

**Sub-GeV Dark Matter  
and  
Collective Excitation Signals**

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2401.11971

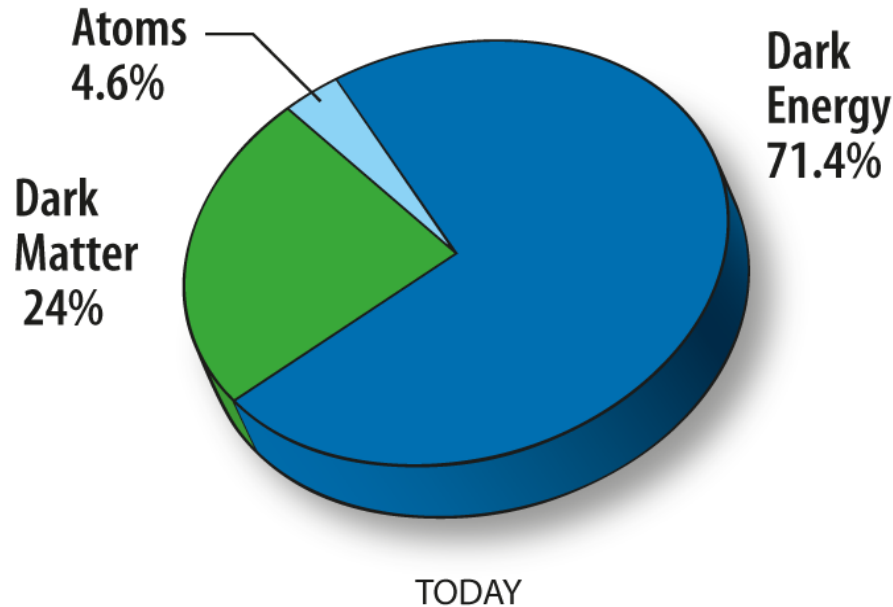
*in collaboration with L. Liang, L. Su, B. Zhu*

# Outline

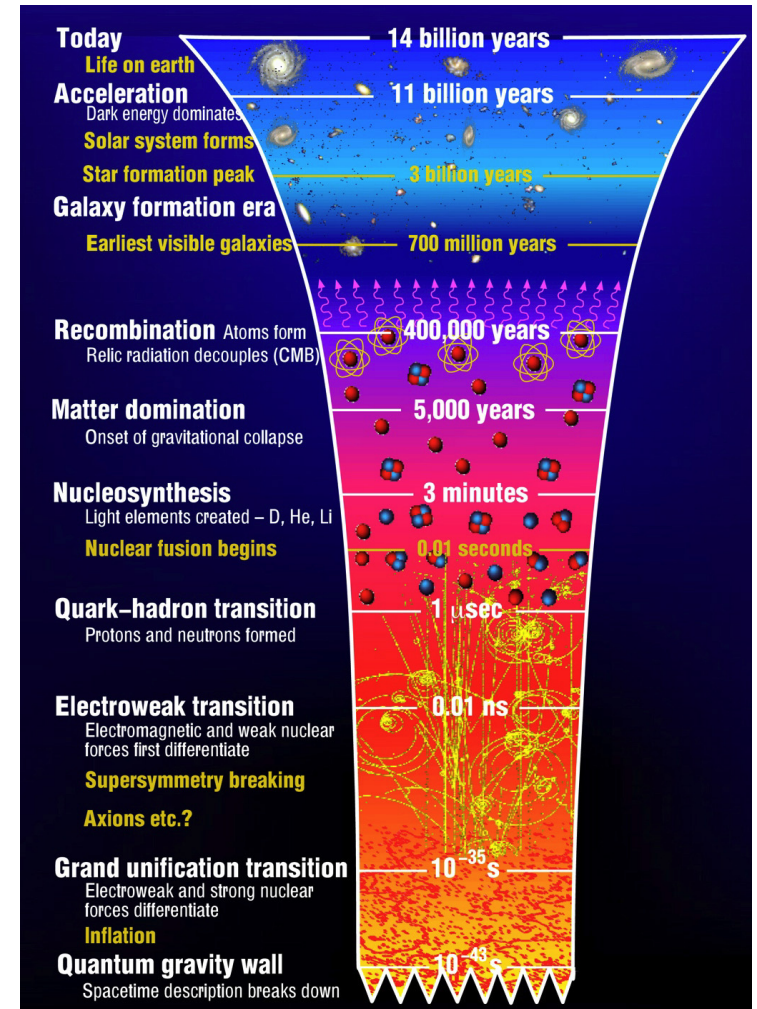
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- **From WIMPs to Light Dark Matter**
- **Sub-GeV DM and Collective Excitations**
- **Conclusions**

# 1. From WIMPs to Light DM



“Dark Matter Paradigm Remains the Only Viable and Simple Framework for the Full Range of Observations”

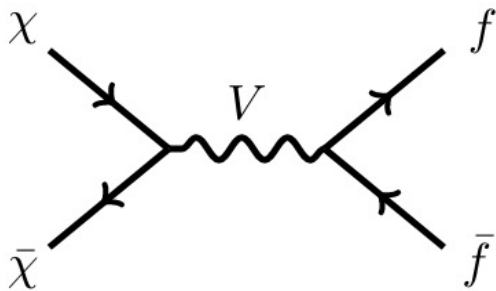


# 1. From WIMPs to Light DM

## Weakly Interacting Massive Particles

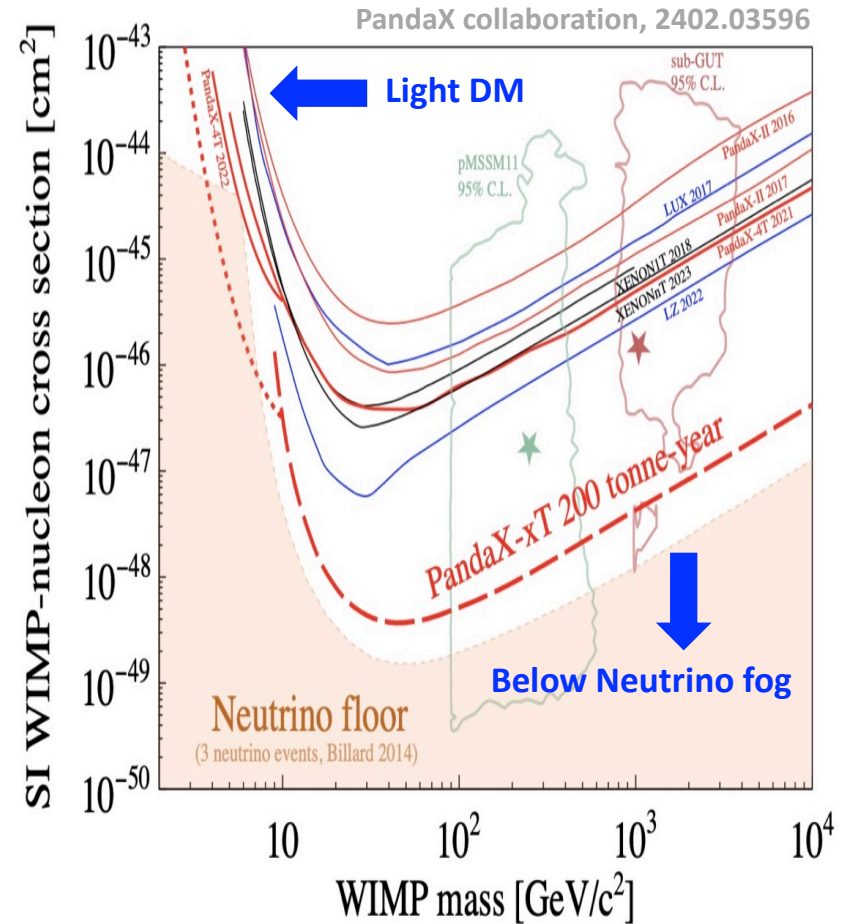
(WIMPs)

$$\Omega_\chi \propto \frac{1}{\langle \sigma v \rangle} \sim \frac{m_\chi^2}{g_\chi^4}$$



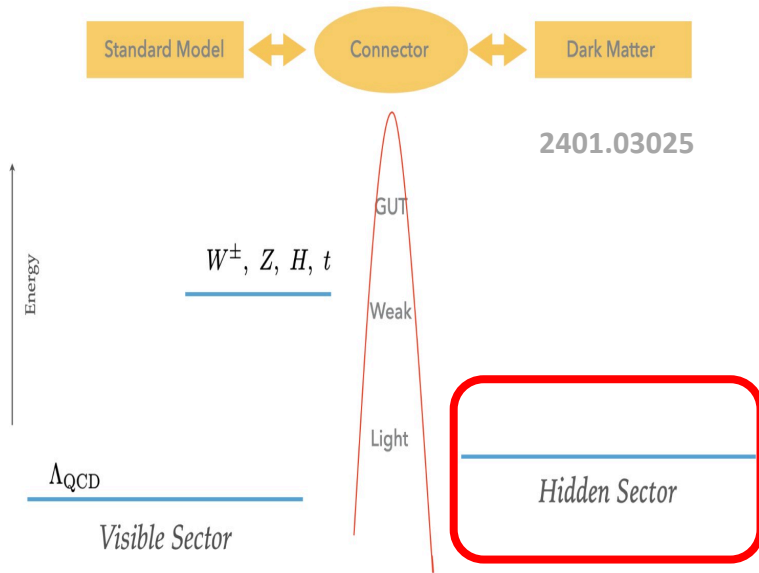
$$m_\chi \sim 100 \text{ GeV}, g_\chi \sim 0.6 \rightarrow \Omega_\chi \sim 0.1$$

## SUPerSYmmetry

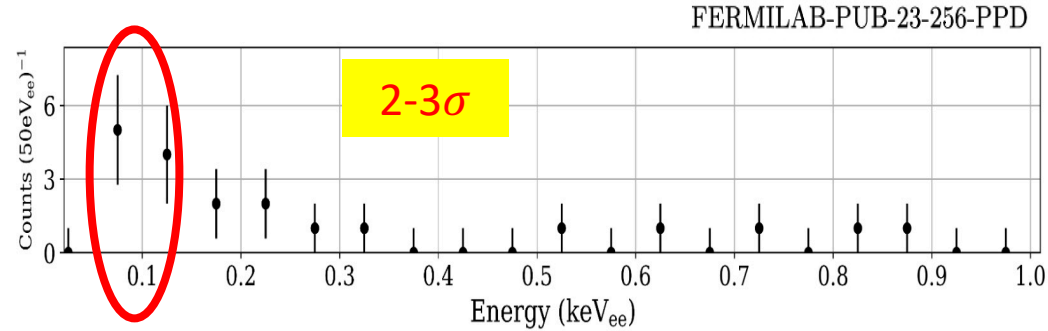


# 1. From WIMPs to Light DM

## Theory (Hidden sector DM)



## Data ( excess-1 )



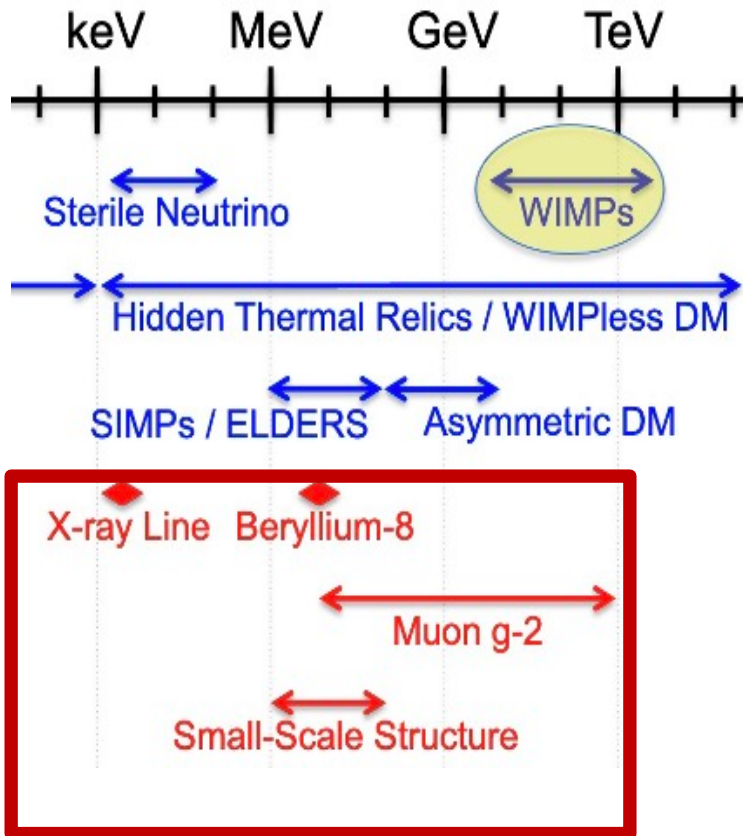
2002.06937

Readout Type	Target	Resolution	Exposure	Threshold	Excess Rate (Hz/kg)	Depth	Reference
Charge ( $E_e$ )	Ge	$1.6 e^-$	80 g · d	0.5 eVee ( $\sim 1e^-$ ) <sup>a</sup>	[20, 100]	.7 km	EDELWEISS [6]
	Si	$\sim 0.2 e^-$	0.18 g · d	1.2 eVee ( $< 1 e^-$ )	[6, 400]	100 m	SENSEI [4]
	Si	$0.1 e^-$	0.5 g · d	1.2 eVee ( $< 1 e^-$ )	[10, 2000]	$\sim 1$ m	CDMS HVeV [3]
	Si	$1.6 e^-$	200 g · d	1.2 eVee ( $\sim 1e^-$ )	$[1 \times 10^{-3}, 7]$	2 km	DAMIC [7]
Energy ( $E_{det}$ )	Ge	18 eV	200 g · d	60 eV	$> 2$	$\sim 1$ m	EDELWEISS [1]
	CaWO <sub>4</sub>	4.6 eV	3600 g · d	30 eV	$> 3 \times 10^{-3}$	.4 km	CRESST-III [2]
	Al <sub>2</sub> O <sub>3</sub>	3.8 eV	0.046 g · d	20 eV	$> 30$	$\sim 1$ m	$\nu$ CLEUS [8]
Photo $e^-$	Xe	6.7 PE ( $\sim 0.25 e^-$ )	15 kg · d	12.1 eVee ( $\sim 14$ PE)	$[0.5, 3] \times 10^{-4}$	.4 km	XENON10 [5, 9]
	Xe	6.2 PE ( $\sim 0.31 e^-$ )	30 kg · yr	$\sim 70$ eVee ( $\sim 80$ PE)	$> 2.2 \times 10^{-5}$	.4 km	XENON100 [5]
	Xe	$< 10$ PE	60 kg · yr	$\sim 140$ eVee ( $\sim 90$ PE)	$> 1.7 \times 10^{-6}$	.4 km	XENON1T [10]
	Ar	$\sim 15$ PE ( $\sim 0.5 e^-$ )	6780 kg · d	50 eVee	$> 6 \times 10^{-4}$	.4 km	Darkside50 [11]

**QCD/QED-like theory,**  
**Remnant from SUSY breaking,**  
**Sector arising from an RS ...**

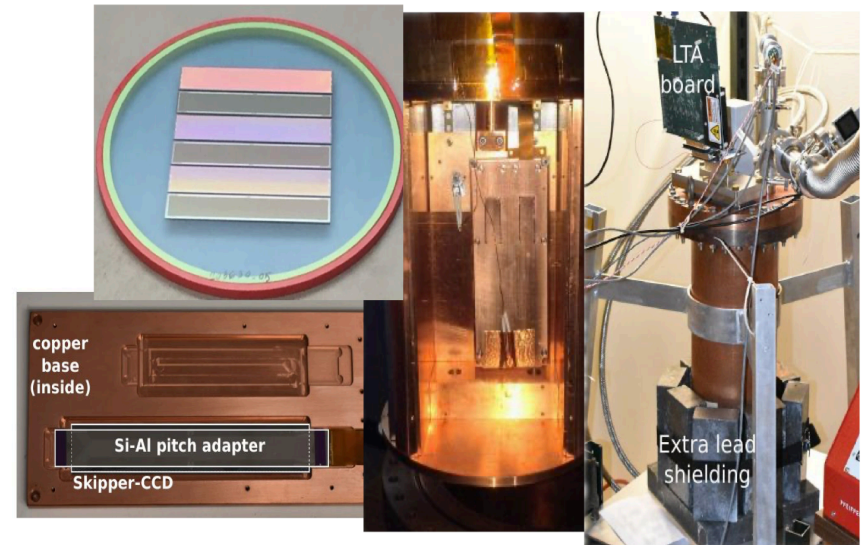
# 1. From WIMPs to Light DM

## Data ( excess-2 )



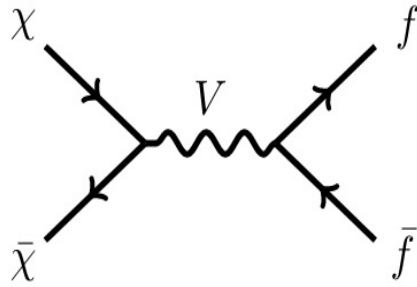
## Technology (Skipper-CCD)

Stephen Holland @ LBNL



charges in pixel,  
accessing to  $1e^-$  signal,  
very low noise

# 2. sub-GeV DM and Collective Excitations



- $m_V > m_\chi \quad \langle \sigma v \rangle \simeq \frac{16\pi\alpha_\chi\alpha_f m_\chi^2}{m_V^4}$



**Light non-SM mediator**

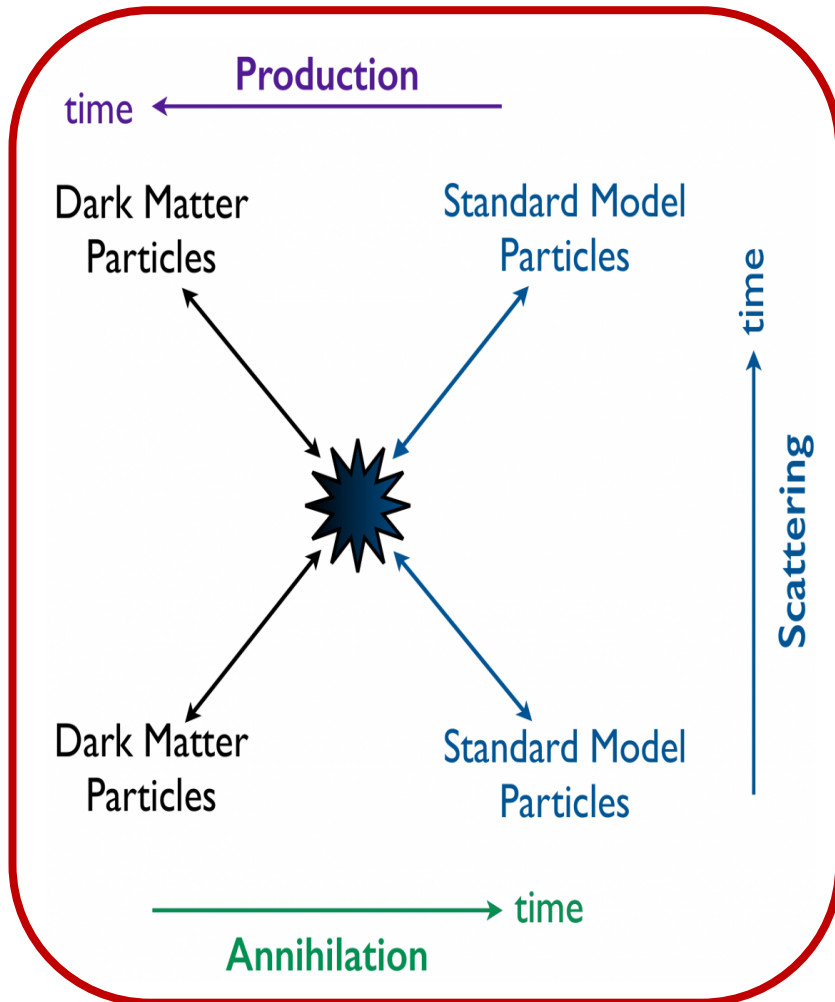
- $m_V < m_\chi \quad \langle \sigma v \rangle \simeq \frac{\pi\alpha_\chi\alpha_f}{m_\chi^2}$

## 1 Relic Density and Models

### Secluded DM or WIMPless DM

- SIMPs [YH, Kuflik, Volansky, Wacker, 2014; YH, Kuflik, Murayama, Volansky, Wacker, 2015]
- ELDERs [Kuflik, Perelstein, Rey-Le Lorier, Tsai, 2016 & 2017]
- Forbidden dark matter [Griest, Seckall, 1991; D'Agnolo, Ruderman, 2015]
- Co-decaying dark matter [Dror, Kuflik, Ng, 2016]
- Co-scattering dark matter [D'Agnolo, Pappadopulo, Ruderman, 2017]
- .....

# 2. sub-GeV DM and Collective Excitations



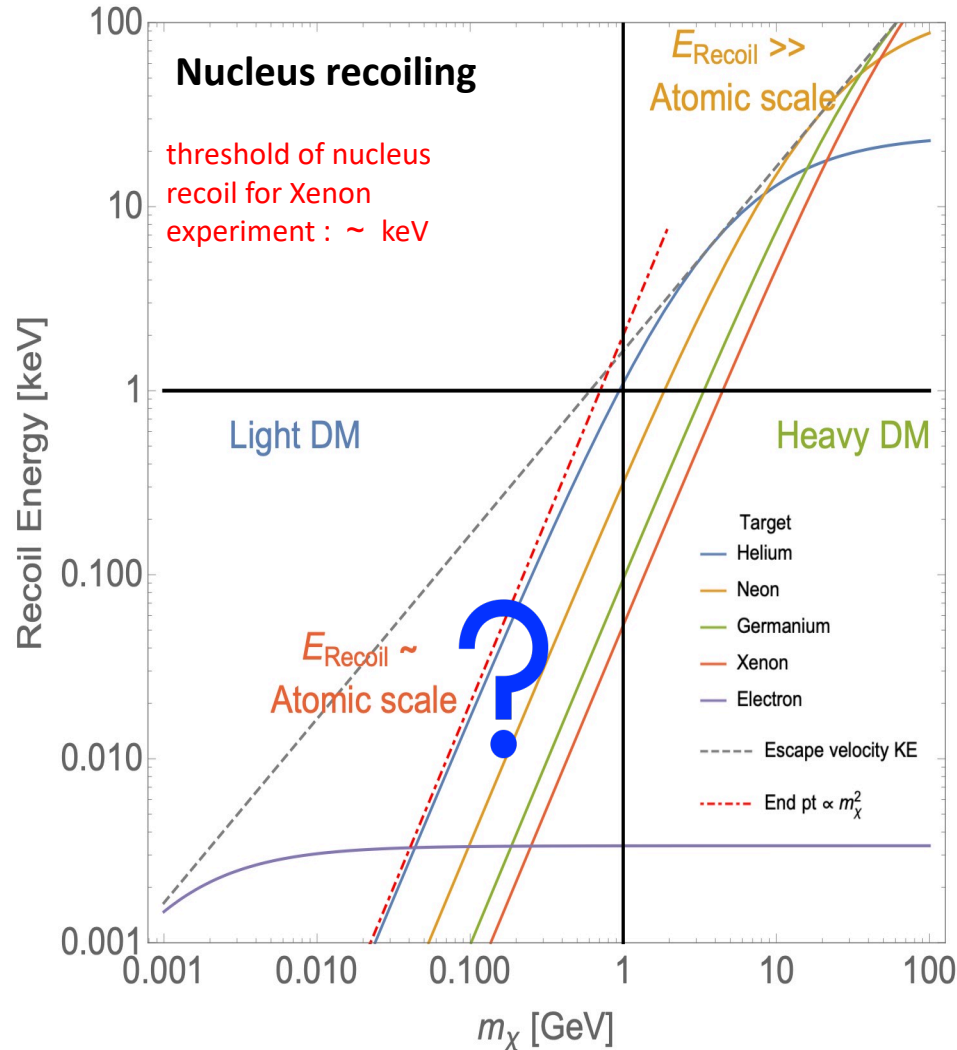
2

## How to test Light DM models

- **BBN, CMB**
- **Large Scale Structure**
- **Stellar Evolution and Cooling**
- **BeamDump**



# 2. sub-GeV DM and Collective Excitations



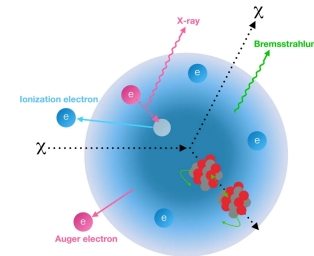
## 1. Boosted Dark Matter



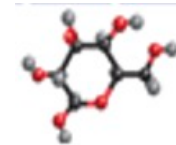
$$V \sim 10^{-3} c$$

$$V \sim c$$

## 2. Ionization Effects



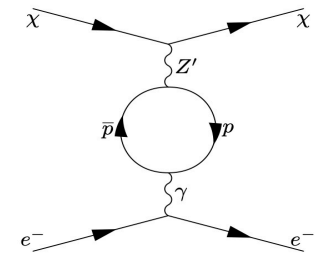
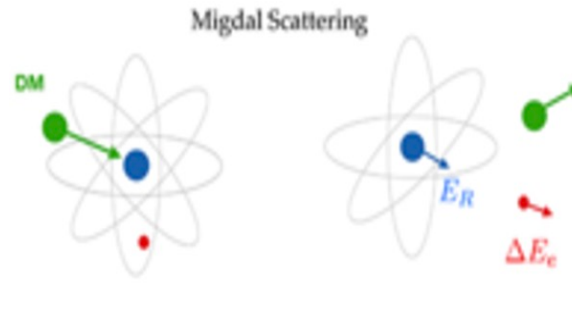
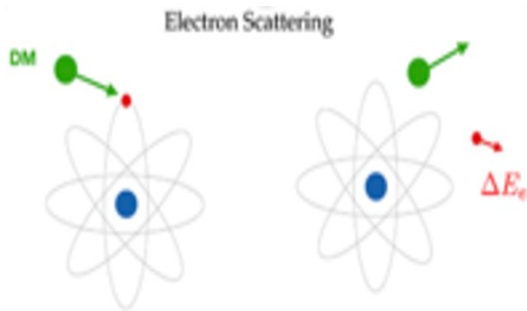
## 3. New Materials



# 2. sub-GeV DM and Collective Excitations

## Ionization signals

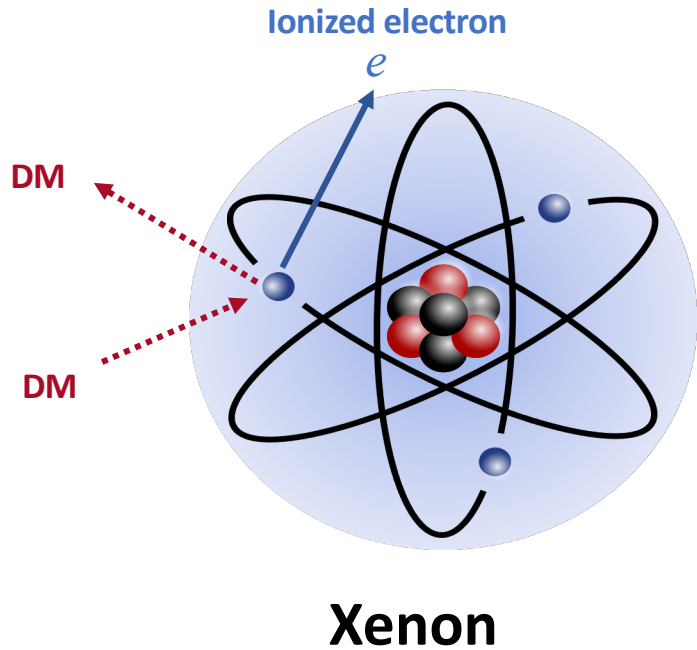
(DM-e, Migdal effect, Loop induced)



Light Dark Matter	Noble Liquid Detector (large volume, high threshold)	Solid detector (small volume, low threshold)
Non-Relativistic	✓	✓
Relativistic	✓	?

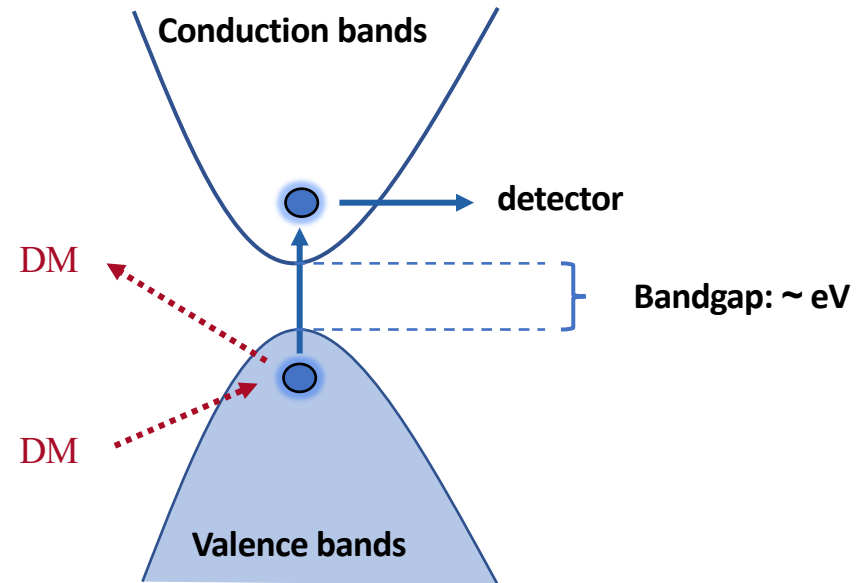
## Collective Excitations

# 2. sub-GeV DM and Collective Excitations



lower threshold,  
lighter dark matter

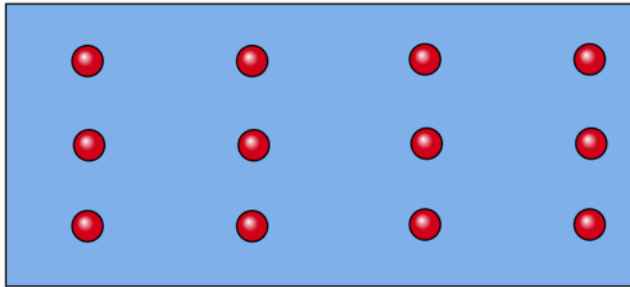
**Collective behavior is crucial in Solid Detector**



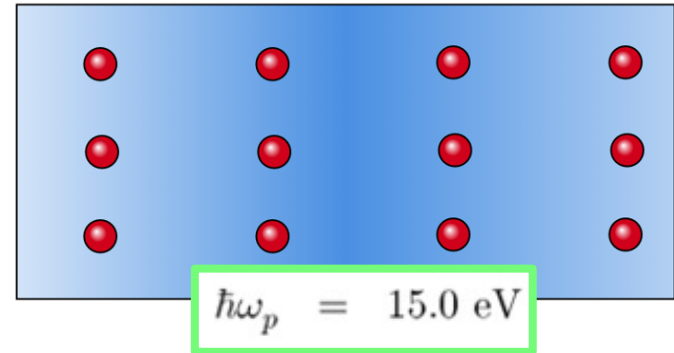
**Semi-conductor**

# 2. sub-GeV DM and Collective Excitations

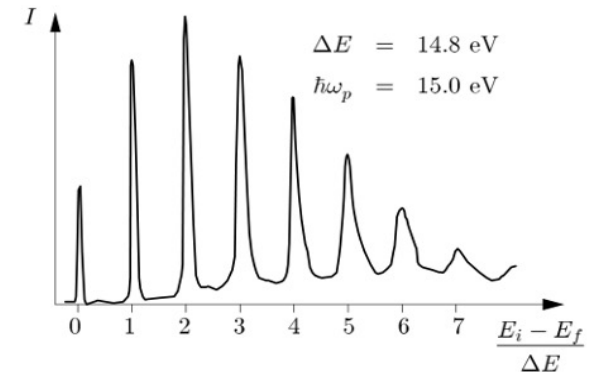
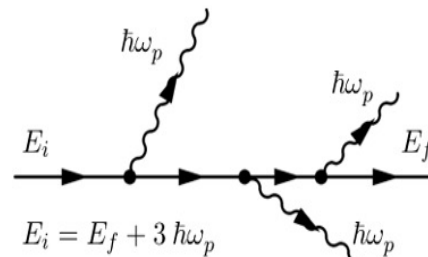
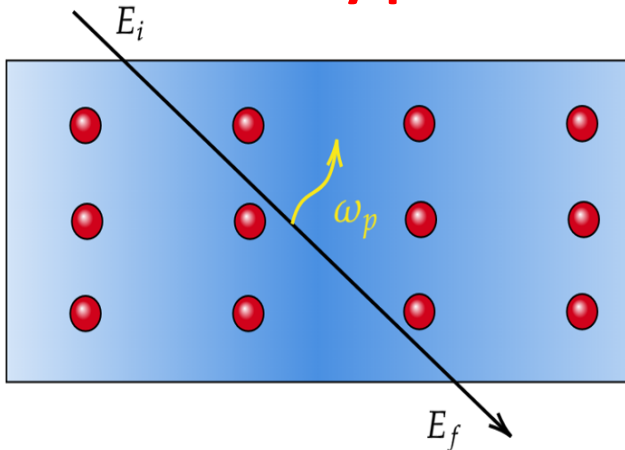
**Valence electrons** are free in metals and semiconductors, and form an electron gas holding the ionic lattice together.



Plasmon is the quantized form of the **electron density fluctuation** in the medium, oscillating with a frequency:  $\omega_p = \sqrt{\frac{4\pi\alpha n_e}{m_e}}$



## How to identify plasmons: Electron Energy Loss Spectroscopy (EELS)



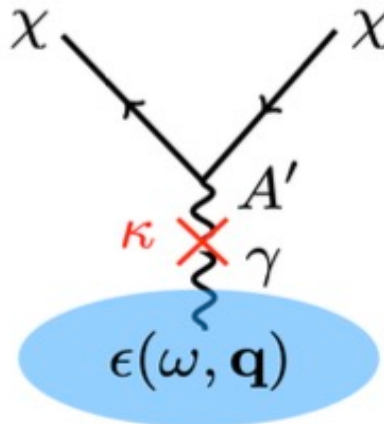
# 2. sub-GeV DM and Collective Excitations

Liang, Su, LW, Zhu, 2401.11971

## ◆ DM-induced Plasmon Signals in Solid Detector

$$R \sim \int d^3\mathbf{v} f(\mathbf{v}) \int d^3\mathbf{q} F_{\text{DM}}^2(\mathbf{q}) S(\mathbf{q}, \omega_{\mathbf{q}})$$

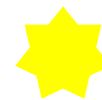
Interactions in CM systems



**Dielectric Function**

**Electronic Structure  
(polarization)**

$$\epsilon(Q, \omega) = \epsilon_1(Q, \omega) + i\epsilon_2(Q, \omega)$$




**Dissipation Processes  
(absorption and scattering)**

# 2. sub-GeV DM and Collective Excitations

Liang, Su, LW, Zhu, 2401.11971

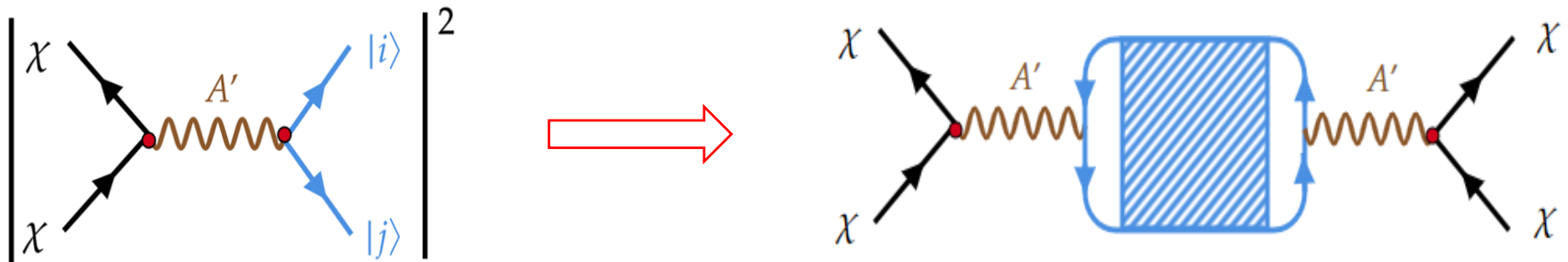
## Benchmark Model :

$$\mathcal{L}_{\text{int}} \supset g_\chi \bar{\chi} \gamma^\mu \chi A'_\mu + g_e \bar{e} \gamma^\mu e A'_\mu$$

 Non-relativistic QFT

$$\mathcal{L}_e^{\text{eff}} \supset g_e A'_0 \psi_e^* \psi_e + \frac{ig_e}{2m_e} \mathbf{A}' \cdot \left( \psi_e^* \vec{\nabla} \psi_e - \psi_e^* \overleftarrow{\nabla} \psi_e \right) + \dots$$

**In-medium effect**



# 2. sub-GeV DM and Collective Excitations

Liang, Su, LW, Zhu, 2401.11971

$$\chi_{\hat{\rho}\hat{\rho}}^r = \text{[Diagram: A rectangle with a cross-hatch pattern inside, representing a density-density correlation function]} \quad \text{Density-Density correlation function (RPA)}$$

$$= \text{[Diagram: A series of circles with diagonal hatching, connected by wavy lines, representing a perturbative expansion of the RPA function]} + \dots$$

$$= \frac{\text{[Diagram: A single hatched circle]} + \dots}{1 - \text{[Diagram: A hatched circle connected to another hatched circle by a wavy line]}} \quad \xrightarrow{\text{blue arrow}} \quad \Pi_e(Q, \omega) \quad \text{Lindhard function}$$

$$= \frac{\Pi_e}{1 - V_e \Pi_e} \quad \xrightarrow{\text{blue arrow}} \quad \frac{1}{\epsilon(Q, \omega)} = 1 + V_e(Q) \chi_{\hat{\rho}\hat{\rho}}^r(Q, \omega) \quad \text{Dielectric Function}$$

## Energy Loss Function (ELF)

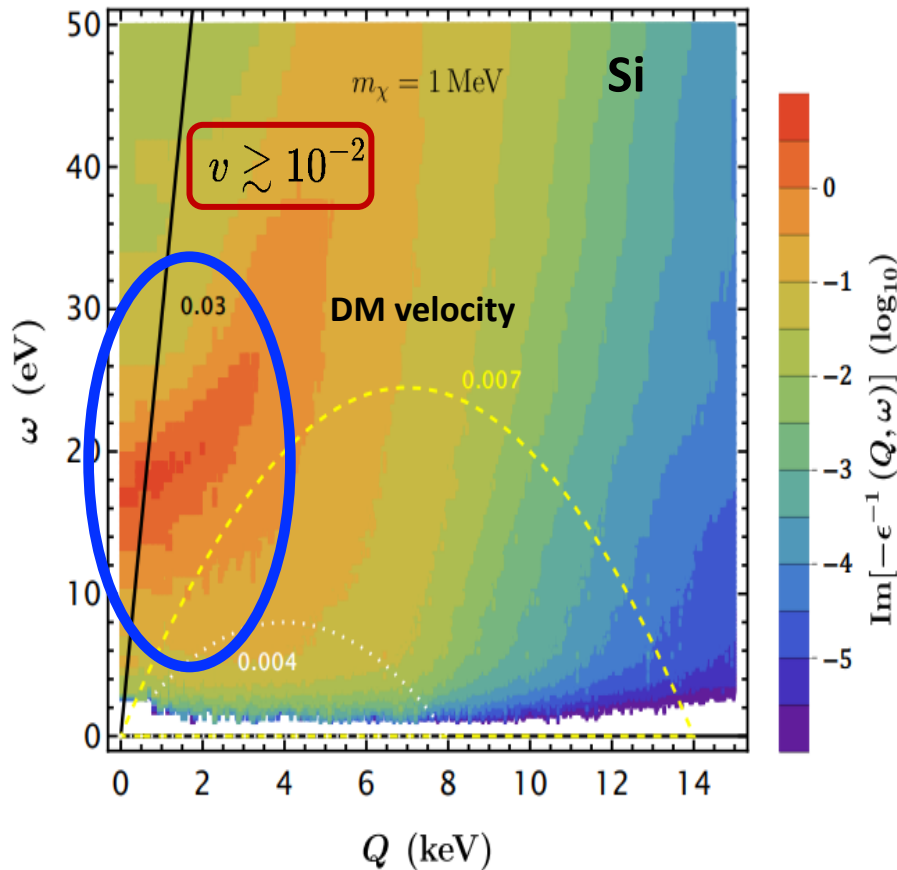
$$\Gamma(\mathbf{p}_\chi) = \int \frac{d^3\mathbf{Q}}{(2\pi)^3} |V(\mathbf{Q}, \omega)|^2 \left[ 2 \frac{Q^2}{e^2} \text{Im} \left( -\frac{1}{\epsilon(\mathbf{Q}, \omega)} \right) \right] \text{Im} \left[ \frac{-1}{\epsilon} \right] = \frac{\text{Im}[\epsilon]}{\text{Re}[\epsilon]^2 + \text{Im}[\epsilon]^2} \frac{1}{\pi} \frac{\frac{1}{2}\Gamma}{(x-x_0)^2 + (\frac{1}{2}\Gamma)^2} \sim \delta(x-x_0)$$

DM scattering rate

Scattering potential

# 2. sub-GeV DM and Collective Excitations

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Low deposited energy

◆ Energy Loss Function:

Density Functional Theory

◆ Resonance (plasmon) :

$|Q| < 5 \text{ keV}, \omega \sim 15 \text{ eV}$

◆ To excite plasmon:

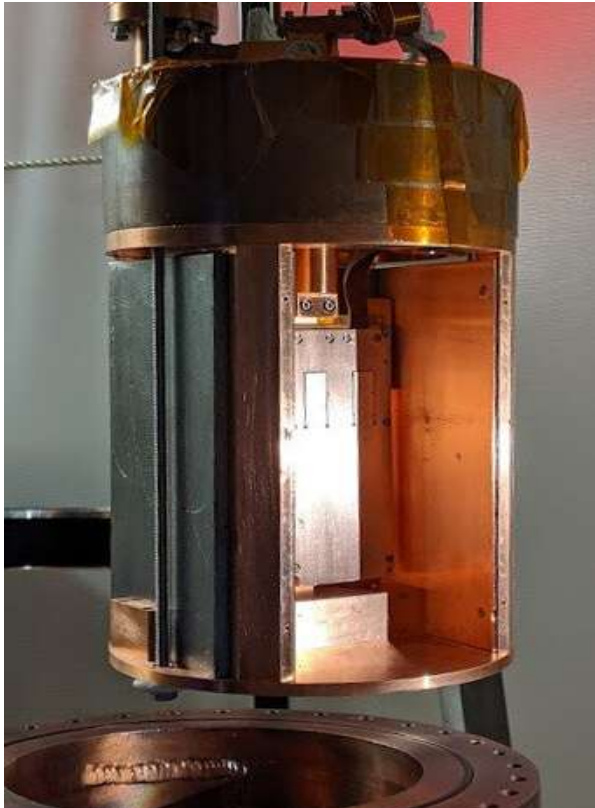
$V_{\min} > q/\omega \sim 10^2$

Relativistic Dark Matter



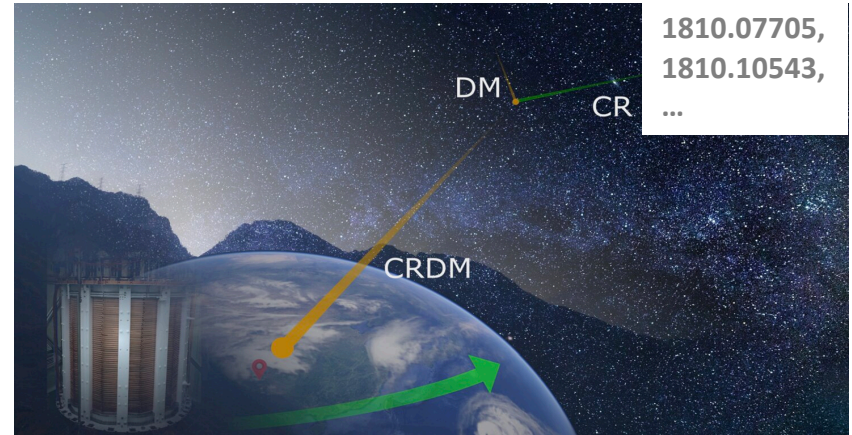
# 2. Light DM and Collective Excitations

## Sensei @ SNOLab

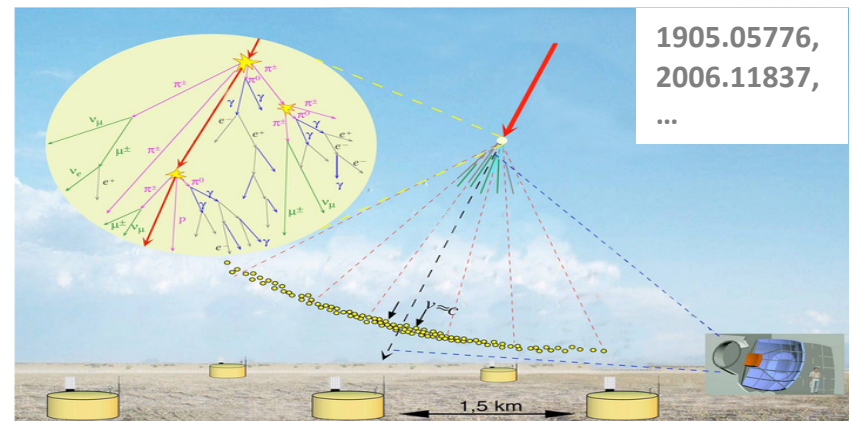


Electron recoil **ionizing 1–6  $e^-$**  in the detector, which corresponds to **(~1.2–20 eV) energy depositions**

## Cosmic ray Boosted DM



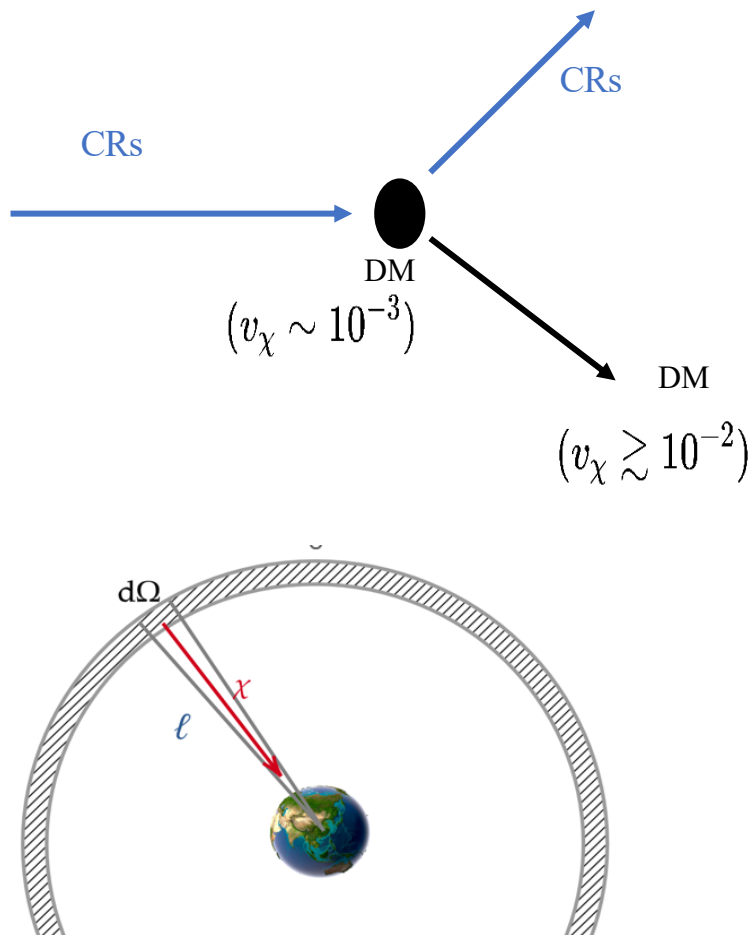
## Atmospheric DM



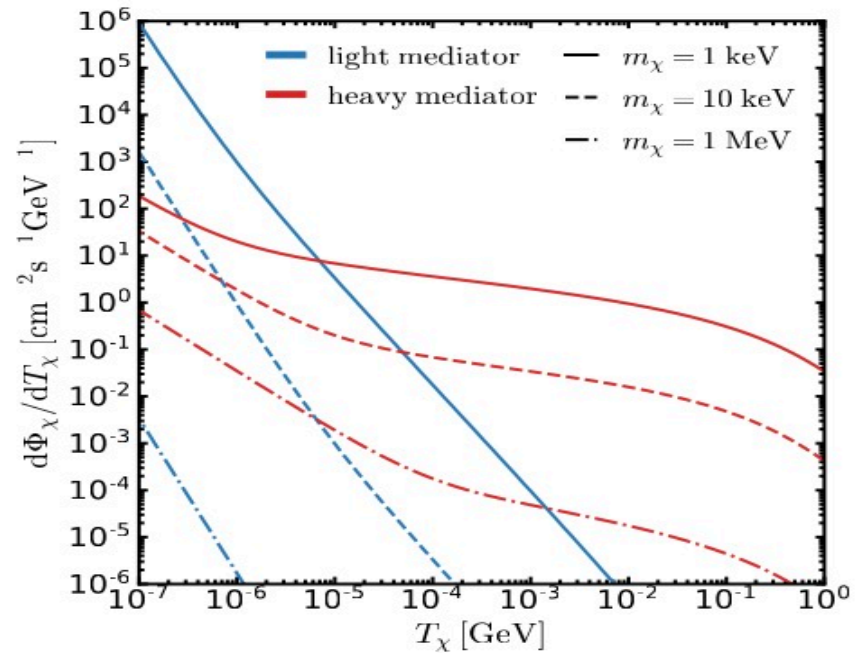
# 2. Light DM and Collective Excitations

Liang, Su, LW, Zhu, 2401.11971

## Cosmic Ray Boosted DM



$$\frac{d\Phi_\chi}{dT_\chi} = \int \frac{d\Omega}{4\pi} \int_{\text{l.o.s.}} \frac{\rho_\chi(\Omega, \ell)}{m_\chi} d\ell \int_{T_e^{\min}}^{+\infty} dT_e \frac{d\sigma_{\chi e}}{dT_\chi} \frac{d\Phi_e}{dT_e}$$

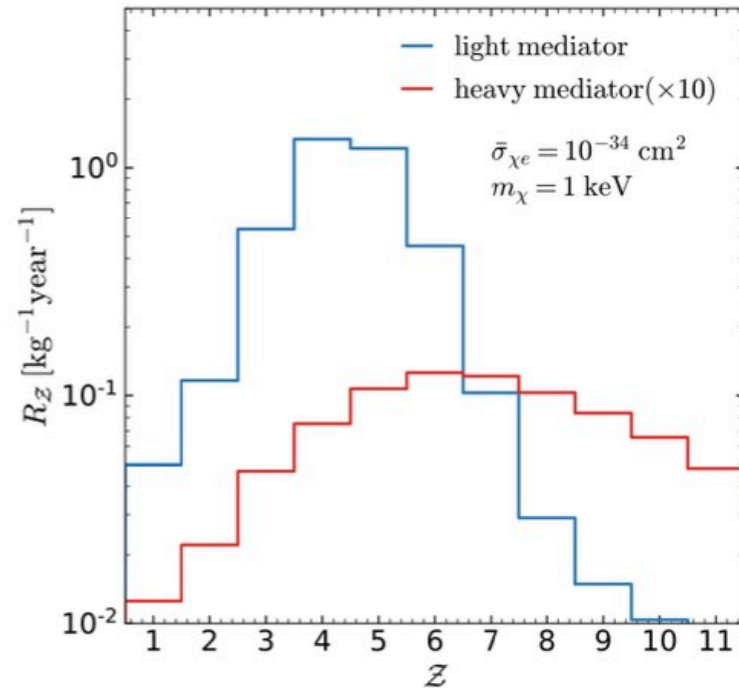
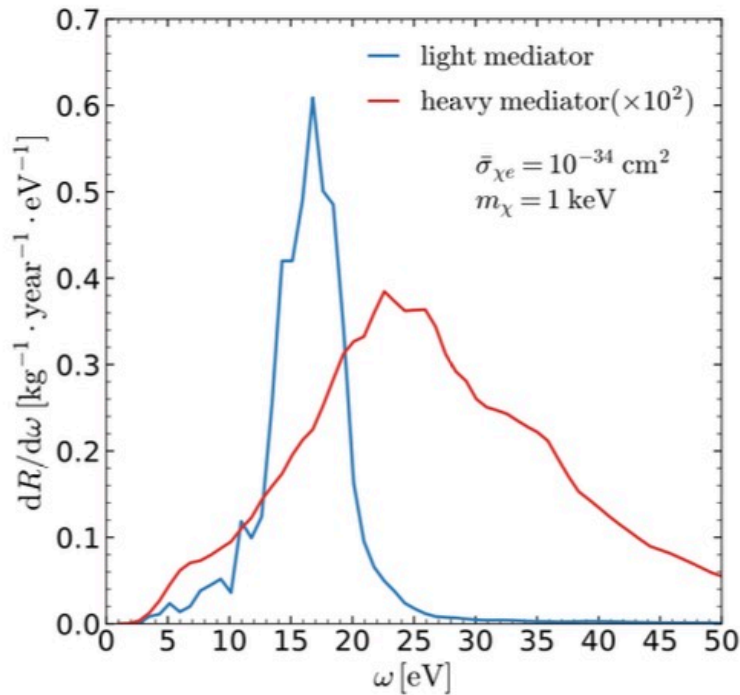


$$R = \frac{1}{\rho_T} \int dT_\chi \int \frac{d\Omega}{4\pi} \frac{d\Phi_\chi}{dT_\chi} \left( \frac{E_\chi}{p_\chi} \right) \Gamma(p_\chi)$$

# 2. Light DM and Collective Excitations

Liang, Su, LW, Zhu, 2401.11971

$$|F_{\text{DM}}(q)|^2 = \frac{(\alpha^2 m_e^2 + m_{A'}^2)^2}{(q^2 + m_{A'}^2)^2} = \begin{cases} 1 & \text{heavy mediator} \\ \frac{(\alpha m_e)^4}{q^4} & \text{light mediator} \end{cases}$$

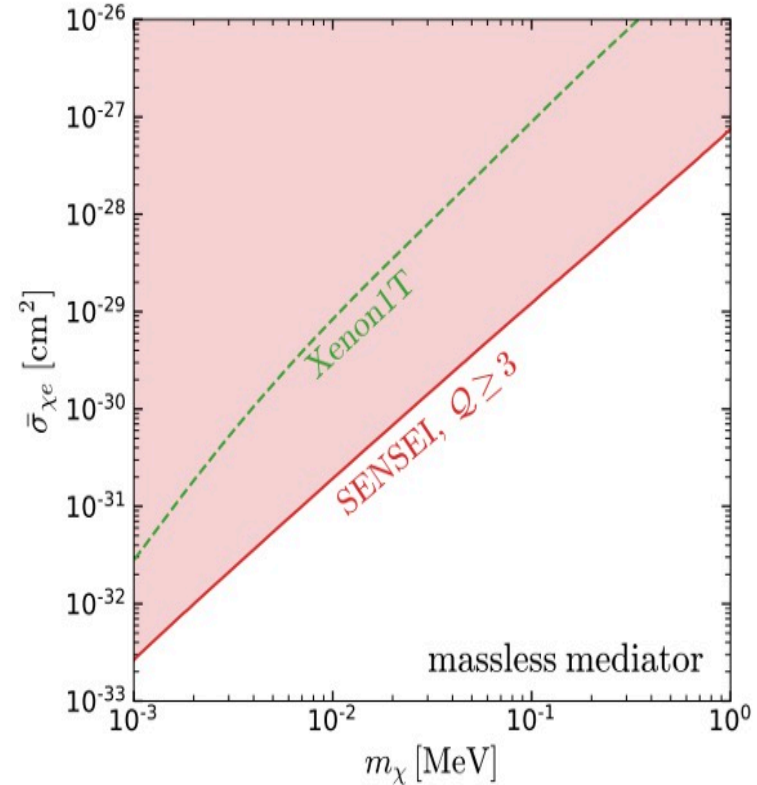


# 2. Sub-GeV DM and Collective Excitations

Liang, Su, WU, Zhu, 2401.11971

2312.13342 SENSEI @ SNOLAB

	All data		
Shape	Expo.	Ev.	Bkgd.
2e2p, h	13.58	10	10.66
2e2p, v	16.21	13	12.65
2e2p, d	17.82	32	25.31
<b>2e, all</b>	<b>46.61</b>	<b>55</b>	<b>48.62</b>
3e2p	30.87	3	0.01
3e3p	26.84	1	0.06
<b>3e, all</b>	<b>57.71</b>	<b>4</b>	<b>0.07</b>
4e2p	19.51	0	0.00
4e3p	27.60	0	0.00
4e4p	15.93	0	0.00
<b>4e, all</b>	<b>63.03</b>	<b>0</b>	<b>0.00</b>
5e, all	65.56	0	0.00
6e, all	67.31	0	0.00
7e, all	68.53	0	0.00
8e, all	69.52	0	0.00
9e, all	70.30	0	0.00
10e, all	70.89	0	0.00



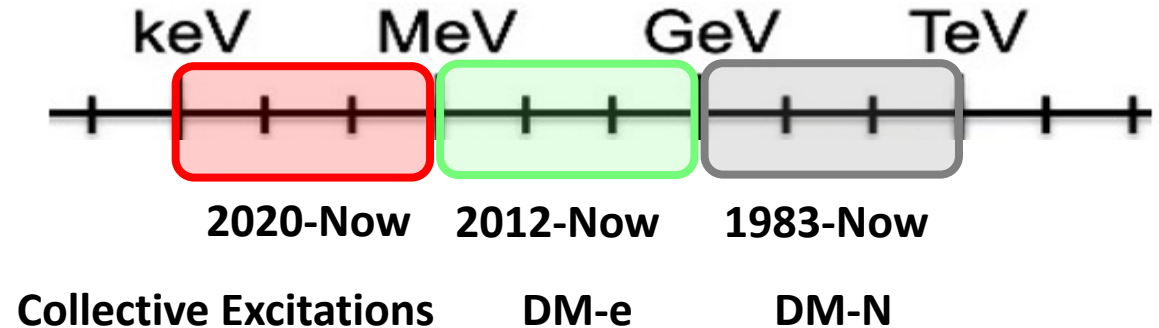
Light Dark Matter	Noble Liquid Detector (large volume, high threshold)	Solid detector (small volume, low threshold)
Non-Relativistic	✓	✓
Relativistic	✓	<b>Collective Excitations</b>

# Conclusions

Dark Matter Particle

Direct Detection

(~ 40 years,  $10^9$  orders)



Multi-disciplines + Multi-messengers

A New Era in the Quest for Dark Matter

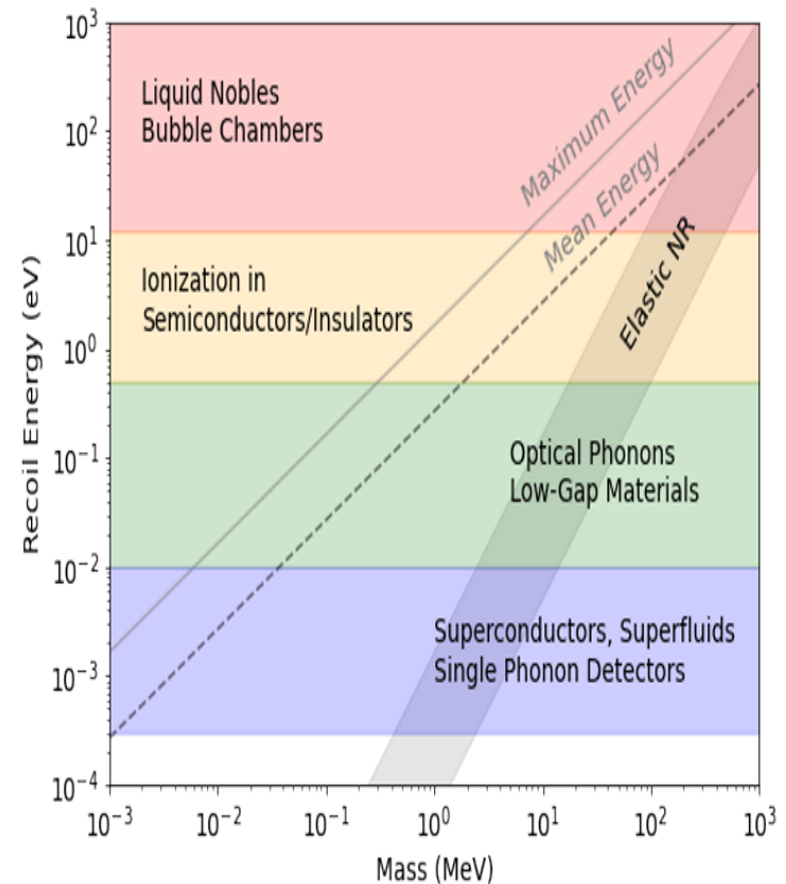


*Thank you very much !*

# Back Up



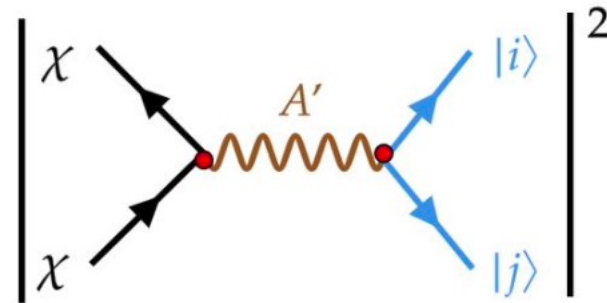
Threshold	Lowest DM Mass Probed	Relevant Techniques, Technologies, and Materials
20 eV	20 MeV (ER) 100 MeV (NR)	Noble Elements (TPCs & SPCs) Solid-State Charge Detectors Phonon Detectors Threshold Detectors
500 meV	1 MeV (ER/NR) 500 meV (Abs.)	Semiconductor Detectors Athermal Phonon Detectors Scintillators NIR Photon Detectors
5 meV	10 keV (CE) 5 meV (Abs.)	Superconductors Low-Gap Materials Athermal Phonon Detectors Polar Materials Superfluids FIR Photon Detectors Magnetic Bubble Chambers Other new ideas



Rouven Essig, *et al.*, arxiv: 2203.08297

NOPP-2024@ IHEP, Beijing

# Back Up



$$\sum_{i,j} p_j \langle j | e^{-i\mathbf{Q}\cdot\hat{\mathbf{x}}} | i \rangle \langle i | e^{i\mathbf{Q}\cdot\hat{\mathbf{x}}} | j \rangle \int d^3x e^{i\mathbf{Q}\cdot\mathbf{x}} u_i^*(\mathbf{x}) u_j(\mathbf{x})$$

$$\hat{\psi}(\mathbf{x}) = \sum_k \hat{a}_k u_k(\mathbf{x})$$

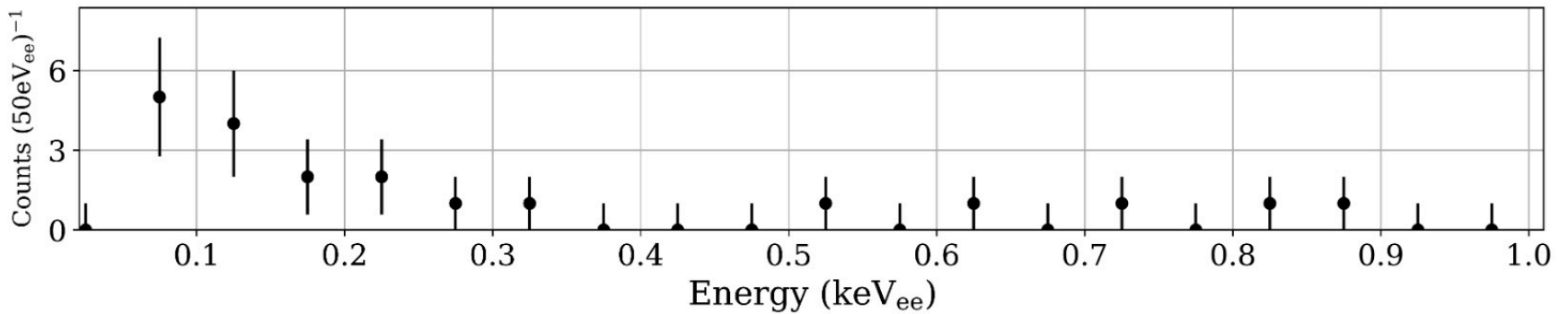
$$\langle i | e^{i\mathbf{Q}\cdot\hat{\mathbf{x}}} | j \rangle = \int d^3x e^{i\mathbf{Q}\cdot\mathbf{x}} \langle i | \hat{\psi}^\dagger(\mathbf{x}) \hat{\psi}(\mathbf{x}) | j \rangle = \int d^3x e^{i\mathbf{Q}\cdot\mathbf{x}} \langle i | \hat{\rho}(\mathbf{x}) | j \rangle$$

$$\int d^3x' d^3x e^{i\mathbf{Q}\cdot(\mathbf{x}-\mathbf{x}')} \sum_j p_j \langle j | \hat{\rho}(\mathbf{x}') \hat{\rho}(\mathbf{x}) | j \rangle \sim \langle \hat{\rho} \hat{\rho} \rangle$$

# Back Up

FERMILAB-PUB-23-256-PPD

## Confirmation of the spectral excess in DAMIC at SNOLAB with skipper CCDs



EXCESS21

EXCESS22

EXCESS22@IDM

EXCESS23@TAUP

excess may corresponds to a WIMP with mass  $\sim 2.5$  GeV/ $c^2$  and a WIMP-nucleon scattering cross section  $\sim 3 \times 10^{-40}$  cm $^2$ .



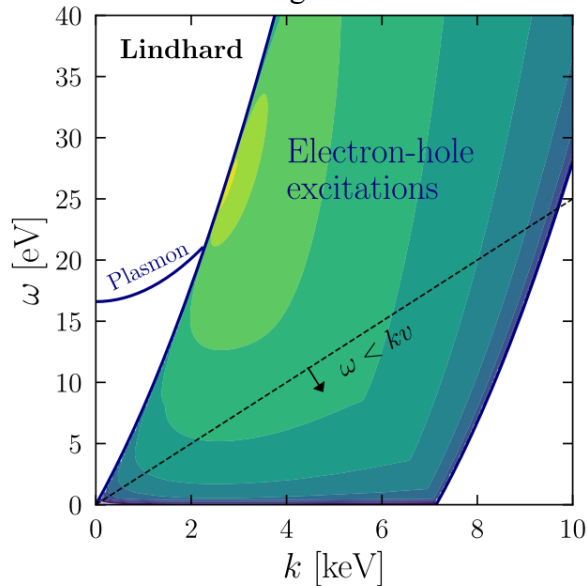


# Back Up



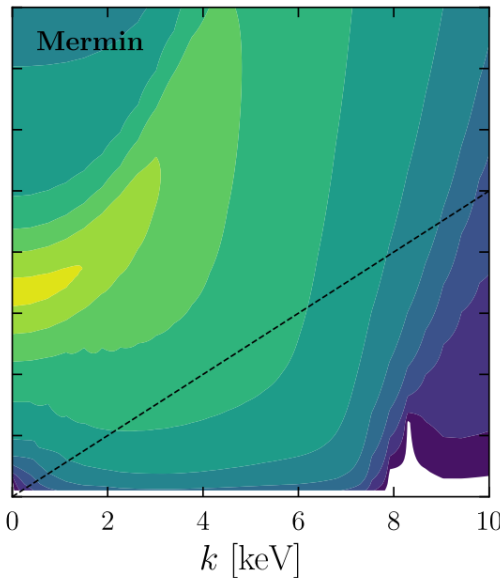
- The calculation of dielectric function:

DarkELF: Including the real and imaginary parts of the dielectric function calculated using the DFT software GPAW.



➤ The Lindhard method:

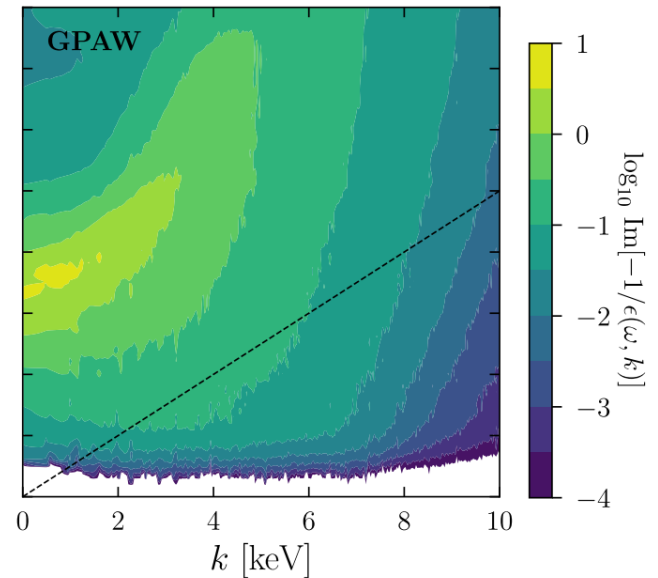
Assuming homogenous material and neglecting all dissipation effects.



➤ The Mermin method:

A generalization of the Lindhard which includes dissipation.

Simon Knapen, Jonathan Kozacuk, and Tongyan Lin, Phys. Rev. D **105**, 015014(2021).



➤ The GPAW method:

It relies on a first principles calculation with package GPAW.

# Back Up

**e.g. Cosmic Ray DM**

$$\frac{d\Phi_\chi}{dT_\chi} \equiv D_{\text{eff}} \frac{\rho_\chi^{\text{local}}}{m_\chi} \int_{T_{\text{CR}}^{\text{min}}}^{\infty} dT_{\text{CR}} \frac{d\sigma}{dT_\chi} \frac{d\Phi_{\text{CR}}^{\text{LIS}}}{dT_{\text{CR}}}$$

differential electronic excitation rate:

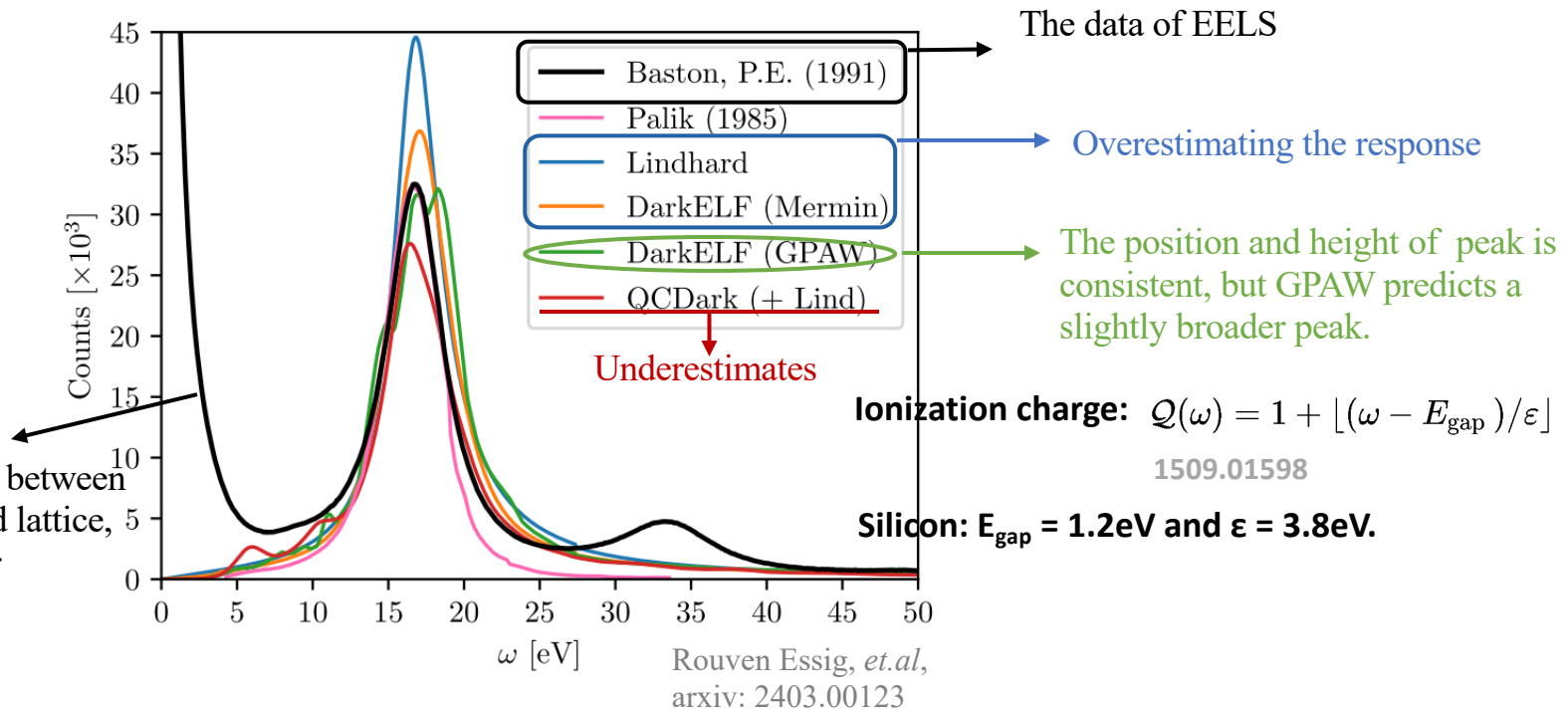
$$\begin{aligned} \frac{dR}{d\omega} &= \frac{1}{M_{\text{target}}} \int dE_\chi \int \frac{d\Omega}{4\pi} \frac{d\Phi_\chi}{dE_\chi} \frac{d\sigma}{d\omega} \\ &= \frac{1}{\rho_T} \frac{\bar{\sigma}_{\chi e}}{4\alpha\mu_{\chi e}^2} \int \frac{d^3\mathbf{Q}}{(2\pi)^3} \int dE_\chi \frac{d\Phi(E_\chi)}{dE_\chi} \left( \frac{(2E_\chi - \omega)^2 - Q^2}{4E_\chi(E_\chi - \omega)} \right) \left( \frac{\alpha^2 m_e^2 + m_{A'}^2}{Q^2 - \omega^2 + m_{A'}^2} \right)^2 \left( \frac{E_\chi(E_\chi - \omega)}{p_\chi^2} \right) \\ &\times Q \text{Im} \left[ \frac{-1}{\epsilon(\mathbf{Q}, \omega)} \right] \Theta \left[ E_\chi - \sqrt{(p_\chi - Q)^2 + m_\chi^2} - \omega \right], \end{aligned}$$

# Back Up

## DM-Electron scattering in semiconductor

- Comparison with EELS data:

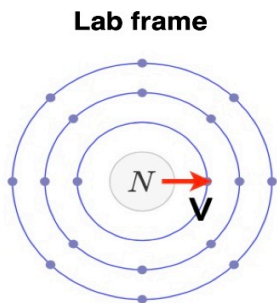
The results an incident electron beam kinetic energy of  $T = 100$  keV



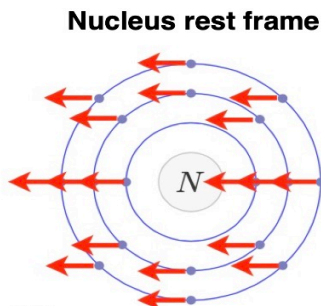
在论文中，作者们使用了几种不同的方法来计算或近似介电函数，包括：

1. **Lindhard模型**：这是一个半经典的模型，用于描述低动量转移下的电子气响应。它提供了一个合理的描述，特别是在低动量转移和小能量损失的区域。
2. **密度泛函理论 (DFT) 和随机相近似 (RPA)**：这些方法可以更精确地计算介电函数，包括电子波函数的量子力学计算。DFT结合RPA可以提供对介电函数实部和虚部的全面描述。
3. **QCDark计算**：这是一个基于第一性原理的计算工具，用于计算暗物质与电子散射的介电函数虚部。
4. **DarkELF工具**：这是一个开源软件包，它提供了使用DFT软件GPAW计算介电函数的实部和虚部的功能，并且可以选择使用Lindhard模型来拟合介电函数。

论文中还提到了如何将不同的计算方法结合起来，以获得更准确的介电函数表达式，特别是通过使用Lindhard模型来校正QCDark计算的介电函数虚部，从而得到一个在低能量区域表现良好的介电函数模型。这种组合模型在描述等离子体激元峰值方面特别有效，这是因为Lindhard模型在这一区域内提供了可靠的描述，并且与QCDark计算的介电函数虚部的更复杂的计算相匹配。



**v**: nucleus velocity after the recoil  
 The electron clouds are no more in the energy eigenstates

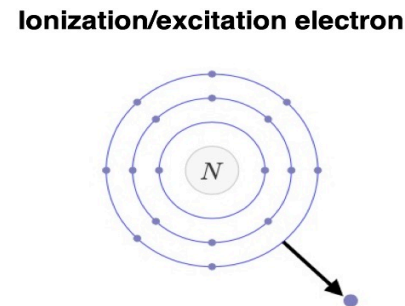


after NR                      before NR

$$|\Phi'_{ec}\rangle = e^{-im_e \sum_i \mathbf{v} \cdot \hat{x}_i} |\Phi_{ec}\rangle$$

Galilei transformation

$x_i$  the position operator of the  $i$ -th electron

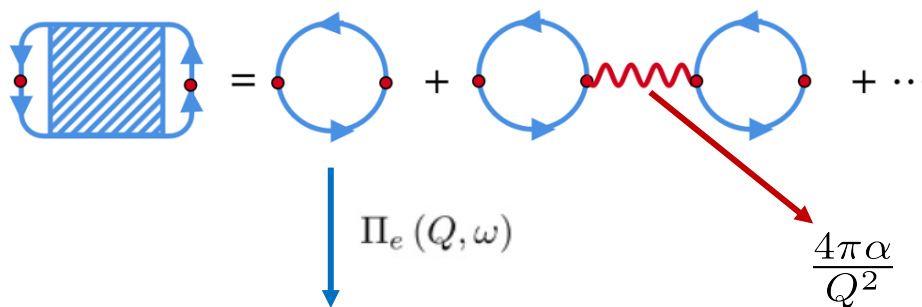


$$\mathcal{P} = |\langle \Phi_{ec}^* | \Phi'_{ec} \rangle|^2$$

ionized/excited energy eigenstate

$$\frac{1}{2} m_\chi v_\chi^2 \gtrsim E_{\text{threshold}} \quad \Rightarrow \quad m_\chi \gtrsim 6 \text{ MeV} \cdot \left( \frac{E_{\text{threshold}}}{15 \text{ eV}} \right)$$

where



$$\frac{1}{V} \sum_{i,j} \frac{|\langle i | e^{i\mathbf{Q} \cdot \hat{x}} | j \rangle|^2}{\varepsilon_i - \varepsilon_j - \omega - i0^+} (f_i - f_j)$$

**Lindhard function**

$$\Gamma(\mathbf{p}_\chi) = \int \frac{d^3\mathbf{Q}}{(2\pi)^3} |V(\mathbf{Q}, \omega)|^2 \left[ 2 \frac{Q^2}{e^2} \text{Im} \left( -\frac{1}{\epsilon(\mathbf{Q}, \omega)} \right) \right]$$

- ▶ Similar Fermi's Golden Rule, but different kinematics

$$Q = |\mathbf{Q}| = |\mathbf{p}_\chi - \mathbf{p}'_\chi|, \quad \omega = E_\chi - E'_\chi = \sqrt{p_\chi^2 + m_\chi^2} - \sqrt{|\mathbf{p}_\chi - \mathbf{Q}|^2 + m_\chi^2}$$

- ▶ Scattering potential

$$|V(\mathbf{Q}, \omega)|^2 = \frac{\pi \bar{\sigma}_{\chi e} \left[ (2E_\chi - \omega)^2 - Q^2 \right]}{4\mu_{\chi e}^2 E_\chi (E_\chi - \omega)} |F_{\text{DM}}(q)|^2,$$

- ▶ Dielectric function remains the same
- ▶ Event rate

$$R = \frac{1}{\rho_T} \int dT_\chi \int \frac{d\Omega}{4\pi} \frac{d\Phi_\chi}{dT_\chi} \left( \frac{E_\chi}{p_\chi} \right) \Gamma(p_\chi)$$