

Spin manipulation in atomic and molecular systems for nuclear spin applications

Chrysovalantis Kannis
with

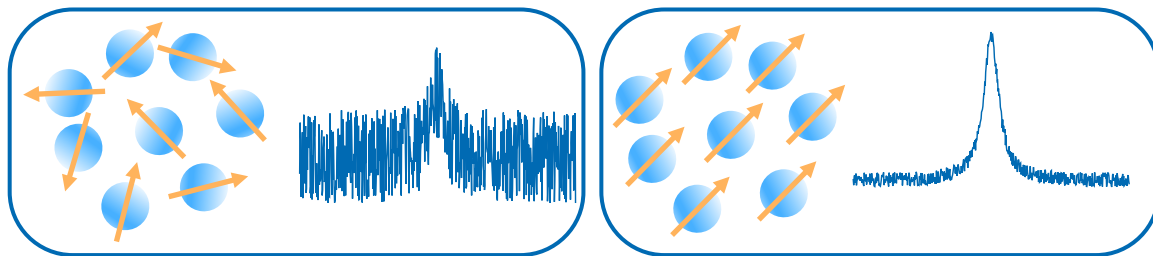
Ralf Engels, Tarek El-Kordy, Nicolas Faatz, Simon J. Pütz,
Vincent Verhoeven, T. Peter Rakitzis, and Markus Büscher

23.09.2025

Why nuclear spin polarization is important

■ Physics and medicine

- NMR/MRI signal enhancement



■ Polarized fusion

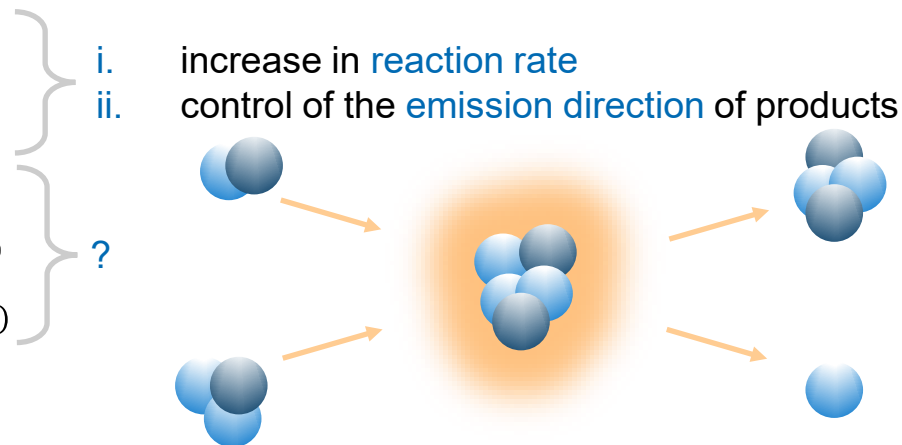
■ Five-nucleon reactions

- $d + t \longrightarrow \alpha (3.5 \text{ MeV}) + n (14.1 \text{ MeV})$
- $d + h \longrightarrow \alpha (3.6 \text{ MeV}) + p (14.7 \text{ MeV})$

■ Four-nucleon reactions

- $d + d \xrightarrow{50\%} t (1.01 \text{ MeV}) + p (3.02 \text{ MeV})$
- $d + d \xrightarrow{50\%} h (0.82 \text{ MeV}) + n (2.45 \text{ MeV})$

- Ralf Engels-Polarized Fusion
(Plenary talk, Thu 25/09, 8:30)



■ Spin and angular momentum in bound systems

- Atomic states: nuclear spin I , electron spin S (with $L = 0$)
- Molecular states: nuclear spins I , rotational angular momentum J (with $S = 0$)

■ Hyperfine regime

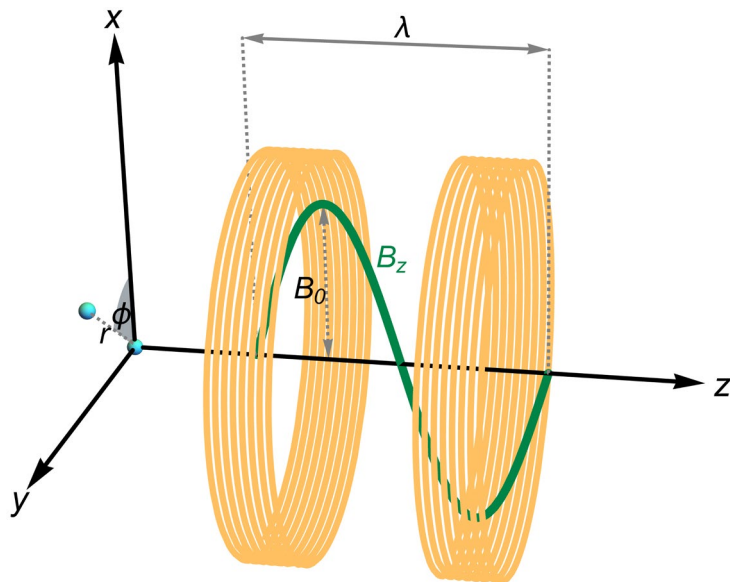
- Interactions between the nuclear spins and residual angular momenta: H_0
- Spin states
 - Uncoupled basis: $|m_S, m_I\rangle, |m_J, m_{I_1}, m_{I_2}\rangle$
 - Coupled basis: $|F, m_F\rangle$ with $\mathbf{F} = \mathbf{S} \otimes \mathbb{1} + \mathbb{1} \otimes \mathbf{I}$ or $|I_t, F, m_F\rangle$ $\mathbf{F} = \mathbf{J} \otimes \mathbb{1} + \mathbb{1} \otimes \mathbf{I}_t$ and $\mathbf{I}_t = \mathbf{I}_1 \otimes \mathbb{1} + \mathbb{1} \otimes \mathbf{I}_2$

■ Density operator formalism

- Ensemble of nonidentically prepared atoms or molecules $\rightarrow$ statistical mixture of states
- Density matrix: diagonal elements $\rightarrow$ populations, off-diagonal elements $\rightarrow$ coherences
- Liouville-von Neumann equation: $i\hbar \frac{\partial \rho}{\partial t} = [H, \rho]$
- Average of observable O : $\langle O \rangle = \text{Tr}(\rho O)$

■ External perturbations

- Interaction with a magnetic field $\mathbf{B}$: $H_B = -\boldsymbol{\mu} \cdot \mathbf{B}$, with total Hamiltonian $H = H_0 + H_B$
- Sinusoidal field, inspired by P. G. Sona [Energ. Nucl. **14**, 295 (1967)]



$$\mathbf{B} = B_z \hat{\mathbf{z}} + B_r \hat{\mathbf{r}}, \text{ with } B_r = -\frac{r}{2} \frac{\partial B_z}{\partial z}$$

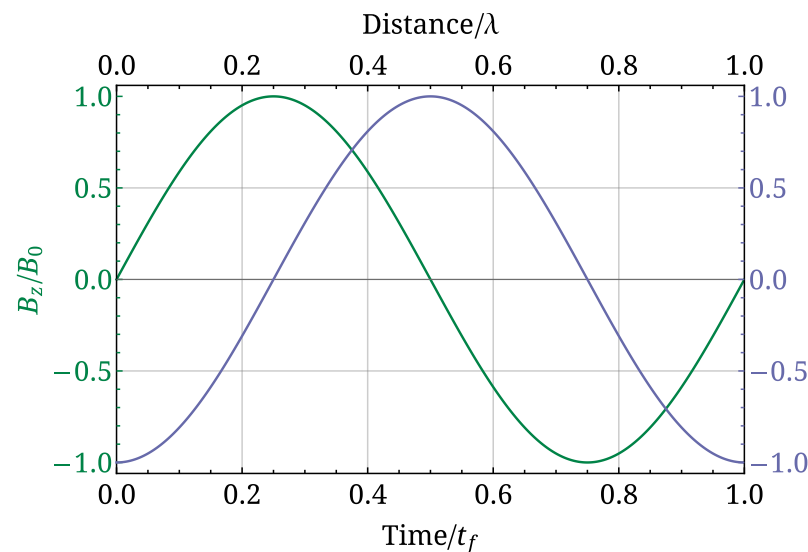
- B_z : static longitudinal component
- Particle motion: along z with constant, nonrelativistic velocity v

$$\text{Laboratory frame: } \mathbf{B} = B_0 \sin\left(\frac{2\pi z}{\lambda}\right) \hat{\mathbf{z}} - B_0 \frac{\pi r}{\lambda} \cos\left(\frac{2\pi z}{\lambda}\right) \hat{\mathbf{r}}$$

$$\text{Particle rest frame: } \mathbf{B} = B_0 \sin\left(\frac{2\pi v t}{\lambda}\right) \hat{\mathbf{z}} - B_0 \frac{\pi r}{\lambda} \cos\left(\frac{2\pi v t}{\lambda}\right) \hat{\mathbf{r}}$$

■ External perturbations

- Interaction with a magnetic field $\mathbf{B}$: $H_B = -\boldsymbol{\mu} \cdot \mathbf{B}$, with total Hamiltonian $H = H_0 + H_B$
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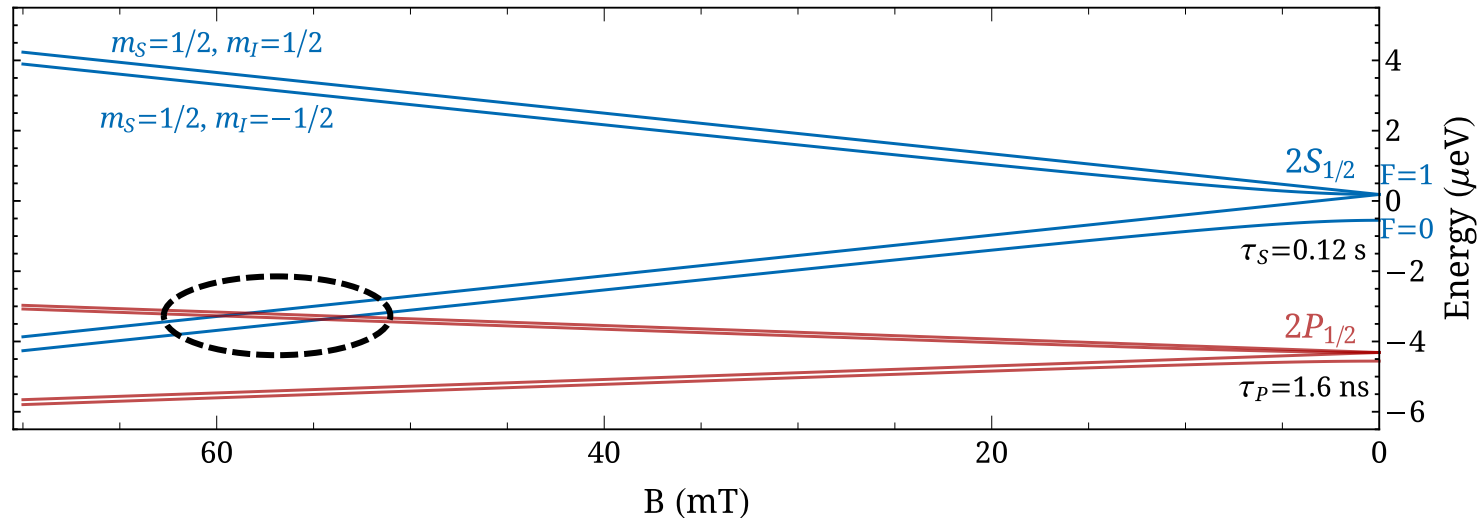
Laboratory frame: $\mathbf{B} = B_0 \sin\left(\frac{2\pi z}{\lambda}\right) \hat{\mathbf{z}} - B_0 \frac{\pi r}{\lambda} \cos\left(\frac{2\pi z}{\lambda}\right) \hat{\mathbf{r}}$

Particle rest frame: $\mathbf{B} = B_0 \sin\left(\frac{2\pi vt}{\lambda}\right) \hat{\mathbf{z}} - B_0 \frac{\pi r}{\lambda} \cos\left(\frac{2\pi vt}{\lambda}\right) \hat{\mathbf{r}}$

$$r \ll \lambda$$

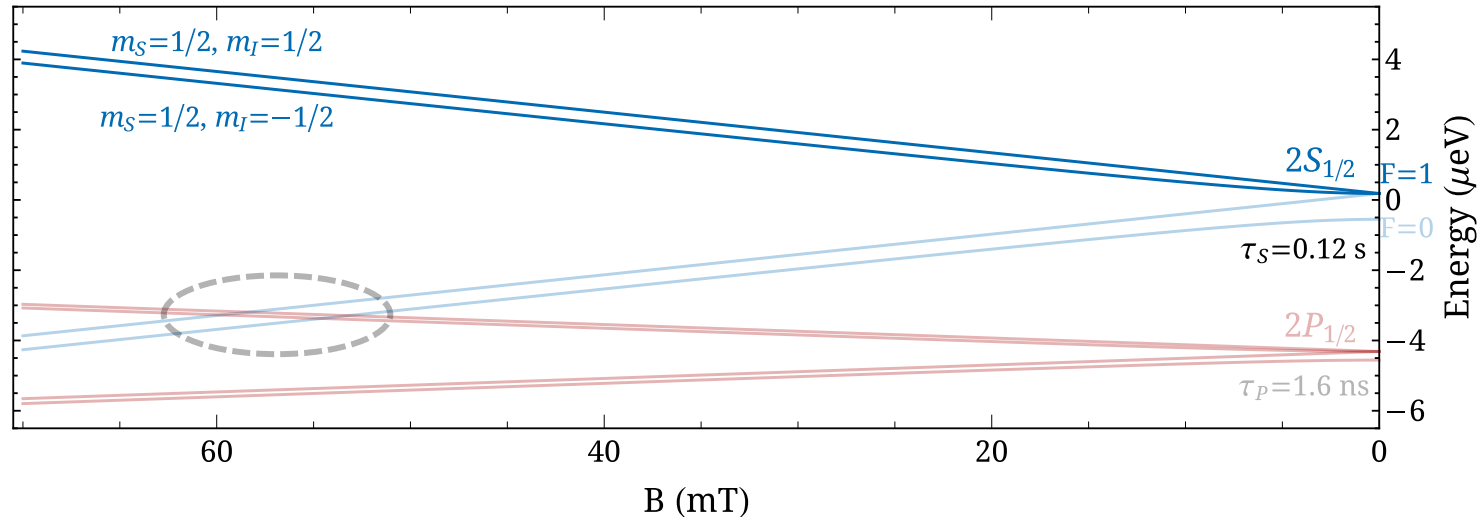
Hydrogen atom

- Ground state: protons capture polarized electrons in an optically pumped, polarized Rb vapor cell (in high magnetic field)
- Metastable state: use of the $2S_{1/2} - 2P_{1/2}$ level crossing (also at high magnetic field)



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- Metastable state: use of the $2S_{1/2} - 2P_{1/2}$ level crossing (also at high magnetic field)



Incoherent initial preparation

Hydrogen atom

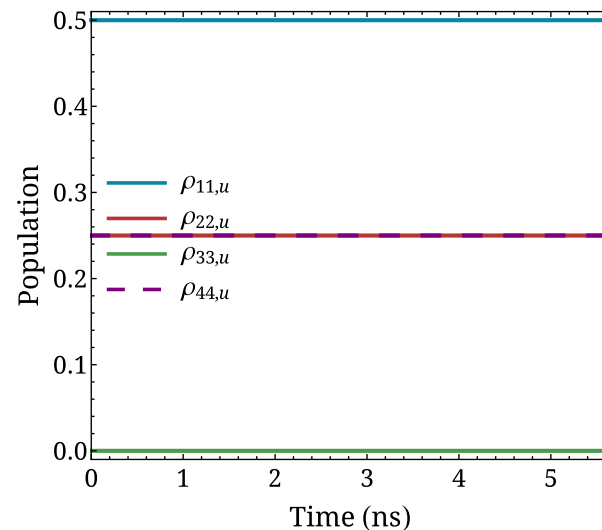
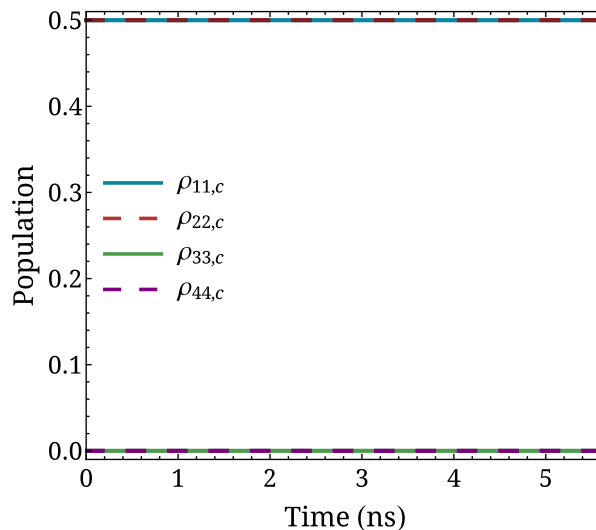
$$\rho_c = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

coupled basis uncoupled basis

	$ F, m_F\rangle$	$ m_S, m_I\rangle$
1.	$ 1, 1\rangle$	$ 1/2, 1/2\rangle$
2.	$ 1, 0\rangle$	$ 1/2, -1/2\rangle$
3.	$ 1, -1\rangle$	$ -1/2, -1/2\rangle$
4.	$ 0, 0\rangle$	$ -1/2, 1/2\rangle$

$$H_0 = \frac{A_{2S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S} \quad \text{with} \quad \frac{A_{2S}}{h} = 177.6 \text{ MHz}$$

$$t_f = \frac{h}{\Delta E_{hfs}} = \frac{h}{A_{2S}} \sim 5.6 \text{ ns}$$



■ Hydrogen atom

$$\rho_c = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

coupled basis uncoupled basis

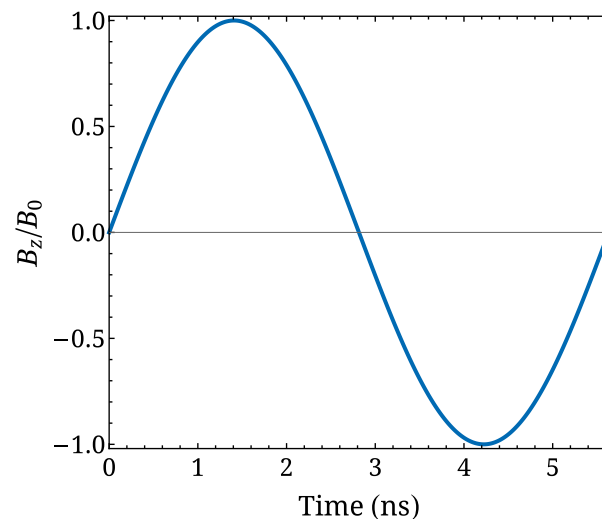
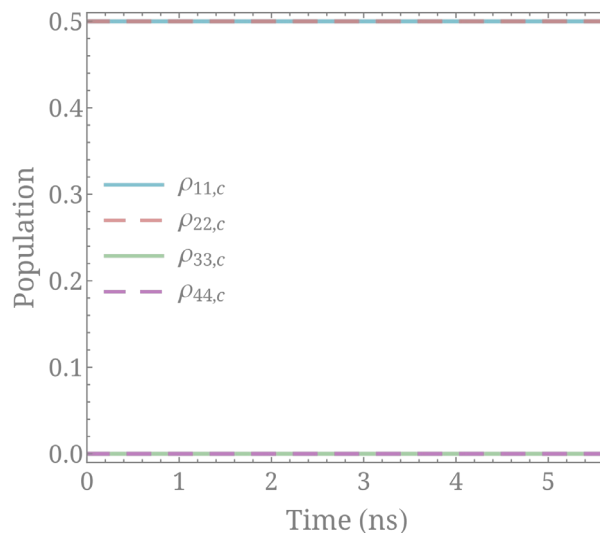
	$ F, m_F\rangle$	$ m_S, m_I\rangle$
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4.	$ 0, 0\rangle$	$ -1/2, 1/2\rangle$

➤ How is population evolution affected by $B_z = B_0 \sin\left(\frac{2\pi\nu t}{\lambda}\right)$?

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$$t_f = \frac{h}{\Delta E_{hfs}} = \frac{h}{A_{2S}} \sim 5.6 \text{ ns}$$

$$\lambda = v t_f$$



Hydrogen atom

$$\rho_c = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

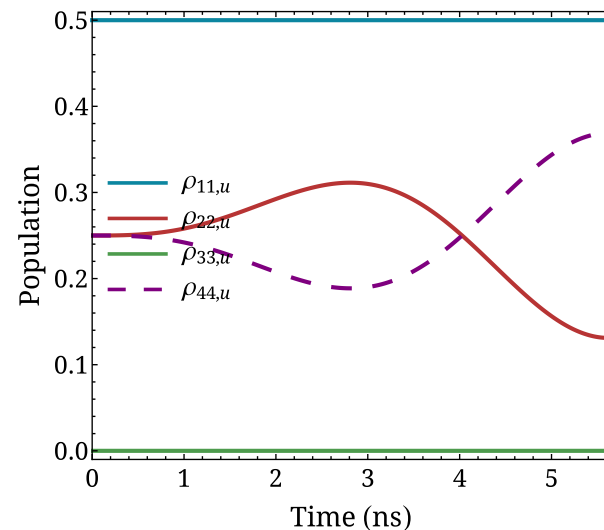
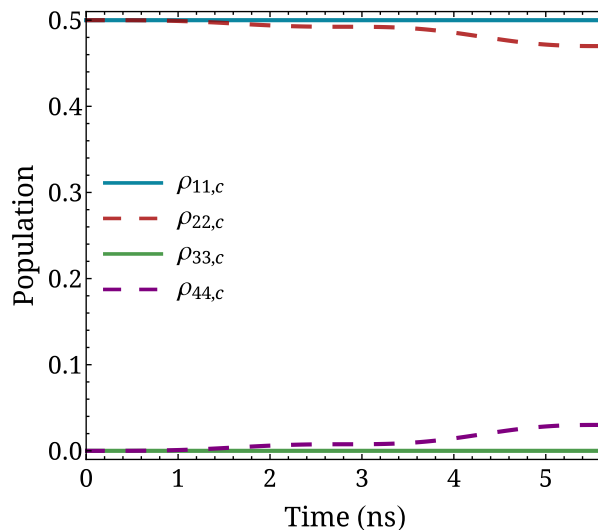
coupled basis uncoupled basis

	$ F, m_F\rangle$	$ m_S, m_I\rangle$
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$$H = H_0 + H_B = \frac{A_{2S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S} - \left(\frac{g_S \mu_B}{\hbar} \mathbf{S} + \frac{g_I \mu_N}{\hbar} \mathbf{I} \right) \cdot \mathbf{B}$$

$B_0 = 1 \text{ mT}$



Hydrogen atom

$$\rho_c = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

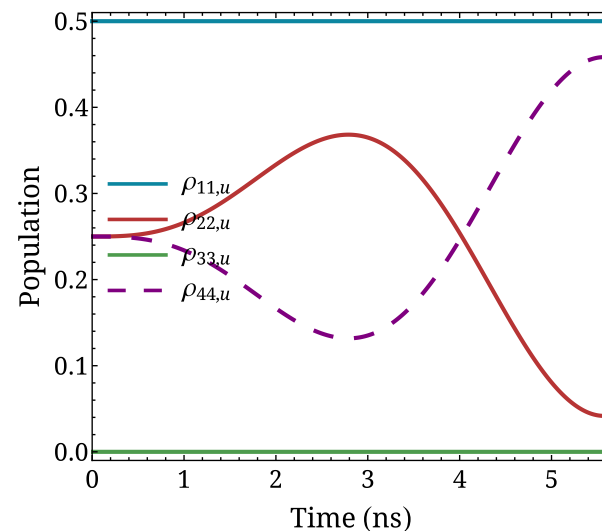
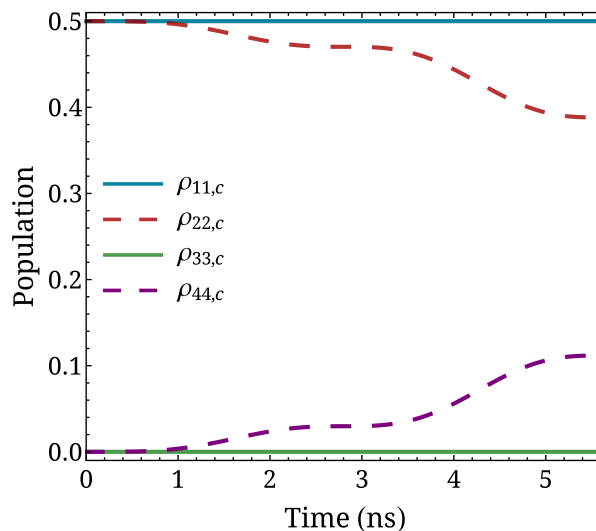
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$B_0 = 2 \text{ mT}$



Hydrogen atom

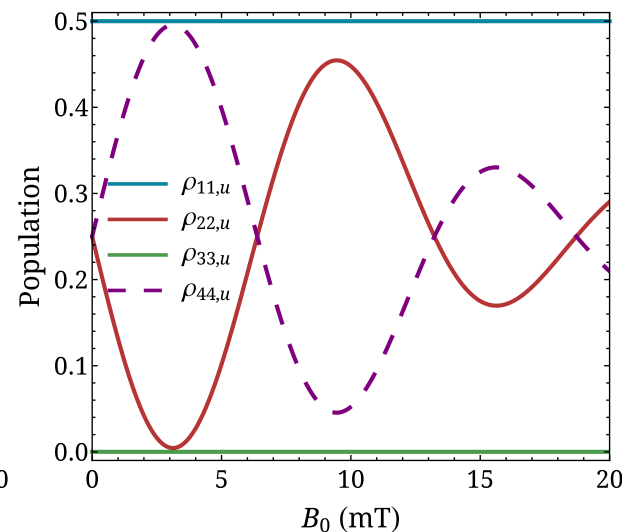
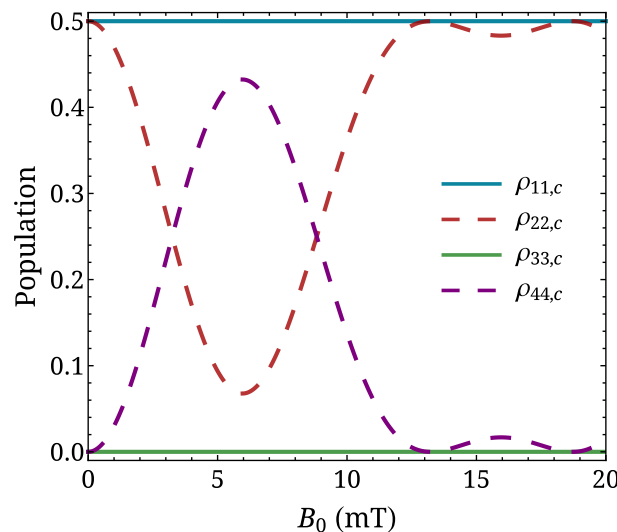
$$\rho_c = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

coupled basis uncoupled basis

	$ F, m_F\rangle$	$ m_S, m_I\rangle$
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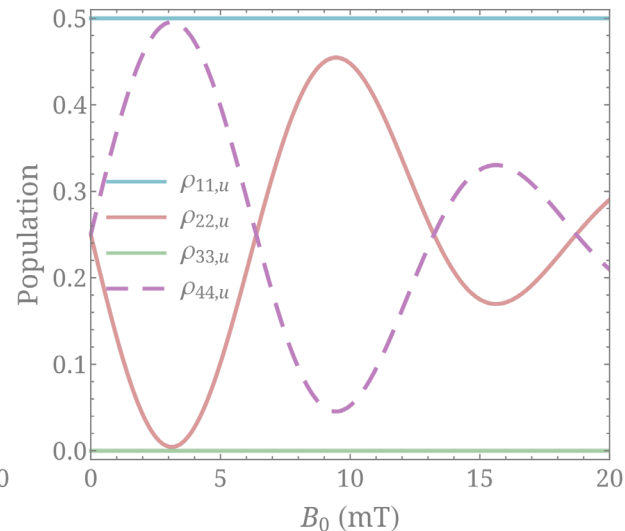
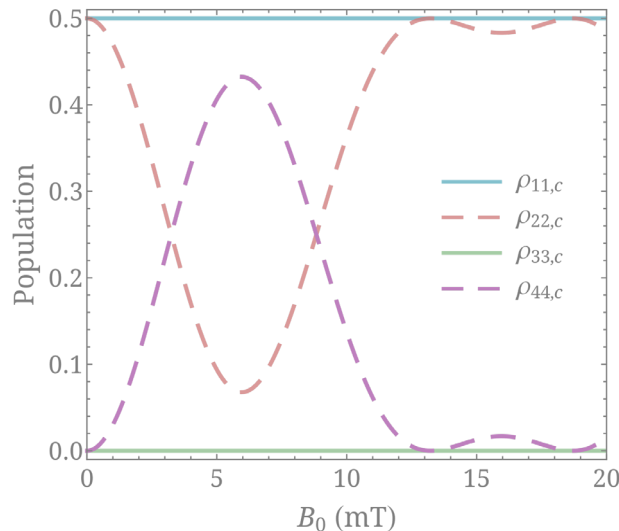
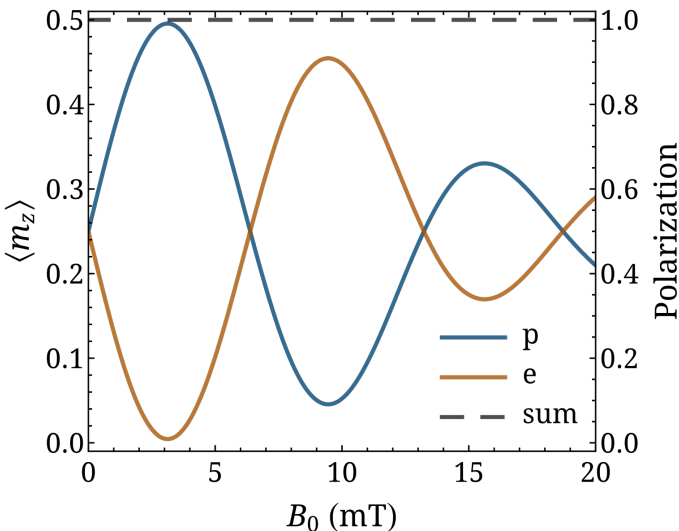
$$H = H_0 + H_B = \frac{A_{2S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S} - \left(\frac{g_S \mu_B}{\hbar} \mathbf{S} + \frac{g_I \mu_N}{\hbar} \mathbf{I} \right) \cdot \mathbf{B}$$



Hydrogen atom

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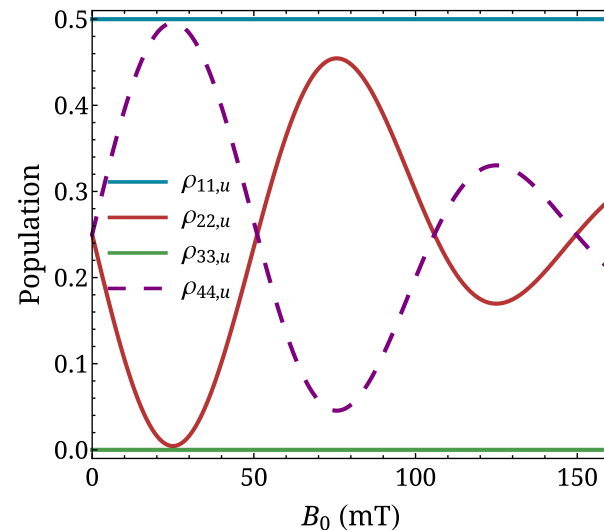
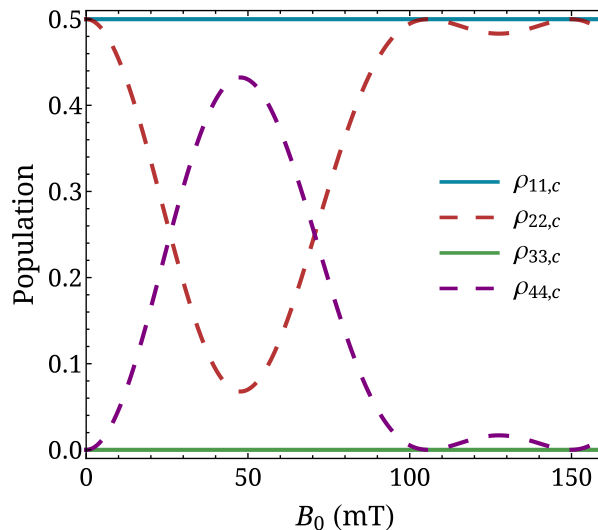
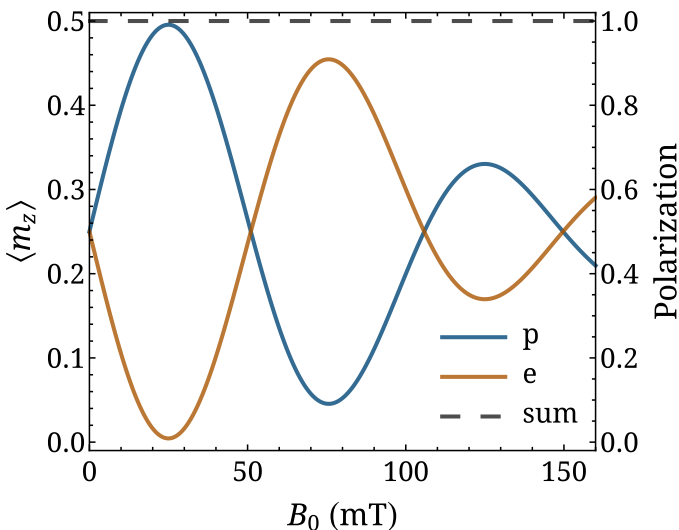


Hydrogen atom

- Ground state ($A_{1S} = 1.42$ GHz)

➤ How is population evolution affected by $B_z = B_0 \sin\left(\frac{2\pi\nu t}{\lambda}\right)$?

$$H = H_0 + H_B = \frac{A_{1S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S} - \left(\frac{g_S \mu_B}{\hbar} \mathbf{S} + \frac{g_I \mu_N}{\hbar} \mathbf{I} \right) \cdot \mathbf{B}$$



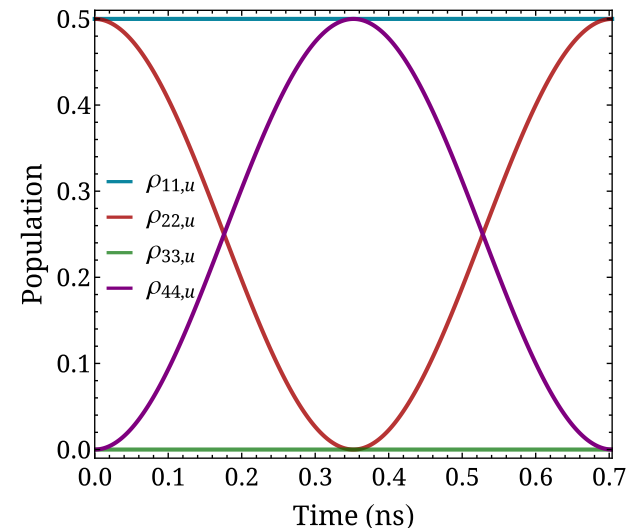
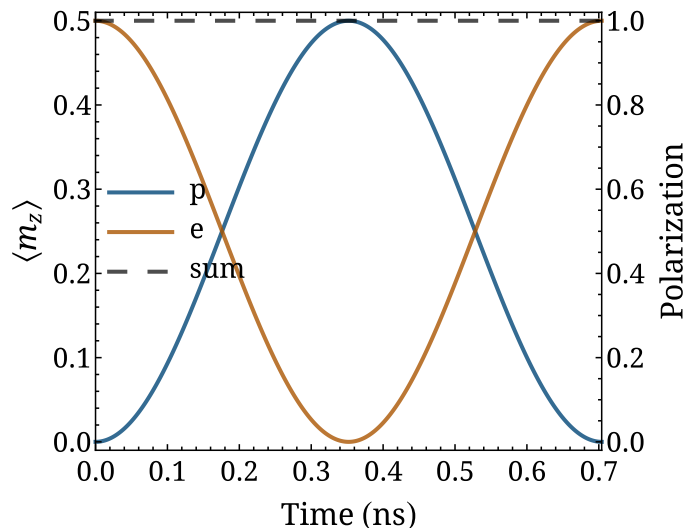
■ Hydrogen atom

- Ground state: Molecular photodissociation of a hydrogen halide
- Example reaction: $\text{HY} + hf \rightarrow \dots \rightarrow \text{H} \left(m_{\text{H}} = \pm \frac{1}{2} \right) + \text{Y} \left(m_{\text{Y}} = \pm \frac{1}{2} \right)$

$$H_0 = \frac{A_{1S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S}$$

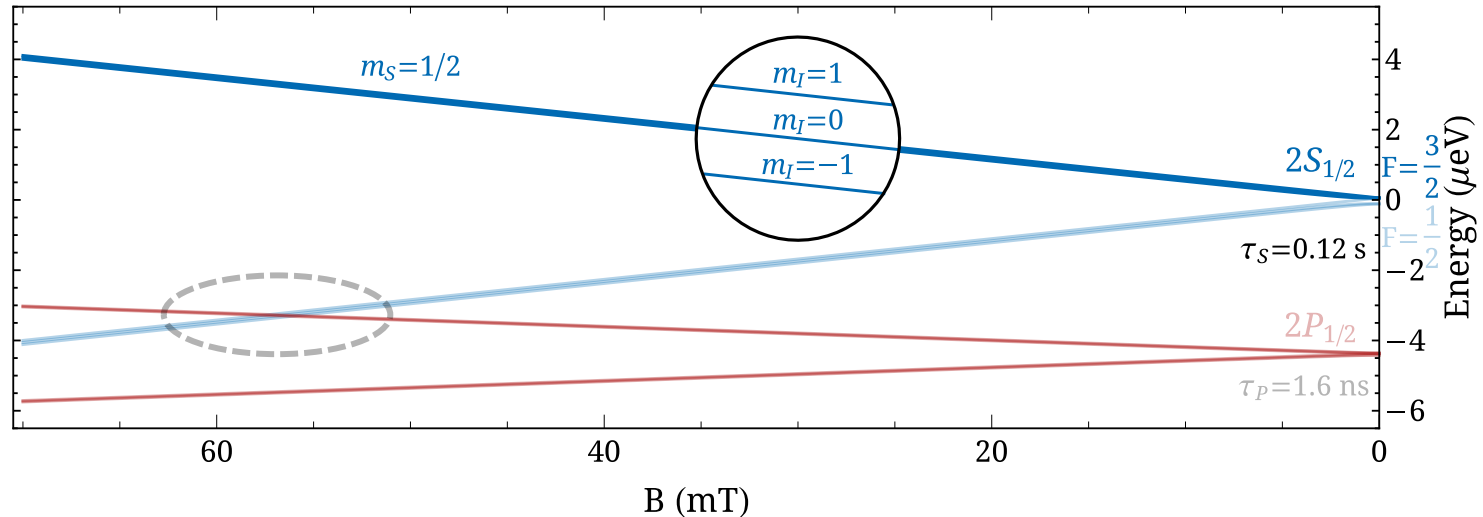
$$\rho_u = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$\rho_c = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/4 & 0 & 1/4 \\ 0 & 0 & 0 & 0 \\ 0 & 1/4 & 0 & 1/4 \end{pmatrix}$$



■ Deuterium atom

- Ground state: similar to hydrogen
- Metastable state: use of the $2S_{1/2} - 2P_{1/2}$ level crossing (at high magnetic field)



■ Deuterium atom

$$\rho_c = \begin{pmatrix} 1/3 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1/3 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1/3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

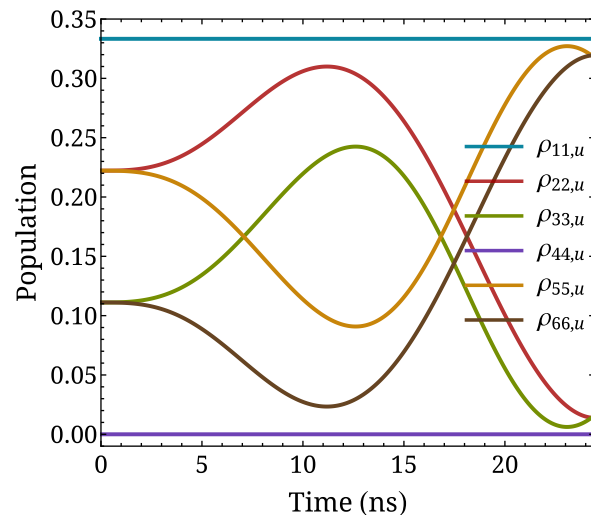
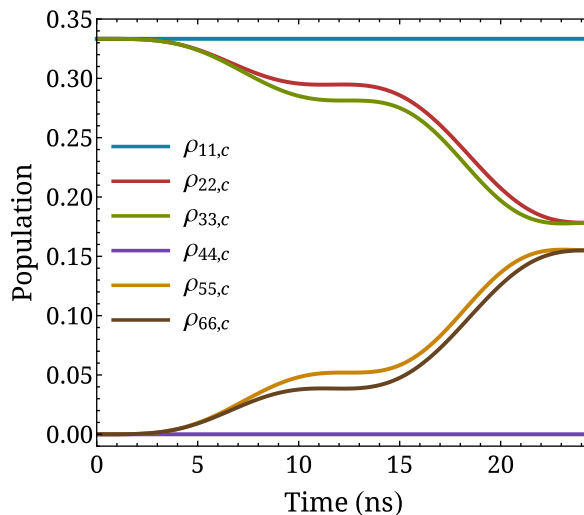
	coupled basis $ F, m_F\rangle$	uncoupled basis $ m_S, m_I\rangle$
1.	$ 3/2, 3/2\rangle$	$ 1/2, 1\rangle$
2.	$ 3/2, 1/2\rangle$	$ 1/2, 0\rangle$
3.	$ 3/2, -1/2\rangle$	$ 1/2, -1\rangle$
4.	$ 3/2, -3/2\rangle$	$ -1/2, -1\rangle$
5.	$ 1/2, -1/2\rangle$	$ -1/2, 0\rangle$
6.	$ 1/2, 1/2\rangle$	$ -1/2, 1\rangle$

$$H = H_0 + H_B = \frac{A_{2S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S} - \left(\frac{g_S \mu_B}{\hbar} \mathbf{S} + \frac{g_I \mu_N}{\hbar} \mathbf{I} \right) \cdot \mathbf{B}$$

$$H_0 = \frac{A_{2S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S} \text{ with } \frac{A_{2S}}{h} = 27.28 \text{ MHz}$$

$$t_f = \frac{h}{\Delta E_{hfs}} = \frac{2h}{3A_{2S}} \sim 24.4 \text{ ns}$$

$$B_0 = 0.76 \text{ mT}$$

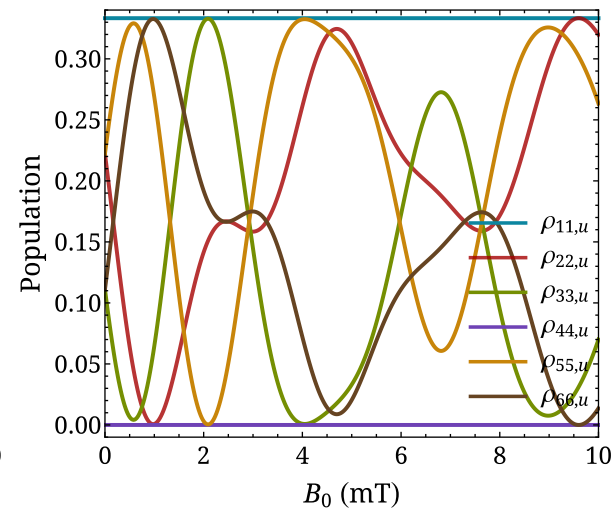
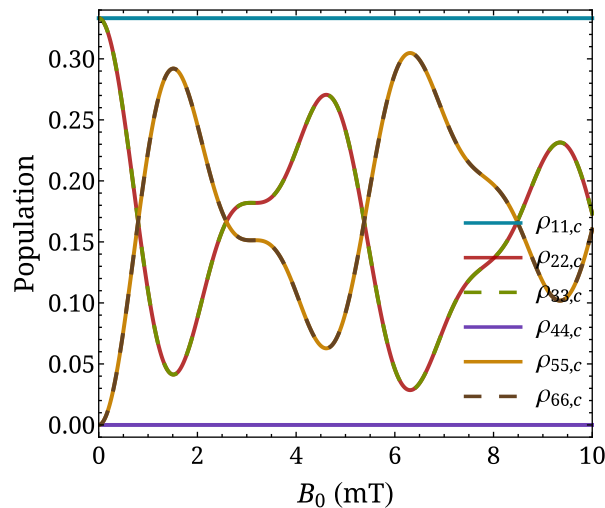
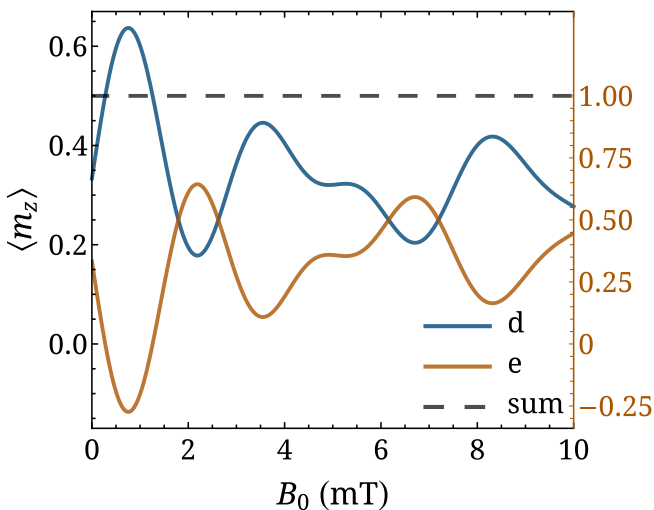


■ Deuterium atom

$$H = H_0 + H_B = \frac{A_{2S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S} - \left(\frac{g_S \mu_B}{\hbar} \mathbf{S} + \frac{g_I \mu_N}{\hbar} \mathbf{I} \right) \cdot \mathbf{B}$$

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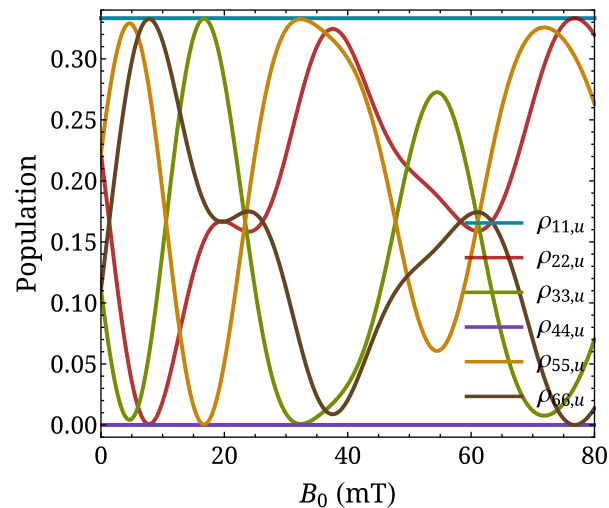
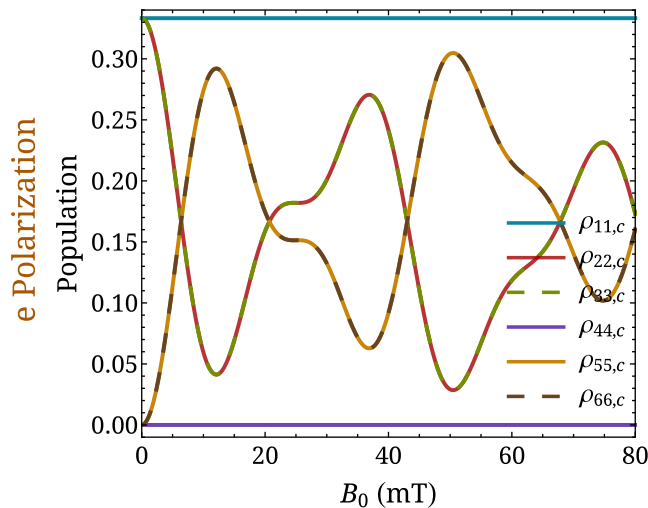
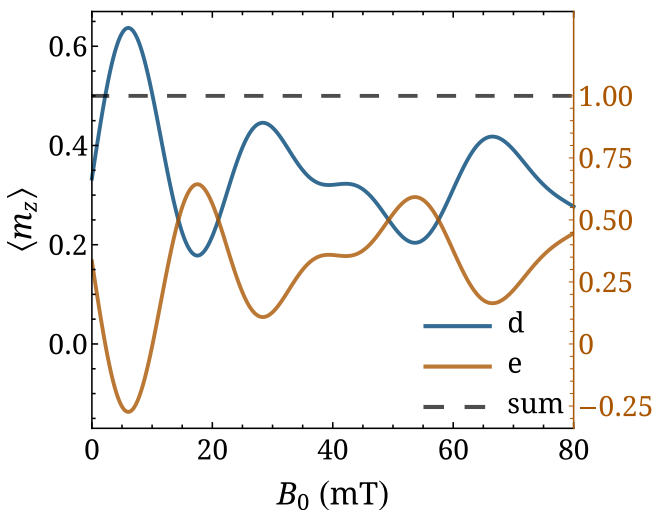


Incoherent initial preparation

- Deuterium atom
- Ground state

$$H = H_0 + H_B = \frac{A_{1S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S} - \left(\frac{g_S \mu_B}{\hbar} \mathbf{S} + \frac{g_I \mu_N}{\hbar} \mathbf{I} \right) \cdot \mathbf{B}$$

$$H_0 = \frac{A_{1S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S} \quad \text{with} \quad \frac{A_{1S}}{h} = 218.26 \text{ MHz} \quad t_f = \frac{h}{\Delta E_{hfs}} = \frac{2h}{3A_{1S}} \sim 3 \text{ ns}$$

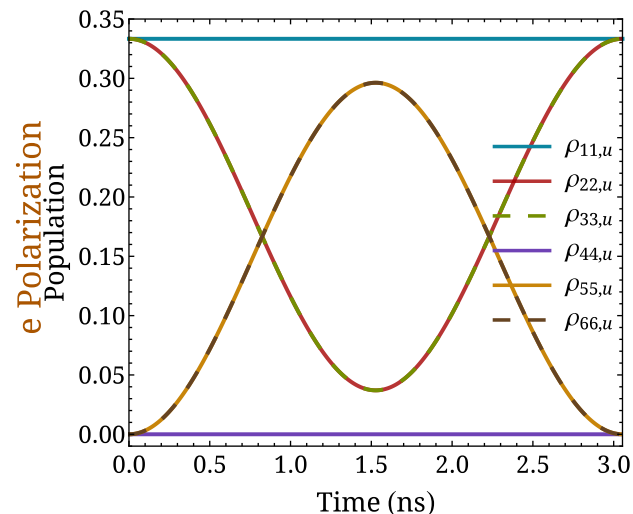
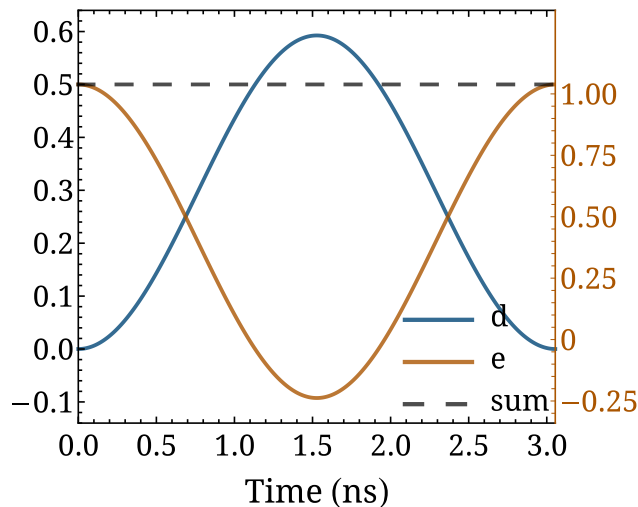


■ Deuterium atom

- Ground state: Molecular photodissociation of a deuterium halide
- Example reaction: $DY + hf \rightarrow \dots \rightarrow D \left(m_D = \pm \frac{1}{2} \right) + Y \left(m_Y = \pm \frac{1}{2} \right)$

$$\rho_u = \begin{pmatrix} 1/3 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1/3 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1/3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

$$\rho_c = \begin{pmatrix} 1/3 & 0 & 0 & 0 & 0 & 0 \\ 0 & 2/9 & 0 & 0 & 0 & \sqrt{2}/9 \\ 0 & 0 & 1/9 & 0 & \sqrt{2}/9 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & \sqrt{2}/9 & 0 & 2/9 & 0 \\ 0 & \sqrt{2}/9 & 0 & 0 & 0 & 1/9 \end{pmatrix} \langle m_z \rangle$$



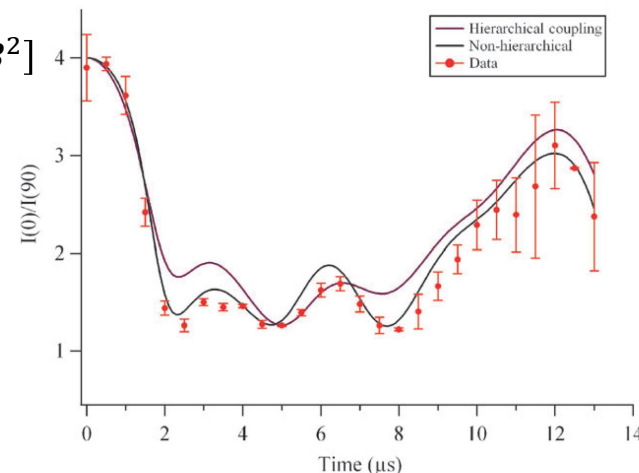
$$H_0 = \frac{A_{1S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S}$$

$$\begin{aligned}
 \blacksquare \quad H_0 &= -\frac{c_p}{\hbar^2} \mathbf{I}_p \cdot \mathbf{J} - \frac{c_d}{\hbar^2} \mathbf{I}_d \cdot \mathbf{J} + \frac{\delta}{\hbar^2} \mathbf{I}_p \cdot \mathbf{I}_d \\
 &+ \frac{5d_1}{(2J-1)(2J+3)\hbar^4} \left[\frac{3}{2} (\mathbf{I}_p \cdot \mathbf{J})(\mathbf{I}_d \cdot \mathbf{J}) + \frac{3}{2} (\mathbf{I}_d \cdot \mathbf{J})(\mathbf{I}_p \cdot \mathbf{J}) - \mathbf{I}_p \cdot \mathbf{I}_d J^2 \right] \\
 &+ \frac{5d_2}{(2J-1)(2J+3)\hbar^4} \left[3(\mathbf{I}_d \cdot \mathbf{J})^2 + \frac{3\hbar^2}{2} (\mathbf{I}_d \cdot \mathbf{J}) - \mathbf{I}_d^2 J^2 \right] \\
 \blacksquare \quad H_B &= -\frac{a'_p}{\hbar} \mathbf{I}_p \cdot \mathbf{B} - \frac{a'_d}{\hbar} \mathbf{I}_d \cdot \mathbf{B} - \frac{b'}{\hbar} \mathbf{J} \cdot \mathbf{B} - \frac{5f'}{3(2J-1)(2J+3)\hbar^2} [3(\mathbf{J} \cdot \mathbf{B})^2 - J^2 \mathbf{B}^2]
 \end{aligned}$$

with

$c_p/h = 85589 \text{ Hz}$, $c_d/h = 13118 \text{ Hz}$, $\delta/h = 43 \text{ Hz}$, $d_1/h = 17764 \text{ Hz}$,
 $d_2/h = -22452 \text{ Hz}$, $a'_p/h = 4257.796 \times 10^4 \text{ Hz/T}$, $f'/h = -2630 \text{ Hz/T}^2$
 $a'_d/h = 653.5832 \times 10^4 \text{ Hz/T}$, and $b'/h = 505.5870 \times 10^4 \text{ Hz/T}$

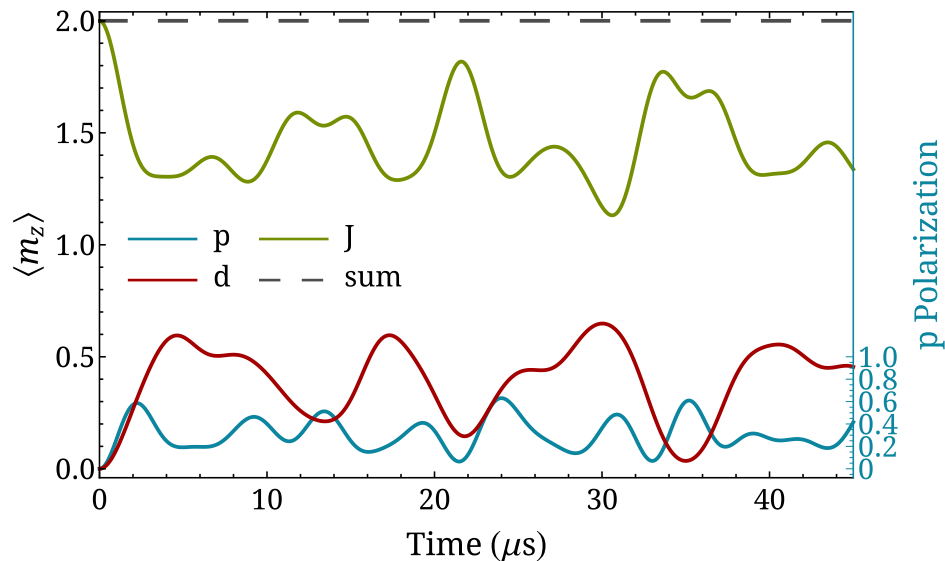
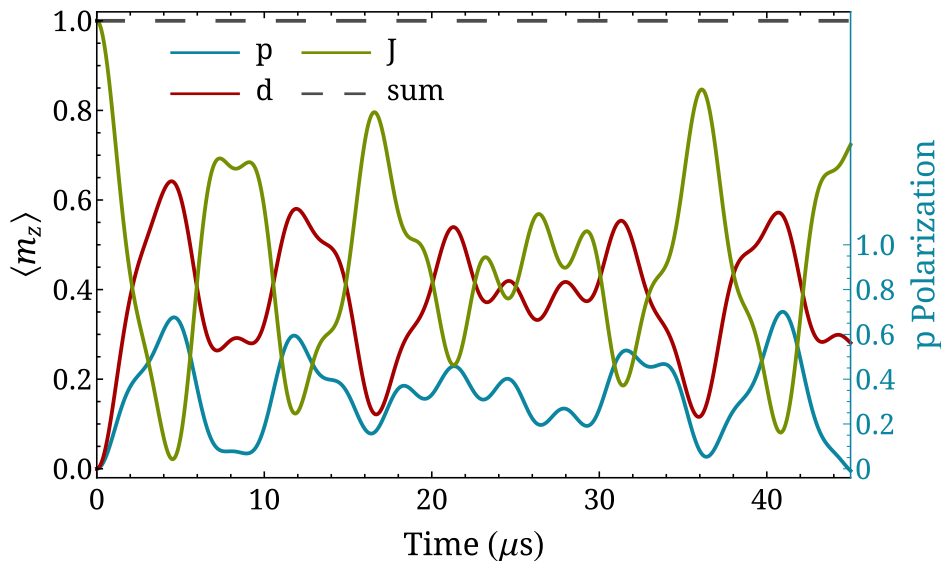
- Coherent initial preparation: $|J, m_J = +J\rangle$
- Transfer of molecular rotational polarization to nuclear spins



N. C.-M. Bartlett *et al.*, Phys. Chem. Chem. Phys. **11**, 142 (2009)

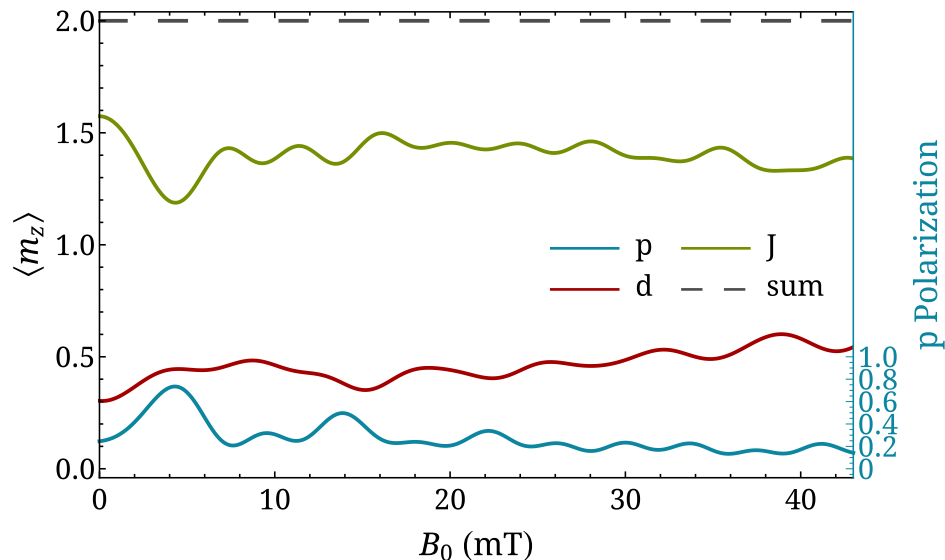
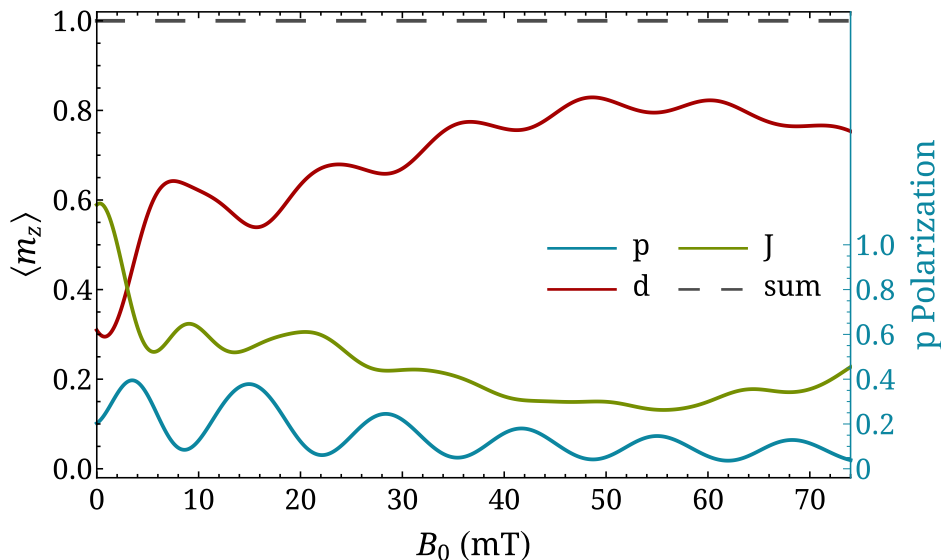
Coherently rotationally state-selected HD

- $J = 1$: proton pol. 70% at 41 μs and **deuteron** pol. **64%** at 4.5 μs
- $J = 2$: **proton** pol. **63%** at 24 μs and deuteron pol. 65% at 30 μs



Coherently rotationally state-selected HD

- $J = 1$: $t_f = \frac{1}{153935 \text{ Hz}} \sim 6.5 \text{ } \mu\text{s}$, **deuteron** pol. **83%** ($B_0 = 48.6 \text{ mT}$)
- $J = 2$: $t_f = \frac{1}{87852 \text{ Hz}} \sim 11.4 \text{ } \mu\text{s}$, **proton** pol. **74%** ($B_0 = 4.3 \text{ mT}$)



- Non-interacting particles following classical trajectories
 - External degrees of freedom (position, momentum): treated classically
 - Internal degrees of freedom (spins, rotational ang. momenta): treated quantum mechanically
 - Nonrelativistic velocities $\frac{v}{c} < 1\%$; beam kinetic energies up to: 46.9 keV (H), 93.8 keV (D), and 140.7 keV (HD)
- Radial magnetic field component B_r neglected ($r \ll \lambda$)
 - B_r induces $\Delta m_F = \pm 1$ transitions, not conserving longitudinal spin polarization
 - B_r also responsible for motional Stark effect
 - $\mathbf{E} = \mathbf{v} \times \mathbf{B} \Rightarrow \mathbf{E} = v B_r \hat{\phi}$
 - Electric dipole interaction does not couple states within the same hyperfine regime, only states of opposite parity
- Magnetic field need not be a perfect sinusoid
 - E.g., a sine-cubed waveform of the same frequency yields similar results with adjusted field strength

- Theory (H)
- Theory (D)
- Incoherent preparation of metastable H (maximum nuclear polarization)
- Coherent preparation of ground-state H (signal from dissociation of HBr at 213 nm)
- Incoherent preparation of metastable D ($2S_{1/2} - 2P_{1/2}$ level crossing)
- Coherent preparation of ground-state D (signal from dissociation of DI at 266 nm)
- Coherently rotationally state-selected HD (maximum nuclear polarization)
- Experimental setup
- Applications

Theory: coupled and uncoupled bases

■ Hydrogen atom: $\frac{1}{2} \otimes \frac{1}{2} = 1 \oplus 0$

	coupled basis $ F, m_F\rangle$	uncoupled basis $ m_S, m_L\rangle$
1.	$ 1, 1\rangle$	$ 1/2, 1/2\rangle$
2.	$ 1, 0\rangle$	$ 1/2, -1/2\rangle$
3.	$ 1, -1\rangle$	$ -1/2, -1/2\rangle$
4.	$ 0, 0\rangle$	$ -1/2, 1/2\rangle$

$$\langle m_{z,p/e} \rangle = \frac{1}{\hbar} \langle S_{z,p/e} \rangle = \frac{1}{\hbar} \text{Tr}(\rho S_{z,p/e})$$

$$S_{z,p,u} = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \text{ and } S_{z,p,c} = Q S_{z,p,u} Q^{-1} = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$|\chi\rangle_c = Q |\chi\rangle_u \text{ with } Q = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1/\sqrt{2} & 0 & 1/\sqrt{2} \\ 0 & 0 & 1 & 0 \\ 0 & 1/\sqrt{2} & 0 & -1/\sqrt{2} \end{pmatrix}$$

$$\text{For } \rho_c = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}, \rho_u = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/4 & 0 & 1/4 \\ 0 & 0 & 0 & 0 \\ 0 & 1/4 & 0 & 1/4 \end{pmatrix}$$

Theory: coupled and uncoupled bases

■ Deuterium atom: $\frac{1}{2} \otimes 1 = \frac{3}{2} \oplus \frac{1}{2}$

	coupled basis $ F, m_F\rangle$	uncoupled basis $ m_S, m_I\rangle$
1.	$ 3/2, 3/2\rangle$	$ 1/2, 1\rangle$
2.	$ 3/2, 1/2\rangle$	$ 1/2, 0\rangle$
3.	$ 3/2, -1/2\rangle$	$ 1/2, -1\rangle$
4.	$ 3/2, -3/2\rangle$	$ -1/2, -1\rangle$
5.	$ 1/2, -1/2\rangle$	$ -1/2, 0\rangle$
6.	$ 1/2, 1/2\rangle$	$ -1/2, 1\rangle$

$$\begin{aligned}\langle m_{z,d/e} \rangle &= \frac{1}{\hbar} \langle S_{z,d/e} \rangle \\ &= \frac{1}{\hbar} \text{Tr}(\rho S_{z,d/e})\end{aligned}$$

$$S_{z,d,u} = \hbar \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 & 0 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 1 \end{pmatrix}$$

$$S_{z,d,c} = Q S_{z,d,u} Q^{-1} = \hbar \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1/3 & 0 & 0 & 0 & -\sqrt{2}/3 \\ 0 & 0 & -1/3 & 0 & -\sqrt{2}/3 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 \\ 0 & 0 & -\sqrt{2}/3 & 0 & -2/3 & 0 \\ 0 & -\sqrt{2}/3 & 0 & 0 & 0 & 2/3 \end{pmatrix}$$

$$|\chi\rangle_c = Q |\chi\rangle_u \text{ with } Q = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & \sqrt{2}/3 & 0 & 0 & 0 & 1/\sqrt{3} \\ 0 & 0 & 1/\sqrt{3} & 0 & \sqrt{2}/3 & 0 \\ 0 & 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & \sqrt{2}/3 & 0 & -1/\sqrt{3} & 0 \\ 0 & 1/\sqrt{3} & 0 & 0 & 0 & -\sqrt{2}/3 \end{pmatrix}$$

$$\text{For } \rho_c = \begin{pmatrix} 1/3 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1/3 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1/3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}, \rho_u = \begin{pmatrix} 1/3 & 0 & 0 & 0 & 0 & 0 \\ 0 & 2/9 & 0 & 0 & 0 & \sqrt{2}/9 \\ 0 & 0 & 1/9 & 0 & \sqrt{2}/9 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & \sqrt{2}/9 & 0 & 2/9 & 0 \\ 0 & \sqrt{2}/9 & 0 & 0 & 0 & 1/9 \end{pmatrix}$$

Hydrogen atom

$$\rho_c = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

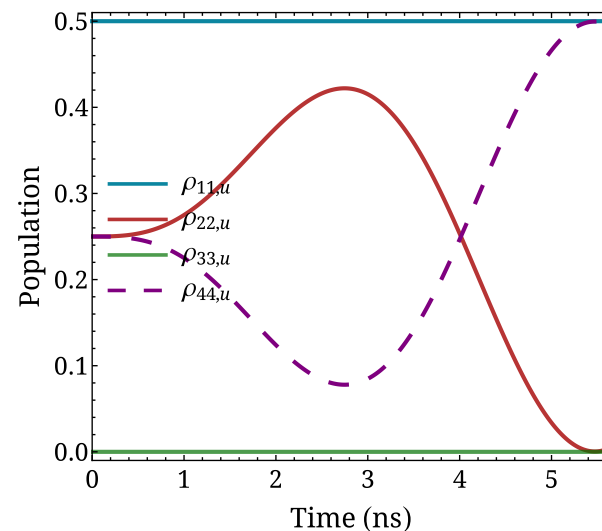
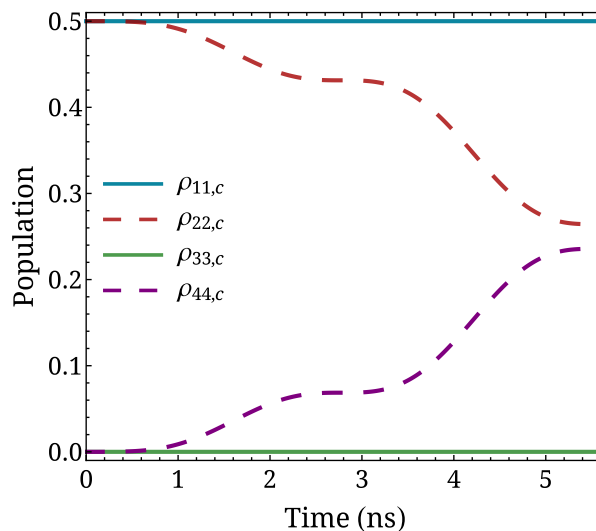
coupled basis uncoupled basis

	$ F, m_F\rangle$	$ m_S, m_I\rangle$
1.	$ 1, 1\rangle$	$ 1/2, 1/2\rangle$
2.	$ 1, 0\rangle$	$ 1/2, -1/2\rangle$
3.	$ 1, -1\rangle$	$ -1/2, -1/2\rangle$
4.	$ 0, 0\rangle$	$ -1/2, 1/2\rangle$

➤ How is population evolution affected by $B_z = B_0 \sin\left(\frac{2\pi\nu t}{\lambda}\right)$?

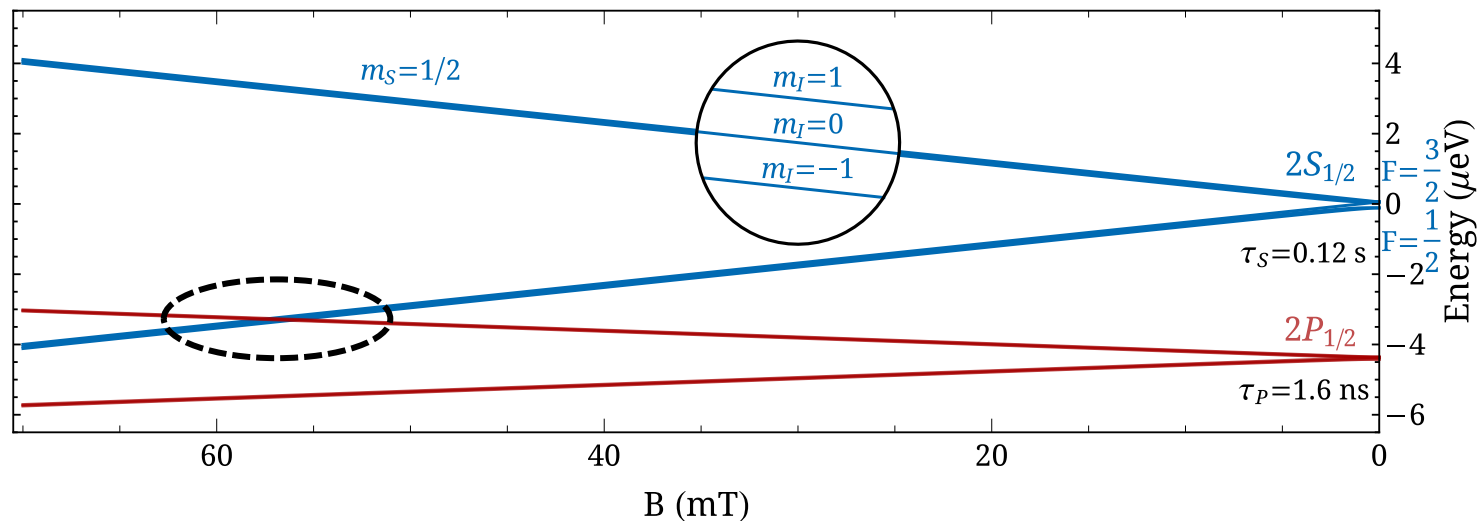
$$H = H_0 + H_B = \frac{A_{2S}}{\hbar^2} \mathbf{I} \cdot \mathbf{S} - \left(\frac{g_S \mu_B}{\hbar} \mathbf{S} + \frac{g_I \mu_N}{\hbar} \mathbf{I} \right) \cdot \mathbf{B}$$

$B_0 = 3.12 \text{ mT}$



■ Deuterium atom

- Ground state: similar to hydrogen
- Metastable state: use of the $2S_{1/2} - 2P_{1/2}$ level crossing (at high magnetic field)

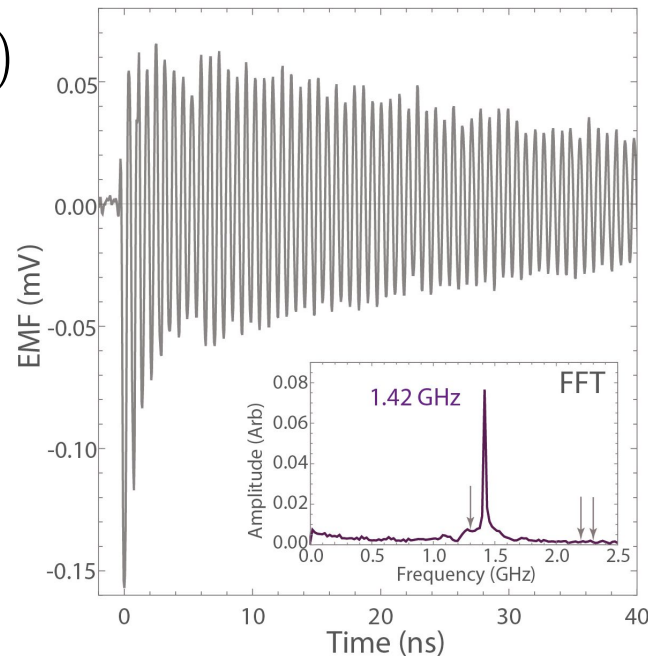
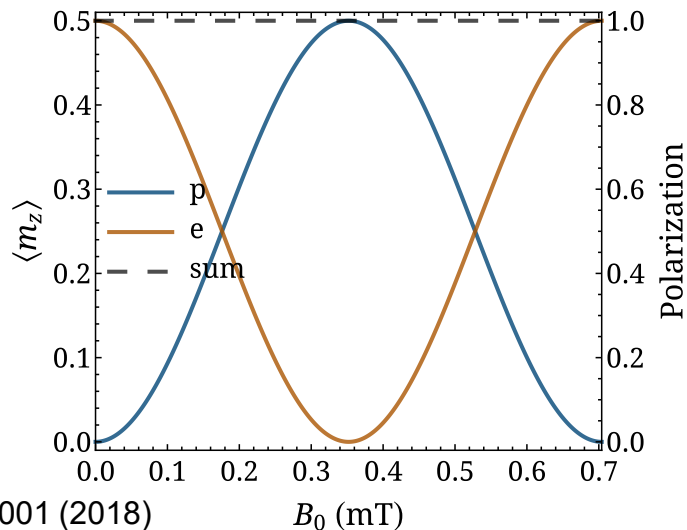


Hydrogen atom

- Ground state: Molecular photodissociation of a hydrogen halide
- Example reaction: $\text{HY} + hf \rightarrow \dots \rightarrow \text{H} \left(m_{\text{H}} = \pm \frac{1}{2} \right) + \text{Y} \left(m_{\text{Y}} = \pm \frac{1}{2} \right)$

$$\rho_u = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix}$$

$$\rho_c = \begin{pmatrix} 1/2 & 0 & 0 & 0 \\ 0 & 1/4 & 0 & 1/4 \\ 0 & 0 & 0 & 0 \\ 0 & 1/4 & 0 & 1/4 \end{pmatrix}$$



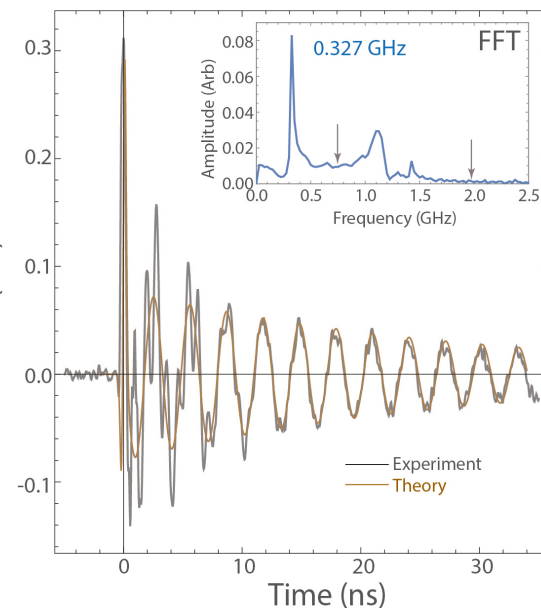
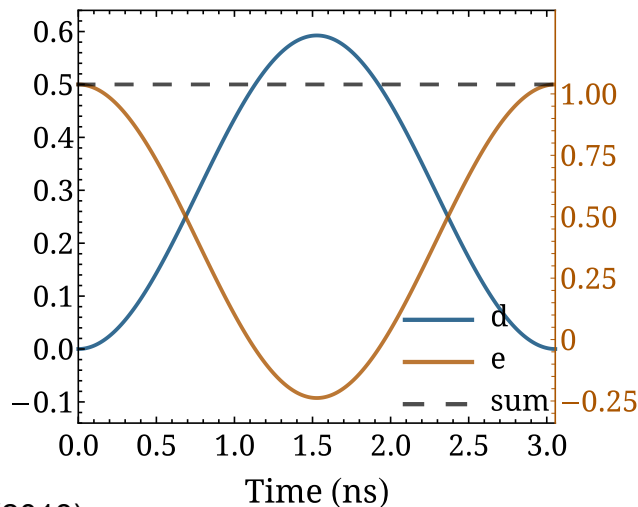
■ Deuterium atom

■ Ground state: Molecular photodissociation of a deuterium halide

■ Example reaction: $DY + hf \rightarrow \dots \rightarrow D \left(m_D = \pm \frac{1}{2} \right) + Y \left(m_Y = \pm \frac{1}{2} \right)$

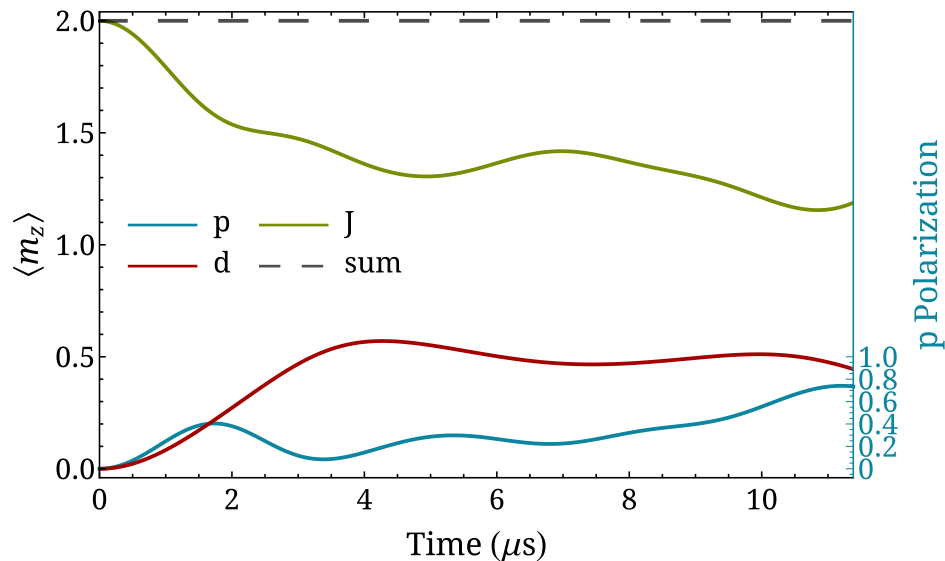
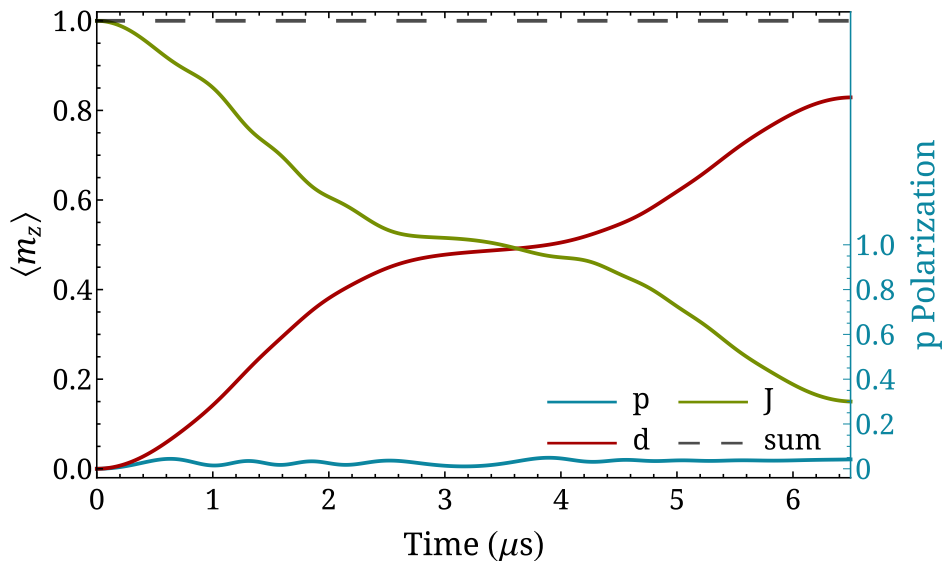
$$\rho_u = \begin{pmatrix} 1/3 & 0 & 0 & 0 & 0 & 0 \\ 0 & 1/3 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1/3 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{pmatrix}$$

$$\rho_c = \begin{pmatrix} 1/3 & 0 & 0 & 0 & 0 & 0 \\ 0 & 2/9 & 0 & 0 & 0 & \sqrt{2}/9 \\ 0 & 0 & 1/9 & 0 & \sqrt{2}/9 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & \sqrt{2}/9 & 0 & 2/9 & 0 \\ 0 & \sqrt{2}/9 & 0 & 0 & 0 & 1/9 \end{pmatrix} \langle m_z \rangle$$

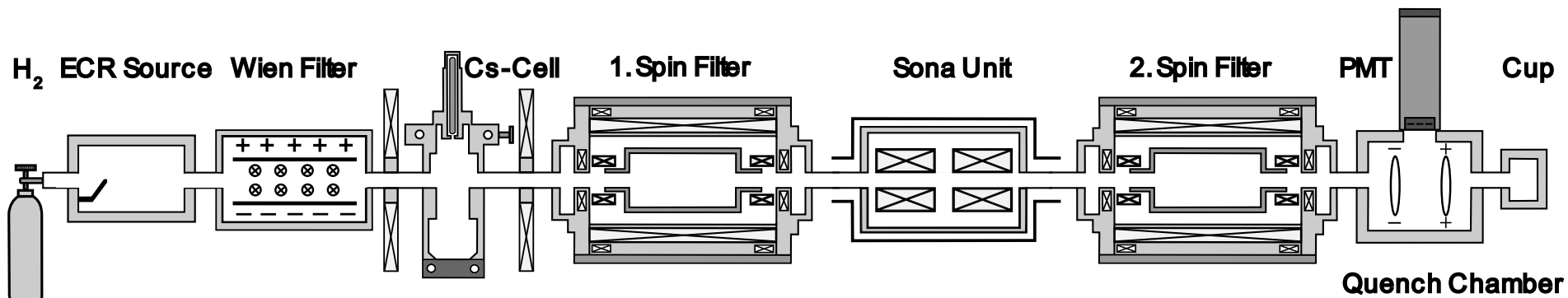


Coherently rotationally state-selected HD

- $J = 1: t_f = \frac{1}{153935 \text{ Hz}} \sim 6.5 \mu\text{s}$, **deuteron** pol. **83%** ($B_0 = 48.6 \text{ mT}$)
- $J = 2: t_f = \frac{1}{87852 \text{ Hz}} \sim 11.4 \mu\text{s}$, **proton** pol. **74%** ($B_0 = 4.3 \text{ mT}$)



Experimental setup



■ Experimental implementation

- Enhances nuclear spin polarization
 - Atoms: $t_f \sim ns$, $\lambda \sim mm$, $B_0 \sim mT$ (nonconventional wire coils)
 - Molecules: $t_f \sim \mu s$, $\lambda \sim m$, $B_0 \sim mT$ (conventional wire coils)
- Induces transitions within the hyperfine regime
 - B_r must be comparable to B_z [R. Engels *et al.*, Eur. Phys. J. D **75**, 257 (2021)]

■ Computational modeling

- Numerical solution of spin dynamics for systems with interacting angular momenta, assuming no entanglement with external degrees of freedom
- Handles arbitrarily varying magnetic fields (not limited to zero-crossings or direction changes)
- Enables Fourier analysis of time-domain observables in periodic fields