

# Hydrodynamic effects on spin polarization in p+Pb collisions



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**26th International Symposium on Spin Physics (SPIN2025)**  
**September 22-26, 2025, Qingdao, Shandong Province**

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Based on: Phys. Rev. C 111 (2025) 4, 044901

# Outline

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## ➤ Introduction

## ➤ Spin polarization at p+Pb collision

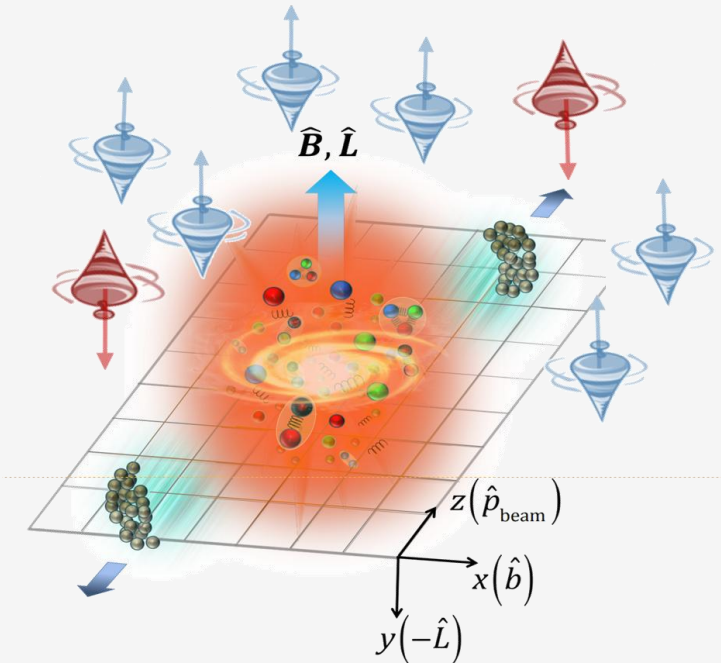
- Numerical Setup
- Results

## ➤ Summary and Outlook

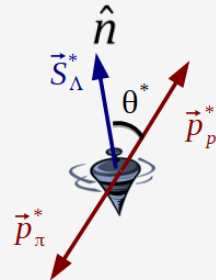
# Global Spin Polarization

## ➤ Spin-Orbital Coupling

STAR, Nature 548, 62 (2017).



Hyperons  
Spin Polarization  $S = \frac{1}{2}$



$$\Lambda \rightarrow p + \pi^-$$



$$\omega = k_B T (\overline{P}_{\Lambda'} + \overline{P}_{\bar{\Lambda}'}) / \hbar \sim 10^{22} s^{-1}$$

**Most vortical fluid !**

Liang, Wang, PRL. 94, 102301 (2005)

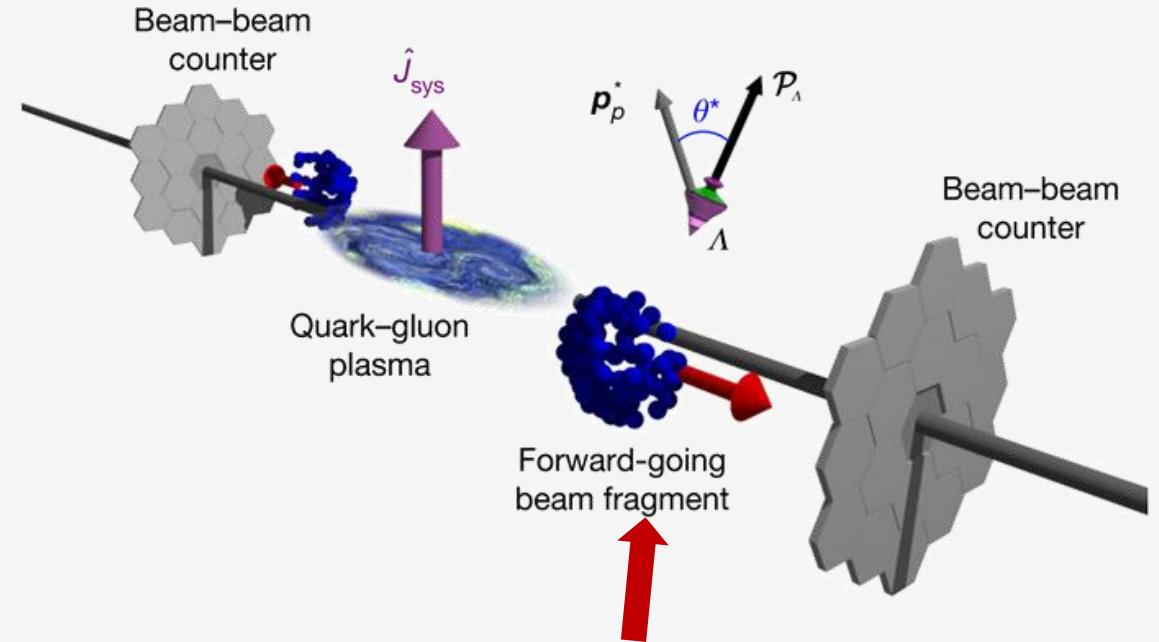
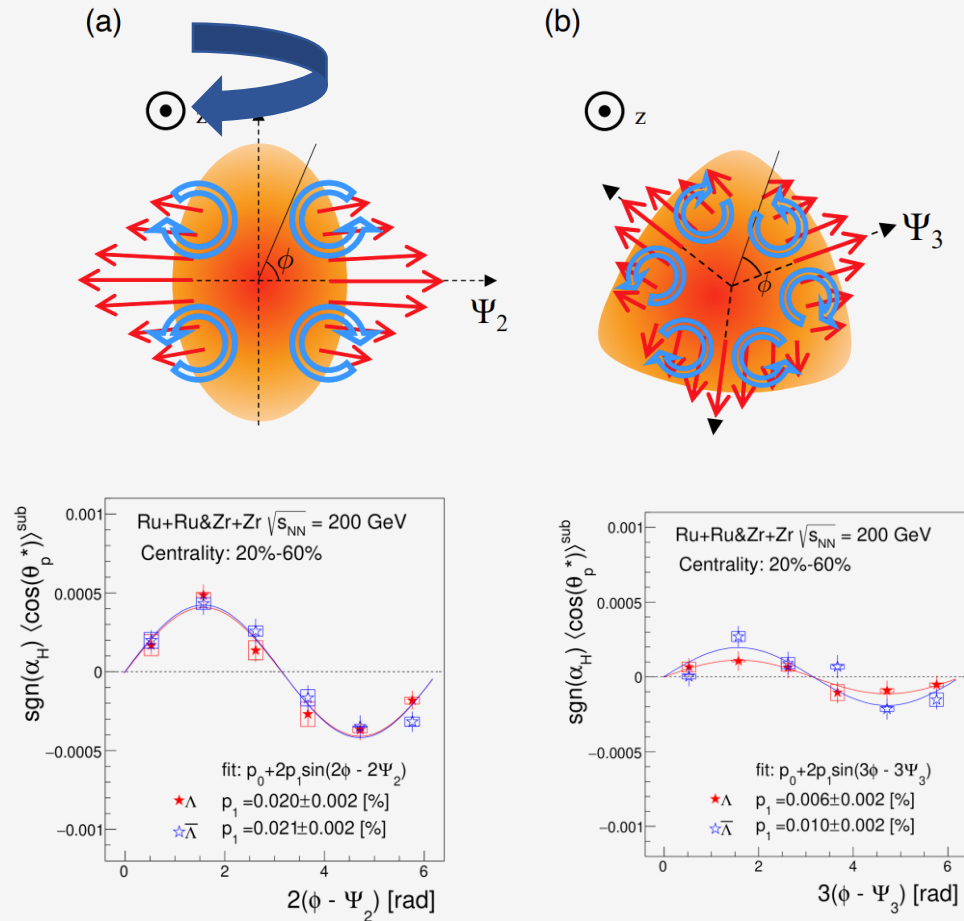
Liang, Wang, PLB 629, 20 (2005)

Becattini, Chandra, Zanna, and Grossi, Annals Phys. (2013).

Becattini and Karpenko, PRL 120, 012302

# Local Spin Polarization

## ➤ Local Vortical Structure



**Beam direction polarization**

STAR, PRL 123, 132301 (2019).  
STAR, PRL 131, 202301 (2023)

# Polarization induced by different sources

## ➤ Spin polarization at local equilibrium

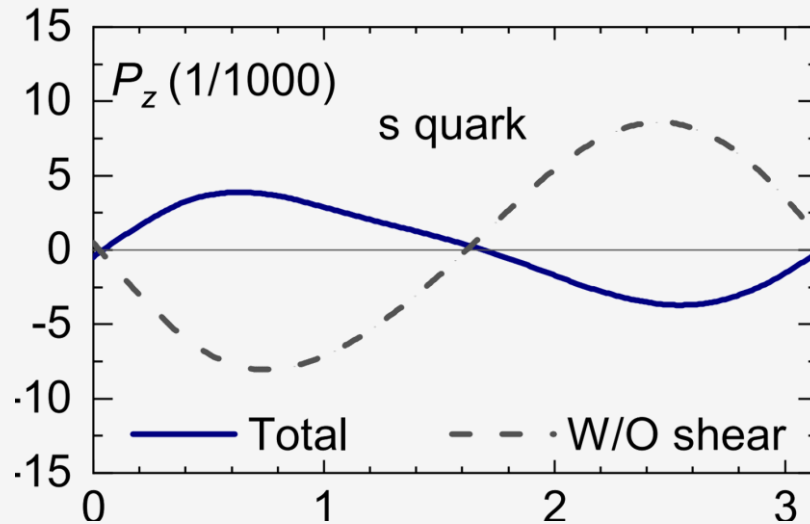
$$\mathcal{S}^\mu(\mathbf{p}) = \mathcal{S}_{\text{thermal}}^\mu + \mathcal{S}_{\text{shear}}^\mu + \mathcal{S}_{\text{accT}}^\mu + \mathcal{S}_{\text{chemical}}^\mu + \mathcal{S}_{\text{EB}}^\mu$$

$$\begin{aligned} \mathcal{S}_{\text{thermal}}^\mu(\mathbf{p}) &= \frac{\hbar}{8m_\Lambda N} \int d\Sigma^\sigma p_\sigma f_V^{(0)} (1 - f_V^{(0)}) \epsilon^{\mu\nu\alpha\beta} p_\nu \partial_\alpha \frac{u_\beta}{T}, & \text{Thermal vorticity} \\ \mathcal{S}_{\text{shear}}^\mu(\mathbf{p}) &= -\frac{\hbar}{4m_\Lambda N} \int d\Sigma \cdot p f_V^{(0)} (1 - f_V^{(0)}) \frac{\epsilon^{\mu\nu\alpha\beta} p_\alpha u_\beta}{(u \cdot p) T} \frac{1}{2} p^\sigma [(\partial_\sigma u_\nu + \partial_\nu u_\sigma) - u_\sigma D u_\nu], & \text{Shear Induced Polarization (SIP)} \\ \mathcal{S}_{\text{accT}}^\mu(\mathbf{p}) &= -\frac{\hbar}{8m_\Lambda N} \int d\Sigma \cdot p f_V^{(0)} (1 - f_V^{(0)}) \frac{1}{T} \epsilon^{\mu\nu\alpha\beta} p_\nu u_\alpha (D u_\beta - \frac{1}{T} \partial_\beta T), & \text{Fluid Acceleration} \\ \mathcal{S}_{\text{chemical}}^\mu(\mathbf{p}) &= \frac{\hbar}{4m_\Lambda N} \int d\Sigma \cdot p f_V^{(0)} (1 - f_V^{(0)}) \frac{1}{(u \cdot p)} \epsilon^{\mu\nu\alpha\beta} p_\alpha u_\beta \partial_\nu \frac{\mu}{T}, & \text{Spin Hall Effect (SHE)} \\ \mathcal{S}_{\text{EB}}^\mu(\mathbf{p}) &= \frac{\hbar}{4m_\Lambda N} \int d\Sigma \cdot p f_V^{(0)} (1 - f_V^{(0)}) \left( \frac{1}{(u \cdot p) T} \epsilon^{\mu\nu\alpha\beta} p_\alpha u_\beta E_\nu + \frac{B^\mu}{T} \right), & \text{EM Field} \end{aligned}$$

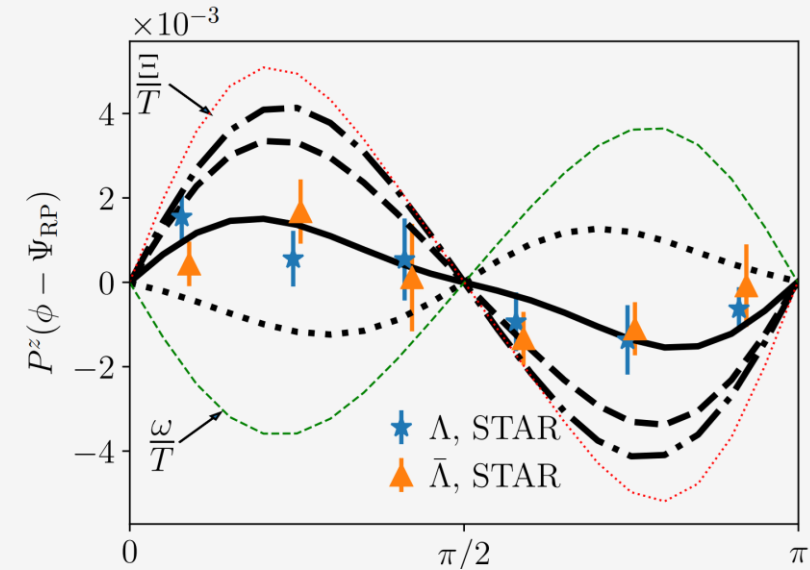
Liu, Yin, PRD 104, 054043 (2021) ; Becattini, Buzzegoli, Palermo, PLB 820, 136519 (2021) ; Liu, Yin, JHEP 07,188 (2021); Hidaka, Pu, Yang, PRD 97, 016004 (2018) ; CY, Pu, Yang, PRC 04, 064901(2021)

# Shear induced polarization(SIP)

## ➤ Hydrodynamic contribution to the local spin polarization



Fu, et al. PRL 127, 142301

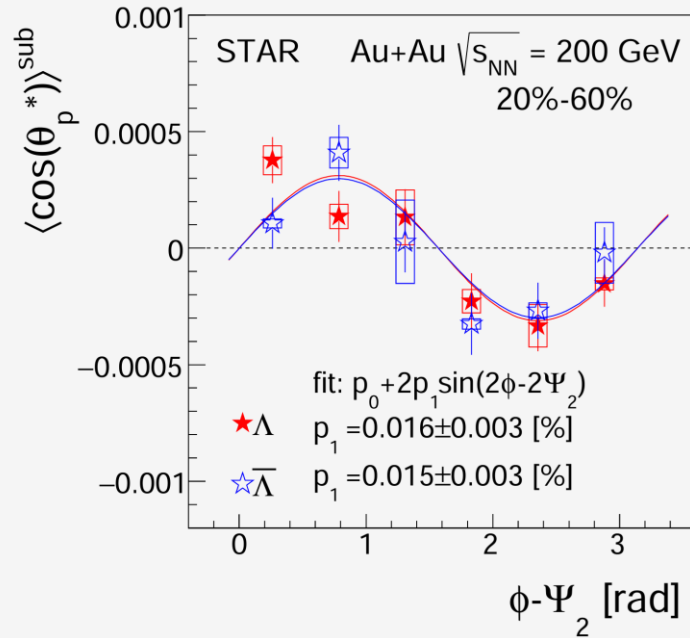


Becattini et al, PRL 127, 272302

Also see: CY, Pu, Yang, PRC 104, 064901.  
Ryu, Jupic, Shen, PRC 104, 054908  
Wu, CY, Qin, Pu, PRC 105 6, 064909  
Fu, Pang, Song, Yin, (2022), 2201.12970.  
.....

**Considering SIP under some assumptions, the theoretical calculations qualitatively/quantitatively agree with the experimental data at 200GeV Au+Au collision.**

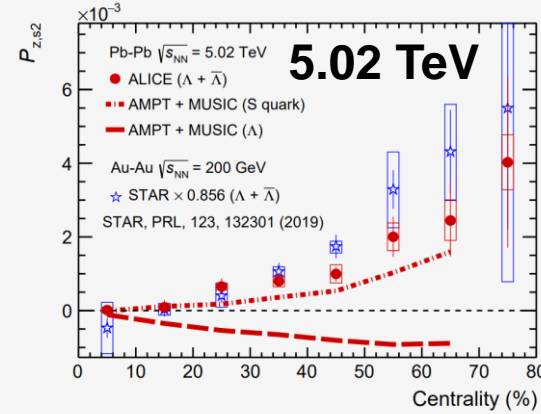
# $P_{2,z}$ across collision energy



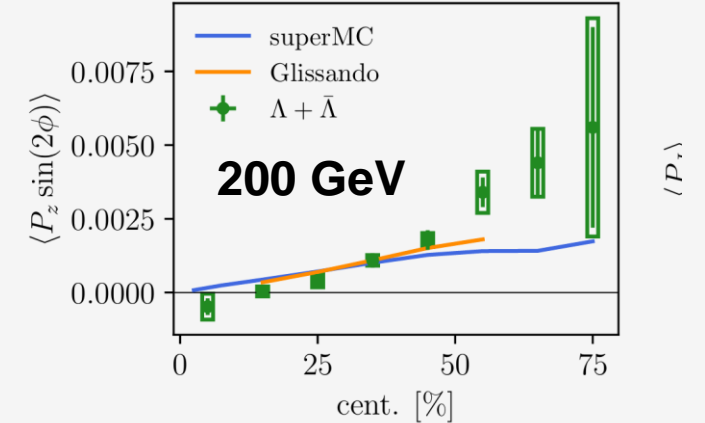
STAR, PRL 123, 132301 (2019).

$$P_{2,z} \equiv \langle P_z \sin[2(\phi - \Psi_2)] \rangle$$

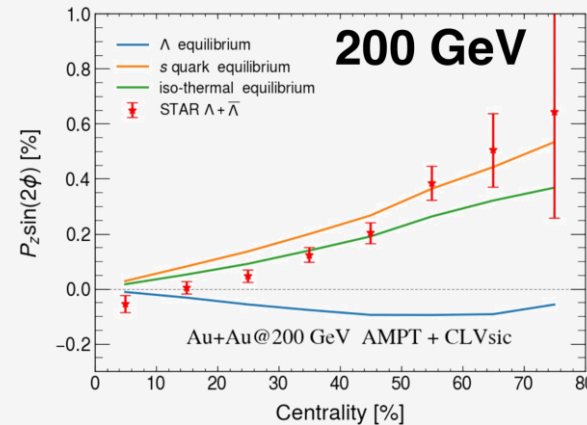
➤ How about hydrodynamic contributions to spin polarization **at p+Pb collisions?**



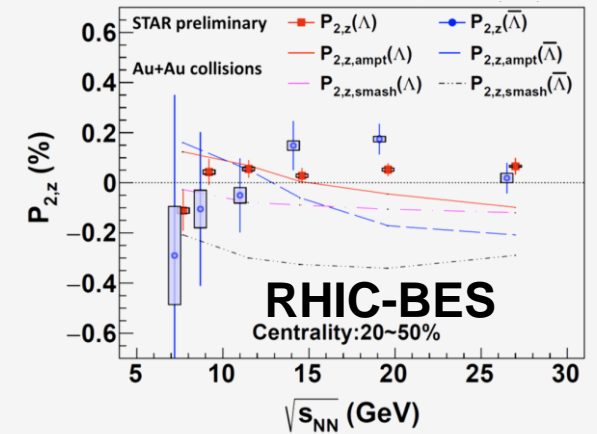
ALICE, PRL 128, 172005(2022)



Palermo, et al. EPJC 84 9, 920 (2024)



arXiv: 2509.00377



PRC 105 6, 064909 (2022)

# Outline

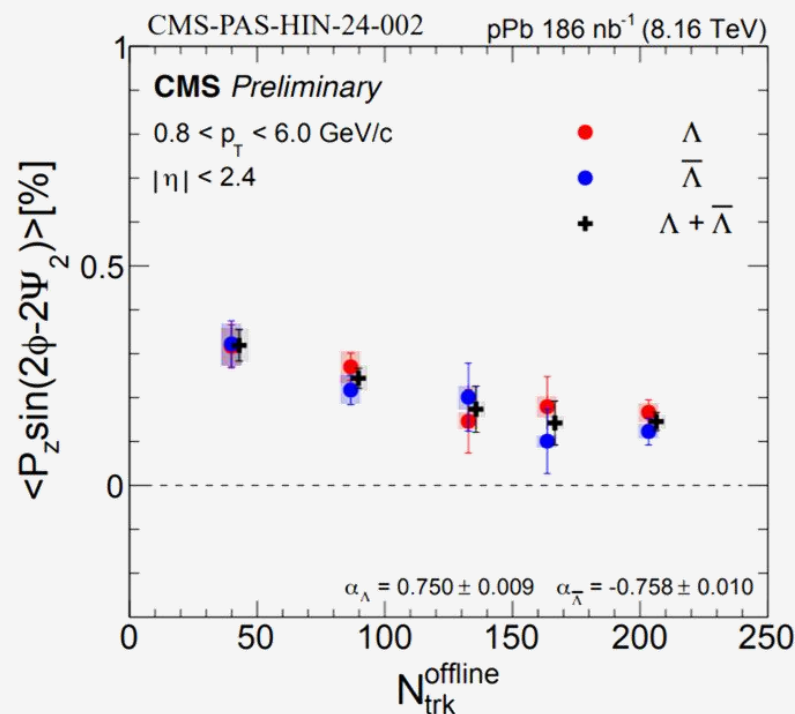
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- Introduction
- **Spin polarization at p+Pb collision**
  - Numerical Setup
  - Results
- Summary and Outlook

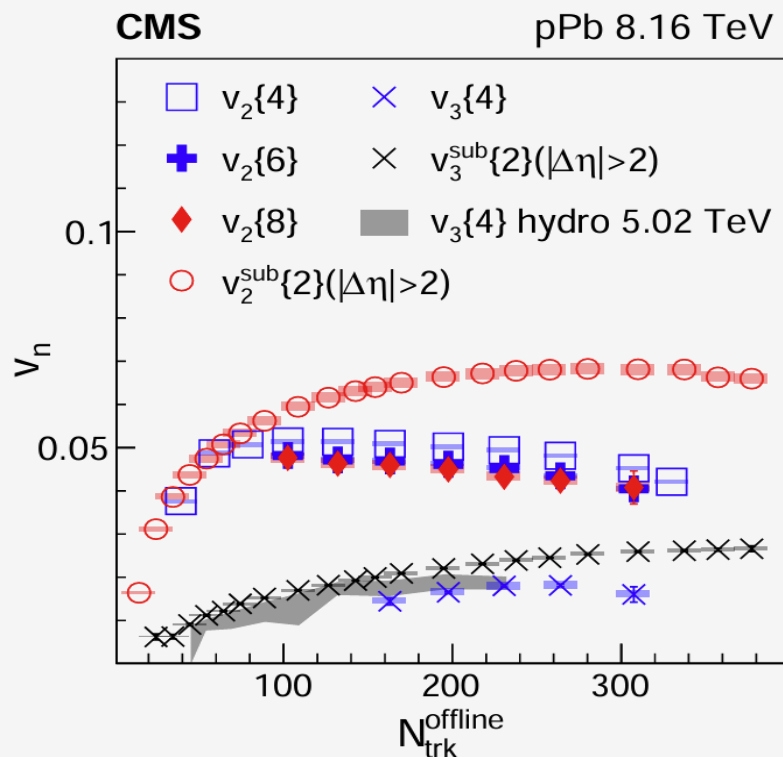


# CMS Measurements

## ➤ Polarization along the beam direction in p+Pb collisions



CMS, arXiv:2502.07898 (2025)



CMS, PRC, 101, 014912 (2020)

- The magnitude of polarization is the same order of magnitude as that in AA collisions
- Its dependence on multiplicity is inconsistent with that of  $v_2$

# Initial Condition

## ➤ Initial condition

W. Ke, J. S. Moreland, J. E. Bernhard, and S. A. Bass, PRC 96, 044912 (2017),

**We implement the parameterized TRENTo-3D model as initial conditions and consider the constituents**

$$T_{A/B}(\mathbf{x}_\perp) = \sum_{i=1}^{N_{A/B}} \frac{1}{n_c} \sum_{q=1}^{n_c} \gamma_q \frac{e^{-(\mathbf{x}_\perp - \mathbf{x}_\perp^i - \mathbf{s}_q)^2 / 2v^2}}{2\pi v^2}$$
$$s(\mathbf{x}_\perp) \propto \left( \frac{T_A^a + T_B^a}{2} \right)^{1/a}$$

**IP-Glasma like entropy deposition with  $a=0$ . For the longitudinal direction,**

$$s(\mathbf{x}_\perp, \eta_s)|_{\tau=\tau_0} = K s(\mathbf{x}_\perp) g(\mathbf{x}_\perp, y) \frac{dy}{d\eta_s},$$

**We construct the function from  $g$  by parameterizing its cumulant generating function.**

# CLVisc Framework

## ➤ CLVisc Framework

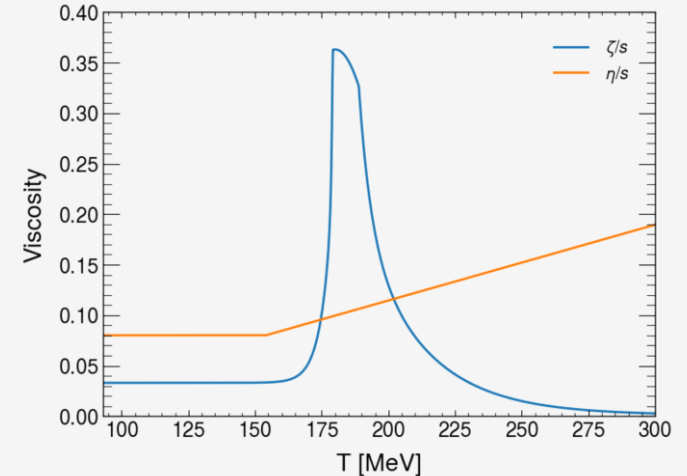
Pang, Wang, and Wang, PRC 86, 024911 .  
Wu, Qin, Pang, and Wang, PRC 105, 034909.

The subsequent evolution of the system is simulated by the 3+1D CLVisc hydrodynamics model.

We just focus on the energy-momentum conservation equations

$$\partial_\mu T^{\mu\nu} = 0,$$

$$\begin{aligned} \tau_\Pi D\Pi + \Pi &= -\zeta\theta - \delta_{\Pi\Pi}\Pi\theta + \lambda_{\Pi\pi}\pi^{\mu\nu}\sigma_{\mu\nu} \\ \tau_\pi \Delta_{\alpha\beta}^{\mu\nu} D\pi^{\alpha\beta} + \pi^{\mu\nu} &= \eta_v\sigma^{\mu\nu} - \delta_{\pi\pi}\pi^{\mu\nu}\theta + \tau_{\pi\pi}\pi^{\lambda\langle\mu}\sigma_{\lambda}^{\nu\rangle} \\ &\quad + \varphi_1\pi_{\alpha}^{\langle\mu}\pi^{\nu\rangle\alpha}. \end{aligned}$$

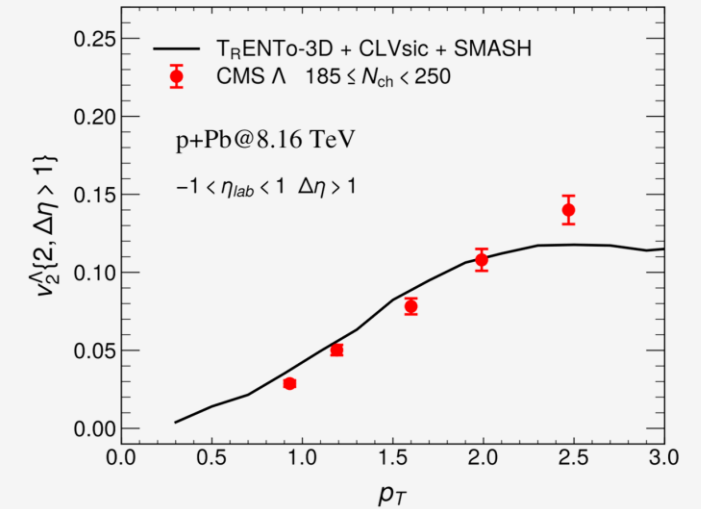
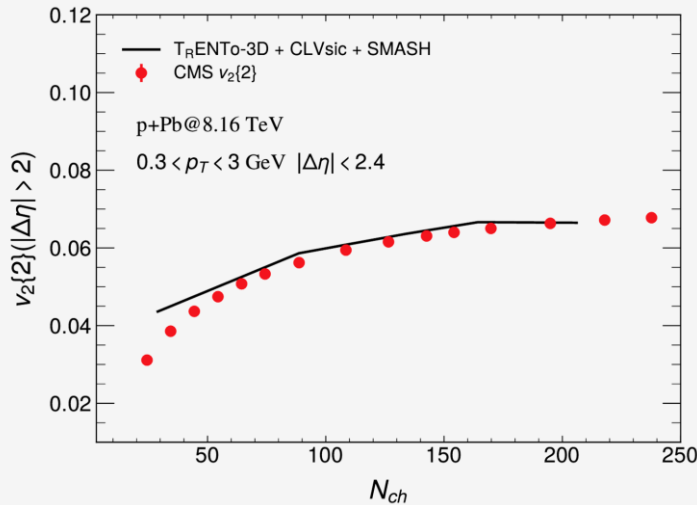
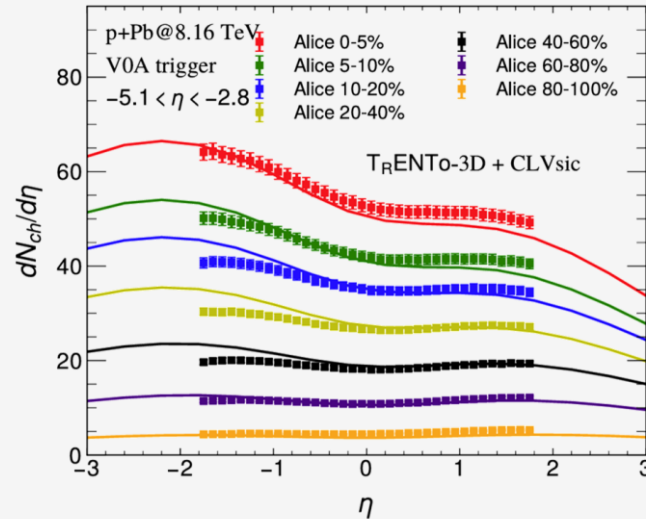


We use the temperature dependent shear and bulk viscosity given by Bayesian parameter estimation in Phys. Rev. C 94, 024907 (2016) .

The equations of state is given by the HotQCD collaboration and freeze-out temperature  $T_f = 154$  MeV.

# Fitting Multiplicity and $v_2$

## ➤ Bulk properties



Multiplicity intervals	$\langle N_{ch} \rangle_{\text{exp}}$	$\langle N_{ch} \rangle_{\text{CLVvisc}}$
[185,250)	203.3	204.2
[150,185)	163.6	164.5
[120,150)	132.7	133.57
[60,120)	86.7	87.7
[3,60)	40	29.3

➤ **CLVisc with Trento-3D initial condition can have a good description of the multiplicity of charged particles and elliptic flow for  $\Lambda$  hyperons**

# Spin Polarization Vector

We follow the modified Cooper-Frye formula to compute the polarization pseudo-vector including the contribution from thermal vorticity and thermal shear tensor and neglect the spin hall effect:

$$\mathcal{S}^\mu(\mathbf{p}) = \mathcal{S}_{\text{thermal}}^\mu(\mathbf{p}) + \mathcal{S}_{\text{th-shear}}^\mu(\mathbf{p}),$$

$$\mathcal{S}_{\text{thermal}}^\mu(\mathbf{p}) = \hbar \int d\Sigma \cdot \mathcal{N}_p \frac{1}{2} \epsilon^{\mu\nu\alpha\beta} p_\nu \varpi_{\alpha\beta},$$

$$\mathcal{S}_{\text{th-shear}}^\mu(\mathbf{p}) = \hbar \int d\Sigma \cdot \mathcal{N}_p \frac{\epsilon^{\mu\nu\alpha\beta} p_\nu n_\beta}{(n \cdot p)} p^\sigma \xi_{\sigma\alpha},$$

$$\text{thermal vorticity:} \quad \varpi_{\alpha\beta} = \frac{1}{2} \left[ \partial_\alpha \left( \frac{u_\beta}{T} \right) - \partial_\beta \left( \frac{u_\alpha}{T} \right) \right],$$

$$\text{thermal shear tensor:} \quad \xi_{\alpha\beta} = \frac{1}{2} \left[ \partial_\alpha \left( \frac{u_\beta}{T} \right) + \partial_\beta \left( \frac{u_\alpha}{T} \right) \right]$$

# Spin Polarization

In the experiment, the polarization of  $\Lambda$  are measured in their own rest frames. Therefore, we express the polarization pseudo vector in the rest frame of  $\Lambda$ , by taking the Lorenz transformation,

$$\vec{P}^*(\mathbf{p}) = \vec{P}(\mathbf{p}) - \frac{\vec{P}(\mathbf{p}) \cdot \vec{p}}{p^0(p^0 + m)} \vec{p}, \quad P^\mu(\mathbf{p}) \equiv \frac{1}{s} \mathcal{S}^\mu(\mathbf{p}),$$

$P_{2z}$  as a function of  $p_T$  and multiplicity are given by

$$P_{2,z}(p_T) = \frac{\int d\eta \int d\phi [\Phi P^{*z} \sin 2(\phi - \Psi_2)]}{\int d\eta \int d\phi \Phi}$$
$$P_{2,z} = \frac{\int p_T dp_T \int d\eta \int d\phi [\Phi P^{*z} \sin 2(\phi - \Psi_2)]}{\int p_T dp_T \int d\eta \int d\phi \Phi}$$

# Different scenarios

We consider three different scenarios:

➤  **$\Lambda$  equilibrium :**

It is assumed that  $\Lambda$  hyperons reach the local (thermal) equilibrium at the freeze-out hyper-surface

➤ **s quark equilibrium:** Fu, et al. PRL 127, 142301

The spin of  $\Lambda$  hyperons is assumed to be carried by the constituent s quark. We take the s quark's mass instead of  $\Lambda$ 's mass in the simulation

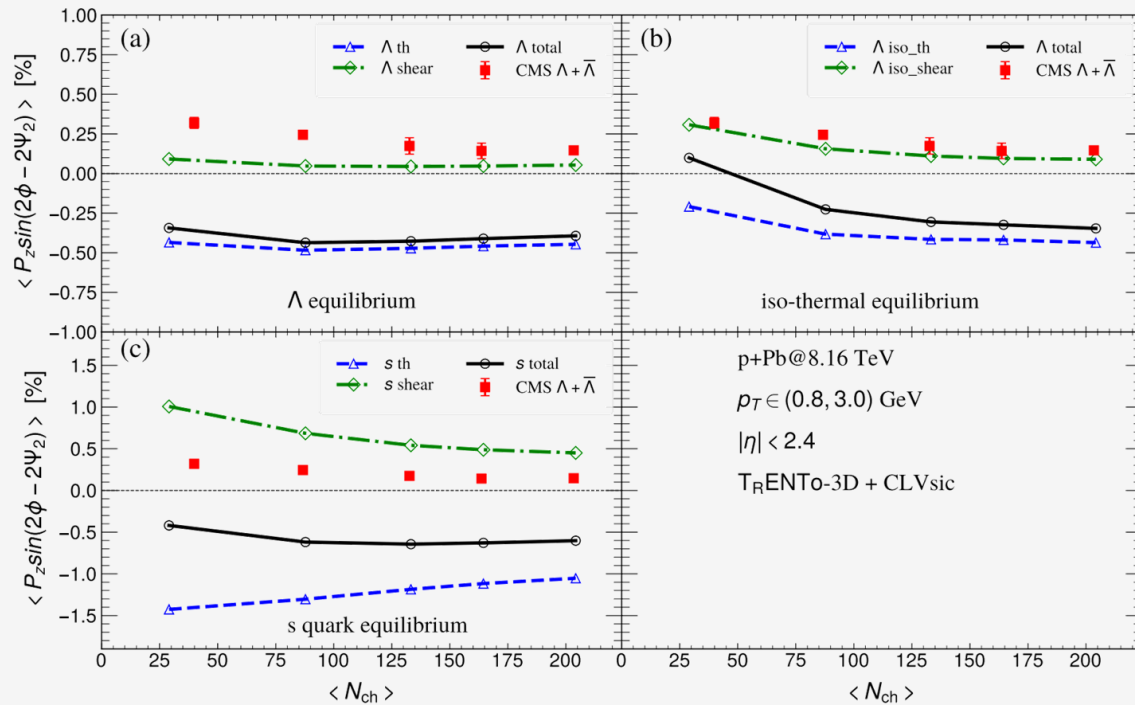
➤ **Iso-thermal equilibrium:** Becattini et al, PRL 127, 272302

The temperature of the system at the freeze-out hyper-surface is assumed to be constant. The time unit vector is taken as fluid velocity for simplicity.

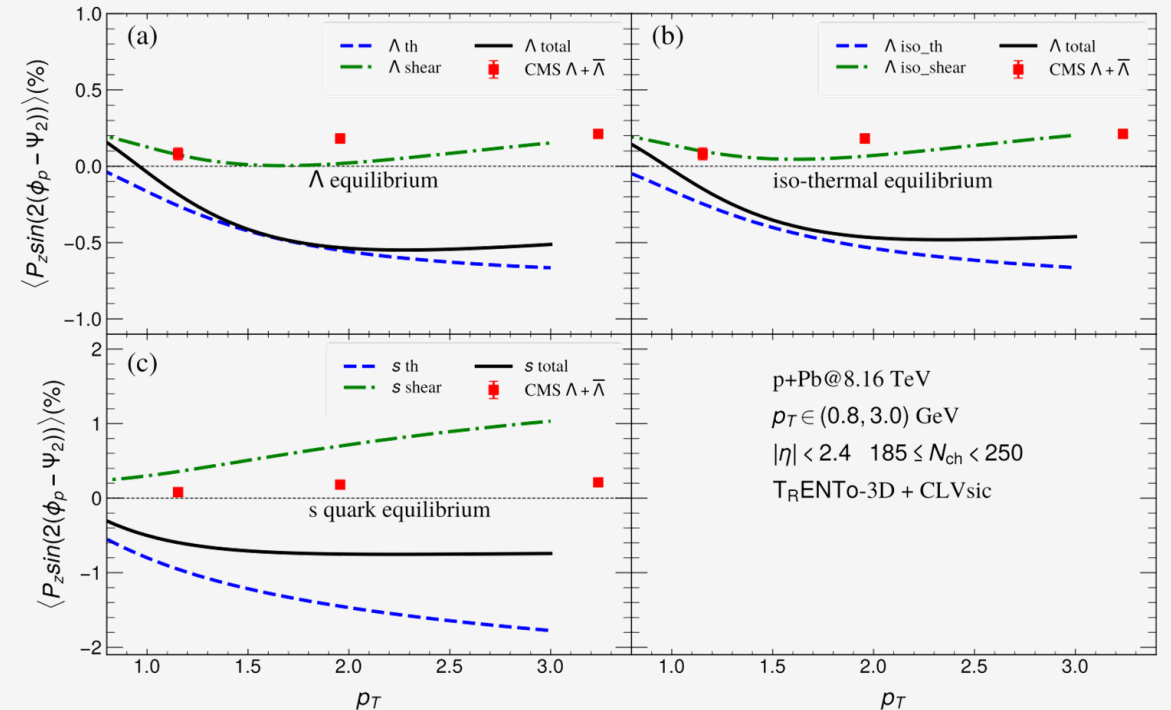
$$\varpi_{\alpha\beta} \rightarrow (\partial_{\alpha}u_{\beta} - \partial_{\beta}u_{\alpha})/(2T)$$

$$\xi_{\alpha\beta} \rightarrow (\partial_{\sigma}u_{\alpha} + \partial_{\alpha}u_{\sigma})/(2T)$$

# Multiplicity (centrality) dependence



Multiplicity dependence

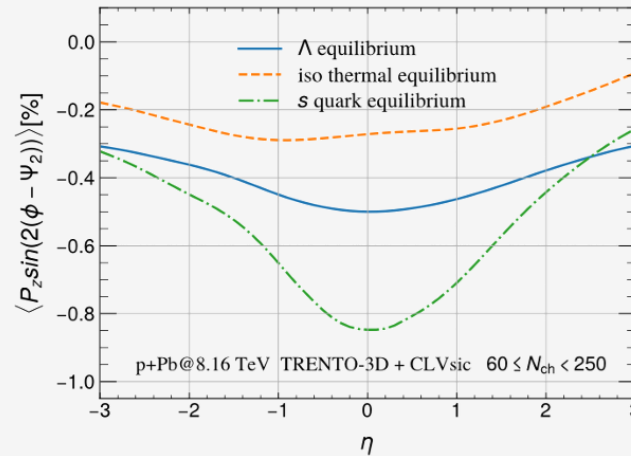
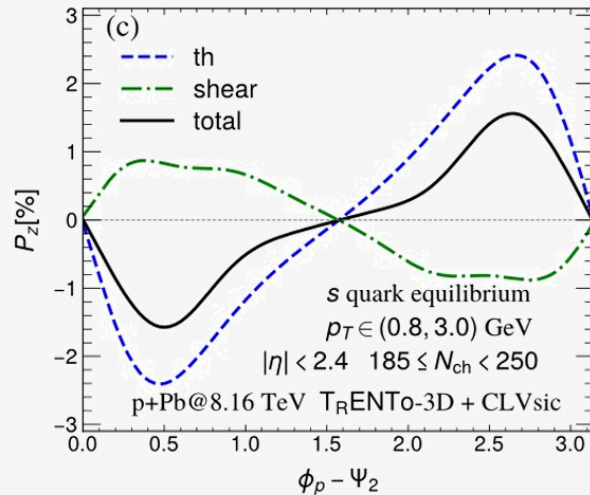
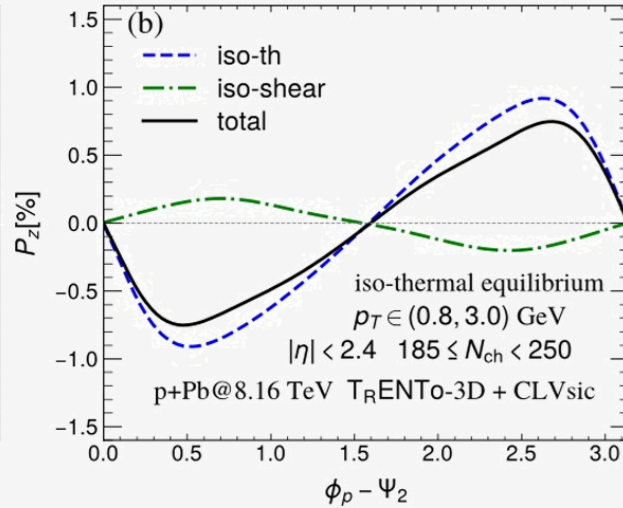
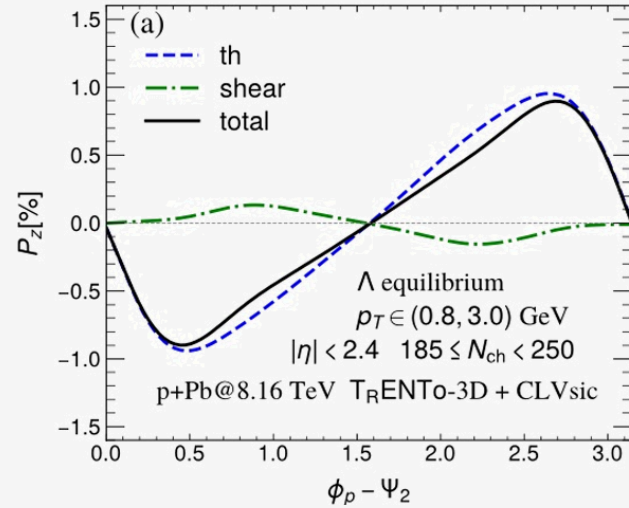


$p_T$  dependence

- Shear induced polarization always gives a positive contribution
- Polarization induced by the thermal vorticity is negative
- The results in the three scenarios are inconsistent with the data from the LHC-CMS experiments.



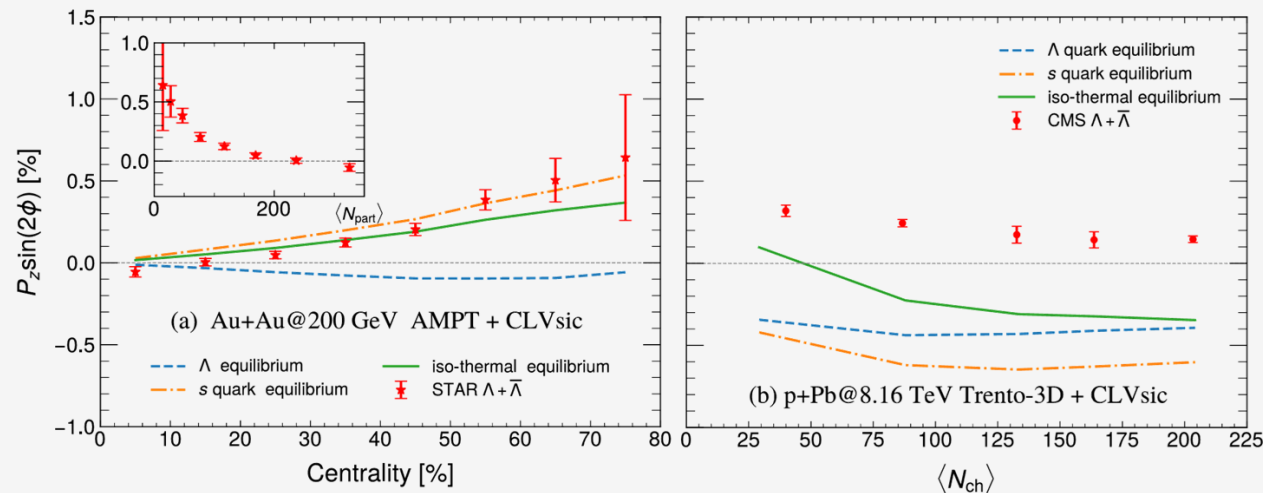
# Azimuthal angle and pseudo-rapidity dependence



➤ The local spin polarization along the beam direction shows a significant deviation from  $\sin(2\phi)$ , mainly due to enhanced event-by-event fluctuations in small systems.

➤  $P_{2z}(\eta)$  is not symmetric under the transformation  $\eta \rightarrow -\eta$ , reflecting the initial-state asymmetry between the proton and the lead nucleus.

# Summary and Outlook

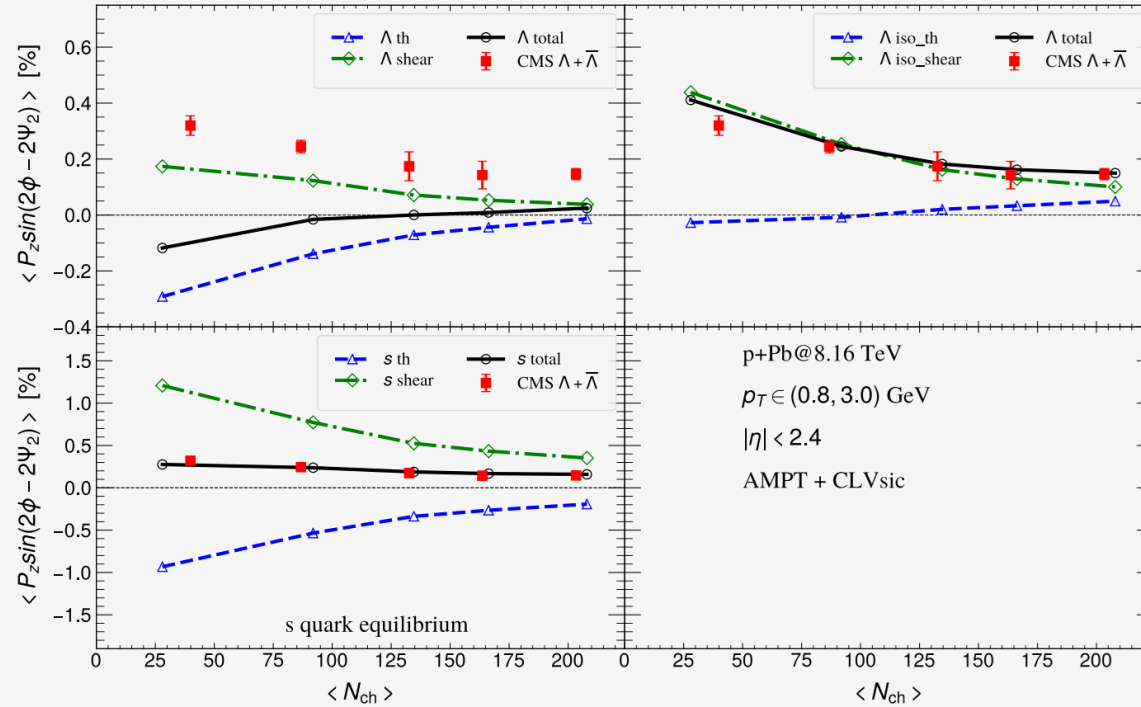


- Shear induced polarization always gives a positive contribution.
- Polarization induced by the thermal vorticity is negative.
- The results from hydrodynamics are inconsistent with the data from CMS.
- New effects need to be considered for the polarization at pPb collisions in the future work.
- We will explore the correlation between  $v_2$  and  $P_{2z}$  across collision system and energy.

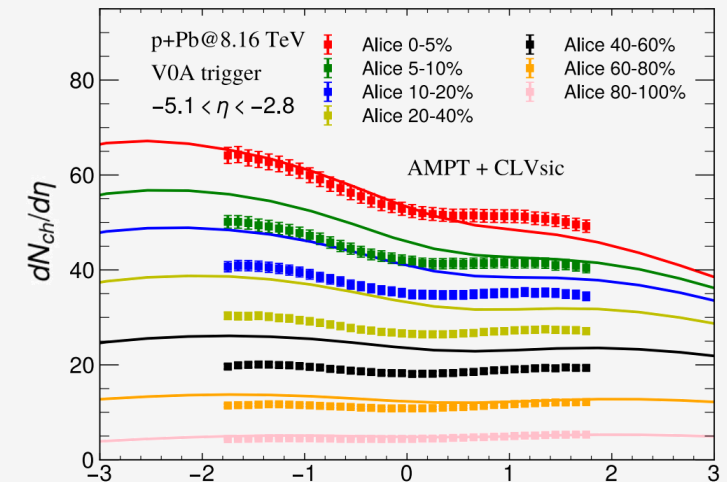
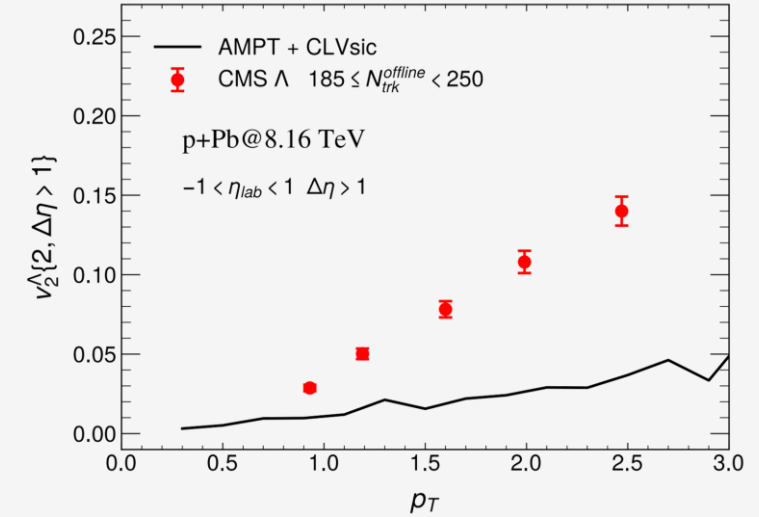
Thanks for your time !

# Back Up

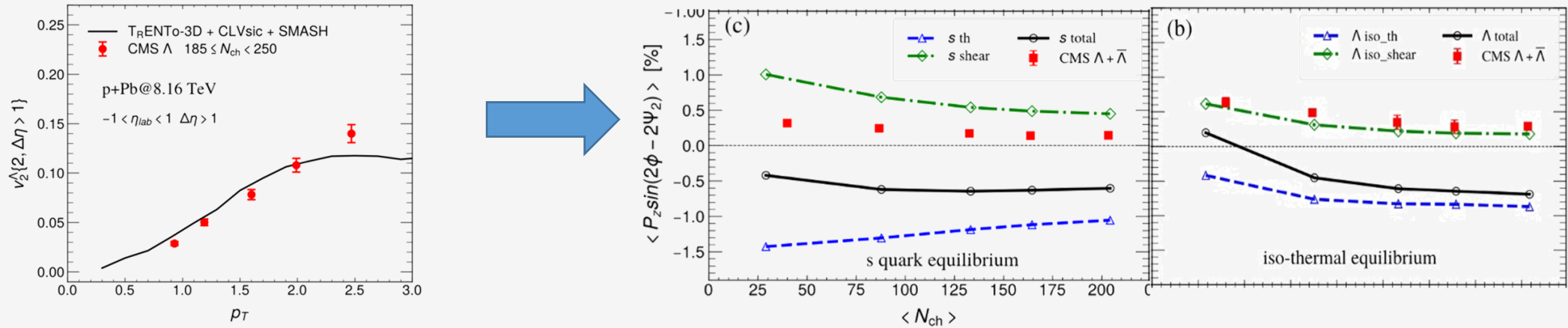
# Test for AMPT initial conditions



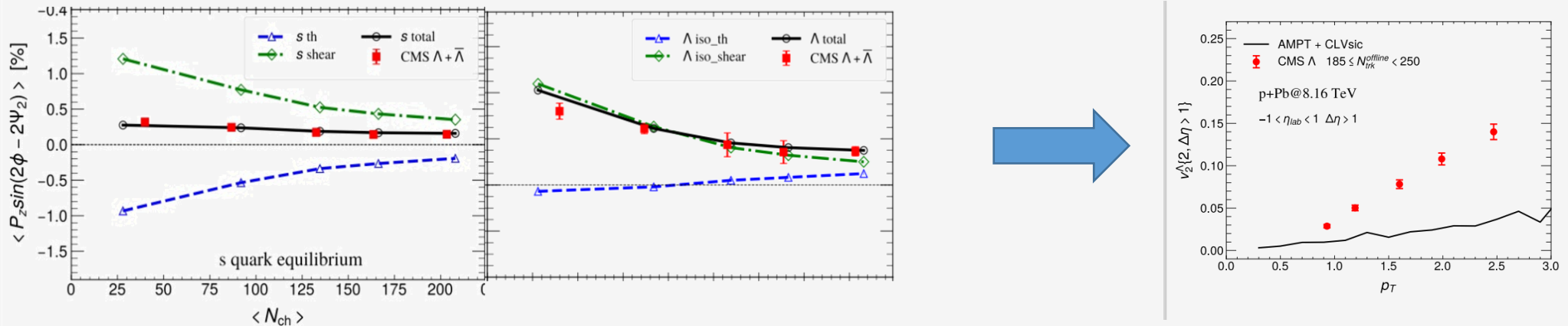
- The parameters can describe spin polarization at the s quark equilibrium and iso-thermal equilibrium can not fit the multiplicity of charged particles and  $v_2$  of  $\Lambda$ .



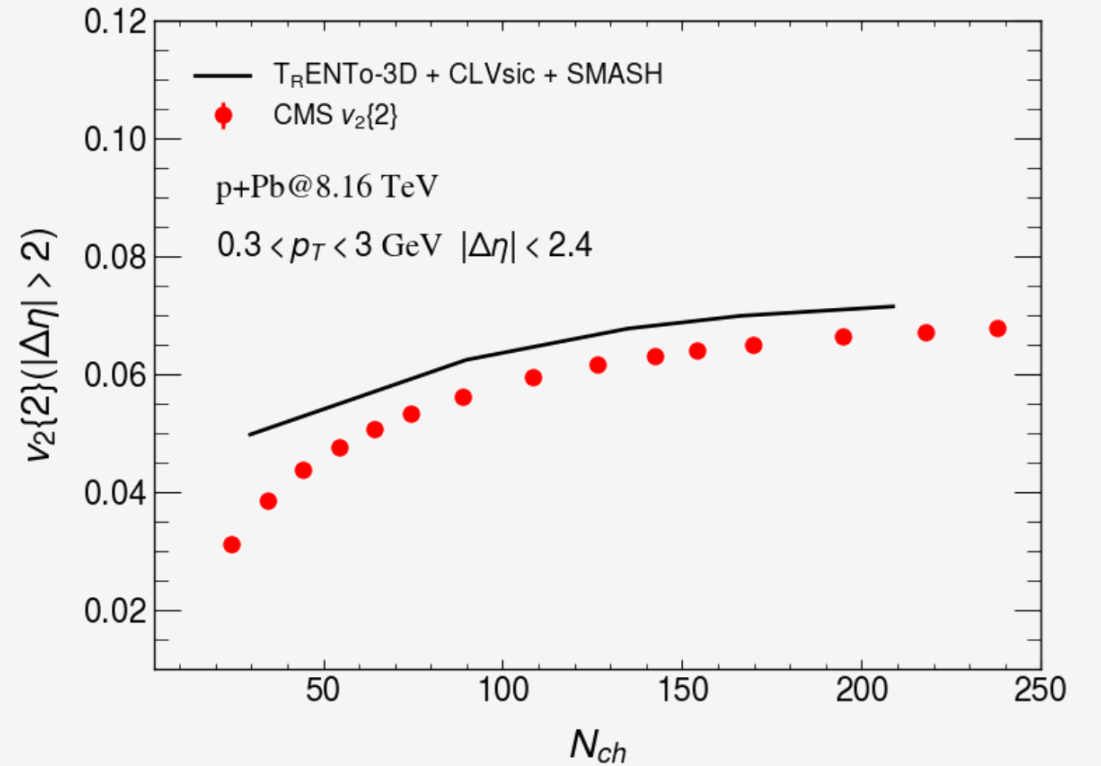
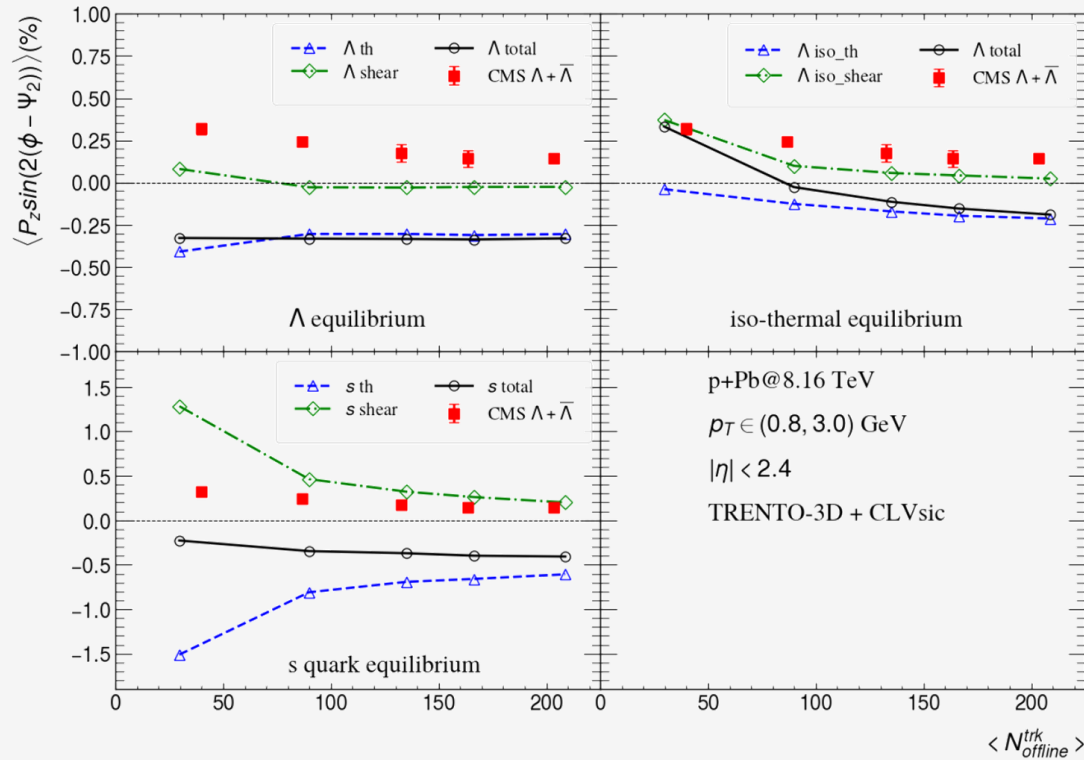
# Different initial conditions



The  $P_{2z}$  of  $\Lambda$  hyperons is not only induced by the  $v_2$  in the p+Pb collisions. New effects need to be considered in the polarization at p+Pb collisions.



# Without bulk viscosity



# Event Plane

The event plane is constructed by the flow vector

$$Q_{n,x} = \sum_i \omega_i \cos(n\phi_i)$$

$$Q_{n,y} = \sum_i \omega_i \sin(n\phi_i)$$

$$\Psi_n = \frac{1}{n} \arctan\left(\frac{Q_{n,y}}{Q_{n,x}}\right)$$

$$0.3 < p_T < 3 \text{ GeV}$$

$$|\eta| < 2.4$$