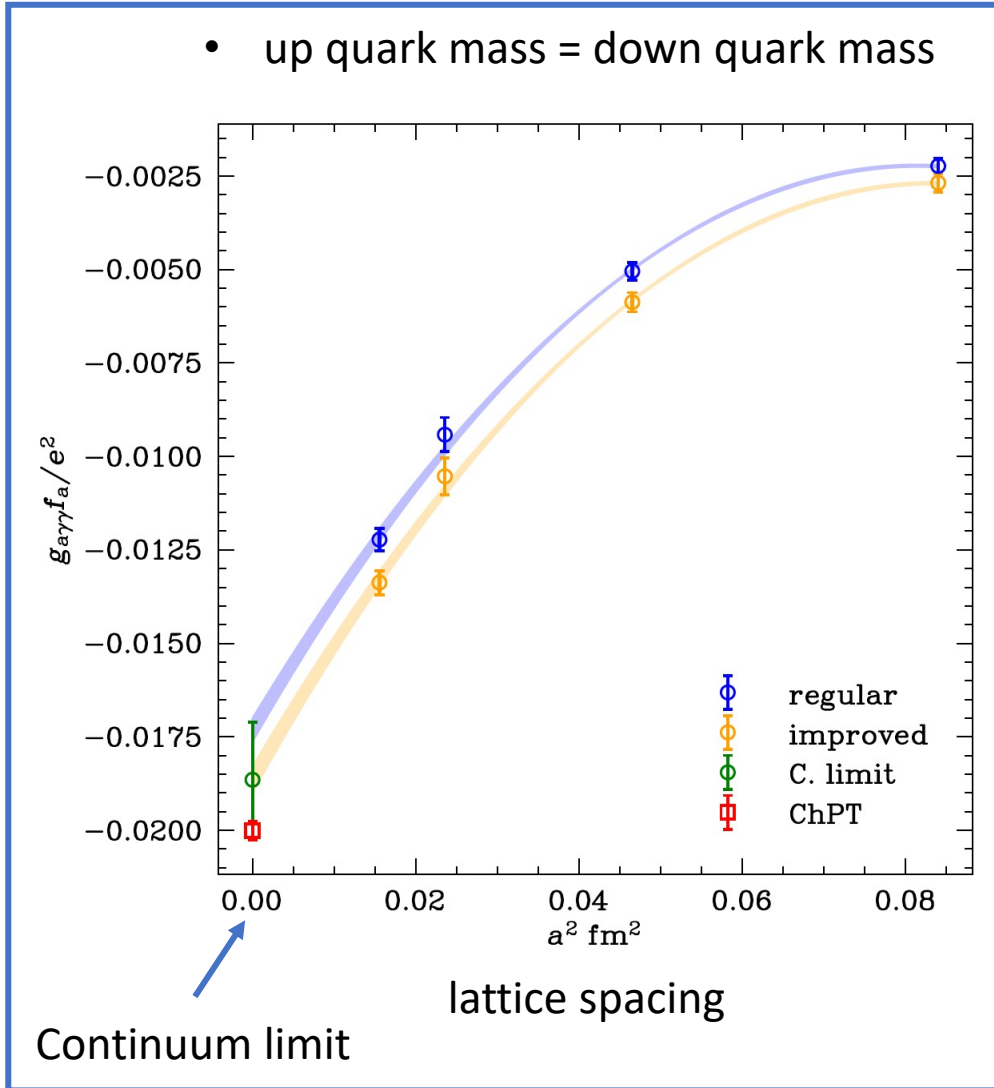


Axion-photon coupling

$$g_{a\gamma\gamma}^{\text{QCD}} f_a = \frac{i}{\Omega} \frac{\partial^2}{\partial\theta\partial(\vec{E} \cdot \vec{B})} \ln Z_{\text{QCD}+\theta+\text{bEM}} \Big|_{\theta=\vec{E}=\vec{B}=0} \quad \Omega \text{ is space-time volume.}$$

This coupling has been numerically evaluated in e-Print: [2312.14660](https://arxiv.org/abs/2312.14660) [hep-lat].

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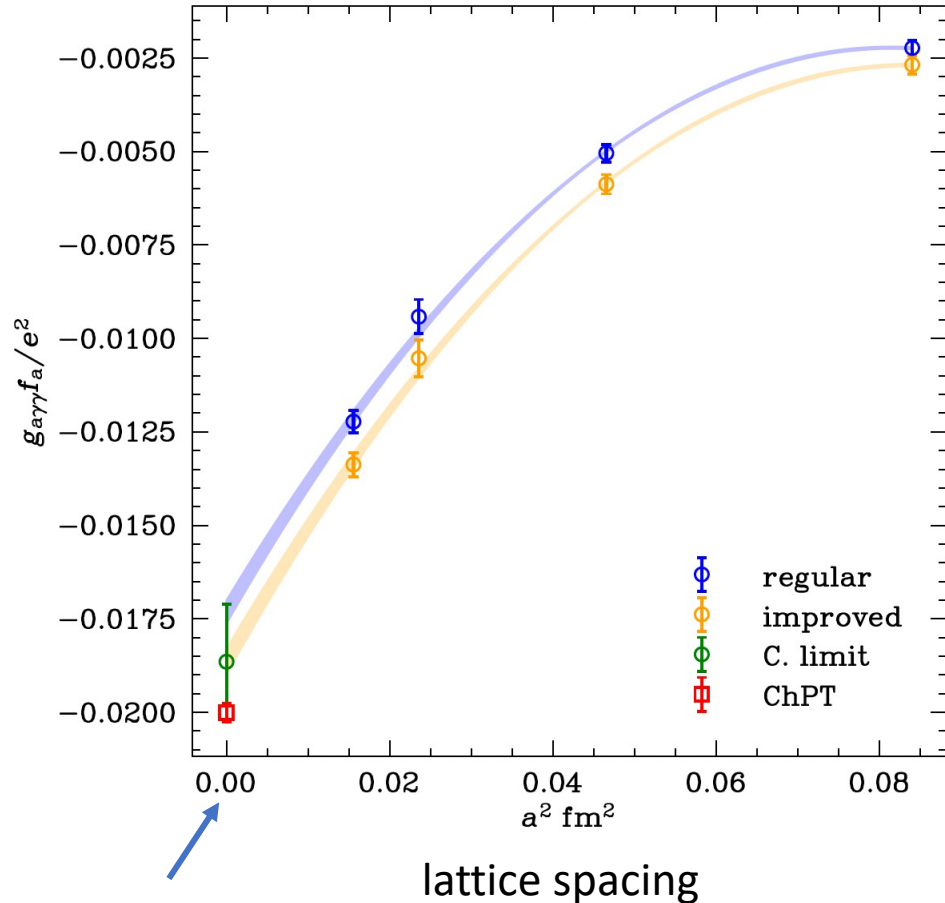
- up quark mass is not equal to down quark mass

$$g_{\gamma\gamma}^{\text{QCD}} f_a / e^2 = -0.023(2).$$

at continuum limit with physical quark mass.

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- Lattice results tell us **magnitude/sign** of axion-photon coupling:

$$g_{a\gamma\gamma}^{\text{QCD}} f_a / e^2 \sim O(-0.02)$$

- Lattice results takes similar value in ChPT.

In this week, I realized that the coupling constant can be evaluated based on model independent way.

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$$Z_{\text{QCD}} = \int [dq d\bar{q}] [dG] \exp \left[i \int d^4x \mathcal{L}_{\text{QCD}} \right]$$



U(1) axial transformation

$$Z_{\text{QCD}} = \int [dq' d\bar{q}'] [dG] \exp \left[i \int d^4x \left(\bar{q}'_L i \gamma^\mu D_\mu q'_L + \bar{q}'_R i \gamma^\mu D_\mu q'_R + \frac{1}{2} \sum_{a=1}^8 \text{tr} [(G_{\mu\nu}^a T_c^a)^2] \right) \right. \\ \left. - \sum_f \left(\bar{q}'_L m_f \exp \left(i \frac{\bar{m}}{m_f} \theta \right) q'_R + \bar{q}'_R m_f \exp \left(-i \frac{\bar{m}}{m_f} \theta \right) q'_L \right) \right]$$

$$- \left(\sum_f \frac{\bar{m}}{m_f} \theta Q_f^2 \right) N_c \frac{e^2}{64\pi^2} \epsilon^{\mu\nu\rho\sigma} F_{\mu\nu} F_{\rho\sigma} \Bigg]$$



$$g_{a\gamma\gamma}^{\text{QCD}} f_a / e^2 = - \left[\frac{\bar{m}}{m_u} q_u^2 + \frac{\bar{m}}{m_d} q_d^2 + \frac{\bar{m}}{m_s} q_s^2 \right] N_c \frac{1}{32\pi^2} 8$$

I got the analytical expression.

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My model independent result is comparable to the lattice evaluation.

$$g_{a\gamma\gamma}^{\text{QCD}} f_a / e^2 = -0.0209$$



$$m_l = 3.45 \text{ MeV}$$

$$m_s = 93.4 \text{ MeV}$$

$$\bar{m} = \left(\sum_f \frac{1}{m_f} \right)^{-1}$$

Model independent result
(My result)

$$g_{a\gamma\gamma}^{\text{QCD}} f_a/e^2 = -0.0209$$

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The coupling comes from **the U(1) anomaly associated with electromagnetic fields**.

Such anomaly is the one-loop exact result. → The coupling is exactly fixed. (Anomaly matching)



Axion-photon coupling is model independent.



Is it really necessary to implement lattice simulations/effective model analysis?
→ Maybe no...

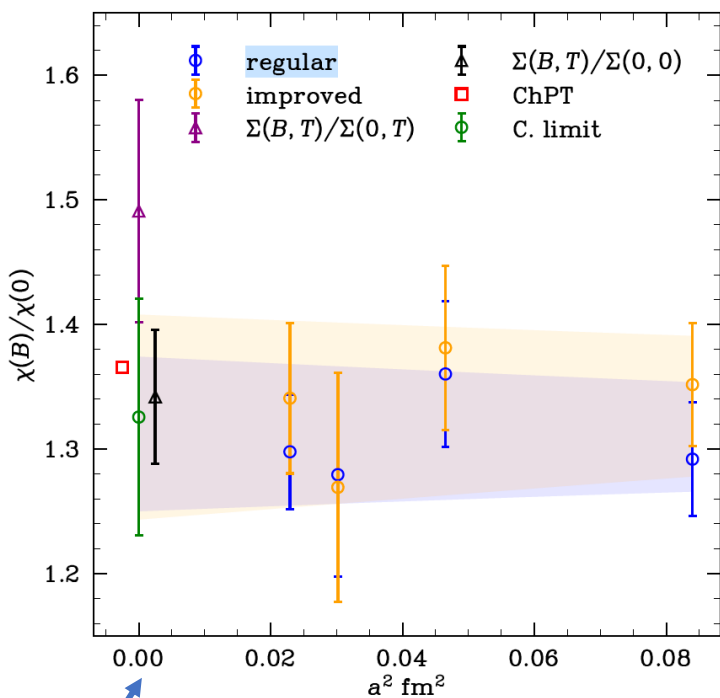
This paper also shows the magnetic effect on the topological susceptibility.

Definition

$$\chi_{\text{top}} = -\frac{1}{\Omega} \left. \frac{\partial^2 \ln Z_{\text{QCD}+\theta}}{\partial \theta^2} \right|_{\theta=0} = \frac{\langle Q_{\text{top}}^2 \rangle_{\text{QCD}}}{\Omega}$$

$$\begin{aligned} \chi_{\text{top}} &= -i \int d^4x \left\langle 0 \left| T \left(\frac{g^2}{64\pi^2} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu}^a G_{\rho\sigma}^a \right) (x) \left(\frac{g^2}{64\pi^2} \epsilon^{\mu\nu\rho\sigma} G_{\mu\nu}^b G_{\rho\sigma}^b \right) (0) \right| 0 \right\rangle \\ &= -\bar{m}^2 \left[\sum_f \frac{1}{m_f} \langle 0 | \bar{q}^f q^f | 0 \rangle + i \int d^4x \left\langle 0 \left| T \left(\sum_f \bar{q}^f i\gamma_5 q^f \right) (x) \left(\sum_{f'} \bar{q}^{f'} i\gamma_5 q^{f'} \right) (0) \right| 0 \right\rangle \right] \end{aligned}$$

Lattice observation



$eB = 0.5 \text{ GeV}^2$
 $T = 113 \text{ MeV}$

$$\frac{\chi_{\text{top}}(eB = 0.5 \text{ GeV}^2)}{\chi_{\text{top}}(B = 0)} > 1$$

Magnetic field enhances the topological susceptibility.



It seems...

Magnetic field is “catalyzer” for the topological property?



But

The lattice data is too limited. → More results are needed to clarify the magnetic property.

Summary

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$$g_{a\gamma\gamma}^{\text{QCD}} f_a/e^2 \sim O(-0.02)$$

I also obtained the coupling based on model independent way.

$$g_{a\gamma\gamma}^{\text{QCD}} f_a/e^2 = -0.0209$$

$$\frac{\chi_{\text{top}}(eB = 0.5\text{GeV}^2)}{\chi_{\text{top}}(B = 0)} > 1$$

Magnetic field enhances the topological susceptibility.

In my future work, I will address this topic using NJL model with Irfan.