



# **Determining heavy meson LCDAs from lattice QCD**

Based on <u>2403.17492</u>, <u>2410.18654</u> et al.

In collaboration with LPC members and J. Xu, S. Zhao, et al.

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> Theoretical framework for the two-step factorization method

- Lattice QCD verification
- Phenomenological discussions
- Summary and prospect

Heavy flavor physics is one of the frontier topics in particle physics:



Current experimental results show deviations from theoretical prediction...



High precision calculations play a crucial role in the search for new physics!

Uncertainties originating from *B* meson LCDAs dominate the primary errors in theoretical calculation.

• For example:  $B \rightarrow \pi$ ,  $K^*$  form factors from LCSRs:

Gao, Lu, Shen, Wang, Wei, PRD 101 (2020) 074035 Cui, Huang, Shen, Wang, JHEP 03 (2023) 140

$$\begin{aligned} \mathcal{V}_{B\to K^*}(0) &= 0.359_{-0.085}^{+0.141} \Big|_{\lambda_B} \stackrel{+0.019}{_{-0.019}} \Big|_{\sigma_1} \stackrel{+0.001}{_{-0.062}} \Big|_{\mu} \stackrel{+0.010}{_{-0.004}} \Big|_{M^2} \stackrel{+0.016}{_{-0.017}} \Big|_{s_0} \stackrel{+0.153}{_{-0.079}} \Big|_{\varphi_{\pm}(\omega)}, \\ f_{B\to\pi}^+(0) &= 0.122 \times \left[ 1 \pm 0.07 \Big|_{S_0^{\pi}} \pm 0.11 \Big|_{\Lambda_q} \pm 0.02 \Big|_{\lambda_E^2/\lambda_H^2} \stackrel{+0.05}{_{-0.066}} \Big|_{M^2} \pm 0.05 \Big|_{2\lambda_E^2+\lambda_H^2} \right]_{\lambda_E^2+\lambda_H^2} \\ &+ 0.06 \Big|_{\mu_h} \pm 0.04 \Big|_{\mu_0} \stackrel{+1.36}{_{-0.56}} \Big|_{\lambda_B} \stackrel{+0.25}{_{-0.43}} \Big|_{\sigma_1,\sigma_2}. \end{aligned}$$

 $\lambda_B$  and  $\sigma_n$ : the first inverse and inverse-log moments,

 $\varphi_B^+$ : uncertainties from different parameterizations of the *B* meson LCDA.

Without reliable B LCDA, it is impossible to discuss precision calculation!

• HQET field  $h_{v}$  and the associated cusp divergence hinder the first-principle calculation of *B* meson LCDA.

$$i\tilde{f}_{H}(\mu)m_{H}\varphi^{+}(\omega,\mu) = \int_{-\infty}^{+\infty} \frac{dt}{2\pi} e^{i\omega n_{+}\cdot vt}$$
$$\times \langle 0|\bar{q}(tn_{+})\not n_{+}\gamma_{5}W_{c}(tn_{+},0)h_{v}(0)|H(v)\rangle$$

- Only available through model parametrizations now, predictions from different models vary significantly.
- This model dependence contribute to the largest theoretical uncertainties in  $B \rightarrow K^*$  form factor:

 $\mathcal{V}_{B\to K^*}(0) = 0.359^{+0.141}_{-0.085} \Big|_{\lambda_B} \Big|_{\sigma_1} \Big|_{\sigma_1} \Big|_{\sigma_1} \Big|_{\sigma_1} \Big|_{\sigma_1} \Big|_{\mu} \Big|_{M^2} \Big|_{M^2} \Big|_{M^2} \Big|_{\sigma_1} \Big|_{\sigma_1} \Big|_{\sigma_2} \Big|_{\sigma_1} \Big|_{\sigma_2} \Big|_{\sigma$ 

Grozin, Neubert, PRD 55, 272 (1997) Braun, Ivanov, Korchemsky, PRD 69, 034014 (2004) Beneke, Braun, Ji, Wei, JHEP 07, 154 (2018)



Gao, Lu, Shen, Wang, Wei, PRD 101 (2020) 074035

#### A two-step factorization method

We propose a **chained factorization formula** to determine both <u>heavy meson QCD</u> <u>LCDA</u> and <u>HQET LCDA</u> from the first-principle:



- LaMET:  $\Lambda_{\rm QCD}$ ,  $m_H \ll P^z$  and integrate out  $P^z$
- bHQET:  $\Lambda_{\text{QCD}} \ll m_H$  and integrate out  $m_H$ 
  - $\Rightarrow$  Introduce a hierarchy  $\Lambda_{\rm QCD} \ll m_H \ll P^z$

#### A two-step factorization method

The hierarchy  $\Lambda_{\rm QCD} \ll m_H \ll P^z$  imposes limitations on lattice calculation:

• Currently, direct simulation of *B* meson is not practical.



- *D* meson can be realized on the lattice;
- <u>Heavy quark flavor symmetry</u> ensures that the HQET LCDA is <u>independent</u> of heavy quark mass;
- $m_H (m_D \text{ or } m_B)$  only contributes to the power corrections.

 A numerical simulation on the finest CLQCD ensemble (a = 0.05187 fm); CLQCD Collaboration, PRD 109, 054507 (2024)

• Simulate the *D* meson quasi DA with  $m_D \simeq 1.92$ GeV, up to  $P^Z \simeq 3.98$ GeV;

$$\tilde{\phi}(x, P^z) = \int \frac{dz}{2\pi} e^{-ixP^z z} \tilde{M}(z, P^z)$$



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• Matching formula in LaMET:

$$\tilde{\phi}(x,P^z) = \int_0^1 C\left(x,y,\frac{\mu}{P^z}\right)\phi(y,\mu) + \mathcal{O}\left(\frac{m_H^2}{(P^z)^2},\frac{\Lambda_{\rm QCD}^2}{(xP^z,\bar{x}P^z)^2}\right)$$

Matching kernel at NLO in  $\alpha_s$ 

Liu, Wang, Xu, QAZ, Zhao, PRD 99, 094036 (2019) Han, Hua, Ji, Lu, Wang, Xu, QAZ, Zhao, 2410.18654



fraction;

LCDA.

Related to the HQET

Ishaq, Jia, Xiong, Yang, PRL125(2020)132001

Beneke, Finauri, Vos, Wei, JHEP 09, 066 (2023)

Peak region: 
$$y \sim \frac{\Lambda_{\text{QCD}}}{m_H}$$
  
• Light quark carries  
small momentum

QCD LCDA of D meson

End-point region:

- LaMET matching kernel suffer large power corrections.
- Lattice QCD predictions fail

Tail region:  $y \sim 1$ 

- Contain only <u>hard-</u> <u>collinear</u> physics, perturbative calculable;
- Suppressed in LCDA.

- Peak region:
  - A multiplicative factorization from QCD LCDA to
     HQET LCDA:

$$\varphi_{\text{peak}}^{+}(\omega,\mu) = \frac{f_{H}}{\widetilde{f}_{H}} \frac{1}{\mathcal{J}_{\text{peak}}} \phi(y,\mu;m_{H}) + \mathcal{O}\left(\frac{\Lambda_{\text{QCD}}}{m_{H}}\right)$$

Beneke, Finauri, Vos, Wei, JHEP 09, 066 (2023)

• Nonperturbative, determined from lattice QCD.



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Beneke, Finauri, Vos, Wei, JHEP 09, 066 (2023)

- Nonperturbative, determined from lattice QCD.
- ➤ Tail region:
  - Perturbative result at 1-loop order:

$$\varphi_{\text{tail}}^{+}(\omega,\mu) = \frac{\alpha_{s}C_{F}}{\pi\omega} \left[ \left(\frac{1}{2} - \ln\frac{\omega}{\mu}\right) + \frac{4\bar{\Lambda}}{3\omega} \left(2 - \ln\frac{\omega}{\mu}\right) \right]$$

Lee, Neubert, PRD 72, 094028 (2005)



#### **Discussions I:** comparison with models

- The model dependence contributes the largest systematic error in the form factors.
- Our result is basically consistent with most of the model estimates, and will also provide a first-principle constrains on the existing models.



• For theoretical calculations, result from first-principles will help to REMOVE the primary uncertainties arising from the model parametrizations.

#### **Discussions II:** Inverse and inverse-logarithmic moments

• Significant uncertainties from  $\lambda_B$  and  $\sigma_1$ :

Gao, Lu, Shen, Wang, Wei, PRD 101 (2020) 074035

$$\mathcal{V}_{B\to K^*}(0) = 0.359^{+0.141}_{-0.085} \Big|_{\lambda_B} \Big|_{\sigma_1} \Big|_{\sigma_1} \Big|_{0.062} \Big|_{\mu} \Big|_{-0.004} \Big|_{M^2} \Big|_{M^2} \Big|_{\sigma_1} \Big|_{s_0} \Big|_{\varphi_{\pm}(\omega)},$$

• A model-independent parametrization form:

$$\varphi^{+}(\omega,\mu) = \sum_{n=1}^{N} c_n \frac{\omega^n}{\omega_0^{n+1}} e^{-\omega/\omega_0}$$
$$= \frac{c_1 \omega}{\omega_0^2} \left[ 1 + c_2' \frac{\omega}{\omega_0} + c_3' \left(\frac{\omega}{\omega_0}\right)^2 + \cdots \right] e^{-\omega/\omega_0},$$

Fit results up to the *N*-th order:

$$N = 1: \ \omega_0 = 0.403(44), \ c_1 = 0.932(73);$$
  

$$N = 2: \ \omega_0 = 0.352(82), \ c_1 = 0.69(37),$$
  

$$c'_2 = 0.17(32);$$
  

$$N = 3: \ \omega_0 = 0.32(15), \ c_1 = 0.63(44),$$
  

$$c'_2 = 0.12(37), c'_3 = 0.04(19).$$



#### **Discussions II:** Inverse and inverse-logarithmic moments

Numerical results of  $\lambda_B$  and  $\sigma_B^{(1)}$  at  $\mu = 1$ GeV:

		$\lambda_B$ (GeV)	$\sigma_{\!B}^{(1)}$
Our results	N=1	0.389(35)	1.63(8)
	N=2	0.393(37)	1.62(7)
	N=3	0.381(59)	1.63(12)
Experiment	Belle 2018	> 0.24	
	Khodjamirian, Mandal, Mannel, 2020	0.383(153)	
	Gao, Lu, Shen, Wang, Wei, 2020	$0.343\substack{+0.064\\-0.079}$	
Other	Lee, Neubert, 2005	0.48(11)	1.6(2)
approach	Braun, Ivanov, Korchemsky, 2004	0.46(11)	1.4(4)
	Grozin, Neubert, 1997	0.35(15)	
	Mandal, Nandi, Ray, 2024	0.338(68)	

#### Discussions III: Impact on $B \rightarrow V$ form factors

An accurate  $\lambda_B$  will significantly improve the prediction for the  $B \rightarrow K^*$  form factors:

Gao, Lu, Shen, Wang, Wei, PRD 101 (2020) 074035					
GLSWW			Our result		
Error of $\mathcal{V}(0)$ :	0.23	$\rightarrow$	0.11		
$\lambda_B$ :	$0.343^{+64}_{-79}$	$\rightarrow$	0.389(35)		



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	GLSWW		Our result		
Error of $\mathcal{V}(0)$ :	0.23	$\rightarrow$	0.11		
$\lambda_B$ :	$0.343^{+64}_{-79}$	$\rightarrow$	0.389(35)		

We are looking forward to a <u>more precise analysis</u> of the form factors and accordingly physical observables.

$$\mathcal{V}_{B\to K^*}(0) = 0.359_{-0.085}^{+0.141} \left|_{\lambda_B}^{+0.019} \right|_{\sigma_1}^{+0.001} \left|_{\sigma_1}^{+0.001} \right|_{\mu} \\ +0.010 \left|_{M^2}^{+0.016} \right|_{M^2}^{+0.016} \left|_{s_0}^{+0.141} \right|_{\sigma_1}^{+0.001} \left|_{\varphi_{\pm}(\omega)}^{+0.001} \right|_{\varphi_{\pm}(\omega)},$$



- ✓ We present a first lattice-implementable method to extract the heavy meson LCDA, and implement it on a CLQCD ensemble.
- ✓ Although the results are preliminary, they can be continually improved.
- The phenomenological implications demonstrate that our results will significantly advance the theoretical studies towards the frontier of high precision.

More importantly, improving the reliability of our results for the next stage:

- How to properly <u>control the power corrections</u> within two step factorization?
- More systematic lattice QCD calculations: more a, larger  $P^{z}$ , ...

Thanks for your attention!