

One-Pont Energy correlator inside jet

An New Jet observable

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Base on arxiv:2507.21613 with Zihao Mi



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A Century of Spin

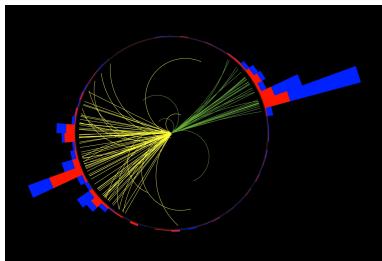


Spin 2025 Shandong China September 21-26 2025

Outline

- Motivation
- One-Point Energy Correlator
- Factorization and Resummation
- Summary and Outlook

Jet Substructure



Hadrons within the jet are detected at **infinity**

Study QCD dynamics and probe
New Physics

from field theory first principles

Energy flow operator

$$\mathcal{E}(n) = \lim_{r \rightarrow \infty} \int d\vec{n} \int_0^\infty dt T_{0\vec{n}}(t, \vec{n}r) r^2$$

ideal detector

Motivation

Energy-Energy Correlators

$$\text{EEC}(\chi) = \frac{1}{\sigma} \sum_{a,b} \int \frac{E_a E_b}{Q^2} d\sigma_{e^+e^- \rightarrow h_a h_b X} \delta(\chi - \theta_{ab})$$

- Separation of energy scales: different slopes in different regimes, and probe different physics.

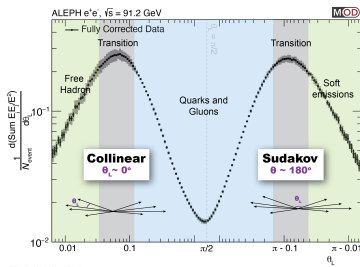
Strong coupling constant

W. Chen, J. Gao, Y. B. Li, Z. Xu, X. Y. Zhang, H. X.

Zhu, [2307.07510](#)

Nucleon structure

X. H. Liu, H. X. Zhu, [2209.02080](#)

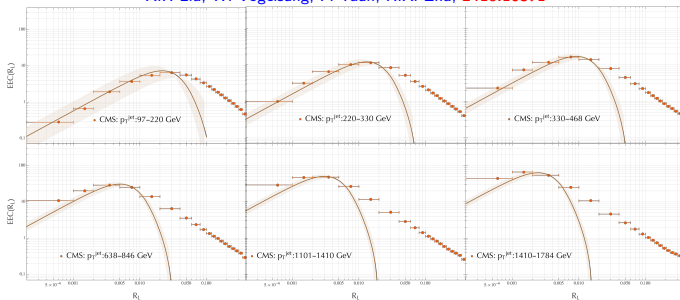


H. Bossi, et al., [2505.11828](#)

TMD Model Description

Near-Side Energy-Energy Correlator

X.H Liu, W. Vogelsang, F. Yuan, H.X. Zhu, **2410.16371**



$$\frac{1}{\sigma_h} \frac{d^3\sigma}{dzd^2p_T} = d_h(z) \int \frac{db}{2\pi} b J_0(p_T b) e^{-S_{\text{Pert}}(b, \mu) - S_{\text{NP}}(b, \mu)}$$

Perturbative Sudakov factor

Non-perturbative Collin-Soper kernel

$S_{\text{Pert}}(b, \mu)$ **calculated by theory**

$S_{\text{NP}}(b, \mu)$ **constrained by experiment**

One-Ponit Energy Correlator

Energy weighted jet cross-section

$$\text{EC}(\chi) = \frac{1}{\sigma_J} \sum_h \int d\sigma \frac{\omega_h}{\omega_J} \delta(\chi - \theta_h)$$

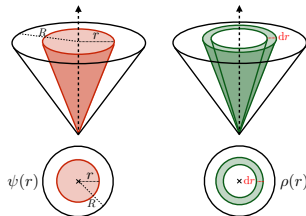
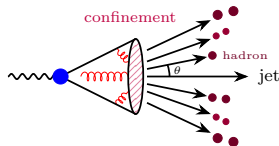
θ_h : angular respect to standard jet axis.

Jet Shape

P. Cal, F. Ringer, W.J. Waalewijn, **1901.06389**

$$\psi(r) = \frac{\sum_{r_i < r} p_{Ti}}{\sum_{r_i < R} p_{Ti}}, \quad \rho(r) = \frac{d\psi(r)}{dr}$$

- fraction of the jet transverse momentum
- Differential mathematically equivalent to EC

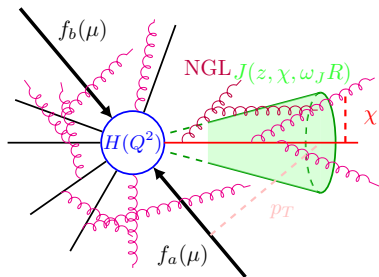


One-Pont Energy Correlator

Inclusive jet cross section

Z.B. Kang, F. Ringer, I. Vitev, **1606.06732**

$$\begin{aligned} & \frac{d\sigma^{pp \rightarrow (\text{jeth})X}}{dp_T d\eta d\chi} \\ &= \sum_{a,b,c} \int_{x_a^{\min}}^1 \frac{dx_a}{x_a} f_a(x_a, \mu) \int_{x_b^{\min}}^1 \frac{dx_b}{x_b} f_b(x_b, \mu) \\ & \times \int_{z^{\min}}^1 \frac{dz}{z^2} H_{ab}^c(\hat{s}, \hat{p}_T, \hat{\eta}, \mu) J_c(z, \chi, \omega_J R, \mu) \end{aligned}$$



TMD Region $\Lambda_{\text{QCD}} \lesssim j_\perp \ll p_T R$

$$\begin{aligned} J_c &= \sum_X \sum_{h \in J} \delta(\theta_h - \chi) \langle 0 | \bar{\chi}_n \delta_{\omega - \bar{n} \cdot P} | (Jh)X \rangle \frac{\omega_h}{\omega_J} \langle (Jh)X | \chi_n | 0 \rangle \\ &= \sum_h \frac{\pi \omega_J^2 \chi}{2} \mathcal{H}(z, \omega_J R, \mu) \int dz_h \int dk_\perp^2 \int d\lambda_\perp^2 \\ & \times \delta^{(2)}(z_h \lambda_\perp + k_\perp - \frac{z_h \omega_J \chi}{2}) \underbrace{z_h D_{h/i}(z_h, k_\perp, \mu, \nu)}_{\text{second order mellin moment}} S_i(\lambda_\perp, \mu, \nu R) \end{aligned}$$

Sum Rule

$$\sum_h \int_0^1 dz \, z \, d_{h/i}(z; \mu) = 1$$

Z.B Kang, X.H Liu, F. Ringer, H.X Xing, **1705.08443**

Resummation Global Logs

RG evolution equation $\Rightarrow \mu_J$

$$\begin{aligned}\frac{d}{d \ln \mu} \mathcal{H}_{i \rightarrow j}(z, \omega_J R, \mu) &= \sum_k \int_z^1 \frac{dz'}{z'} \gamma_{ik} \left(\frac{z}{z'}, \frac{\omega_J \tan(R/2)}{\mu}; \alpha_s(\mu) \right) \mathcal{H}_{k \rightarrow j}(z', \omega_J R, \mu) \\ \frac{d}{d \ln \mu} \ln D_{h/i}(z_h, b, \mu, \nu) &= \gamma_D \left(\frac{\nu}{\omega_J}; \alpha_s(\mu) \right) \\ \frac{d}{d \ln \mu} \ln S(b, \mu, \nu) &= \gamma_S \left(\frac{\nu \tan(R/2)}{\mu}; \alpha_s(\mu) \right) \\ \frac{d}{d \ln \nu} \ln D_{h/i}(z_h, b, \mu, \nu) &= -\gamma_R \left(\frac{\mu}{\mu_b}; \alpha_s(\mu) \right) \\ \frac{d}{d \ln \nu} \ln S(b, \mu, \nu) &= \gamma_R \left(\frac{\mu}{\mu_b}; \alpha_s(\mu) \right)\end{aligned}$$

Central scale choice

$$\mu_b = \frac{2e^{-\gamma_E}}{b}$$

$$\mu_D \sim \mu_b,$$

$$\mu_S \sim \mu_b,$$

$$\mu_J \sim p_T R,$$

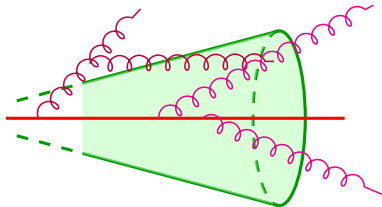
$$\mu_H \sim p_T,$$

Small-R resummation: TimeLike DGLAP $\Rightarrow \mu_H$

$$\mu \frac{d}{d \mu} J_i(z, \chi, \omega_J R, \mu) = \frac{\alpha_s}{2\pi} \sum_j \int_z^1 \frac{dz'}{z'} P_{ji} \left(\frac{z}{z'} \right) J_j(z, \chi, \omega_J R, \mu).$$

Resummation Non-Global Logs

Non-global observable



Wide angle soft radiation

T Becher, B D. Pecjak, D.Y Shao, [1503.06090](#)

Small-R

same as the hemisphere mass

M. Dasgupta, G.P. Salam, [hep-ph/0104277](#)

M. Dasgupta, K. Khelifa-Kerfa, S. Marzani, M.

Spannowsky, [1207.1640](#)

$$U_{\text{NG}}(\mu_b, \mu_J) = \exp \left[-C_A C_a \frac{\pi^2}{3} u^2 \frac{1 + (au)^2}{1 + (bu)^c} \right]$$

$$u = \int_{\mu_b}^{\mu_J} \frac{d\mu}{\mu} \frac{\alpha_s(\mu)}{2\pi} = \frac{1}{\beta_0} \ln \left[\frac{\alpha_s(\mu_b)}{\alpha_s(\mu_J)} \right]$$

LL at large N_C limit

$a = 0.85, b = 0.86, c = 1.33$

Non-perturbative Collin-Soper kernel

Non-perturbative corrections

- Rapidity anomalous dimension
- TMD matrix elements

$$R(b, \mu, \nu) \Rightarrow R(b, \mu, \nu) \times \exp \left(g_K(b) \ln \frac{\nu}{\nu_S} \right)$$

$$SJ_i(b, \mu_0, \nu_0) = S^{\text{pert}} J_i^{\text{OPE}}(b, \mu_0, \nu_0) j_i(b)$$

LO approximation model function

H.T. Li, Y. Makris, I. Vitev, [2102.05669](#)

$$j_i|_{\text{LO}}(b) = \sum_a \int_0^1 dy \, y \underbrace{\frac{di/a(y, 1\text{GeV})}{\text{DSS Fragmentation Dataset}}}_{e^{-0.042b^2/y^2}} \underbrace{\frac{D_{i/a}^{\text{NP}}(y, b)}{e^{-0.042b^2/y^2}}}_{\Rightarrow} \begin{cases} j_q(b) = e^{-0.59b-0.03b^2} \\ j_g(b) = e^{-0.17b-0.09b^2} \end{cases}.$$

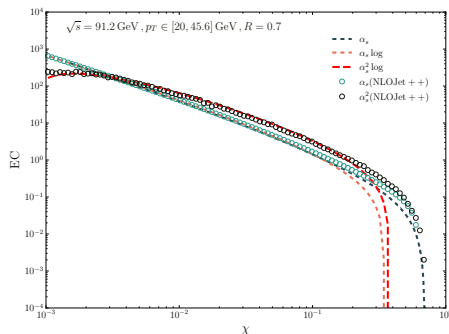
b_* prescription

$$b_{\text{max}} = 1.5$$

$$b_* = \frac{b}{\sqrt{1 + b^2/b_{\text{max}}^2}}$$

Stimulations Fixed Order

Comparison with NLOJET++ e^+e^- annihilation



theoretical

- α_s (W/O Factorization)
- $\alpha_s \text{Log}$
- $\alpha_s^2 \text{Log}$
- anti-kt jet algorithm

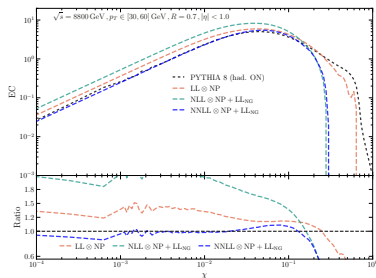
NLOJET++

- α_s
- α_s^2

$$\begin{aligned}
 J_q(z, \chi, \omega_J, \mu) = & \delta(1-z) \frac{\alpha_s}{2\pi^2} \frac{1}{\mu^2} \frac{\pi \omega_J^2 (\tan \chi + \tan^3 \chi)}{2} \\
 & \times \left(\frac{\mu^2}{(\frac{\omega_J \tan \chi}{2})^2} \right) \left\{ -\frac{3}{2} + \frac{3 \tan \chi}{2 \tan \frac{R}{2}} - \frac{3 \tan^2 \chi}{8 \tan^2 \frac{R}{2}} - \log \left(\frac{\tan^2 \chi}{4 \tan^2 \frac{R}{2}} \right) \right\}
 \end{aligned}$$

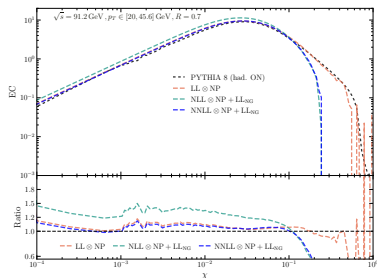
Stimulations Resumamtion

Comparison with Pythia



Theoretical

- NNLL for global logs
- LL for non-global logs
- NP corrections

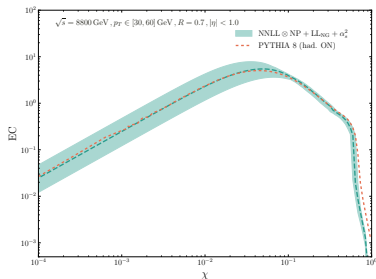
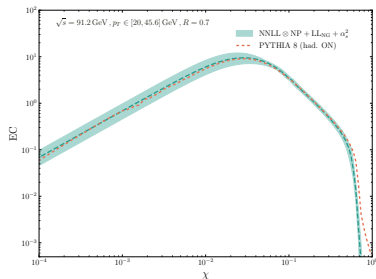


Pythia

- Hadronization
- Decay

Match

Combine fixed order and resumamtion



Match

- Resumamtion for $\lesssim 0.1$
- Fixed order for $\gtrsim 0.1$

Similar behavior as EEC

- Free hadron gas
- Confinement
- Perturbative

Summary And Outlook

Summary

- First Proposed and defined the one-point energy correlator inside jet in TMD region
- Developed a factorization framework and performed resummation of global and non-global logarithms
- NLO EC jet function without TMD Factorization

Outlook

- Similar QCD scaling behavior as EEC
- NLL non-global
- Track-Based EC for the measurement
- Nucleon structure using Spin-dependent EC

C. H Chang, H. Chen, X. H. Liu, D. Simmons-Duffin, F. Yuan, H. X. Zhu, [2507.15923](#)

M. S. Gao, Z. B. Kang, W. C. Li, D. Y. Shao, [2509.15809](#)

Thank you for listening!