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Unifying the study of leading and subleading twist PDFs within Dyson-Schwinger equations.

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Outline

- A brief introduction to $f(x)$ and $e(x)$.
- A new approach to $f(x)$ and $e(x)$ through Dyson-Schwinger equations.
- Some new results.
- Summary & Outlook.

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f(x) & e(x)

$$f(x) = \int \frac{d\xi^-}{4\pi} e^{ixP^+\xi^-} \langle P | \bar{\psi}(0) \gamma^+ \psi(\xi) | P \rangle \Big|_{\xi^+=0, \xi_\perp=0} \quad e(x) = \frac{P^+}{M} \int \frac{d\xi^-}{4\pi} e^{ixP^+\xi^-} \langle P | \bar{\psi}(0) I_4 \psi(\xi) | P \rangle \Big|_{\xi^+=0, \xi_\perp=0}$$

$$\Phi_{ij}(k, P, S) = \int d^4\xi e^{ik \cdot \xi} \langle PS | \bar{\psi}_j(0) \psi_i(\xi) | PS \rangle.$$

$$\Phi_{ij}(x) = \int \frac{d^4k}{(2\pi)^4} \Phi_{ij}(k, P, S) \delta(x - k^+/P^+)$$

$$\Phi(x) = \underbrace{\frac{1}{2} \{ f(x) \not{P} + \lambda_N \Delta f(x) \gamma_5 \not{P} + \Delta_T f(x) \not{P} \gamma_5 \not{S}_\perp \}}_{\text{Twist-2}} + \frac{M}{2} \underbrace{\left\{ e(x) + g_T(x) \gamma_5 \not{S}_\perp + \frac{\lambda_N}{2} h_L(x) \gamma_5 [\not{P}, \not{n}] \right\}}_{\text{Twist-3}} + \dots$$

Hadron Polarization	0	L	T
Twist 2	$f(x)$	$\Delta f(x)$	$\Delta_T f(x)$
Twist 3	$e(x)$	$h_L(x)$	$g_T(x)$

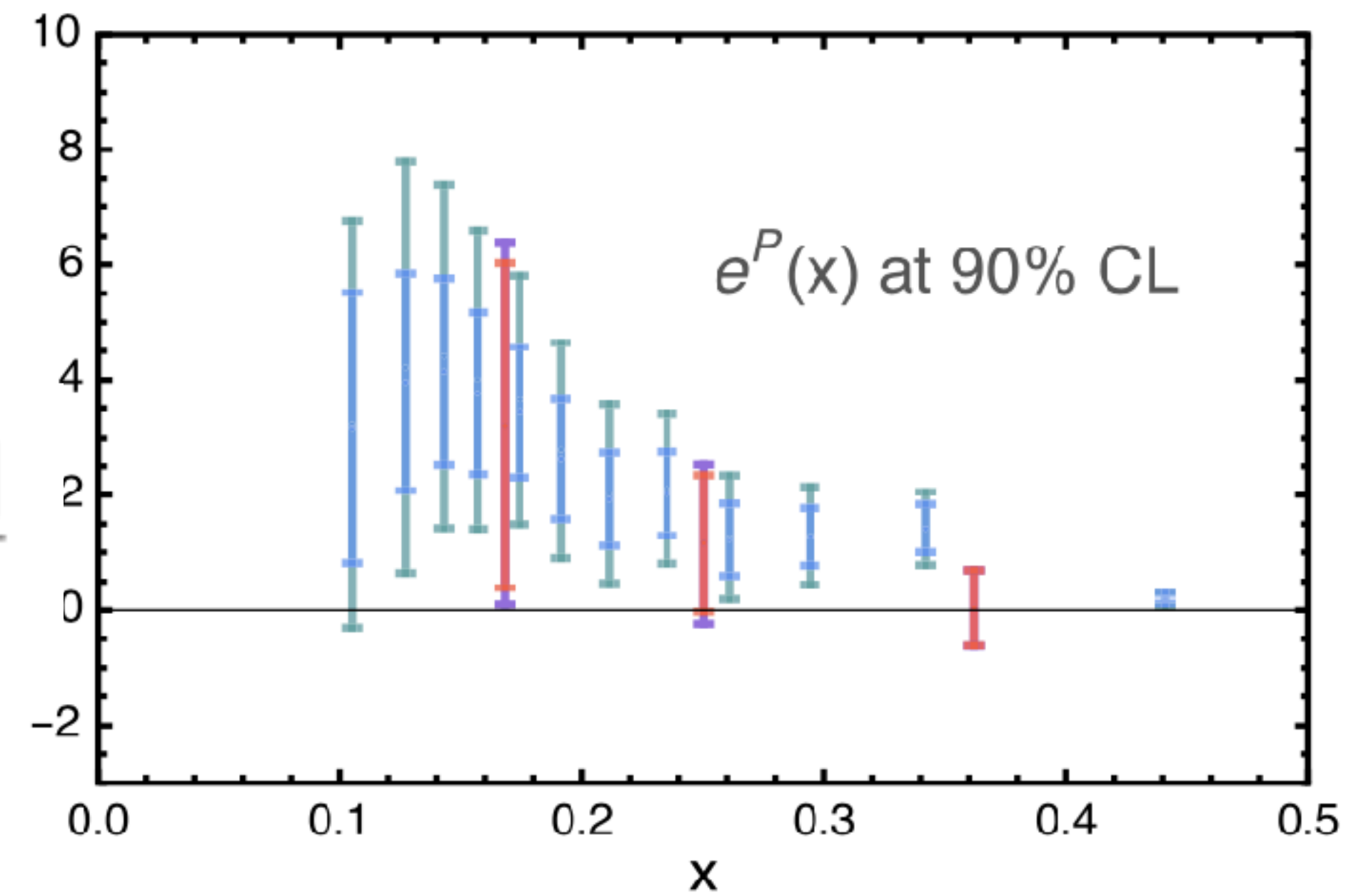
$e(x)$ in experiment.

- Dihadron SIDIS — beam single-spin asymmetry

$$e(\ell) + p(P) \rightarrow e'(\ell') + \pi^+(P_1) + \pi^-(P_2) + X,$$

$$A_{LU}^{\sin \phi_R} = \frac{\sum_q e_q^2 [x e^q(x, Q^2) H_{1,sp}^{\leftarrow, q}(z, M_h, Q^2) + \frac{M_h}{zM} f_1^q(x, Q^2) \tilde{G}_{sp}^{\leftarrow, q}(z, M_h, Q^2)]}{\sum_q e_q^2 f_1^q(x, Q^2) D_{1,ss+pp}^q(z, M_h, Q^2)}$$

Aurore Courtoy et al, PRD 106, 014027 (2022)
CLAS Collaboration, PRL 126, 152501 (2021)

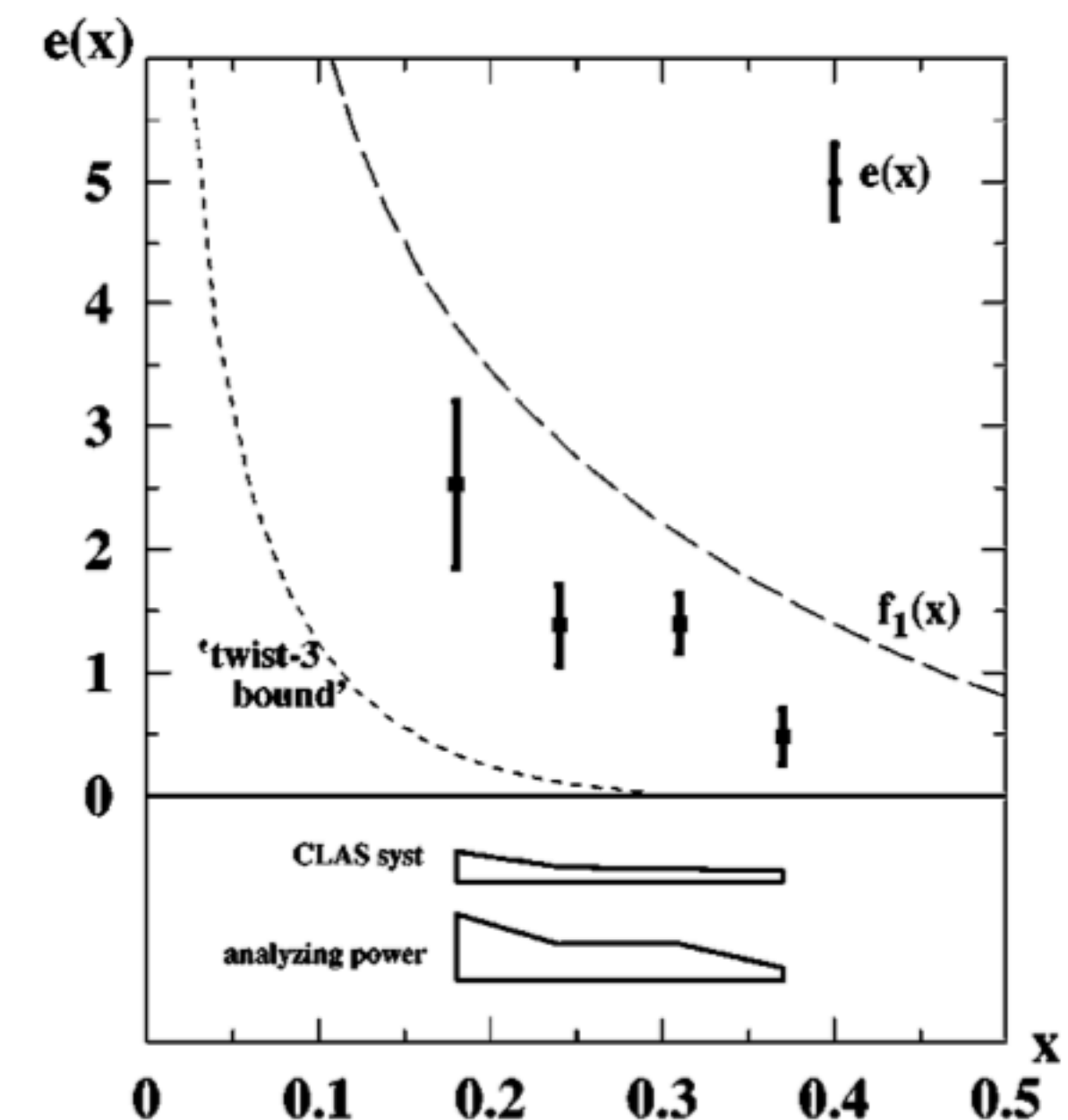


- Single-hadron SIDIS — beam single-spin asymmetry

$$ep \rightarrow e' h X$$

$$A_{LU}^{\sin \phi}(x) = \frac{1}{\langle z \rangle \sqrt{1 + \langle \mathbf{P}_{N\perp}^2 \rangle / \langle \mathbf{k}_{\perp}^2 \rangle}} \frac{\int dy 4y \sqrt{1-y} M_N / Q^5 \sum_a e_a^2 x^2 e^a(x) \langle H_1^{\perp a} \rangle}{\int dy (1 + (1-y)^2) / Q^4 \sum_b e_b^2 x f_1^b(x) \langle D_1^b \rangle}$$

A. Efremov, et al, PRD 67, 114014 (2003)



$e(x)$ decomposition through QCD's EOM

$$e_q(x) = \frac{P \cdot n}{M_h} \int \frac{d\lambda}{4\pi} e^{ixP \cdot n\lambda} \langle P | \bar{\psi}_q(0) [0, \lambda n] \psi_q(\lambda n) | P \rangle$$

$$\begin{aligned} \bar{\psi}(0)[0, z]\psi(z) = & \underbrace{\bar{\psi}(0)\psi(0)}_{\text{red}} + \\ & + \frac{1}{2} \int_0^1 du \int_0^u dv \bar{\psi}(0) \sigma^{\alpha\beta} z_\beta [0, vz] g G_{\alpha\nu}(vz) z^\nu [vz, uz] \psi(uz) - \\ & \underbrace{-im_q \int_0^1 du \bar{\psi}(0) \not{z}[0, uz] \psi(uz)}_{\text{green}} - \\ & - \frac{i}{2} \int_0^1 du \left(\bar{\psi}(0)(i\not{D} - m_q) \not{z}[0, uz] \psi(uz) + \bar{\psi}(0) \not{z}[0, uz](i\not{D} - m_q)\psi(uz) \right). \end{aligned}$$

$$e_q(x) = \underbrace{\frac{\langle h | \bar{\psi}_q(0) \psi_q(0) | h \rangle}{2M_h}}_{\text{red}} \delta(x) + \underbrace{\frac{m_q}{xm_h} f_q(x)}_{\text{green}} + \underbrace{e_q^{\text{tw}3}(x)}_{\text{blue}}$$

Novel features of twist-3 PDF!

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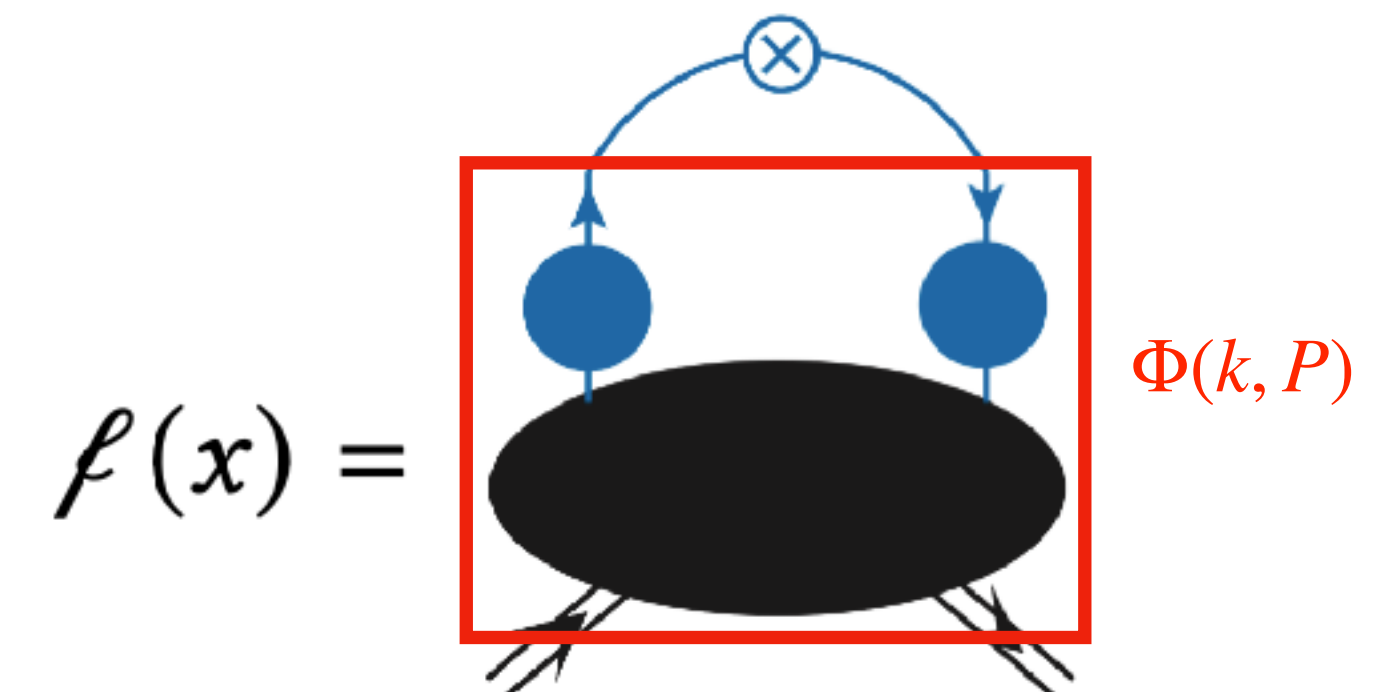
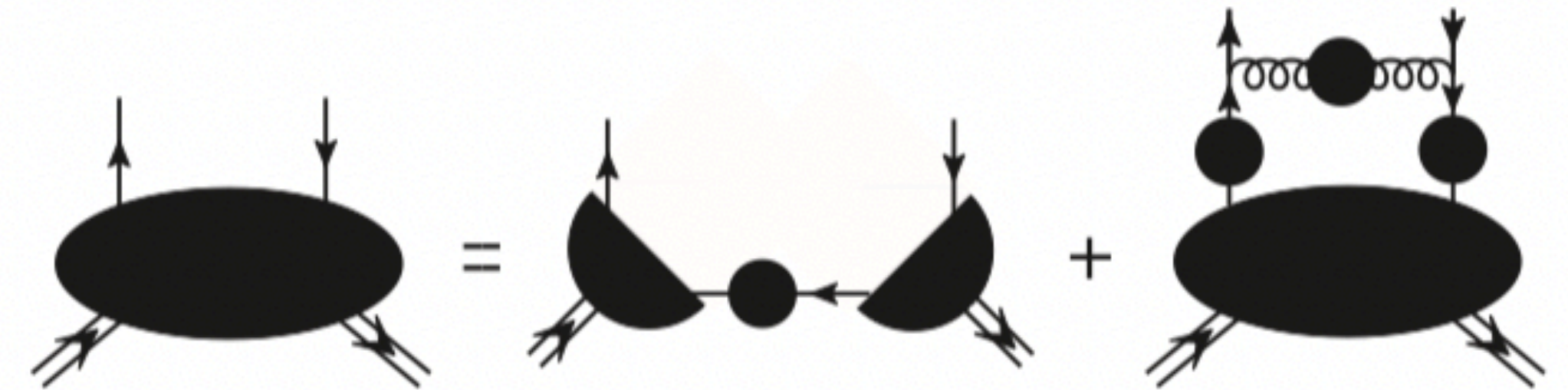
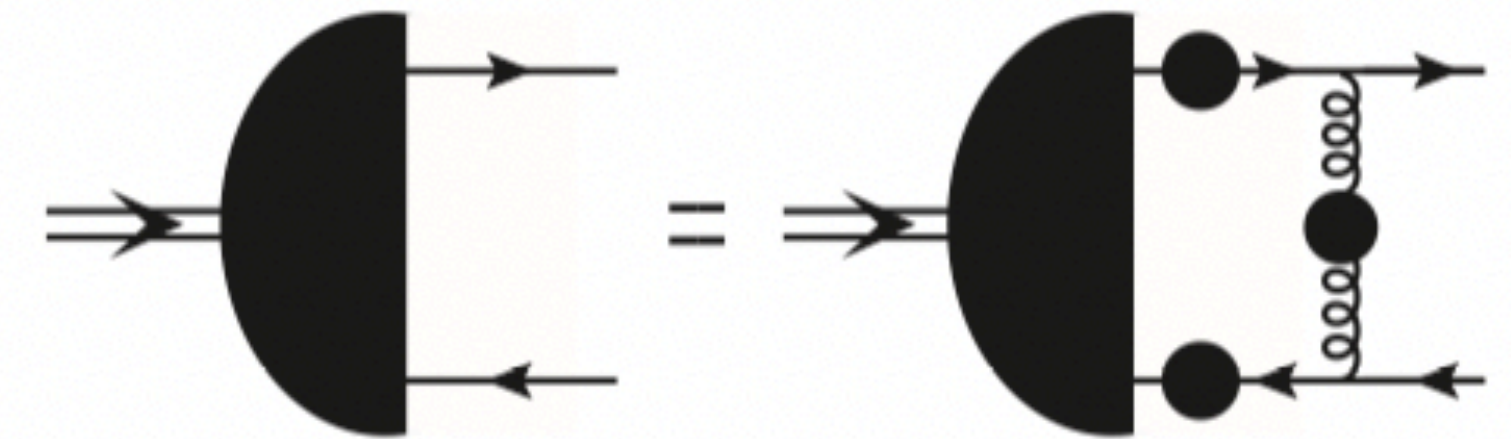
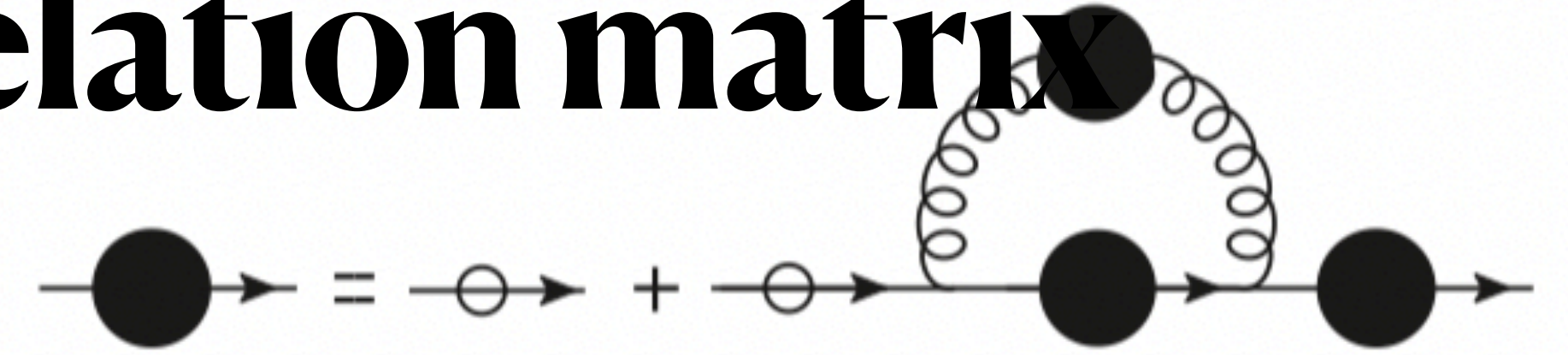
RL-DSEs for quark-quark correlation matrix

- DSEs are EOMs of Green functions.
- Gap equation: quark propagator's DSE.
- Bethe-Salpeter equation: DSE of meson- $\rightarrow q\bar{q}$ vertex.
- We introduce a new DSE aimed for **quark-quark correlation matrix**, which is the **mother function** to various leading and subleading-twist PDFs.

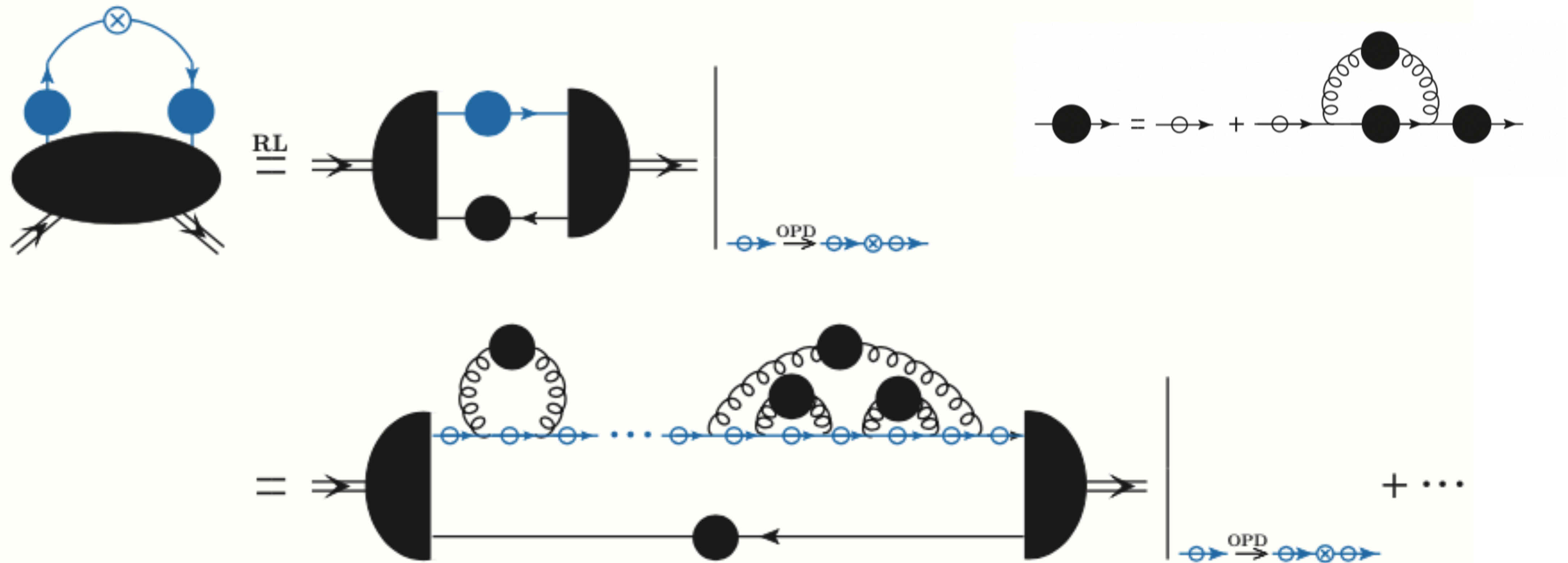
$$\Phi_{ij}(k, P) = \int d^4\xi \, e^{ik \cdot \xi} \langle P | \bar{\psi}_j(0) \psi_i(\xi) | P \rangle.$$

- The twist-2 and -3 PDFs can be simultaneously extracted from Φ

$$\ell(x) = \frac{1}{2} \int \frac{d^4k}{(2\pi)^4} \delta(k_\eta^+ - xP^+) \text{Tr} [\Phi(k, P) \Gamma_\ell]$$



Link Pert. and Nonpert. objects

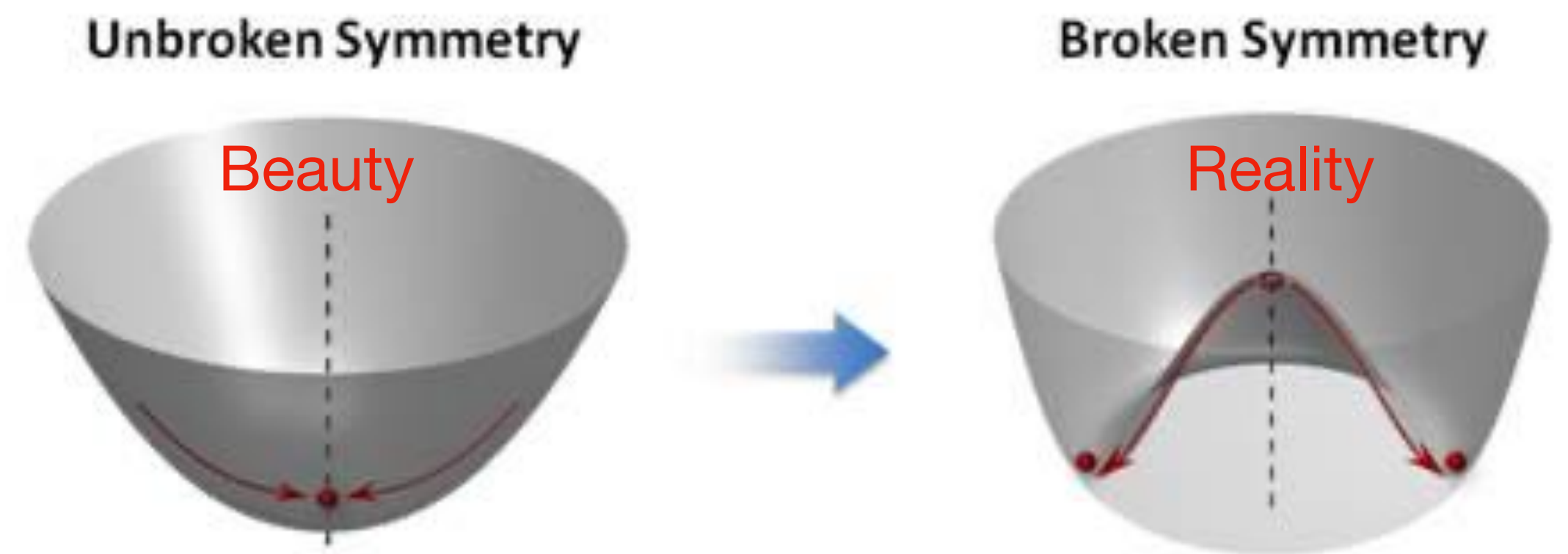


- Every bare quark propagator in the dressed quark propagator is probed by the light-cone vertex concerned.

Symmetries preserved

- Poincare symmetry (Continuum QFT+Regularization)
 - translation, boost, rotation
- Chiral symmetry. (Axial-vector Ward Identity)
 - $m_q \sim 0$ MeV
 - Spontaneous(dynamical) breaking.
 - Evidence: Goldstone boson, etc.
- Global $U(1)$ symmetry. (Vector Ward Identity)
 - Baryon number conservation(quark number sum rule)

$$SU_V(3) \otimes \cancel{SU_A(3)}$$



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Mellin moments

$$\ell(x) = \frac{1}{2} \int \frac{d^4 k}{(2\pi)^4} \delta(k_\eta^+ - xP^+) \text{Tr} [\Phi(k, P) \Gamma_\ell]$$

$$\langle x^n \rangle = \int dx x^n \ell(x)$$

- Compute the Mellin-moments of $f(x)$ and $e(x)$.
- Reconstruct the PDFs by fitting the moments with parameterized ansatz.

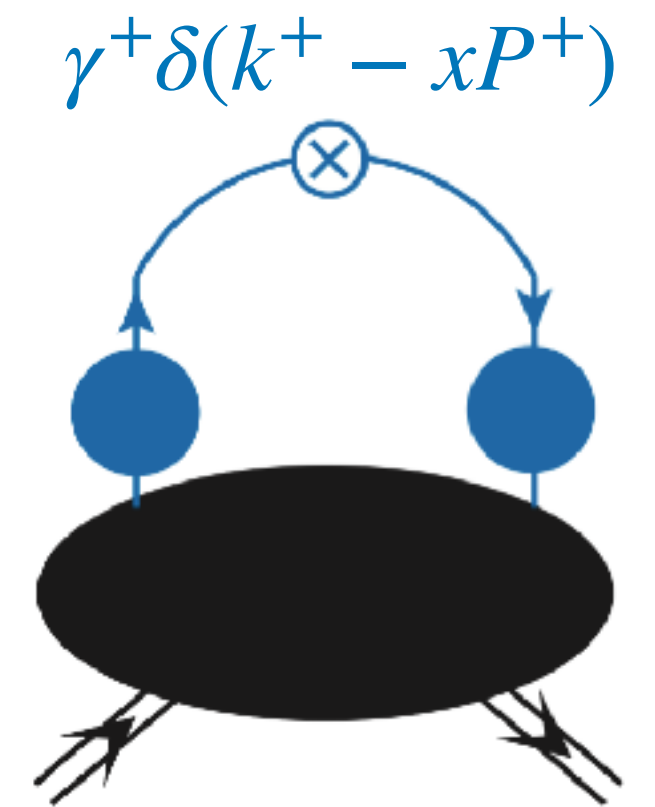
n	0	1	2	3	4	5	6	7	8
$\langle x^n \rangle_{f(x)}$	1.00	0.372	0.202	0.129	0.0910	0.0680	0.0531	0.0428	0.0358
$\langle x^n \rangle_{e(x)}$	6.45	0.113	0.303	0.278	0.231	0.191	0.159	0.134	0.115
$\langle x^n \rangle_{e_{\text{mass}}^{\text{Max}}(x)}$		0.113	0.0419	0.0228	0.0146	0.0103	0.00773	0.00604	0.00488
$\langle x^n \rangle_{e_{\text{tw3}}(x)}$		0	0.261	0.255	0.216	0.181	0.151	0.128	0.110

TABLE I. Computed Mellin moments of pion PDFs.

Pion $f(x)$ at hadronic scale

$$f(x) = \int \frac{d\xi^-}{4\pi} e^{ixP^+\xi^-} \langle P | \bar{\psi}(0) \gamma^+ \psi(\xi) | P \rangle \Big|_{\xi^+=0, \xi_\perp=0}$$

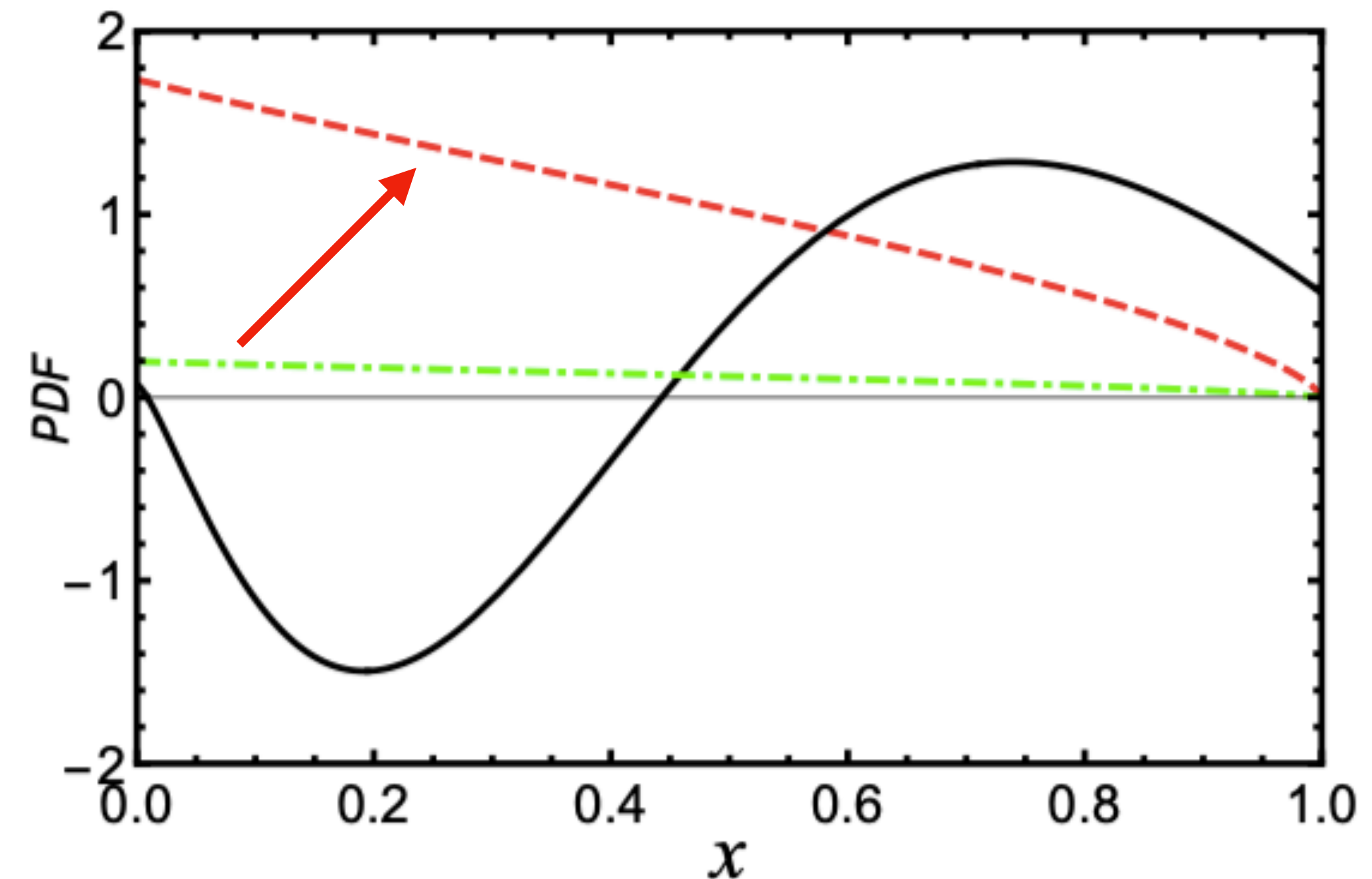
$$f(x) =$$



n	0	1	2	3	4	5	6	7	8
$\langle x^n \rangle_{f(x)}$	1.00	0.372	0.202	0.129	0.0910	0.0680	0.0531	0.0428	0.0358

$$\langle x^n \rangle = \int dx x^n f(x)$$

- $\langle x^0 \rangle_f \equiv \int dx x^0 f(x) = 1$: quark number sum rule.
- $\langle x^1 \rangle_f < 0.5$: gluon carries away 26% momentum fraction.
- The RL-DSE is associated with a low scale, where sea quark distributions are absent.



Pion $e(x)$ at initial scale

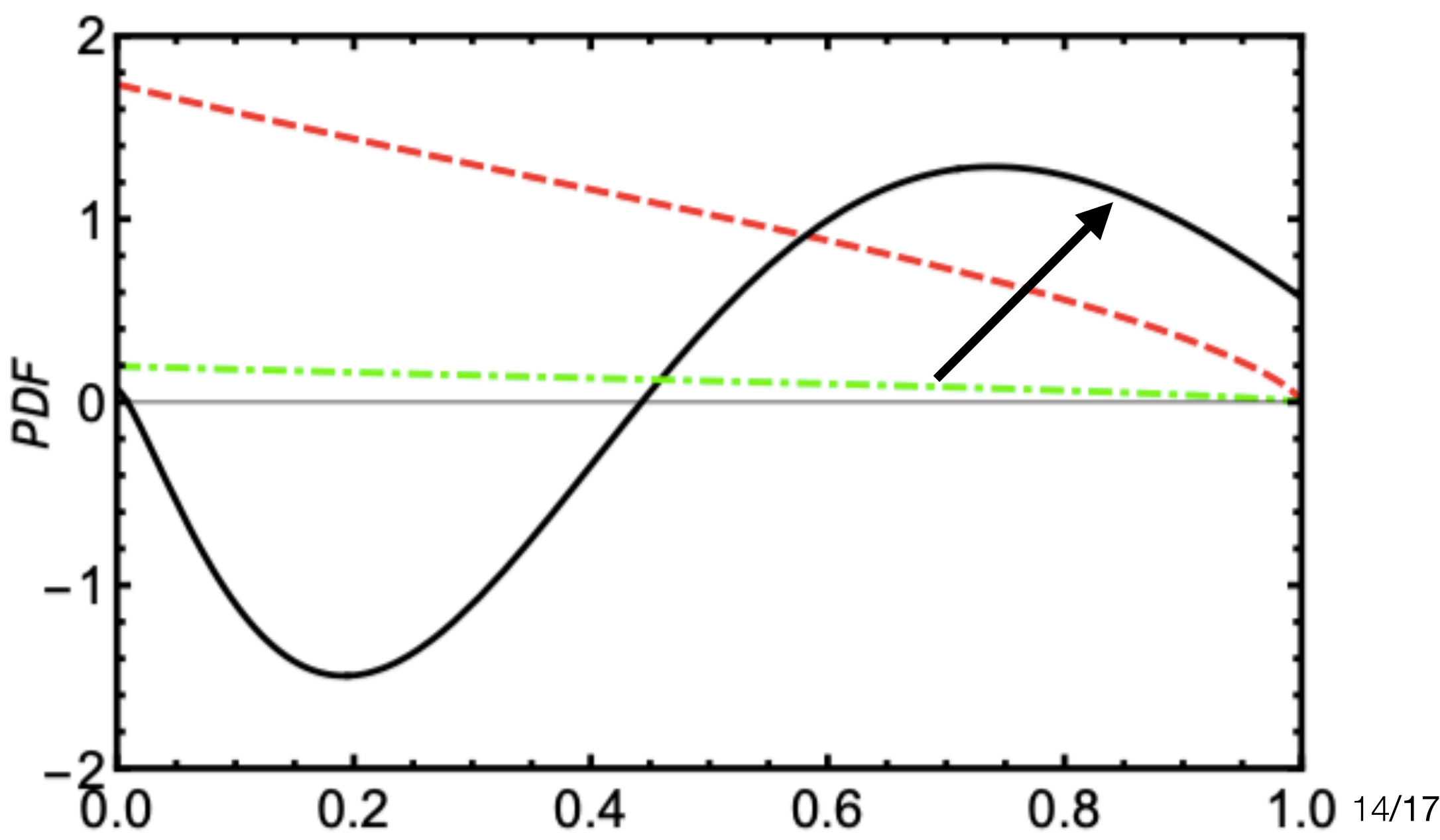
$$e_q(x) = \frac{\langle h | \bar{\psi}_q(0) \psi_q(0) | h \rangle}{2M_h} \delta(x) + \frac{m_q}{xm_h} f_q(x) + e_q^{\text{tw}3}(x)$$

$$e(x) = \frac{P^+}{M} \int \frac{d\xi^-}{4\pi} e^{ixP^+\xi^-} \langle P | \bar{\psi}(0) I_4 \psi(\xi) | P \rangle \Big|_{\xi^+=0, \xi_\perp=0}$$

n	0	1	2	3	4	5	6	7	8
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$\langle x^n \rangle_{e_{\text{mass}}^{\text{Max}}(x)}$		0.113	0.0419	0.0228	0.0146	0.0103	0.00773	0.00604	0.00488
$\langle x^n \rangle_{e_{\text{tw}3}(x)}$		0	0.261	0.255	0.216	0.181	0.151	0.128	0.110

- Strong indication of $\delta(x)$ is found! (Zero mode)
- Strong suppression of first moment is found! (Chiral symmetry) $\langle x \rangle_e = m_q/m_h$
- Significant pure twist-3 contribution is found! Arising from $qA\bar{q}$ correlation.

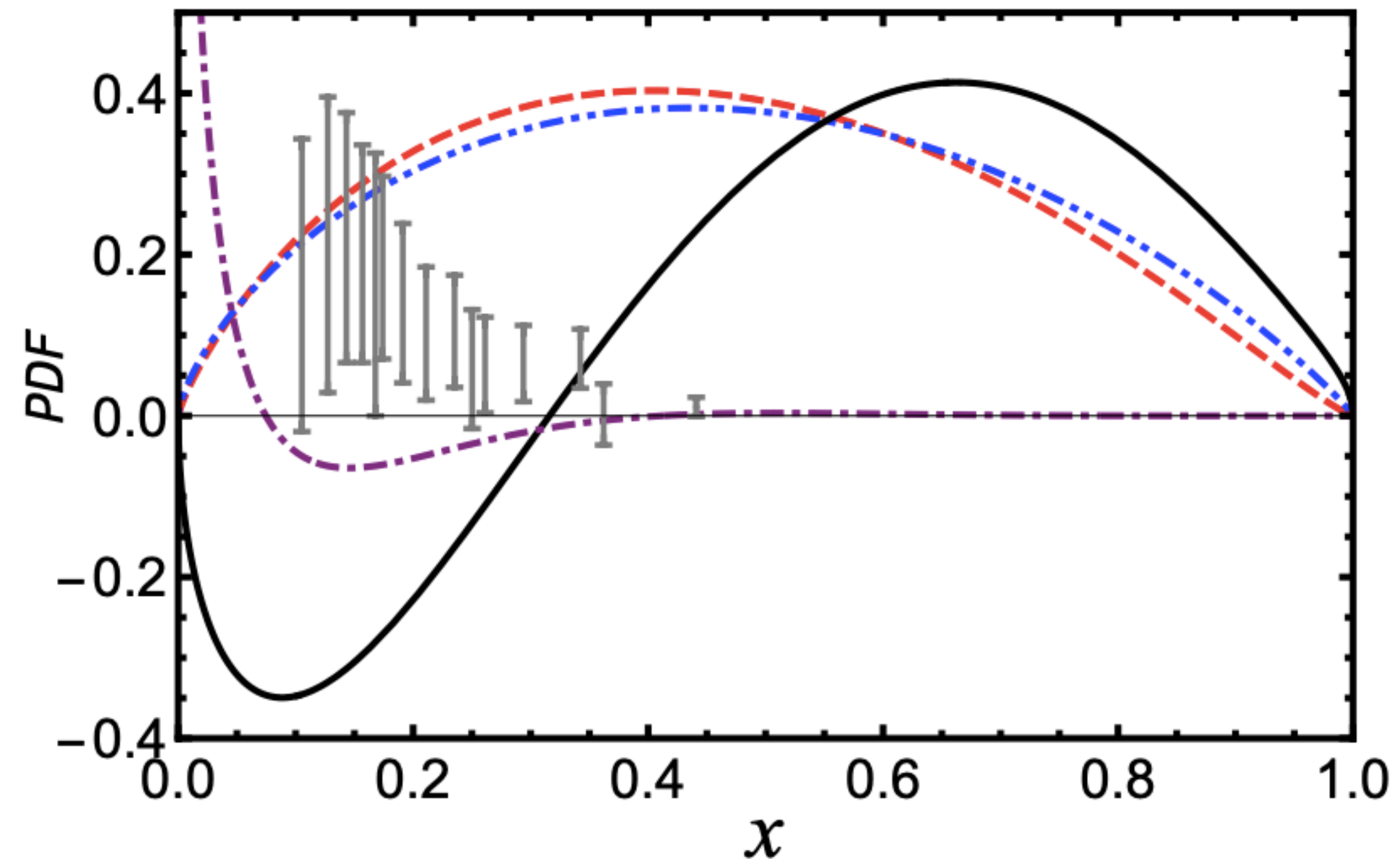
by zero modes in the light-cone formalism [11,16–19]. The scalar PDF has been studied in nonperturbative models for hadron structure [20–28]. In most quark models, the dominant contribution is found to be the mass contribution, which originates from the equation of motions for free



Evolved PDFs and Implications for Nucleon

- Evolution of twist-3 PDFs is rather complicated: mixing between two and three partons
- In large N_C limit the $e(x)$ evolution becomes simpler.
- The $f(x)$ and $e(x)$ are evolved to experimental scale ~ 1 GeV.
- $f(x)$ compares well with JAM global fit.
- The node of $e(x)$ persists under evolution.
- We notice hints of similar pattern in the case of nucleon.

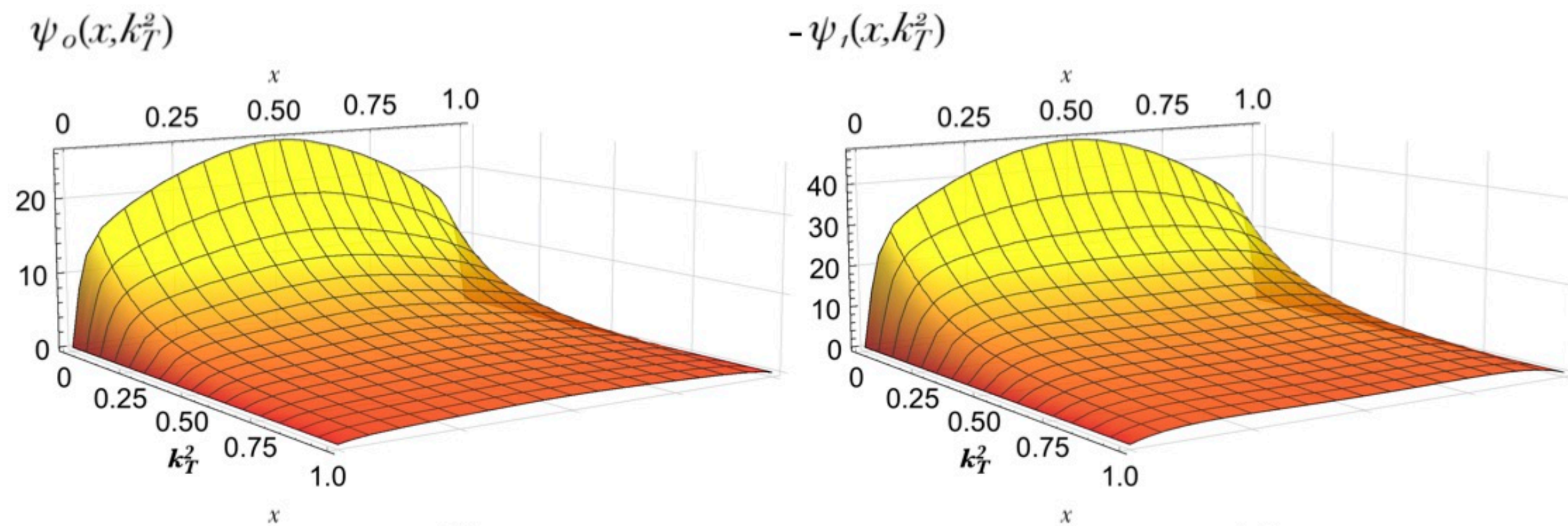
$$\langle x \rangle_e = m_q/m_h$$



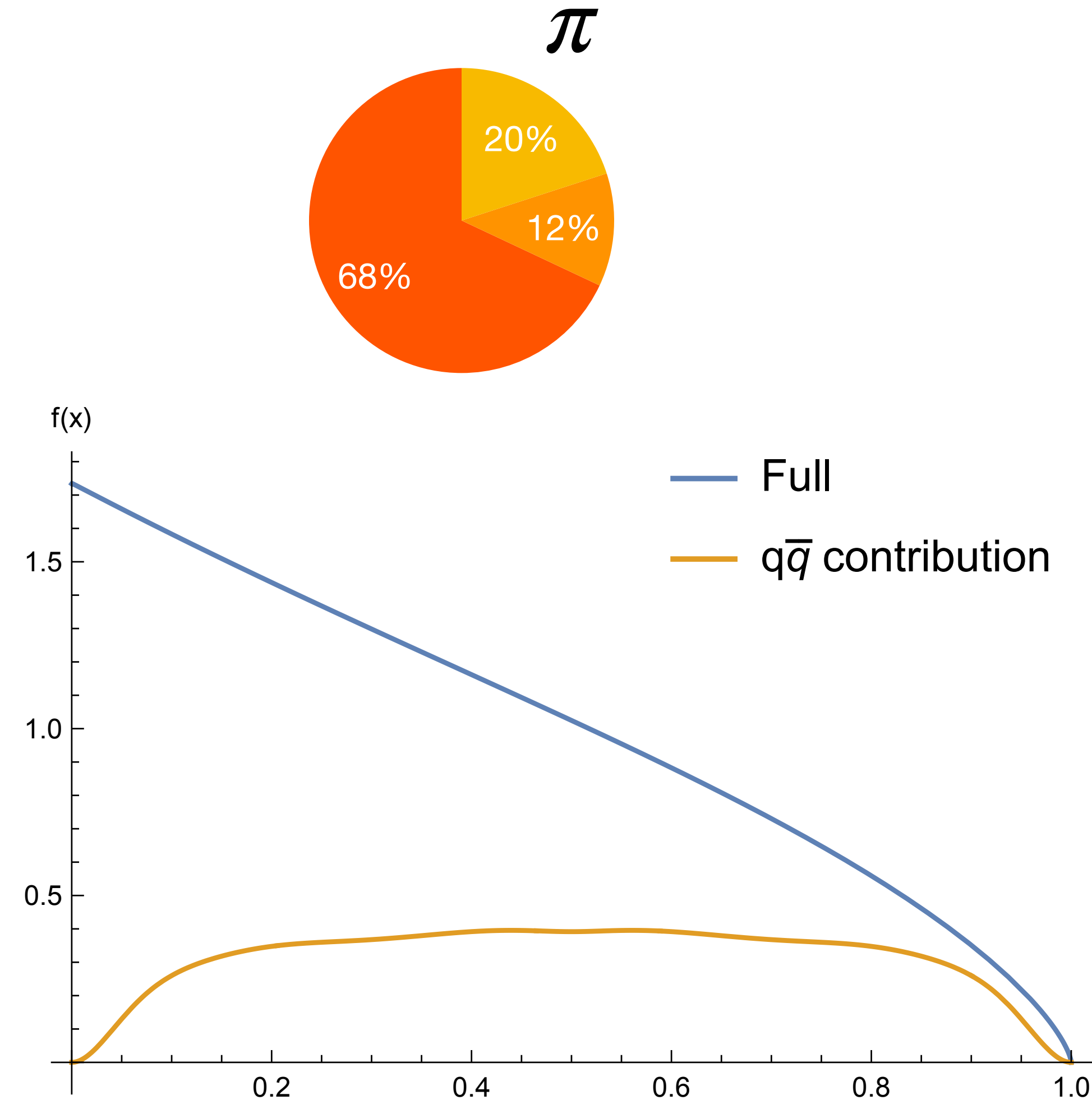
Origin: Higher Fock-states

$$|M\rangle = \phi_2 |q\bar{q}\rangle + \phi_3 |q\bar{q}g\rangle + \phi_4 |q\bar{q}gg\rangle + \dots$$

$$\phi_i(x, \vec{k}_T) \sim \int dk^- dk^+ \delta(xP^+ - k^+) \text{Tr}[\Gamma_i \chi(k, P)]$$



- The $q\bar{q}$ -LFWFs only constitute about 30% to the total normalization.
- Many gluon components reside in pion.



Phys.Rev.Lett. 122 (2019) 8, 082301
 Phys.Rev.D 101 (2020) 7, 074014
 Phys.Rev.D 104 (2021) 9, 094016

Conclusion

- A new Φ -DSE is introduced to render a unified study of $f(x)$ and $e(x)$.
- $e(x)$ contains a Dirac- δ function, through nonperturbative dynamics.
- $e(x)$ is dominated by pure twist-3 part.
- $e(x)$ contains a node.

Outlook

- More PDFs
- TMD PDFs
- Nucleon

Thank you for your attention!