

NUCLEON STRUCTURE FROM BASIS LIGHT-FRONT QUANTIZATION : STATUS AND PROSPECTS



Chandan Mondal

Institute of Modern Physics, Chinese Academy of Sciences, Huizhou, China

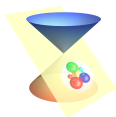


26th International
Symposium on Spin Physics
A Century of Spin

BLFQ Collaboration

**Siqi Xu, Yiping Liu, Sreeraj Nair, Jiangshan Lan
Xingbo Zhao (IMPCAS, China),
Yang Li (USTC, China) and James P. Vary (ISU, USA)**

Qingdao, China, September 23, 2025



Overview



Introduction

Basis Light-Front Quantization (BLFQ)

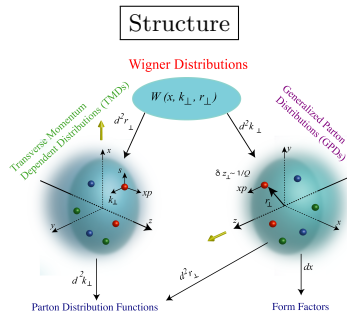
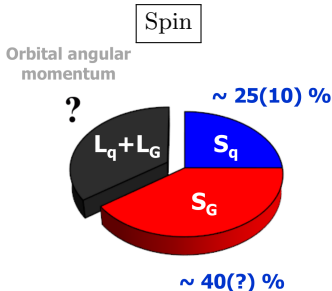
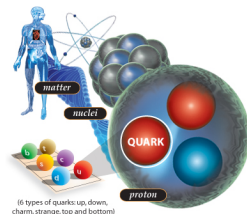
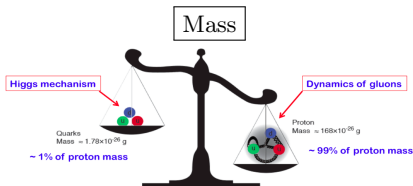
Properties of proton

Conclusions

(Satvir Kaur, 9/23/25, 2:25 PM : Hidden-Color Effects in Deuteron Structure)

(Xingbo Zhao, 9/24/25, 8:30 AM : Proton Spin Structure from a Light-Front Hamiltonian Approach)

Fundamental Properties: Mass, Spin and Structure



Ideal facilities : EIC & EicC

Nonperturbative Approaches



Lagrangian formalism

Euclidean space-time

correlators: $\langle \mathcal{O}(x_1, \dots, x_n) \rangle$

$$\langle \mathcal{O} \rangle = \int \mathcal{D}\psi \mathcal{O} \exp \left(-S_E[\psi] \right)$$

—

Lattice QCD, Dyson-Schwinger, FRG

Hamiltonian formalism

Minkowski space-time

wave functions: $|\psi_h\rangle$

$$H|\psi_h\rangle = E_h|\psi_h\rangle$$

$$i \frac{\partial}{\partial t} |\psi_h(t)\rangle = H |\psi_h(t)\rangle$$

DLCQ, BLFQ, Tamm-Dancoff, RGPEP



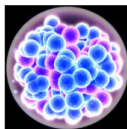
Schrödinger Equation $H|\psi\rangle = E|\psi\rangle$



nonrelativistic

few-body

quant, AMO



nonrelativistic

many-body

nucl, quant chem



relativistic

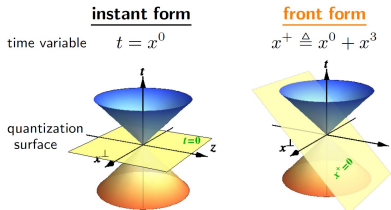
many-body

hadron, QFT

¹ Adopted from Yang Li

Basis Light-Front Quantization (BLFQ)

A computational framework for solving relativistic many-body bound state problems in quantum field theories



$$P^- P^+ |\Psi\rangle = M^2 |\Psi\rangle$$

- $P^\pm \equiv P^0 \pm P^3$: light-front momentum (Hamiltonian)
- First-principle / effective Hamiltonian as input
- Access to mass (M) and LFWFs

- Evaluate observables

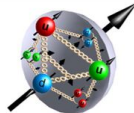
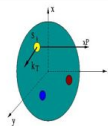
$$O \sim \langle \Psi | \hat{O} | \Psi \rangle$$

GOAL

Light-front wave-functions

3D imaging

Proton spin



FFs

GPDs

TMDs...

¹ Vary, Honkanen, Li, Maris, Brodsky, Harindranath, *et. al.*, Phys. Rev. C 81, 035205 (2010).



- Fock expansion of baryonic bound states:

$$|\text{Proton}\rangle = \psi_{(3q)}|qqq\rangle + \psi_{(3q+1g)}|qqqg\rangle + \psi_{(3q+q\bar{q})}|qqqq\bar{q}\rangle + \dots,$$

Solution proposed by BLFQ

Discrete basis and their direct product

Truncation

2D HO $\phi_{nm}(p^\perp)$ in the transverse plane

$$\sum_i (2n_i + |m_i| + 1) \leq N_{\max}$$

Plane-wave in the longitudinal direction

$$\sum_i k_i = K, \quad x_i = \frac{k_i}{K}$$

Light-front helicity state for spin d.o.f.

$$\sum_i (m_i + \lambda_i) = M_J$$

$$\alpha_i = (k_i, n_i, m_i, \lambda_i)$$

$$|\alpha\rangle = \otimes_i |\alpha_i\rangle$$

Fock sector truncation

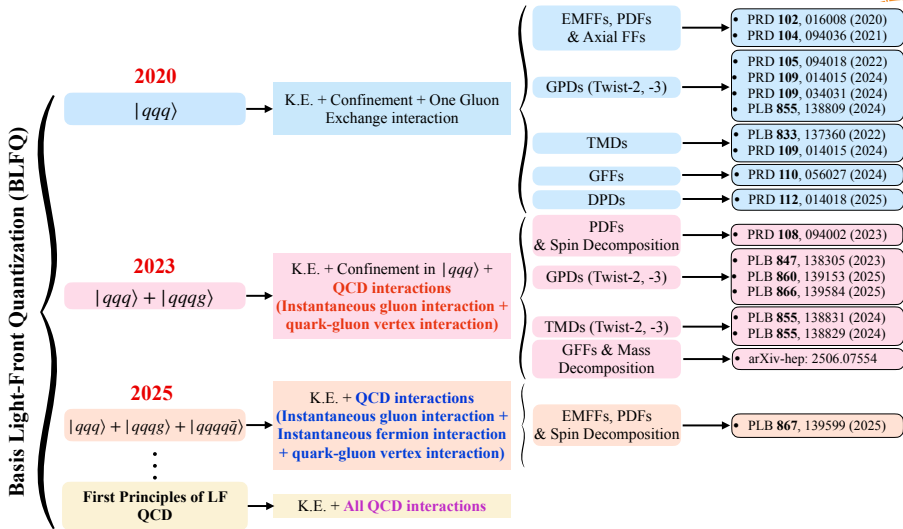
$$\Lambda_{\text{IR}} \sim \frac{b}{\sqrt{N_{\max}}} \quad \text{and} \quad \Lambda_{\text{UV}} \sim b\sqrt{N_{\max}}, \quad (b : \text{HO energy scale})$$

Large $N_{\max}$ and $K \rightarrow$ High UV cutoff & low IR cutoff

- Exact factorization between center-of-mass motion and intrinsic motion

¹Vary, Honkanen, Li, Maris, Brodsky, Harindranath, *et. al.*, Phys. Rev. C 81, 035205 (2010).

BLFQ Road-map toward First Principles



Proton with One Dynamical Gluon



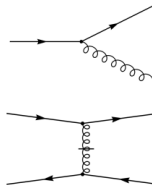
$$P^+ P^- |\Psi\rangle = M^2 |\Psi\rangle$$

$$|\text{proton}\rangle = \psi_{uud} |uud\rangle + \psi_{uudg} |uudg\rangle$$

QCD Interaction:

$$P^- = P_{\text{QCD}}^- + P_C^-$$

$$P_{\text{QCD}}^- = \int dx^- d^2x^\perp \left\{ \frac{1}{2} \bar{\psi} \gamma^+ \frac{m_0^2 + (i\partial^\perp)^2}{i\partial^+} \psi \right. \\ \left. - \frac{1}{2} A_a^i [m_g^2 + (i\partial^\perp)^2] A_a^i + g_s \bar{\psi} \gamma_\mu T^a A_a^\mu \psi \right. \\ \left. + \frac{1}{2} g_s^2 \bar{\psi} \gamma^+ T^a \psi \frac{1}{(i\partial^+)^2} \bar{\psi} \gamma^+ T^a \psi \right\},$$



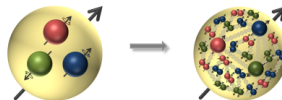
Confinement only in leading Fock:

$$P_C^- P^+ = \frac{\kappa^4}{2} \sum_{i \neq j} \left\{ \{ \vec{r}_{ij\perp}^2 - \frac{\partial_{x_i}(x_i x_j \partial_{x_j})}{(m_i + m_j)^2} \} \right\}$$

Parameters:

Truncation: Nmax=9, K=16.5

HO parameters: b=0.7GeV, b_{inst}=3GeV



m_u	m_d	m_g	κ	m_f	g
0.31GeV	0.25GeV	0.50GeV	0.54GeV	1.80GeV	2.40

¹ S. Xu, CM, X. Zhao, Y. Li, J. P. Vary, Phys.Rev.D 108 (2023) 094002.

² Brodsky, Teramond, Dosch and Erlich, Phys. Rep. 584, 1 (2015).

³ Li, Maris, Zhao and Vary, Phys. Lett. B (2016); M. Burkardt, Phys. Rev. D 58, 096015 (1998).

Unpolarized PDFs

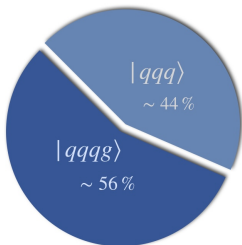
Fock expansion:

$$|\text{Proton}\rangle = a |uud\rangle + b |uudg\rangle + \dots$$

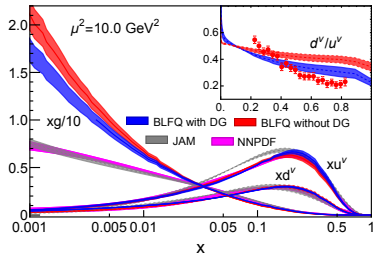
Light-front effective Hamiltonian :

$$H_{\text{eff}} = \sum_a \frac{\vec{p}_{\perp a}^2 + m_a^2}{x_a} + H_{\text{confinement}} + H_{\text{vertex}} + H_{\text{inst}}$$

Fock Sector Decomposition



$$\text{Model scale } \mu_0^2 = 0.24 \pm 0.01 \text{ GeV}^2$$



Diagonalizing $H_{\text{eff}} \Rightarrow$ LF wavefunction
 $\Rightarrow$ Initial PDFs $\Rightarrow$ Scale evolution

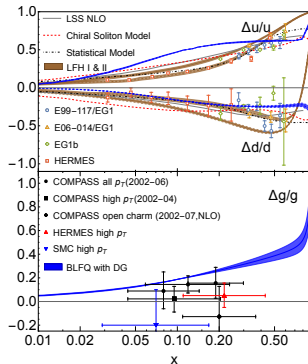
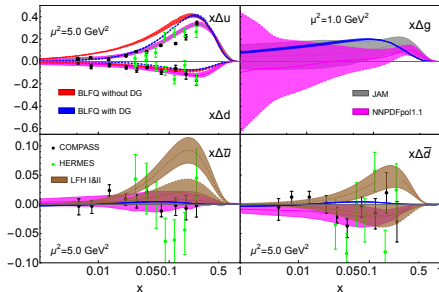
- $\langle x \rangle_u = 0.261 \pm 0.005$,
 $\langle x \rangle_d = 0.109 \pm 0.005$ at 10 GeV^2 .

¹ BLFQ Collaboration, Phys.Rev.D 108 (2023) 094002.





Helicity PDFs



- **Quark spin:** $\frac{1}{2}\Sigma_u = 0.438 \pm 0.004$, $\frac{1}{2}\Delta\Sigma_d = -0.080 \pm 0.002$.
- **Gluon spin:** $\Delta G = 0.131 \pm 0.003$, (PHENIX: $\Delta G^{[0.02,0.3]} = 0.2 \pm 0.1$).

¹ BLFQ: PRD 108 (2023) 094002; LFH: PRL 124 (2020), 082003; PHENIX: PRL 103 (2009) 012003.

BLFQ Results for Spin Decomposition

Fock expansion:

$$| \text{Proton} \rangle = a | uud \rangle + b | uudg \rangle + \dots$$



Quark and gluon helicities :

$$\Delta \Sigma_q = \int dx \Delta q(x)$$

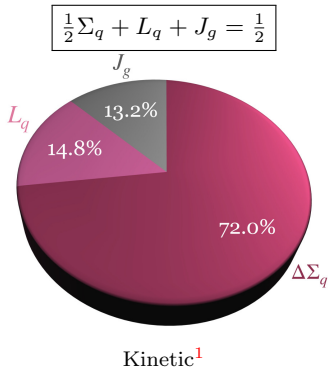
$$\Delta \Sigma_g = \int dx \Delta G(x)$$

Total AM² :

$$J_i = \frac{1}{2} \int dx x [H_i(x, 0, 0) + E_i(x, 0, 0)]$$

Kinetic OAM² :

$$L_q = \frac{1}{2} \int dx [x \{ H_q(x, 0, 0) + E_q(x, 0, 0) \} - \tilde{H}_q(x, 0, 0)]$$



¹S. Xu, CM, X. Zhao, Y. Li, J. P. Vary, Phys.Rev.D 108, 094002 (2023).

²X. Ji, Phys.Rev.Lett. 78, 610 (1997).

BLFQ Results for Spin Decomposition

Fock expansion:

$$| \text{Proton} \rangle = a | uud \rangle + b | uudg \rangle + \dots$$



Quark and gluon helicities :

$$\Delta\Sigma_q = \int dx \Delta q(x)$$

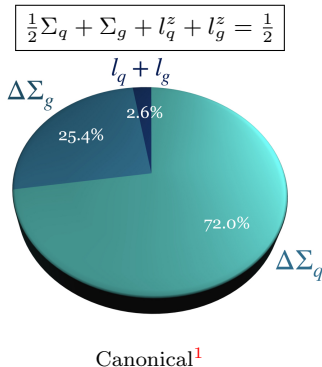
$$\Delta\Sigma_g = \int dx \Delta G(x)$$

Total AM :

$$J_i = \frac{1}{2} \int dx x [H_i(x, 0, 0) + E_i(x, 0, 0)]$$

Canonical OAM :

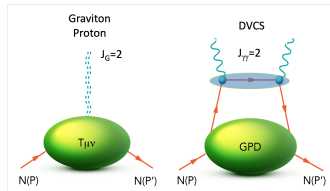
$$l_i^z = - \int dx d^2\vec{p}_\perp \frac{\vec{p}_\perp^2}{M^2} F_{1,4}^i(x, 0, \vec{p}_\perp^2, 0, 0)$$



¹ S. Xu, CM, X. Zhao, Y. Li, J. P. Vary, Phys.Rev.D 108, 094002 (2023).

² Jaffe and Manohar, Nucl.Phys.B 337, 509–546 (1990).

Proton Gravitational Form Factors



- Parametrization of matrix element in terms of GFFs

$$\begin{aligned} \langle P' | T_i^{\mu\nu}(0) | P \rangle = \bar{U}' \bigg[& -B_i(q^2) \frac{\bar{P}^\mu \bar{P}^\nu}{M} + (A_i(q^2) + B_i(q^2)) \frac{1}{2} (\gamma^\mu \bar{P}^\nu + \gamma^\nu \bar{P}^\mu) \\ & + C_i(q^2) \frac{q^\mu q^\nu - q^2 g^{\mu\nu}}{M} + \bar{C}_i(q^2) M g^{\mu\nu} \bigg] U \end{aligned}$$

- Momentum sum rule : $\sum_i A^i(0) = 1$
- Gravitomagnetic moment sum rule : $\sum_i B^i(0) = 0$
- Spin sum rule: $J^i = \frac{1}{2} [A^i(0) + B^i(0)]$
- $4C(q^2) = D(q^2)$ provides shear forces and the pressure distributions

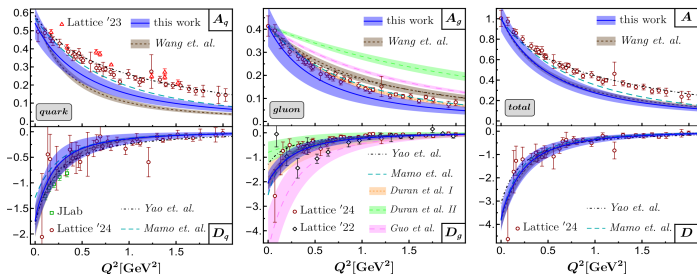
¹Burkert *et. al.*: Rev. Mod. Phys. 95, 041002 (2023)

²Ji, Phys. Rev. Lett. 78, 610 (1997)



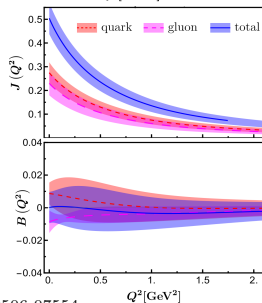
BLFQ Results for Proton GFFs

$$|\text{Nucleon}\rangle = \psi_{(3q)}|qqq\rangle + \psi_{(3q+1g)}|qqqg\rangle$$



- QCD EMT : $T^{\mu\nu} = T_q^{\mu\nu} + T_g^{\mu\nu}$,
where $T_q^{\mu\nu} = i\bar{\psi}\gamma^{\{\mu}\mathcal{D}^{\nu\}}\psi$ and
 $T_g^{\mu\nu} = -F^{\mu\lambda a}F_{\lambda a}^{\nu} + \frac{1}{4}g^{\mu\nu}(F_{\lambda\sigma a})^2$
- $A(Q^2)$ and $B(Q^2)$: T^{++} component
- $D(Q^2) = 4C(Q^2)$: T^{ij} components

$$\sum_i A^i(0) = 1 \text{ and } \sum_i B^i(0) = 0$$



¹ S. Nair, CM, Siqi Xu, Xingbo Zhao, James P. Vary, arXiv: 2506.07554

Effective Hamiltonian with Dynamical Gluon and Sea Quarks

Fock expansion:

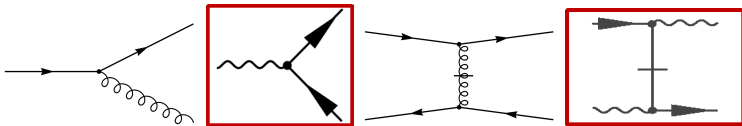
$$| \text{Proton} \rangle = a | uud \rangle + b | uudg \rangle + c_1 | uud\bar{u} \rangle + c_2 | uudd\bar{d} \rangle + c_3 | uuds\bar{s} \rangle + \dots$$



Light-front QCD Hamiltonian :

$$H_{\text{LF}} = \sum_a \frac{\vec{p}_{\perp a}^2 + m_a^2}{x_a} + \cancel{H_{\text{confinement}}} + H_{\text{vertex}} + H_{\text{inst}}$$

$$H_{\text{vertex}} + H_{\text{inst}} = g_s \bar{\psi} \gamma_\mu T^a A_\mu^\mu \psi + \frac{1}{2} g_s^2 \bar{\psi} \gamma^+ T^a \psi \frac{1}{(i\partial^+)^2} \bar{\psi} \gamma^+ T^a \psi \\ + \frac{1}{2} g_s^2 \bar{\psi} \gamma^\mu A_\mu \frac{\gamma^+}{(i\partial^+)} A_\nu \gamma^\nu \psi$$



¹ Brodsky, Pauli, and Pinsky, Phys. Rep. 301, 299 (1998).

² Siqi Xu, Yiping Liu, CM, *et. al.*, PLB 867, 139599 (2025)

Fock Sector Decomposition

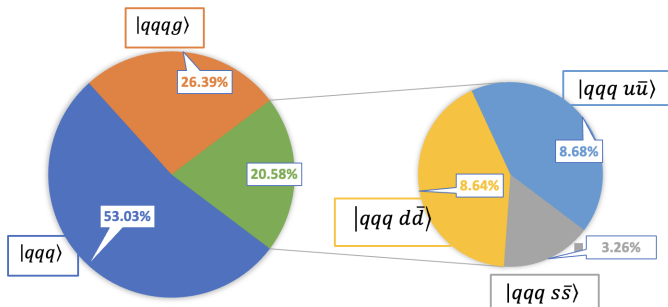
$$|P_{\text{proton}}\rangle \rightarrow |qqq\rangle + |qqqg\rangle + |qqqu\bar{u}\rangle + |qqqd\bar{d}\rangle + |qqqs\bar{s}\rangle$$

Truncation parameter: $N_{\text{max}} = 7$ and $K_{\text{max}} = 16$



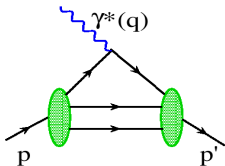
m_u	m_d	m_f	g	b	b_{inst}
0.99 GeV	0.94 GeV	5.9 GeV	3.0	0.6 GeV	2.7 GeV

In five quark Fock sector, we use current quark mass



² Siqi Xu, Yiping Liu, CM, *et. al.*, PLB 867, 139599 (2025)

Proton EM Form Factors



Sach's form factors

$$G_E(q^2) = F_1(q^2) - \frac{q^2}{4M^2} F_2(q^2),$$

$$G_M(q^2) = F_1(q^2) + F_2(q^2).$$

- EM current: $J^\mu = \bar{\psi} \gamma^\mu \psi$
- $\langle p'; \uparrow | J^+(0) | p; \uparrow (\downarrow) \rangle \sim F_{1(2)}(q^2)$
- Two FFs: $F_{1(2)}(q^2 = -Q^2)$

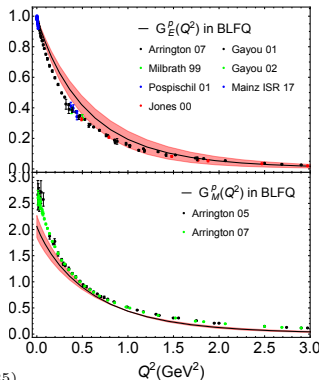
Proton Radii

$$\langle r_E^2 \rangle = -6 \frac{dG_E(Q^2)}{dQ^2} \Big|_{Q^2=0},$$

$$\langle r_M^2 \rangle = -\frac{6}{G_M(0)} \frac{dG_M(Q^2)}{dQ^2} \Big|_{Q^2=0}.$$

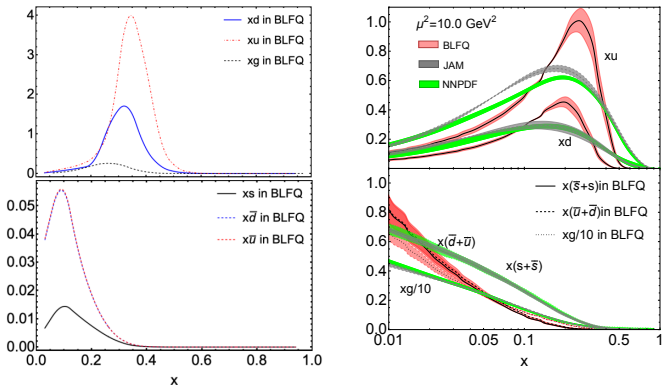
$$\sqrt{\langle r_E^2 \rangle} = 0.72 \pm 0.05 \text{ (} 0.840^{+0.003}_{-0.002} \text{) fm}$$

$$\sqrt{\langle r_M^2 \rangle} = 0.73 \pm 0.02 \text{ (} 0.849^{+0.003}_{-0.003} \text{) fm}$$



¹ Siqu Xu, Yiping Liu, CM, *et. al.*, PLB 867, 139599 (2025)

Unpolarized PDFs

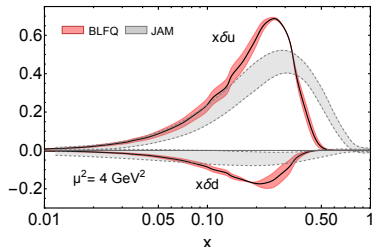
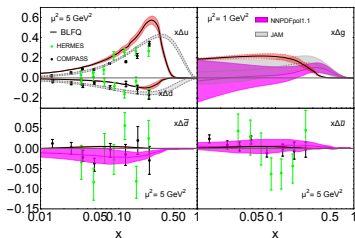


Diagonalizing $H_{\text{LFQCD}} \Rightarrow \text{LFWFs} \Rightarrow \text{Initial PDFs} \Rightarrow \text{Scale evolution}$

- Model scale $\mu_0^2 = 0.22 \pm 0.02 \text{ GeV}^2$, $\langle x \rangle_{u+d} = 0.37 \pm 0.01$ 10 GeV^2 .
- Longitudinal excitations is challenging, in absence of confining potential.

¹ Siqi Xu, Yiping Liu, CM, *et. al.*, PLB 867, 139599 (2025)

Helicity and Transversity PDFs



- Gluon spin $\Delta G = 0.29 \pm 0.03^1$ for $x_g \in [0.05, 0.2]$ at 10 GeV^2
- NNPDF² analysis: $\Delta G = 0.23(6)$; lattice QCD³: $\Delta G = 0.251(47)(16)$
- Tensor Charges: $\delta u = 0.81 \pm 0.08, \delta d = -0.22 \pm 0.01^1$
- JAM⁴ analysis: $\delta u = 0.71(2), \delta d = -0.200(6)$; lattice QCD⁵: $\delta u = 0.784(28), \delta d = -0.204(11)$.

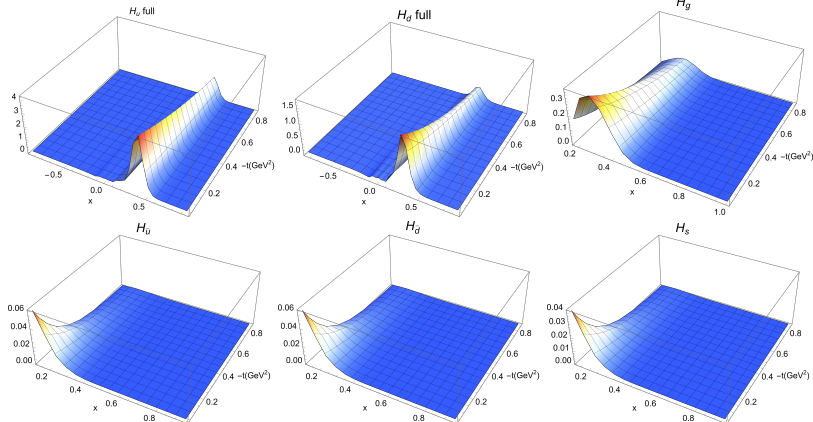
¹ Siqi Xu, Yiping Liu, CM, *et. al.*, PLB 867, 139599 (2025)

² E. R. Nocera, *et. al.* (NNPDF), Nuclear Physics B 887, 276 (2014)

³ Y.-B. Yang, *et. al.* (Lattice) Phys. Rev. Lett. 118, 102001 (2017)

⁴ C. Cocuzza, *et. al.* (JAM), Phys. Rev. Lett. 132, 091901 (2024)

⁵ R. Gupta, *et. al.* (Lattice), Phys. Rev. D 98, 091501 (2018)

GPDs at Skewness $\xi = 0.1$ 

- u and d GPDs from $-1 < x < 1$; exhibit similar distributions.
- At $\xi = 0.1$, DGLAP region dominates
- Gluon and sea quarks GPDs in DGLAP region

¹Y. Liu, S. Xu, CM, *et. al.*, coming soon...

Towards Full BLFQ

Fock expansion:

Xingbo Zhao: 9/24/25, 8:30 AM



$$|\text{Proton}\rangle = |uud\rangle + |uudg\rangle + |uudq\bar{q}\rangle + |uudgg\rangle + |uudq\bar{q}g\rangle + |uudggg\rangle$$

$$\begin{aligned} \hat{P}_{\text{LFQCD}}^- &= \frac{1}{2} \int dx^- d^2x^\perp \bar{\psi} \gamma^+ \frac{(i\partial^\perp)^2 + m^2}{i\partial^+} \psi + A^{ia} (i\partial^\perp)^2 A^{ia} \\ &+ g_s \int dx^- d^2x^\perp \bar{\psi} \gamma_\mu A^{\mu a} T^a \psi \\ &+ \frac{g_s^2}{2} \int dx^- d^2x^\perp \bar{\psi} \gamma_\mu A^{\mu a} T^a \frac{\gamma^+}{i\partial^+} (\gamma_\nu A^{\nu b} T^b \psi) \\ &+ \frac{g_s^2}{2} \int dx^- d^2x^\perp \bar{\psi} \gamma^+ T^a \psi \frac{1}{(i\partial^+)^2} (\bar{\psi} \gamma^+ T^a \psi) \\ &- g_s^2 \int dx^- d^2x^\perp i f^{abc} \bar{\psi} \gamma^+ T^c \psi \frac{1}{(i\partial^+)^2} (i\partial^+ A^{\mu a} A_\mu^b) \\ &+ g_s \int dx^- d^2x^\perp i f^{abc} i\partial^\mu A^{\nu a} A_\mu^b A_\nu^c \\ &+ \frac{g_s^2}{2} \int dx^- d^2x^\perp i f^{abc} i f^{ade} i\partial^+ A^{\mu b} A_\mu^c \frac{1}{(i\partial^+)^2} (i\partial^+ A^{\nu d} A_\nu^e) \\ &- \frac{g_s^2}{4} \int dx^- d^2x^\perp i f^{abc} i f^{ade} A^{\mu b} A^{\nu c} A_\mu^d A_\nu^e. \end{aligned}$$

¹ Siqi Xu, et. al., work in progress



Outlook

Current status

Full QCD interaction

Deuteron calculation
 $|qqq\ qqq\rangle + |qqq\ qqq\ g\rangle$

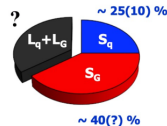
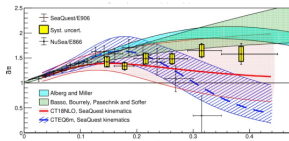
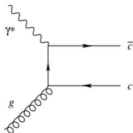
Fock sector expansion
 $|qqq\ q\bar{q}\ g\rangle$ and $|qqq\ ggg\rangle$

EMC effect

Intrinsic charm

Sea asymmetry

Origin of spin and mass



Conclusions



- BLFQ: A non-perturbative approach based on light-front QCD Hamiltonian
- **LF Hamiltonian** $\Rightarrow$ **Wavefunctions** $\Rightarrow$ **Observables**
- $|qqq\rangle + |qqqg\rangle$ ($P^- = P_{\text{QCD}}^- + P_C^-$) $\Rightarrow$ Provides good description of data/global fits and lattice QCD results for various observables
- $|qqq\rangle + |qqqg\rangle + |qqqq\bar{q}\rangle$, ($P^- = P_{\text{QCD}}^-$) $\Rightarrow$ Provides qualitative description of data/global fits for mass, spin, EMFFs, PDFs, axial and tensor charges ...

Outlook

- Expand Fock sectors Xingbo Zhao: 9/24/25, 8:30 AM
- Include four-gluon interactions in the Hamiltonian

Thank You