



## **Collision Energy Dependence of the Light- and** Hyper-Nuclei Directed Flow in 3.0-4.5 GeV Au+Au **Collisions at RHIC-STAR**

Chengdong Han

Institute of Modern Physics, CAS chdhan@impcas.ac.cn

2024/11/03



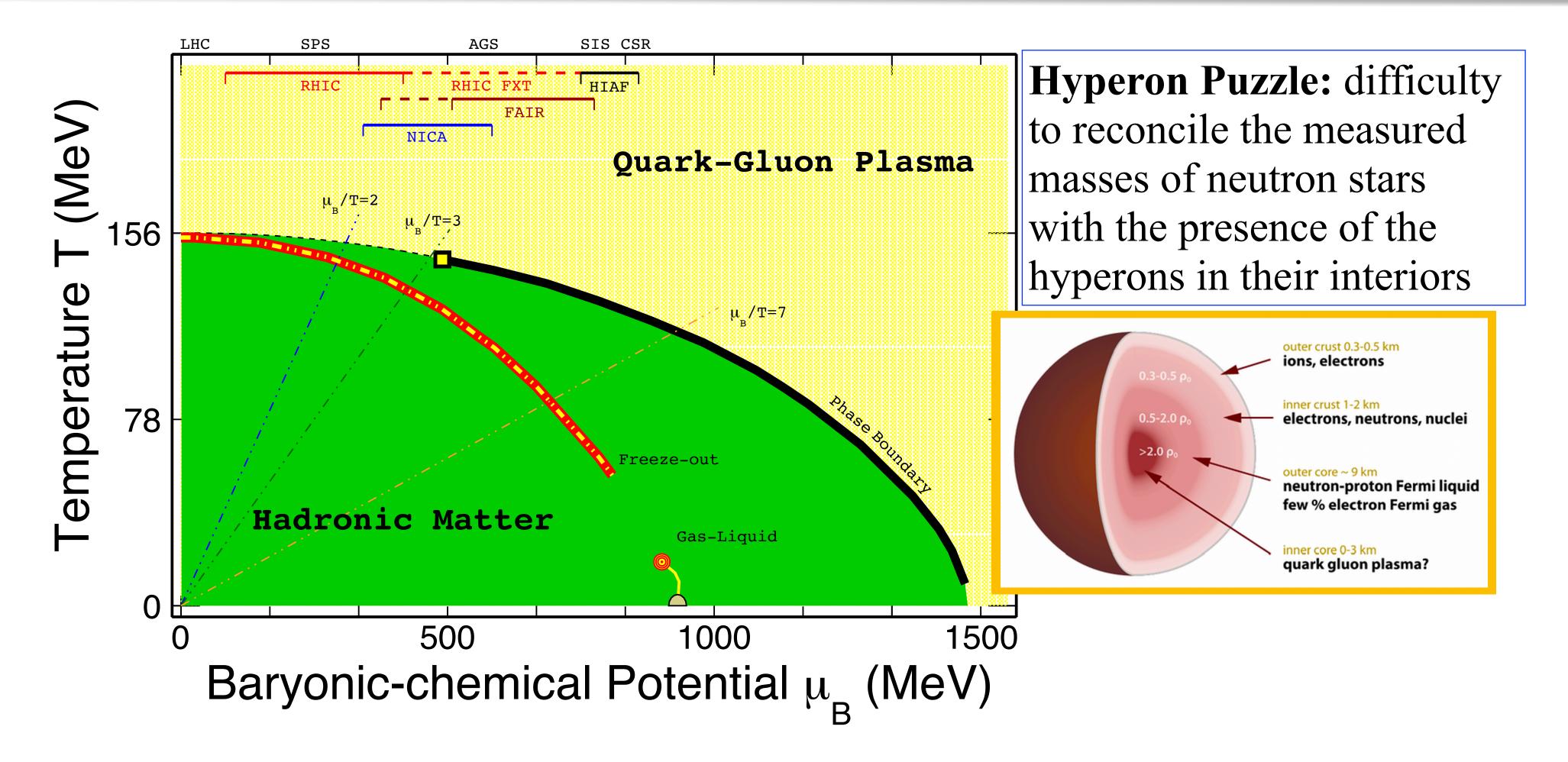
The 1st International Workshop on Physics at High Baryon Density (PHD2024)



- 1. Introduction
- 2. Dataset and Particle Reconstruction
- 3. Light- and Hyper-Nuclei Collective Flow i. Light- and Hyper-Nuclei Directed Flow  $v_1$ ii. Mass and Energy Dependence of  $v_1$
- 4. Summary and Outlook



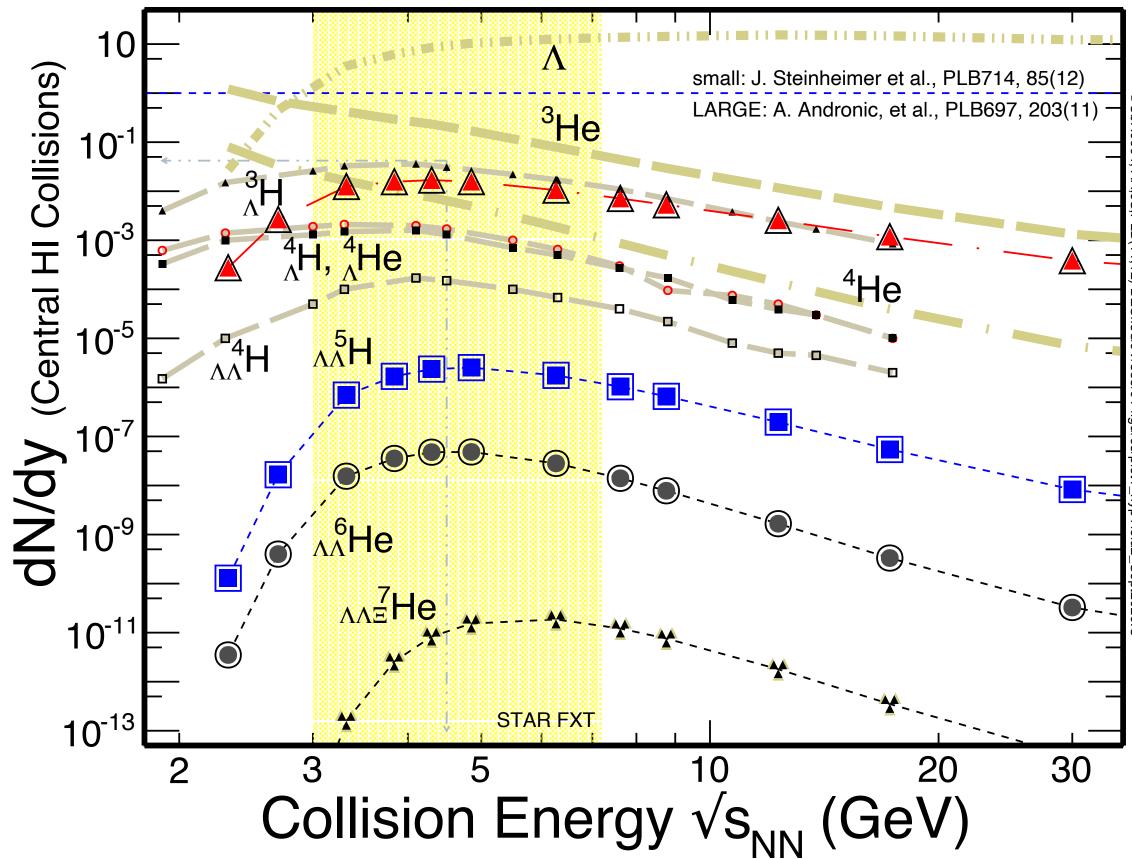
## High-Energy Nuclear Collisions and QCD Phase Diagram



1) RHIC beam energy scan  $\rightarrow$  search for 1<sup>st</sup>-order phase transition and QCD critical point 2) Baryon-baryon interaction (e.g. N-N, Y-N)  $\rightarrow$  inner structure of compact stars



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[1] A. Andronic *et al.* Phys.Lett.B 697, 203 (2011) [2] J. Steinheimer *et al.* Phys.Lett.B 714, 85 (2012)

### **Light- and Hyper-nuclei Productions**

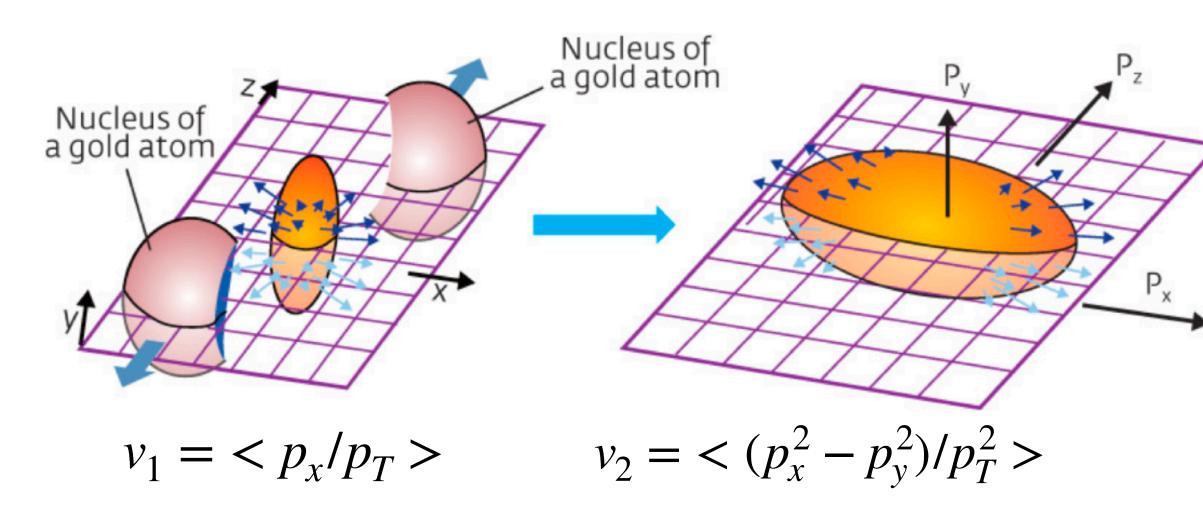
- 1) Light- and Hyper-Nuclei are abundantly produced at high baryon density region
- 2) Light-Nuclei carry information about local baryon density fluctuations at freeze-out; offers insights on the Final State Interaction: N-N
- 3) Hyper-Nuclei provide access to the hyperon-nucleon interaction: Y-N
- Collective flow is sensitive to the Equation-of-State of nuclear matter





### **Collective Flow**

#### Heavy ion collisions: Initial spatial anisotropy → Pressure gradient → Anisotropic flow



- The initial pressure gradient of the collision system is directly 1) 2) Fit  $dN/d(\phi - \Psi_n)$  distribution in rapidity bins to extract related to the magnitude of  $v_{n}$ , which is a sensitive observable observed flow coefficients  $v_n^R$ for studying EoS
- 2) Collectivity of light- and hyper-nuclei in heavy-ion collisions at high baryon density regions is important for understanding their formation mechanism

[1] H. Masui et al., Nucl. Instrum. Methods Phys. Res. A 833, 181 (2016)

[2] A. M. Poskanzer and S. A. Voloshin, Phys. Rev. C 58, 1671 (1998)

#### **Collective Flow analysis**

The n<sup>th</sup> order coefficient of the fourier expansion of the azimuthal distribution in the momentum space

$$\boldsymbol{E}\frac{\boldsymbol{d}^{3}N}{\boldsymbol{d}p^{3}} = \frac{1}{2\pi}\frac{\boldsymbol{d}^{2}N}{\boldsymbol{p}_{T}\boldsymbol{d}\boldsymbol{p}_{T}\boldsymbol{d}\boldsymbol{y}}\left(1 + \sum_{1}^{\infty} 2v_{n}\cos\left[\boldsymbol{n}\left(\boldsymbol{\phi}-\boldsymbol{\psi}_{RP}\right)\right]\right)$$

— v<sub>2</sub> Elliptic flow  $-v_1$  Directed flow

#### Analysis steps with event plane method

Signal extraction for a given  $\phi - \Psi_n$  bin: 1)

$$N^{R}(\phi - \psi_{n}) = \int dM \frac{1}{R_{n}} \frac{dN}{d(\phi - \psi_{n})}$$

3) Correct  $v_n^R$  with signal number weighted EP resolution  $\langle \frac{1}{2} \rangle$  n 1  $R_n$  $\langle v_n \rangle = \langle v_n^R \rangle \langle \frac{1}{R} \rangle$  $<\frac{1}{R_{i}}>=\frac{\sum_{i}N_{i}*<\frac{1}{R_{n}}^{\prime\prime}>}{R_{i}}$ 

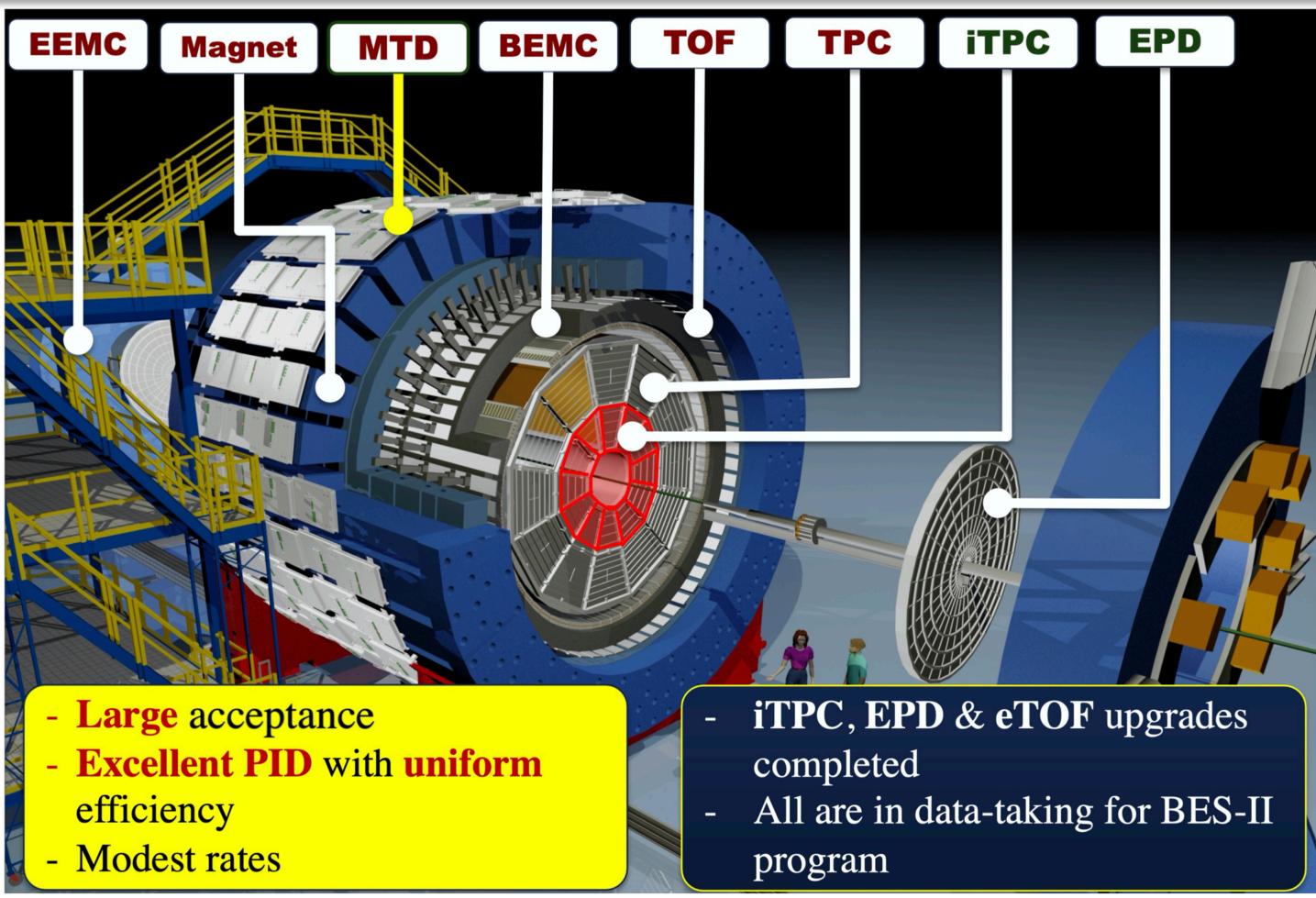








### **STAR Detector System**



- 1) Enlarged rapidity acceptance
- 2) Improved particle identification
- 3) Enhanced event plane resolution

### Major Upgrades in BES-II:

#### **iTPC:**

- Improves dE/dx
- Extends  $\eta$  coverage from 1.0 to 1.5
- Lowers p<sub>T</sub> cut-in from 125 to 60 MeV/c
- Ready in 2019

#### eTOF:

- Forward rapidity coverage
- PID at  $\eta = -1.1$  to -1.6
- Ready in 2019

#### **EPD:**

- Improves trigger
- Event plane measurements
- Ready in 2018

[1] iTPC: <u>https://drupal.star.bnl.gov/STAR/starnotes/</u>. public/sn0619. [2] eTOF: STAR and CBM eTOF group, arXiv: 1609.05102. [3] EPD: J. Adams, et al. NIM A968, 163970 (2020)





### **Dataset and Event Plane Reconstruction**

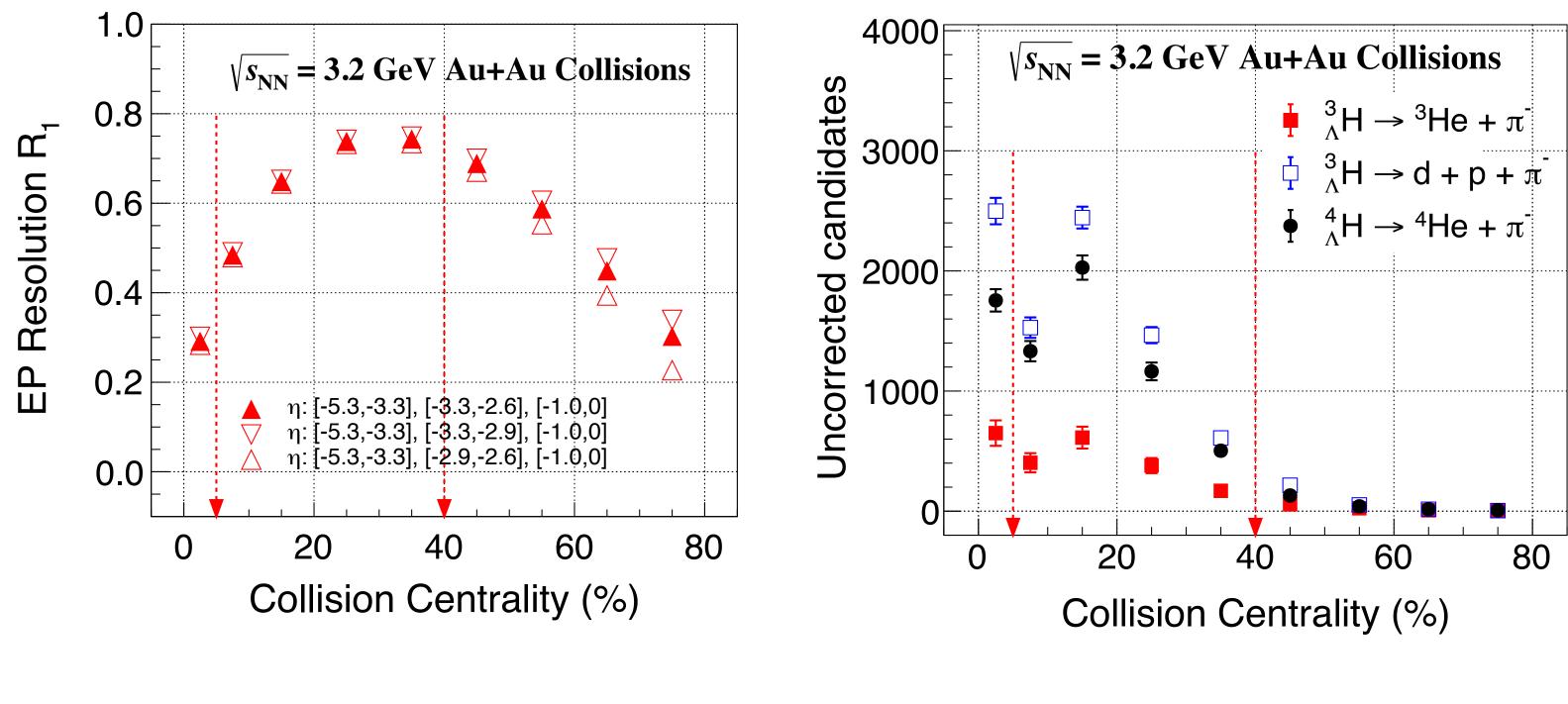
DataSet	√SNN = <b>3.0 GeV</b> (2018) (ytarget = -1.04)	3.2 GeV (2019) (y <sub>target</sub> = -1.14)	3.5  GeV (2020) ( $y_{target} = -1.25$ )	3.9 GeV (2020) (y <sub>target</sub> = -1.37)	4.5 GeV (2020 (y <sub>target</sub> = -1.52
Events	~260 M	~210 M	~115 M	~120 M	~120 M

#### • Event plane reconstruction

- EP reconstruction: Q vector method
- Re-center and shift calibration
- EP resolution: three sub-events method

The EP resolution is determined as:

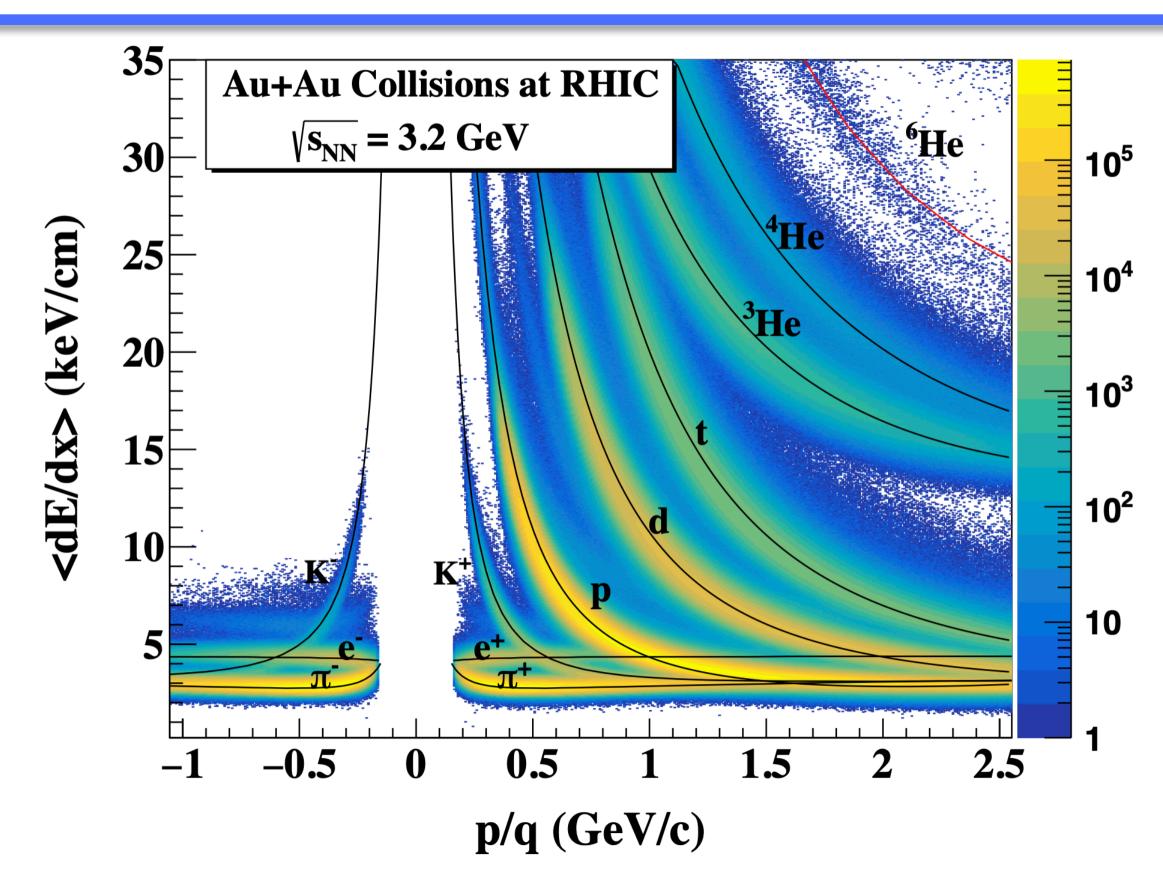
$$R_{1} = \langle \cos(\Psi_{1}^{a} - \Psi_{r}) \rangle$$
$$= \sqrt{\frac{\langle \cos(\Psi_{1}^{a} - \Psi_{1}^{b}) \rangle \langle \cos(\Psi_{1}^{a} - \Psi_{1}^{c}) \rangle}{\langle \cos(\Psi_{1}^{b} - \Psi_{1}^{c}) \rangle}}$$



▶ 5-40% centrality bin used in this analysis

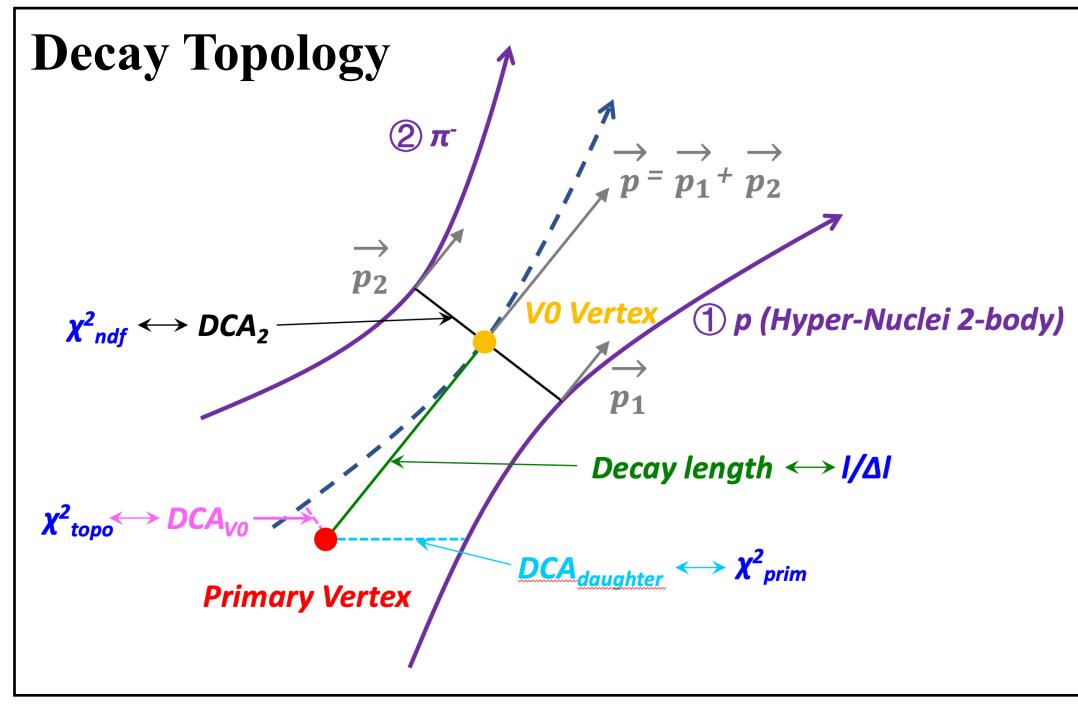


### **Particle Identification and Topological Selection**



- Good particle identification capability based on TPC and TOF
- providing a full set of the particle parameters together with their uncertainties
- 3)

[1] Gorbunov and I. Kisel, Reconstruction of decayed particles based on the Kalman filter. CBM-SOFT-note-2007-003, 7 May 2007 [2] Ivan Kisel. Event Topology Reconstruction in the CBM Experiment. J. Phys. Conf. Ser. **1070(1)**, 012015 (2018)

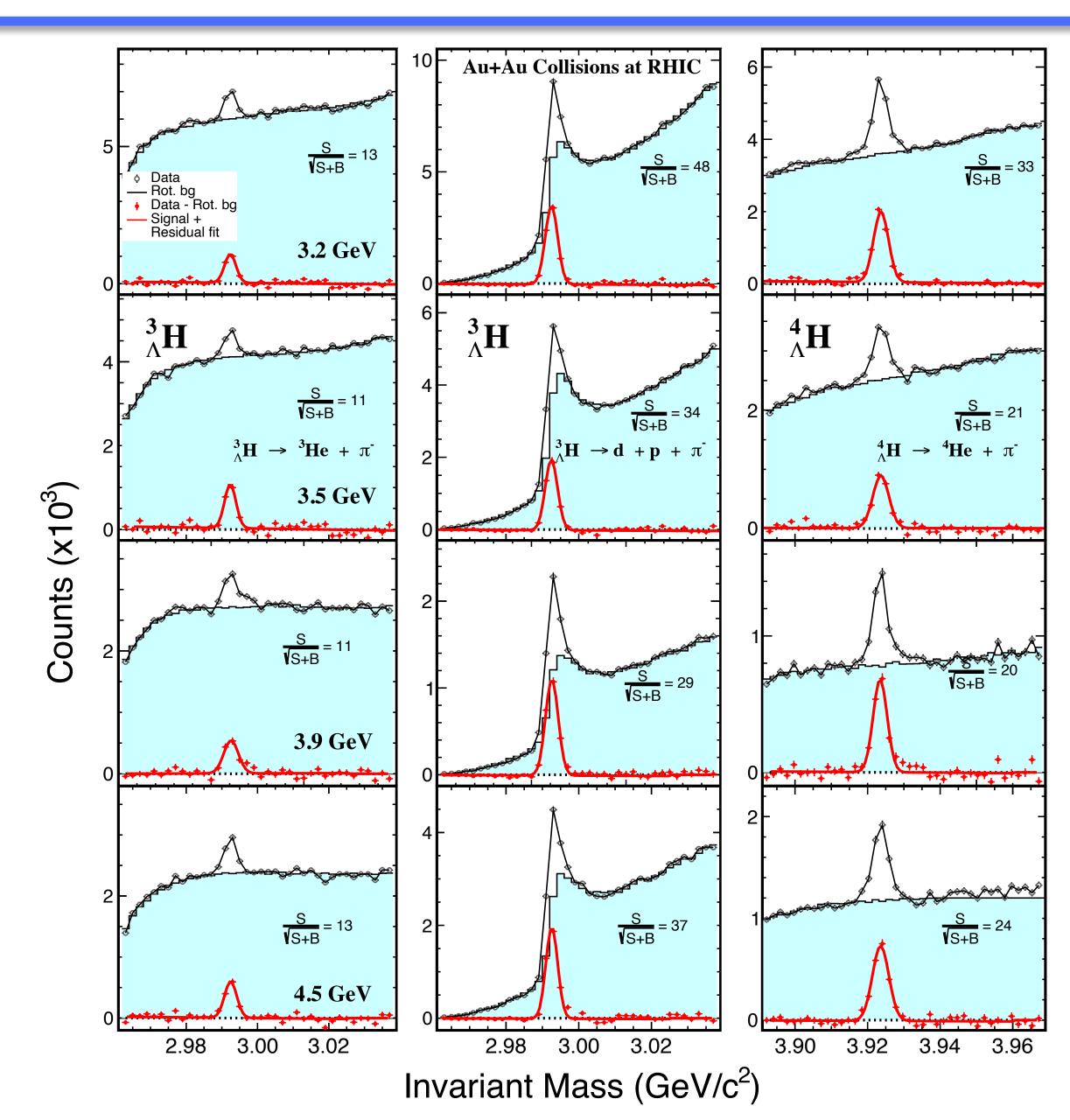


The hyper-nuclei reconstruction with KFParticle package based on the Kalman filter method

Decay topology tremendously helped on particle identification and background suppression



### Hyper-Nuclei Reconstruction

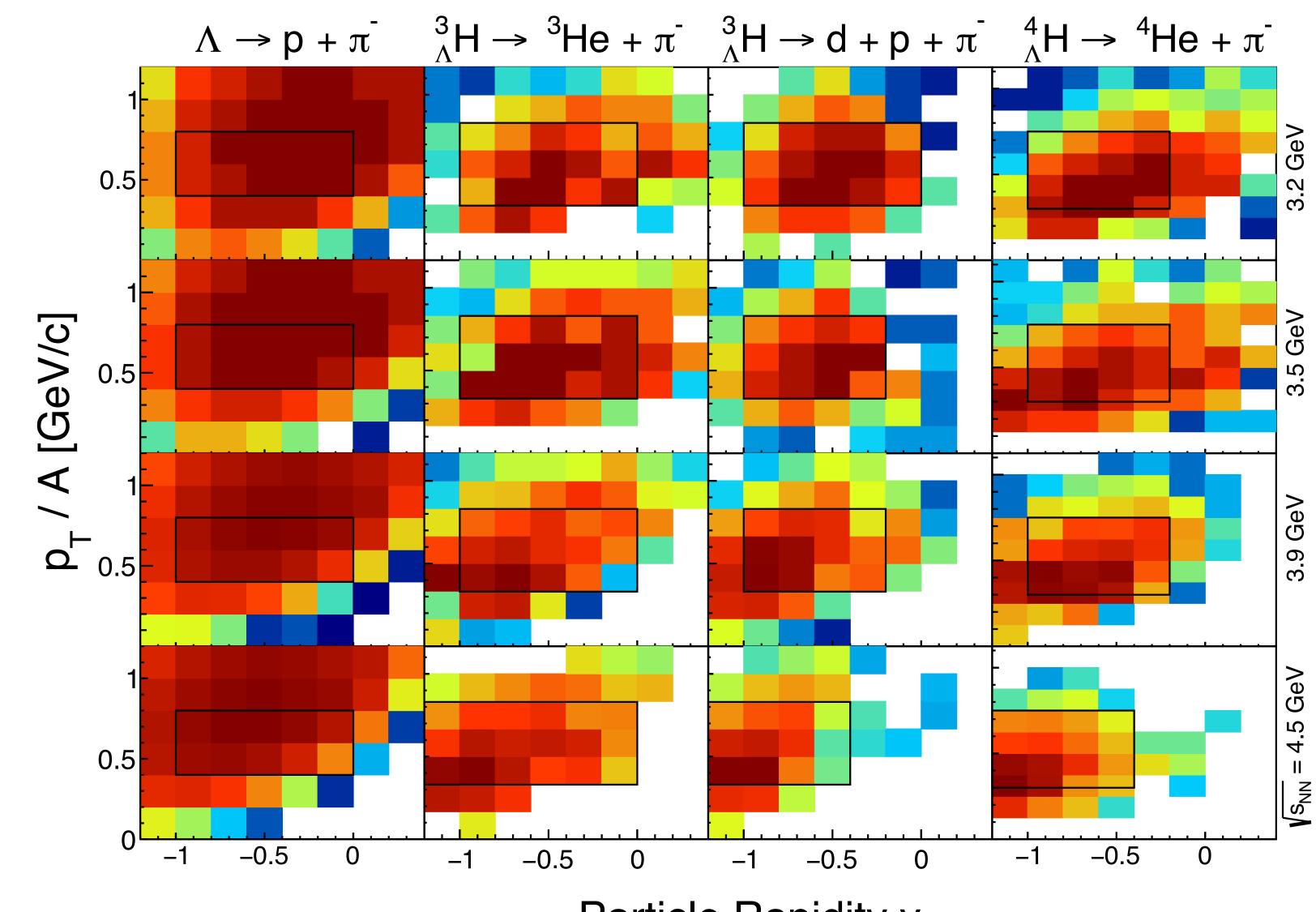


- 1) Topologically reconstructed  ${}^{3}{}_{\Lambda}$ H 2-body decay (left panel), 3-body decay (middle panel) and  ${}^{4}{}_{\Lambda}$ H (right panel) from 3.2, 3.5, 3.9 and 4.5 GeV Au+Au collisions
- Background subtracted distributions are shown as red symbols. The significances of the mass peaks are also indicated
- Obvious hyper-nuclei signals can be observed with the reconstructed invariant mass distributions





### Acceptance of A, ${}^{3}_{\Lambda}$ H and ${}^{4}_{\Lambda}$ H of 3.2, 3.5, 3.9, 4.5 GeV



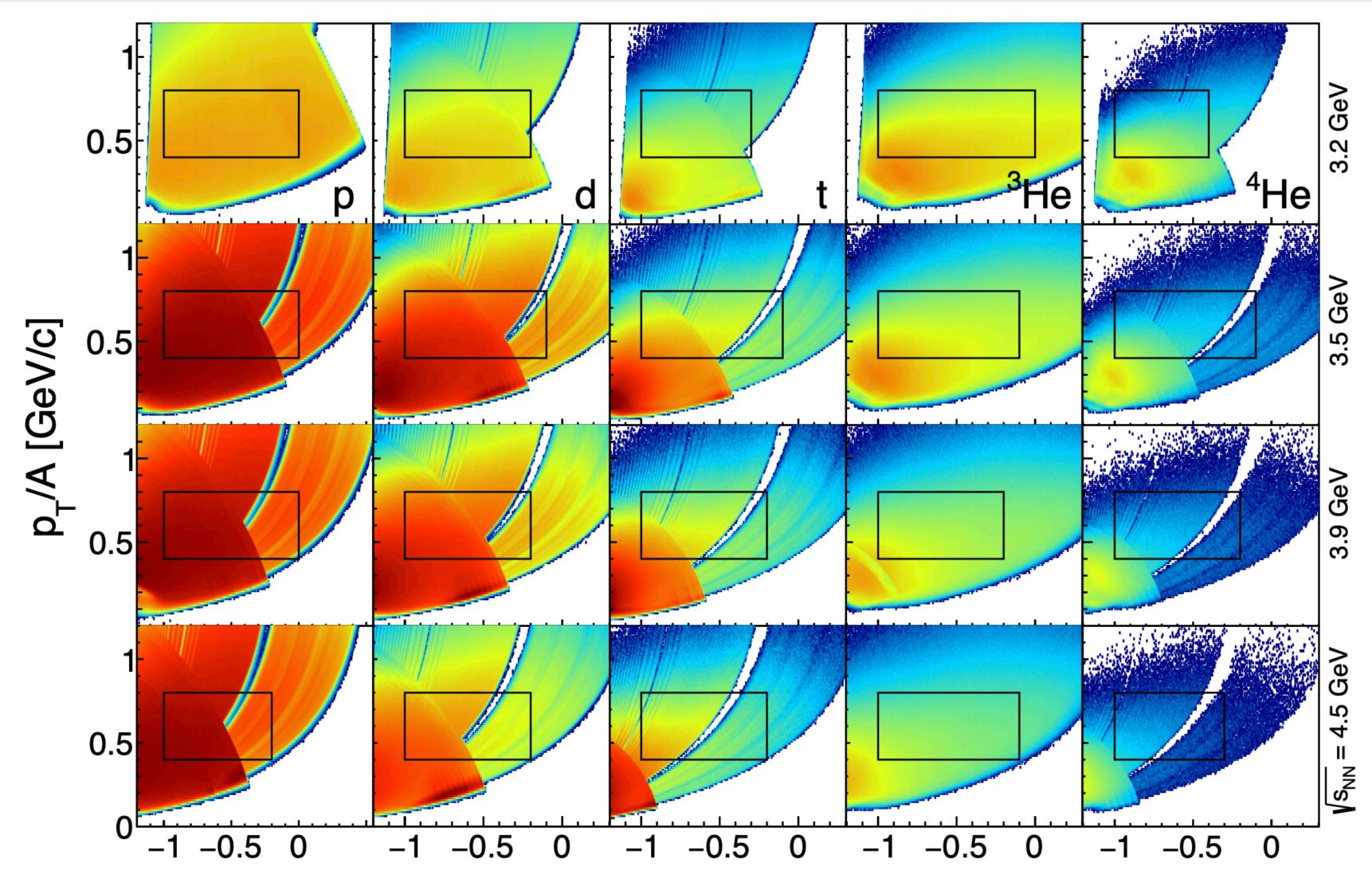
Particle Rapidity y

Directed flow of hypernuclei are calculated within the selected  $0.4 \leq$  $p_T / A \leq 0.8 \text{ GeV/c range}$ as indicated by the boxes





### Acceptance of p, d, t, <sup>3</sup>He and <sup>4</sup>He of 3.2, 3.5, 3.9, 4.5 GeV



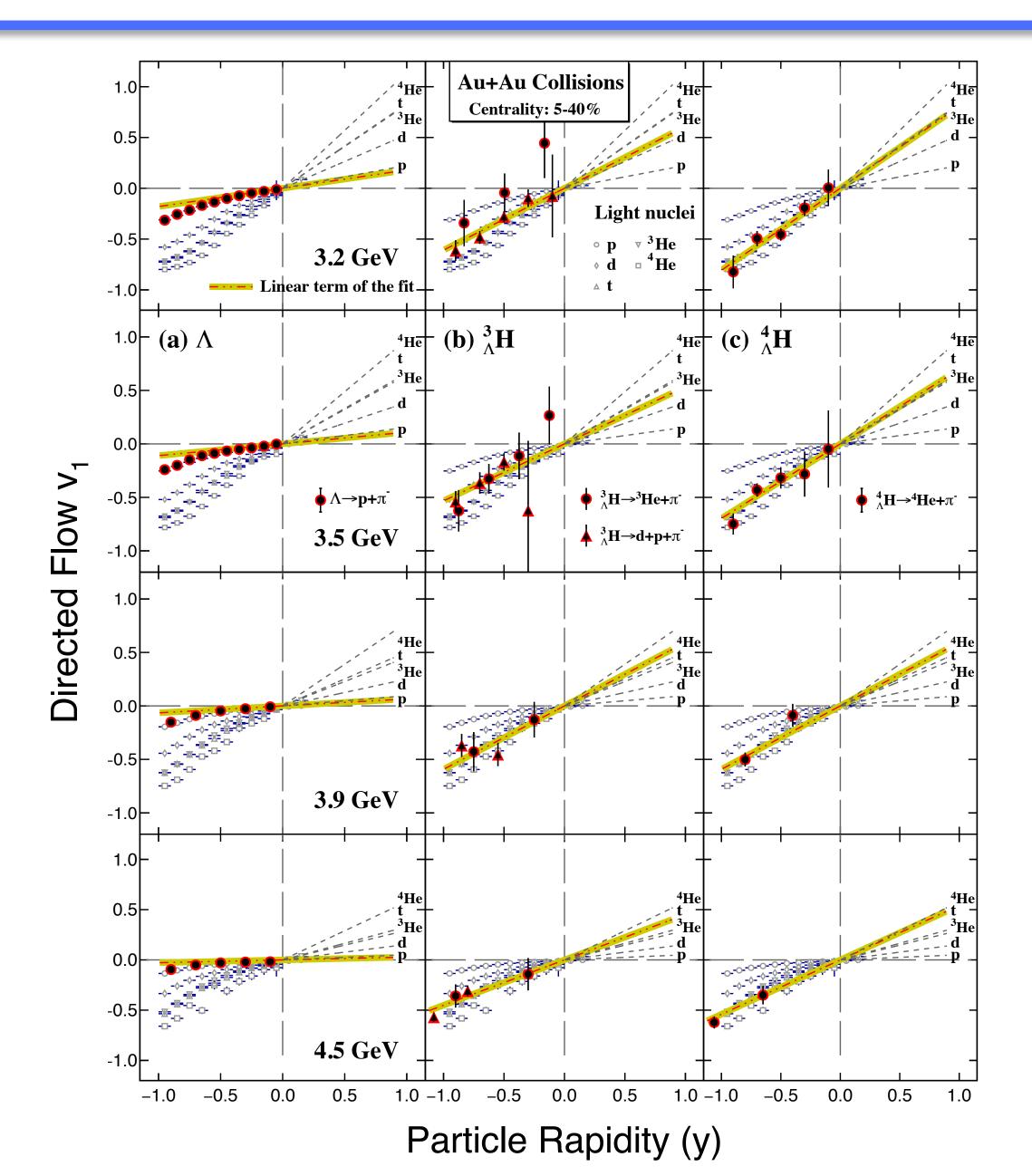
Particle Rapidity y

Directed flow of lightnuclei are calculated within the selected 0.4 < $p_T/A < 0.8 \text{ GeV/c range}$ as indicated by the boxes





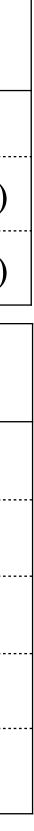
### Light- and Hyper-Nuclei Directed Flow v<sub>1</sub>



The v<sub>1</sub> slope is obtained by fitting the v<sub>1</sub>(y) distribution with a polynomial function, where  $p_0$  is the mid-rapidity slope  $(v_1)^s = dv_1/dy\Big|_{v=0}$ 

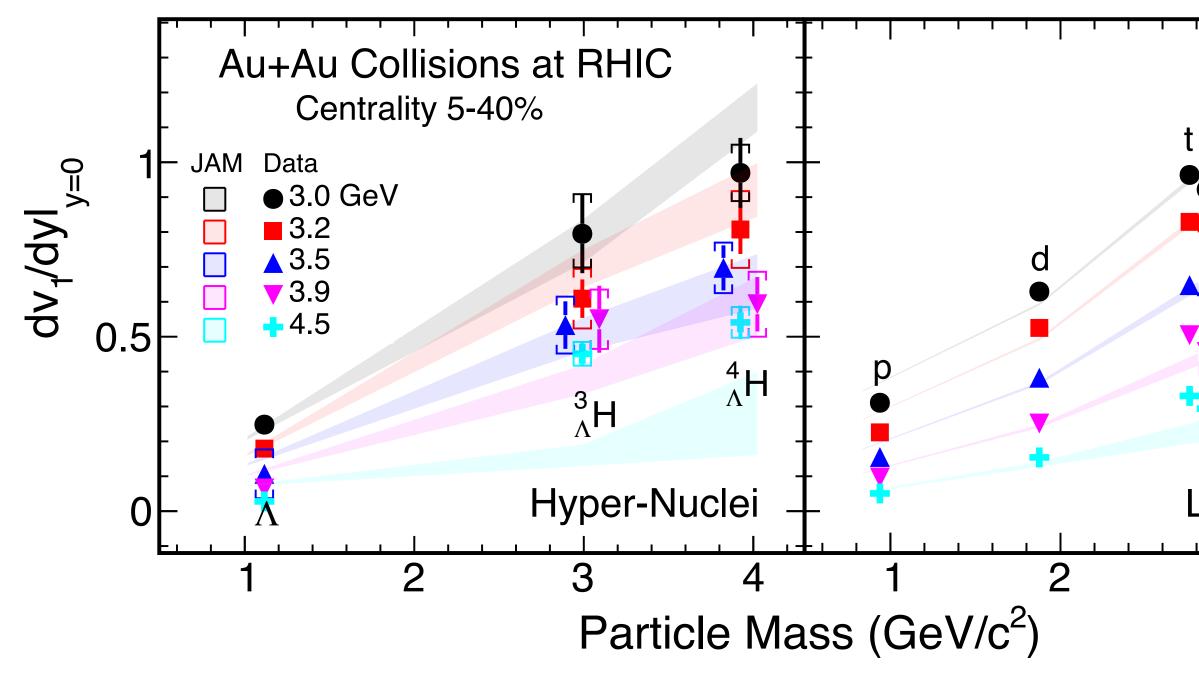
Hyper- Nuclei	Fitting Function	У	p <sub>T</sub> /A	
Λ	$v_1(y) = (v_1)^s \cdot y + p_1 \cdot y_3$	(-1.0, 0.0)	(0.4, 0.8)	
<sup>3</sup> <sub>A</sub> H	$\mathbf{v}_1(\mathbf{y}) = (\mathbf{v}_1)^{\mathbf{s}} \cdot \mathbf{y}$	(-1.0, 0.0)	(0.33, 0.83)	
$^{4}\Lambda H$	$\mathbf{v}_1(\mathbf{y}) = (\mathbf{v}_1)^{\mathbf{s}} \cdot \mathbf{y}$	(-1.0, 0.0)	(0.30, 0.75)	
Light- Nuclei	Fitting Function	У	p <sub>T</sub> /A	
р	$\mathbf{v}_1(\mathbf{y}) = (\mathbf{v}_1)^{\mathbf{s}} \cdot \mathbf{y} + \mathbf{p}_1 \cdot \mathbf{y}_3$	(-1.0, 0.0)	(0.4, 0.8)	
d	$\mathbf{v}_1(\mathbf{y}) = (\mathbf{v}_1)^{\mathbf{s}} \cdot \mathbf{y} + \mathbf{p}_1 \cdot \mathbf{y}_3$	(-1.0, -0.2)	(0.4, 0.8)	
t	$\mathbf{v}_1(\mathbf{y}) = (\mathbf{v}_1)^{\mathbf{s}} \cdot \mathbf{y} + \mathbf{p}_1 \cdot \mathbf{y}_3$	(-1.0, -0.3)	(0.4, 0.8)	
<sup>3</sup> He	$\mathbf{v}_1(\mathbf{y}) = (\mathbf{v}_1)^{\mathbf{s}} \cdot \mathbf{y} + \mathbf{p}_1 \cdot \mathbf{y}_3$	(-1.0, 0.0)	(0.4, 0.8)	
<sup>4</sup> He	$\mathbf{v}_1(\mathbf{y}) = (\mathbf{v}_1)^{\mathbf{s}} \cdot \mathbf{y} + \mathbf{p}_1 \cdot \mathbf{y}_3$	(-1.0, -0.4)	(0.4, 0.8)	

[1] M.S. Abdallah *et al.*, (STAR Collaboration), Phys. Lett. B 827, 136941 (2022)
[2] B. E. Aboona *et al.*, (STAR Collaboration), Phys. Rev. Lett. 130, 211301(2023)





### **Particle Mass Dependence**

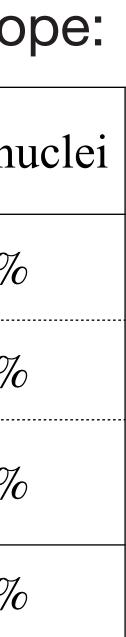


- 1) process for the light- and hyper-nuclei production
- afterburner calculations

<sup>4</sup> He	Systematic	c uncertai	inties fo	r v1 slo
t <sup>•</sup> <sup>3</sup> He	Major source	$^{3}\Lambda H$	$^{4}\Lambda H$	light-n
	EP resolution	4%	4%	4%
	Efficiency	2%	2%	2%
Light-Nuclei	Topological cuts / PID cuts	12%	11%	5%
3 4	Total	13%	12%	6%

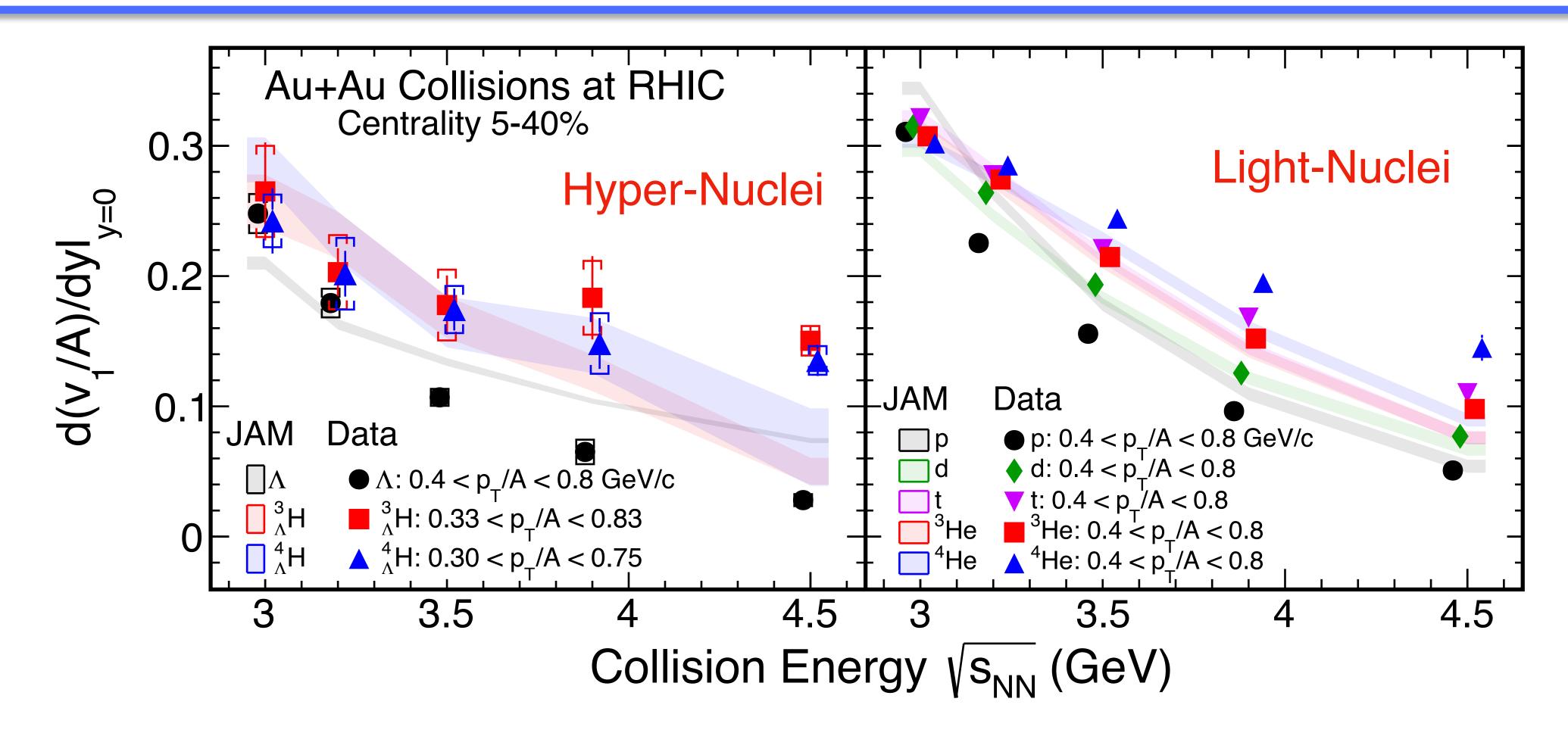
At given energy, for both light- and hyper-nuclei, it seems that the slopes of mid-rapidity  $v_1$  are scaled with atomic mass number A or/and particle mass, implying coalescence is the dominant

2) The feature is also reproduced by transport model JAM (mean field  $\kappa$ = 380 MeV) with coalescence





## **Collision Energy Dependence**



1) As the collision energy increases, the  $v_1$  slope of light- and hyper-nuclei decreases 2) Hadronic transport model (JAM mean field  $\kappa$ = 380 MeV) plus coalescence afterburner calculations are consistent with observed energy dependence



## **Summary and Outlook**

- experiment at RHIC
- measured  $v_1$  slopes of light- and hyper-nuclei at mid-rapidity
- 3) Calculations of hadronic transport model plus coalescence afterburner qualitatively in these heavy ion collisions

### **Outlook:**

2 billion events for 3.0 GeV Au + Au collisions at RHIC-STAR will significantly enhance the precision and help us to further constrain coalescence parameters for both light- and hyper-nuclei at the high density region

1) Studied the collision energy dependence of the directed flow  $v_1$  for both light- and hypernuclei in  $\sqrt{s_{NN}}$  = 3.0, 3.2, 3.5, 3.9, 4.5 GeV Au + Au collisions measured by the STAR

2) An approximate atomic mass number scaling and energy dependence are observed in the

reproduced the observed dependences for hyper-nuclei as well as light-nuclei implying coalescence process dominate the underlying production mechanism for those light clusters







### 中国科学院近代物理研究所

Institute of Modern Physics, Chinese Academy of Sciences

# Thank you for your attention!



