Degenerate Oscillation in Neutron Star

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Fu, **SFG**, Guo & Wang [arXiv:2405.08591]



July 21, 2024 New Opportunities



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改道研究近 **Tsung-Dao Lee Institute**





- Degenerate Oscillation
- Neutron-Antineutron Oscillation in NS
- Neutron Star Cooling & GUT

Neutrino Oscillation





[Kayser, https://arxiv.org/abs/hep-ph/0506165]

$$\nu_{\alpha} = \sum_{i} \bigcup_{\alpha i} \nu_{i} \rightarrow \left[\sum_{i} \bigcup_{\alpha i} e^{i(\mathsf{E}_{i}\mathsf{t} - \vec{\mathsf{P}}_{i} \cdot \vec{\mathsf{x}})} \nu_{i} \right] = \left[\sum_{i} \bigcup_{\alpha i} \mathsf{P}_{i} \bigcup_{\beta i}^{\dagger} \nu_{\beta} \right] \equiv \sum_{\beta} \mathsf{A}_{\alpha\beta} \nu_{\beta}$$

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v Oscillation in Matter

Tsung-Dao Lee Institute

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Neutrino oscillations in matter

L. Wolfenstein

Carnegie-Mellon University, Pittsburgh, Pennsylvania 15213 (Received 6 October 1977; revised manuscript received 5 December 1977)

The effect of coherent forward scattering must be taken into account when considering the oscillations of neutrinos traveling through matter. In particular, for the case of massless neutrinos for which vacuum oscillations cannot occur, oscillations can occur in matter if the neutral current has an off-diagonal piece connecting different neutrino types. Applications discussed are solar neutrinos and a proposed experiment involving transmission of neutrinos through 1000 km of rock.

 $=\frac{\mathbf{MM}^{\dagger}}{2E_{\nu}}\pm\mathbf{V}$



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Collective Neutrino Oscillation



$$\mathrm{i}\frac{\mathrm{d}}{\mathrm{d}x}\psi = \mathsf{H}\psi$$

$$\begin{split} \mathsf{H} &= \frac{\Delta m^2}{2E} \mathsf{B} + \lambda \mathsf{L} + \mathsf{H}_{\nu\nu} \\ \mathsf{B} &= \mathsf{U} \left(\frac{1}{2} \mathrm{diag}[-1, 1] \right) \mathsf{U}^{\dagger} = \frac{1}{2} \begin{bmatrix} -\cos 2\theta_{\mathbf{v}} & \sin 2\theta_{\mathbf{v}} \\ \sin 2\theta_{\mathbf{v}} & \cos 2\theta_{\mathbf{v}} \end{bmatrix} \\ &= \frac{\Delta m^2}{2E} \mathsf{B} + \sqrt{2} G_{\mathrm{F}} n_e \mathsf{L} + \sqrt{2} G_{\mathrm{F}} \int \mathrm{d}^3 \mathbf{p}' (1 - \hat{\mathbf{p}} \cdot \hat{\mathbf{p}}') (\rho_{\mathbf{p}'} - \bar{\rho}_{\mathbf{p}'}) \\ &[\rho_{\mathbf{p}'}(t, \mathbf{x})]_{\alpha\beta} = \sum_{\nu'} n_{\nu', \mathbf{p}'}(t, \mathbf{x}) \langle \nu_{\alpha} | \psi_{\nu', \mathbf{p}'}(t, \mathbf{x}) \rangle \langle \psi_{\nu', \mathbf{p}'}(t, \mathbf{x}) | \nu_{\beta} \rangle, \\ &[\bar{\rho}_{\mathbf{p}'}(t, \mathbf{x})]_{\beta\alpha} = \sum_{\bar{\nu}'} n_{\bar{\nu}', \mathbf{p}'}(t, \mathbf{x}) \langle \bar{\nu}_{\alpha} | \psi_{\bar{\nu}', \mathbf{p}'}(t, \mathbf{x}) \rangle \langle \psi_{\bar{\nu}', \mathbf{p}'}(t, \mathbf{x}) | \bar{\nu}_{\beta} \rangle, \end{split}$$

Duan, Fuller & Qian [arXiv:1001.2799]

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Degeneracy with Boltzmann Eq.







$$\mathbb{C}[f] = \int rac{d^3 p_1}{(2\pi)^3 2E_1} \int rac{d^3 p_2}{(2\pi)^3 2E_2} \int rac{d^3 p_3}{(2\pi)^3 2E_3} \int rac{d^3 p_4}{(2\pi)^3 2E_4} (2\pi)^4 \delta^4(p_1 + p_2 - p_3 - p_4)
onumber \ |\mathcal{M}|^2 \{f_1 f_2 [1 \pm f_3] [1 \pm f_4] - f_3 f_4 [1 \pm f_1] [1 \pm f_2] \}$$

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Degeneracy is described by phase space factors: $1\pm f$

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$$|N\rangle \equiv \frac{(a^{\dagger})^{N}}{\sqrt{N!}}|0\rangle$$





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$$a^{\dagger}|N\rangle = \sqrt{N+1}|N+1\rangle$$
$$a|N\rangle = \sqrt{N}|N-1\rangle$$

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$$a|N\rangle = \sqrt{N}|N-1\rangle$$

$$a^{\dagger}a|N\rangle = N|N\rangle$$

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 $1 - f_{\nu_1}$

VS





 $1 - f_{\nu_1}$

 $|\Omega\rangle = f(\mathbf{x}, \mathbf{p})$

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VS





 $|\Omega\rangle \quad f(\mathbf{x}, \mathbf{p})$ $N \equiv \int f(\mathbf{x}, \mathbf{p}) d^3 \mathbf{x} d^3 \mathbf{p}$

 $1 - f_{\nu_1}$

Shao-Feng Ge [gesf@sjtu.edu.cn]

VS





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$$\begin{split} &|\Omega\rangle \quad f(\mathbf{x},\mathbf{p}) \\ &N \equiv \int f(\mathbf{x},\mathbf{p}) d^3 \mathbf{x} d^3 \mathbf{p} \\ &n_{\mathbf{p}} \equiv \int f(\mathbf{x},\mathbf{p}) d^3 \mathbf{x} \end{split}$$

 $1 - f_{\nu_1}$

VS

 $a_{\mathbf{p}}^{\dagger}$

 $a_{\mathbf{p}}$

 $1 - f_{\nu_1}$

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VS





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 $a^{\dagger}_{\mathbf{p}}$

 $a_{\mathbf{p}}$

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 $|\Omega
angle = f(\mathbf{x}, \mathbf{p})$

 $\hat{n}_{\mathbf{p}} = a_{\mathbf{p}}^{\dagger} a_{\mathbf{p}}$

 $N \equiv \int f(\mathbf{x}, \mathbf{p}) d^3 \mathbf{x} d^3 \mathbf{p}$

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VS





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 $|\Omega\rangle = f(\mathbf{x}, \mathbf{p})$ $N \equiv \int f(\mathbf{x}, \mathbf{p}) d^3 \mathbf{x} d^3 \mathbf{p}$ $n_{\mathbf{p}} \equiv \int f(\mathbf{x}, \mathbf{p}) d^3 \mathbf{x}$ $a_{\mathbf{p}}^{\dagger} \quad a_{\mathbf{p}}$ $\hat{n}_{\mathbf{p}} = a_{\mathbf{p}}^{\dagger} a_{\mathbf{p}}$ $a_{\mathbf{p}}^{\dagger}a_{\mathbf{p}}|\Omega\rangle = \int d^{3}\mathbf{x}f(\mathbf{x},\mathbf{p})|\Omega\rangle$





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专改道研究听



专改道研究听



$$\psi_i \equiv \sum_{\alpha} U^*_{\alpha i} \psi_{\alpha}$$



$$\psi_i \equiv \sum_{\alpha} U^*_{\alpha i} \psi_{\alpha} \qquad \psi_i \sim au e^{-ip_i \cdot x} + b^{\dagger} v e^{ip_i \cdot x}$$

SFG, Chui-Fan Kong, Pedro Pasquini [2310.04077]



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SFG, Chui-Fan Kong, Pedro Pasquini [2310.04077]

Fermion oscillation needs to involve 3 parts:

$$\mathcal{M}_{\beta\alpha} \equiv \mathcal{M}_d \left[\sum_i U_{\beta i} U_{\alpha i}^* e^{i p_i \cdot (x-y)} \langle \Omega | a_{p i} a_{p i}^{\dagger} | \Omega \rangle \right] \mathcal{M}_p$$

Spinors u & v combined into M_p & M_d

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Spinors u & v combined into M_p & M_d

Pauli blocking factor already appears in amplitude!

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Degenerate Oscillation in QFT

$$\mathcal{M}_{\beta\alpha} = \mathcal{M}_d \left[\sum_i U_{\beta i} U_{\alpha i}^* e^{i p_i \cdot (x-y)} (1-f_i) \right] \mathcal{M}_p$$

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Degenerate Oscillation in QFT

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Events detected w/o oscillation

$$N_{\alpha}(x=y) \propto \sum_{\beta} \left| \sum_{i} U_{\beta i} U_{\alpha i}^{*} (1-f_{i}) \right|^{2}$$
Degenerate Oscillation in QFT

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$$N_{\alpha \to \beta}(x-y) \propto \left| \sum_{i} U_{\beta i} U_{\alpha i}^* e^{i p_i \cdot (x-y)} (1-f_i) \right|^2$$

Degenerate Oscillation in QFT



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Fraction of events in β flavor

$$P_{\alpha\beta}(x-y) \equiv N_{\alpha\to\beta}(x-y)/N_{\alpha}(x=y)$$

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GUT & Baryon Number Violation





Table 1 $\,$

GUT model	Is $N - \bar{N}$ observable?	Implications
(NON SUSY)		
SU(5)	No	$\Delta(B-L) = 0$
$SU(2)_L \times SU(2)_R \times SU(4)_c$	Yes	$M_c \simeq 10^5 { m GeV}$
Minimal $SO(10)$	No	
E_6	No	
(SUSY GUT)		
$[SU(3)]^3$	Yes	Induced breaking of R-parity
SO(10)	No	

Table Caption: This table summarizes the observability of neutron-anti-neutron oscillation in various GUT models.

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Being neutral, neutron can have Majorana mass term:

 $\mathbf{H} \approx \begin{pmatrix} H_{11} & \delta m \\ \delta m & H_{22} \end{pmatrix}$



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Mixing between neutron & antineutron

$$\binom{n_1}{n_2} = \operatorname{U} \binom{n}{\bar{n}} \qquad \operatorname{U} = \begin{pmatrix} \cos\theta & \sin\theta \\ -\sin\theta & \cos\theta \end{pmatrix} \qquad \tan 2\theta = \frac{2\delta m}{H_{22} - H_{11}}$$



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$$P_{n\bar{n}} = \frac{c^2 s^2 (1 - f_1)^2 + c^2 s^2 (1 - f_2)^2}{c^2 (1 - f_1)^2 + s^2 (1 - f_2)^2}$$
$$-\frac{2(1 - f_1)(1 - f_2)c^2 s^2 \cos(\Delta Et)}{c^2 (1 - f_1)^2 + s^2 (1 - f_2)^2}$$

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$$-\frac{2(1 - f_1)(1 - f_2)c^2 s^2 \cos(\Delta Et)}{c^2 (1 - f_1)^2 + s^2 (1 - f_2)^2} \quad \Delta E \equiv \sqrt{(H_{11} - H_{22})^2 + 4\delta m^2}$$

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reduces to the usual one with $f_i \rightarrow 0$

$$P_{n\bar{n}} = 4c^2 s^2 \sin^2\left(\frac{\Delta Et}{2}\right)$$



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in a dense neutron environment $f_1 \rightarrow 1, f_2 \rightarrow 0$



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$$s \ll 1 - f_1 \ll 1 \qquad P_{n\bar{n}} \to \frac{s^2}{(1 - f_1)^2} \ll 1$$

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 $P_{n\bar{n}} \to c^2 \sim 1$ **X**
 $s \ll 1 - f_1 \ll 1$ $P_{n\bar{n}} \to \frac{s^2}{(1 - f_1)^2} \ll 1$ **V**

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 $\approx c^2$

$$P_{n\bar{n}} \approx \frac{s^2}{(1-f_1)^2} \qquad \qquad P_{nn}$$

The degeneracy effect already appears at zero distance!



$$P_{n\bar{n}} \approx \frac{s^2}{(1-f_1)^2} \qquad \qquad P_{nn} \approx c^2$$

The degeneracy effect already appears at zero distance!

Standing fraction of antineutron in neutron star.



$$P_{n\bar{n}} \approx \frac{s^2}{(1-f_1)^2} \qquad \qquad P_{nn} \approx c^2$$

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Standing fraction of antineutron in neutron star.

$$f_1(\mathbf{p}) = \frac{1}{e^{\frac{\varepsilon_n(\mathbf{p})-\mu}{T}} + 1}$$



$$P_{n\bar{n}} \approx \frac{s^2}{(1-f_1)^2} \qquad \qquad P_{nn} \approx c^2$$

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Standing fraction of antineutron in neutron star.

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The degeneracy effect already appears at zero distance! Standing fraction of antineutron in neutron star.

$$f_1(\mathbf{p}) = \frac{1}{e^{\frac{\varepsilon_n(\mathbf{p})-\mu}{T}} + 1} \approx f_n$$
$$R(\mathbf{p}) \equiv \frac{f_{\bar{n}}}{f_n} = \frac{P_{n\bar{n}}}{P_{nn}} \approx \frac{\tan^2 \theta}{[1 - f_1(\mathbf{p})]^2}$$



$$P_{n\bar{n}} \approx \frac{s^2}{(1-f_1)^2} \qquad \qquad P_{nn} \approx c^2$$

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$$P_{n\bar{n}} \approx \frac{s^2}{(1-f_1)^2} \qquad \qquad P_{nn} \approx c^2$$

The degeneracy effect already appears at zero distance! Standing fraction of antineutron in neutron star.

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Note that only neutron is in thermal equilibrium!

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$$f_1(\mathbf{p}) = \frac{1}{e^{\frac{\varepsilon_n(\mathbf{p})-\mu}{T}} + 1} \qquad \varepsilon_n(\mathbf{p}) = \sqrt{\mathbf{p}^2 + m^2}$$

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Degenerate fermi gas:

$$k_F = (3\pi^2 n_n)^{1/3}$$

1



Fu, SFG, Guo & Wang [arXiv:2405.08591]

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1



$$f_1(\mathbf{p}) = \frac{1}{e^{\frac{\varepsilon_n(\mathbf{p}) - \mu}{T}} + 1} \qquad \varepsilon_n(\mathbf{p}) = \sqrt{\mathbf{p}^2 + m^2}$$

$$k_F = (3\pi^2 n_n)^{1/3}$$
 $\mu = \sqrt{k_F^2 + m^2}$



Fu, SFG, Guo & Wang [arXiv:2405.08591]

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Degenerate fermi gas:

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$$f_{\bar{n}} \approx \frac{s^2}{[1 - f_1(p)]^2} f_1$$

$$n_{\bar{n}} = \frac{m^2 T}{2\pi^2} e^{\frac{\mu}{T}} \left[2K_2 \left(\frac{m}{T}\right) + e^{\frac{\mu}{T}} K_2 \left(\frac{2m}{T}\right) \right] s^2$$



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Huge enhancement

$$\frac{\mu - m}{T} \sim \mathcal{O}(10^4) \qquad \qquad e^{\frac{\mu - m}{T}} \sim 10^{\mathcal{O}(1000)}$$

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Antineutron Annihilation



During a single annihilation, nucleon number reduces by 2

$$\mathrm{d}N = -2\langle \sigma v \rangle n_n n_{\bar{n}} \mathrm{d}t \mathrm{d}V$$

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$$n_n = 10^{38} \,\mathrm{cm}^{-3}$$

 $\langle \sigma v \rangle \approx 10^{-15} \,\mathrm{cm}^3/\mathrm{s}$

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NS Cooling Data





Potekhin, Zyuzin, Yakovlev, Beznogov & Shibanov MNRAS 496, 5052-5071 (2020) [arXiv:2006.15004]





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$$R \approx \left(\frac{T}{T_0}\right)^{\frac{3}{2}} e^{\frac{2(\mu-m)}{T} - \frac{2(\mu_0 - m)}{T_0}} R_0 \lesssim 10^{-43}$$

 $\log_{10} R_0 \lesssim -\mathcal{O}(10^4) \qquad R_0 \equiv R(T_0, n_0)$ $T_0 \equiv 9.9 \times 10^9 \text{K}$ $n_0 \equiv 5.8 \times 10^{38} \text{cm}^{-3}$



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Heating power

$$2Rn_n^2 \langle \sigma v \rangle Vm \sim 10^{24} \mathrm{W}$$

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Cooling power

$$2Rn_n^2 \langle \sigma v \rangle Vm \sim 10^{24} \mathrm{W}$$

$$4\pi R^2 \sigma T^4 \sim 10^{21} \,\mathrm{W}$$



 $\mathbf{H} \approx \begin{pmatrix} H_{11} & \delta m \\ \delta m & H_{22} \end{pmatrix}$





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 $s \sim \frac{\delta m}{\Lambda H}$

 $\delta m \sim \frac{\Lambda_{\rm QCD}^6}{M_V^5} \qquad \Lambda_{\rm QCD} \sim 180 \,{\rm MeV}$



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$$\Delta H = \mathcal{O}(\mathrm{MeV}) \sim \mathcal{O}(\mathrm{GeV})$$



$$\begin{split} \mathbf{H} &\approx \begin{pmatrix} H_{11} & \delta m \\ \delta m & H_{22} \end{pmatrix} \qquad s \sim \frac{\delta m}{\Delta H} \\ \delta m &\sim \frac{\Lambda_{\rm QCD}^6}{M_X^5} \qquad \Lambda_{\rm QCD} \sim 180 \, {\rm MeV} \\ \Delta H &= \mathcal{O}({\rm MeV}) \sim \mathcal{O}({\rm GeV}) \end{split}$$

 $s^2 \lesssim 10^{-\mathcal{O}(10^4)}$



$$H \approx \begin{pmatrix} H_{11} & \delta m \\ \delta m & H_{22} \end{pmatrix} \qquad s \sim \frac{\delta m}{\Delta H} \\ \delta m \sim \frac{\Lambda_{\rm QCD}^6}{M_X^5} \qquad \Lambda_{\rm QCD} \sim 180 \,\,{\rm MeV} \\ \Delta H = \mathcal{O}({\rm MeV}) \sim \mathcal{O}({\rm GeV}) \\ s^2 \lesssim 10^{-\mathcal{O}(10^4)} \qquad \Longrightarrow \qquad M_X > 10^{\mathcal{O}(1000)} \,{\rm GeV}$$



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$$s^2 \lesssim 10^{-\mathcal{O}(10^4)} \qquad \Longrightarrow \qquad M_X > 10^{\mathcal{O}(1000)} \,{\rm GeV}$$

which is far beyond the Planck scale!

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• Degenerate Oscillation

 Consistent picture of degeneracy in external & intermediate state

Neutron-Antineutron Oscillation in NS

- **1.** Concrete realization with n-nbar oscillation
- **2.** Standing fraction of antineutron

Neutron Star Cooling & GUT

- **1. Degeneracy enhancement**
- 2. Very strong constraint



大波道研究町 Tsung-Dao Lee Institute

Thank You

Superfluidity in NS



