



26th International Symposium on Spin Physics

A Century of Spin

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Physics-informed neural networks for angular momentum conservation in computational relativistic spin hydrodynamics

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1. *Background*

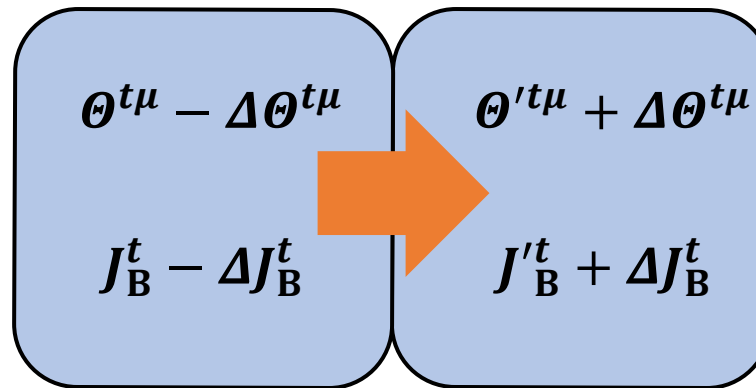
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Numerical Challenges in Angular Momentum Conservation

- Finite Volume Method with Godunov's scheme is widely used in computational hydrodynamics, which enables solving continuity equations ($\partial_\mu \Theta^{\mu\nu} = 0, \partial_\mu J_B^\mu = 0$) while preserving associated conservation laws within numerical accuracy.



- However**, angular momentum conservation is generally not guaranteed in this scheme...

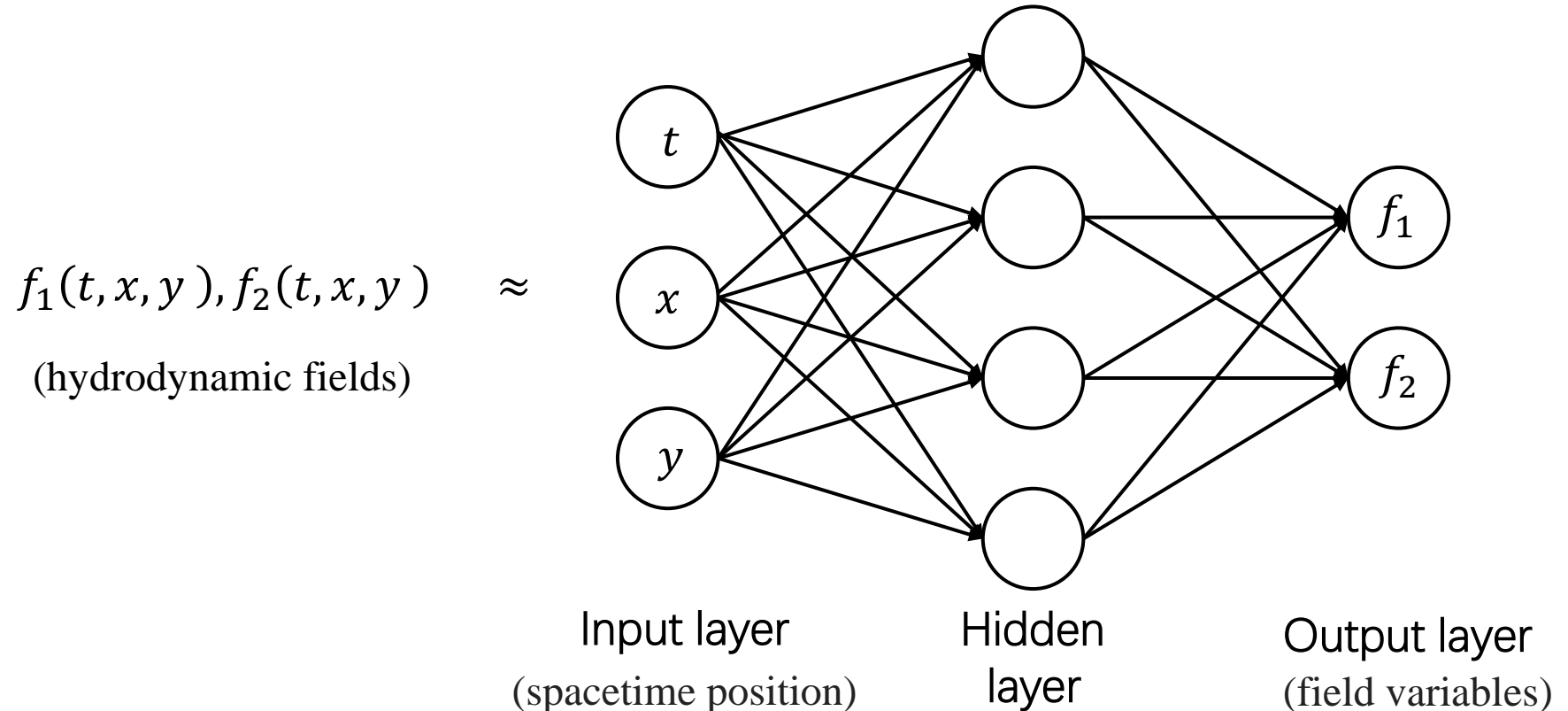
→ **Exploration of complementary numerical strategies**

Physics-Informed Neural Networks (PINNs)

M. Raissi, et.al. (2019).

- Neural network approximate solutions of given partial differential equation (PDE): $\mathcal{F}[f(t, \vec{x})] = 0$
- Loss function = PDE + boundary condition (B.C.) + optional constraint
- Active development of applications in hydrodynamics

M. Raissi, et.al. (2020); Z. Mao, et.al. (2020); X. Jin, et.al. (2021); S. Cai, et. al. (2021).

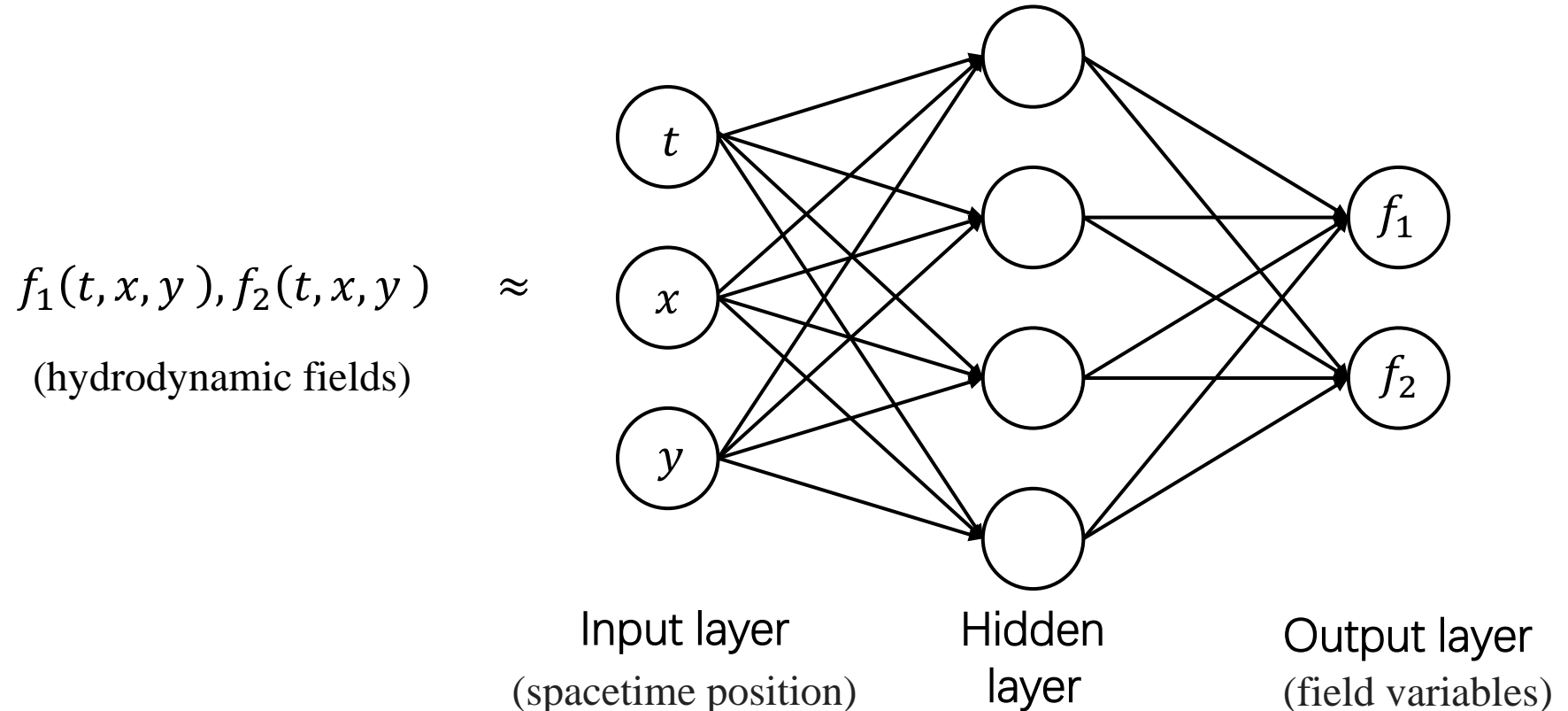


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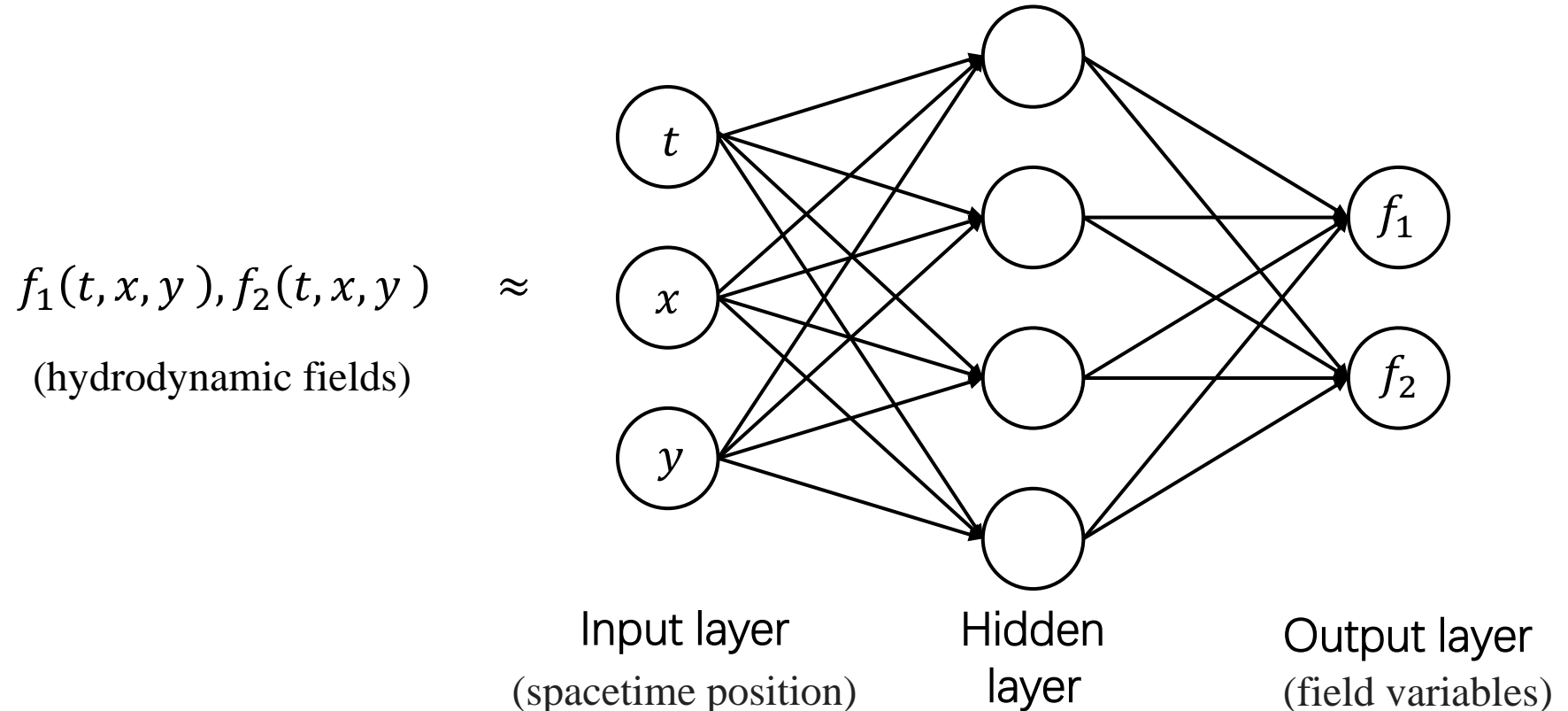


Physics-Informed Neural Networks (PINNs)

M. Raissi, et.al. (2019).

- Neural network approximate solutions of given partial differential equation (PDE): $\mathcal{F}[f(t, \vec{x})] = 0$
- Loss function = PDE + boundary condition (B.C.) + **Angular momentum (AM) conservation**
- Active development of applications in hydrodynamics

M. Raissi, et.al. (2020); Z. Mao, et.al. (2020); X. Jin, et.al. (2021); S. Cai, et. al. (2021).



This Study

- *In this study*, we build a PINNs-based framework for relativistic spin hydrodynamics
- As a proof of concept, we consider a simple setup: 2nd order relativistic spin fluid confined in a cylinder, with symmetry assumptions that reduce the problem to two dimensions (t, r) .

- ✓ Cylindrical symmetry
- ✓ Translational symmetry along the z -direction
- ✓ Parity symmetry with respect to the z -axis

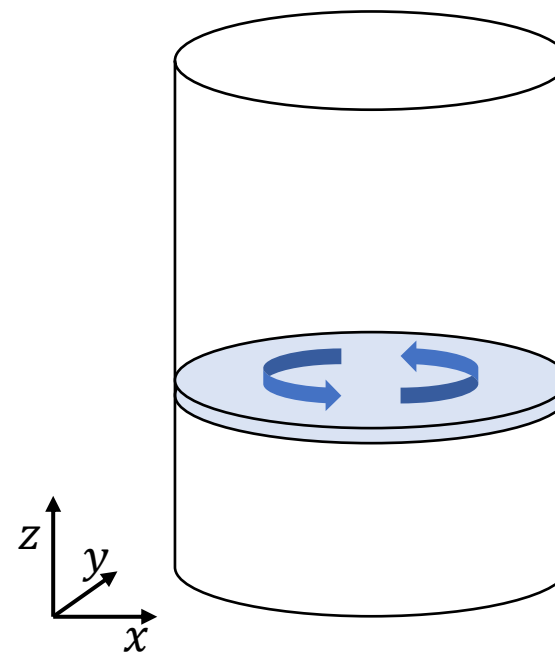


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Relativistic Spin Hydrodynamic Equation

W. Florkowski, et.al. (2018).
X.-G. Huang (2024). [Review]

Hydrodynamic Variables

$$e(t, r), \quad u^r(t, r), \quad u^\theta(t, r), \quad S^z(t, r)(\equiv \Sigma^{txy}), \quad \phi^{r\theta}(t, r)$$

Tensor Decompositions (Landau frame and Totally antisymmetric pseudo gauge)

$$\Theta^{\mu\nu} = eu^\mu u^\nu + P\Delta^{\mu\nu} + \phi^{\mu\nu}$$

$$J^{\mu\nu\xi} = (x^\nu \Theta^{\mu\xi} - x^\xi \Theta^{\mu\nu}) + \Sigma^{\mu\nu\xi}$$

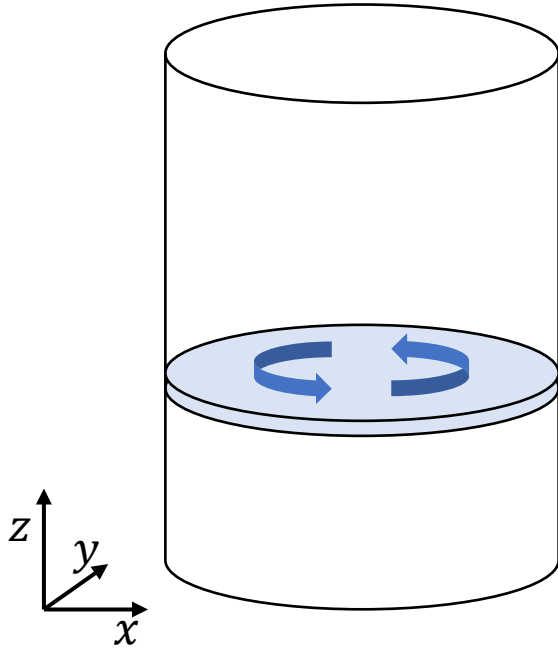
Hydrodynamic Equations (Continuity Equations and Constitutive Relations)

$$\nabla_\mu \Theta^{\mu t} = 0, \quad \nabla_\mu \Theta^{\mu x} = 0, \quad \nabla_\mu \Theta^{\mu y} = 0, \quad \nabla_\mu J^{\mu xy} = 0$$

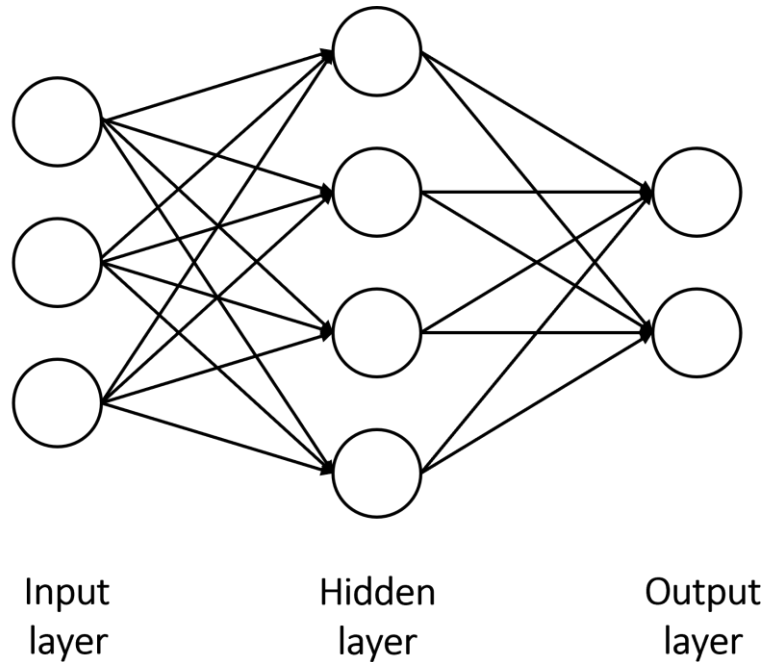
$$\tau_\phi \Delta^r_\alpha \Delta^\theta_\beta D\phi^{xy} = -2\gamma \rho^{xy} - \phi^{xy} - \frac{2\tau_\phi}{3} (\nabla \cdot u) \phi^{xy}$$

(Muller-Israel-Stewart Relaxation)

I. Muller (1967); W. Israel (1976); W. Israel and J. M. Stewart (1976,1979)



Specific Design of PINNs



Input (2 neurons) : t, r

Output (5 neurons) : NN_1, NN_2, \dots, NN_5

e. g.
$$e(t, r) \approx e(t = 0, r) + NN_1(t, r) - NN_1(t = 0, r)$$

Loss Function : $L(\psi) = L_{\text{Hydro}}(\psi) + L_{\text{B.C.}}(\psi) + L_{\text{AM}}(\psi)$

e. g. simplest loss form

$$L(\psi) = \sum_i \langle R_{\text{Hydro},i} \rangle^2 + \sum_i \langle R_{\text{B.C.},i} \rangle^2 + \sum_i \langle R_{\text{AM},i} \rangle^2$$

$\langle \rangle$: spacetime average

Residual

- $R_{\text{Hydro},i}$: hydrodynamic equations, e. g. $\nabla_\mu \Theta^{\mu\nu}(t, r)$
- $R_{\text{B.C.},i}$: Dirichlet boundary condition, e. g. $\partial_r u^r(t, r = R)$
- $R_{\text{AM},i}$: local AM conservation, $\nabla_\mu J^{\mu xy}(t, r)$
 global AM conservation, $\int dr r J^{txy}(t, r) - \int dr r J^{txy}(t = 0, r)$

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 - 3-2. Application
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Verification: Setup

Calculation Setup

- Hidden layer: 1
- Neuron (hidden layer): 250
- Collocation points: 75,000
- Activation function: tanh
- Optimizer: Adam

Physics Setup

- Initial condition

$$e(t = 0, r) = \bar{e} ,$$

$$u^r(t = 0, r) = 0 ,$$

$$u^\theta(t = 0, r) = \delta_1 \cdot \sin^4 \left(\frac{\pi r}{R} \right) ,$$

$$S^z(t = 0, r) = 0 ,$$

$$\phi^{r\theta}(t = 0, r) = 0 ,$$

- Boundary condition

Dirichlet type:

All variables for $r = 0$ and $u^r, S^z, \phi^{r\theta}$ for $r = R$

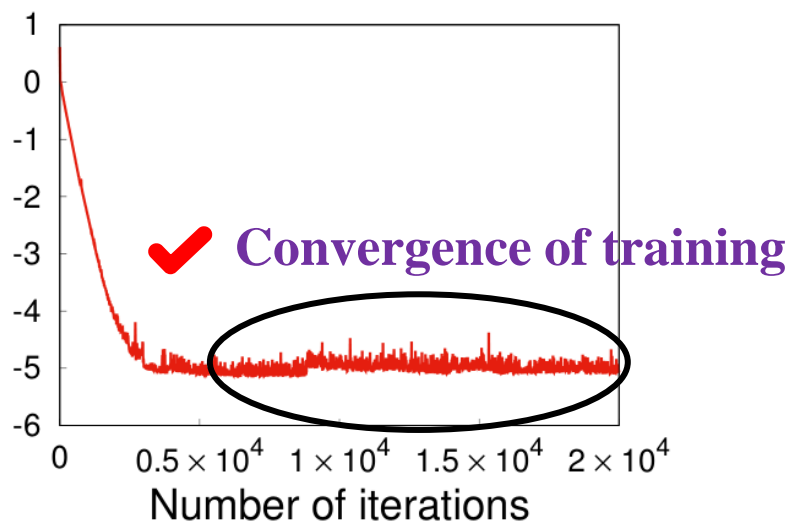
Neumann: $\partial_r u^r, \partial_r \phi^{r\theta}$ for $r = R$

- Parameters

\bar{e}	1
R	1
t_{\max}	0.15
γ_ϕ	2
τ_ϕ	2
δ_1	0.2

Verification: Results

Loss function
 $L(\psi)$

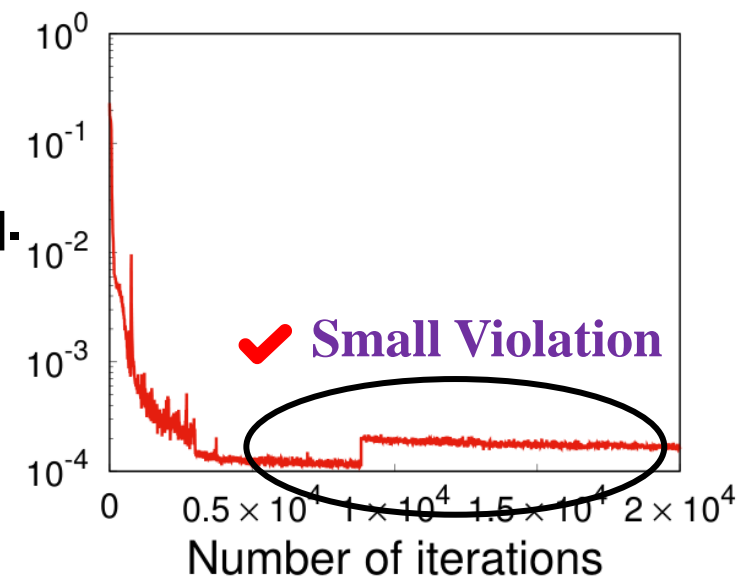


$$\begin{aligned} e(t=0, r) &= \bar{e}, \\ u^r(t=0, r) &= 0, \\ u^\theta(t=0, r) &= \delta_1 \cdot \sin^4\left(\frac{\pi r}{R}\right), \\ S^z(t=0, r) &= 0, \\ \phi^{r\theta}(t=0, r) &= 0, \end{aligned}$$

\bar{e}	1
R	1
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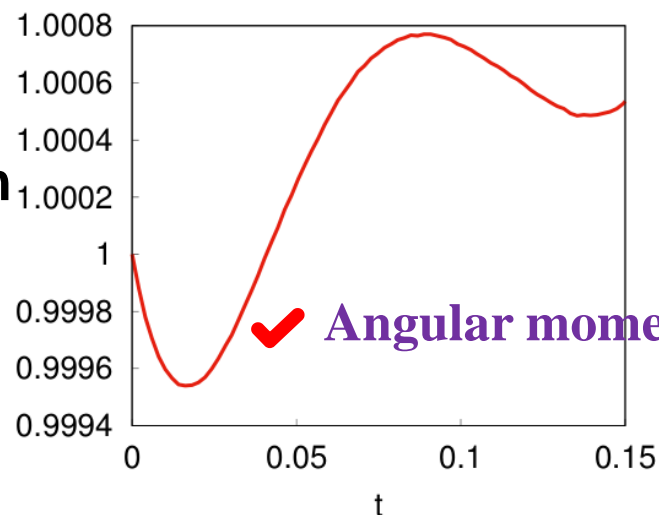
Residual for Hydro eq.

$$\sum_i \frac{\int dr (2\pi r) |R_{\text{Hydro}, i}|^2}{\text{Volume}}$$



Net angular momentum

$$\frac{\int dr r J^{txy}(t, r)}{\int dr r J^{txy}(t=0, r)}$$



Application: Setup

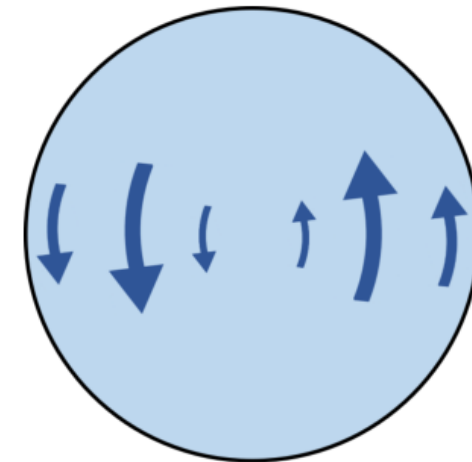
Case 1. Orbital Initial Condition

- ✓ Motivated by QGP
in high-energy heavy-ion collisions

Z.-T. Liang and X.-N. Wang (2005).

$$\begin{aligned}e(t=0, r) &= \bar{e}, \\u^r(t=0, r) &= 0, \\u^\theta(t=0, r) &= \delta_1 \cdot \sin^4\left(\frac{\pi r}{R}\right), \\S^z(t=0, r) &= 0, \\\phi^{r\theta}(t=0, r) &= 0.\end{aligned}$$

\bar{e}	1
R	1
t_{\max}	0.4
γ_ϕ	0, 2
τ_ϕ	γ_ϕ
δ_1	0.2



t = 0

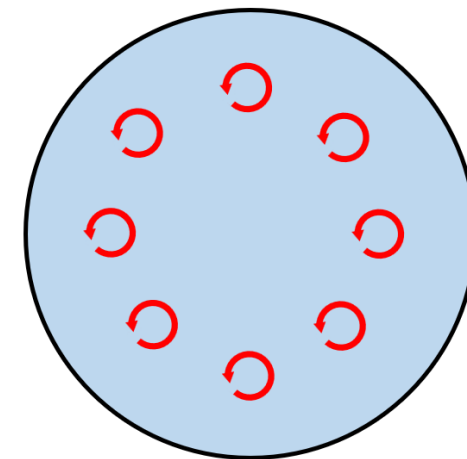
Case 2. Spin Initial Condition

- ✓ Motivated by Einstein-de Haas effect

A. Einstein and W. De Haas (1915).

$$\begin{aligned}e(t=0, r) &= \bar{e}, \\u^r(t=0, r) &= 0, \\u^\theta(t=0, r) &= 0, \\S^z(t=0, r) &= \delta_2 \cdot \sin^4\left(\frac{\pi r}{R}\right), \\\phi^{r\theta}(t=0, r) &= 0.\end{aligned}$$

\bar{e}	1
R	1
t_{\max}	0.4
γ_ϕ	2
τ_ϕ	2
δ_2	0.2

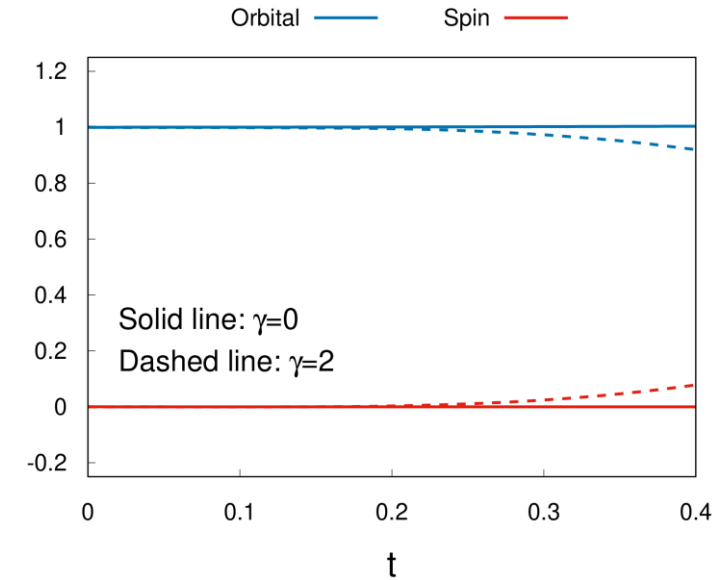


t = 0

Application: Net AM

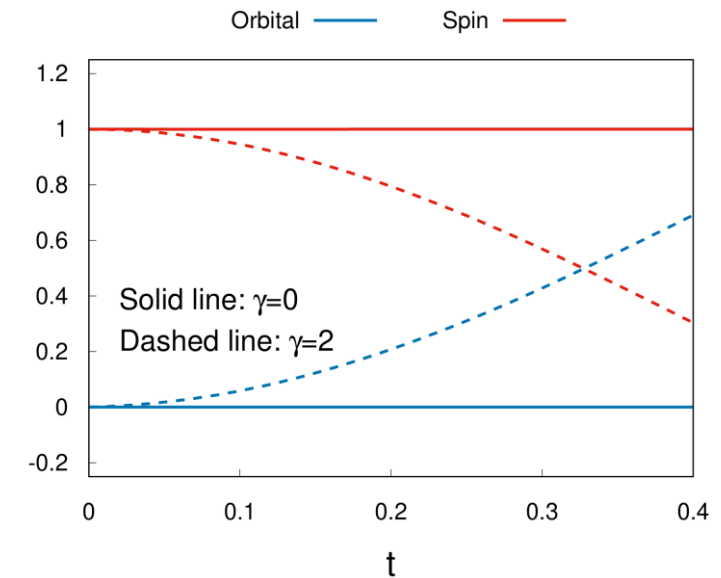
Case 1. Orbital Initial Condition

✓ Orbital-to-Spin AM conversion at $\gamma = 2$



Case 2. Spin Initial Condition

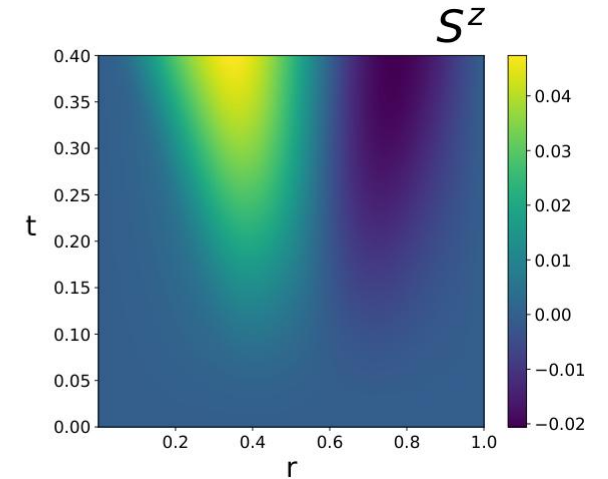
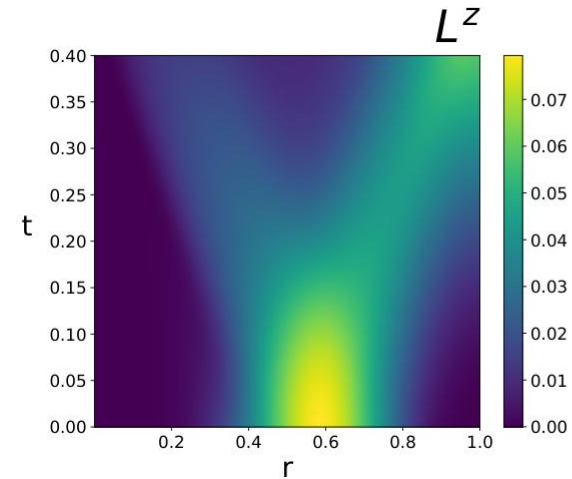
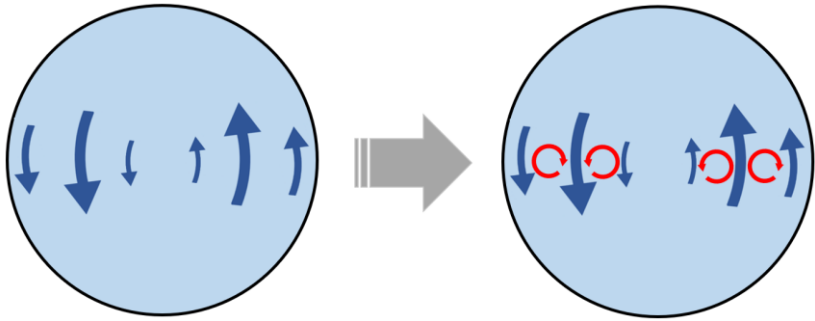
✓ Spin-to-Orbital AM conversion at $\gamma = 2$



Application: Heatmap of AM

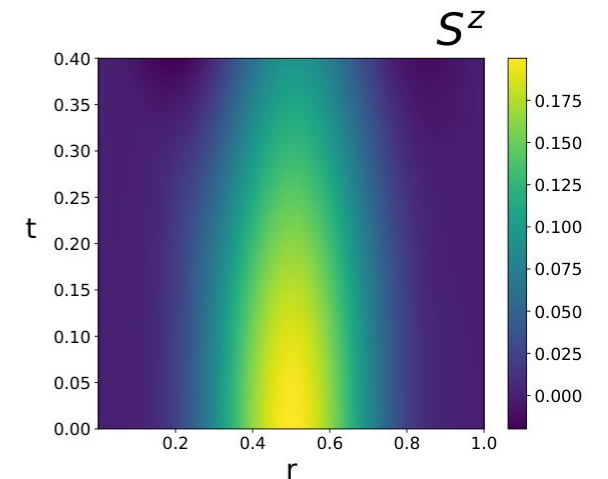
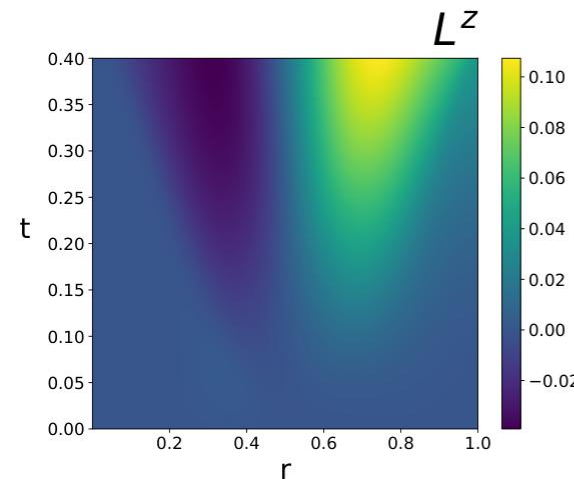
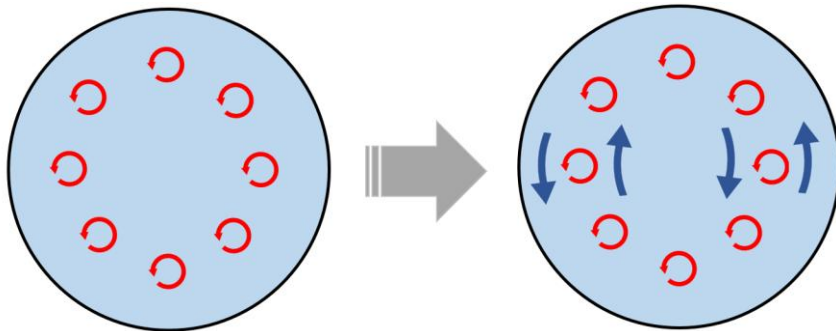
Case 1. Orbital Initial Condition

✓ Generated spin has opposite signs
for $r > 0.5R$ and $r < 0.5R$



Case 2. Spin Initial Condition

✓ Generated orbital spin has opposite signs
for $r > 0.5R$ and $r < 0.5R$



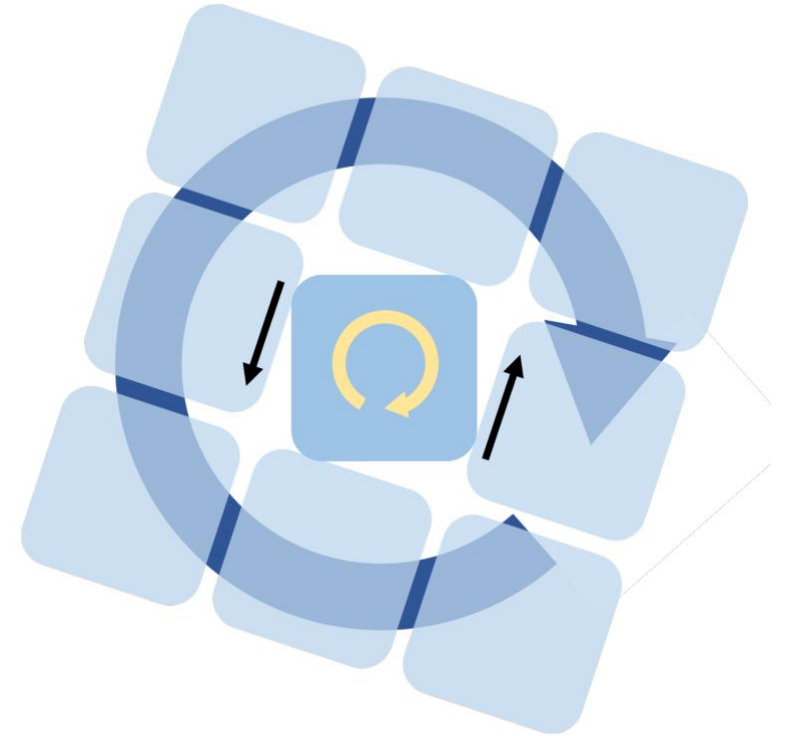
Application: Discussion

- **Rotational viscous correction** is generated by the difference between **transversely projected thermal vorticity** (ϖ_{\perp}) and **spin potential** ($\omega^{\mu\nu}$)

$$\tau_{\phi} \Delta^r_{\alpha} \Delta^{\theta}_{\beta} D\phi^{\alpha\beta} = -2\gamma \rho^{r\theta} - \phi^{r\theta} - \frac{2\tau_{\phi}}{3} \theta \phi^{r\theta}$$

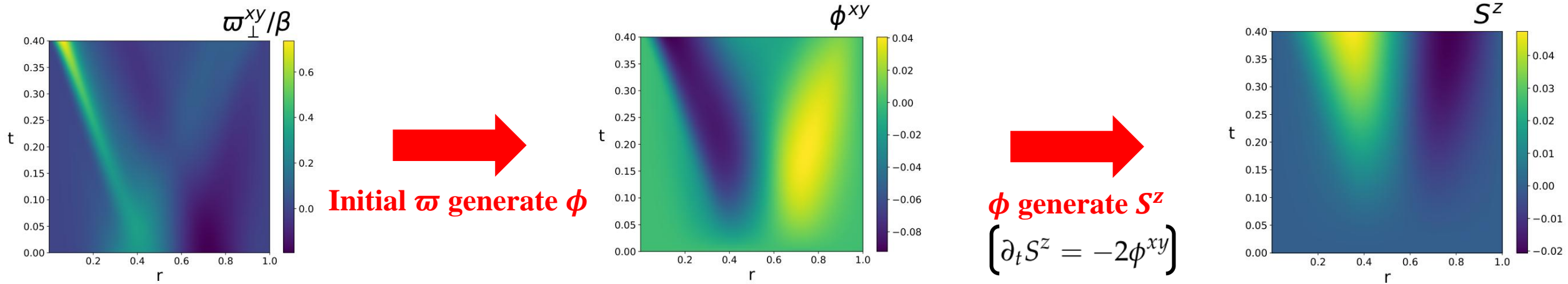
source term

$$\rho^{\mu\nu} = \Delta^{\mu\alpha} \Delta^{\nu\beta} \left(\partial_{[\alpha} \beta_{\beta]} - 2\omega_{\alpha\beta} \right) = \beta^{-1} \omega_{\perp}^{\mu\nu} - 2\omega^{\mu\nu}$$



Application: Conversion Process

Case 1. Orbital Initial Condition



Case 2. Spin Initial Condition

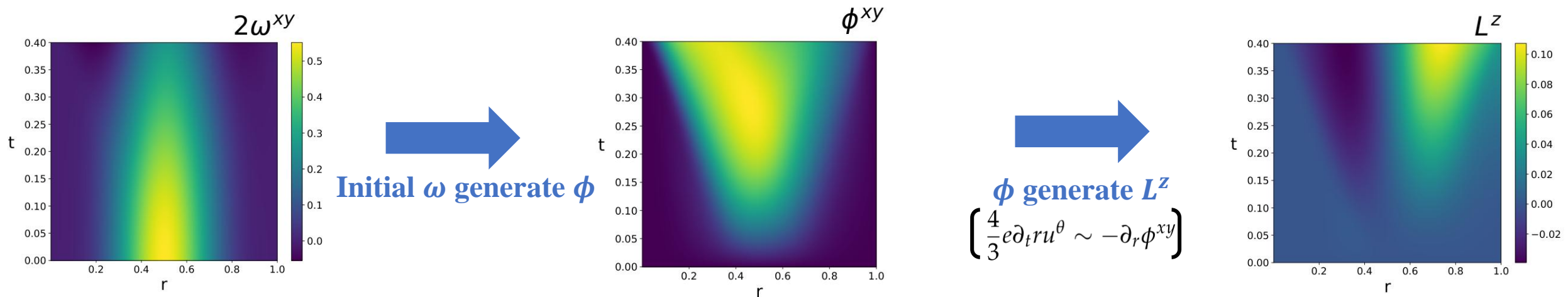


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Summary & Outlook

Summary

- We have developed a PINNs-based framework for relativistic spin hydrodynamics.
- Angular momentum conservation is found to be successfully enforced in the training process.
- Our proof-of-concept simulations have demonstrated orbital–spin angular momentum conversion, which is driven by the rotational effect.
- Our demonstration has illustrated that the rotational viscous effect in second-order viscous hydrodynamics can be interpreted as a friction between a rotating fluid cell and the surrounding rotational flow.

Outlook

- Compare with linear analyses to deepen the understanding of nonlinear effects
- Examine system behavior when treating spin potential and thermal vorticity at zeroth order, or when incorporating second-order effects such as those induced by a boosted heat current
- Apply the framework to the analysis of high-energy heavy-ion collisions (long-term goal)