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Probing 2HDM-I light Higgs in the top-pairassociated diphoton channel

Yabo Dong (董亚博)

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on behalf of Dr. Kun Wang and Prof. Jingya Zhu



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Background and motivation



Background and motivation



THDM-I and Parameter scan

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THDM-I and Parameter scan



THDM-I and Parameter scan





Collider simulations



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Collider simulations



the cut flow at HL-LHC with L=300 fb⁻¹

Cute	$\operatorname{Signal}(\sigma \times L)$	$Background(\sigma \times L)$						
Outs	$t\bar{t}h$	$tar{t}\gamma$	$t\bar{t}H$	tjH	$bar{b}\gamma\gamma$	$t\bar{t}\gamma\gamma$	$tj\gamma\gamma$	$W j j \gamma \gamma$
Initial	165	1536000	168	23	2293800	3354	5253	65400
Basic cut	14.5	4577.3	17.25	0.712	848.7	200.0	89.61	268.1
Mass cut	11.03	46.08	0.023	0.0	0.0	6.47	2.67	4.58
Energy cut	10.33	0.0	0.023	0.0	0.0	5.8	1.91	3.27

$$S = \sqrt{2L\left[(S+B)ln\left(1+\frac{S}{B}\right)-S\right]}$$

required luminosity to get 2σ or 5σ significance

2σ	5σ
137 fb ⁻¹	859 fb ⁻¹
69 fb ⁻¹	430 fb ⁻¹
6.8 fb ⁻¹	43 fb ⁻¹
	2σ 137 fb ⁻¹ 69 fb ⁻¹ 6.8 fb ⁻¹





1. The data of Higgs direct search strongly constrain the model parameter α , the samples with $-0.9 \leq \alpha \leq 1.3$ are excluded. Constraints from $B \rightarrow X_s \gamma$, require tan $\beta \geq 2.6$. Muon g-2 need 20 m_A - 0.5 < $|m_{12}| < 20 m_A + 0.5$ TeV.

2. The Br(h $\rightarrow \gamma\gamma$) and reduced top-Higgs coupling are mainly affected by α and tan β , having max values 0.2 and 0.6. The cross section $\sigma(pp \rightarrow t\bar{t}h(\rightarrow \gamma\gamma))$ at HL-LHC peaks at 0.6 fb when sin α is close to 1 and tan β is near 2.7.

3. The backgrounds of $t\bar{t}\gamma\gamma$, $tj\gamma\gamma$, and $Wjj\gamma\gamma$ are difficult to eliminate. The samples with $\alpha \approx -1.2$ can hardly be covered all at 2σ level by the diphoton channel, even with increasing collider energy or integrated luminosity.

4. For the surviving samples, to get a significance of 2σ and 5σ at the HL-LHC require at least about 137 fb⁻¹ and 859 fb⁻¹. The same significance can be achieved at the HE-LHC with 69 fb⁻¹ and 430 fb⁻¹ and at the FCC-hh with 6.8 fb⁻¹ and 43 fb⁻¹.







	Type I	Type II	Lepton-specific	Flipped
ξ_h^u	$\cos \alpha / \sin \beta$	$\cos \alpha / \sin \beta$	$\cos \alpha / \sin \beta$	$\cos \alpha / \sin \beta$
ξ_h^d	$\cos \alpha / \sin \beta$	$-\sin \alpha / \cos \beta$	$\cos \alpha / \sin \beta$	$-\sin \alpha / \cos \beta$
ξ_h^ℓ	$\cos \alpha / \sin \beta$	$-\sin \alpha / \cos \beta$	$-\sin \alpha / \cos \beta$	$\cos \alpha / \sin \beta$
ξ^u_H	$\sin \alpha / \sin \beta$	$\sin \alpha / \sin \beta$	$\sin \alpha / \sin \beta$	$\sin \alpha / \sin \beta$
ξ^d_H	$\sin \alpha / \sin \beta$	$\cos \alpha / \cos \beta$	$\sin \alpha / \sin \beta$	$\cos \alpha / \cos \beta$
ξ_H^ℓ	$\sin \alpha / \sin \beta$	$\cos \alpha / \cos \beta$	$\cos \alpha / \cos \beta$	$\sin \alpha / \sin \beta$
ξ^u_A	\coteta	\coteta	$\cot eta$	$\cot eta$
ξ^d_A	$-\cot\beta$	an eta	$-\cot\beta$	an eta
ξ^{ℓ}_{A}	$-\cot\beta$	an eta	$\tan\beta$	$-\cot\beta$

arXiv: 1106.0034

from vacuum stability, tree-level unitarity, Z decays into hadrons, and the Δr parameter. The ratio was analyzed for the neutral CP-even Higgs h with 110 GeV $\leq M_h \leq 150$ GeV (which is the region where an SM-like light Higgs boson h can be discovered at the LHC in $h \to \gamma\gamma$), and with $M_{H^{\pm}} = 100$ GeV for the charged Higgs boson. The maximal possible enhancement differs for the potentials considered due to different interference scenarios. For the IDM the maximal enhancement was found to be around +70%, for V_A around -20%. The results for V and V_B were vather similar, being around 70% and 60%, respectively. This is larger than what was found for

1001.1759, 1706.07414



 $v^{2}\lambda_{1} = \frac{m_{H}^{2}c_{\alpha}^{2} + m_{h}^{2}s_{\alpha}^{2} - m_{12}^{2}\tan\beta}{c_{\beta}^{2}}, \qquad v^{2}\lambda_{2} = \frac{m_{H}^{2}s_{\alpha}^{2} + m_{h}^{2}c_{\alpha}^{2} - m_{12}^{2}/\tan\beta}{s_{\beta}^{2}},$ $v^{2}\lambda_{3} = \frac{(m_{H}^{2} - m_{h}^{2})s_{\alpha}c_{\alpha} + 2m_{H^{\pm}}^{2}s_{\beta}c_{\beta} - m_{12}^{2}}{s_{\beta}c_{\beta}}, \qquad v^{2}\lambda_{4} = \frac{(m_{A}^{2} - 2m_{H^{\pm}}^{2})s_{\beta}c_{\beta} + m_{12}^{2}}{s_{\beta}c_{\beta}},$ $v^{2}\lambda_{5} = \frac{-m_{A}^{2}s_{\beta}c_{\beta} + m_{12}^{2}}{s_{\beta}c_{\beta}}. \qquad (A.$ (A.1)

$$\begin{split} m_{11}^2 &= m_{12}^2 \tan\beta - \frac{1}{2} v^2 \left(\lambda_1 c_\beta^2 + (\lambda_3 + \lambda_4 + \lambda_5) s_\beta^2 \right) \,, \\ m_{22}^2 &= m_{12}^2 \cot\beta - \frac{1}{2} v^2 \left(\lambda_2 s_\beta^2 + (\lambda_3 + \lambda_4 + \lambda_5) c_\beta^2 \right) \,. \end{split}$$

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