Common origin of dark matter and leptogenesis in $U(1)_{B-L}$

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Outline



2 The Model

The global U(1)_{B-L} scenario

4 The local $U(1)_{B-L}$ scenario



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Motivation

Solving simultaneously three sharp BSM questions:

- Baryon asymmetry
- Dark matter
- Neutrino mass

Who can connect them together? \rightarrow sterile neutrinos *N* What is the mass scale of connector? \rightarrow TeV (colliders favor) What symmetry do they follow? $\rightarrow U(1)_{B-L}$ What is the corresponding mechanism?

- Resonant leptogenesis
- Sterile neutrino portal WIMP
- Type-I seesaw

4 **A** N A **B** N A **B** N

The Model

Extending $\phi = (v_{\phi} + \tilde{\rho} + i\eta)/\sqrt{2}$, Majorana sterile neutrinos *N* and DM χ

| | q_L | U _R | d _R | L_L | e _R | Ν | χ | Н | ϕ |
|--------------|----------------|----------------|----------------|----------------|----------------|----|----|---------------|--------|
| $SU(2)_L$ | 2 | 1 | 1 | 2 | 1 | 1 | 1 | 2 | 1 |
| $U(1)_Y$ | $+\frac{1}{6}$ | $+\frac{2}{3}$ | $-\frac{1}{3}$ | $-\frac{1}{2}$ | -1 | 0 | 0 | $\frac{1}{2}$ | 0 |
| $U(1)_{B-L}$ | $+\frac{1}{3}$ | $+\frac{1}{3}$ | $+\frac{1}{3}$ | -1 | -1 | -1 | -1 | 0 | 2 |
| Z_2 | + | + | + | + | + | + | _ | + | + |

Table: Relevant particle contents and corresponding charge assignments.

The relevant Yukawa interaction:

$$\mathcal{L} \supset -\frac{\lambda_{N}}{\sqrt{2}}\phi N^{T}CN - \frac{\lambda_{\chi}}{\sqrt{2}}\phi \chi^{T}C\chi - y_{\nu}\bar{L}_{L}HP_{R}N + h.c..$$
(1)

The free parameters are $\{m_{\rho}, m_{\eta,Z'}, m_N, m_{\chi}, \tilde{m}, v_{\phi}, \theta, \varepsilon_{\text{CP}}\}$ Fixing $\tilde{m} = 5 \times 10^{-11} \text{ GeV}$.

The Model



Figure: The dominant Feynman diagrams for *N* and χ annihilation. (a)-(c) are universal for the global and the local scenarios. In (d), *Z'* mediated \rightarrow local, ρ mediated \rightarrow global

The relevant Boltzmann equations in global scenario:

$$\frac{dY_{N}}{dz} = -\frac{z}{sH(m_{N})} \left(\left(\frac{Y_{N}}{Y_{N}^{eq}} - 1 \right) (\gamma_{N \to HL} + 2\gamma_{NL \to qt} + 4\gamma_{Nt \to qL}) \right)$$

$$- \frac{z}{sH(m_{N})} \left(\left(\left(\frac{Y_{N}}{Y_{N}^{eq}} \right)^{2} - 1 \right) (2\gamma_{NN \to \rho\rho, \eta\eta} + 2\gamma_{NN \to \rho\eta, h\eta} + 2\gamma_{NN \to VV}) \right)$$

$$+ \frac{z}{sH(m_{N})} \left(\left(\left(\frac{Y_{\chi}}{Y_{\chi}^{eq}} \right)^{2} - \left(\frac{Y_{N}}{Y_{N}^{eq}} \right)^{2} \right) 2\gamma_{\chi\chi \to NN}, \quad (2)$$

$$\frac{dY_{\chi}}{dz} = -\frac{z}{sH(m_{N})} \left(\left(\left(\frac{Y_{\chi}}{Y_{\chi}^{eq}} \right)^{2} - 1 \right) (2\gamma_{\chi\chi \to \rho\rho, \eta\eta} + 2\gamma_{\chi\chi \to \rho\eta, h\eta} + 2\gamma_{\chi\chi \to VV}) \right)$$

$$- \frac{z}{sH(m_{N})} \left(\left(\left(\frac{Y_{\chi}}{Y_{\chi}^{eq}} \right)^{2} - \left(\frac{Y_{N}}{Y_{N}^{eq}} \right)^{2} \right) 2\gamma_{\chi\chi \to NN}, \quad (3)$$

$$\frac{dY_{B-L}}{dz} = \frac{z}{sH(m_{N})} \left(\varepsilon_{CP} \left(\frac{Y_{N}}{Y_{N}^{eq}} - 1 \right) - \frac{Y_{B-L}}{2Y_{L}^{eq}} \right) \gamma_{N \to HL}$$

$$- \frac{z}{sH(m_{N})} \frac{Y_{B-L}}{Y_{L}^{eq}} \left(\frac{Y_{N}}{Y_{N}^{eq}} + 2\gamma_{NL \to qL} \right), \quad (4)$$

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Figure: The evolutions of various abundances Y_i .

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Figure: Y_i as a function of m_ρ at present.

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Figure: Perturbation constraints and exceeding ε_{CP} .

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Figure: Branching ratio of Higgs invisible decay.

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October 11,2024 10/20



$$\sigma_{\rm SI} = \frac{C^2 (\sin 2\theta)^2 m_h^4 m_{\chi}^4}{4 \pi v_H^2 v_{\phi}^2 (m_h + m_{\chi})^2} \\ \times \left(\frac{1}{m_h^2} - \frac{1}{m_{\rho}^2}\right)^2, (6)$$

LZ (present) limit: $\sigma_{\rm SI}\gtrsim 10^{-46}{\rm cm}^2$

LZ (future) limit:: $\sigma_{\rm SI}\gtrsim 10^{-47}{\rm cm}^2$

LZ (future): Most allowed samples.

Figure: Spin-independent scattering cross section.

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October 11,2024 11/20



Resonance: p wave $\langle \sigma v \rangle \lesssim 10^{-28} \ {\rm cm}^3/{\rm s}$

Non resonance: s wave $\langle \sigma v \rangle \simeq 2 \times 10^{-26} \text{ cm}^3/\text{s}$

Non resonance: $\langle \sigma v \rangle$ of $\chi \chi \rightarrow \rho \eta \rightarrow W^+ W^- \eta$ is around $10^{-27} \text{cm}^3/\text{s}$.

Figure: Constraints from the indirect detections.

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The relevant Boltzmann equations in local scenario:

$$\frac{dY_{N}}{dz} = -\frac{z}{sH(m_{N})} \left(\frac{Y_{N}}{Y_{N}^{eq}} - 1 \right) \left(\gamma_{N \to HL} + 2\gamma_{NL \to qt} + 4\gamma_{Nt \to qL} \right)
- \frac{z}{sH(m_{N})} \left(\left(\frac{Y_{N}}{Y_{N}^{eq}} \right)^{2} - 1 \right) \left(2\gamma_{NN \to \rho\rho, Z'Z'} + 2\gamma_{NN \to \rhoZ', hZ'} + 2\gamma_{NN \to SMSM} \right)
+ \frac{z}{sH(m_{N})} \left(\left(\frac{Y_{\chi}}{Y_{\chi}^{eq}} \right)^{2} - \left(\frac{Y_{N}}{Y_{N}^{eq}} \right)^{2} \right) 2\gamma_{\chi\chi \to NN}, \qquad (8)
\frac{dY_{\chi}}{dz} = -\frac{z}{sH(m_{N})} \left(\left(\frac{Y_{\chi}}{Y_{\chi}^{eq}} \right)^{2} - 1 \right) \left(2\gamma_{\chi\chi \to \rho\rho, Z'Z'} + 2\gamma_{\chi\chi \to \rhoZ', hZ'} + 2\gamma_{\chi\chi \to SMSM} \right)
- \frac{z}{sH(m_{N})} \left(\left(\frac{Y_{\chi}}{Y_{\chi}^{eq}} \right)^{2} - \left(\frac{Y_{N}}{Y_{N}^{eq}} \right)^{2} \right) 2\gamma_{\chi\chi \to NN}, \qquad (9)
\frac{dY_{B-L}}{dz} = \frac{z}{sH(m_{N})} \left(\varepsilon_{CP} \left(\frac{Y_{N}}{Y_{N}^{eq}} - 1 \right) - \frac{Y_{B-L}}{2Y_{L}^{eq}} \right) \gamma_{N \to HL}
- \frac{z}{sH(m_{N})} \frac{Y_{B-L}}{Y_{L}^{eq}} \left(\frac{Y_{N}}{Y_{N}^{eq}} + 2\gamma_{NL \to qL} \right), \qquad (10)$$

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October 11,2024 13/20



Fixing $m_N = 1000, m_\rho = 500$ $m_{Z'} = 7000 \varepsilon_{\rm CP} = 0.1$

 $\begin{array}{l} \text{Solid} {\rightarrow} v_{\phi} = 4000 \\ \text{Dashed} {\rightarrow} v_{\phi} = 5000 \\ \text{Dot-dashed} {\rightarrow} v_{\phi} = 6000 \end{array}$

Figure: The evolutions of various abundances Y_i .

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Figure: Y_i as a function of m_ρ at present.

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Figure: Perturbation constraints and exceeding ε_{CP} .

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16/20



Figure: The current LEP-II and prospective ILC limits

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October 11.2024 17/20



Figure: Constraints from the direct detection experiment.

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Resonance: p wave $\langle \sigma v \rangle \lesssim 10^{-27} \text{ cm}^3/\text{s}$

Non resonance: s wave $\langle \sigma v \rangle \simeq 2 \times 10^{-26} \text{ cm}^3/\text{s}$

Promising channel: $\chi\chi \rightarrow \rho Z' \rightarrow W^+ W^- \bar{t}t.$ (a) and (b) with $m_\chi \gtrsim 5800 \text{ GeV}$

Figure: Constraints from the indirect.

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19/20

Global scenario

Relic density: resonance and non resonance Higgs invisible decay:future CEPC-non resonance Direct detection:current and future LZ Indirect detection: unpromising

Local scenario

Relic density: resonance and non resonance Higgs invisible decay:unpromising Direct detection:current and future LZ Indirect detection: promising in non resonance in (a) and (b)

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