



## Neutrino Physics at Future Colliders

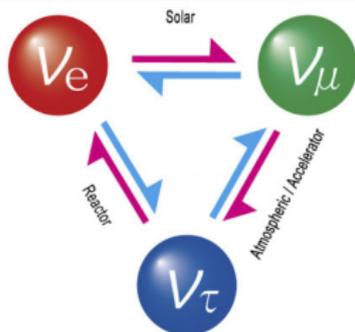
**Bhupal Dev**

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*Washington University in St. Louis*



# Neutrino Oscillations Primer



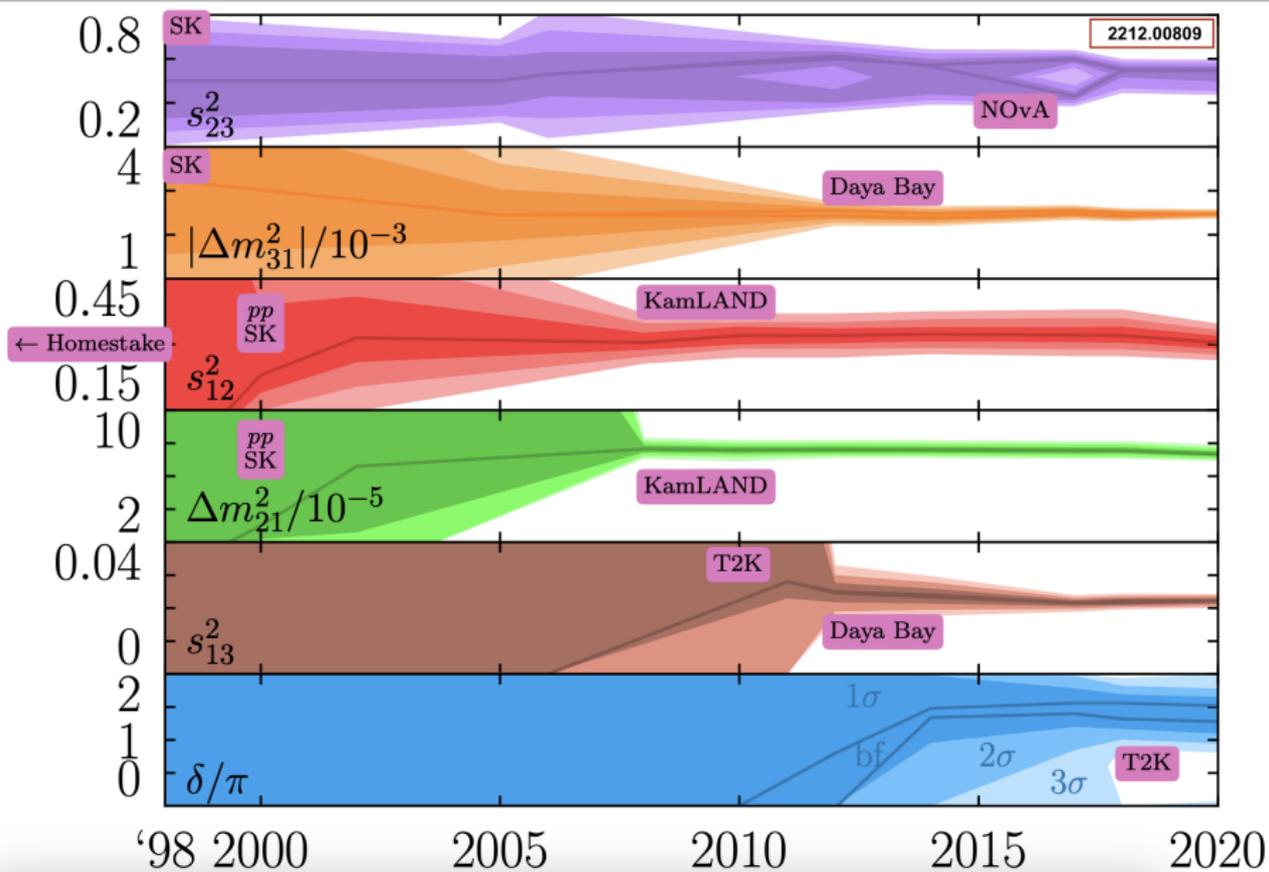
$$|\nu_i\rangle = U_{\alpha i} |\nu_\alpha\rangle$$

$$\begin{aligned}
 U &= \begin{pmatrix} 1 & & \\ & c_{23} & s_{23} \\ & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & s_{13}e^{-i\delta} \\ & 1 \\ -s_{13}e^{i\delta} & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} \\ -s_{12} & c_{12} \\ & & 1 \end{pmatrix} \\
 &= \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{13}s_{23}e^{i\delta} & c_{12}c_{23} - s_{12}s_{13}s_{23}e^{i\delta} & c_{13}s_{23} \\ s_{12}s_{23} - c_{12}s_{13}c_{23}e^{i\delta} & -c_{12}s_{23} - s_{12}s_{13}c_{23}e^{i\delta} & c_{13}c_{23} \end{pmatrix}
 \end{aligned}$$

$$P_{\alpha\beta} = \delta_{\alpha\beta} - 4 \sum_{i>j} \Re(U_{\alpha i}^* U_{\beta i} U_{\alpha j} U_{\beta j}^*) \sin^2 \left( \frac{\Delta m_{ij}^2 L}{4E} \right) \pm 8J \prod_{i>j} \sin \left( \frac{\Delta m_{ij}^2 L}{4E} \right)$$

**Requires nonzero neutrino masses and mixing  $\implies$  Physics beyond the Standard Model**

# Oscillation Parameters: Current Status



[Snowmass NF01 Topical Group Report, [2212.00809](#)]

## Future Prospects and Wish-List

		Atmospheric Mass Ordering	$\theta_{23}$ Octant	$\sin \delta \neq 0$ for 50% of $\delta$
JUNO	Optimistic	2030: $3\sigma$	-	-
	Conservative	2030: $2.5\sigma$	-	-
DUNE	Optimistic	2030: $5\sigma$	2036: $3\sigma$ , 2040: $5\sigma$	2035: $3\sigma$ , 2039: $5\sigma$
	Conservative	2032: $5\sigma$	2040: $2\sigma$	2037: $3\sigma$
HK	Optimistic	2033: $5\sigma$	2033: $5\sigma$	2029: $3\sigma$ , 2032: $5\sigma$
	Conservative	2032: $3\sigma$	2034: $3\sigma$	2029: $3\sigma$ , 2037: $5\sigma$
IceCube	Optimistic	2030: $3\sigma$ , 2033: $4\sigma$	-	-
	Conservative	2033: $2\sigma$	-	-
KM3NeT	Optimistic	2026: $3\sigma$ , 2029 $5\sigma$	-	-
	Conservative	2030: $3\sigma$ , 2032: $4\sigma$	-	-

[Snowmass NF01 Topical Group Report, [2212.00809](#)]

**What else do we want to learn about neutrinos?**

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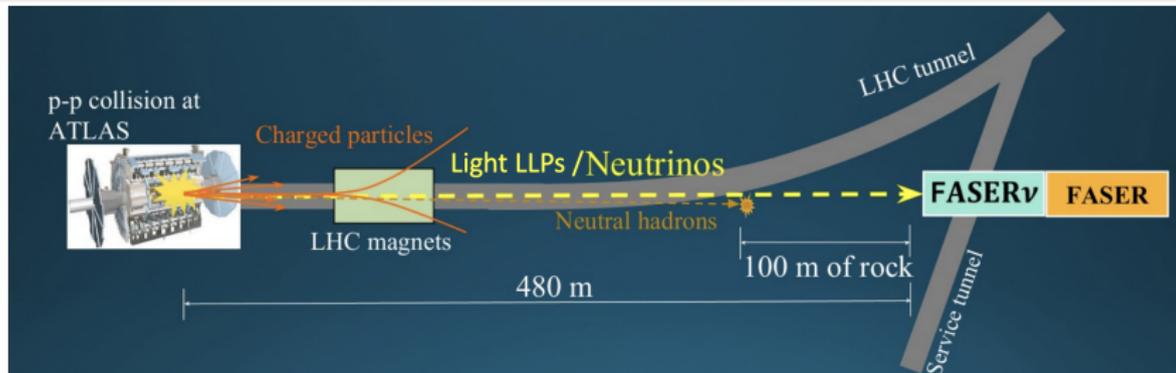
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## What else do we want to learn about neutrinos?

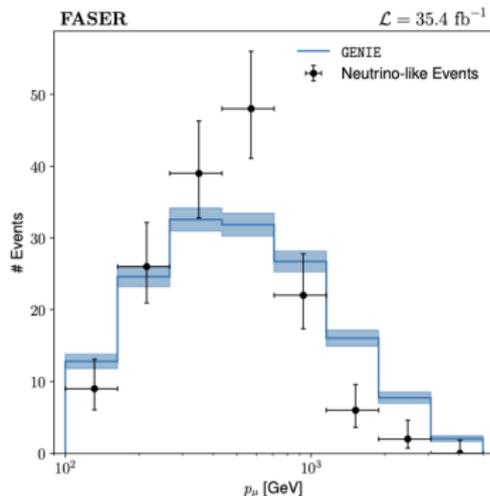
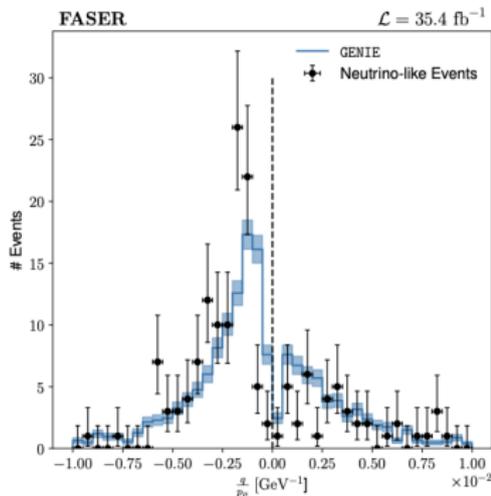
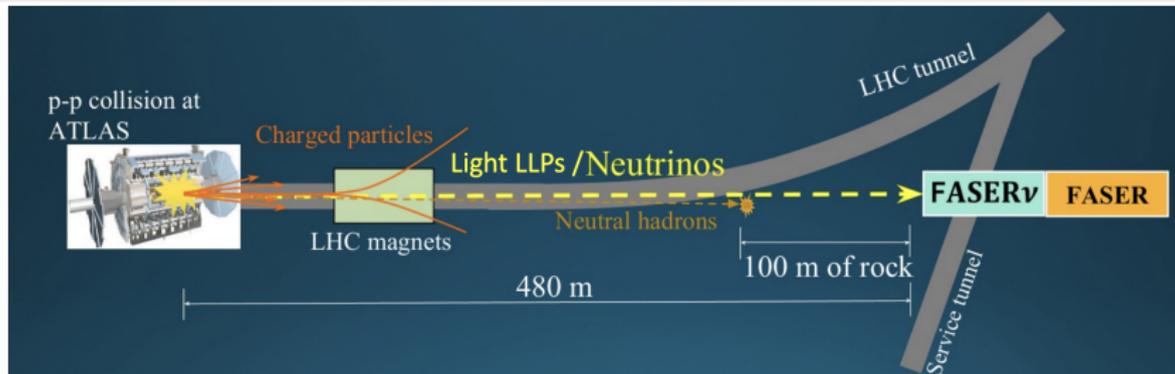
- How do they get mass? Dirac or Majorana (or something else)?
- Do they have more than 3 species? Must be ‘sterile’! [LEP, [hep-ex/0509008](#) (Phys. Rep.)]
- Do they have nonstandard interactions with matter or with themselves?
- Is the low-energy  $\delta_{CP}$  related to baryogenesis?
- Any connection to dark matter?

**High-energy colliders can help us address some of these questions.**

# Seeing the 'Invisible' at LHC

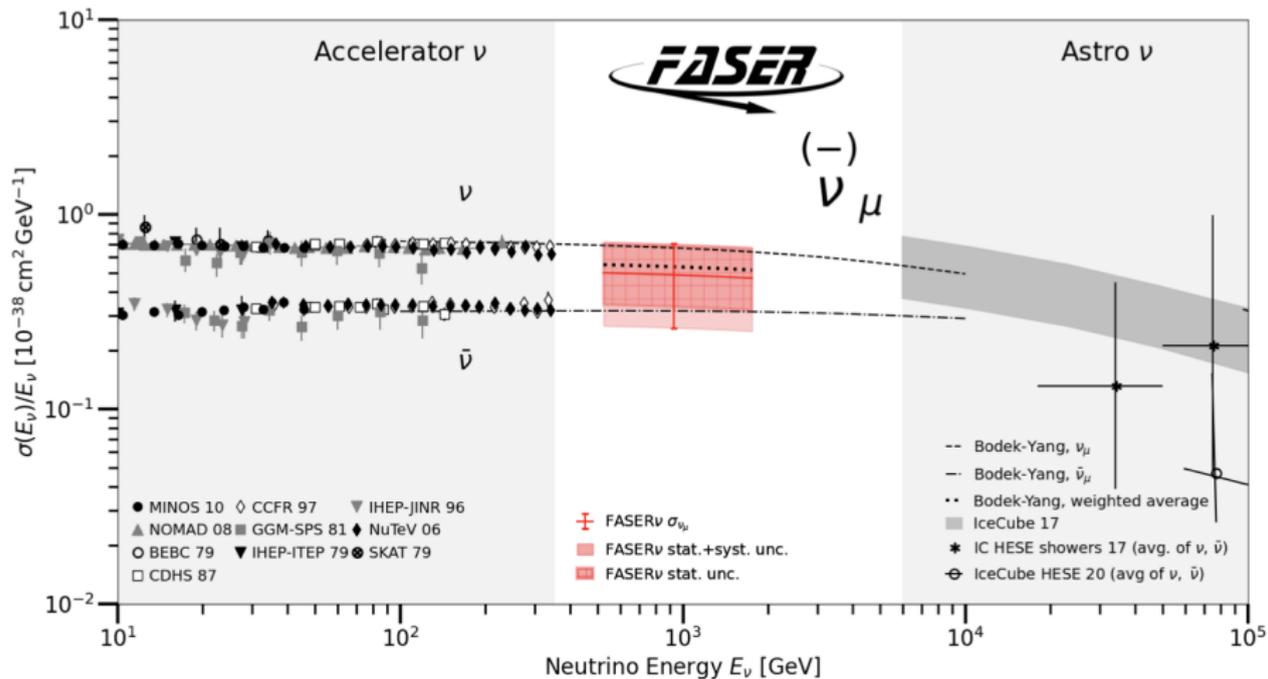


# Seeing the 'Invisible' at LHC



[FASER Collaboration, [2303.14185 \(PRL '23\)](#)]

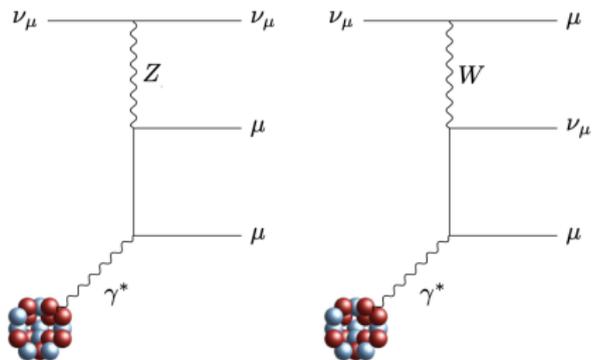
# Cross Section Measurements



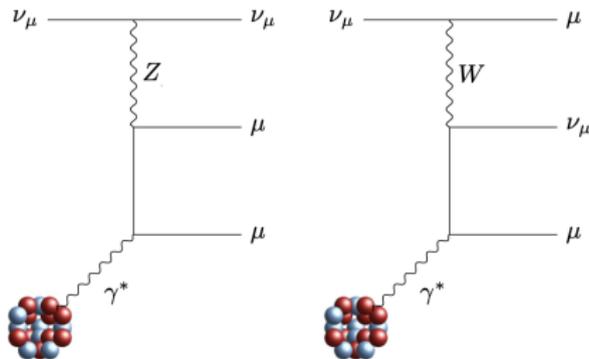
[FASER Collaboration, [2403.12520](#) (PRL '24)]

VERY important for validating neutrino-nucleus cross section models at different energies.

# Rare SM Processes



# Rare SM Processes

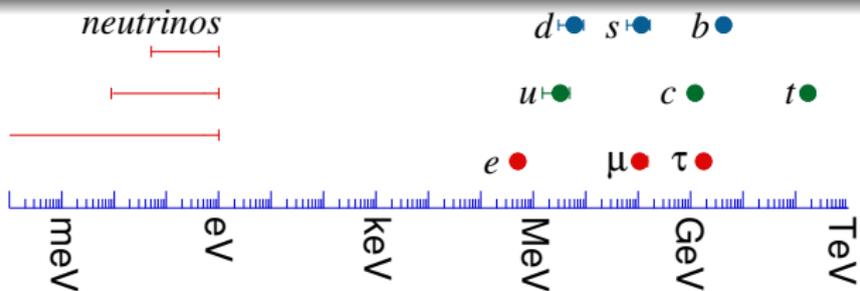


Name	Mass [tons]	Target nucleus	On(Off)-Axis	$L_{\text{int,TeV}\nu \rightarrow \mu\mu}^{-1} \times 10^{17} [\text{cm}^{-1}]$	Neutrino Tridents, $\nu N \rightarrow \nu N' \ell^+ \ell^-$						
					$\mu^+ \mu^-$	$\mu^+ \mu^-_{f_s=0.5}$	$e^+ e^-$	$\tau^+ \tau^-$	$e^\pm \mu^\mp$	$e^\pm \tau^\mp$	$\mu^\pm \tau^\mp$
Run 3 (150 fb <sup>-1</sup> )											
FASER $\nu$	1.1	W	On	252	0.22	0.54	0.24	0.0029	0.83	0.035	0.060
SND@LHC	0.83	W	Off	252	0.024	0.06	0.03	0.0002	0.10	0.004	0.004
HL-LHC (3 ab <sup>-1</sup> )											
FASER $\nu 2$	20	W	On	252	40	97	44	0.51	150	6.3	10
AdvSND@LHC (Far)	5	W	Off	252	2.2	5.3	2.7	0.02	9.0	0.3	0.4
FLArE	10	LAr	On	8.56	4.5	11	4.5	0.07	16	0.7	1.2
FLArE-100	100	LAr	On	8.56	26	63	27	0.37	91	4.1	6.8
NuTeV-like (Fe)	95	Fe	On	65.4	21	52	22	0.29	76	3.4	5.5
NuTeV-like (Pb)	135	Pb	On	154	48	116	57	0.45	190	7.0	10

[Altmannshofer, Makela, Sarkar, Trojanowski, Xie, Zhou, 2406.16803;

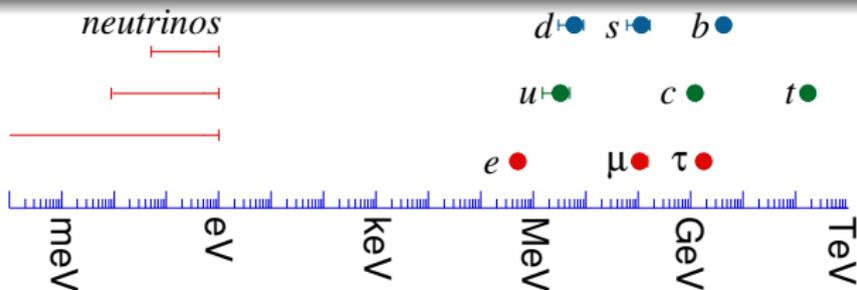
see also Bigaran, BD, Lopez Gutierrez, Machado, 2406.20067]

# Probing Neutrino Mass Mechanism



Perhaps something beyond the standard Higgs mechanism...

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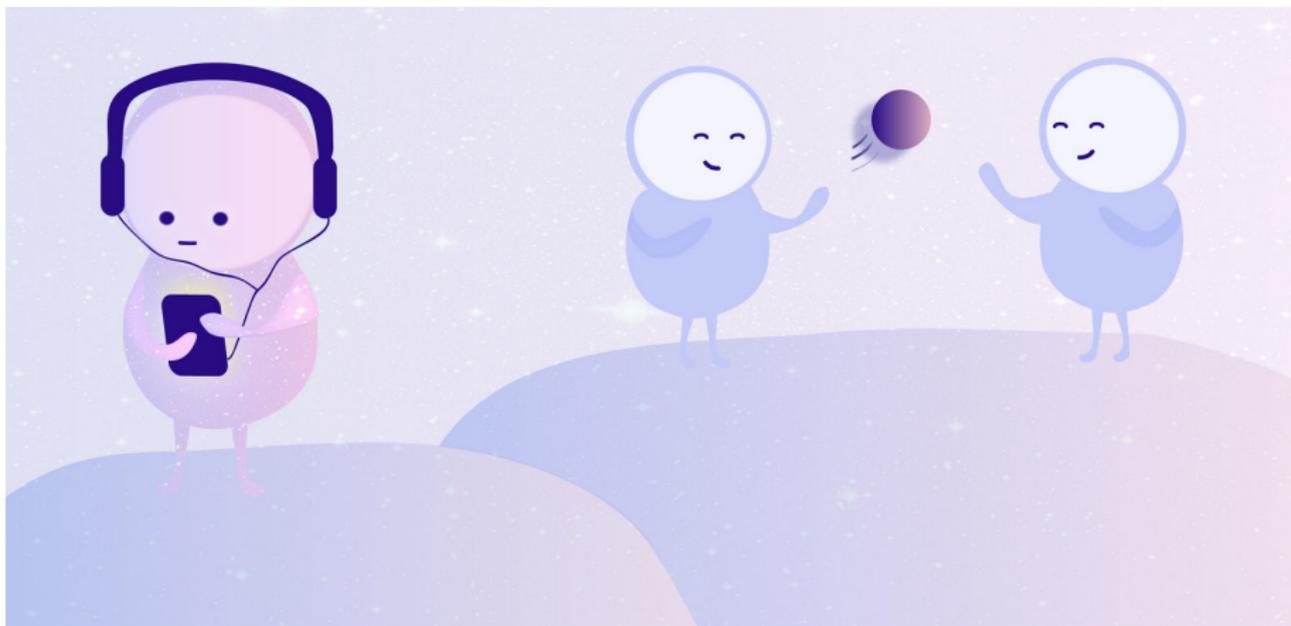
Perhaps something beyond the standard Higgs mechanism...

Can we probe the origin of neutrino mass at colliders?

- New fermions, gauge bosons, and/or scalars – **messengers of neutrino mass**.
- Rich phenomenology, both at hadron and lepton colliders, for messenger scale  $\lesssim \mathcal{O}(\text{few TeV})$ . [Deppisch, BD, Pilaftsis, [1502.06541](#); Cai, Han, Li, Ruiz, [1711.02180](#)]
- Complementarity with low-energy lepton number/ flavor violation searches.
- Possible connections to other sectors (e.g. NSI and oscillation physics, anomalies, baryogenesis, [dark matter](#)).

## SM-singlet Fermions

(aka sterile neutrinos/heavy neutrinos/heavy neutral leptons/right-handed neutrinos)



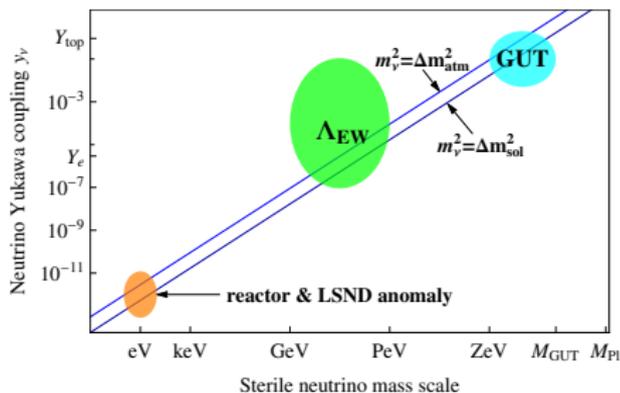
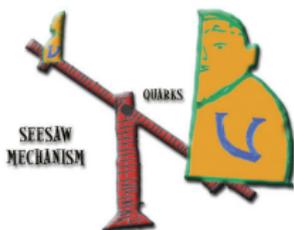
# Motivated from Type-I Seesaw

[Minkowski (PLB '77); Mohapatra, Senjanović (PRL '80); Yanagida '79; Gell-Mann, Ramond, Slansky '79; Glashow '80]

- SM-singlet **Majorana** fermions ( $N$ ):

$$-\mathcal{L} \supset Y_\nu \bar{L} \phi^c N + \frac{1}{2} M_N \bar{N}^c N + \text{H.c.}$$

- After EWSB,  $m_\nu \simeq -M_D M_N^{-1} M_D^T$ , where  $M_D = v Y_\nu$ .



[Antusch, Cazzato, Fischer, [1612.02728](#)]

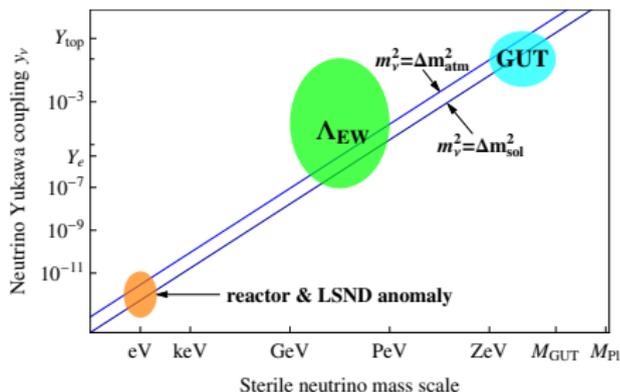
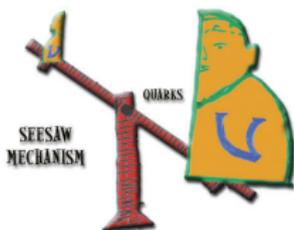
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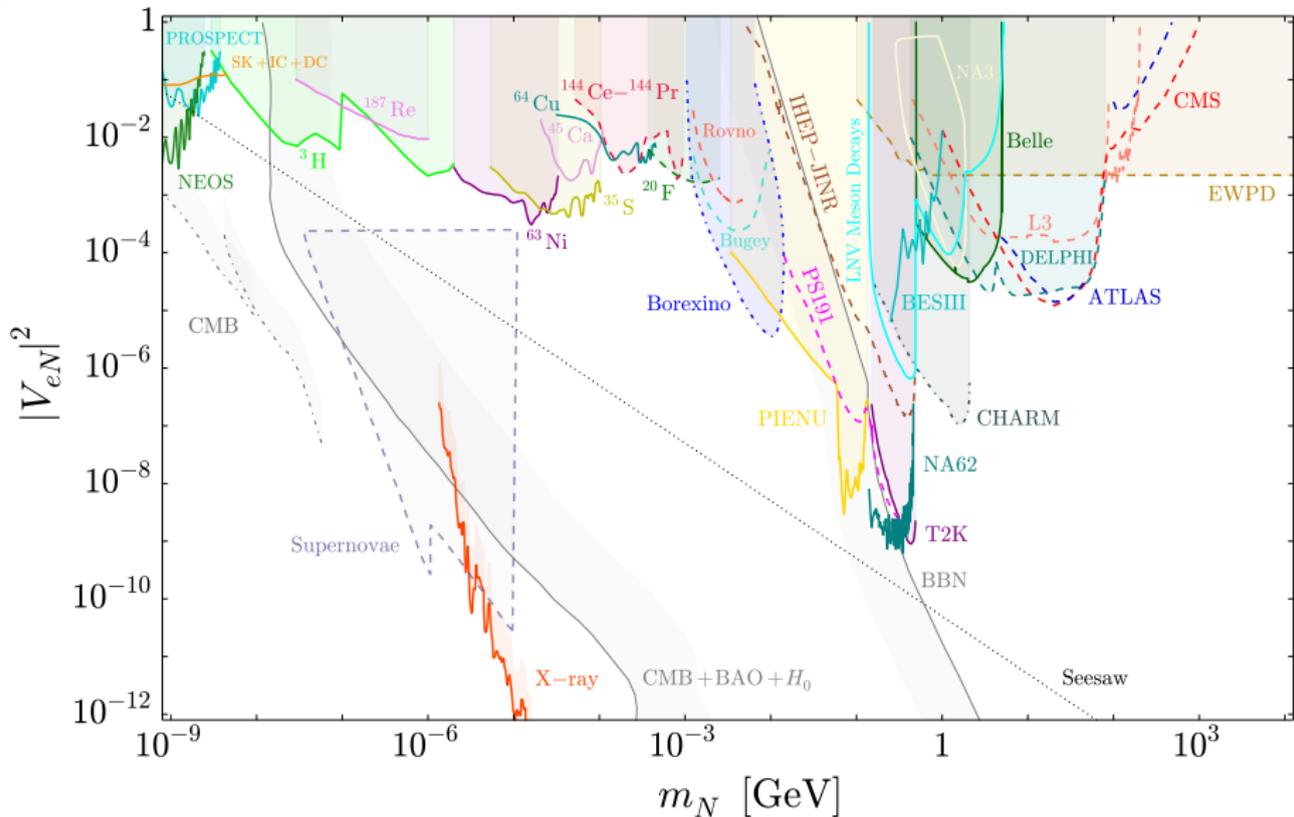
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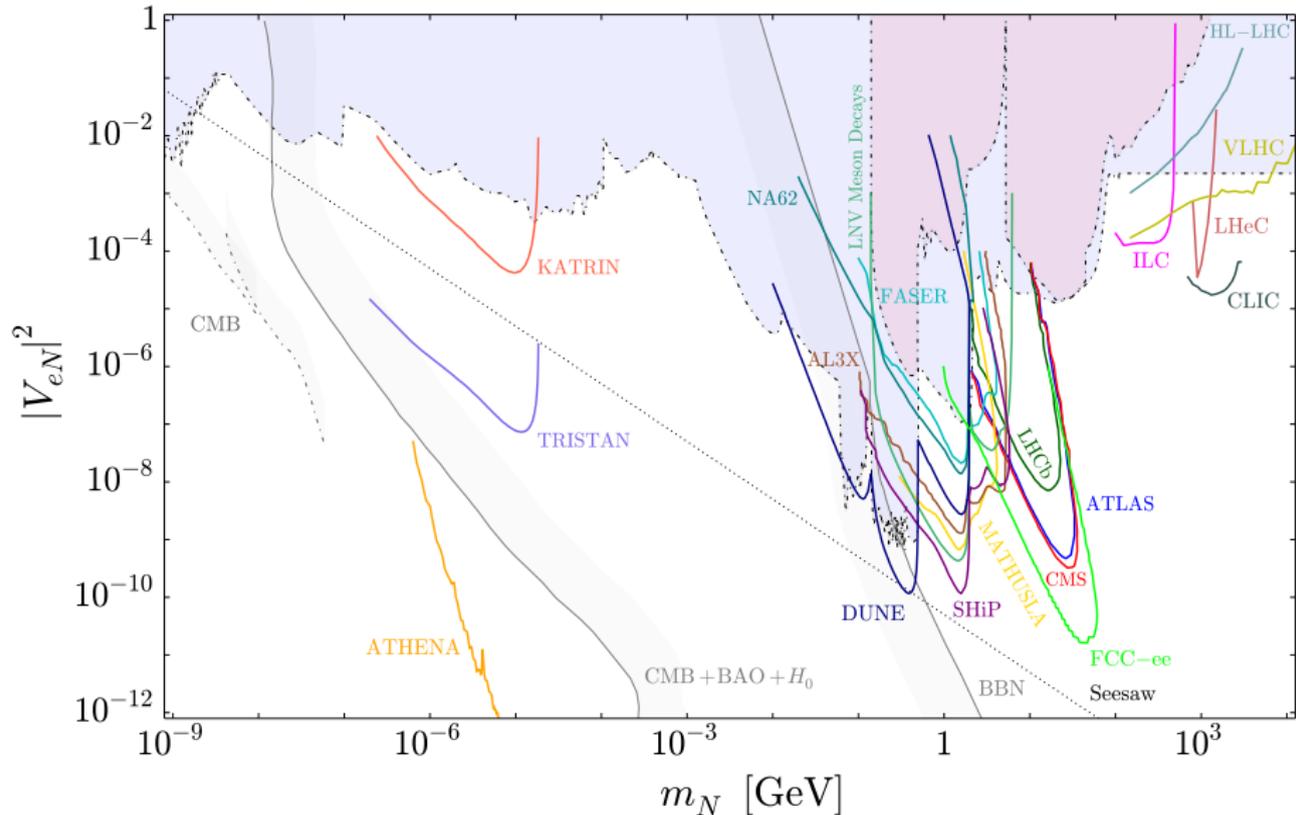
- Each  $N_i$  corresponds to  $m_{\nu_i} \neq 0$ . Need at least 2 (3 is better).
- Naturalness of Higgs mass suggests  $M_N \lesssim 10^7$  GeV.** [Vissani ([hep-ph/9709409](#)); Farina, Pappadopulo, Strumia ([1303.7244](#)); Clarke, Foot, Volkas ([1502.01352](#)); Bambhaniya, BD, Goswami, Khan, Rodejohann ([1611.03827](#))]

# Summary of Current Constraints on Sterile Neutrinos



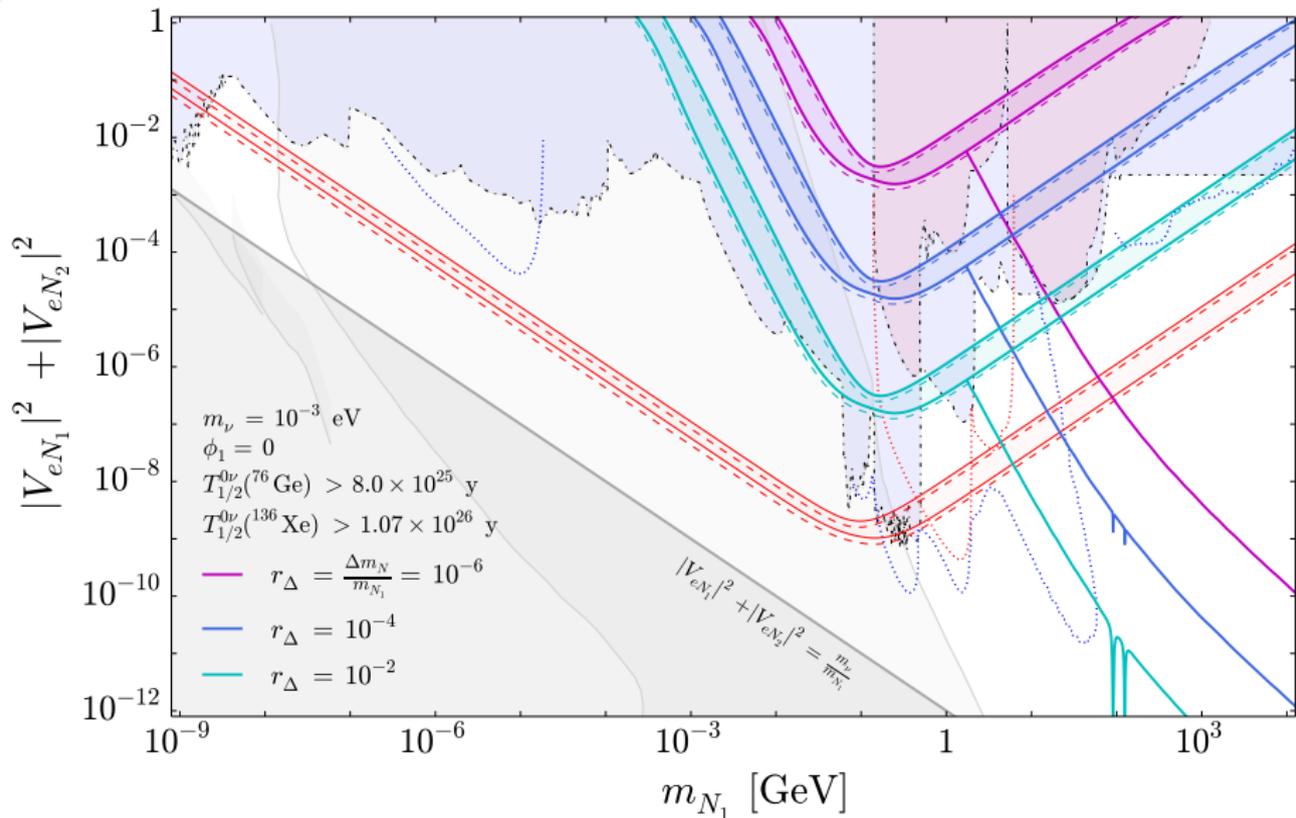
Bolton, Deppisch, BD, [1912.03058](https://arxiv.org/abs/1912.03058) (JHEP '20); see <http://sterile-neutrino.org> for regular updates and public code

# Future Prospects



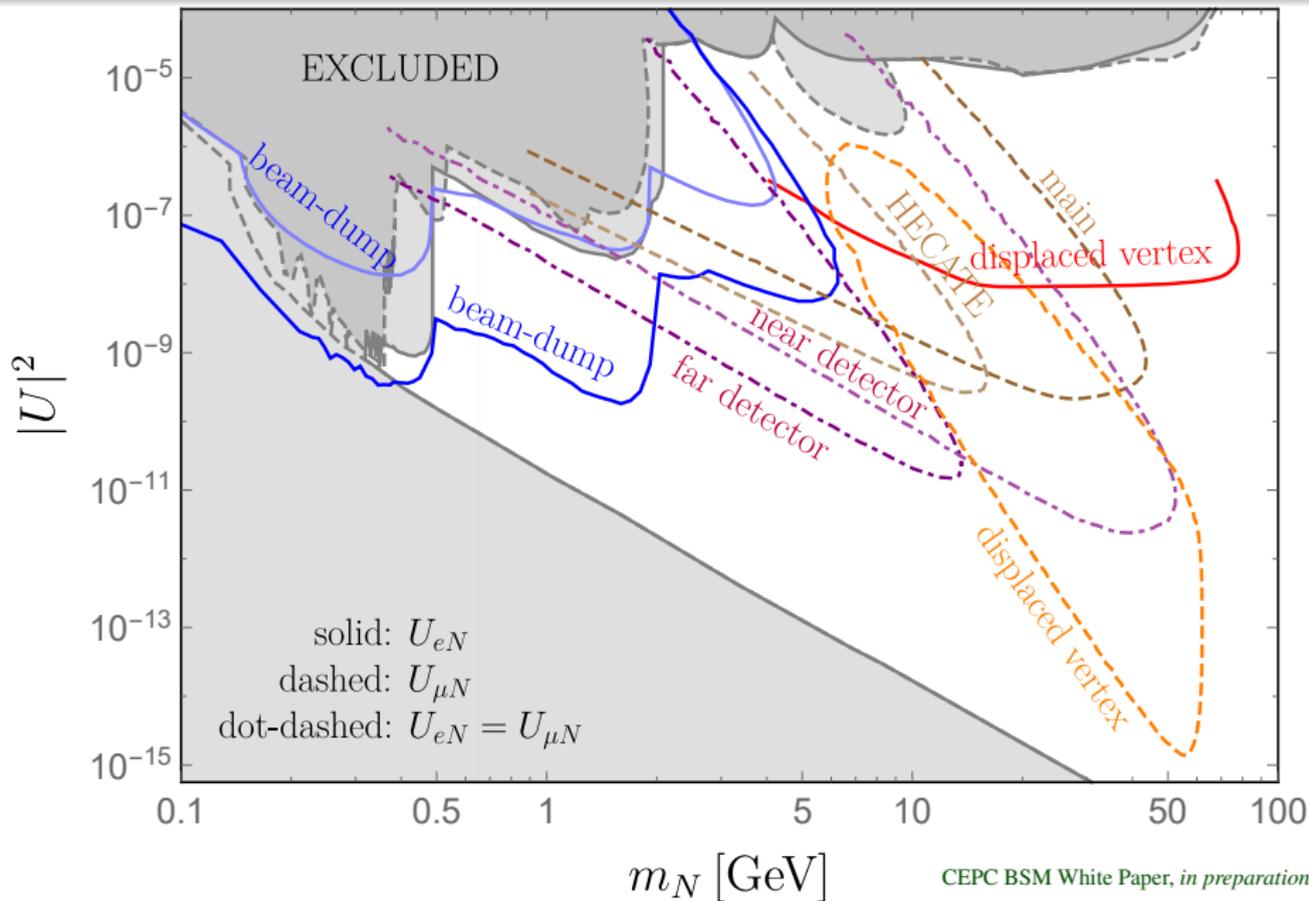
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# Neutrinoless Double Beta Decay



Bolton, Deppisch, BD, [1912.03058](#) (JHEP '20); see also Hernandez, Jones-Perez, Suarez-Navarro, [1810.07210](#) (EPJC '19)

# Future Prospects at CEPC



see previous talk by M. Drewes

## New Gauge Bosons

$(W', Z')$

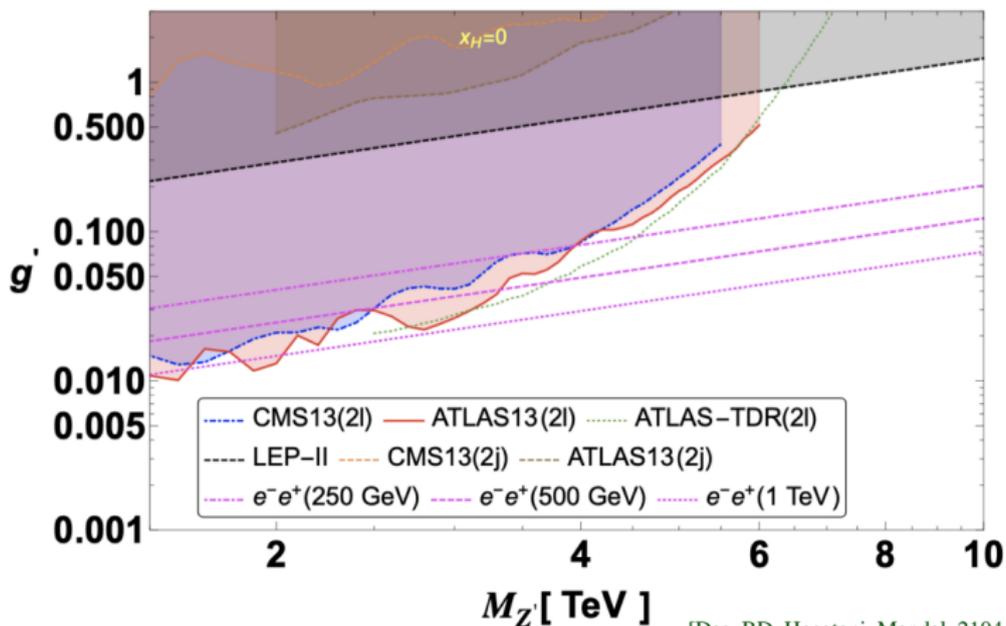


# $U(1)_X$ at Future Colliders

Gauge group	$q_L^i$	$u_R^i$	$d_R^i$	$\ell_L^i$	$e_R^i$	$N_R^i$	$H$	$\Phi$
$SU(3)_C$	<b>3</b>	<b>3</b>	<b>3</b>	<b>1</b>	<b>1</b>	<b>1</b>	<b>1</b>	<b>1</b>
$SU(2)_L$	<b>2</b>	<b>1</b>	<b>1</b>	<b>2</b>	<b>1</b>	<b>1</b>	<b>2</b>	<b>1</b>
$U(1)_Y$	1/6	2/3	-1/3	-1/2	-1	0	1/2	0
$U(1)_X$	$\frac{1}{6}x_H + \frac{1}{3}x_\Phi$	$\frac{2}{3}x_H + \frac{1}{3}x_\Phi$	$-\frac{1}{3}x_H + \frac{1}{3}x_\Phi$	$-\frac{1}{2}x_H - x_\Phi$	$-x_H - x_\Phi$	$-x_\Phi$	$-\frac{x_H}{2}$	$2x_\Phi$

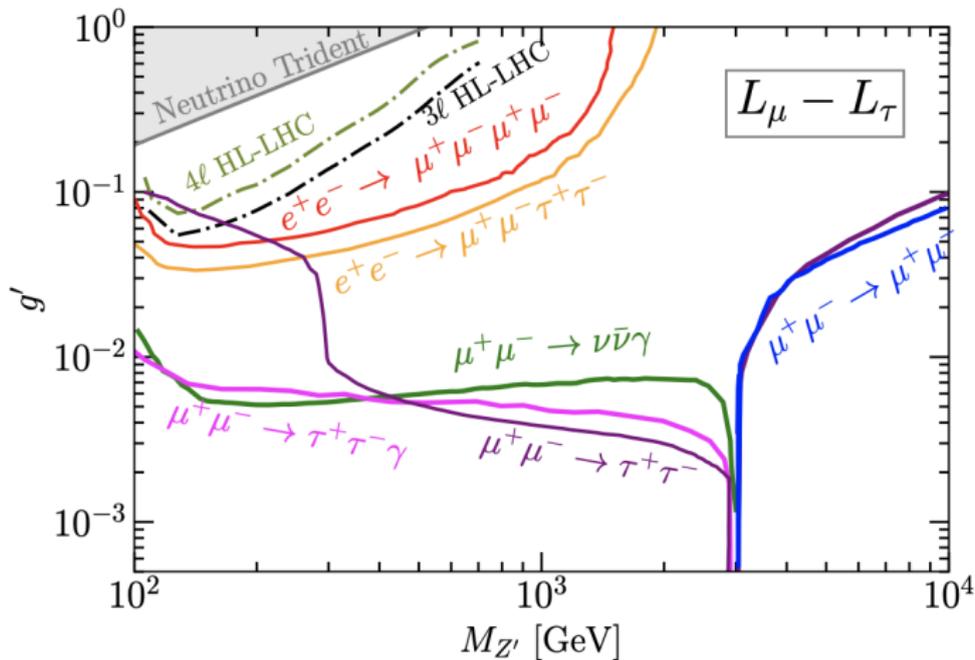
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$SU(2)_L$	<b>2</b>	<b>1</b>	<b>1</b>	<b>2</b>	<b>1</b>	<b>1</b>	<b>2</b>	<b>1</b>
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$U(1)_X$	$\frac{1}{6}x_H + \frac{1}{3}x_\Phi$	$\frac{2}{3}x_H + \frac{1}{3}x_\Phi$	$-\frac{1}{3}x_H + \frac{1}{3}x_\Phi$	$-\frac{1}{2}x_H - x_\Phi$	$-x_H - x_\Phi$	$-x_\Phi$	$-\frac{x_H}{2}$	$2x_\Phi$



# $U(1)_{L_\alpha - L_\beta}$ at Future Colliders

Gauge group	$L_e$	$L_\mu$	$L_\tau$	$e_R$	$\mu_R$	$\tau_R$	$H$	$\Phi$
$SU(3)_c$	1	1	1	1	1	1	1	1
$SU(2)_L$	2	2	2	1	1	1	2	1
$U(1)_Y$	$-\frac{1}{2}$	$-\frac{1}{2}$	$-\frac{1}{2}$	-1	-1	-1	$\frac{1}{2}$	0
$U(1)_{L_\mu - L_\tau}$	0	1	-1	0	1	-1	0	2

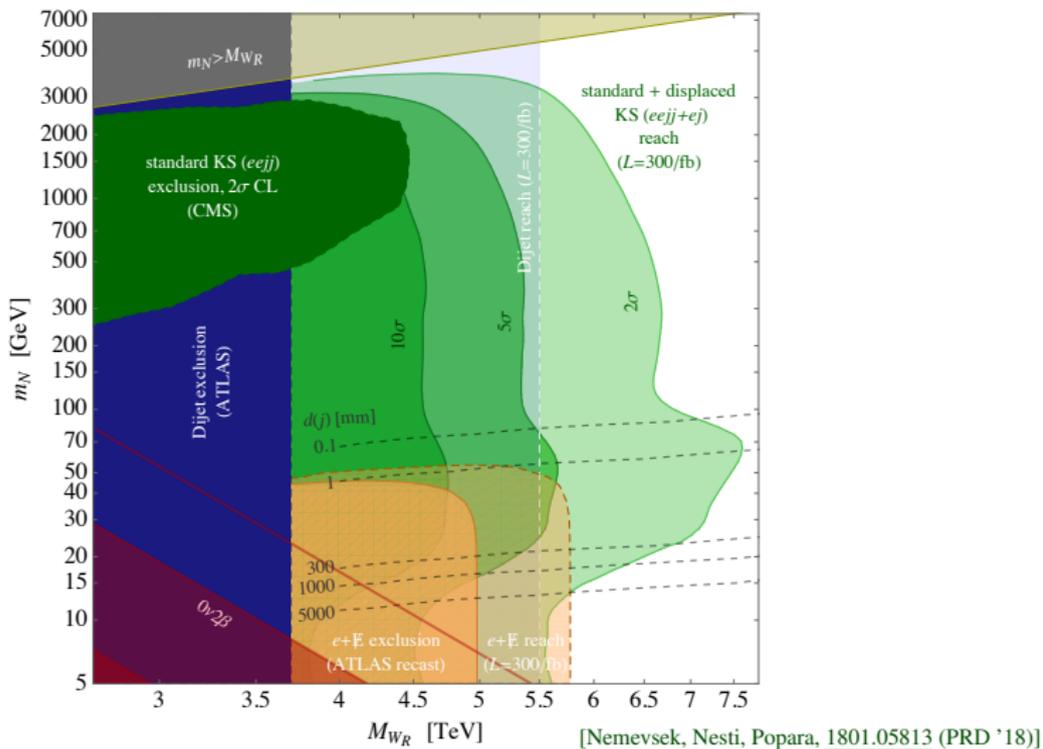


$$SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

- Parity restoration at high scale. [Mohapatra, Pati (PRD '75); Senjanovic, Mohapatra (PRD '75)]
- A natural UV-completion of seesaw. [Mohapatra, Senjanovic (PRD '81)]
- New contributions to collider signals. [Keung, Senjanovic (PRL '83); Chen, BD, Mohapatra, [1306.2342](#) (PRD '13)]

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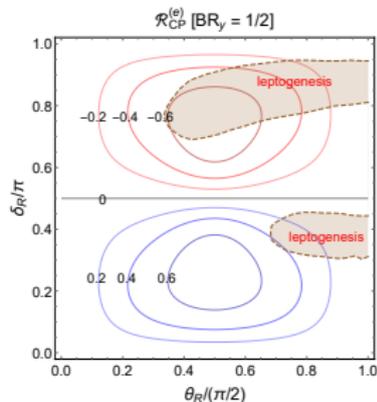
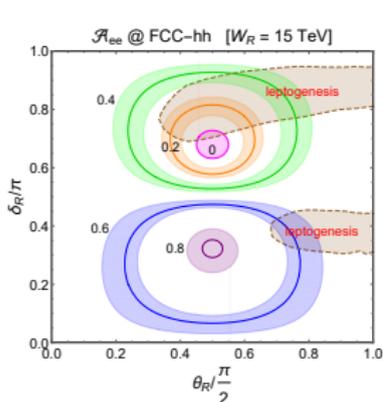
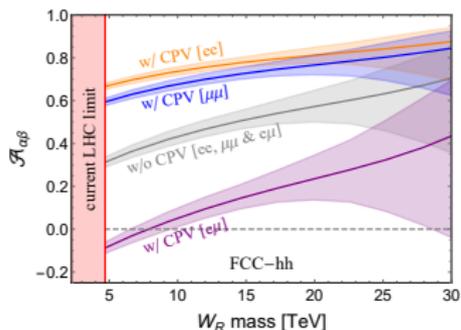
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# CP Violation in the RHN Sector

$$\begin{pmatrix} N_e \\ N_\mu \end{pmatrix} = \begin{pmatrix} \cos \theta_R & \sin \theta_R e^{-i\delta_R} \\ -\sin \theta_R e^{i\delta_R} & \cos \theta_R \end{pmatrix} \begin{pmatrix} N_1 \\ N_2 \end{pmatrix}.$$

$$\mathcal{A}_{\alpha\beta} \equiv \frac{\mathcal{N}(\ell_\alpha^+ \ell_\beta^+) - \mathcal{N}(\ell_\alpha^- \ell_\beta^-)}{\mathcal{N}(\ell_\alpha^+ \ell_\beta^+) + \mathcal{N}(\ell_\alpha^- \ell_\beta^-)}; \quad \mathcal{R}_{\text{CP}}^{(\ell)} \equiv \frac{\frac{\sigma(pp \rightarrow W_R^+ \rightarrow \ell^+ \ell^+ jj)}{\sigma(pp \rightarrow W_R^+ \rightarrow e^+ \mu^+ jj)} - \frac{\sigma(pp \rightarrow W_R^- \rightarrow \ell^- \ell^- jj)}{\sigma(pp \rightarrow W_R^- \rightarrow e^- \mu^- jj)}}{\frac{\sigma(pp \rightarrow W_R^+ \rightarrow \ell^+ \ell^+ jj)}{\sigma(pp \rightarrow W_R^+ \rightarrow e^+ \mu^+ jj)} + \frac{\sigma(pp \rightarrow W_R^- \rightarrow \ell^- \ell^- jj)}{\sigma(pp \rightarrow W_R^- \rightarrow e^- \mu^- jj)}}.$$

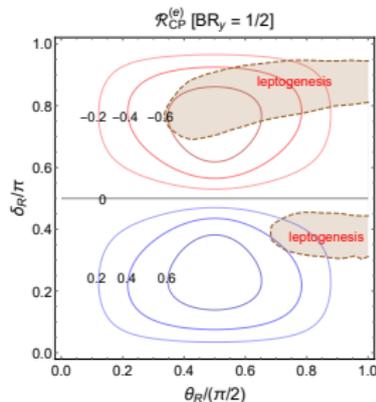
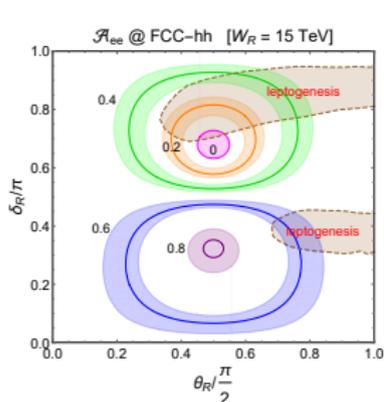
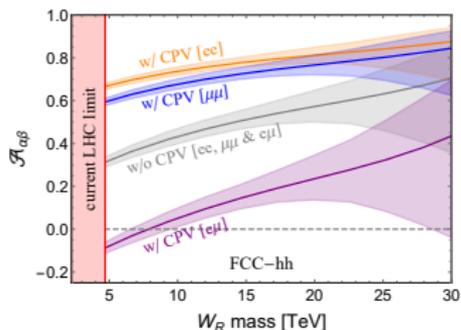


[BD, Mohapatra, Zhang, [1904.04787](#) (JHEP '19)]

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[BD, Mohapatra, Zhang, [1904.04787](#) (JHEP '19)]

- Lower bound on  $M_{W_R} \gtrsim 15$  TeV from leptogenesis.

[Frere, Hambye, Vertongen ([0806.0841](#)); BD, Lee, Mohapatra ([1503.04970](#))]

- FCC-hh will provide a direct collider test of thermal leptogenesis in LRSM.

## New Scalars



$$\phi(\mathbf{2}, \mathbf{2}, 0) = \begin{pmatrix} \phi_1^0 & \phi_2^+ \\ \phi_1^- & \phi_2^0 \end{pmatrix}, \quad \delta_R(\mathbf{1}, \mathbf{3}, 2) = \begin{pmatrix} \frac{\delta_R^+}{\sqrt{2}} & \delta_R^{++} \\ \delta_R^0 & -\frac{\delta_R^+}{\sqrt{2}} \end{pmatrix}$$

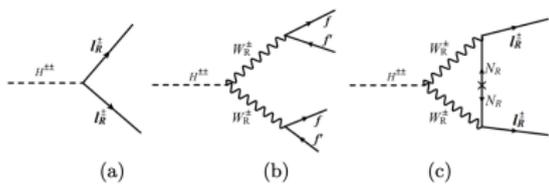
- $\langle \delta_R^0 \rangle \equiv v_R$  gives rise to RH Majorana neutrino masses  $\implies$  type-I seesaw.
- 8 (14 if  $\delta_L$  included) physical Higgs bosons: Rich phenomenology.

[Gunion, Grifols, Mendez, Kayser, Olness ([PRD '89](#)); Chao, Luo, Xing, Zhou ([PRD '08](#)); Fileviez Perez, Han, Huang, Li, Wang ([PRD '08](#)); Kanemura, Yagyu, Yokoya ([PLB '13](#)); Bambhaniya, Chakraborty, Gluza, Kordiaczyńska, Szafron ([JHEP '14](#)); BD, Mohapatra, Zhang ([JHEP '16](#)); Babu, Jana ([PRD '17](#)); BD, Ramsey-Musolf, Zhang ([PRD '18](#)); BD, Zhang ([JHEP '18](#)); ...]

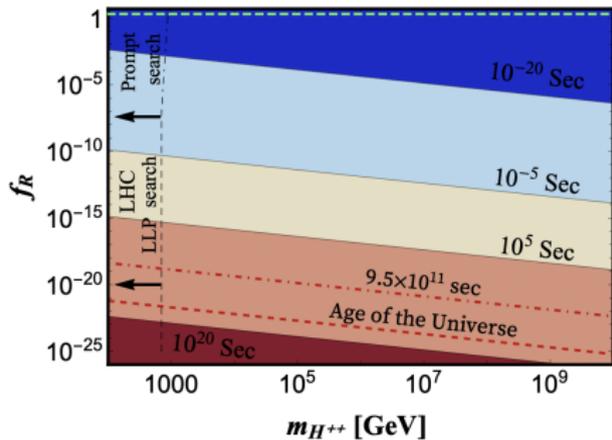
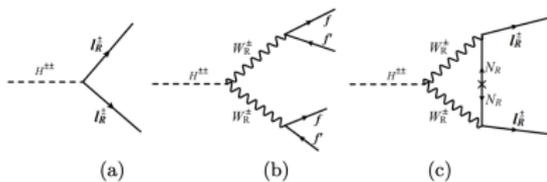
$$\phi(\mathbf{2}, \mathbf{2}, 0) = \begin{pmatrix} \phi_1^0 & \phi_2^+ \\ \phi_1^- & \phi_2^0 \end{pmatrix}, \quad \delta_R(\mathbf{1}, \mathbf{3}, 2) = \begin{pmatrix} \frac{\delta_R^+}{\sqrt{2}} & \delta_R^{++} \\ \delta_R^0 & -\frac{\delta_R^+}{\sqrt{2}} \end{pmatrix}$$

- $\langle \delta_R^0 \rangle \equiv v_R$  gives rise to RH Majorana neutrino masses  $\implies$  type-I seesaw.
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- But FCNC constraints require the bidoublet scalars ( $H_1^0, A_1^0, H_1^\pm$ ) to be very heavy  $\gtrsim 15$  TeV. [An, Ji, Mohapatra, Zhang (NPB '08); Bertolini, Maiezza, Nesti (PRD '14; PRD '20)] **Need FCC-hh.**
- Doubly-charged scalars ( $H^{\pm\pm}$ ) constrained to be  $\gtrsim 900$  (700) GeV from prompt (displaced) multilepton searches. [ATLAS, 2211.07505 (EPJC '23)]
- Neutral component ( $H_3^0$ ) is hadrophobic and can be much lighter!
- Can even be a dark matter candidate (albeit highly fine-tuned).  
[Nemevsek, Senjanovic, Zhang, 1205.0844 (JCAP '12)]

# Doubly-charged Scalar from $SU(2)_R$ Triplet

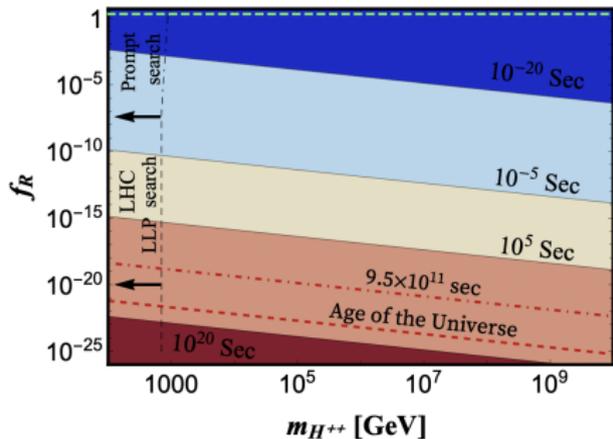
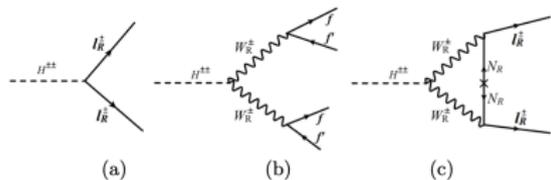


# Doubly-charged Scalar from $SU(2)_R$ Triplet



- Implications for HSCP searches at LHC.

# Doubly-charged Scalar from $SU(2)_R$ Triplet

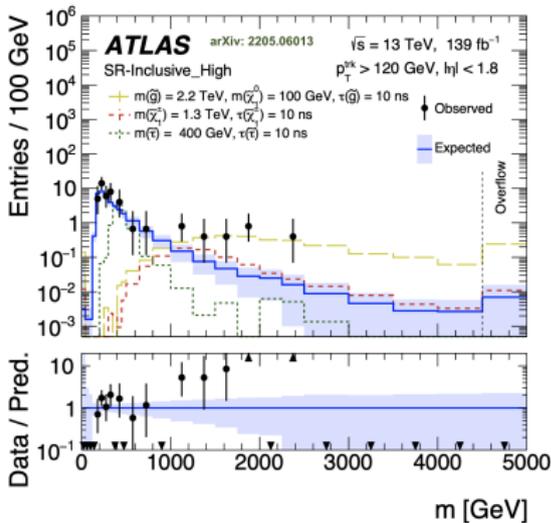


- Implications for HSCP searches at LHC.

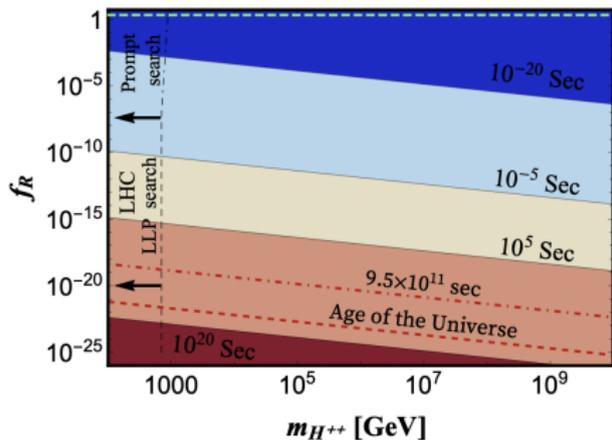
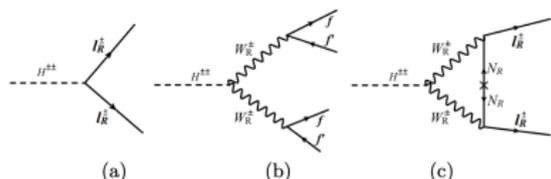
- Can explain the ATLAS  $dE/dx$  excess!

Akhmedov, BD, Jana, Mohapatra, [2401.15145](#) (PLB '24);

see also Giudice, McCullough, Teresi, [2205.04473](#) (JHEP '22)



# Doubly-charged Scalar from $SU(2)_R$ Triplet

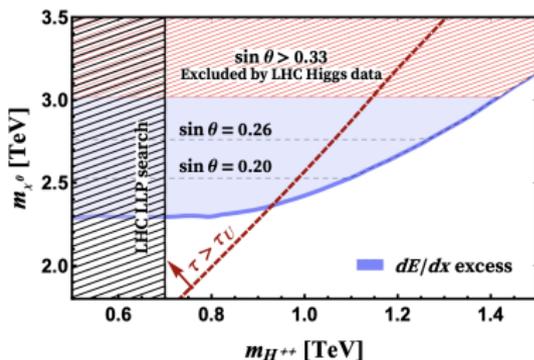
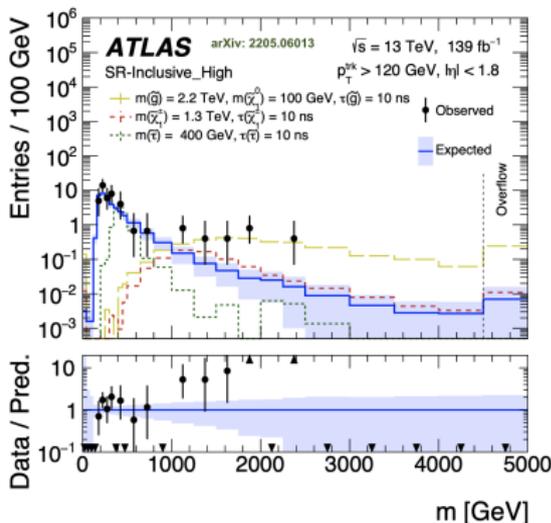


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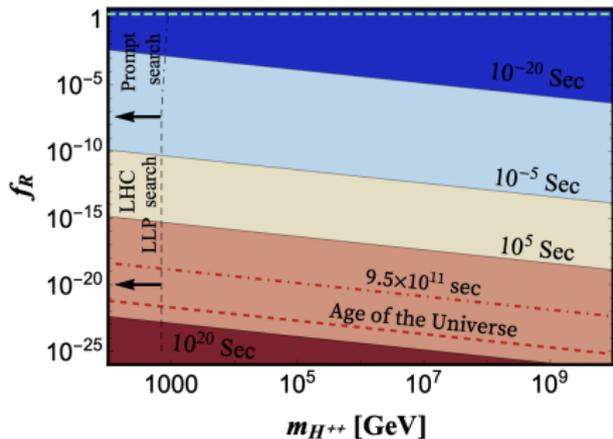
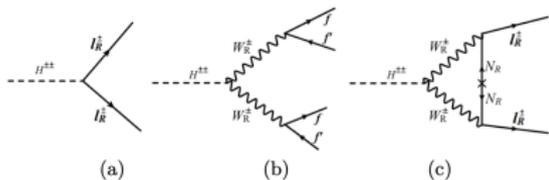
- Can explain the ATLAS  $dE/dx$  excess!

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# Doubly-charged Scalar from $SU(2)_R$ Triplet

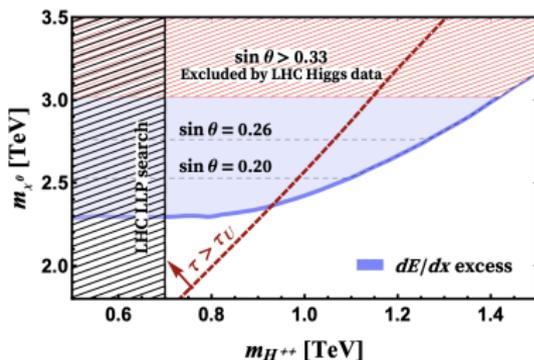
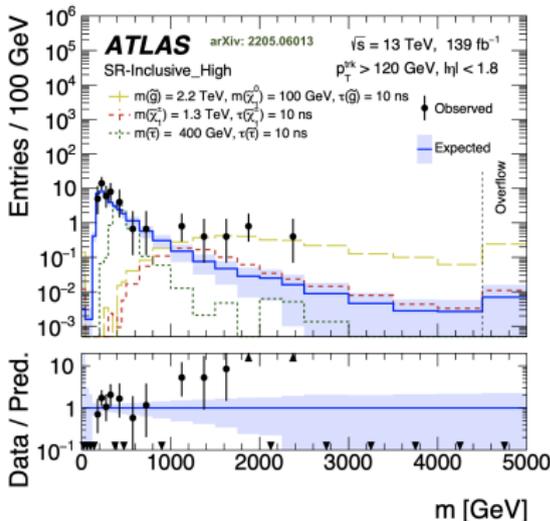


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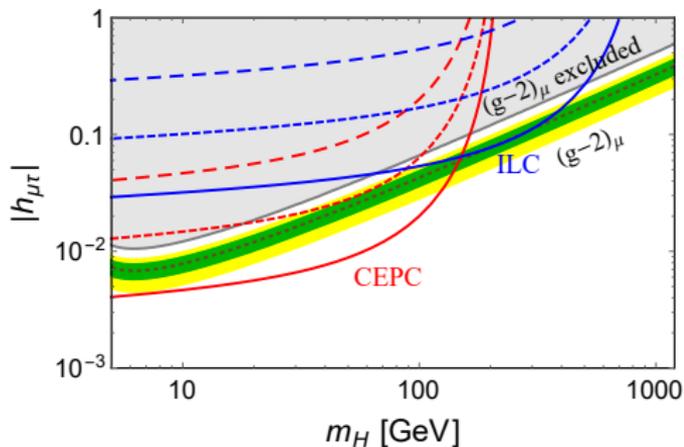
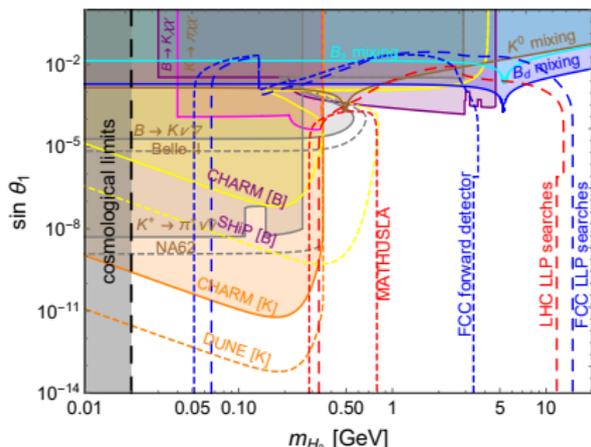
see also Giudice, McCullough, Teresi, [2205.04473](#) (JHEP '22)



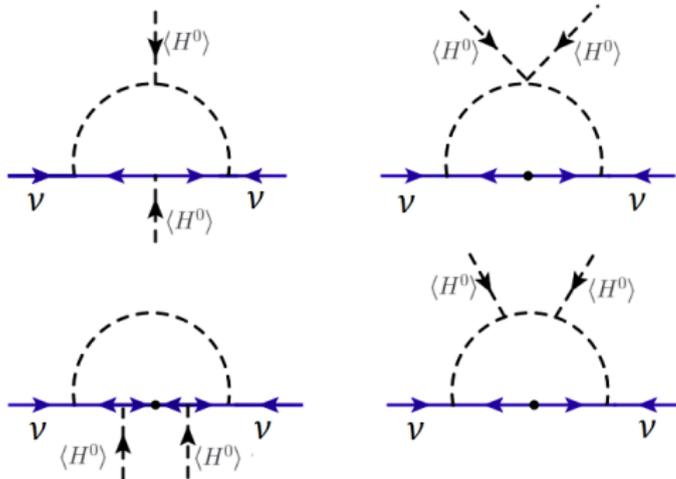
Catalyzed nuclear fusion to solve energy problem. [Akhmedov, [2109.13960](#) (PRD '21)]

# Neutral Scalar from $SU(2)_R$ Triplet

- Hadrophobic and allowed to be light (down to sub-GeV scale) by current constraints.
- Suppressed couplings to SM particles (either loop-level or small mixing).
- Necessarily long-lived at the LHC, with displaced vertex signals.
- Clean LFV signals at future lepton colliders.



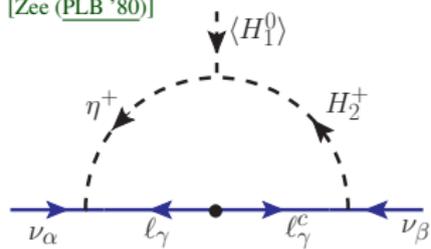
## Radiative Models (One-loop Only)



[Zee ([PLB '80](#)); Ma ([PRL '98](#)); Babu, Leung ([NPB '01](#)); de Gouvêa, Jenkins ([PRD '08](#)); Bonnet, Hirsch, Ota, Winter ([JHEP '12](#)); Cai, Clarke, Schmidt, Volkas ([JHEP '14](#)); Babu, BD, Jana, Thapa, ([JHEP '20](#)); Wang, Zhang, Zhou ([JHEP '23](#)); ...]

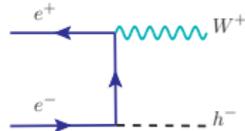
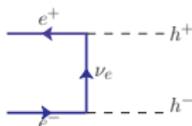
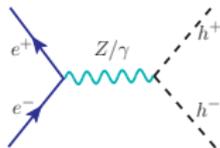
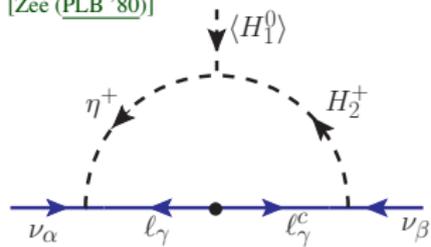
# Singlet Charged Scalar in Zee Model

[Zee (PLB '80)]



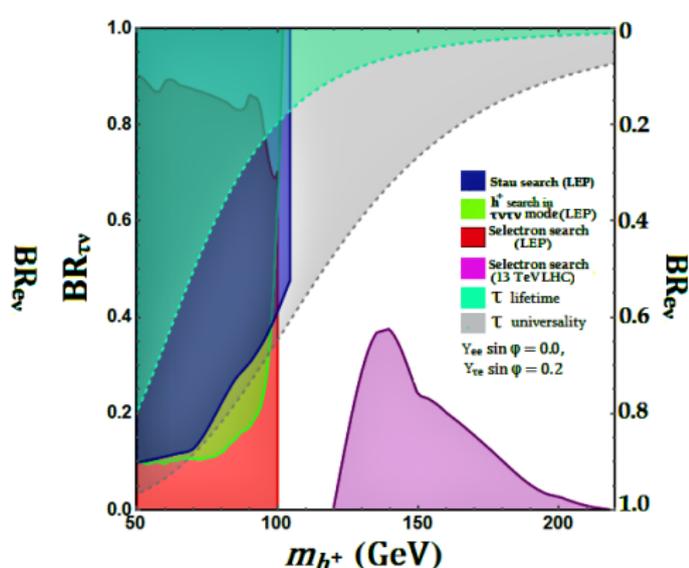
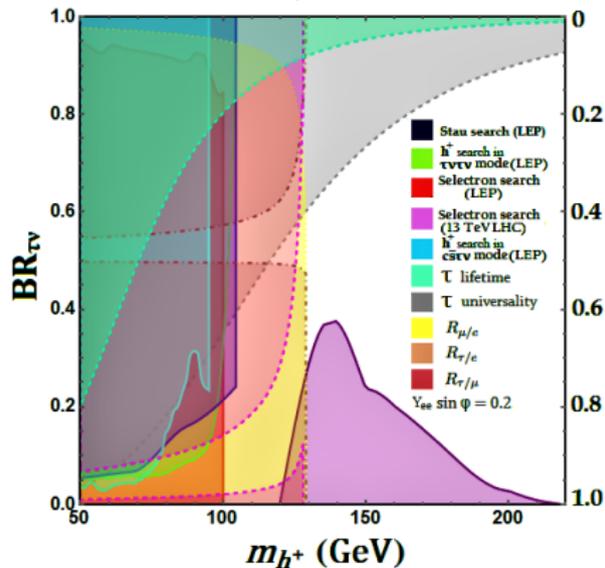
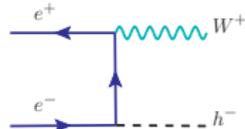
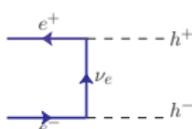
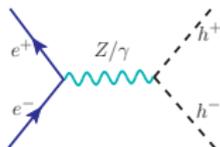
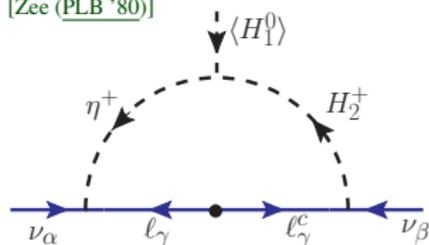
# Singlet Charged Scalar in Zee Model

[Zee (PLB '80)]



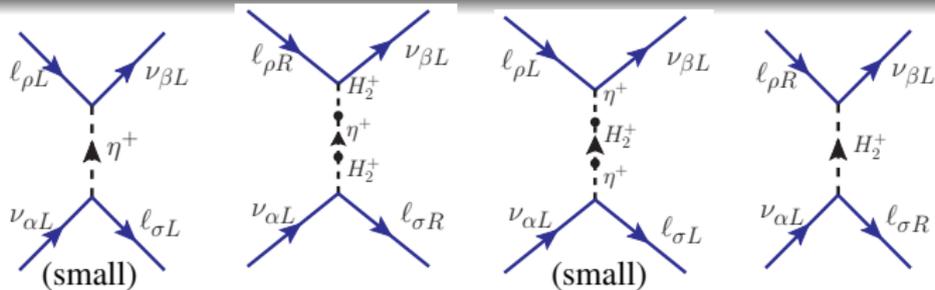
# Singlet Charged Scalar in Zee Model

[Zee (PLB '80)]



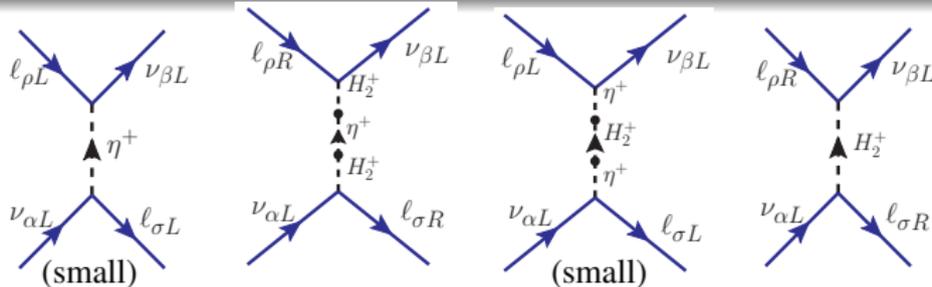
[Babu, BD, Jana, Thapa, [1907.09498](#) (JHEP '20)]

# Nonstandard Neutrino Interactions

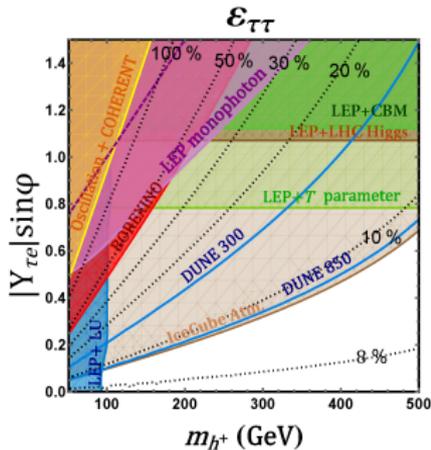
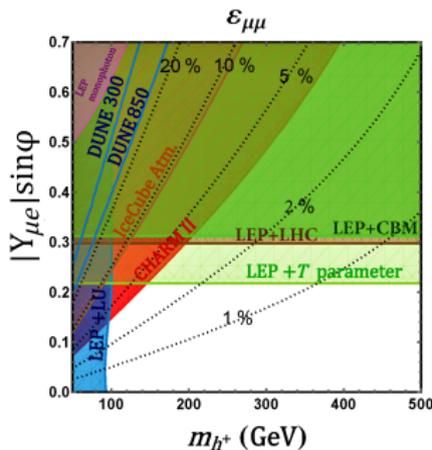
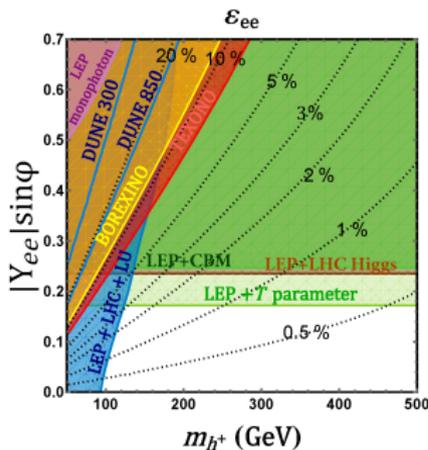


$$\varepsilon_{\alpha\beta} \equiv \varepsilon_{\alpha\beta}^{(h^+)} + \varepsilon_{\alpha\beta}^{(H^+)} = \frac{1}{4\sqrt{2}G_F} Y_{\alpha e} Y_{\beta e}^* \left( \frac{\sin^2 \varphi}{m_{h^+}^2} + \frac{\cos^2 \varphi}{m_{H^+}^2} \right).$$

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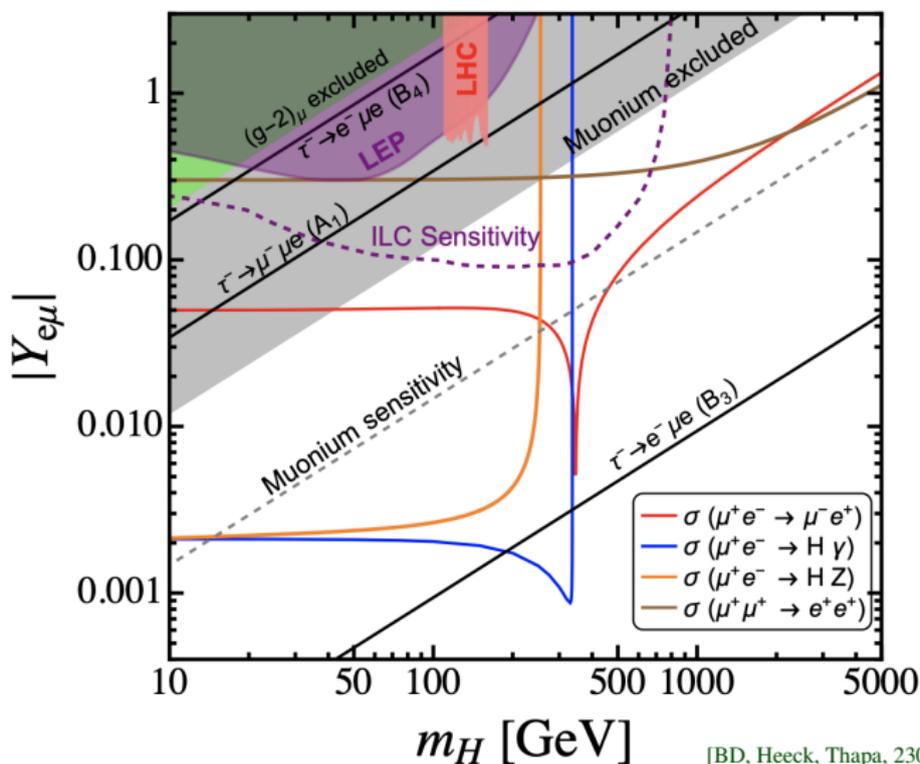
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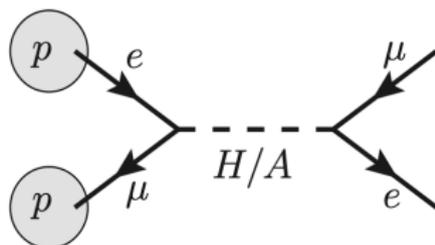


[Babu, BD, Jana, Thapa, [1907.09498](#) (JHEP '20)]

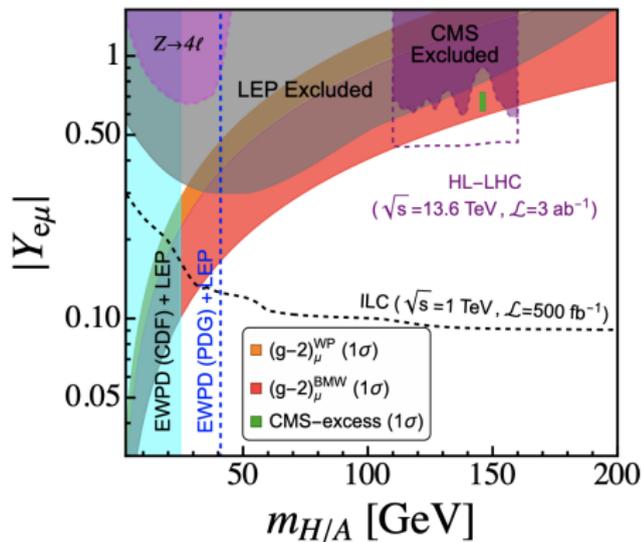
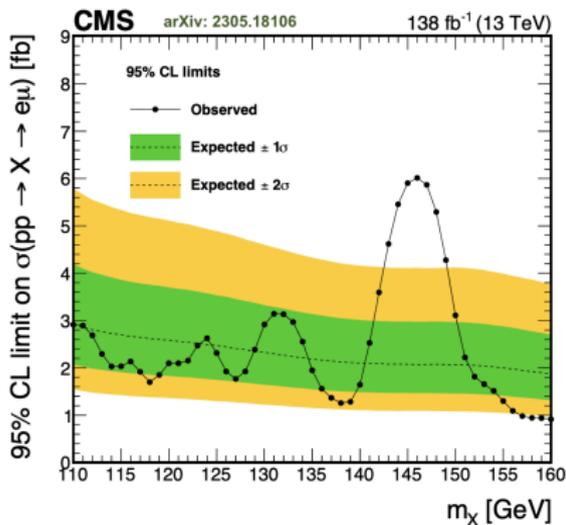
# Extra Neutral Scalars in Zee Model

- Must have LFV couplings (to fit neutrino oscillation data).
- Stringent cLFV constraints. But depend on Yukawa texture.
- Future lepton colliders could provide an independent test.





Using the lepton PDF of proton. [Buonocore, Nason, Tramontano, Zanderighi, [2005.06477](#) (JHEP '20)]



[Afik, BD, Thapa, [2305.19314](#) (PRD '24)]

# Conclusions

- Understanding the neutrino mass mechanism will provide key insights into the BSM world.
- Current and future colliders provide an ideal testing ground for (sub) TeV-scale neutrino mass models.
- Can probe the messenger particles (new fermions/gauge bosons/scalars) in a wide range of open parameter space.
- Complementary to the low-energy precision measurements.
- Important implications for current experimental anomalies.