

Future High Precision Measurements on the Spin Structure g_1 at EIC

Win Lin (林婉)

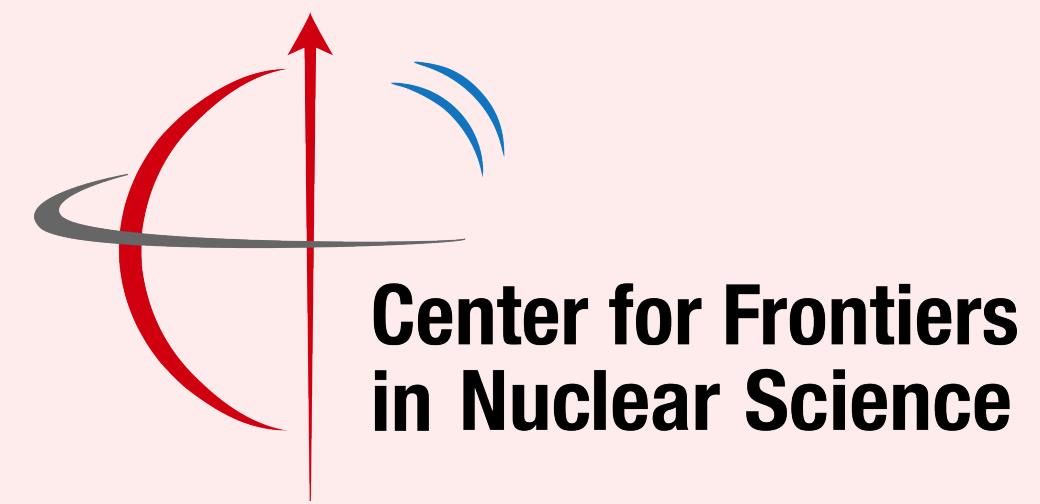
Stony Brook University

SPIN 2025

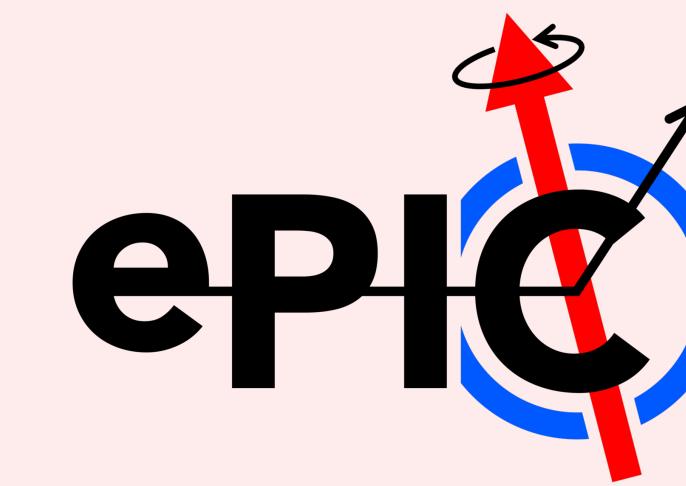
09/24/2025



Stony Brook
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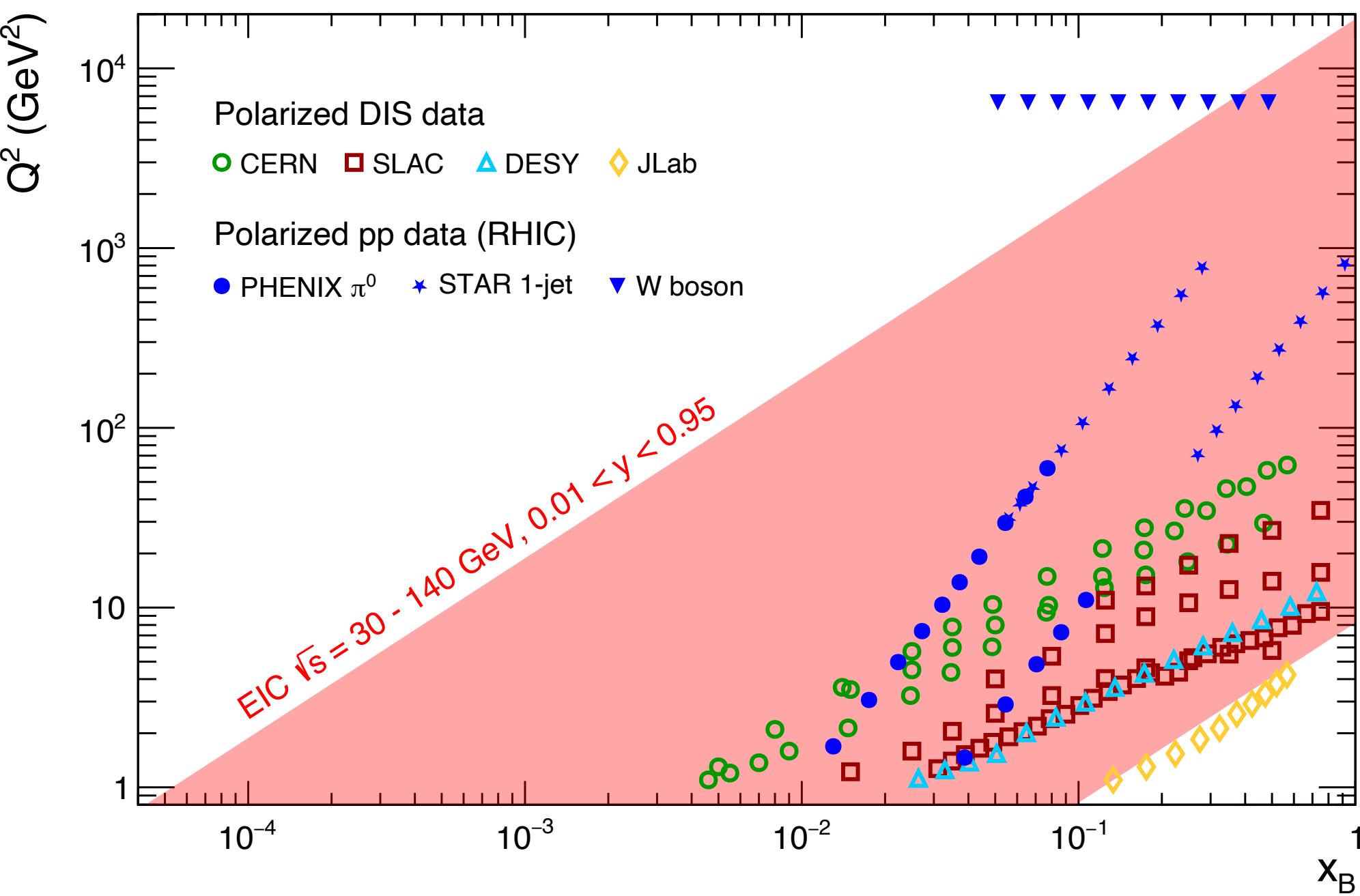
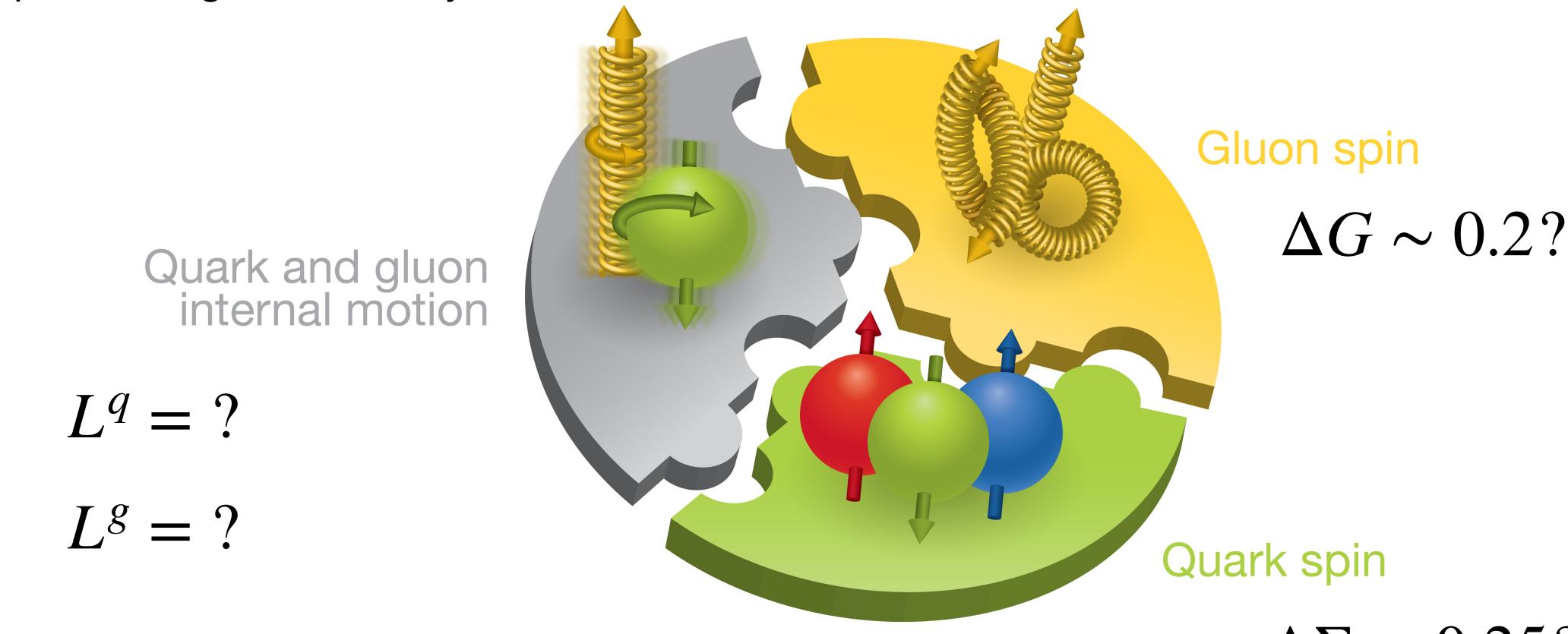


Center for Frontiers
in Nuclear Science



The Nucleon Spin Puzzle

<https://doi.org/10.1103/PhysRevLett.113.012001>



<https://doi.org/10.48550/arXiv.1212.1701>

Spin composition (Jaffe-Manohar sum):

$$\frac{1}{2} = \frac{1}{2} \Delta \Sigma + \Delta G + L^q + L^g$$

Quark-Parton Model:

$$\Gamma_1^{p(n)} = \int_0^1 g_1^{p(n)} dx = \frac{1}{12} \left[+(-)a_3 + \frac{1}{3}a_8 \right] + \frac{1}{9}a_0$$

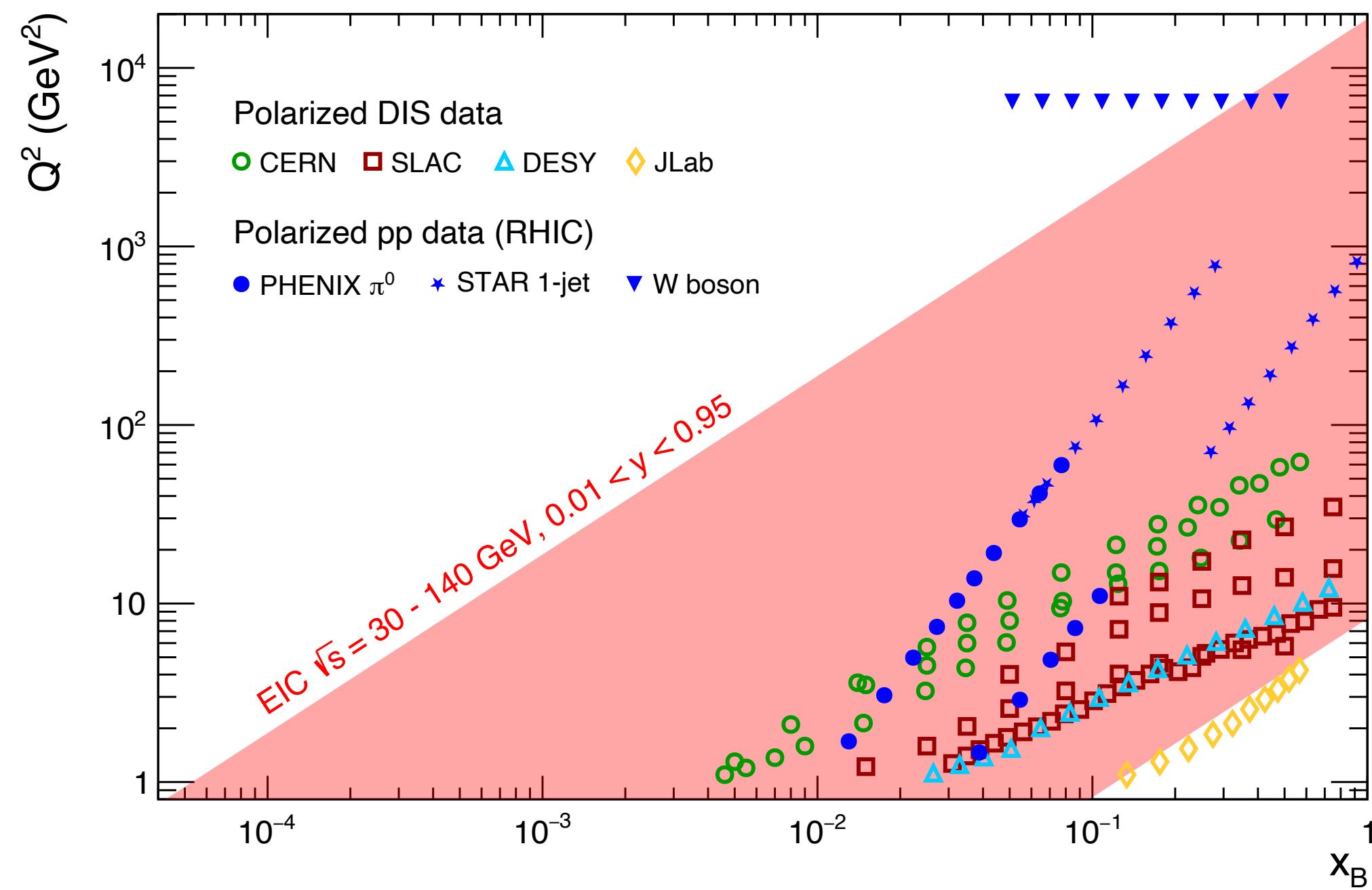
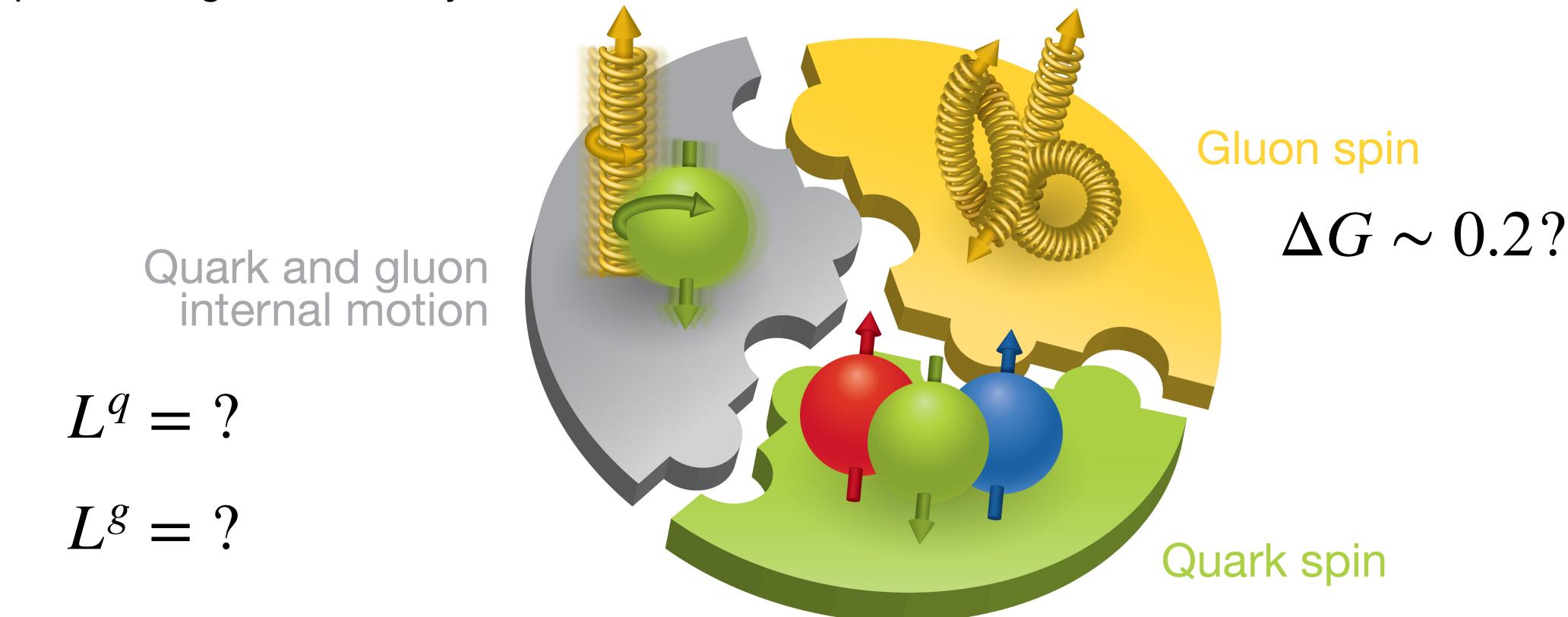
$$a_0 = \Delta u + \Delta d + \Delta s \equiv \boxed{\Delta \Sigma}$$

$$a_3 = \Delta u - \Delta d = \left| \frac{g_A}{g_V} \right|$$

$$a_8 = \Delta u + \Delta d - 2\Delta s$$

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Perturbative QCD:

$$g_1(x, t) = \frac{1}{2} \sum_{k=1}^{n_f} \frac{e_k^2}{n_f} \int_x^1 \frac{dy}{y} \left[C_q^S \left(\frac{x}{y}, \alpha_s(t) \right) \boxed{\Delta\Sigma(y, t)} \right.$$

$$\left. + 2n_f C_g \left(\frac{x}{y}, \alpha_s(t) \right) \boxed{\Delta G(y, t)} \right]$$

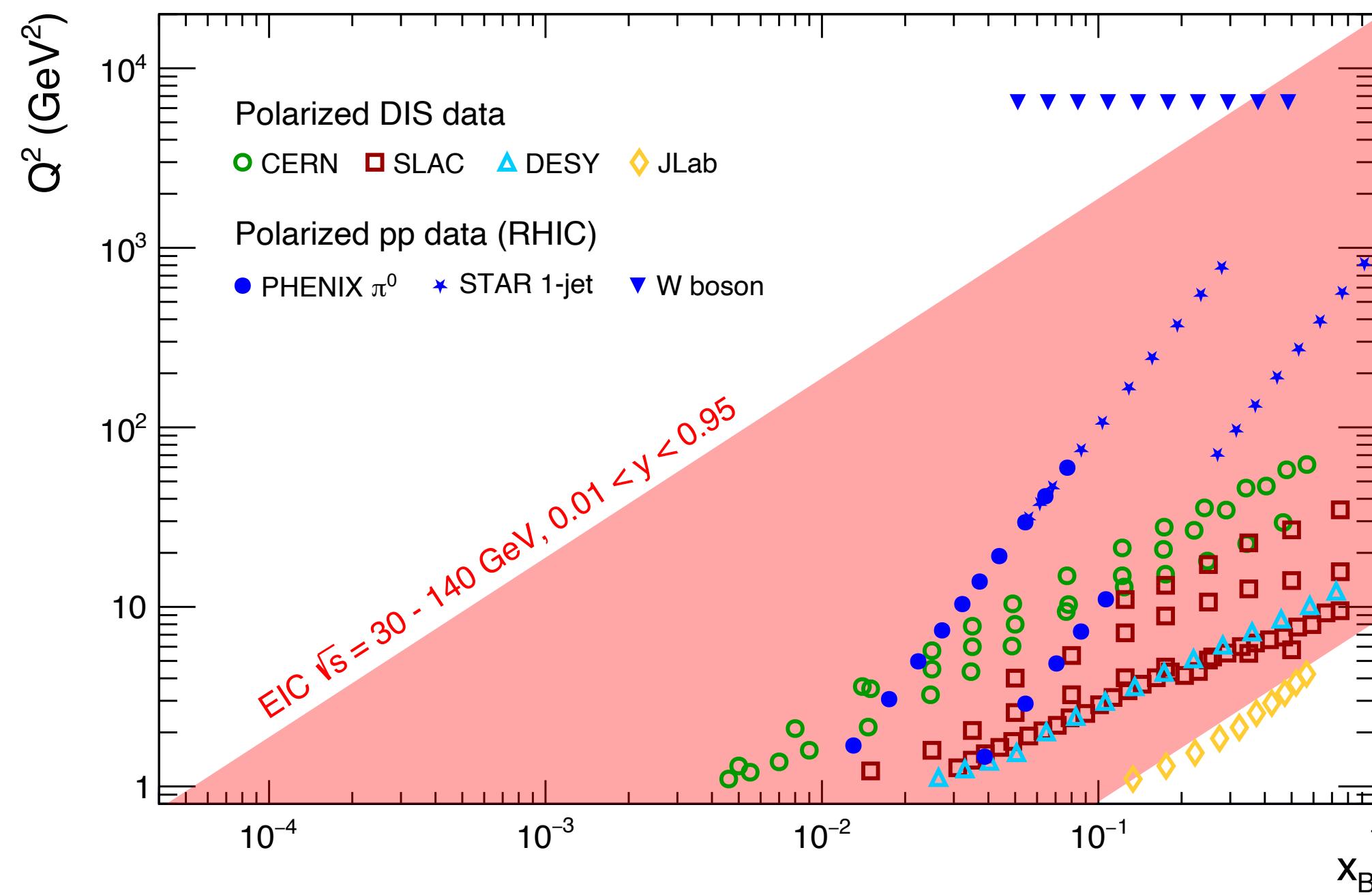
$$+ C_q^{\text{NS}} \left(\frac{x}{y}, \alpha_s(t) \right) \Delta q_{\text{NS}}(y, t) \Big]$$

Follow DGLAP evolution

g_1 constrains the contributions through Q^2 dependence

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Follow DGLAP evolution

g_1 constrains the contributions through Q^2 dependence
Need wide range data in x and Q^2 !

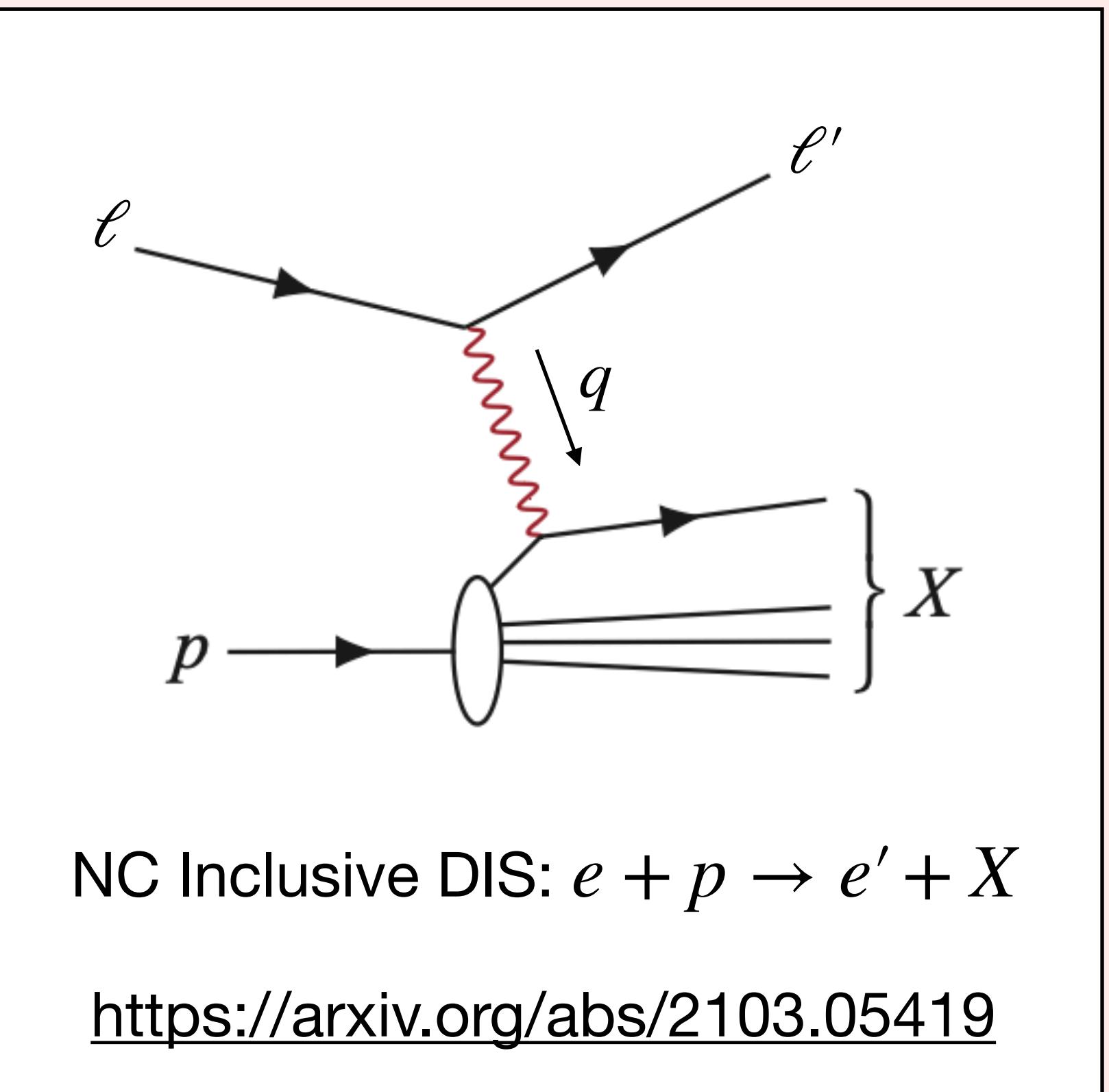
Probing Spin Structure with DIS

Deep Inelastic Scattering is a strong tool for probing nucleons and nuclei structure!

$$\frac{d^3\sigma(\beta)}{dQ^2dx d\phi} = \boxed{\frac{d^3\sigma_0}{dQ^2dx d\phi}} - \boxed{\frac{d^3\Delta\sigma(\beta)}{dQ^2dx d\phi}}$$

$$\frac{d^2\sigma}{dQ^2dx} = \frac{4\pi\alpha^2}{Q^4} \frac{1}{x} \left[xy^2 F_1(x, Q^2) + \left(1 - y - \frac{Mxy}{2E} F_2(x, Q^2) \right) \right]$$

$$\begin{aligned} \frac{d^3\Delta\sigma(\beta)}{dQ^2dx d\phi} &= \frac{4\alpha^2}{Q^2} y \left\{ \cos\beta \left[\left(1 - \frac{y}{2} - \frac{\gamma^2 y^2}{4} \right) g_1(x, Q^2) - \frac{\gamma^2 y}{4} g_2(x, Q^2) \right] \right. \\ &\quad \left. - \cos\phi \sin\beta \frac{\sqrt{Q^2}}{\nu} \left(1 - y - \frac{\gamma^2 y^2}{4} \right)^{\frac{1}{2}} \left[\frac{y}{2} g_1(x, Q^2) + g_2(x, Q^2) \right] \right\} \end{aligned}$$



Double Spin Asymmetry

$$\frac{d^3 \Delta\sigma(\beta)}{dQ^2 dx d\phi} = \frac{4\alpha^2}{Q^2} y \left\{ \cos\beta \left[\left(1 - \frac{y}{2} - \frac{\gamma^2 y^2}{4} \right) g_1(x, Q^2) - \frac{\gamma^2 y}{4} g_2(x, Q^2) \right] \right.$$

$$\left. - \cos\phi \sin\beta \frac{\sqrt{Q^2}}{\nu} \left(1 - y - \frac{\gamma^2 y^2}{4} \right)^{\frac{1}{2}} \left[\frac{y}{2} g_1(x, Q^2) + g_2(x, Q^2) \right] \right\}$$

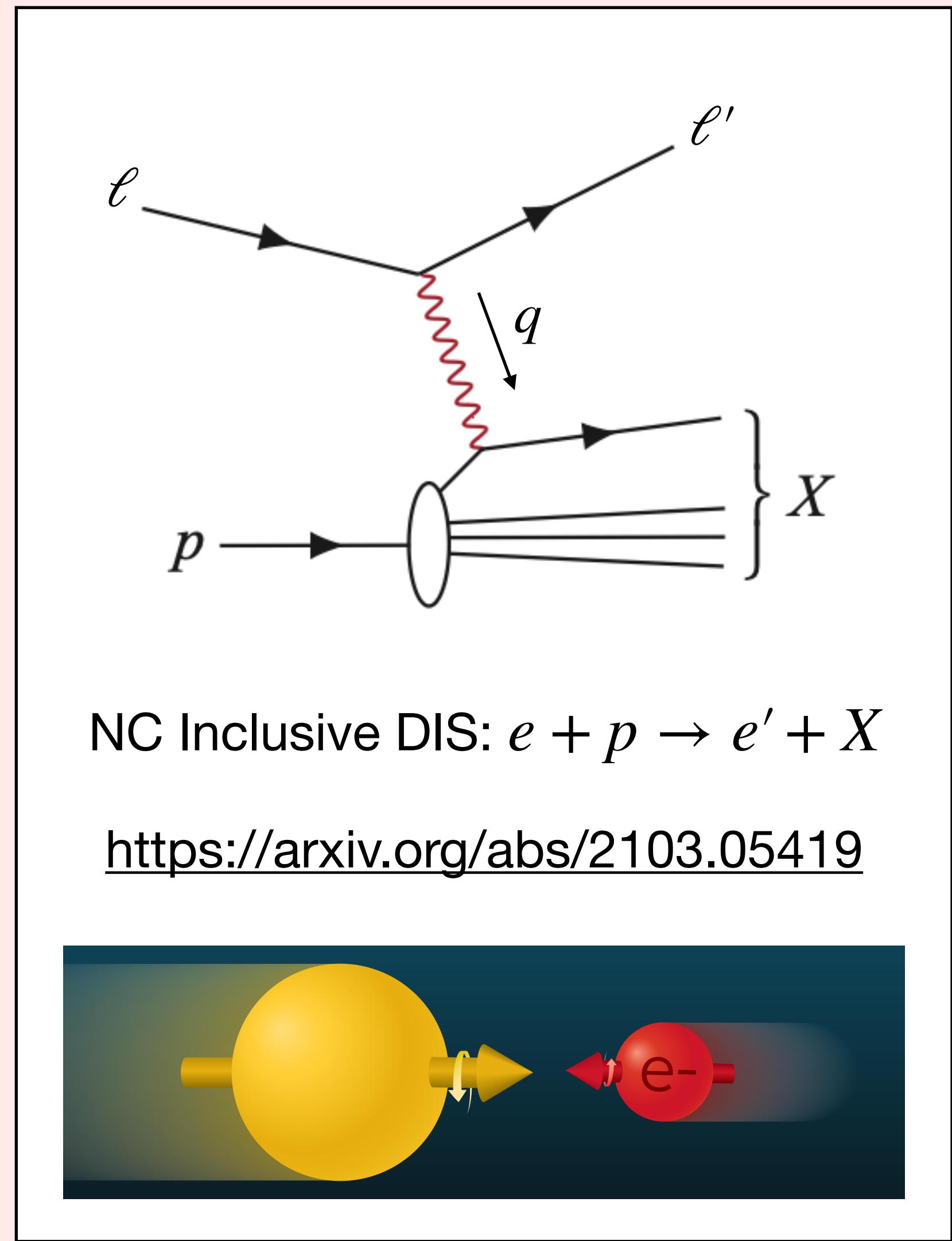
- Direct measurement is more challenging
- Easier to measure via cross section ratio
- Systematic is largely canceled out

$$g_1 = \frac{F_2}{2x(1+R)} (A_1 + \gamma A_2)$$

$$A_1(x, Q^2) \equiv \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{A_{||}}{D(1 + \eta\xi)} - \frac{\eta A_{\perp}}{d(1 + \eta\xi)}$$

$$A_{||} = \frac{\sigma_{\downarrow\uparrow} - \sigma_{\uparrow\uparrow}}{\sigma_{\downarrow\uparrow} + \sigma_{\uparrow\uparrow}}$$

$$A_{\perp} = \frac{\sigma_{\downarrow\Rightarrow} - \sigma_{\uparrow\Rightarrow}}{\sigma_{\downarrow\Rightarrow} + \sigma_{\uparrow\Rightarrow}}$$



Double Spin Asymmetry

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$$A_1(x, Q^2) \equiv \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{A_{\parallel}}{D(1 + \eta\xi)} - \frac{\eta A_{\perp}}{d(1 + \eta\xi)}$$

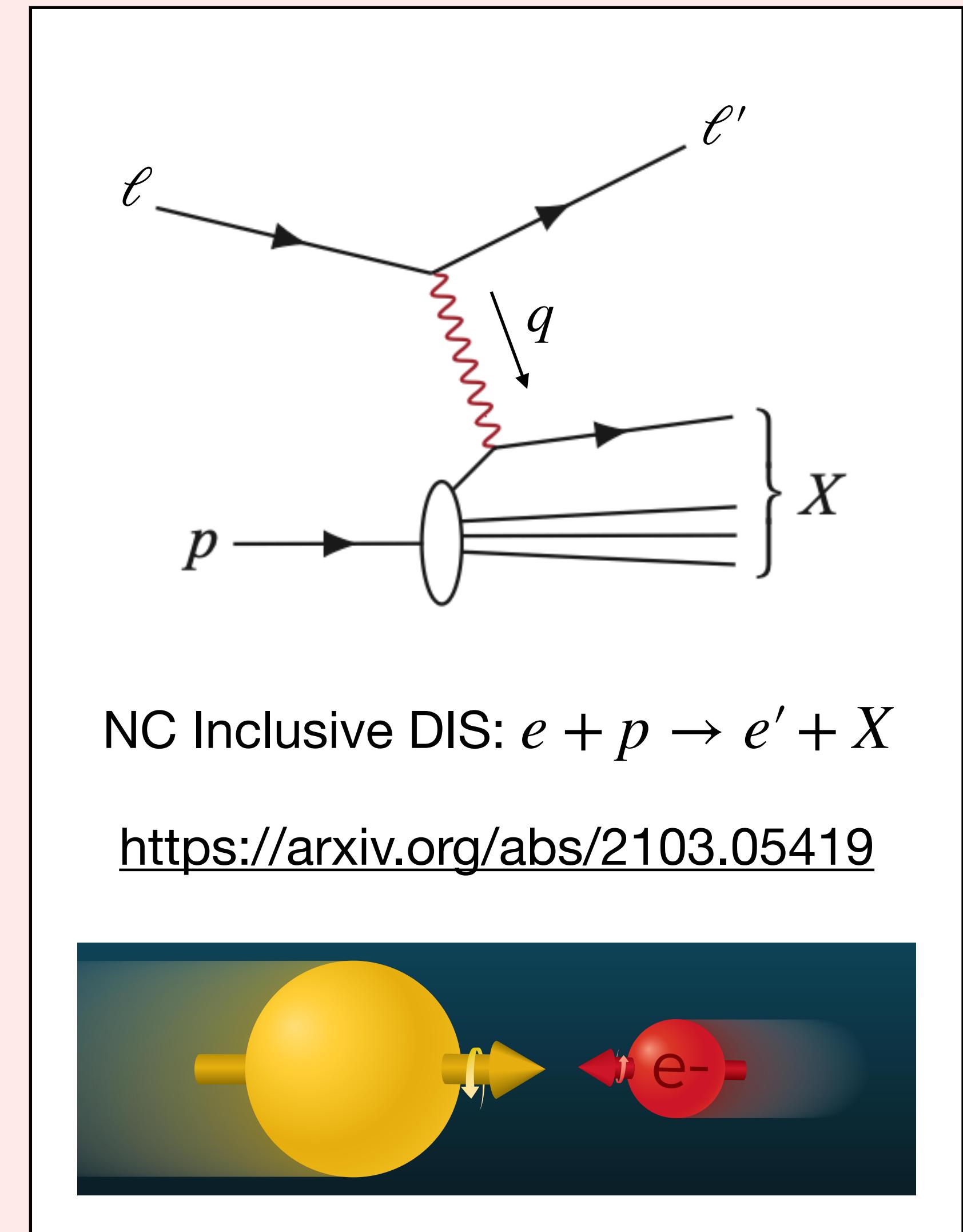
p & γ^* spins
anti-aligned p & γ^* spins
aligned

$$A_{\parallel} = \frac{\sigma_{\downarrow\uparrow} - \sigma_{\uparrow\uparrow}}{\sigma_{\downarrow\uparrow} + \sigma_{\uparrow\uparrow}}$$

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$$A_{\perp} = \frac{\sigma_{\downarrow\Rightarrow} - \sigma_{\uparrow\Rightarrow}}{\sigma_{\downarrow\Rightarrow} + \sigma_{\uparrow\Rightarrow}}$$

p & e spins
perpendicular



Double Spin Asymmetry

$$A_1(x, Q^2) \equiv \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{A_{\parallel}}{D(1 + \eta\xi)} - \frac{\eta A_{\perp}}{d(1 + \eta\xi)}$$

p & γ^* spins
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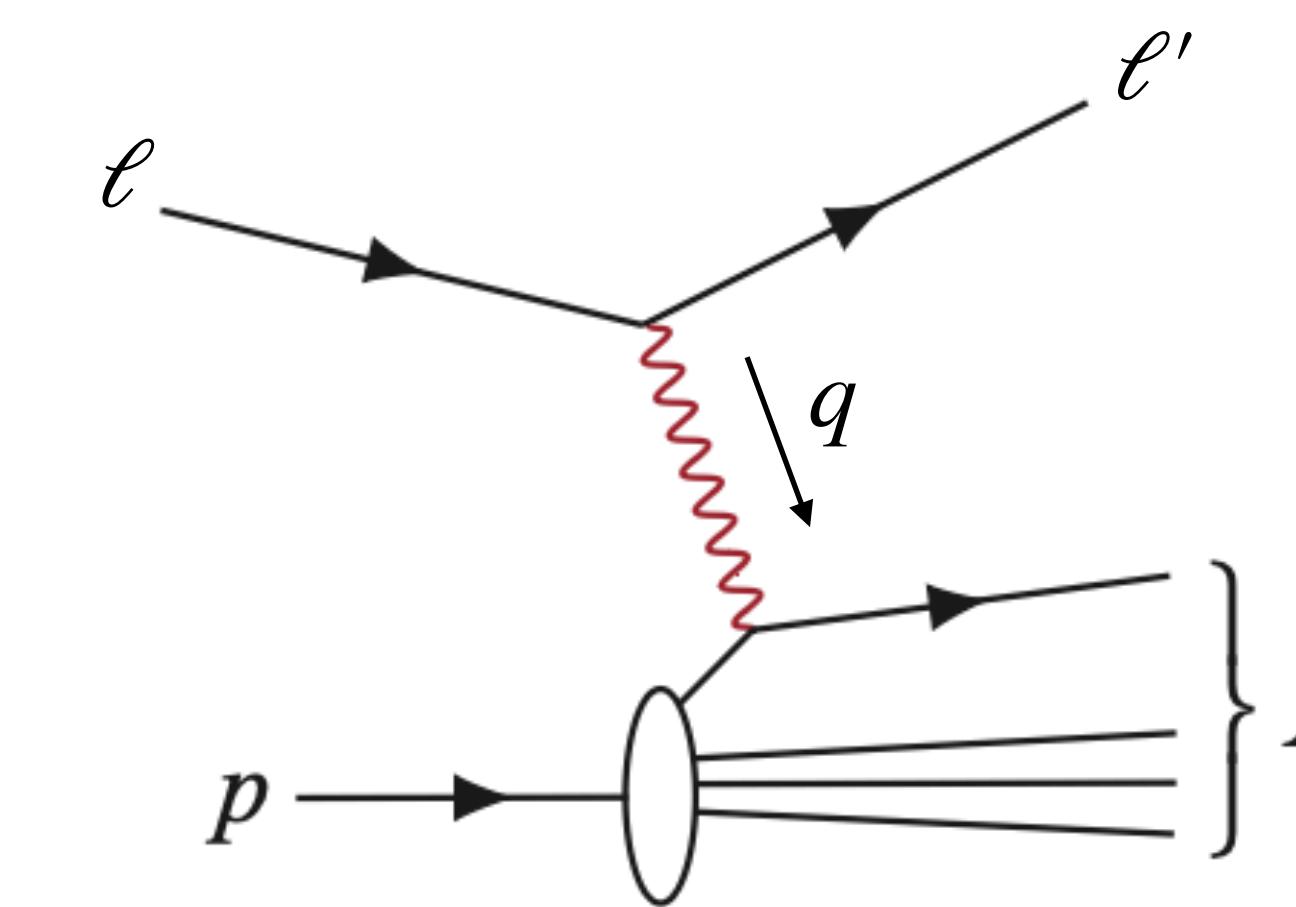
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p & e spins
anti-aligned

p & e spins
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$$A_{\perp} = \frac{\sigma_{\downarrow\Rightarrow} - \sigma_{\uparrow\Rightarrow}}{\sigma_{\downarrow\Rightarrow} + \sigma_{\uparrow\Rightarrow}}$$

p & e spins
perpendicular



NC Inclusive DIS: $e + p \rightarrow e' + X$
<https://arxiv.org/abs/2103.05419>

$$D = \frac{y(2-y)(2+\gamma^2 y)}{2(1+\gamma^2)y^2 + (4(1-y)-\gamma^2 y^2)(1+R)}$$

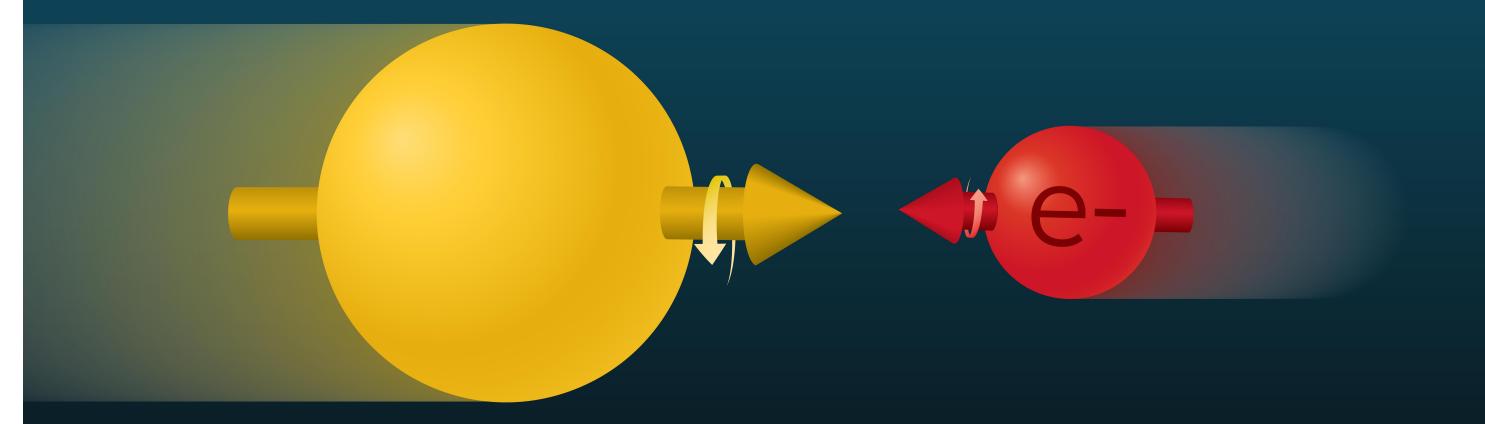
$$d = \frac{D\sqrt{4(1-y)-\gamma^2 y^2}}{2-y}$$

$$R \equiv \frac{\sigma_L}{\sigma_T} = \frac{F_L}{F_2 - F_L}$$

$$\gamma^2 = \frac{4M^2 x^2}{Q^2}$$

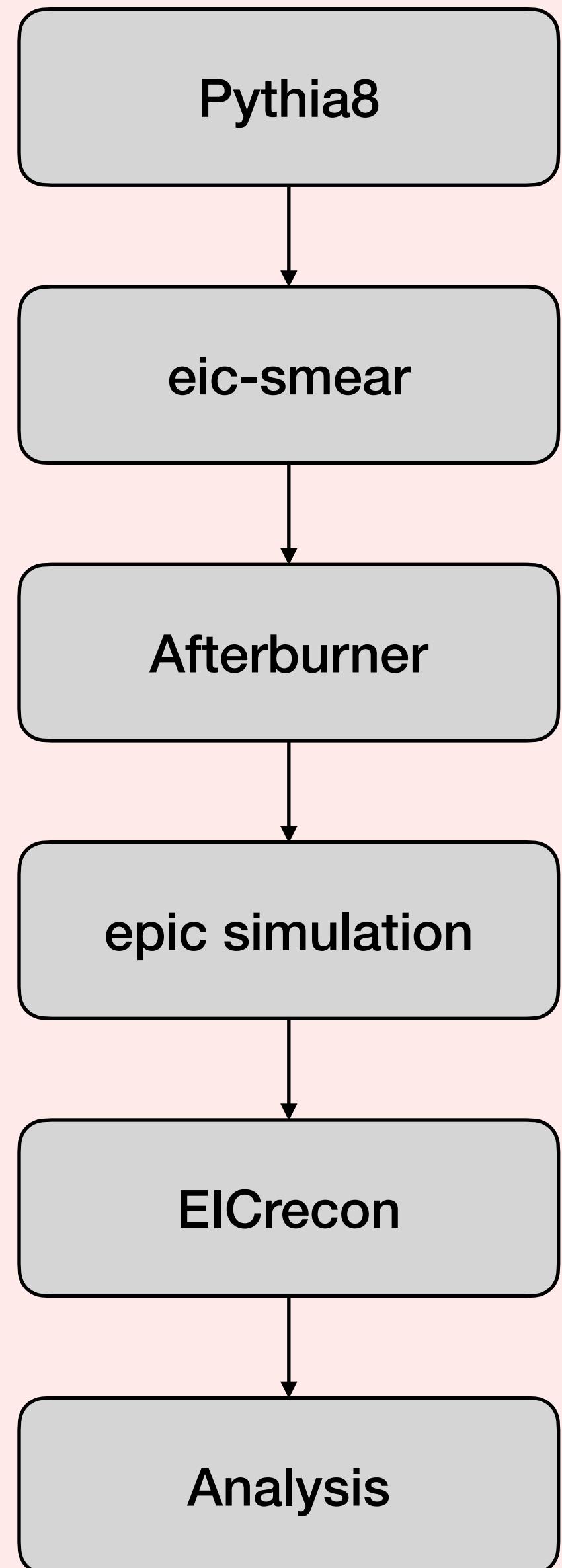
$$\eta = \frac{4(1-y)-\gamma^2 y^2}{(2-y)(2+\gamma^2 y)}$$

$$\xi = \frac{\gamma(2-y)}{2+\gamma^2 y}$$

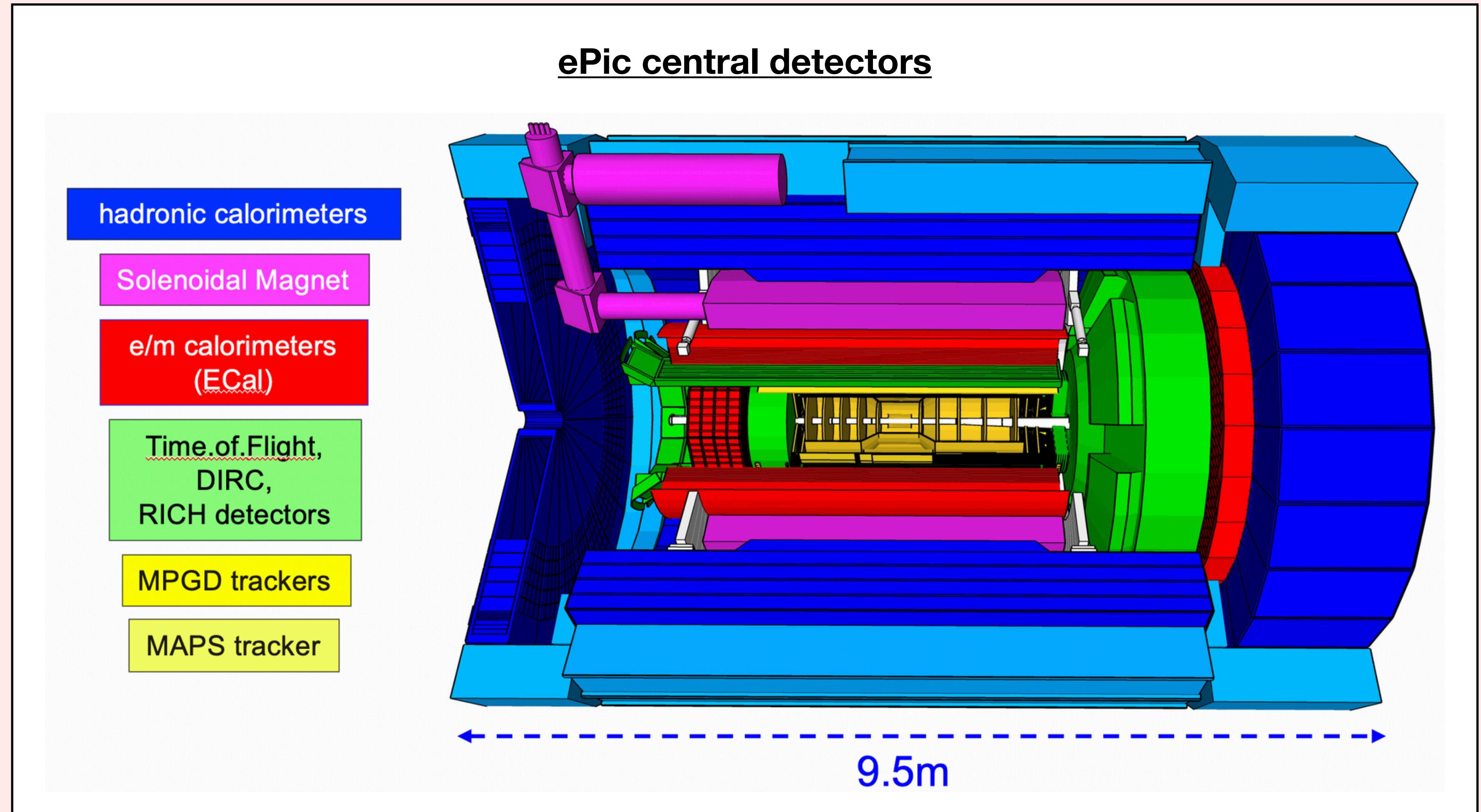


Projection study analysis chain for g_1^p

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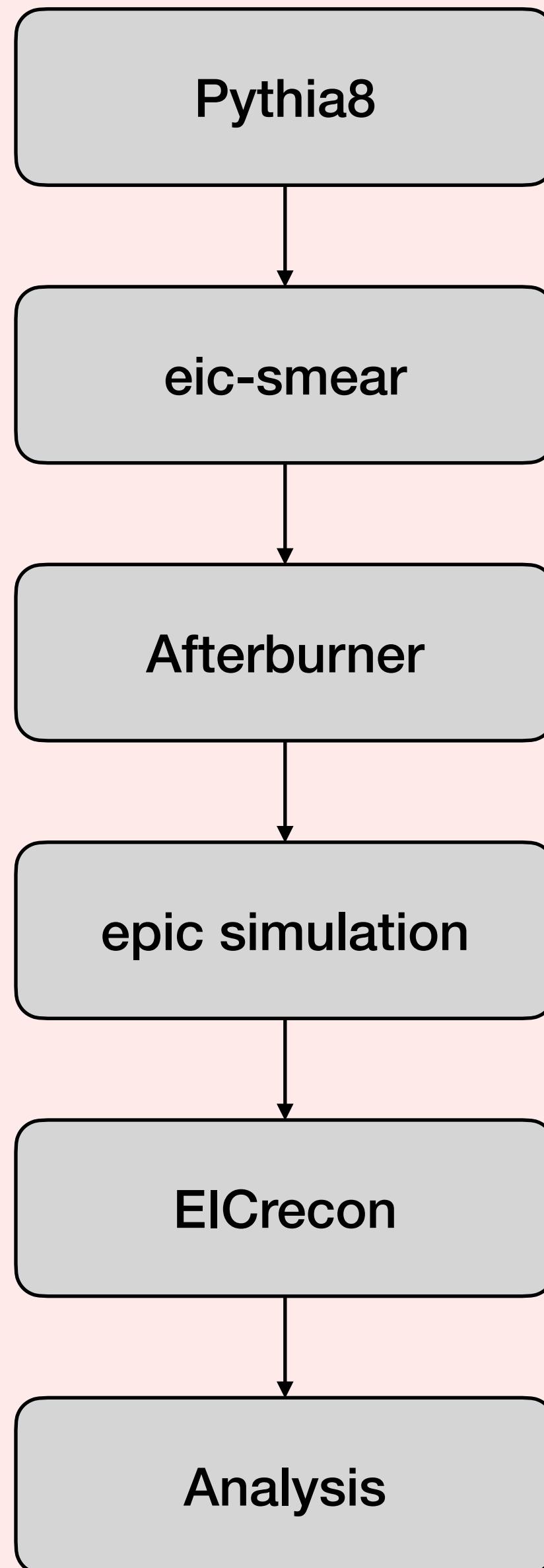


- ▶ Simulation with full ePIC detector setup
- ▶ Beam crossing angle and smearing included

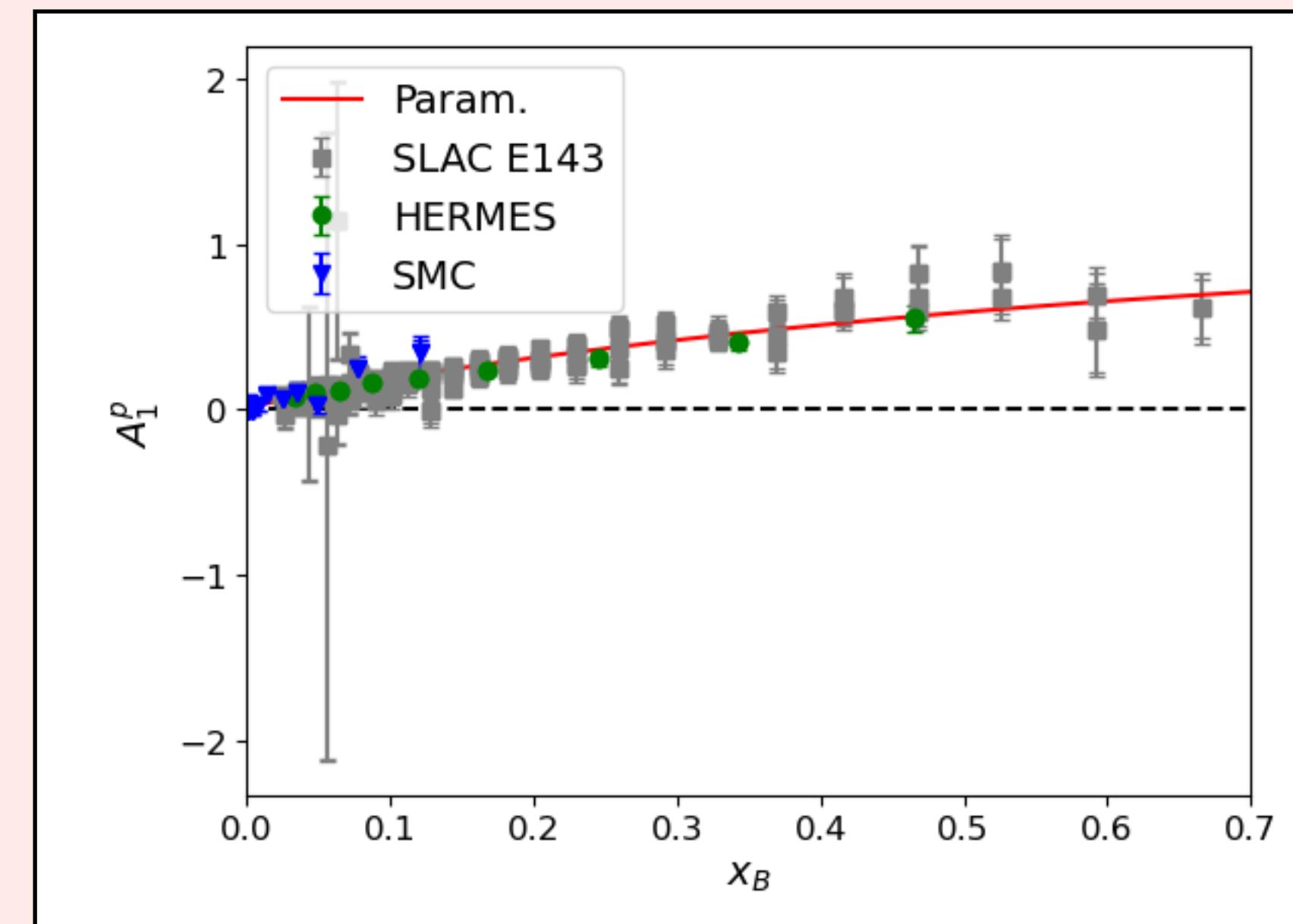


Projection study analysis chain for g_1^p

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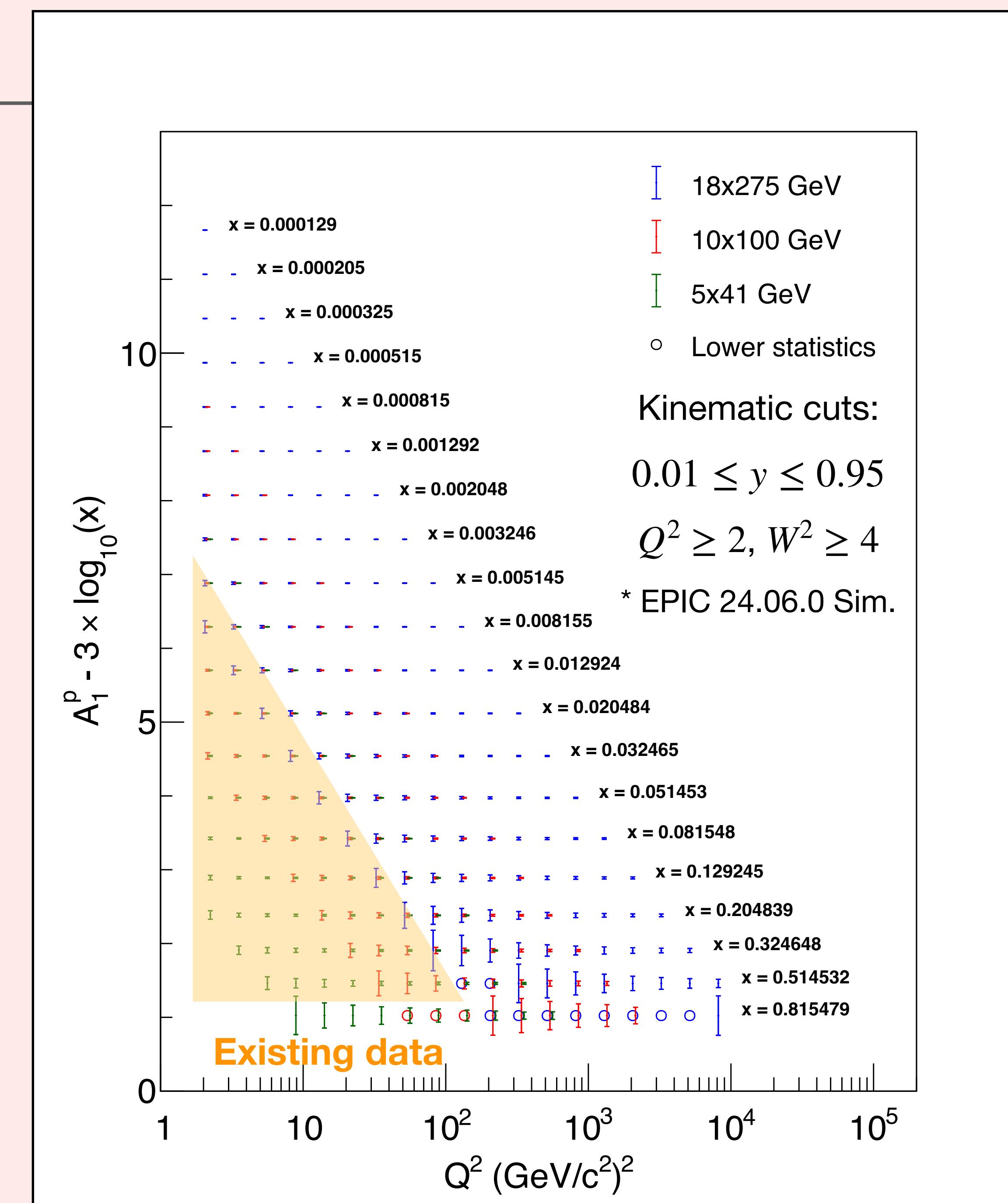
- ▶ Simulation with full ePIC detector setup
- ▶ Beam crossing angle and smearing included
- ▶ Generated events are non-polarized
- ▶ Central A_1^p is calculated from JLab global fit ([Doi: 10.1103/PhysRevC.70.065207](https://doi.org/10.1103/PhysRevC.70.065207))
- ▶ Statistical uncertainty is estimated by $\delta A_{\parallel,\perp} \approx \frac{1}{\sqrt{N} P_e P_N}$



- Parameterization at $Q^2 = 2.88 \text{ GeV}^2$
- Data points are at various Q^2 with majority $< 5 \text{ GeV}^2$

Projected A_1^P at EIC

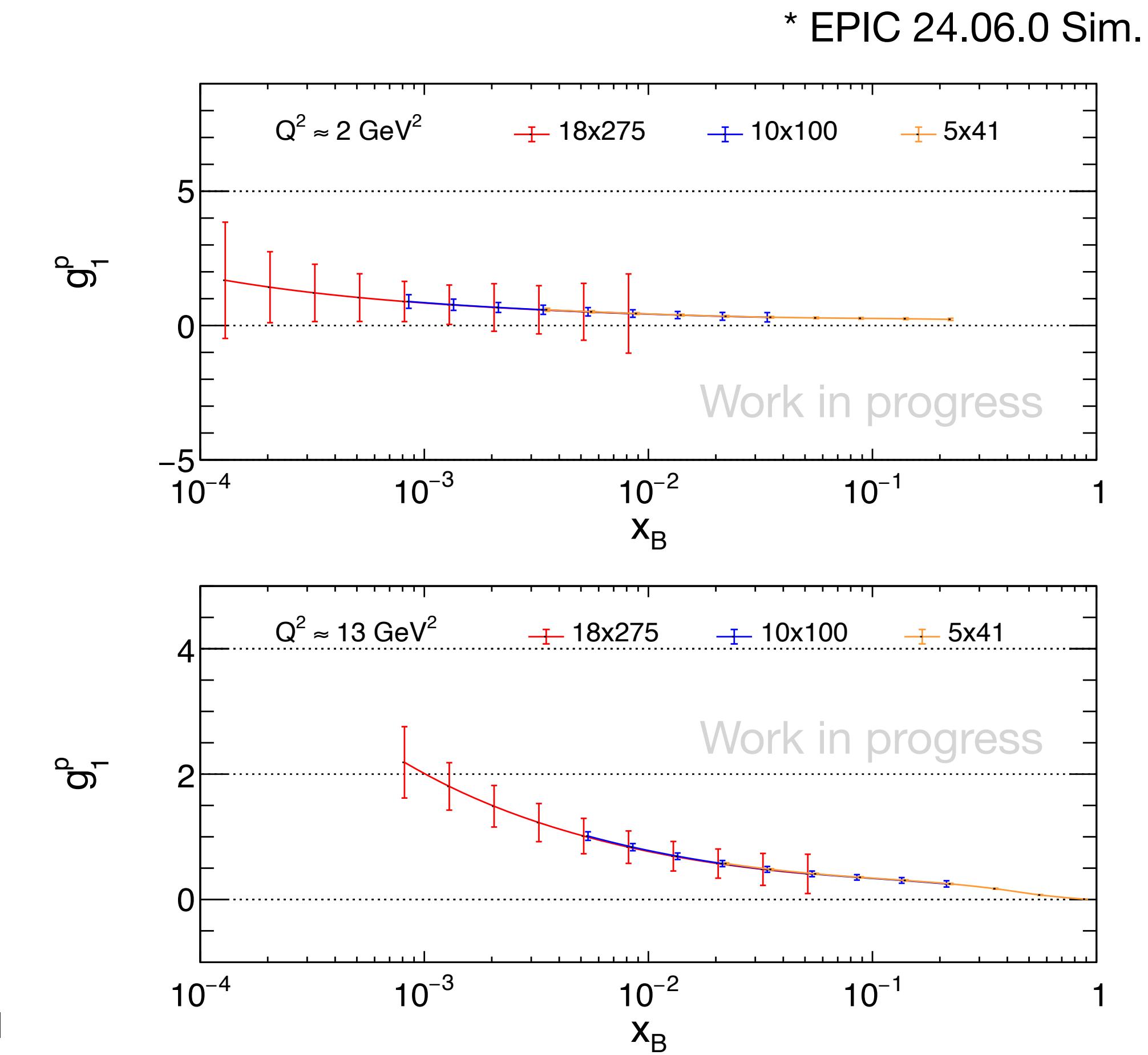
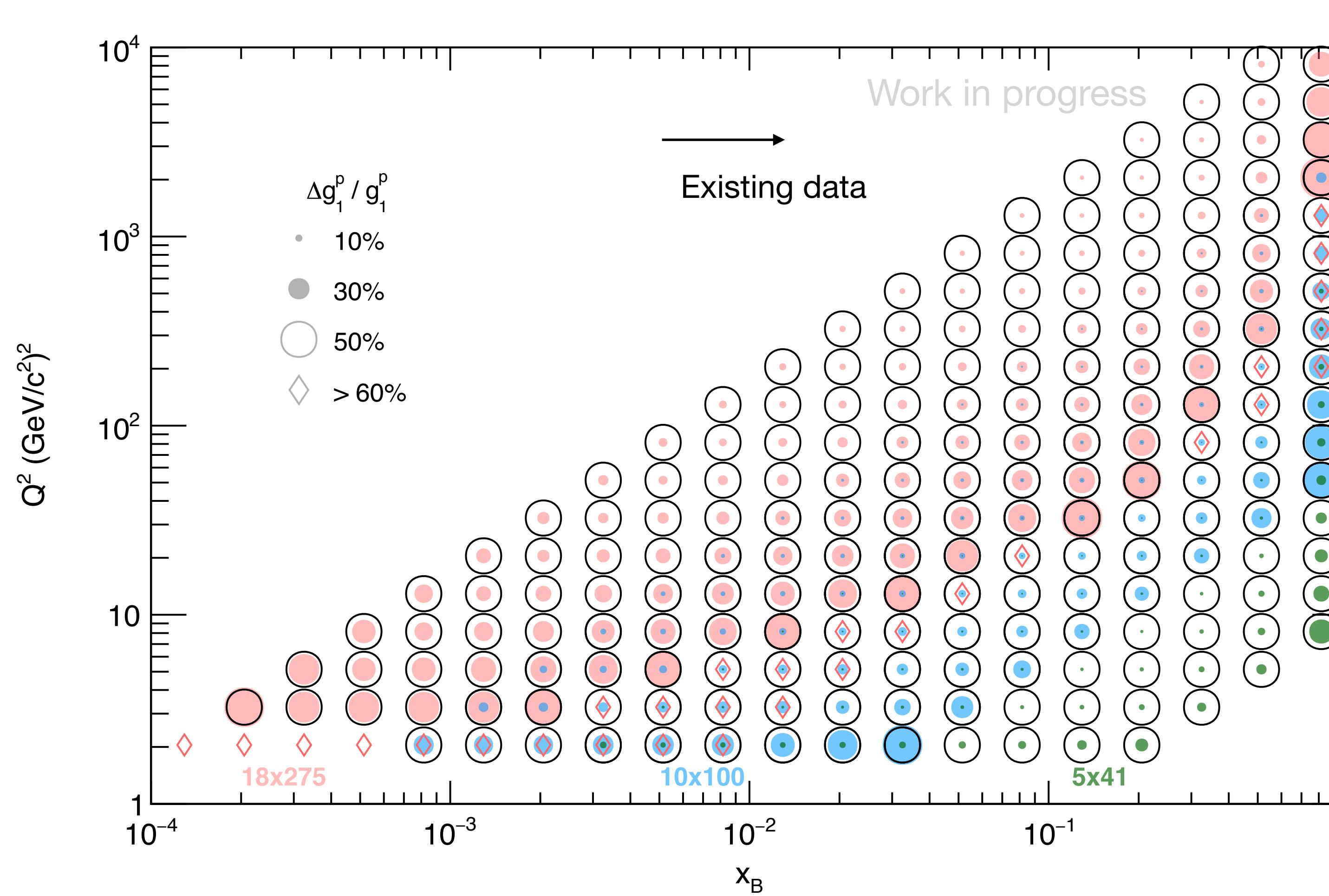
- $A_1(x, Q^2) \equiv \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{A_{||}}{D(1 + \eta\xi)} - \frac{\eta A_{\perp}}{d(1 + \eta\xi)}$
- $\delta A_{||, \perp} = \frac{1}{\sqrt{N} P_e P_N}$
- $\mathcal{L} = 10 \text{ fb}^{-1}, P_e = P_p = 70 \%$
- Data split evenly between $A_{||}$ and A_{\perp}
- R calculated from [https://doi.org/10.1016/S0370-2693\(99\)00244-0](https://doi.org/10.1016/S0370-2693(99)00244-0)
- Statistical uncertainty only, correction not yet applied



Projected g_1^p at EIC

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- $A_1 \approx g_1/F_1$ with F_1 calculated from [JAM22](#)
- Statistical uncertainties only

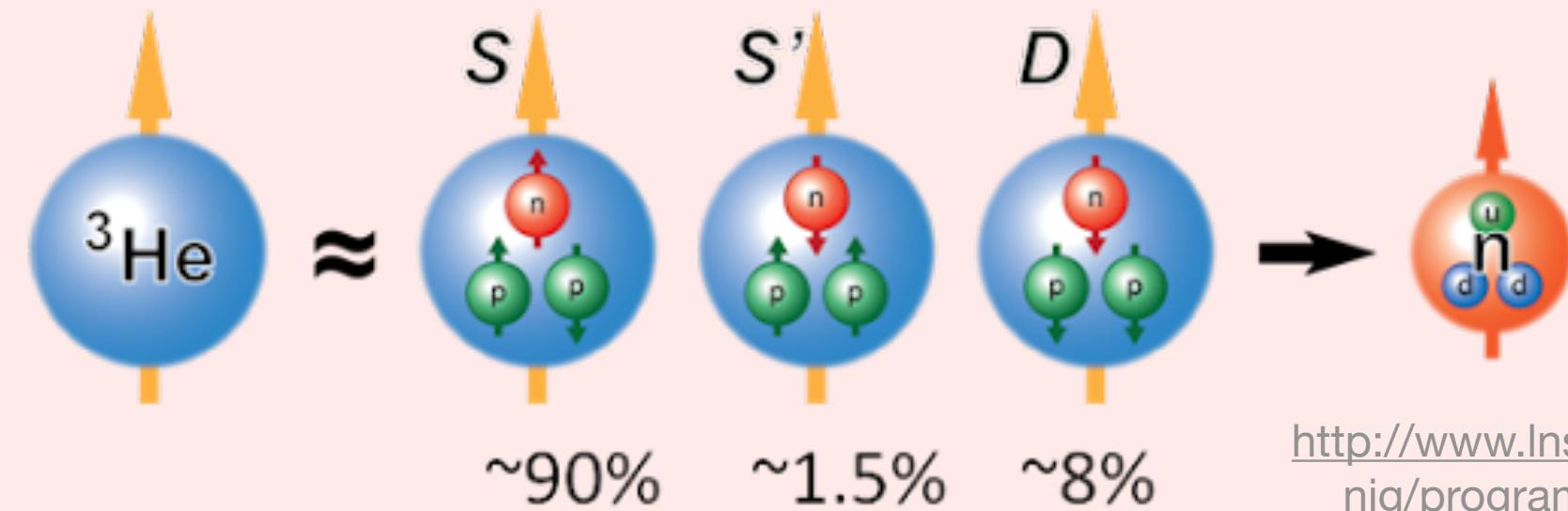


A_1^n from $e^3\text{He}$ DIS:

What about A_1^n ?

- Can be extracted from $A_1^{^3\text{He}} = P_n \frac{F_2^n}{F_{^3\text{He}}} A_1^n + 2P_p \frac{F_2^p}{F_{^3\text{He}}} A_1^p$

$$P_n = 0.86 \pm 0.02 \quad P_p = -0.028 \pm 0.004$$



<http://www.lns.mit.edu/nig/programs.html>

<https://doi.org/10.1103/PhysRevC.64.024004>

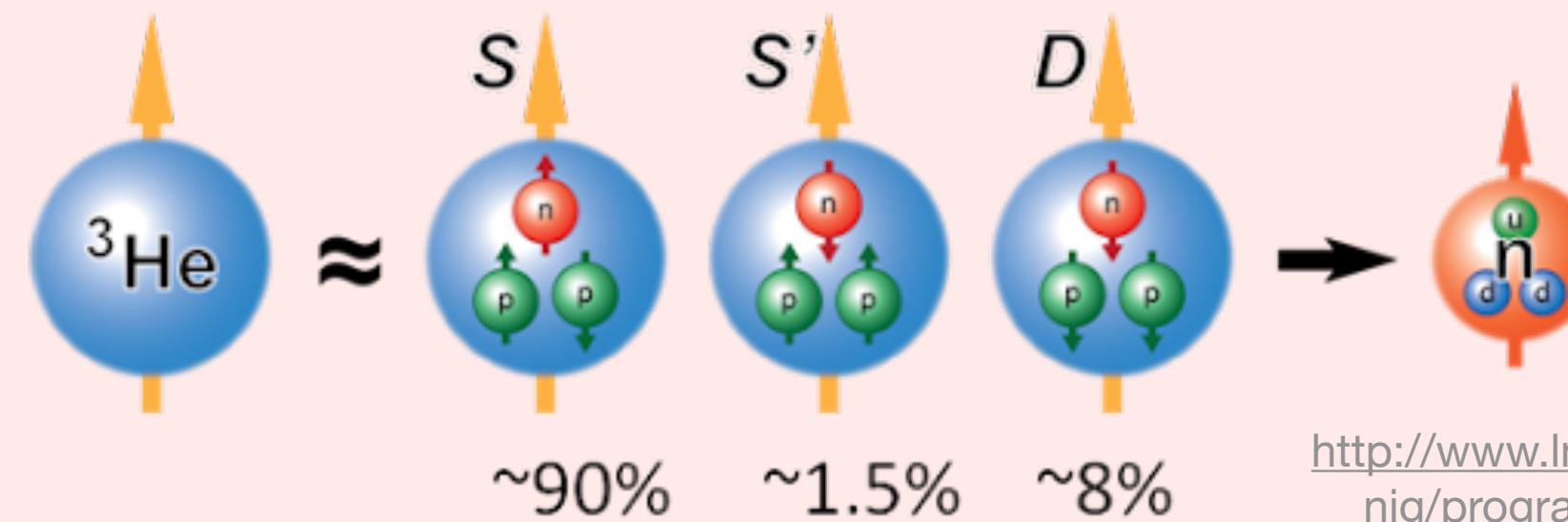
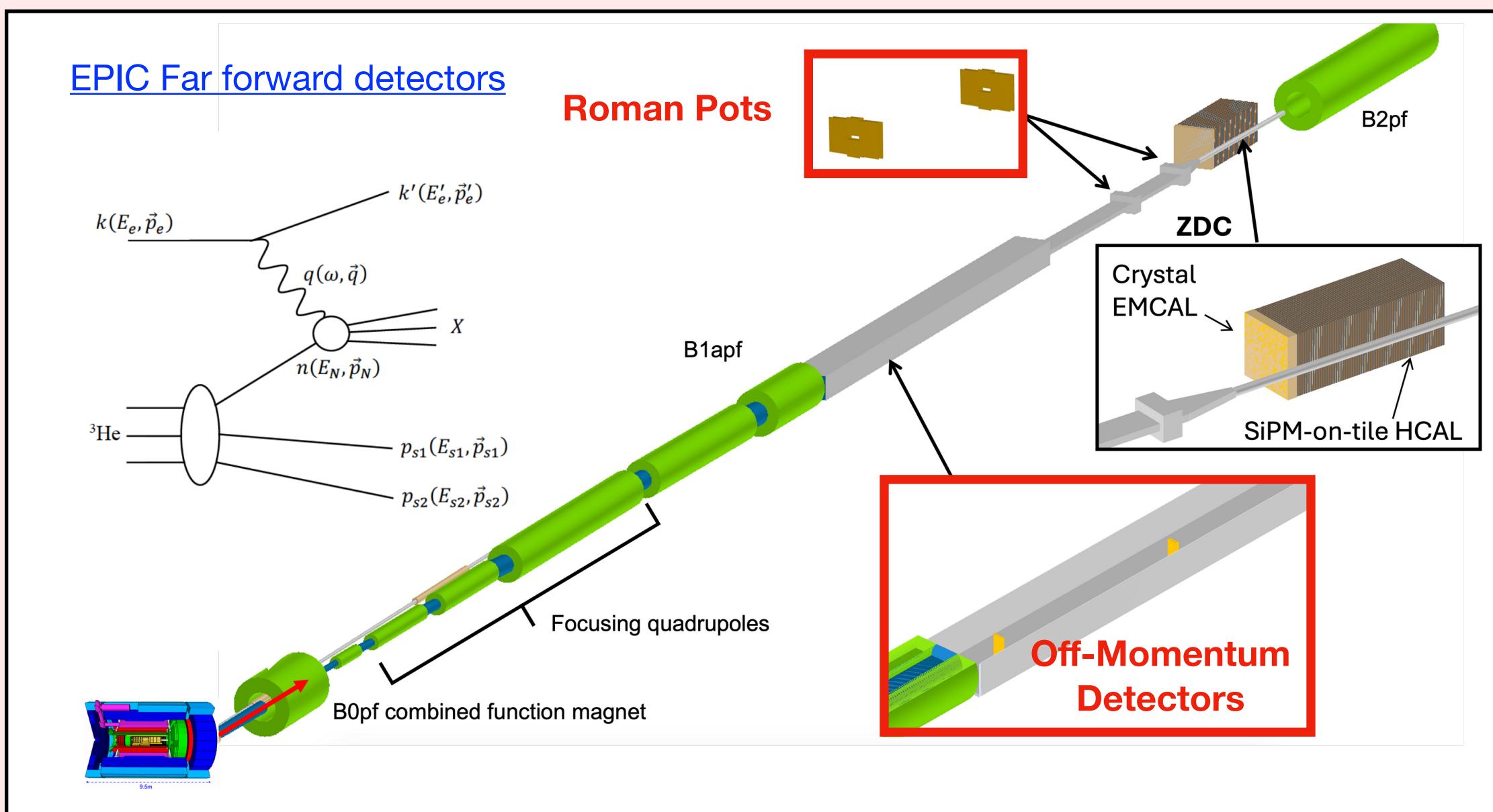
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- Or **directly** measured double spectator tagging:



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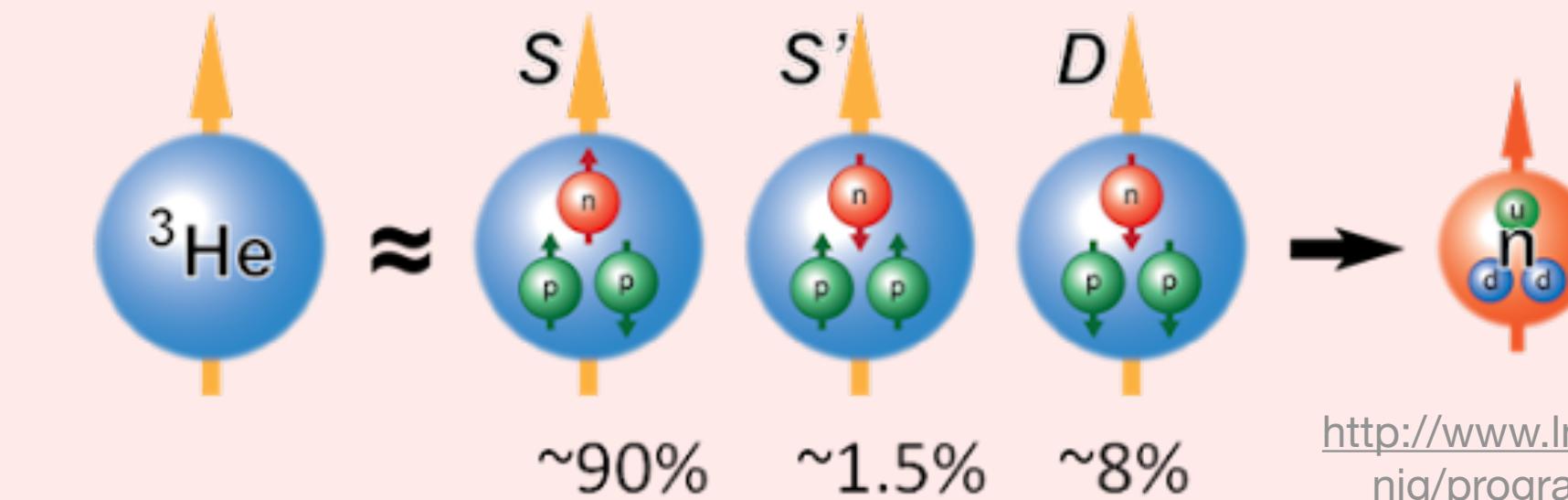
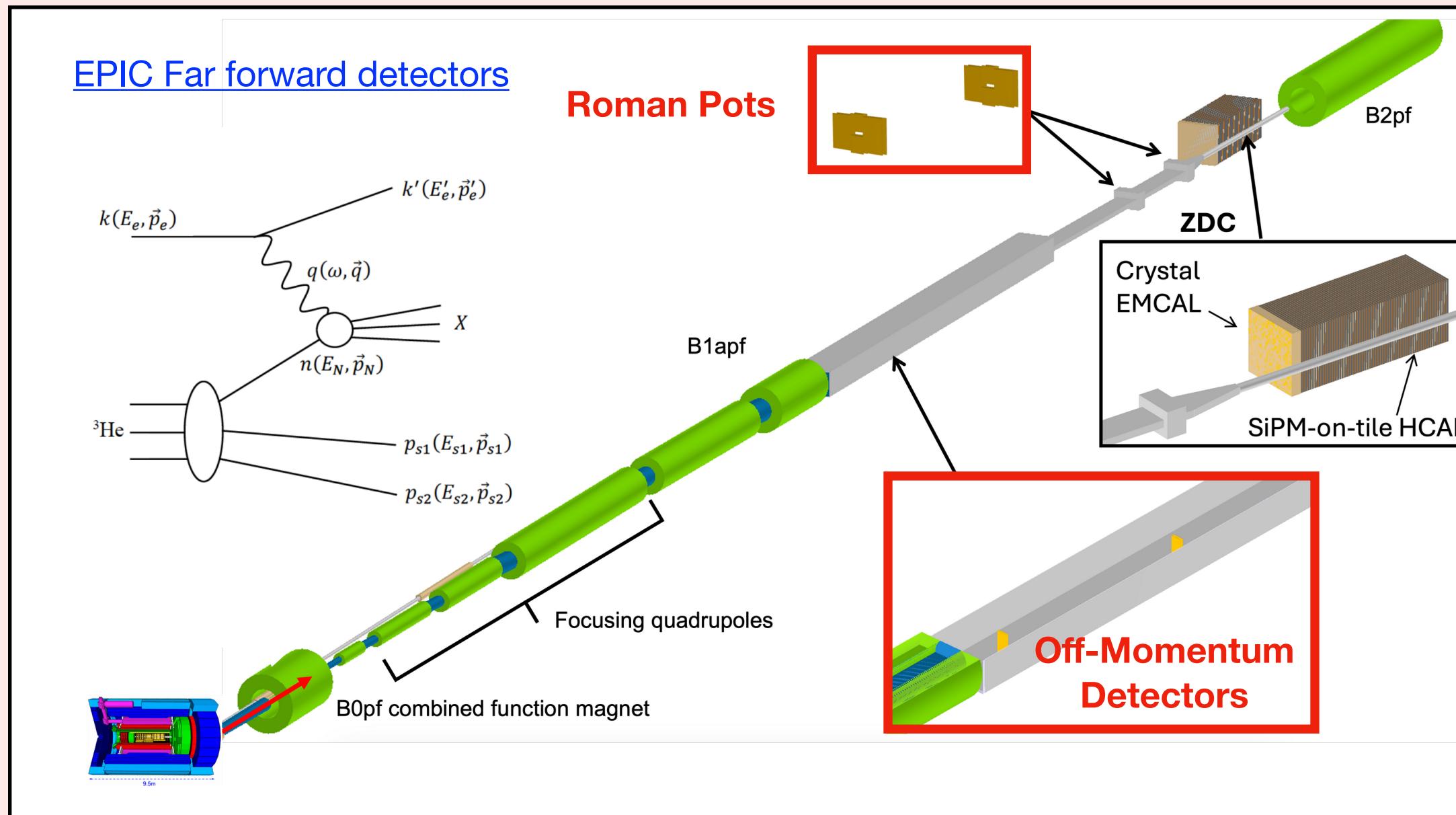
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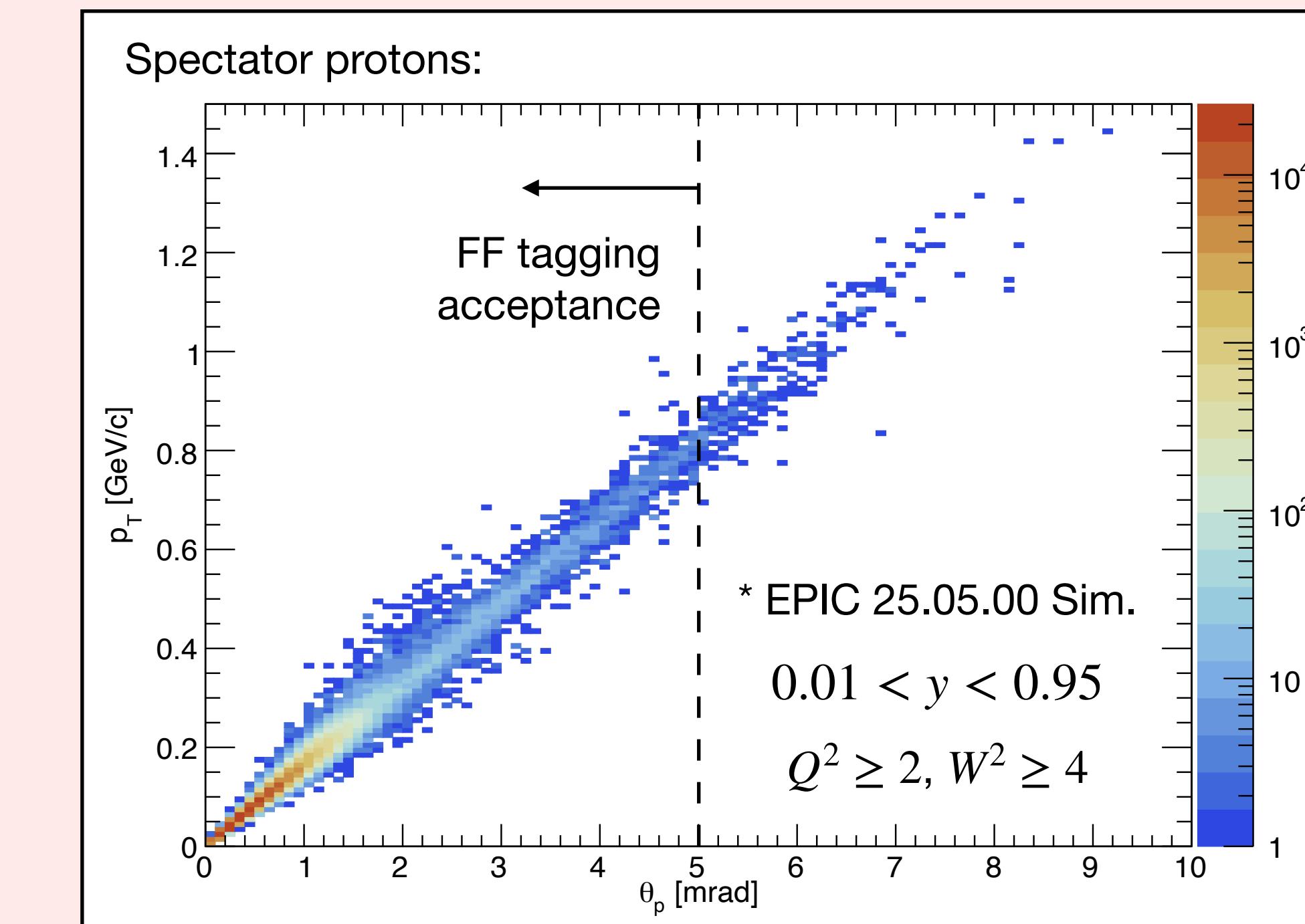
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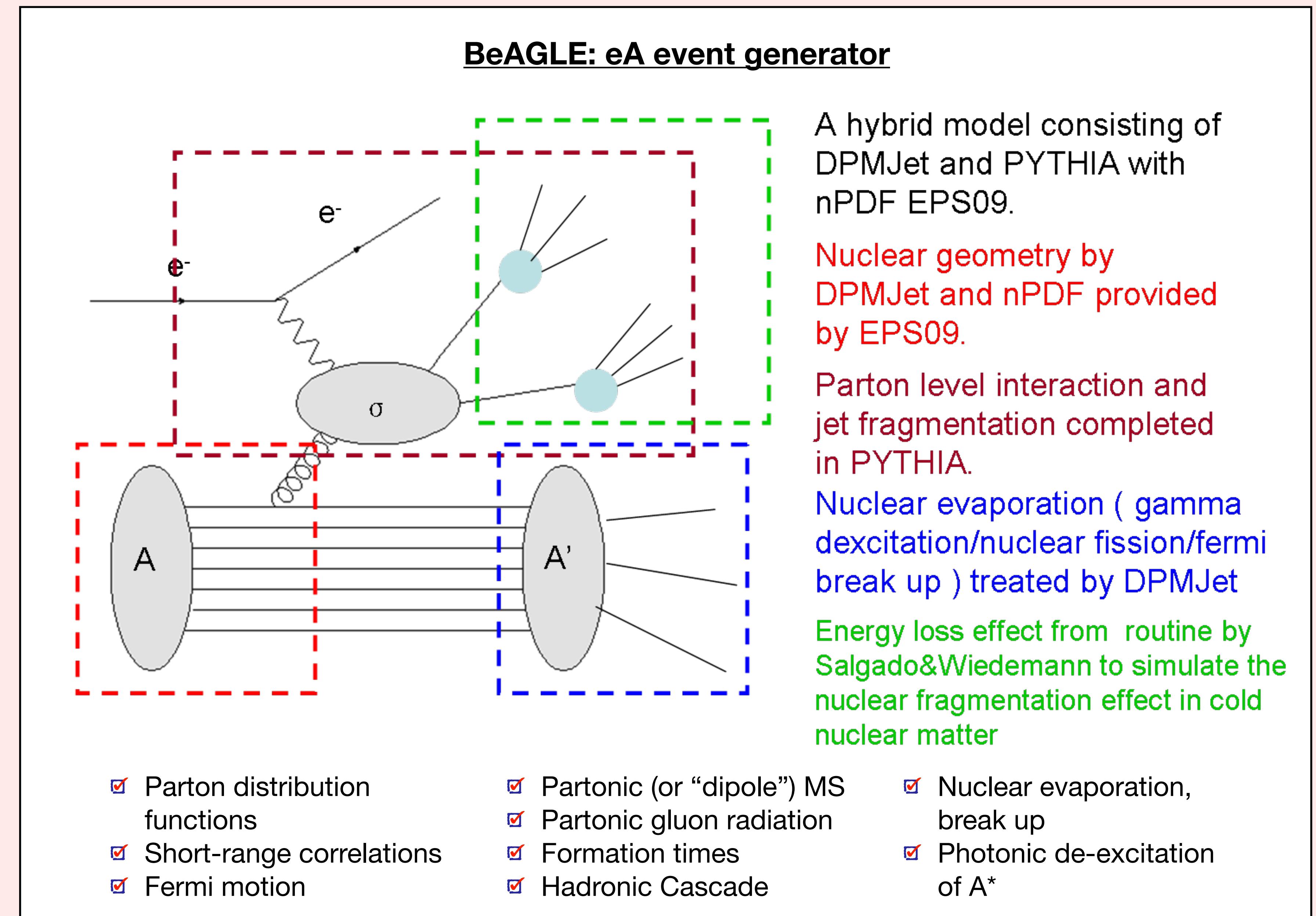
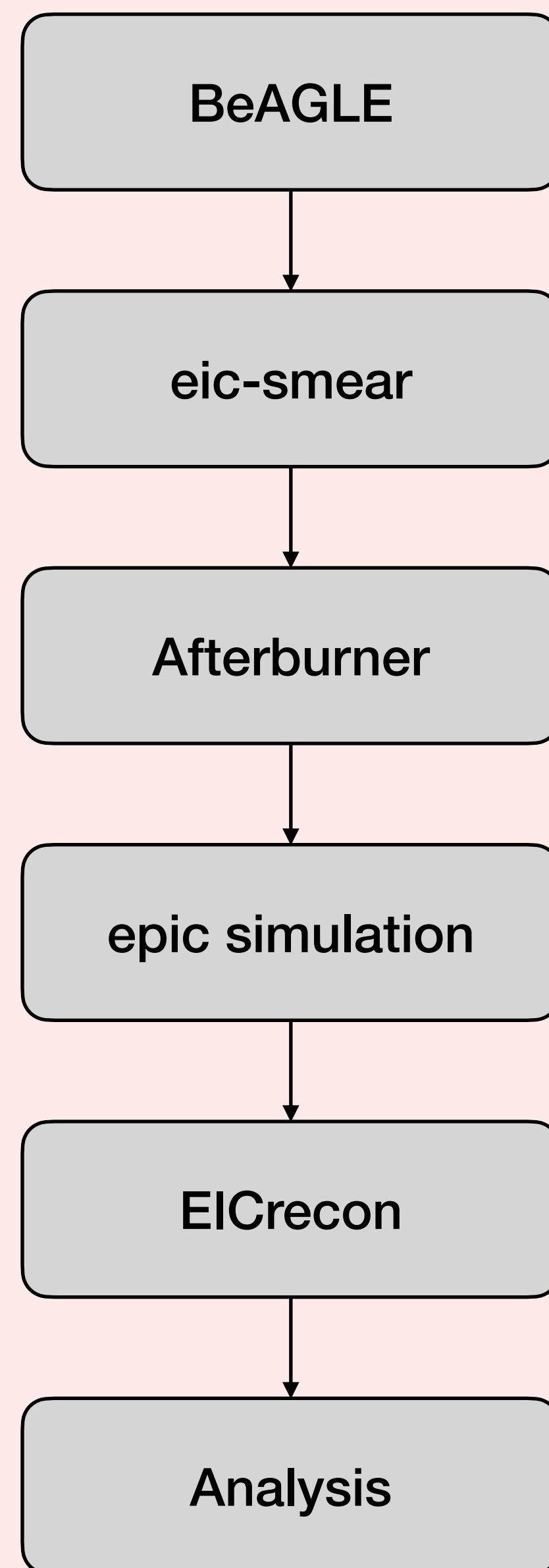


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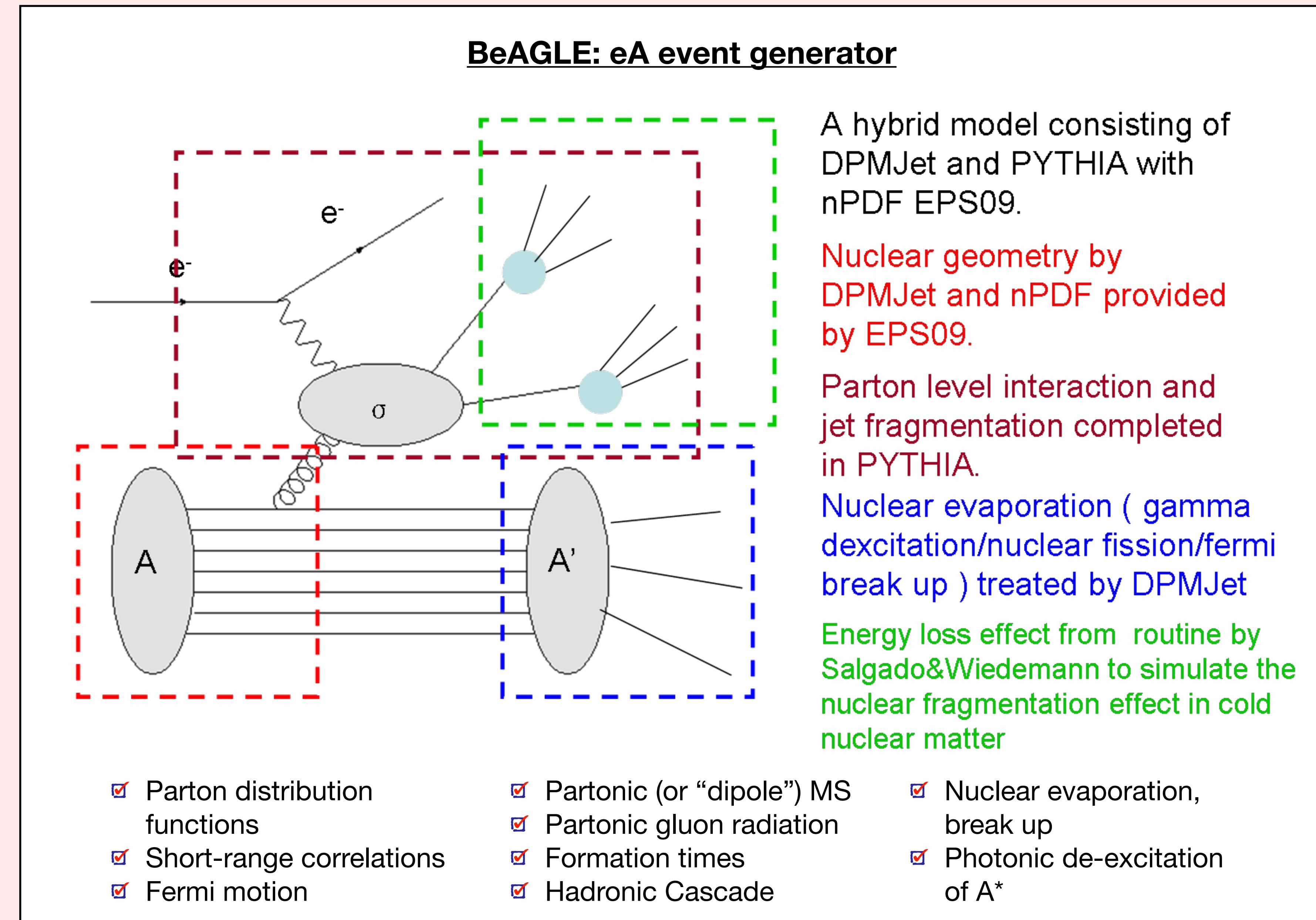
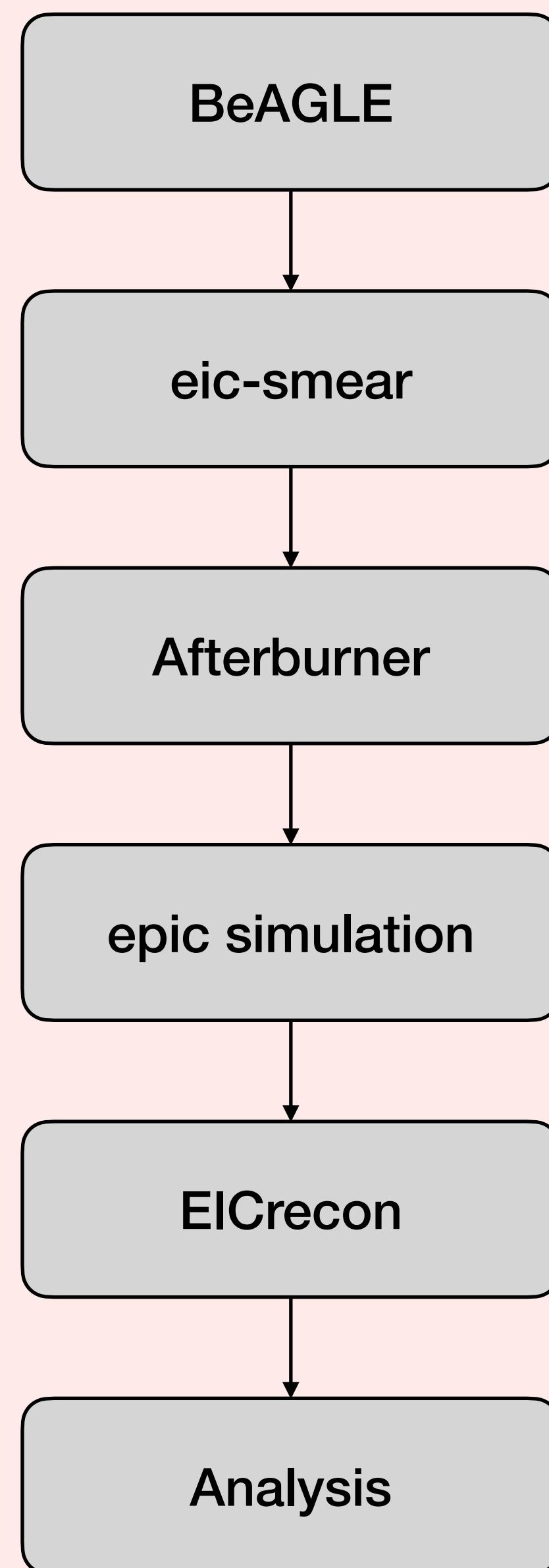
Projection study analysis chain for g_1^n

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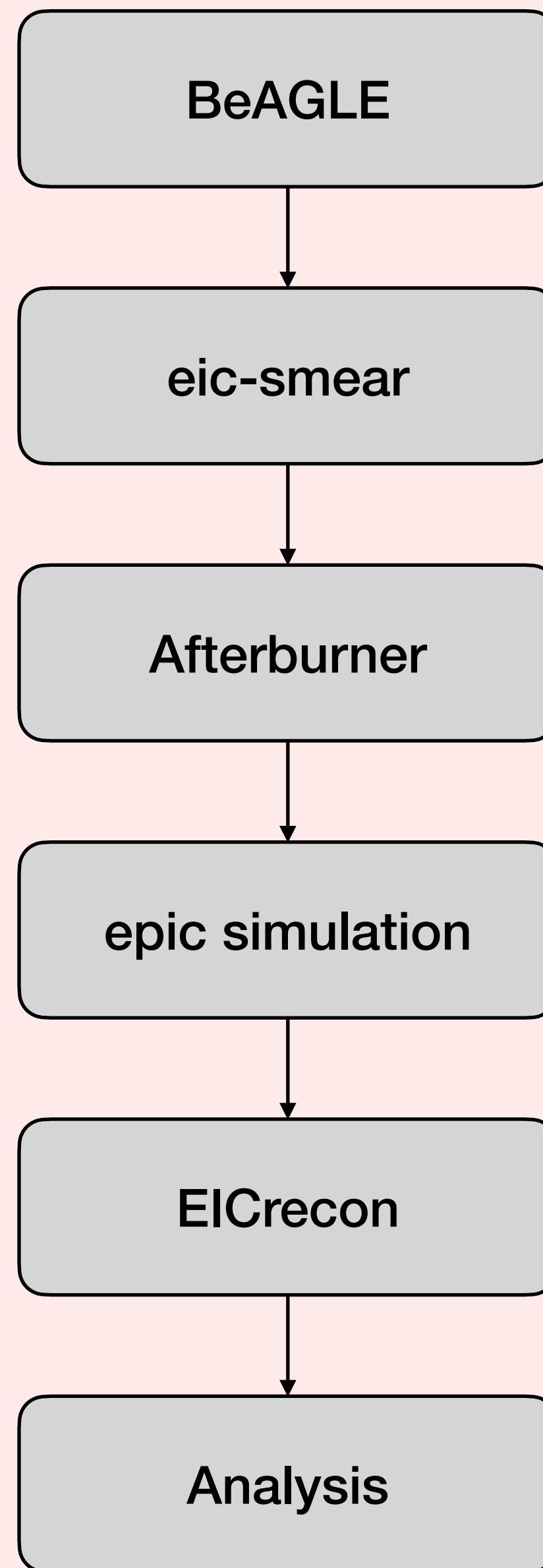
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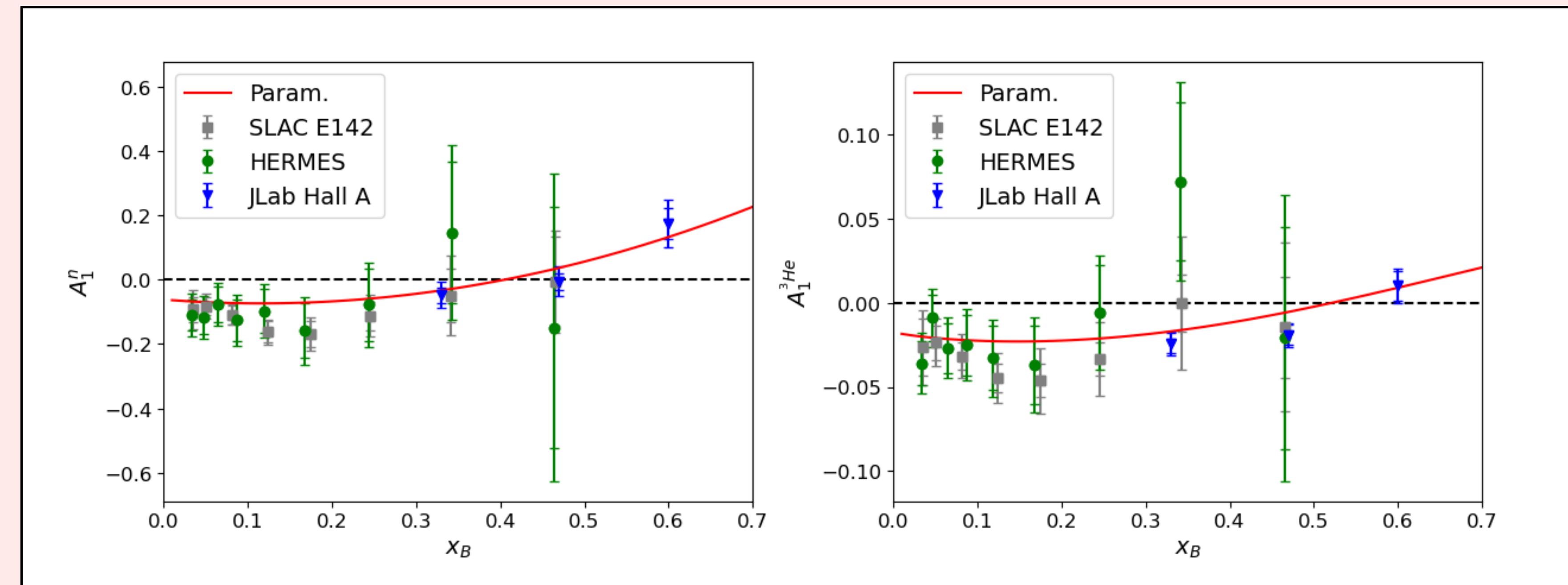


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18



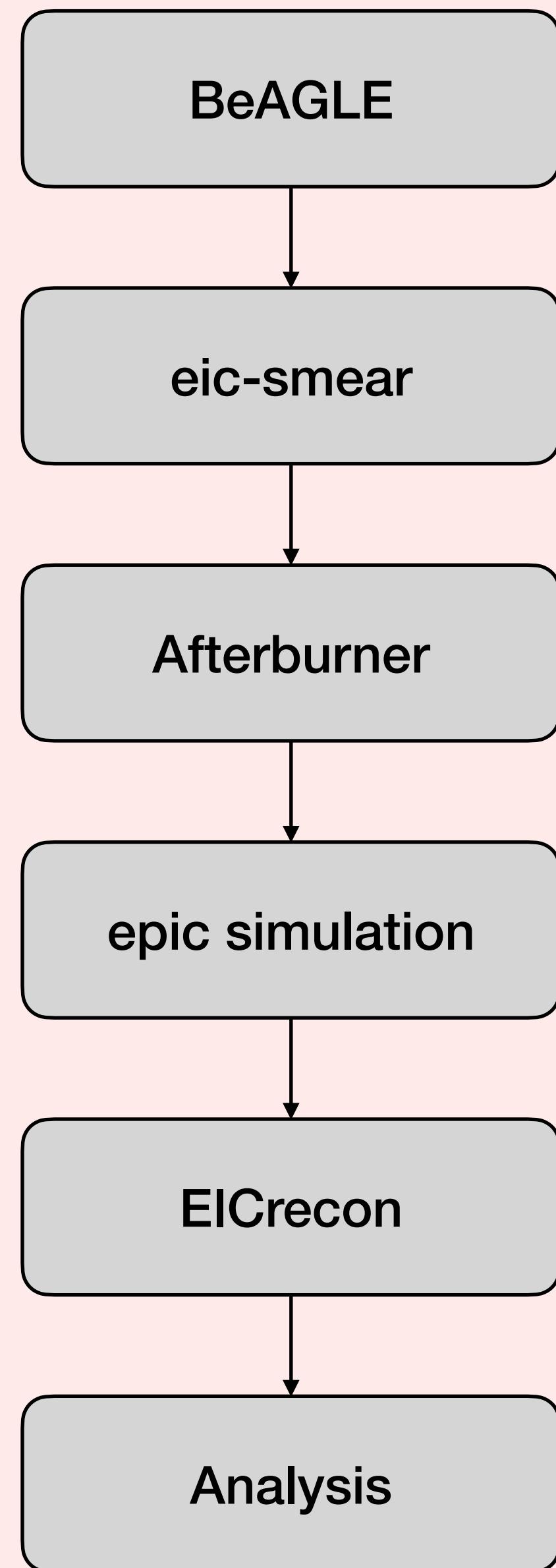
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- all F_2 's are taken from [JAM22](#)



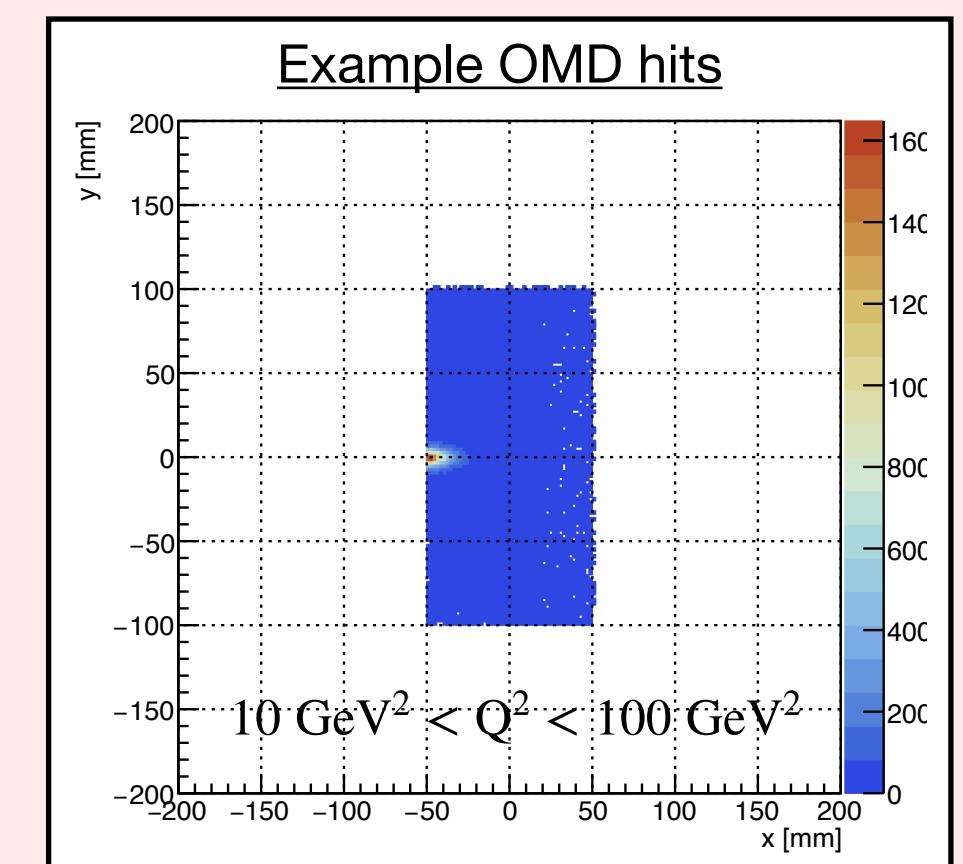
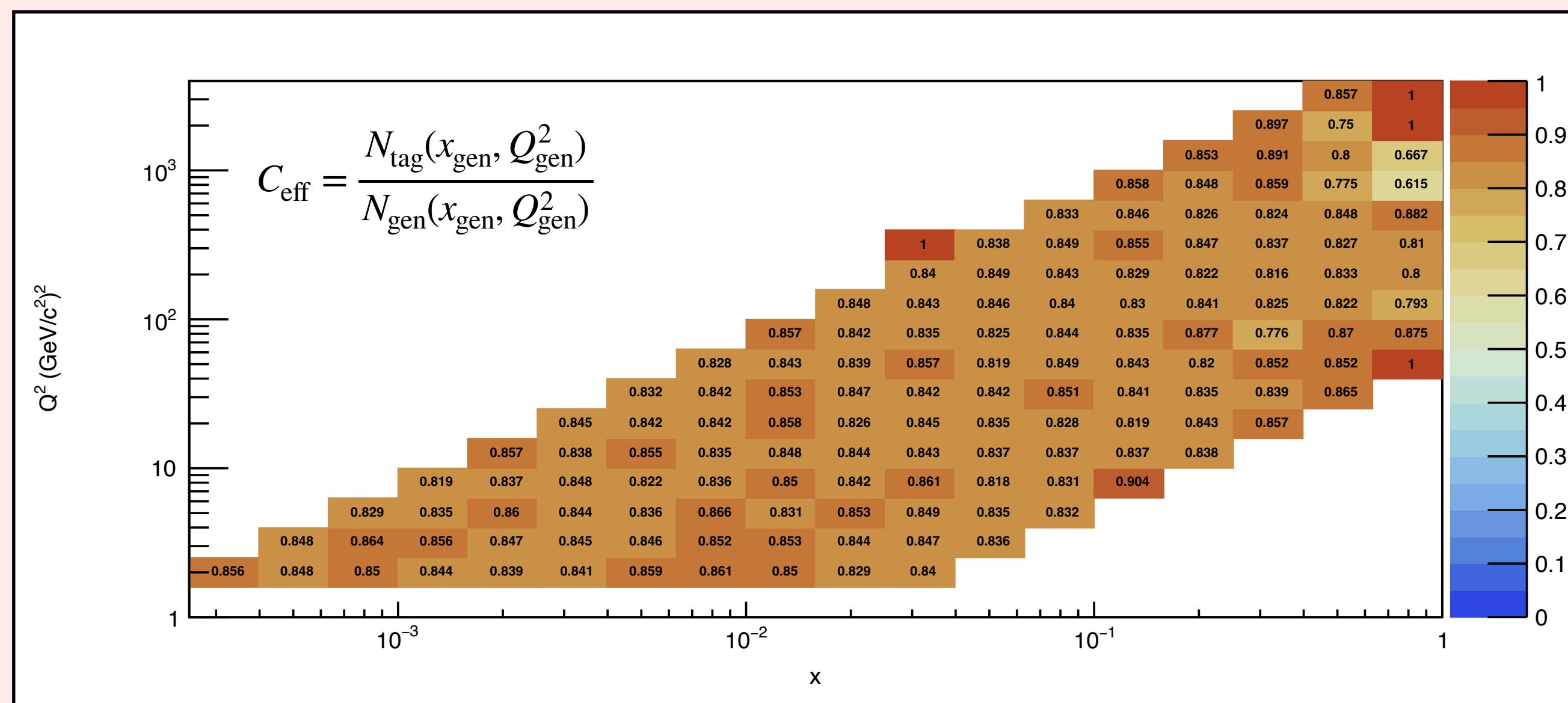
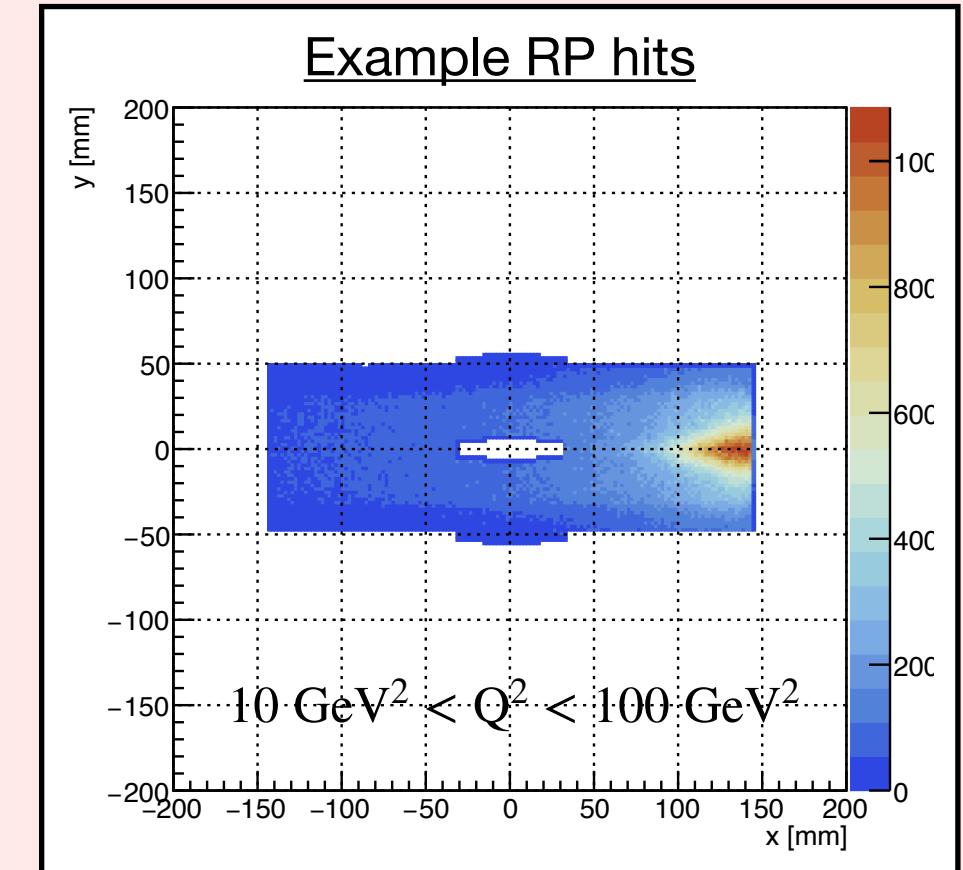
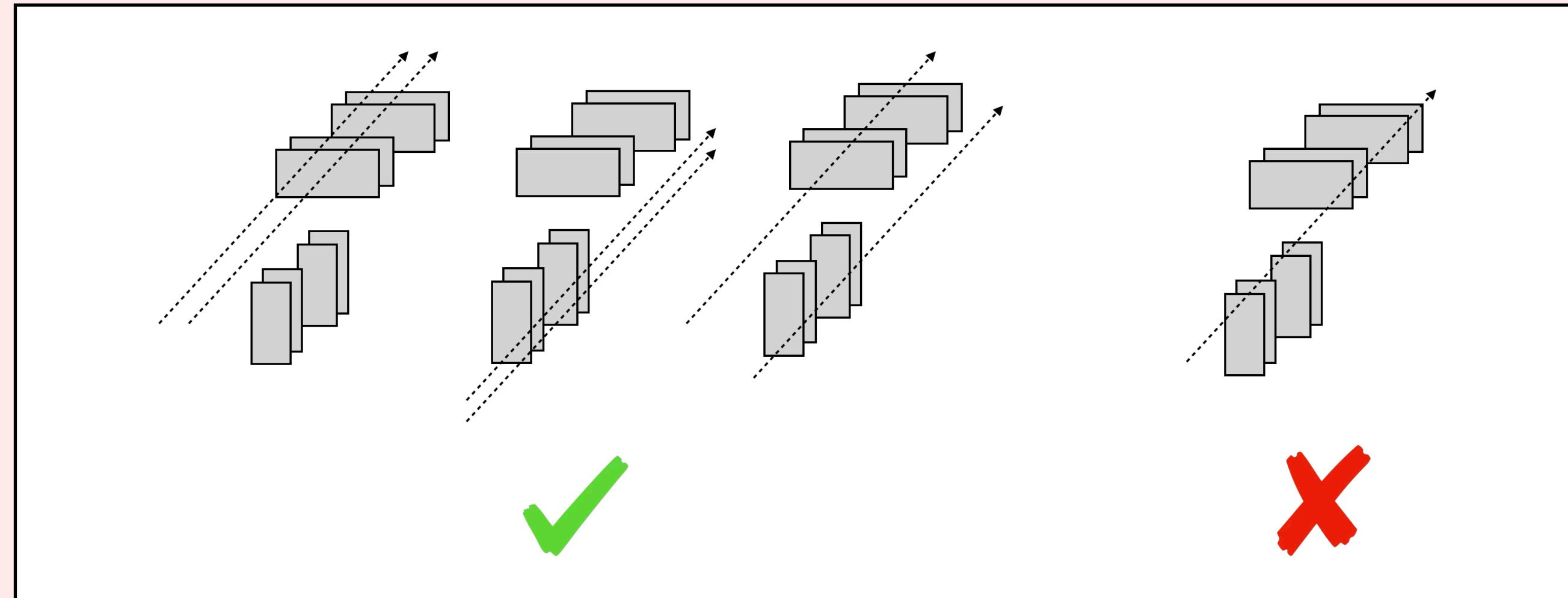
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Projection study analysis chain for g_1^n

19

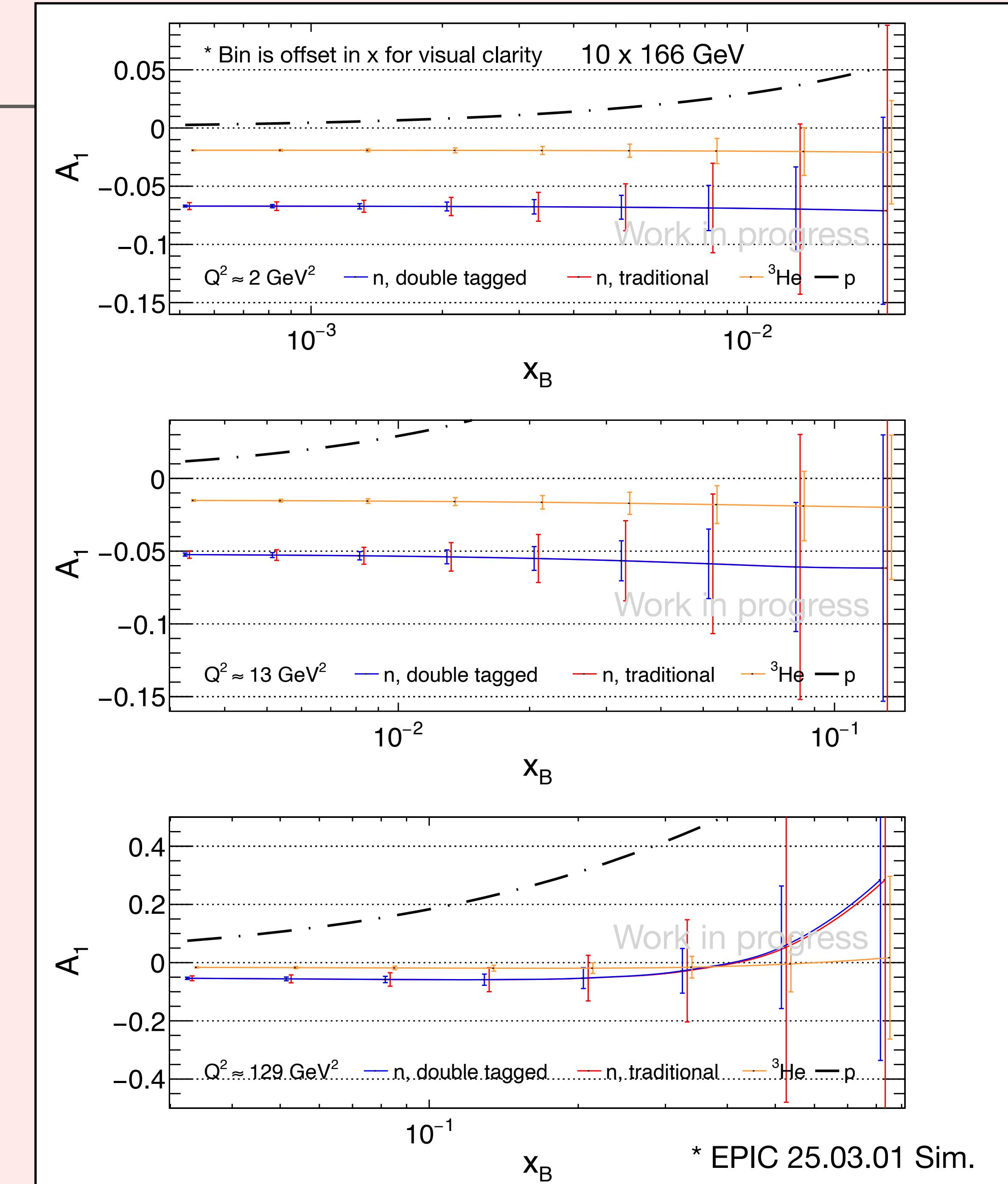


- Currently using a simple tracking algorithm based on hit per plane, $C_{\text{eff}} \sim 83\%$



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- $\mathcal{L} = 8.65 \text{ fb}^{-1}$, $P_e = P_n = 70\%$
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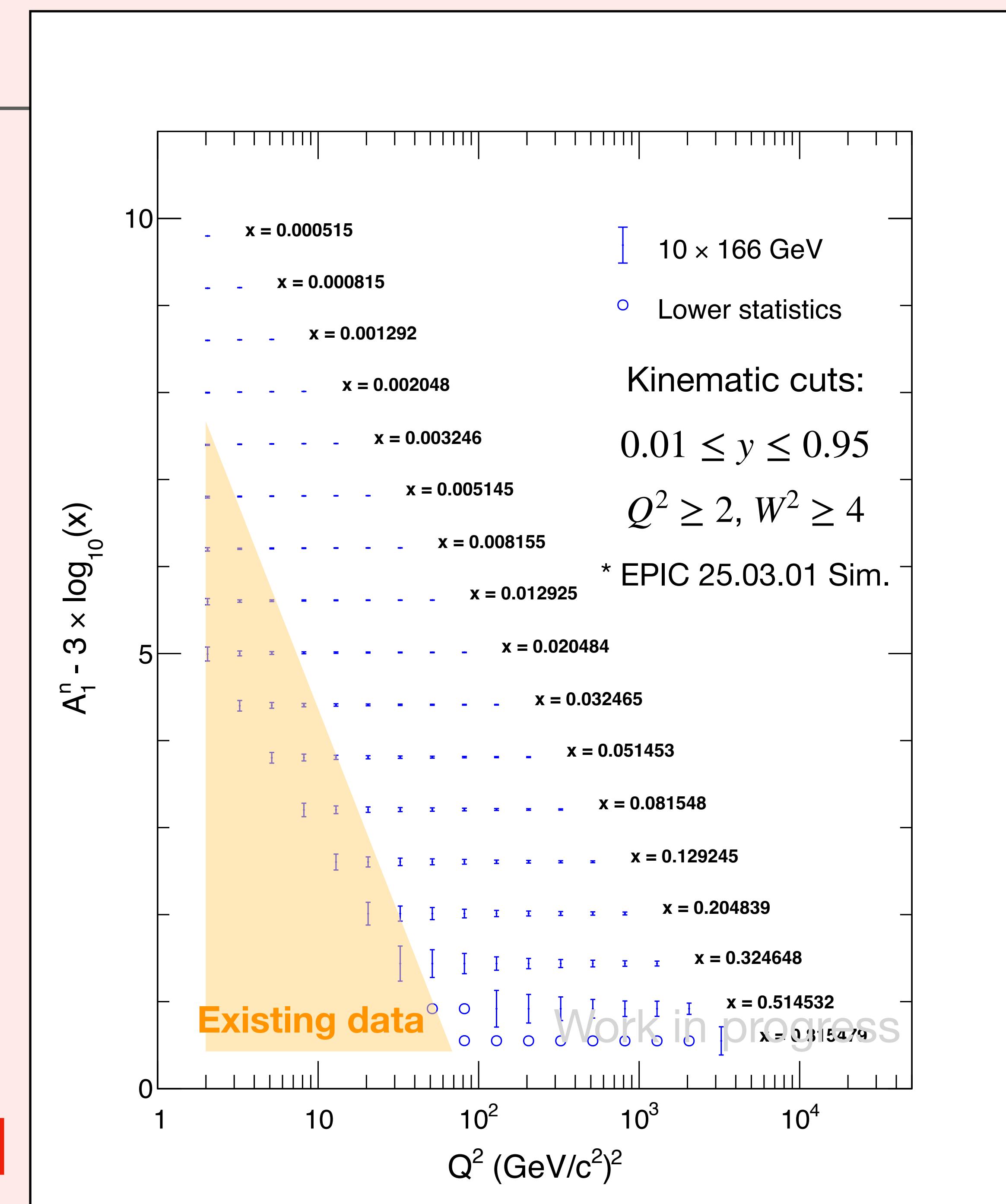
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EIC Early Science

	Species	Energy (GeV)	Luminosity/year (fb ⁻¹)	Electron polarization	p/A polarization
YEAR 1	e+Ru or e+Cu	10 x 115	0.9	NO (Commissioning)	N/A
YEAR 2	e+D e+p	10 x 130	11.4 4.95 - 5.33	LONG	NO TRANS
YEAR 3	e+p	10 x 130	4.95 - 5.33	LONG	TRANS and/or LONG
YEAR 4	e+Au e+p	10 x 100 10 x 250	0.84 6.19 - 9.18	LONG	N/A TRANS and/or LONG
YEAR 5	e+Au e+ ³ He	10 x 100 10 x 166	0.84 8.65	LONG	N/A TRANS and/or LONG

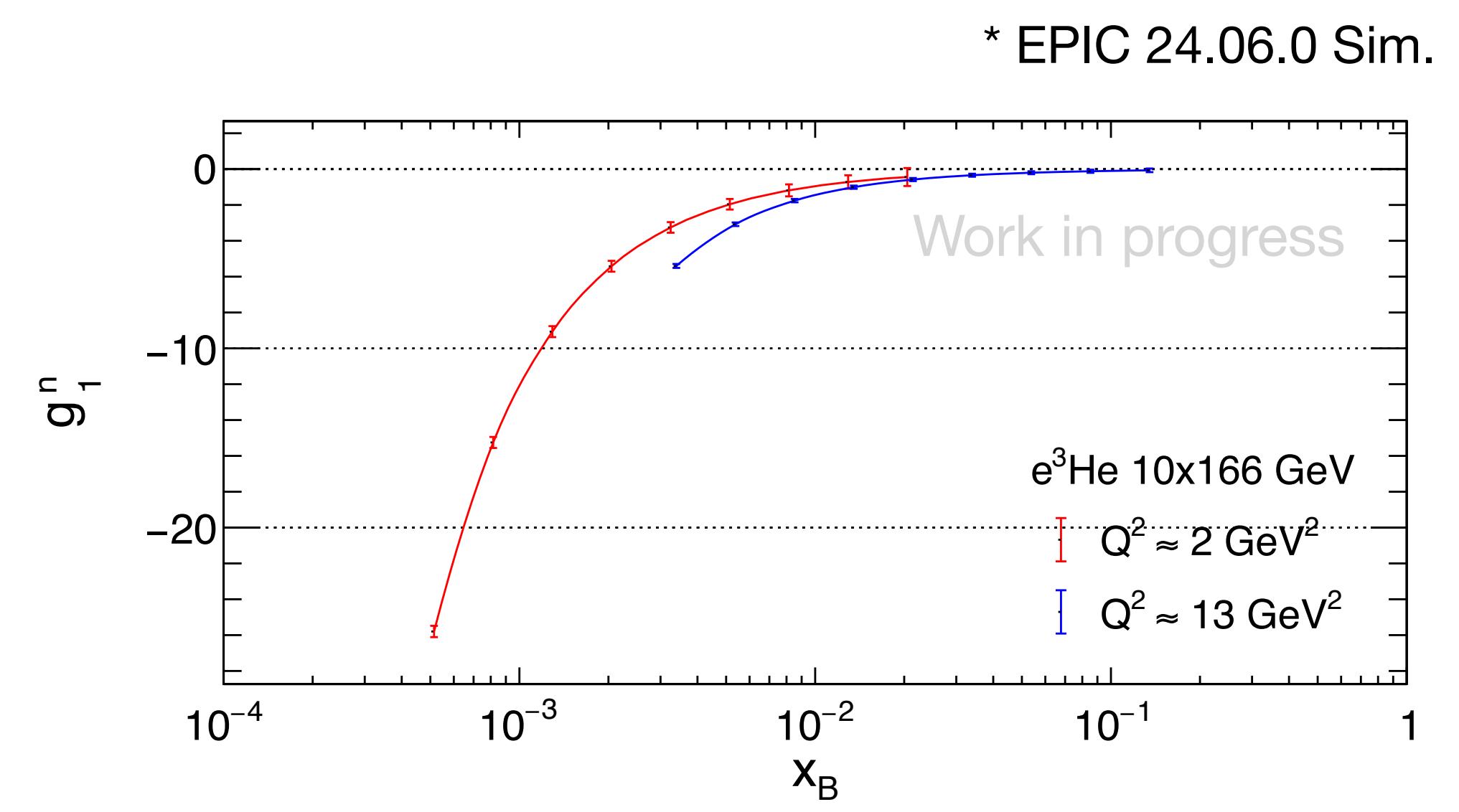
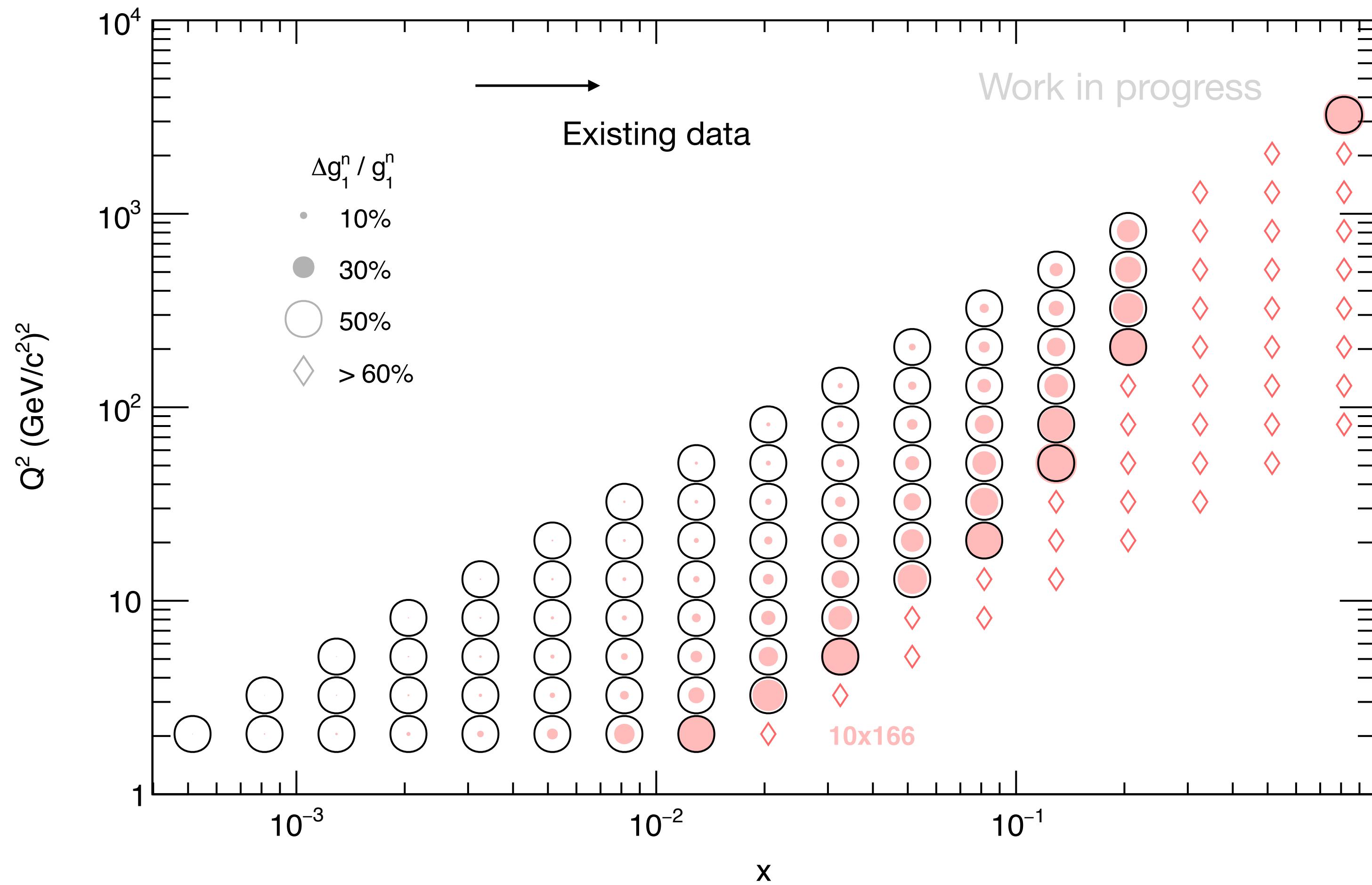
Note: the eA luminosity is per nucleon



Projected g_1^n at EIC Early Science

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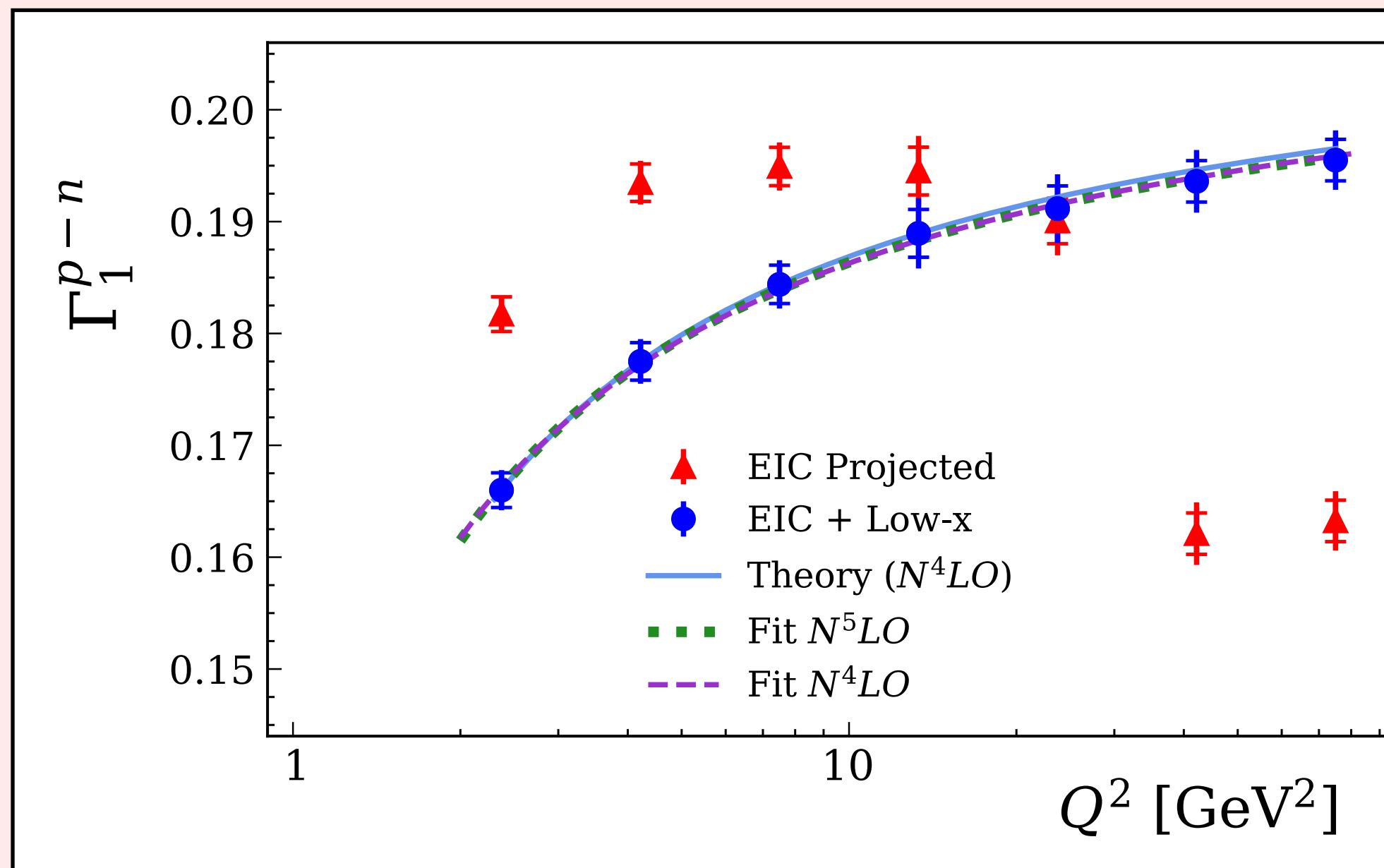
- $A_1 \approx g_1/F_1$ with F_1 calculated from [JAM22](#)
- Statistical uncertainties only



More energies for e and ^3He are available after year 5 of EIC running

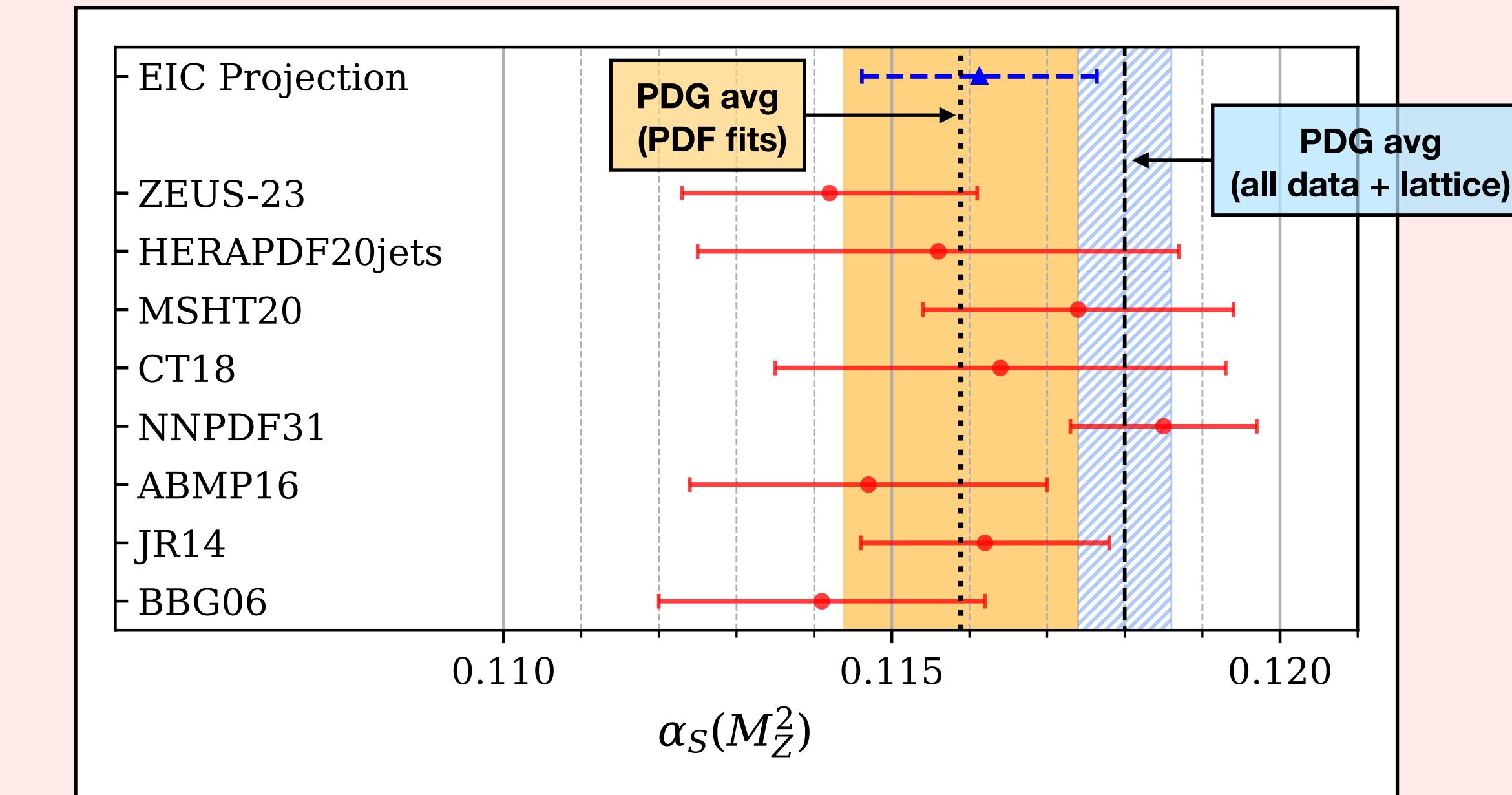
α_s and Bjorken Sum

$$\begin{aligned}\Gamma_1^{p-n}(\alpha_s) &= \Gamma_1^{p-n}(Q^2) = \int_0^{1^-} (g_1^p(Q^2) - g_1^n(Q^2)) dx = \sum_{n>0} \frac{\mu_{2n}^{p-n}(\alpha_s)}{Q^{2n-2}} \\ &= \frac{g_A}{6} \left[1 - \frac{\alpha_s(Q^2)}{\pi} - 3.58 \left(\frac{\alpha_s(Q^2)}{\pi} \right)^2 - 20.21 \left(\frac{\alpha_s(Q^2)}{\pi} \right)^3 - 175.7 \left(\frac{\alpha_s(Q^2)}{\pi} \right)^4 - \mathcal{O}((\alpha_s)^5) \right] \quad \text{at finite } Q^2\end{aligned}$$



Assume $\mathcal{L} = 10 \text{ fb}^{-1}$ for each setting:

5x41, 10x100, 18x275 ep DIS
5x41, 10x100, 18x166 en DIS



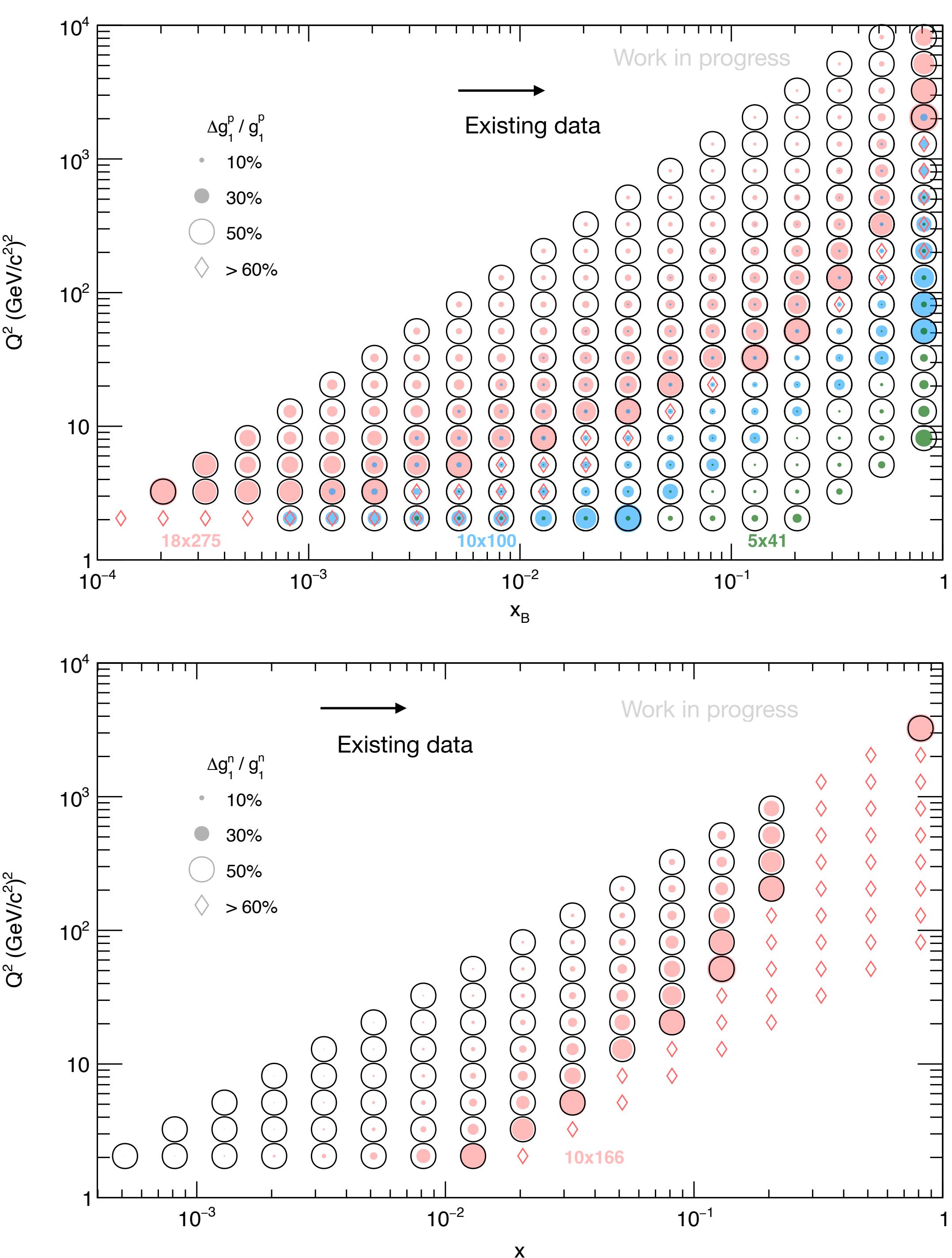
<https://doi.org/10.1103/PhysRevD.110.074004>

$\Delta \alpha_s(M_Z)/\alpha_s(M_Z) = 1.3 \%$

Summary

- New g_1 measurement at EIC will help investigate the quark and gluon spin contributions to the nucleon spin
- Both g_1^p and g_1^n will be measured; double tagging method will help constrain model uncertainty and test light nuclei model
- EIC will be able to measure g_1 at low x region which has not been explored
- g_1^p and g_1^n lead to Bjorken sum which will provide a new high precision measurement on the α_s

Thank you!



Slides++ : Double Spin Asymmetry

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$$\frac{d^3 \Delta\sigma(\beta)}{dQ^2 dx d\phi} = \frac{4\alpha^2}{Q^2} y \left\{ \cos \beta \left[\left(1 - \frac{y}{2} - \frac{\gamma^2 y^2}{4} \right) g_1(x, Q^2) - \frac{\gamma^2 y}{4} g_2(x, Q^2) \right] \right.$$

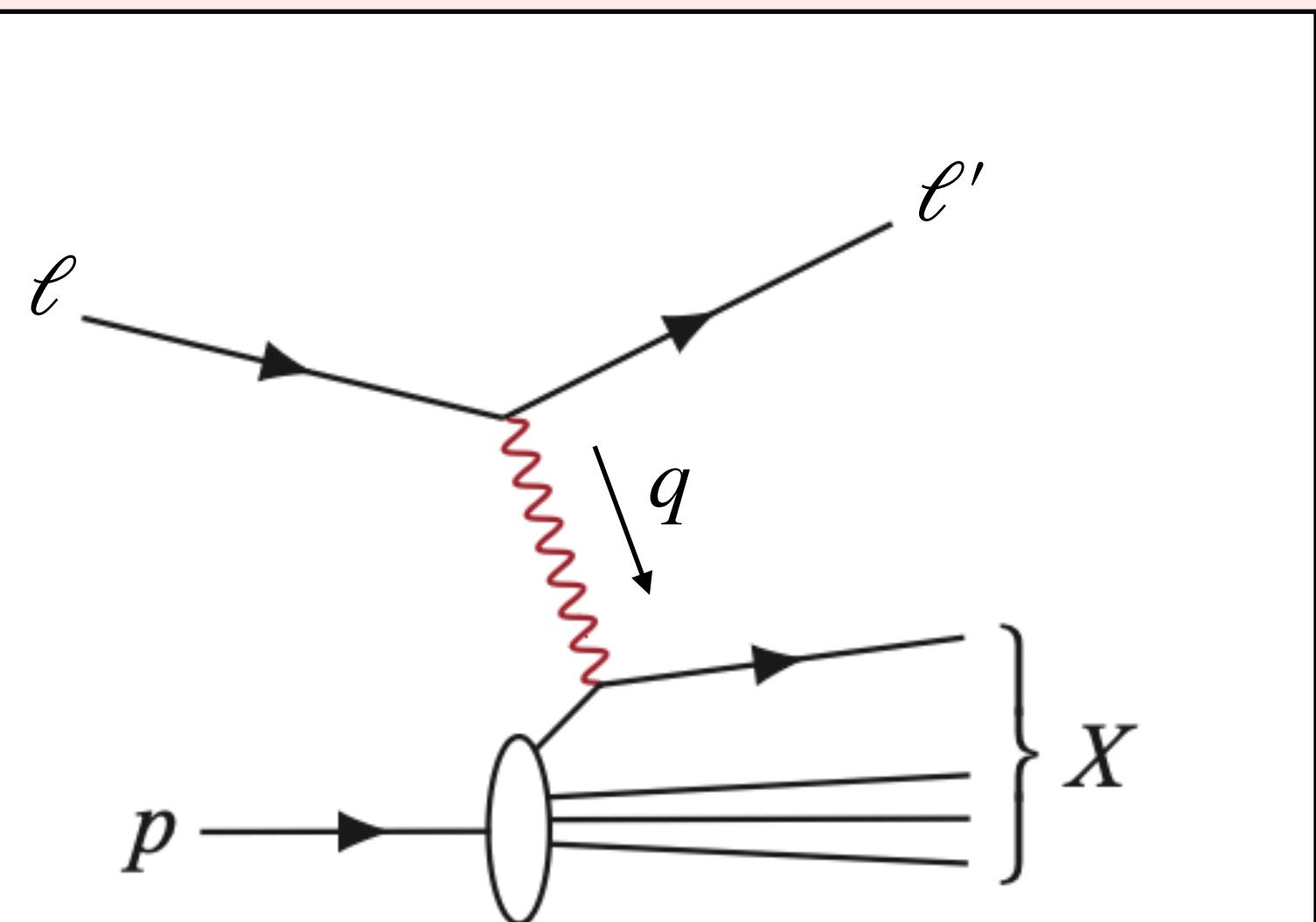
$$\left. - \cos \phi \sin \beta \frac{\sqrt{Q^2}}{\nu} \left(1 - y - \frac{\gamma^2 y^2}{4} \right)^{\frac{1}{2}} \left[\frac{y}{2} g_1(x, Q^2) + g_2(x, Q^2) \right] \right\}$$

- Direct measurement is more challenging
- Easier to measure via cross section ratio
- Systematic is largely canceled out

$$g_1 = \frac{F_2}{2x(1+R)} (A_1 + \gamma A_2)$$

$$A_1(x, Q^2) \equiv \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{A_{||}}{D(1 + \eta\xi)} - \frac{\eta A_{\perp}}{d(1 + \eta\xi)}$$

$$A_2(x, Q^2) \equiv \frac{2\sigma_{LT}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{\xi A_{||}}{D(1 + \eta\xi)} - \frac{A_{\perp}}{d(1 + \eta\xi)}$$



NC Inclusive DIS: $e + p \rightarrow e' + X$

<https://arxiv.org/abs/2103.05419>

