

Dark matter searches at the CEPC

Peng-Fei Yin

Key laboratory of particle astrophysics, IHEP, CAS

Zheng Zhou, 2024.08.31

In collaboration with X. J. Bi, L. Q. Gao, J. W. Wang, Q. F. Xiang, Q. S. Yan, and Z. H. Yu

Outline

- ⊕ *Introduction*
- ⊕ *Direct production of DM at the CEPC*
- ⊕ *Indirect searches through precision measurements at the CEPC*
- ⊕ *Summary*

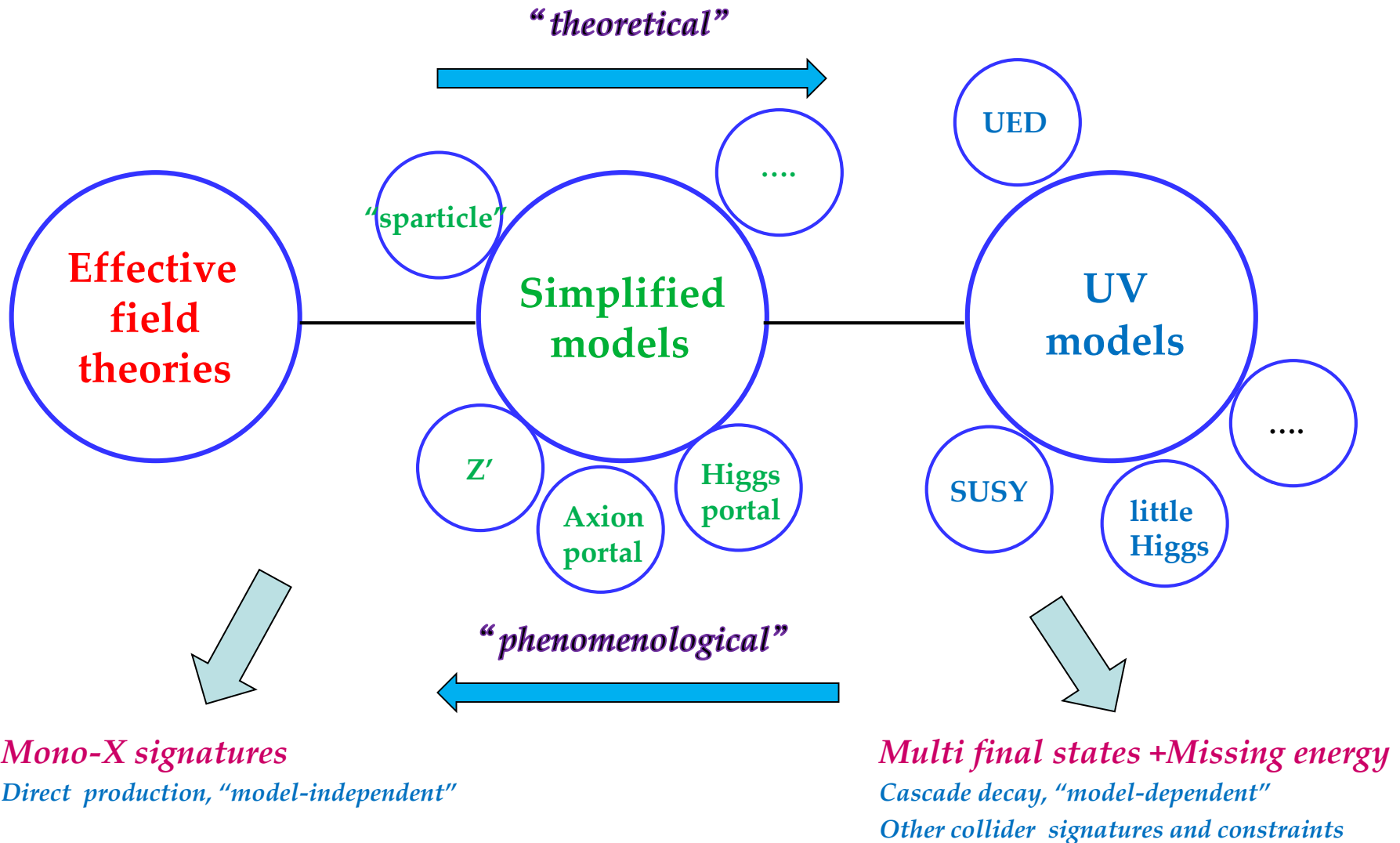
DM searches at the e^+e^- colliders

- ✦ *Search for light DM **directly produced**, e.g. $m < 120 \text{ GeV}$, at the CEPC*
Search for interactions between the DM and electrons/EW gauge bosons/Higgs
limited by the low CM energies of e^+e^- colliders
- ✦ *Probe DM and other relevant particles in new physics models (t-channel annihilation mediator, charged particles in multiplets...)*
Indirect searches through loop effects, e.g. in the Higgs production and decay
- ✦ *Precise measurement rather than discovery*
Full missing energy can be reconstructed
*possible to **precisely measure** the mass, spin, and other quantum numbers of DM*

DM searches at the e^+e^- colliders

- ✦ *Advantage: No large QCD background; precise beam energy; larger luminosity polarized beams....*
Precisely measurements of DM properties are possible
Investigate the production mechanism of DM by using these results
- ✦ *Disadvantage: Low CM energy....*
difficult to directly detect heavy particles in the BSM
- ✦ *Complementary to searches at hadron colliders*
- ✦ *Complementary to direct and indirect DM detection experiments*
- ✦ *Even if a new neutral, stable, and weakly interacting particle is discovered at colliders, we should ascertain its potential as a constituent of dark matter within the Universe*

Direct signatures: theoretical approach

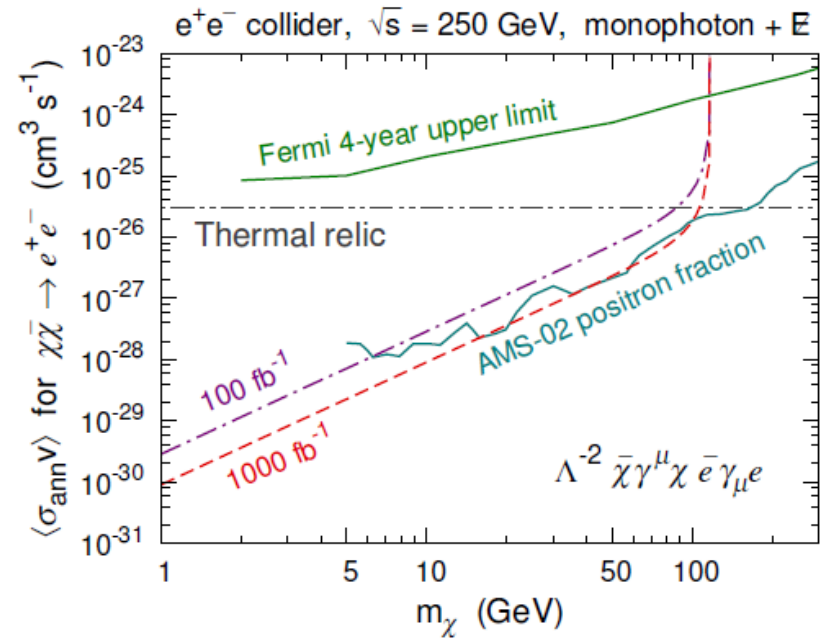
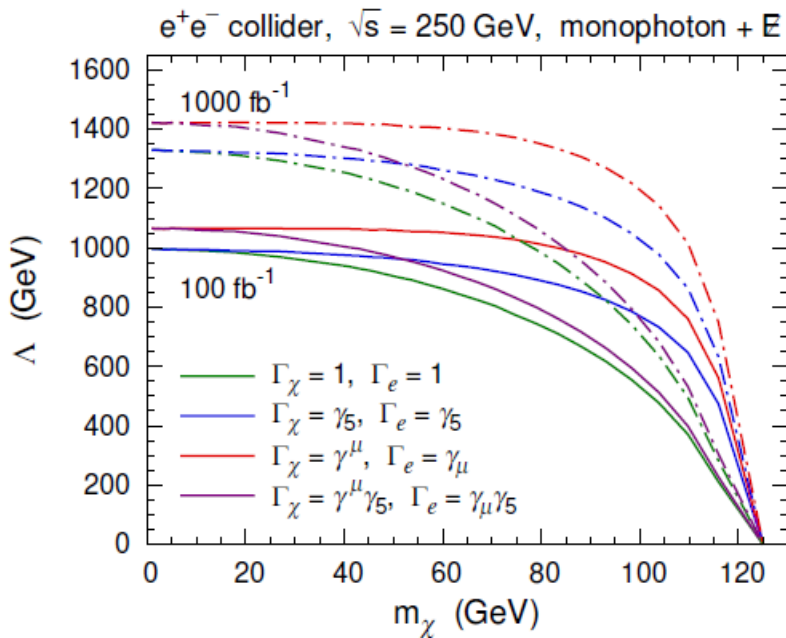


Mono-photon searches for EFTs

- ⊕ Consider the simplest EFTs

$$\mathcal{O}_e = \frac{1}{\Lambda^2} \bar{\chi} \Gamma_\chi \chi \bar{e} \Gamma_e e \quad \Gamma_\chi, \Gamma_e \in \{1, \gamma_5, \gamma^\mu, \gamma^\mu \gamma_5, \sigma^{\mu\nu}\}$$

- ⊕ Photon is emitted from the initial state radiation



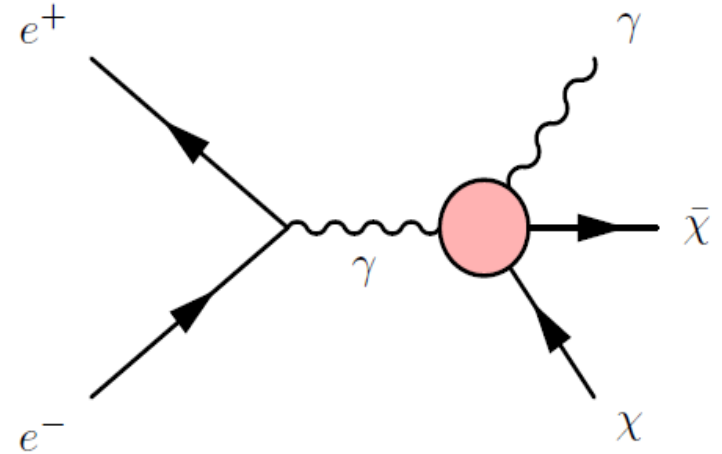
Gamma-ray line and mono-photon

- ✦ *Gamma-ray line signature is a critical evidence of DM annihilation/decay*
- ✦ *Search for corresponding mono-photon signals at future e^+e^- colliders*
- ✦ *Effective operator*

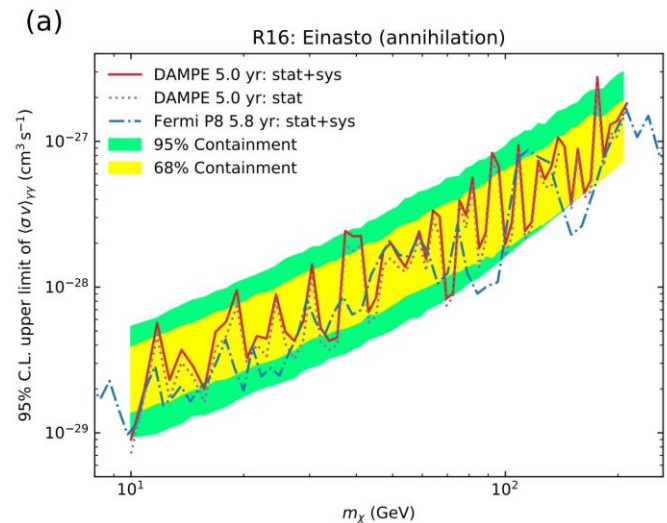
$$\mathcal{O}_F = \frac{1}{\Lambda^3} \bar{\chi} i\gamma_5 \chi F_{\mu\nu} \tilde{F}^{\mu\nu}$$

Cut scale for a ~ 100 GeV DM and a detectable annihilation cross section is $\sim \text{TeV}$

$$\langle \sigma_{\text{ann}} v \rangle_{\chi\bar{\chi} \rightarrow 2\gamma} \simeq \frac{4m_\chi^4}{\pi\Lambda^6} = 10^{-27} \text{ cm}^3 \text{ s}^{-1} \left(\frac{m_\chi}{130 \text{ GeV}} \right)^4 \left(\frac{1272 \text{ GeV}}{\Lambda} \right)^6$$



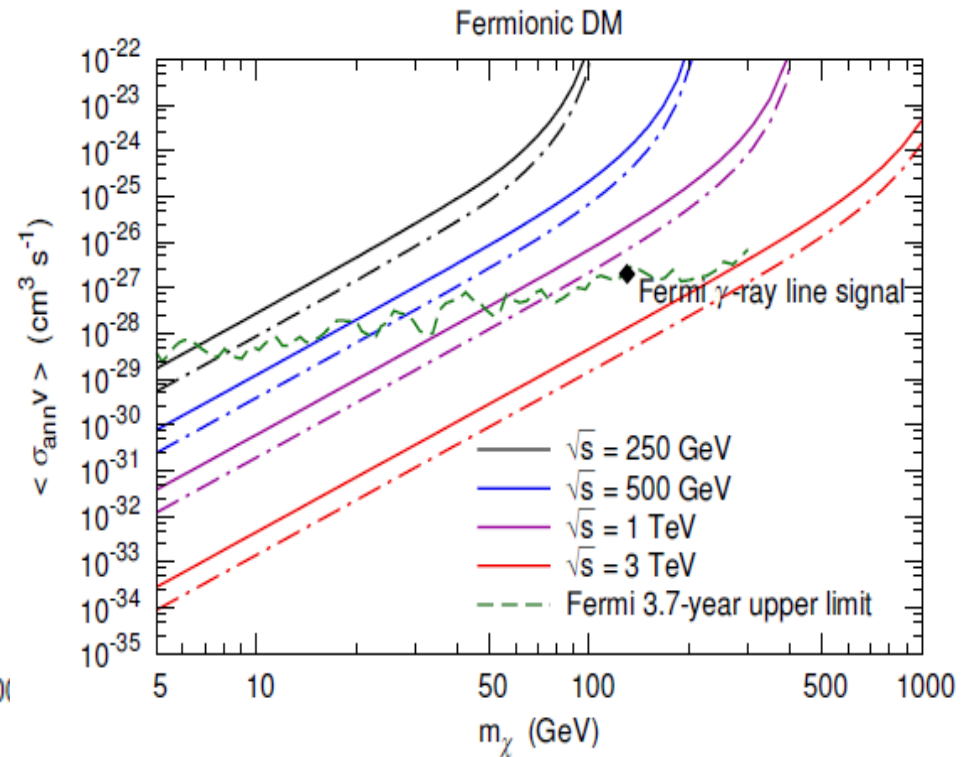
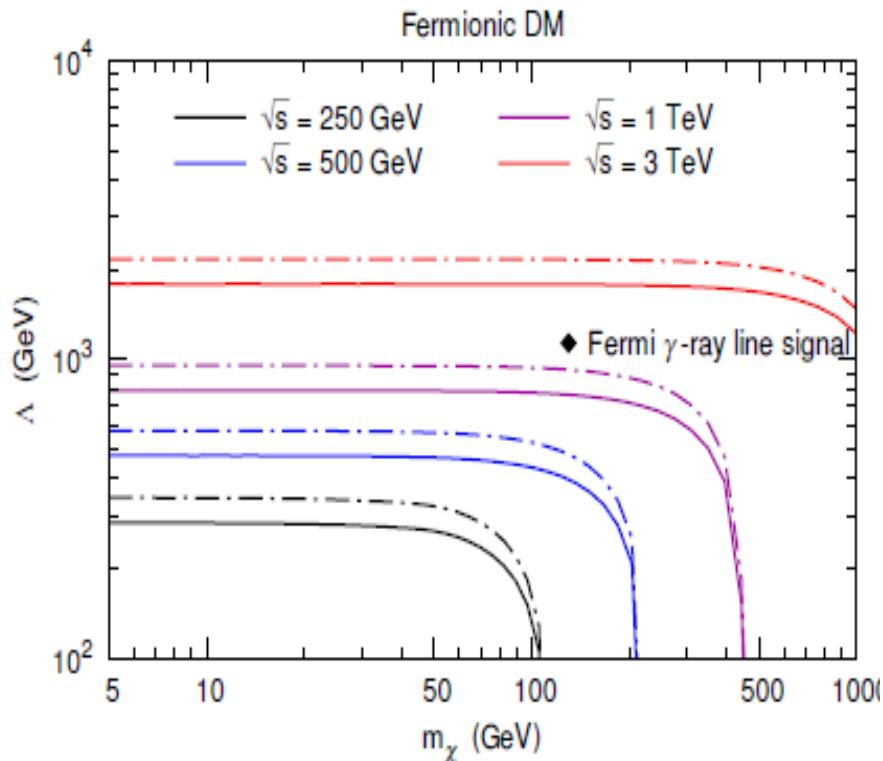
Yu, Yan, Yin, 1307.5740



DAMPE, 2112.08860

Gamma-ray line and mono-photon

- ⊕ Consider possible e^+e^- colliders with several CM energies
- ⊕ 3σ reaches for mass scale and annihilation cross section
- ⊕ Require large luminosities ($>100 \text{ fb}^{-1}$)



Mono-Z signatures for EFTs

⊕ *DM can interact with both the photon and Z boson*

⊕ *Consider effective operators*

$$\mathcal{O}_{F1} = \frac{1}{\Lambda_1^3} \bar{\chi} \chi B_{\mu\nu} B^{\mu\nu} + \frac{1}{\Lambda_2^3} \bar{\chi} \chi W_{\mu\nu}^a W^{a\mu\nu}$$

$$\supset \bar{\chi} \chi (G_{ZZ} Z_{\mu\nu} Z^{\mu\nu} + G_{AZ} A_{\mu\nu} Z^{\mu\nu})$$

$$\mathcal{O}_{F2} = \frac{1}{\Lambda_1^3} \bar{\chi} i\gamma_5 \chi B_{\mu\nu} \tilde{B}^{\mu\nu} + \frac{1}{\Lambda_2^3} \bar{\chi} i\gamma_5 \chi W_{\mu\nu}^a \tilde{W}^{a\mu\nu}$$

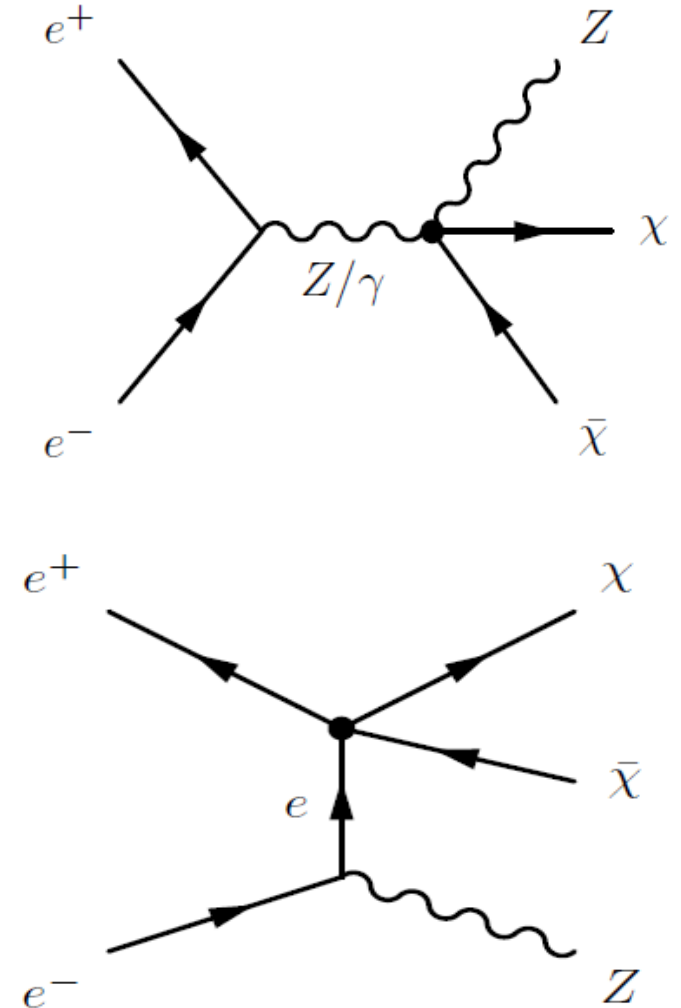
$$\supset \bar{\chi} i\gamma_5 \chi (G_{ZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu} + G_{AZ} A_{\mu\nu} \tilde{Z}^{\mu\nu})$$

$$\mathcal{O}_{FH} = \frac{1}{\Lambda^3} \bar{\chi} \chi (D_\mu H)^\dagger D_\mu H \rightarrow \frac{m_Z^2}{2\Lambda^3} \bar{\chi} \chi Z_\mu Z^\mu$$

⊕ *Z boson can also come from initial state radiation*

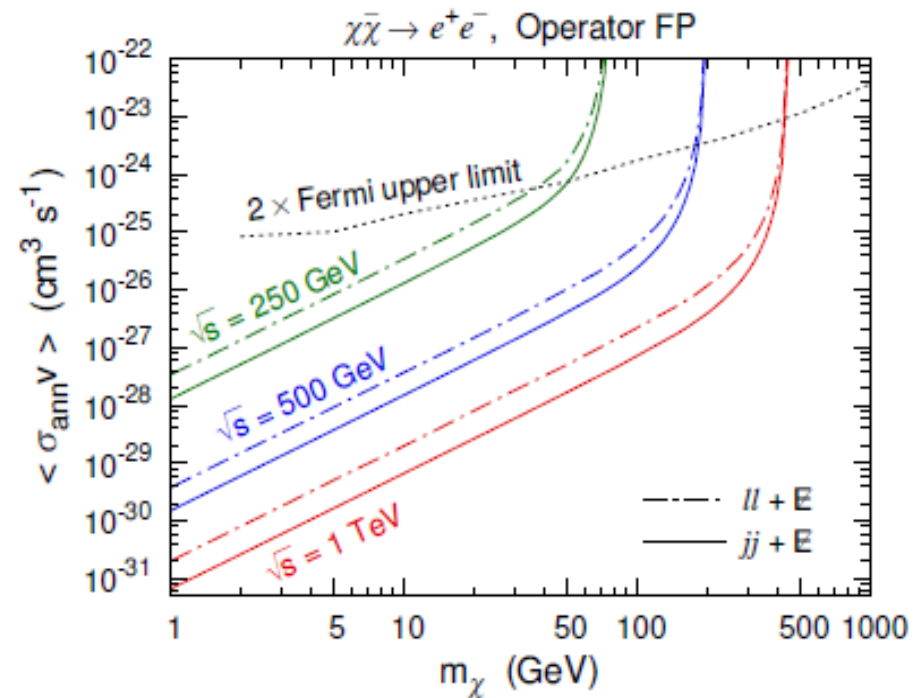
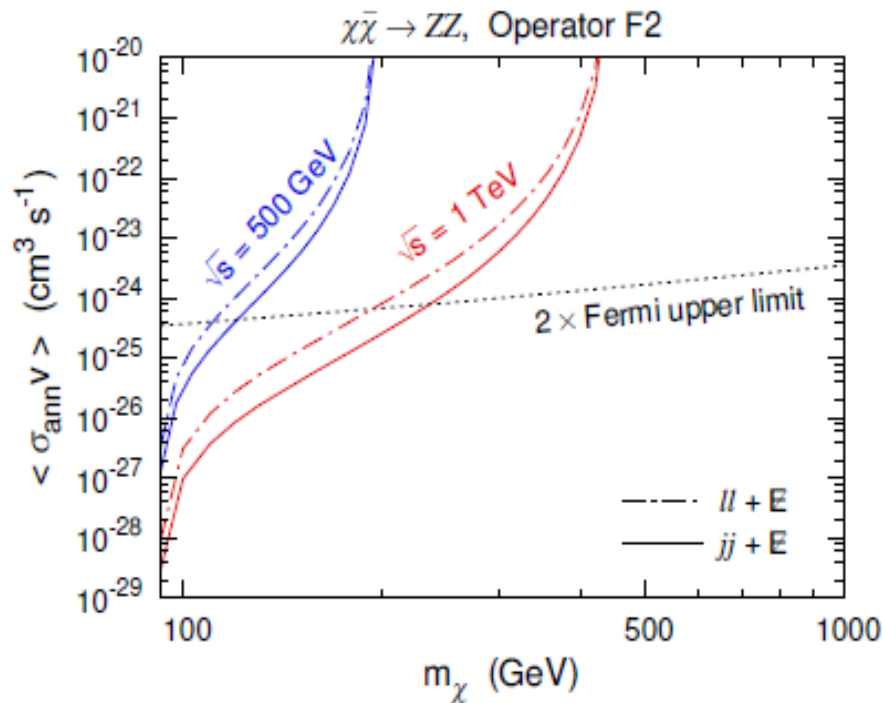
$$\mathcal{O}_{FP} = \frac{1}{\Lambda^2} \bar{\chi} \gamma_5 \chi \bar{e} \gamma_5 e,$$

$$\mathcal{O}_{FA} = \frac{1}{\Lambda^2} \bar{\chi} \gamma^\mu \gamma_5 \chi \bar{e} \gamma_\mu \gamma_5 e$$

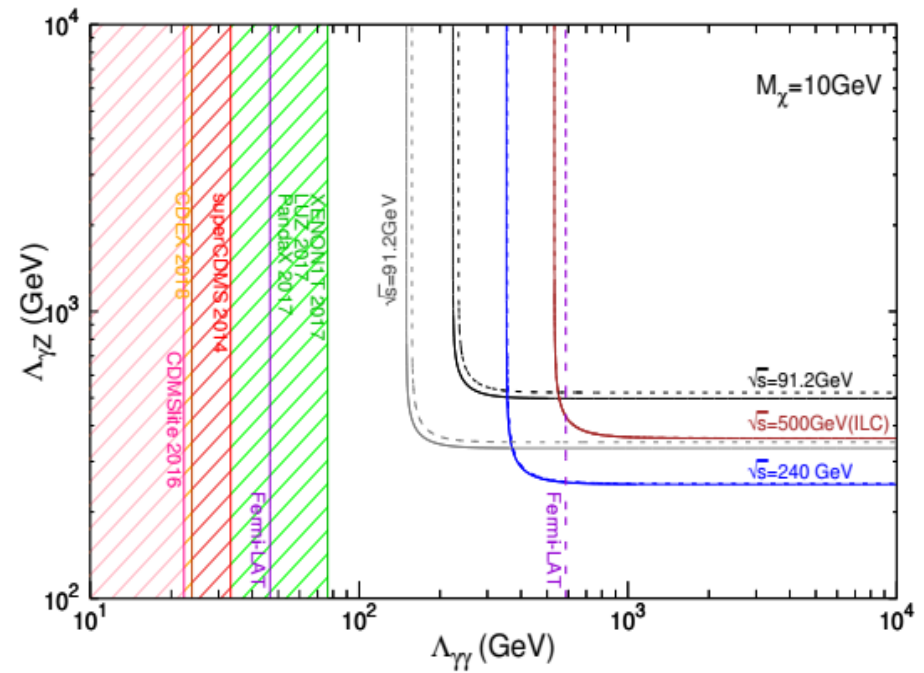
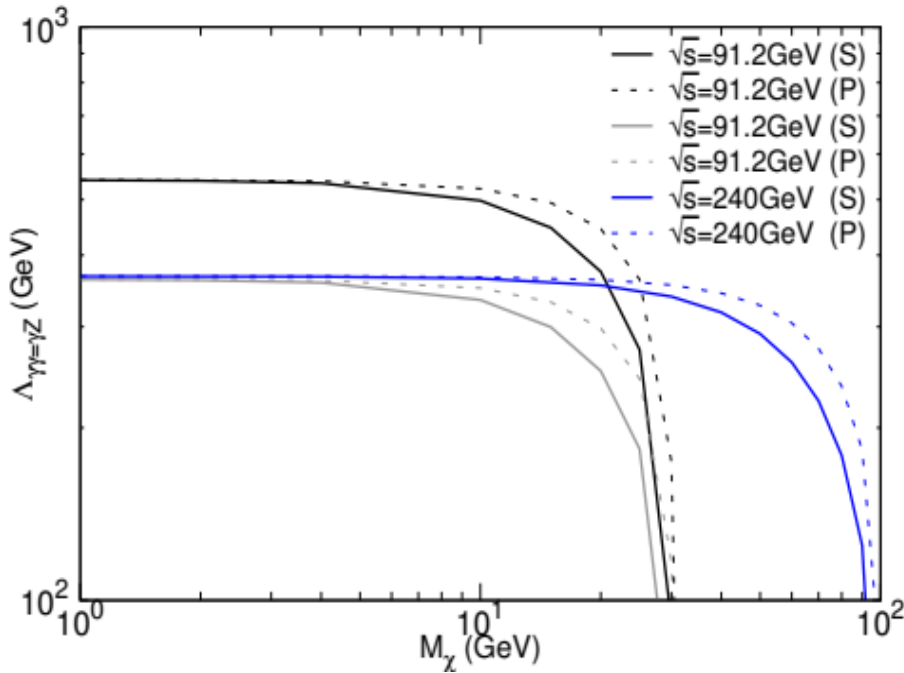


Mono-Z signatures for EFTs

- ✦ 3σ reaches for interactions between the DM and gauge bosons/electrons
- ✦ Assume: 1000 fb^{-1} of data; $L=L_1=L_2$
- ✦ Compare with the limits from Fermi-LAT dwarf galaxy observations



Signatures at the Z-pole



Jin, Gao, 1712.02140

$$\mathcal{L}_S \supset \frac{1}{\Lambda_{\gamma\gamma}^3} \bar{\chi} \chi A^{\mu\nu} A_{\mu\nu} + \frac{1}{\Lambda_{\gamma Z}^3} \bar{\chi} \chi A^{\mu\nu} Z_{\mu\nu},$$

$$\mathcal{L}_P \supset \frac{1}{\Lambda_{\gamma\gamma}^3} \bar{\chi} i \gamma_5 \chi A^{\mu\nu} \tilde{A}_{\mu\nu} + \frac{1}{\Lambda_{\gamma Z}^3} \bar{\chi} i \gamma_5 \chi A^{\mu\nu} \tilde{Z}_{\mu\nu}$$

- ✦ *The high luminosity Z-pole run offers a unique opportunity to search for DM particles*

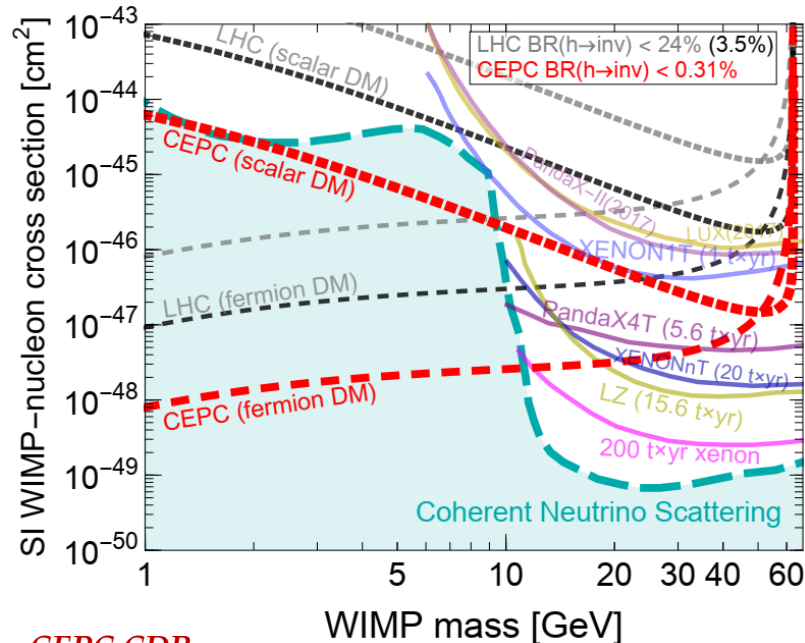
Simplified models: DM couples to mediators

- Assume that DM couples to the SM particles through some SM mediators

$$\mathcal{L} = -hJ_h, \quad J_h = \frac{1}{\sqrt{2}} \left[\sum_f y_f \bar{f}f + \bar{\psi}_{\text{DM}}(y_{\text{DM}} + iy_{\text{DM}}^P \gamma_5)\psi_{\text{DM}} + \frac{\lambda_{\text{DM}} v}{2} s_{\text{DM}}^2 \right]$$

$$\mathcal{L} = -Z_\mu J_\mu^Z, \quad J_\mu^Z = \frac{g_2}{\cos \theta_W} \left[\sum_f [\bar{f} \gamma_\mu (g_V^f + \gamma_5 g_A^f) f] + \sum_s g_s [s^* (i\partial_\mu s) - (i\partial_\mu s^*) s] \right]$$

- Searches for invisible Higgs/Z decays are useful to probe DM

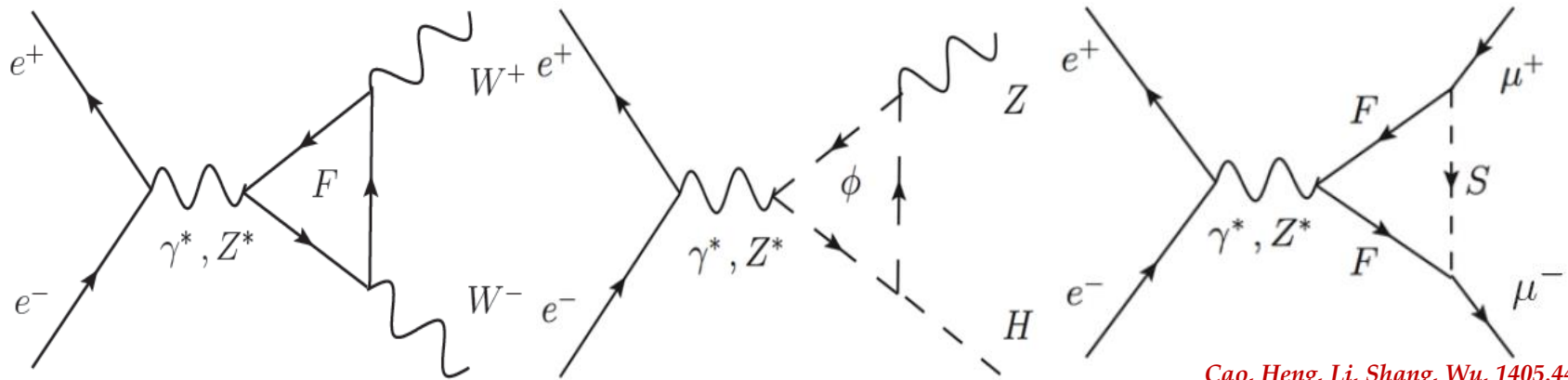


CEPC CDR

- For discussions on more portals and UV models, referred to talks in other sessions, such as dark sector and SUSY.

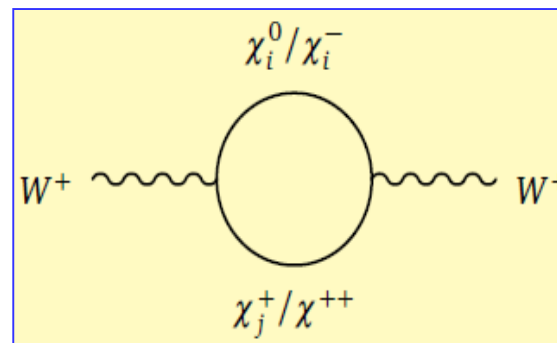
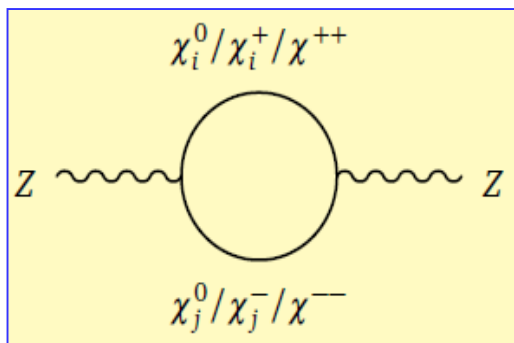
Probing DM at CEPC via loop effects

- ✦ New particles give **corrections** to SM processes, which can be measured at a level of $\sim O(0.1\%)$ at lepton colliders



Cao, Heng, Li, Shang, Wu, 1405.4489
 Cao, Li, Yan, Zhang, Zhang, 1604.07536
 Liu, Wu, 1705.02534
 Xiang, Bi, Yin, Yu, 1707.03094

- ✦ Also give **oblique corrections** to gauge boson propagators
 Search for new particles via the precision measurements and global fitting



Fedderke, Lin, Wang, 1506.05465
 Cai, Yu, Zhang, 1611.02186, 1705.07921

DM models with additional EW multiplets

- ✦ Consider some simple Fermionic DM models containing new *electro-weak multiplets*: few new particles and no new mediator

- ✦ Add one high-dimensional representation: *minimal DM model*

Cirelli et al, hep-ph/0512090

- ✦ The model can also contain *a vector-like fermion and a Z_2 symmetry stabilizing DM*. But no coupling to the SM Higgs : no mass contribution from EWSB, degenerate mass spectrum...

DM models with additional EW multiplets

- ✦ *Add two types of vector EW Fermionic multiplets may be an economical option with a rich phenomenology. We consider the models:*

SDFDM: one singlet + two doublet Weyl spinors

DTFDM: two doublet + one triplet Weyl spinors

TQFDM: two quadruplet + one triplet Weyl spinors

Xiang, Bi, Yin, Yu, 1707.03094

Wang, et. al., 1711.05622

Models	Gauge eigenstates	Mass eigenstates
Singlet-Doublet	$S, \begin{pmatrix} D_1^0 \\ D_1^- \end{pmatrix}, \begin{pmatrix} D_2^+ \\ D_2^0 \end{pmatrix}$	$\chi_1^0, \chi_2^0, \chi_3^0$ χ^\pm
Doublet-Triplet	$\begin{pmatrix} D_1^0 \\ D_1^- \end{pmatrix}, \begin{pmatrix} D_2^+ \\ D_2^0 \end{pmatrix}, \begin{pmatrix} T^+ \\ T^0 \\ -T^- \end{pmatrix}$	$\chi_1^0, \chi_2^0, \chi_3^0$ χ_1^\pm, χ_2^\pm

- ✦ *Analogous to some well-studied DM models, such as **SUSY DM***
SDFDM-> Bino-Higgsino in MSSM; Singlino-Higgsino in NMSSM
DTFDM-> Higgsino-Wino in MSSM

$$\Delta L = M_1 \tilde{B} \tilde{B} + M_2 \tilde{W} \tilde{W} + \mu \tilde{H}_u \tilde{H}_d$$

$$+ \sqrt{2} \kappa_1 h^\dagger \tilde{W} \tilde{H}_u + \sqrt{2} \kappa_2 h \tilde{W} \tilde{H}_d + \frac{\kappa'_1}{\sqrt{2}} h^\dagger \tilde{B} \tilde{H}_u + \frac{\kappa'_2}{\sqrt{2}} h \tilde{B} \tilde{H}_d$$

Arkani et. al, 1511.06495

Singlet-Doublet Fermionic Model

- ⊕ Introduce one weyl singlet and two doublets
Gauge invariant Lagrangian

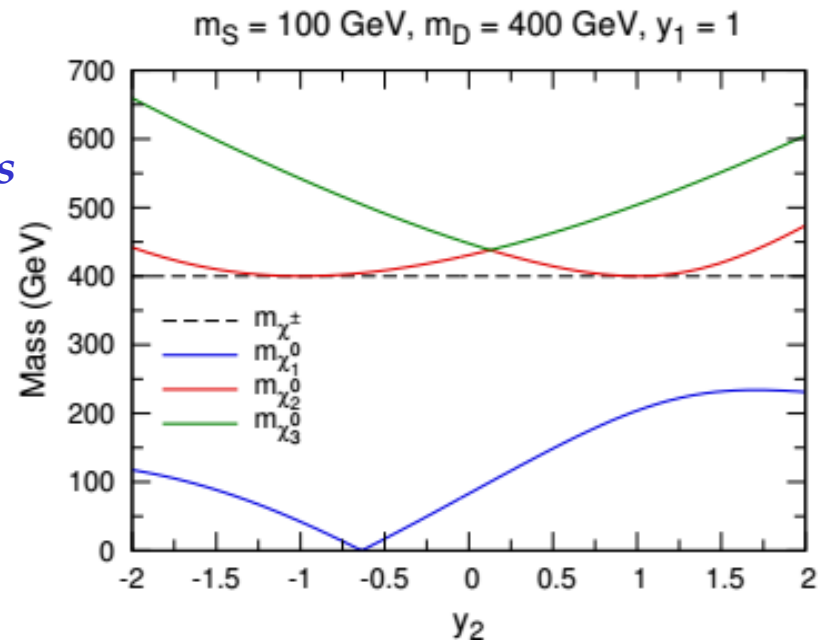
$$\mathcal{L}_S = iS^\dagger \bar{\sigma}^\mu \partial_\mu S - \frac{1}{2}(m_S S S + \text{h.c.}),$$

$$\mathcal{L}_D = iD_1^\dagger \bar{\sigma}^\mu D_\mu D_1 + iD_2^\dagger \bar{\sigma}^\mu D_\mu D_2 - (m_D \epsilon_{ij} D_1^i D_2^j + \text{h.c.})$$

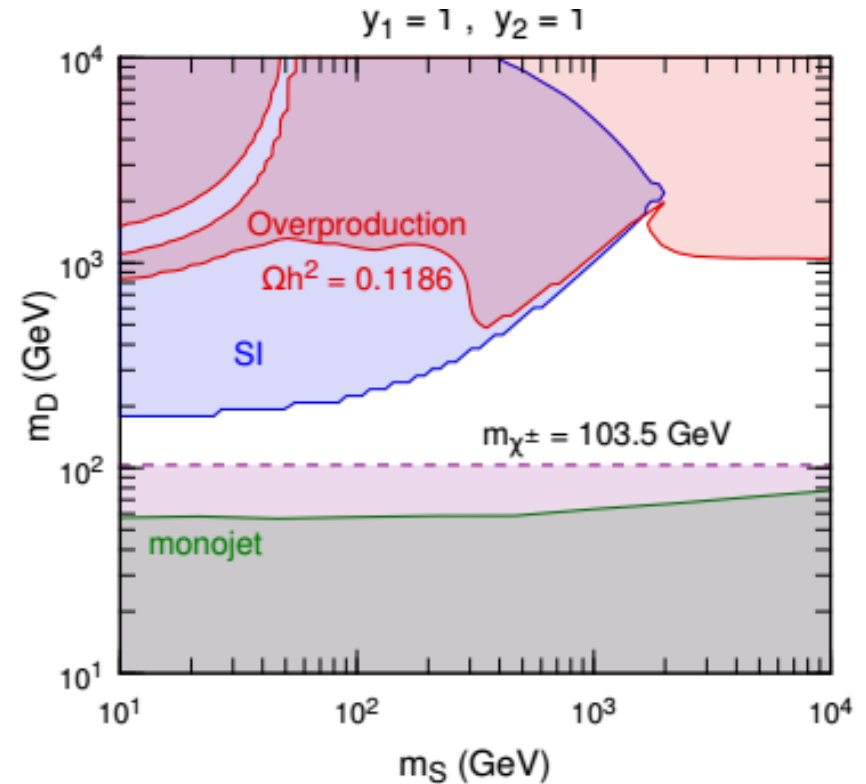
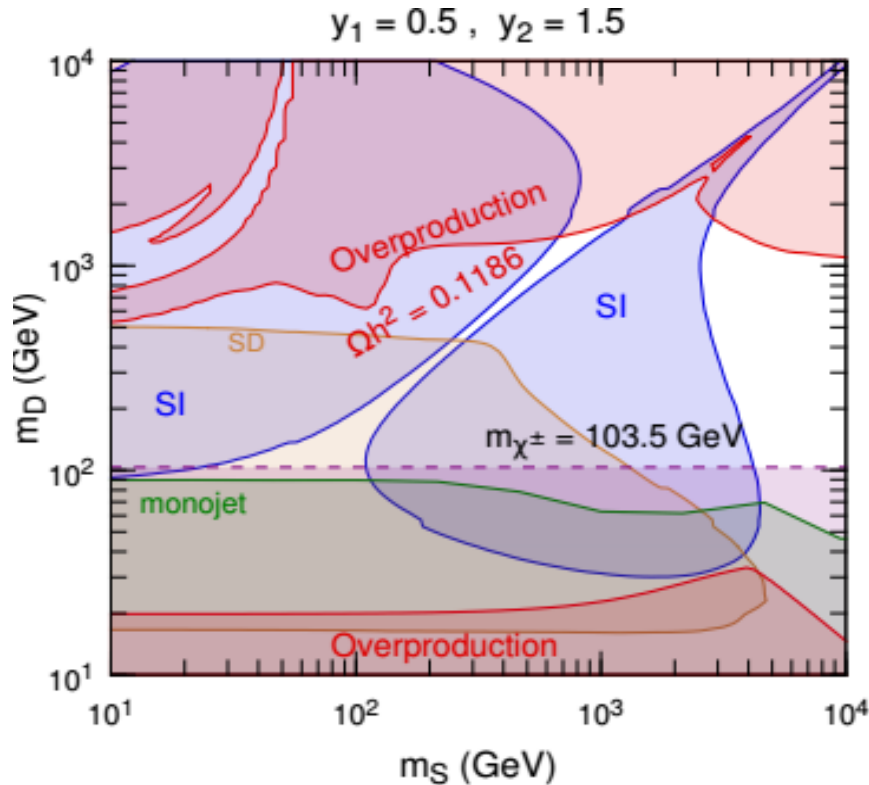
Yukawa coupling

$$\mathcal{L}_{\text{yuk}} = y_1 S D_1^i H_i - y_2 S D_2^i H_i^\dagger + \text{h.c.}$$

- ⊕ Four new parameters: two mass parameters m_S , m_D and two Yukawa couplings y_1 , y_2
- ⊕ DM is the lightest mass eigenstate in the neutral sector after EWSB

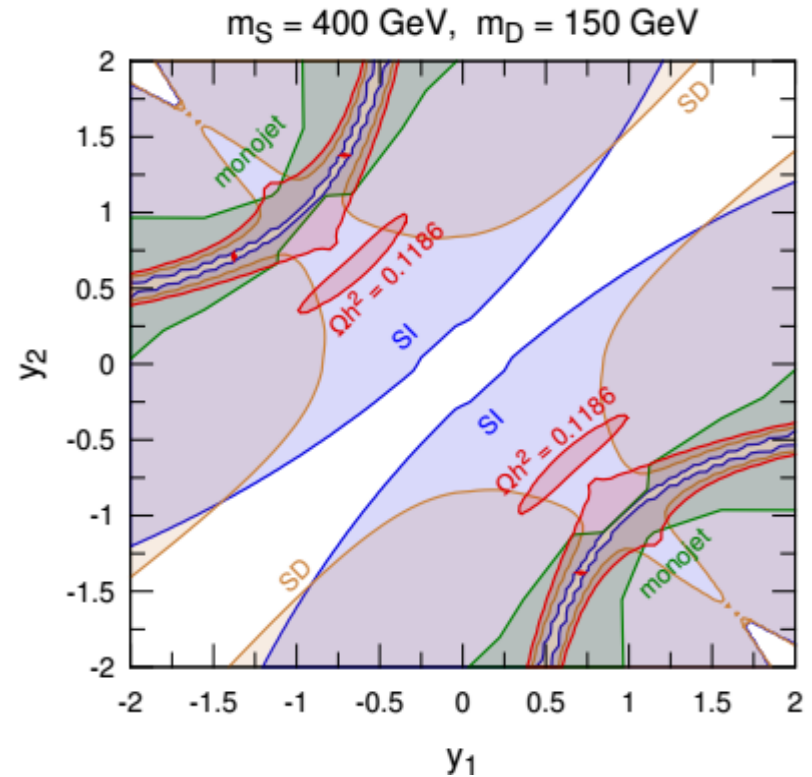
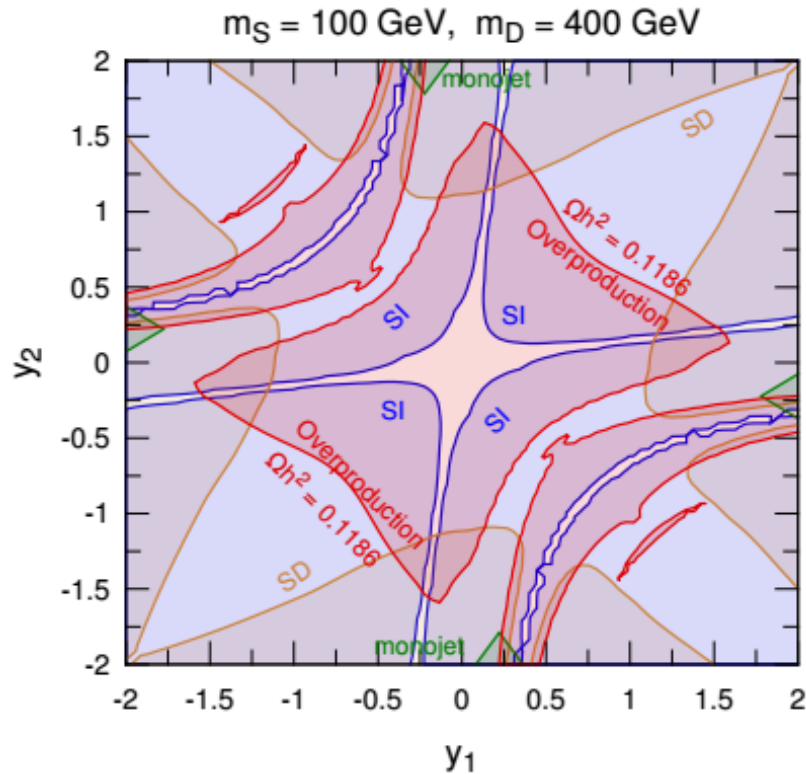


Constraints from DM searches



- ⊕ *DM relic density: take into account the coannihilation effects*
- ⊕ *Direct detection: Spin-independent (SI) and Spin-dependent (SD) constraints (PANDAX)*
- ⊕ *LHC: mono-jet limits (ATLAS) on the production of dark sector particles*
- ⊕ *LEP: require the mass of new charge particle is smaller than 103.5 GeV*

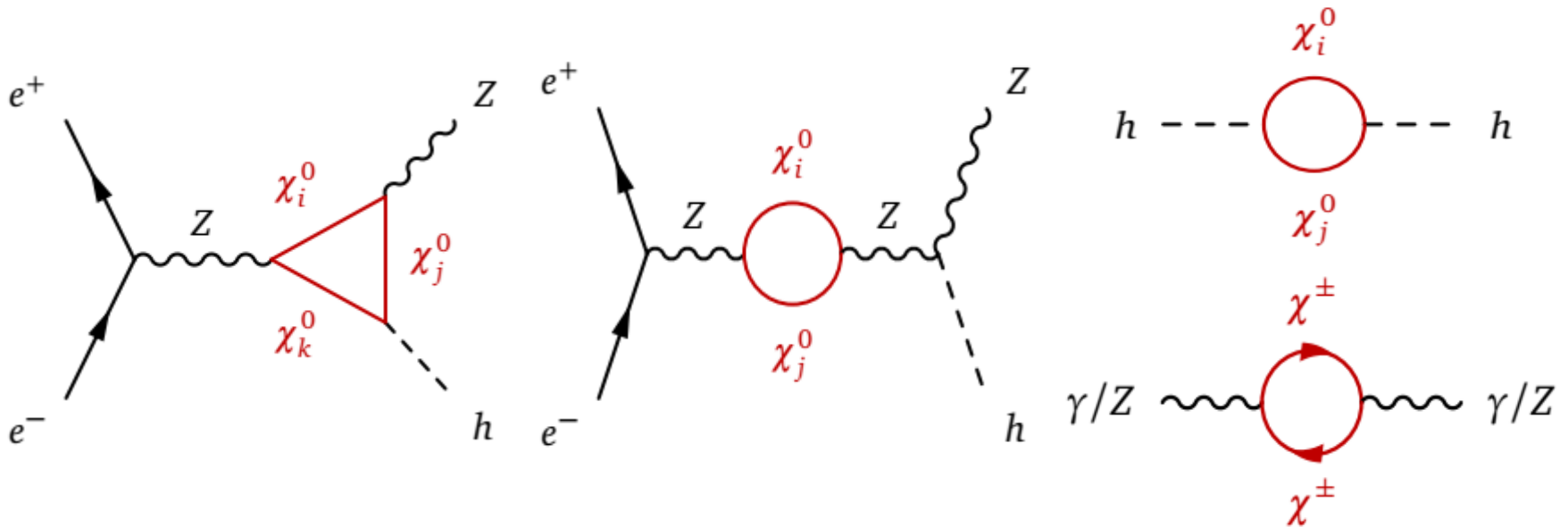
Constraints from DM searches



- ⊕ *Parameter space is stringently constrained by direct detection, except for some regions where the DM couplings to Z and Higgs are suppressed*

Probing DM via measurements of Zh production

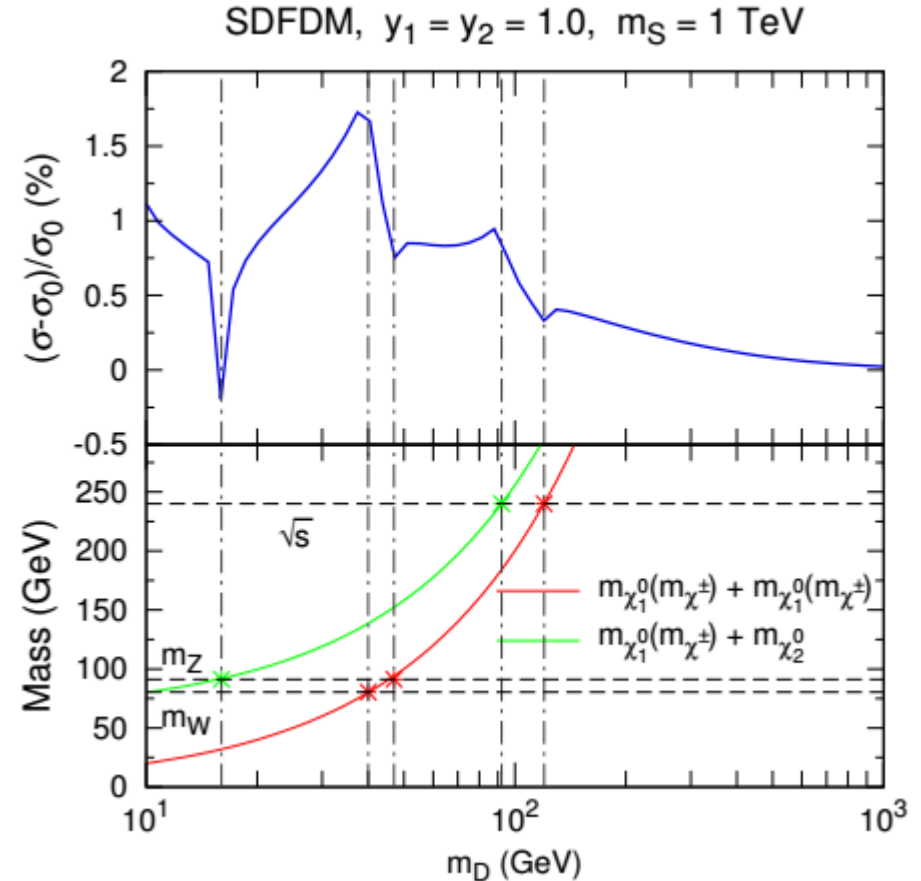
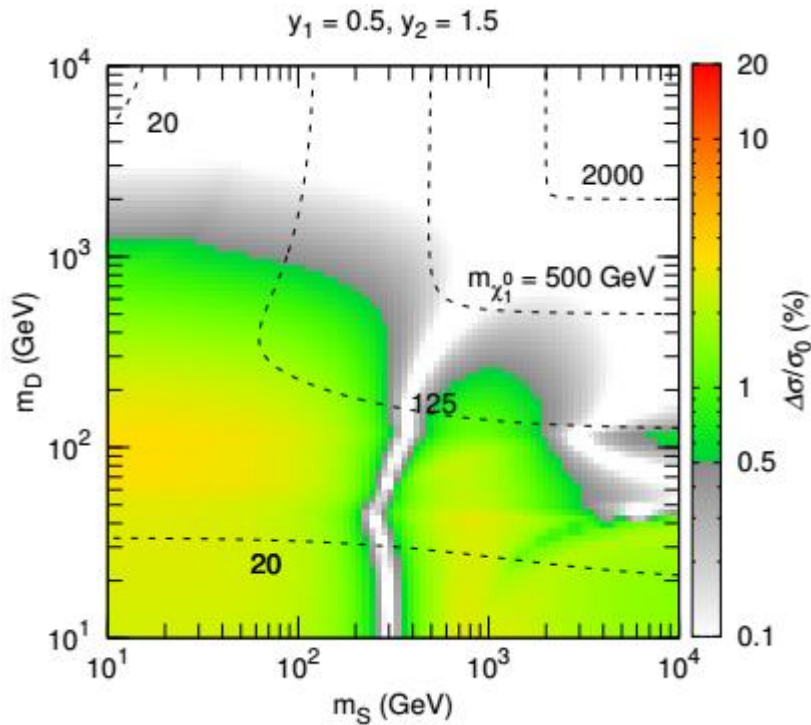
- ✦ We consider the corrections of new particles to the associated **Z-Higgs production**
This process is affected by both the gauge and Yukawa interactions



calculate corrections to the Zh production cross section ($\Delta\sigma/\sigma = (\sigma_{NP} - \sigma_{SM})/\sigma_{SM}$)

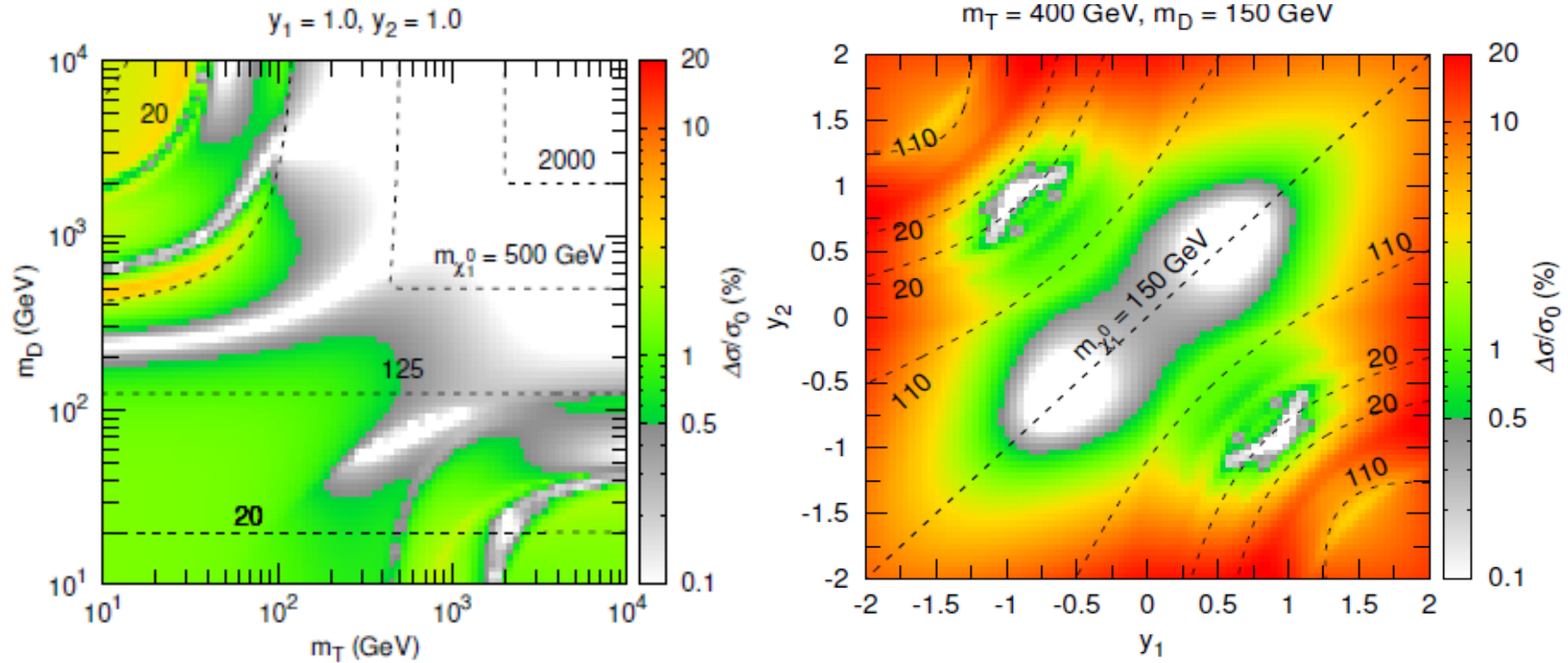
- ✦ It is possible to measure $\Delta\sigma/\sigma|_{Zh}$ at a level of $\sim 0.5\%$ at CEPC (with $5ab^{-1}$ of data)

Probing DM via measurements of Zh production



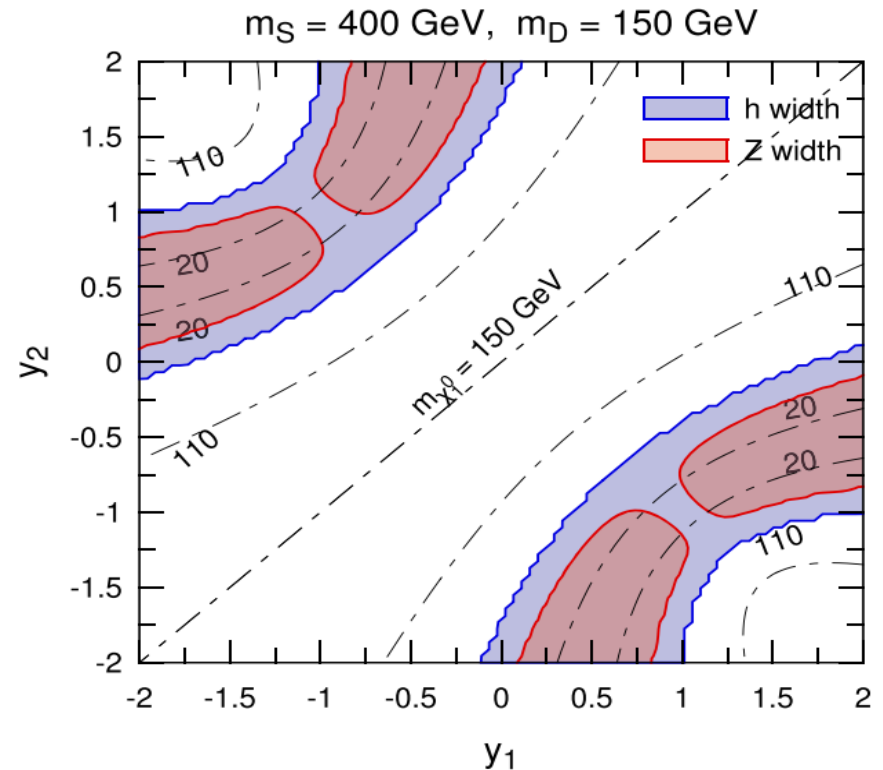
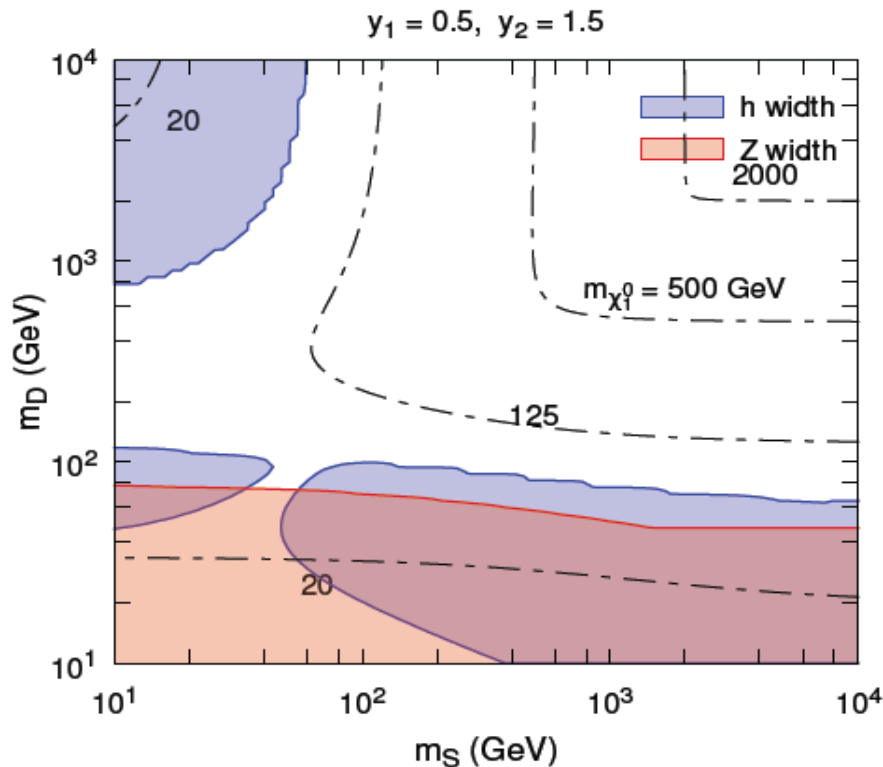
- ⊕ For $y_1=0.5$ and $y_2=1.5$, CEPC can probe to $m_{\chi} \sim 200$ GeV
- ⊕ When the dark sector particles in loops are close to their mass shells, their contributions vary dramatically

Probing DM via measurements of Zh production



Searches of Higgs and Z decays

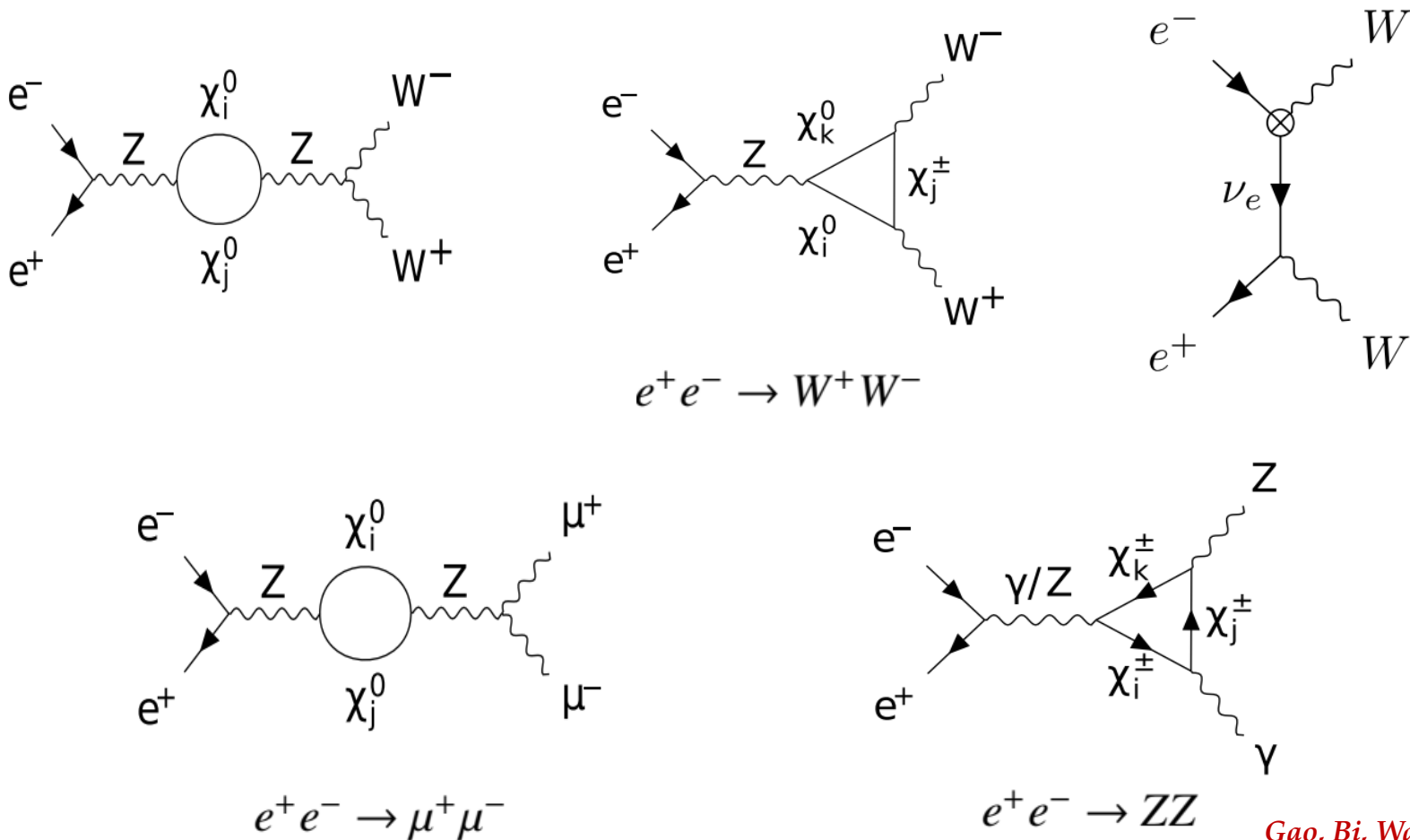
- ✦ If new charged particles couple to the Higgs, they could modified the **Higgs decay width to di-photons** via triangle loops. $\Delta\Gamma/T$ can be test to be a level of $\sim 9\%$
- ✦ If the kinematics is allowed, the Higgs and Z could decay into DM particles. Such **invisible decays** are constrained by the relevant searches at colliders



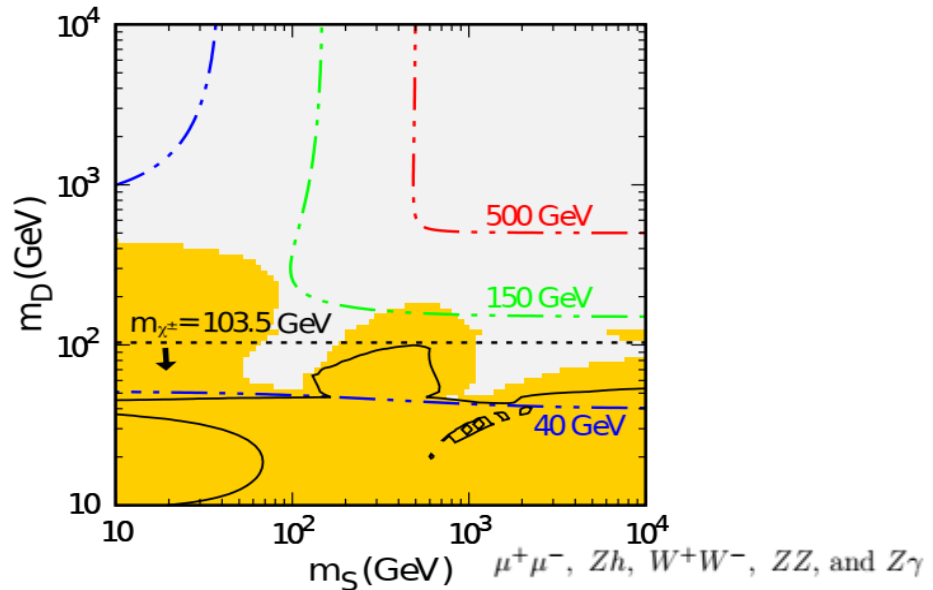
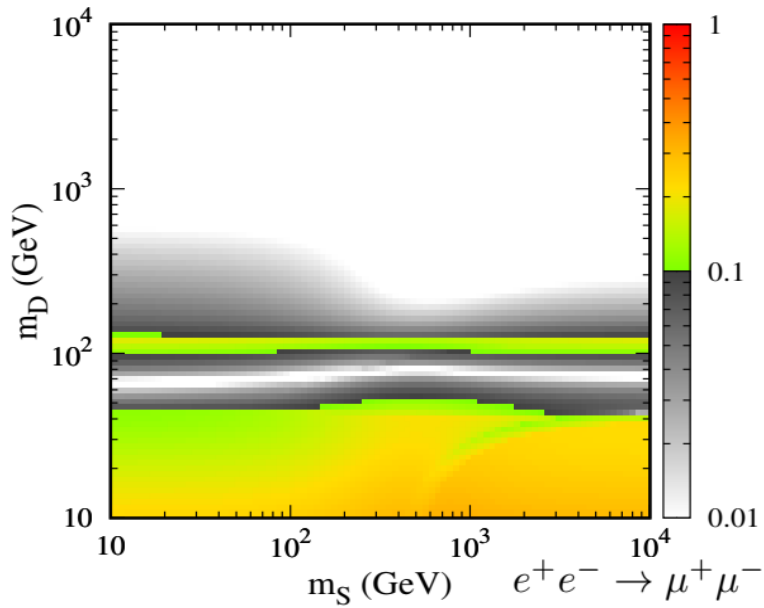
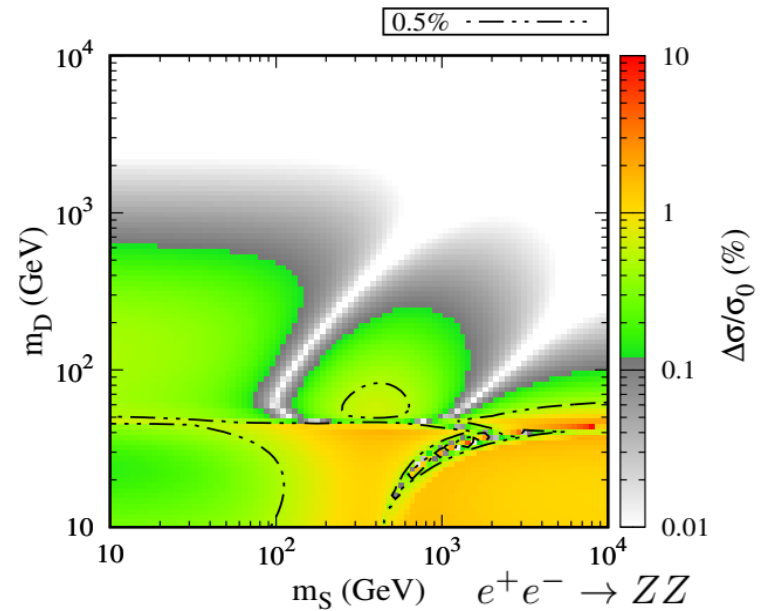
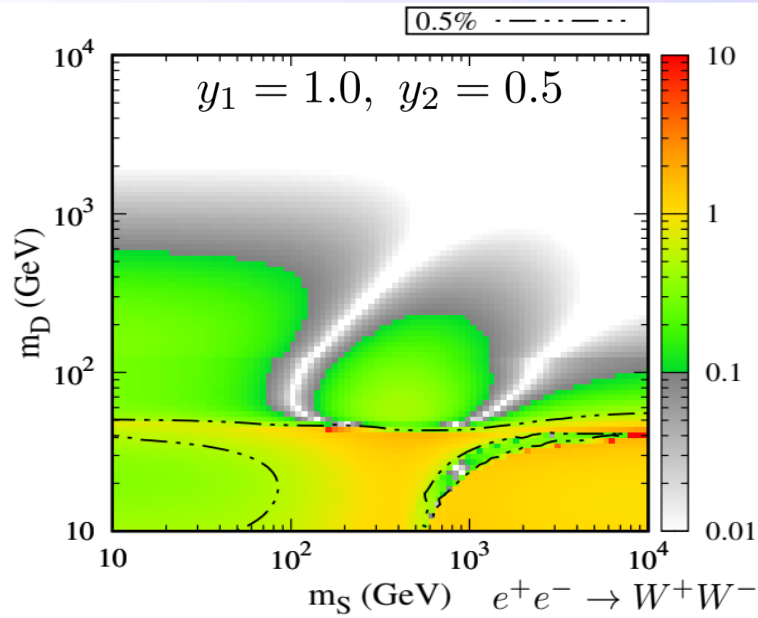
Take $G(Z \rightarrow inv) < 2$ MeV (LEP) and $G(h \rightarrow inv)/G_h < 2.8\%$ (for CEPC $5 ab^{-1}$)

Searches via precision measurements

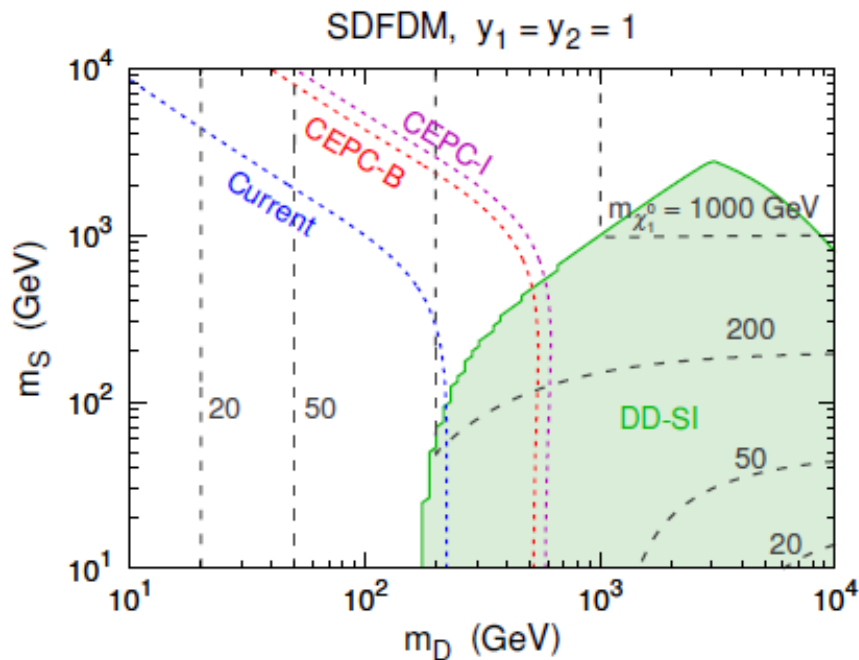
- It is possible to explore these particles through loop effects using the precision measurements for other EW processes at the CEPC



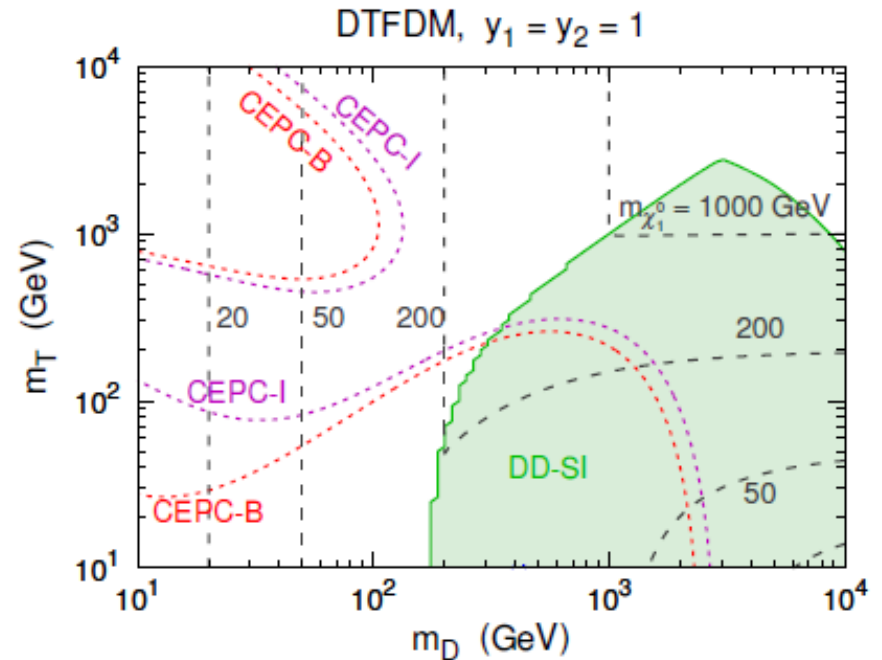
Searches via precision measurements



Oblique parameters: fermionic DM



Singlet and Doublet Fermions



Doublet and Triplet Fermions

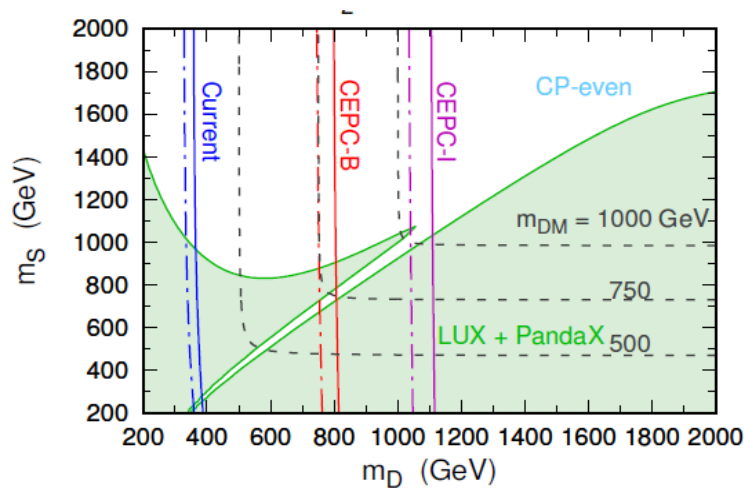
Cai, Yu, Zhang, 1611.02186

- ⊕ After EWSB, dark sector fermions contribute to the electroweak oblique parameters S , T , and U

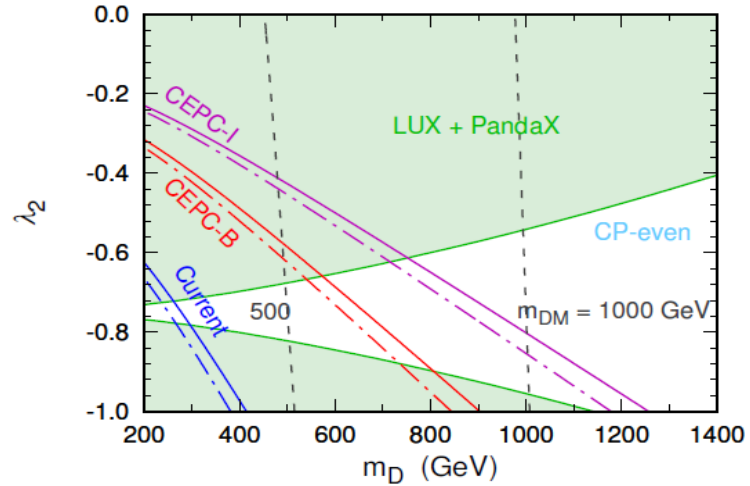
Fedderke, Lin, Wang, 1506.05465

- ⊕ Current: current precision for EW oblique parameters
- CEPC-B: CEPC baseline precision
- CEPC-I: CEPC precision with improvements of m_Z , G_Z , and m_t measurement

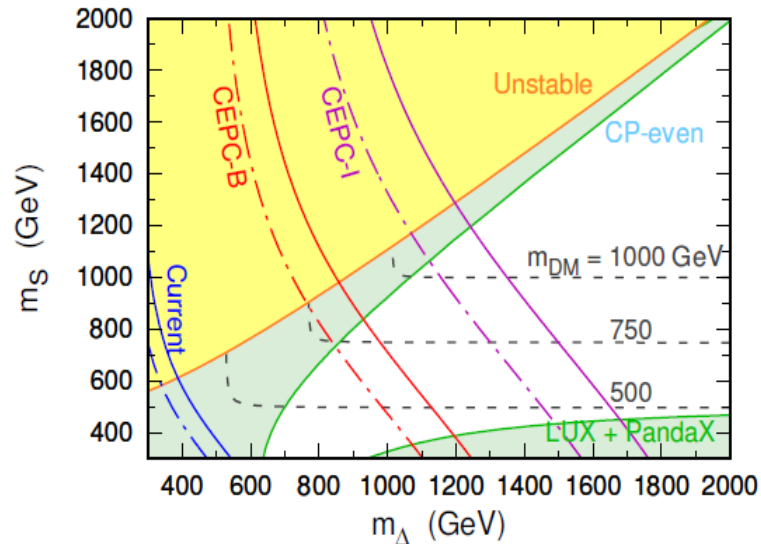
Oblique parameters: scalar DM



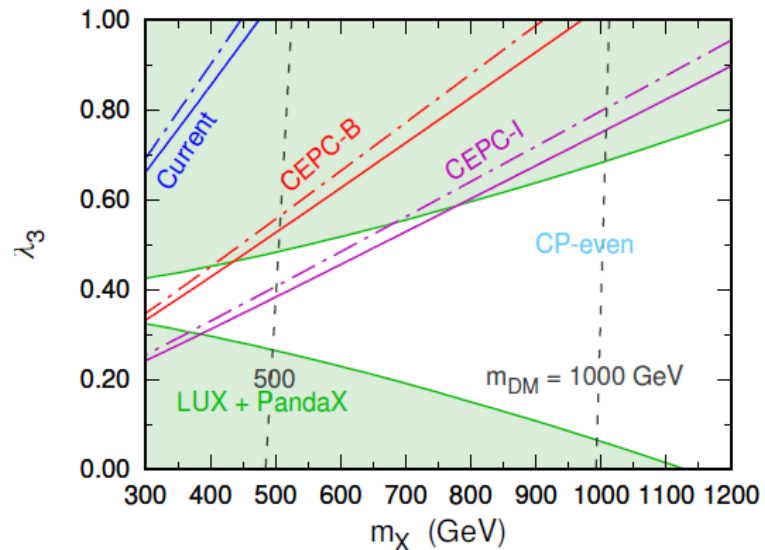
Singlet and Doublet Scalars



Doublet Scalar

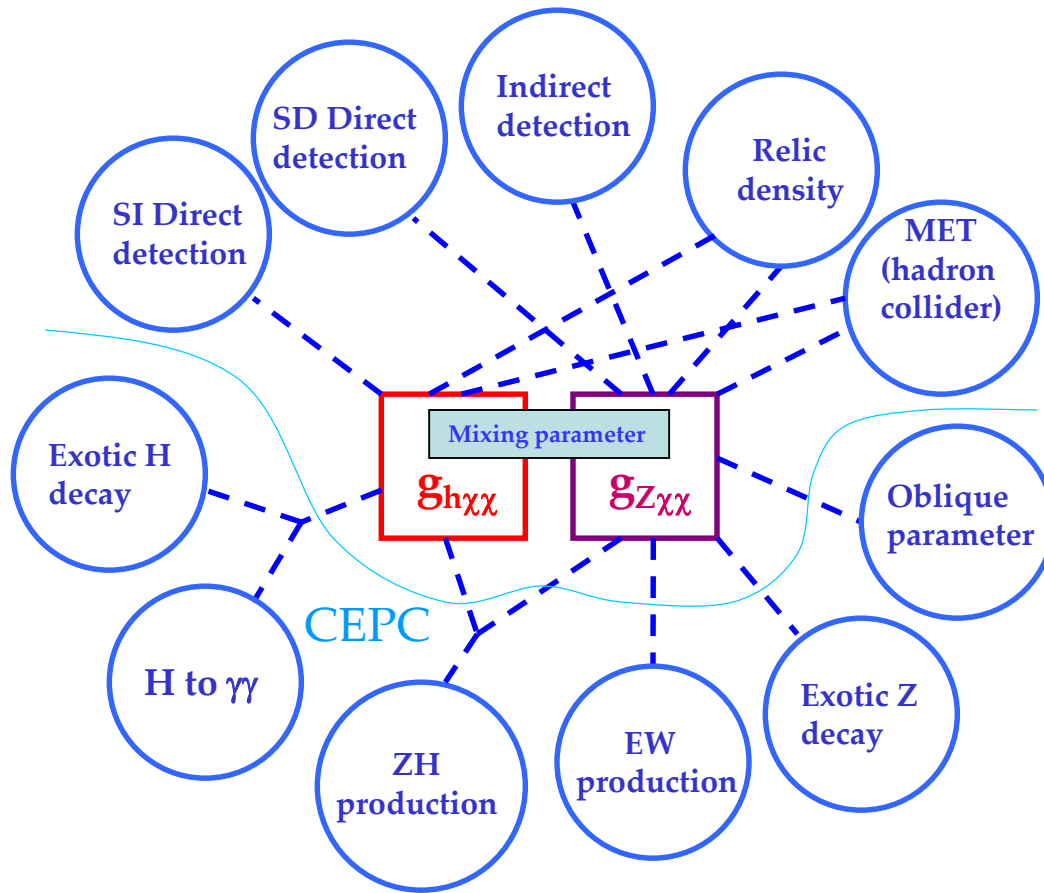


Singlet and Triplet Scalars



Quadruplet Scalar

Summary



✦ *It is possible to probe new particles directly produced or via loop effects at CEPC*

✦ *We consider a kind of DM models containing additional EW multiplets, and focus on their corrections to the EW processes at the CEPC*

✦ *The significant signatures at the CEPC require moderate interactions connecting new particles to Higgs and Z bosons. These interactions can be constrained by collider and DM detection experiments*

Thank you